

# Gravity, Spin Models and Continuous Phase Transitions

**Umut Gürsoy**

**(CERN)**

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**+ ongoing**

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- Paramagnet  $T > T_c$  to ferromagnet transition  $T < T_c$  as the system cools down.
- The  $U(1)$  XY model in 2D (Kosterlitz-Thouless model) or 3D and the “ $O(3)$  quantum rotor” in 3D, the Hubbard model. etc  $\Rightarrow$  **canonical models for super-fluidity/super-conductivity.**

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- **Non-trivial critical exponents** at  $T_c$  only computable by Monte-Carlo for  $D > 2$ .

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- Sound-speed of Goldstone mode:  $\vec{m} \Leftrightarrow |\vec{m}| e^{i\psi}$  then

$$F_L \sim \int |\vec{M}|^2 (\delta\psi)^2$$

In the **MFA**  $c_\psi \sim |\vec{M}|^2 \sim |T - T_c|$



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- Gauge theories  $\Rightarrow$  GR!
- A new approach to holographic super-fluids/super-conductors

# Lattice gauge theory and Spin-models

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- Any LGT with *arbitrary* gauge group  $G$  in  $d$ -dimensions with arbitrary *adjoint matter*
- Integrate out gauge invariant states  $\Rightarrow$  generate effective theory for the Polyakov loop
- $Z_{LGT}(P; T) \sim Z_{SpM}(\vec{s}; T^{-1})$
- Ferromagnetic spin model  $\mathcal{H} = -J \sum_{\langle ij \rangle} \vec{s}_i \cdot \vec{s}_j + \dots$   
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- **Inversion of temperature:**  
Deconf. (high T) phase in LGT  $\Leftrightarrow$  Ordered (low T) phase of SpM  
Conf. (low T) phase in LGT  $\Leftrightarrow$  Disordered (high T) phase of SpM

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Lagrangian of LGT: electric  $U_{\vec{r},0}$  and magnetic  $U_{\vec{r},i}$  link variables



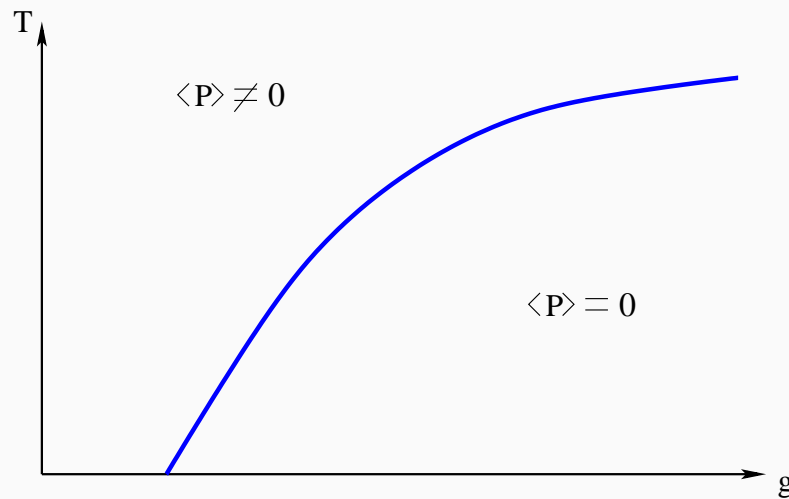
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$P \propto \prod_{n=0}^{N_t-1} U_{\vec{r}+n\hat{t},0}$  is the order parameter

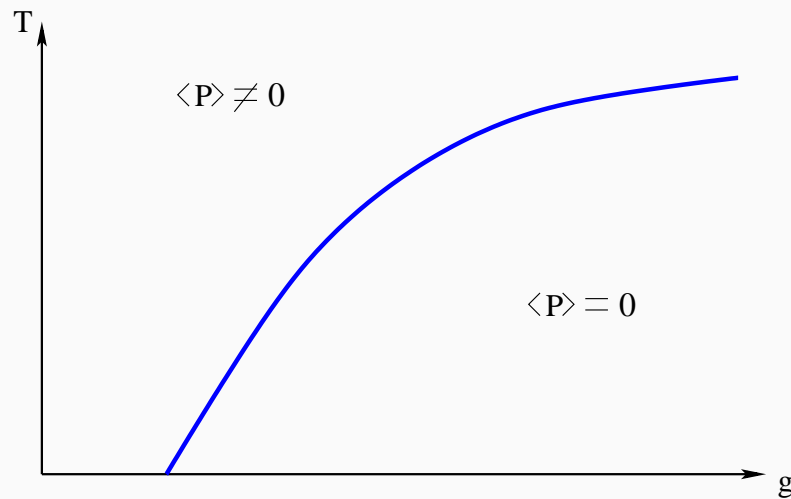
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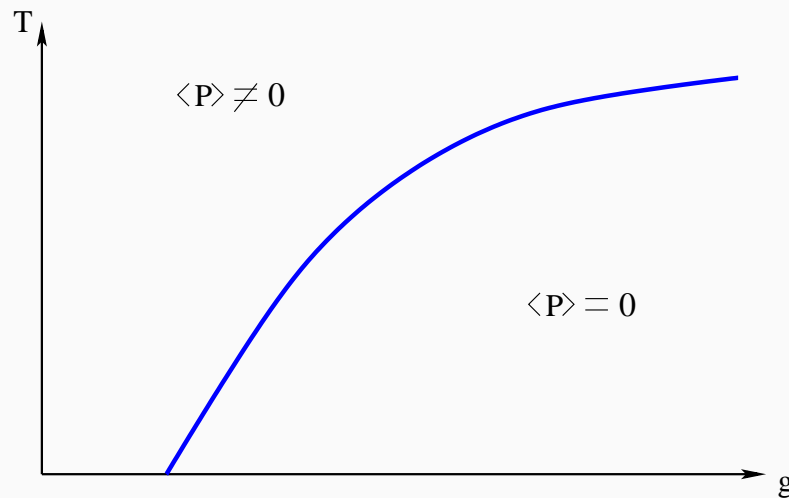
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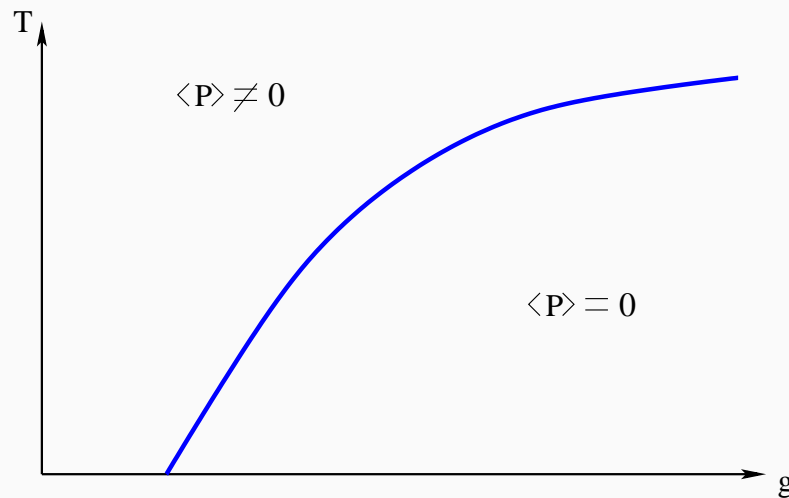
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- No long-range magnetic fluctuations  $\Rightarrow$  integrate out  $U_{\vec{r},j}$
- The resulting theory  $\mathcal{L}[P]$  describes long-range fluctuations at criticality
- Polyakov '78; Susskind '79: Can be mapped onto a spin-model with  $P \Leftrightarrow \vec{s}$  (explicitly shown in the limit  $g \gg 1$ )

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- **Some examples:**
  1. Pure  $SU(2)$  in  $d = 4$  second order transition with  $Z_2$  (Ising) critical exponents,
  2.  $SU(N)$  with  $N > 4$ ,  $d > 3$   
Spin model with  $Z_N$  fixed point:  $d = 4$  non-trivial  $U(1)$  XY model exponents,  $d > 4$  mean-field exponents.

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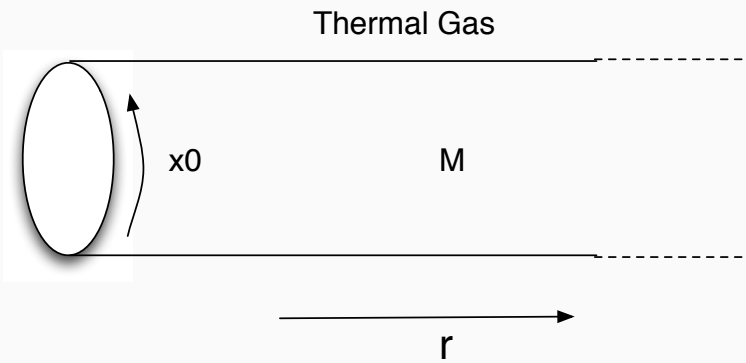
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 $SU(N)$  with  $N \rightarrow \infty$ ,  
Spin model with  $Z_N \rightarrow U(1)$  fixed point.

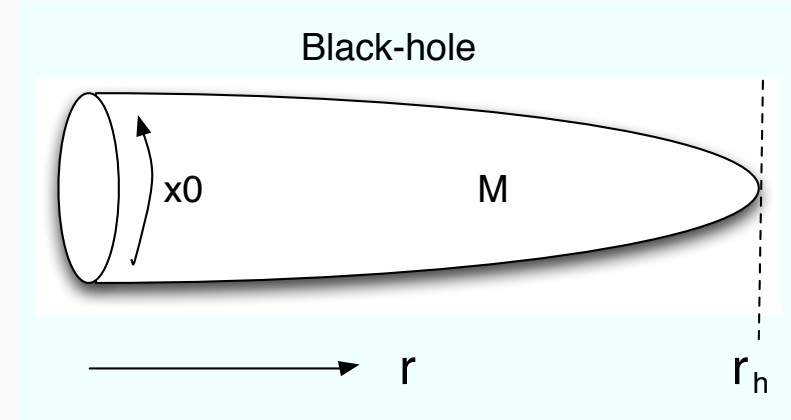


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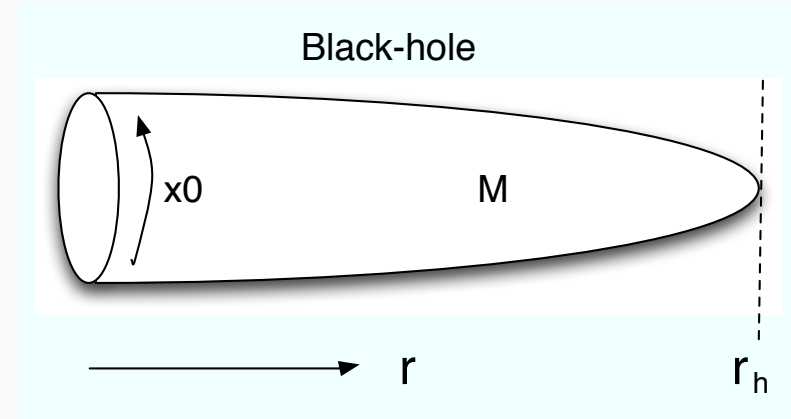
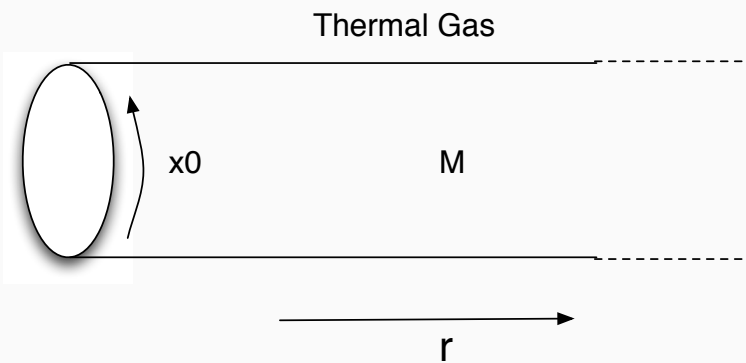


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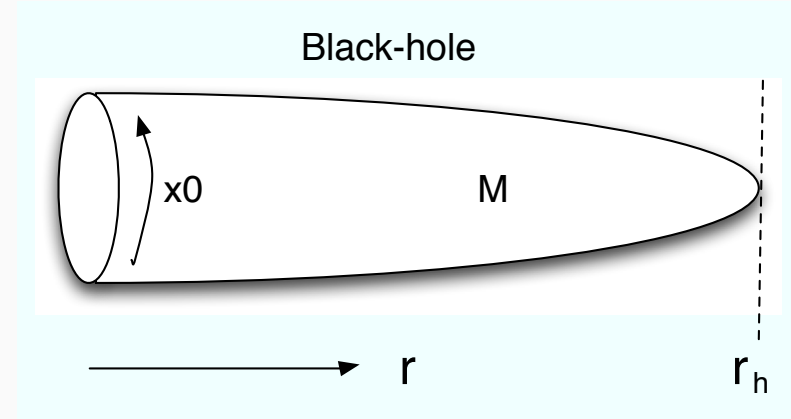
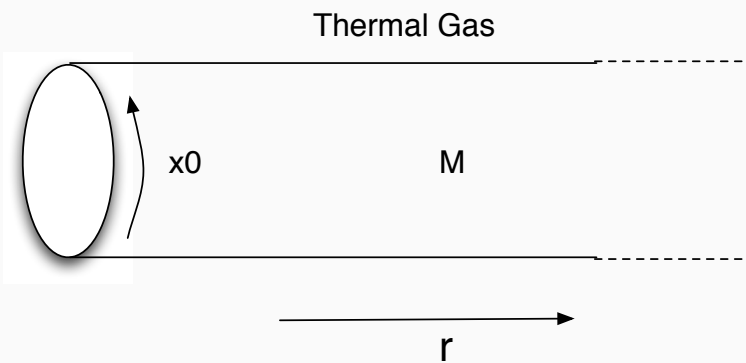


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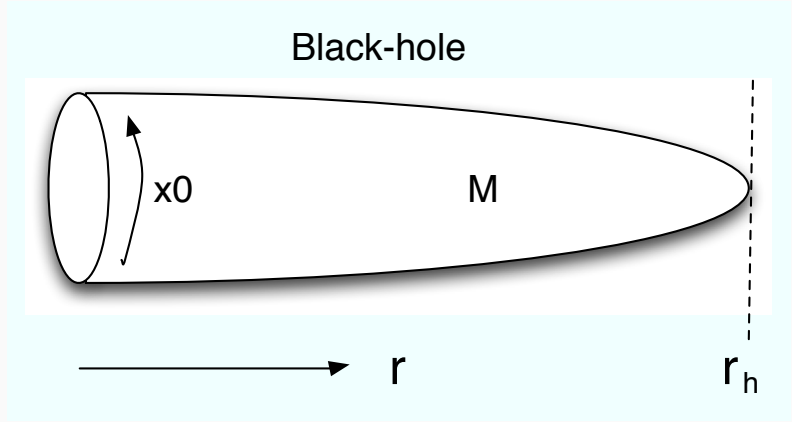
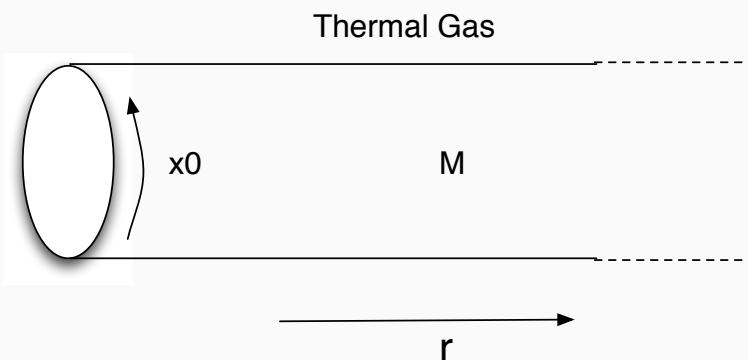


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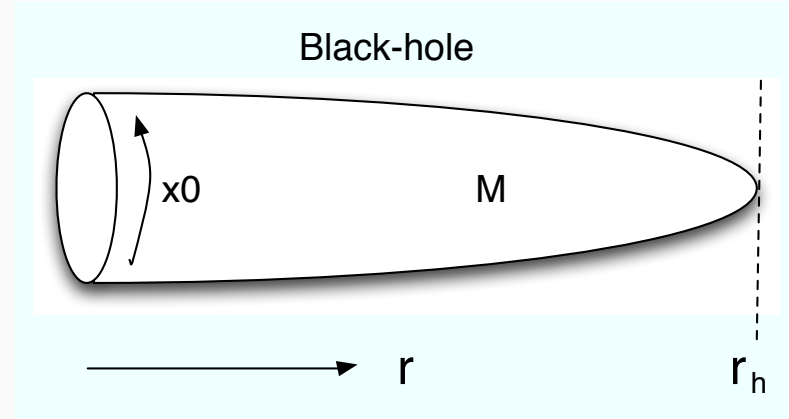
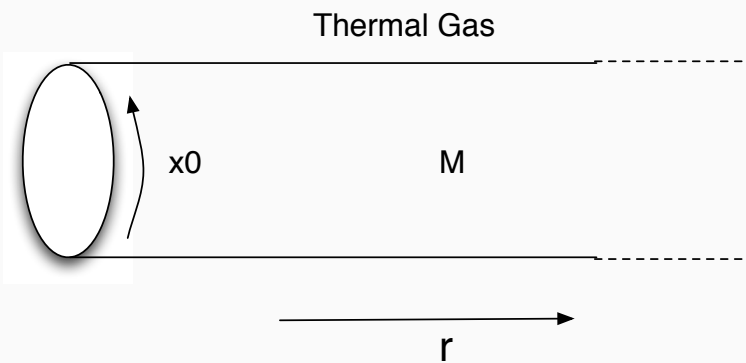


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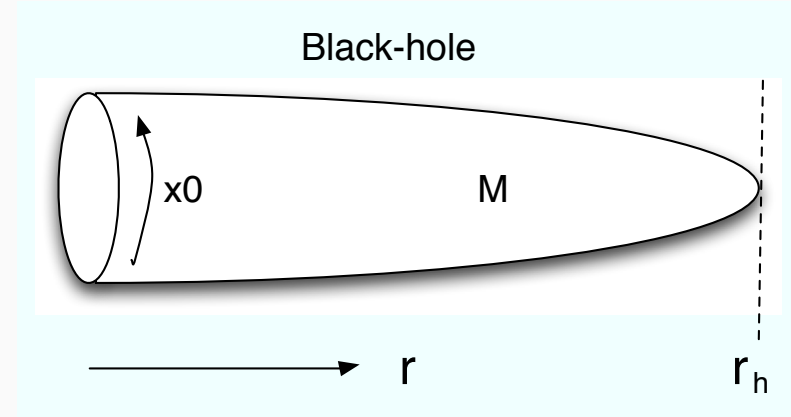
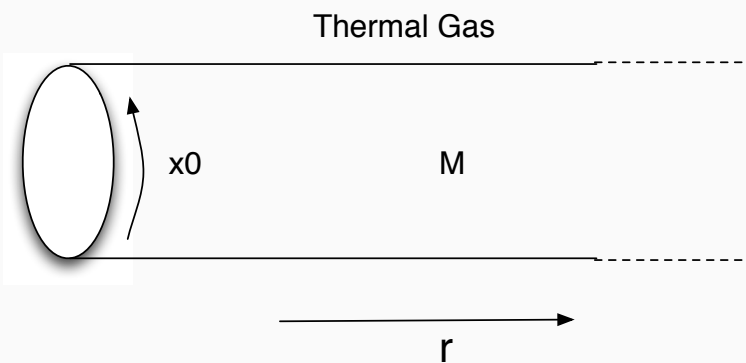


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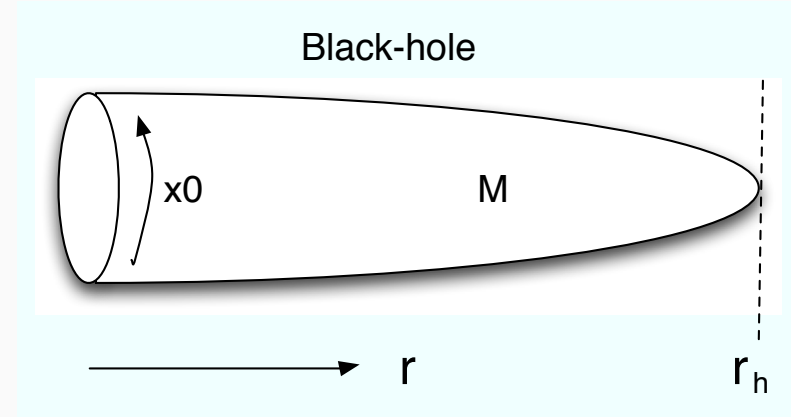
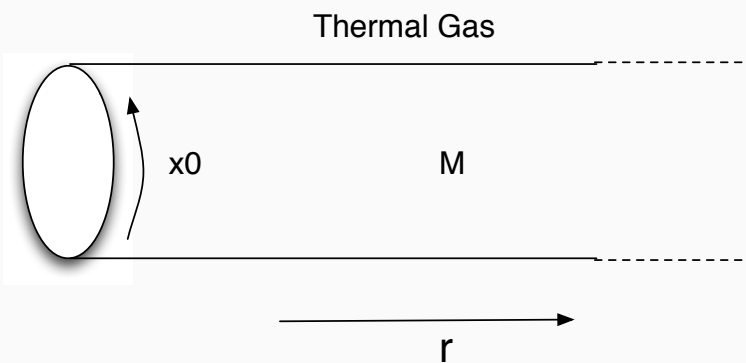


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- Fluctuations  $\delta\Psi \Leftrightarrow$  Goldstone mode in the dual spin-model



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**Continuous Hawking-Page  $\Leftrightarrow$  Normal-to-superfluid transition**

**GRAVITY/SPIN-MODEL CORRESPONDENCE**

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Look for solutions of the type:

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Requirements for a **second order Hawking-Page transition**:



# Continuous HP in dilaton-Einstein U.G. '10

Specify

$$\mathcal{S} \propto N^2 \int d^{d+1}x \sqrt{-g} \left( R - \xi(\partial\Phi)^2 + V(\Phi) - \frac{1}{12} e^{-\frac{8}{d-1}\Phi} (dB)^2 \right)$$

Look for solutions of the type:

$$ds_{TG}^2 = b_0^2(r) (dr^2 + dt^2 + dx_{d-1}^2)$$

$$ds_{BH}^2 = b^2(r) \left( \frac{dr^2}{f(r)} + f(r)dt^2 + dx_{d-1}^2 \right)$$

Requirements for a **second order Hawking-Page transition**:

- i.) There is a finite  $T_c$  at which:
- ii.)  $\Delta F(T_c) = 0$ . **TG(BH)** dominates for  $T < T_c$  ( $T > T_c$ ).
- iii.)  $\Delta S(T_c) = 0$
- iv.) Make sure that this happens between the thermodynamically favored BH and TG branches.

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- $n$ th order transition  $\Delta F \sim t^n$ :

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when  $V_{sub}(\Phi) = e^{-\kappa\Phi}$ , with  $\kappa = \sqrt{\frac{\zeta(d-1)}{n-1}}$  for  $n \geq 2$
- **BKT** scaling  $\Delta F \sim e^{-ct^{-\frac{1}{\alpha}}}$ :  
when  $V_{sub}(\Phi) = \Phi^{-\alpha}$ .

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- **Exact solution to string theory to all orders in  $\ell_s$ !**

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**One has to take into account  $\alpha'$  corrections.**
- Can be done because this regime is governed by a **linear-dilaton CFT** on the world-sheet!



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A consistent truncation of IIB with single scalar! [Pilch-Warner '00](#)

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$$ds_{TG}^2 = e^{-\frac{4}{3}\Phi_0} \frac{\cosh^{\frac{2}{3}}\left(\frac{3r}{2\ell}\right)}{\sinh^2\left(\frac{3r}{2\ell}\right)} (dt^2 + dx_{d-1}^2 + dr^2),$$

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! 2nd order HP at  $T_c$  happens in a sub-dominant branch

! Background in the string frame is NOT linear-dilaton

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- One finds  $c_\psi^2 \propto e^{-\sqrt{V_\infty} r_h} \sim (T - T_c)$ .
- **Second sound indeed vanishes at  $T_c$  with the mean-field exponent!**

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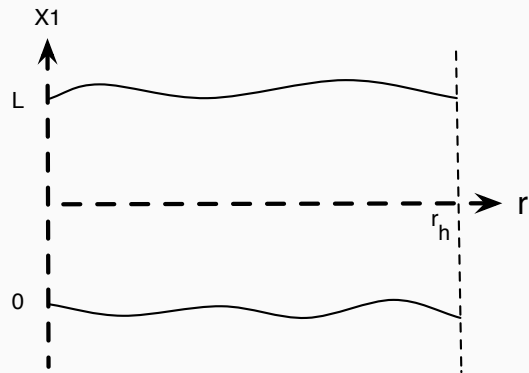
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**Mean-field scaling in the magnetization!**

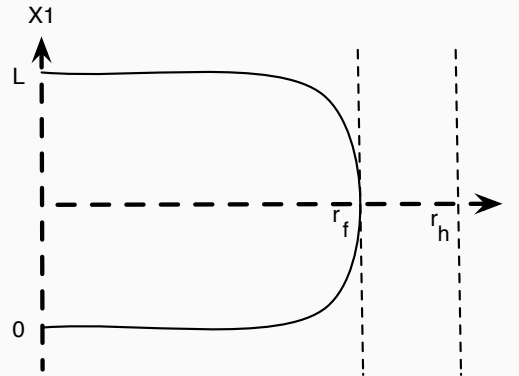


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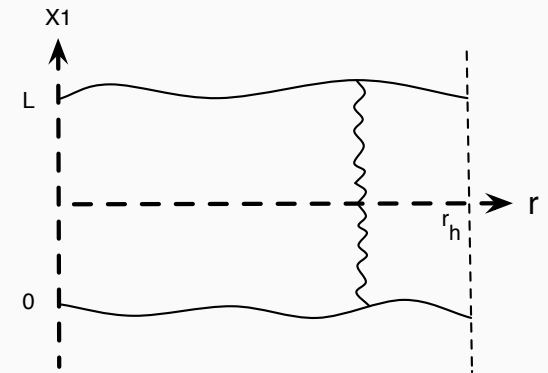
Three types of paths: D. Bak, A. Karch, L. Yaffe '07



(a)



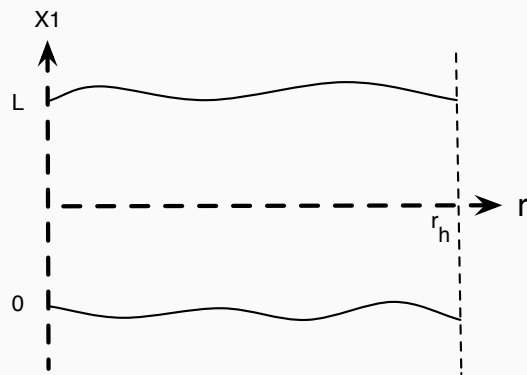
(b)



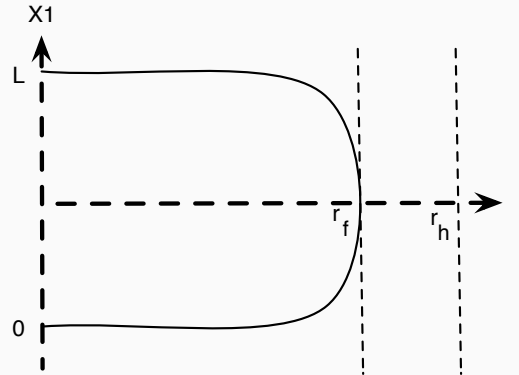
(c)

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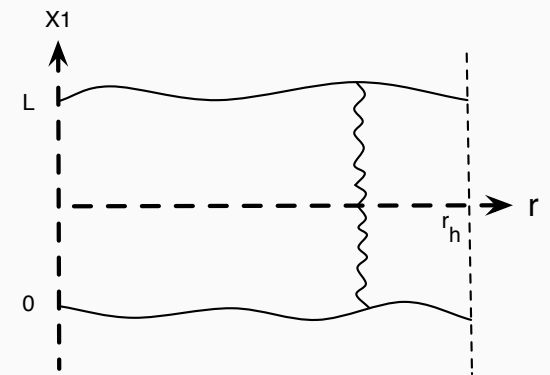
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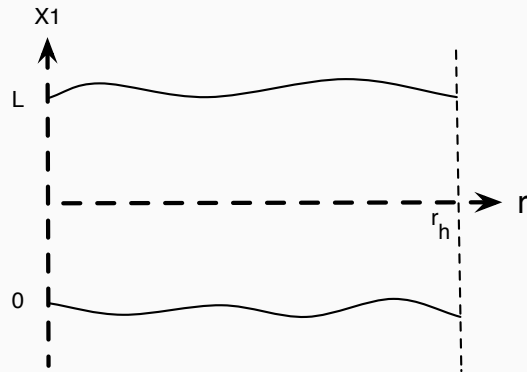


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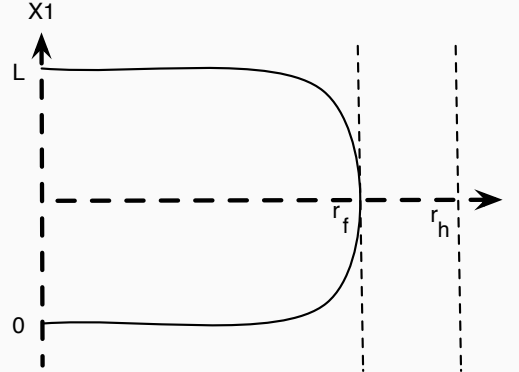
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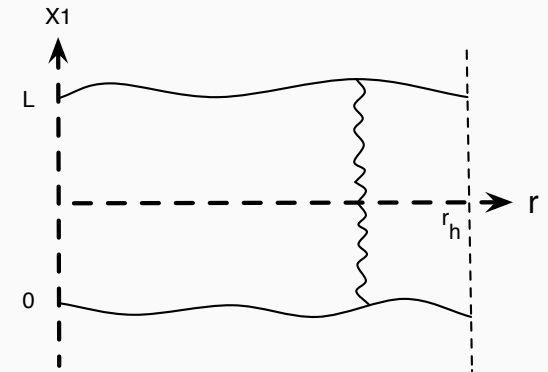
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(b)  $S_{F1} \rightarrow m_T L + \dots$

$\langle \vec{m}(L) \cdot \vec{m}(0) \rangle_b \sim e^{-m_T L + \dots}$  for  $L \gg 1$ .

# Two-point function, cont'ed

(c) bulk exchange diagrams:

$$\langle \vec{m}_{\parallel}(L) \cdot \vec{m}_{\parallel}(0) \rangle_c \propto \langle \text{Re}P[L] \text{Re}P[0] \rangle \sim \frac{e^{-m+L}}{L^{d-3}}$$

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$CT^-$  include a zero-mode:  $m_- = 0$  as  $\psi = \int_M B$  is modulus:

**Goldstone mode!**

Correct qualitative behavior:  $\langle \vec{m}_{\parallel}(L) \cdot \vec{m}_{\parallel}(0) \rangle \sim \frac{e^{-m_+L} + e^{-m_T L}}{L^{d-3}}$

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# Two-point function, cont'ed

(c) bulk exchange diagrams:

$$\langle \vec{m}_{\parallel}(L) \cdot \vec{m}_{\parallel}(0) \rangle_c \propto \langle \text{Re}P[L] \text{Re}P[0] \rangle \sim \frac{e^{-m_+L}}{L^{d-3}}$$

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$m_+$  minimum of the  $CT^+$  modes:  $G_{\mu\nu}, \Phi, \dots$

$m_-$  minimum of the  $CT^-$  modes:  $B_{\mu\nu}, \dots$

Spectrum analysis U.G., Kiritsis, Nitti '07:  $CT^+$  bounded from below for any T.

$CT^-$  include a zero-mode:  $m_- = 0$  as  $\psi = \int_M B$  is modulus:

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Precisely the expected behavior from the XY model,

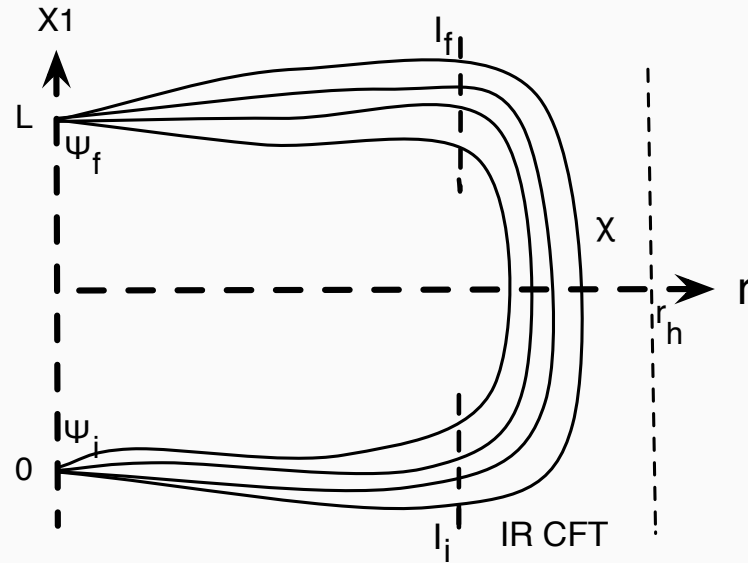
**with**  $\xi_{\parallel}^{-1} \rightarrow \min(m_T, m_+)$  for  $L \gg 1$ .



# Correlation length $\xi$

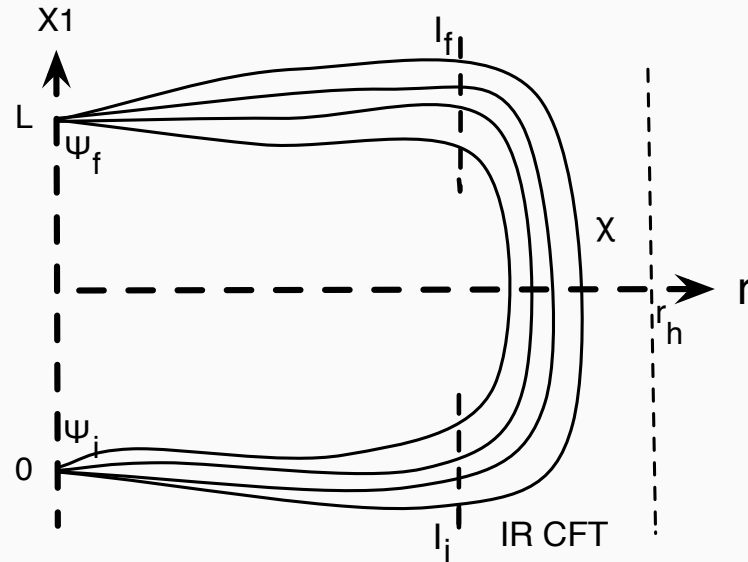
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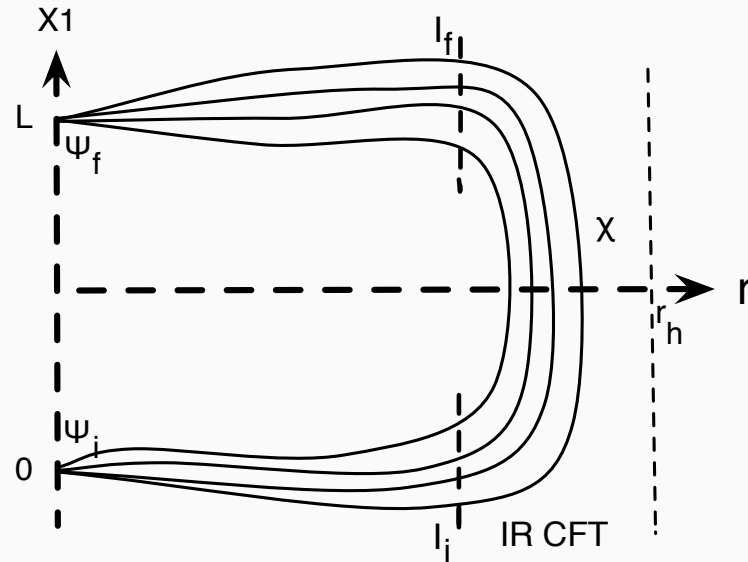
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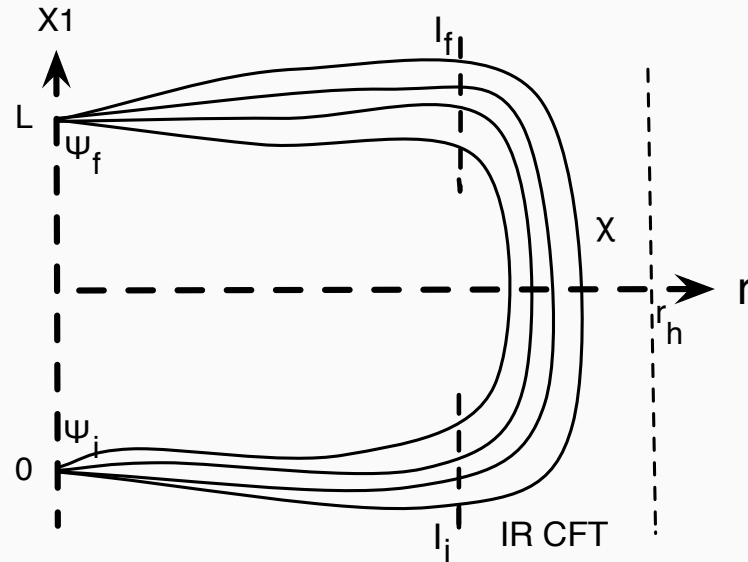
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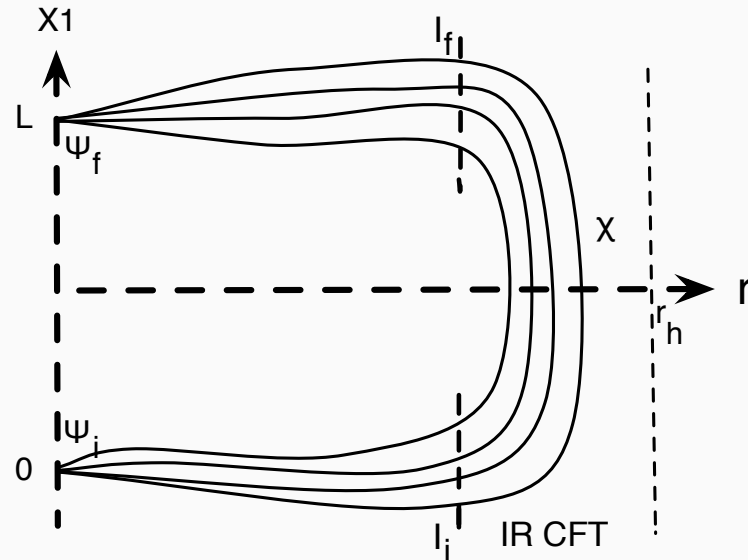
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Mean-field scaling again!

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- Probe strings  $\Leftrightarrow$  spin fluctuations
- Exponents of quantities controlled by bulk fluctuations, suppressed by  $1/N \Rightarrow$  mean-field
- Exponents of stringy quantities  $\Rightarrow$ , controlled by  $\alpha' \Rightarrow$  beyond mean-field possible

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THANK YOU !

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- Then condition **ii)** is automatic.

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$$V(\Phi) \rightarrow V_\infty e^{2\sqrt{\frac{\xi}{d-1}}\Phi} (1 + V_{sub}(\Phi)), \quad \Phi \rightarrow \infty$$