

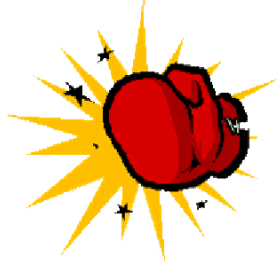
Exposing Nuclear Burning During Radius Expansion X-ray Bursts

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Outline

- I. Punch-line
- II. Brief overview of X-ray bursts
- III. Convection during X-ray bursts
- IV. Ash ejection by winds from super-Eddington bursts
- V. Composition of ejected ashes & detectability (Chandra, XMM, Con-X?)
- VI. Conclusions + Future plans



Punch-line

Ashes of nuclear burning are ejected in the **radiative wind** of photospheric radius expansion bursts and **exposed**

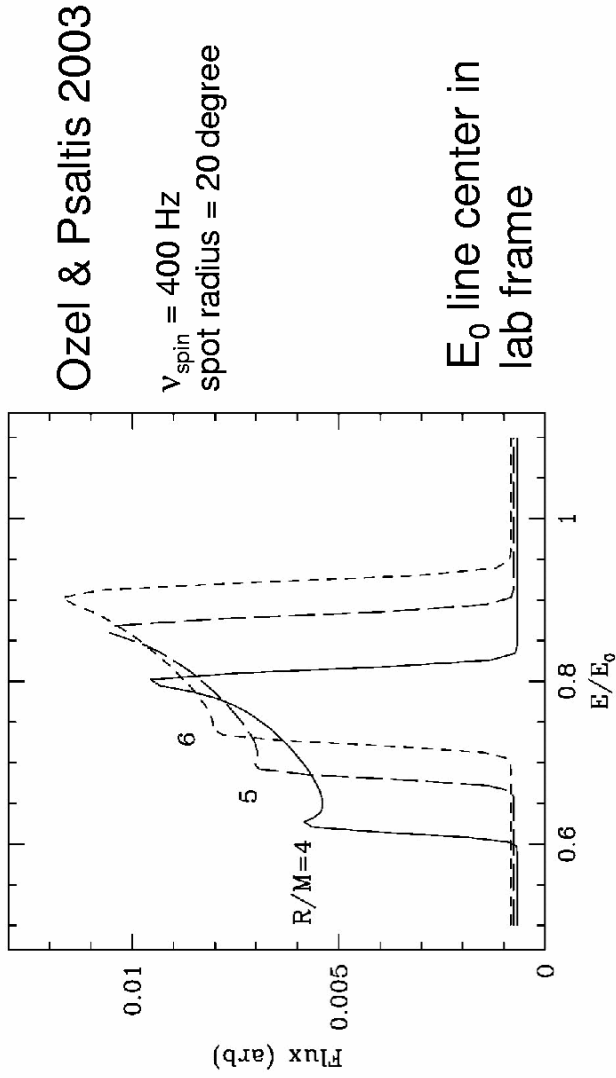
➡ Detect spectral lines of ashes

Our study motivated by numerical results from Woosley, Heger et al.:

➡ they showed ashes can be brought high up in atmosphere by convection.

Why is this interesting?

- By detecting spectral lines from ashes:
 1. Can measure gravitational redshift
 $z = \Delta\lambda / \lambda$ if lines formed near neutron star surface...measure NS **equation of state**.
 2. Probe **nuclear processing** even if lines not formed near surface

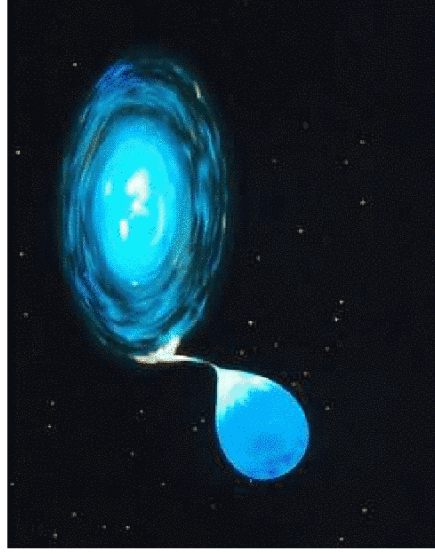


Cottam, Paerels, & Mendez 2002 report detection of redshifted H-like and He-like Fe absorption lines from bursts in EXO 0748-676... $z \sim 0.35$

Fe from accretion material **not** ashes of burning

Overview of X-Ray Bursts

Low-Mass X-ray Binary



Neutron star: $M \sim 1.3 - 2 M_{\text{sun}}$

Donor star: two possible types

1. $\sim 0.5 M_{\text{sun}}$ main sequence star... H/He accretion OR
2. He white dwarf or He core red giant... He accretion

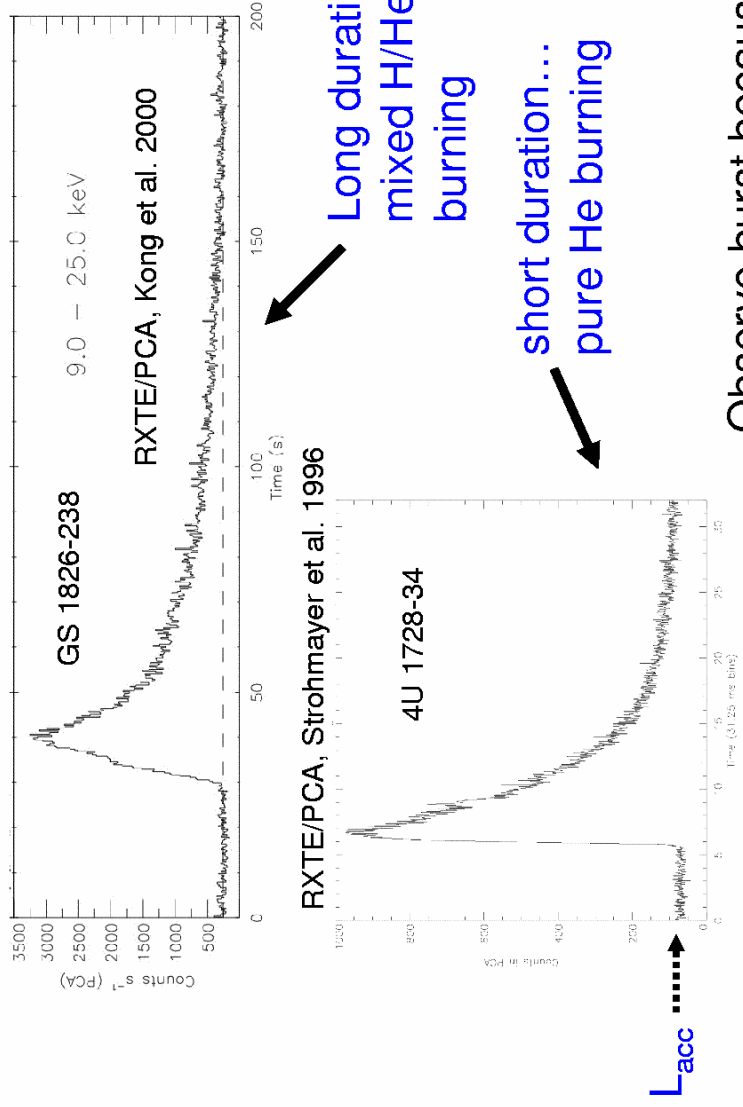
$P_{\text{orb}} \sim \text{hour to days}$

$$\dot{M} \sim 10^{-10} - 10^{-8} M_{\text{sun}} \text{ yr}^{-1}$$

$$L_{\text{acc}} = \frac{GM_{\text{NS}} \dot{M}}{R_{\text{NS}}} \sim 10^{36} - 10^{38} \text{ erg s}^{-1}$$

$kT \sim \text{keV}$
X-ray emitter





Long duration...
mixed H/He
burning

short duration...
pure He burning

Observe burst because
nuclear fuel accumulated for
hrs - days, burned rapidly in
seconds - minutes

$E_G \sim 200$ MeV/nucleon
 $E_{nuc} \sim$ few MeV/nucleon

NS Atmosphere during an X-ray burst

Accreting matter $v \sim c/3$

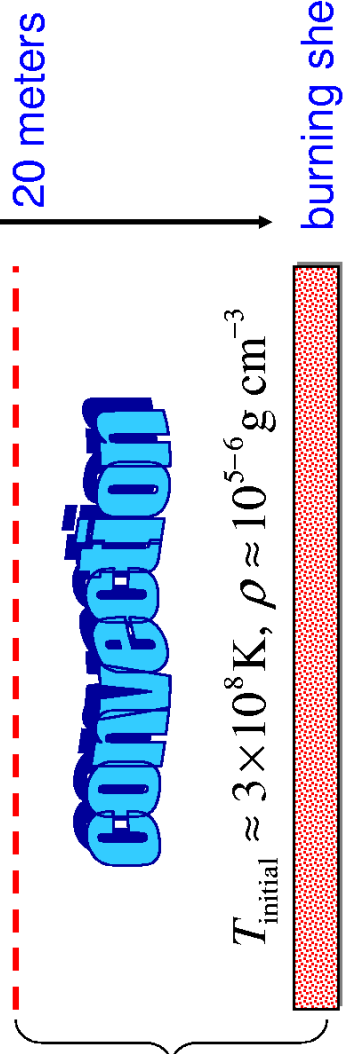
photosphere $T \approx 10^7$ K, $\rho \approx 1$ g cm⁻³

radiative layer

convection

$T_{initial} \approx 3 \times 10^8$ K, $\rho \approx 10^{5-6}$ g cm⁻³

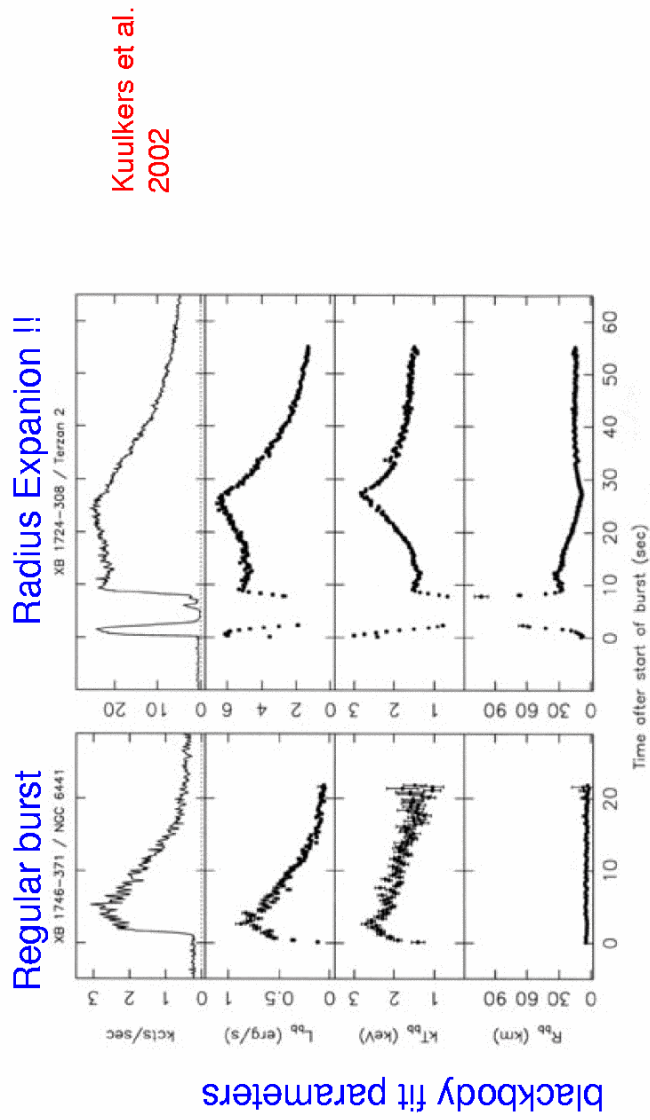
Mixed
with
ashes



Burst properties sensitive to $\dot{M} = \frac{\dot{M}}{4\pi R^2}$

- Hot CNO burning during fuel accumulation... **H**
He →
- $t_{\text{burn H}} \sim 1 \text{ day}$ $t_{\text{ignition mass}} \sim \text{hr}$ ($\dot{m}/\dot{m}_{\text{Edd}}$)
- If $\dot{m} < \dot{m}_{\text{Edd}}/24$ → pure He layer in burning shell ($Y_{\text{he}} \sim 1$)
- Since $\epsilon_{3\alpha} \propto (Y_{\text{he}})^3$ low accretion rate yield more energetic bursts

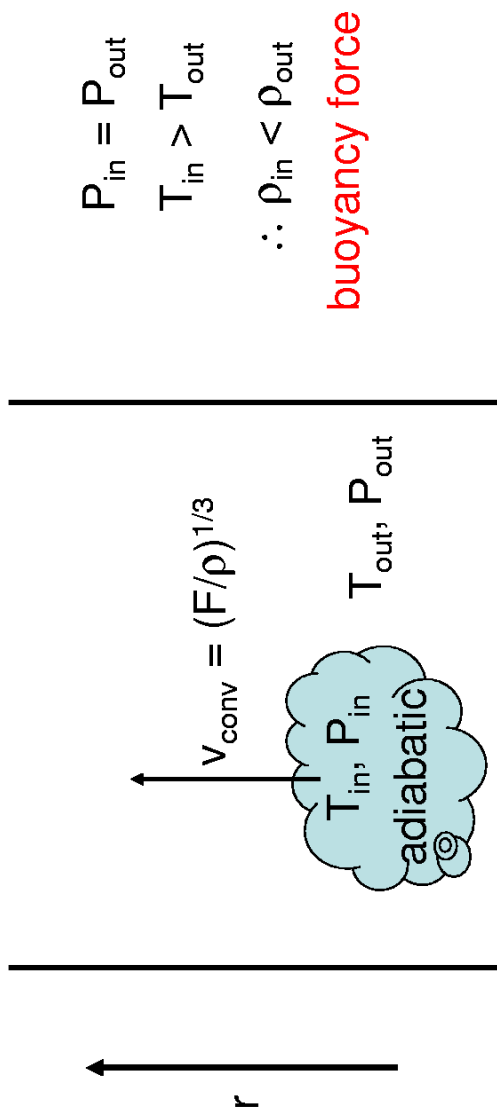
$\dot{m} < \dot{m}_{\text{Edd}}/24$ → **Super-Eddington luminosity** which drives a radiative wind (photospheric **radius expansion**)



$$M_{\text{wind}}/M_{\text{accreted}} \sim E_{\text{nuclear}}/E_{\text{gravity}} \sim 2 / 200 \sim 1\%$$

We show that convection reaches low enough pressures that **wind from RE burst** can **eject ashes**.

Convection during X-ray Bursts



Convectively unstable if $\left| \left(\frac{dT}{dr} \right)_{in} \right| < \left| \left(\frac{dT}{dr} \right)_{out} \right|$

Convection during X-ray Bursts

- Why is atmosphere convective?

$$t_{\text{thermal}}(\text{at base}) = C_P T_b y_b / F = (10^{8+8+8} / 10^{24}) \sim 1\text{s}$$

$$t_{\text{dynamical}}(\text{at base}) = H_P / C_s = (y_b / \rho_b g)^{1/2} = (10^8 / 10^{6+14})^{1/2} \sim 10^{-6} \text{ s}$$

where $y = P / g =$ column depth (g cm^{-2})

- Convection subsonic:

$$V_{\text{conv}} \sim (F / \rho)^{1/3} = (10^{24} / 10^6)^{1/3} = 10^6 \text{ cm/s}$$

$$C_s = (g y / \rho)^{1/2} = (10^{14+8} / 10^6)^{1/2} = 10^8 \text{ cm/s}$$

Thermal profile in convective region very nearly follows an adiabat: $T \propto y^n$, $n \sim 2/5$

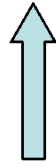
The extent of the convective region

- Upper bound: Joss 1977

$$S \propto \ln(T^{5/2} / P) = \ln(T^{5/2} / g y)$$

Photosphere: ($[10^7]^{5/2} / 10^{14+0}$) ~ 3000

Burning shell: ($[10^9]^{5/2} / 10^{14+8}$) ~ 3



Entropy barrier of photosphere prevents convection from reaching surface

The extent of the convective region

- **Order of Magnitude estimate:** find where entropy of radiative region equals entropy of convective region

$$S \propto \ln(T^{5/2} / P) = \ln(T^{5/2} / g y) \\ = \ln(3) \text{ in convective region}$$

In **radiative region** temperature profile follows

$$T(y) \sim T_{\text{surface}} (y_{\text{surface}} / y)^{1/4}$$

So extent of convective region y_c :

$$3 \sim T_{\text{surface}}^{5/2} (y_s / y_c)^{5/8} / g y_c$$

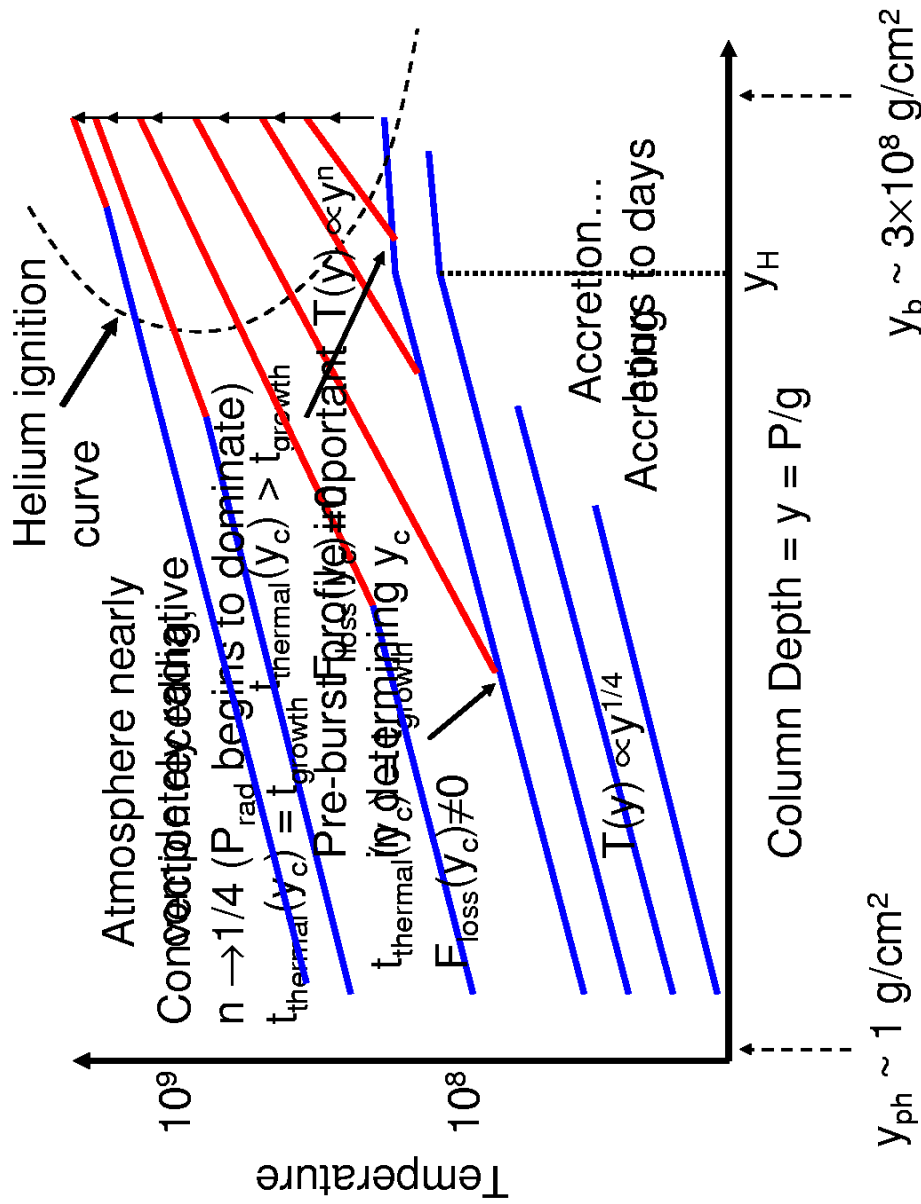
$$y_c = (T_{\text{surface}} / 3)^{8/13} y_{\text{surface}}^{5/13} \\ = (10^7 / 3)^{8/13} = 10^4 \text{ g/cm}^2$$

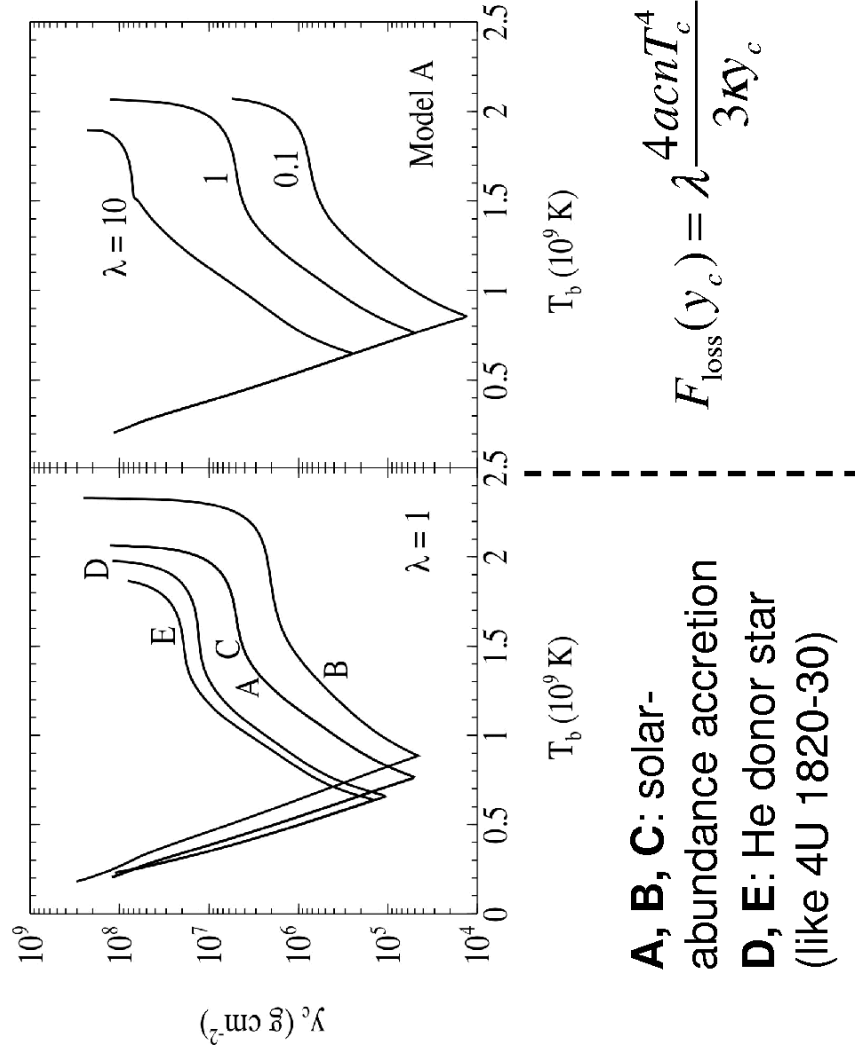
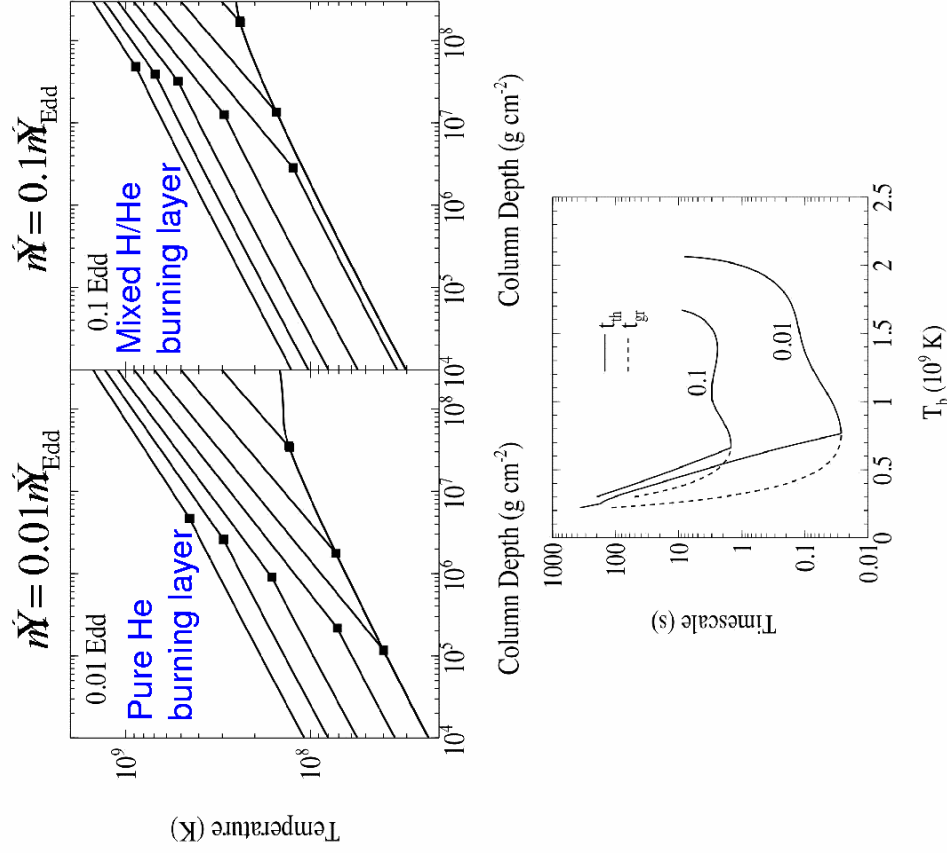
Pretty good estimate... but ignores heating up of radiative region, etc., and doesn't give evolution

The extent of the convective region

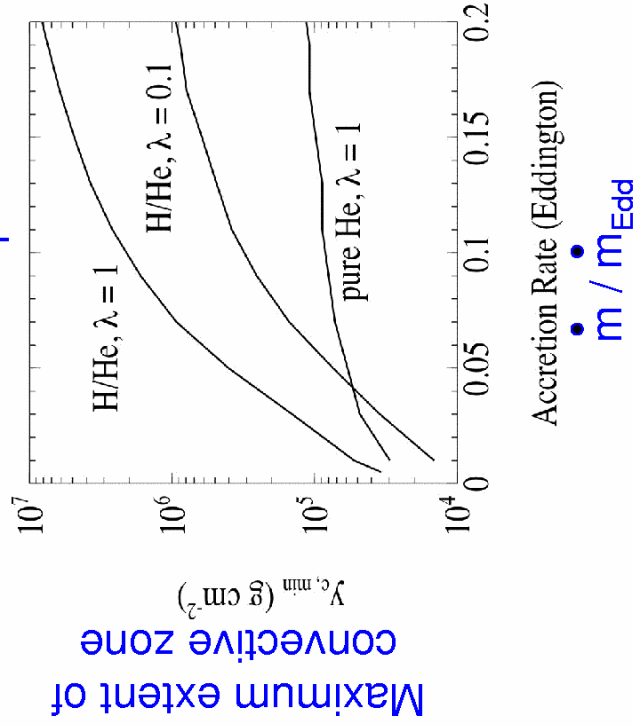
- Solve the **entropy equation**: $T \frac{dS}{dt} = \frac{dF}{dy} + \epsilon_{\text{nuclear}}$
 - Convective region **adiabatic**: $T(y) = T_b (y/y_b)^{n(y)}$
 - Use $T \frac{dS}{dt} = C_P \frac{dT}{dt}$ and integrate over convective region:
- $$\frac{dT_b}{dt} = \frac{\int_{y_c}^{y_b} \epsilon_{\text{nuclear}} dy + F_{\text{crust}} - F_{\text{loss}}(y_c)}{\int_{y_c}^{y_b} C_P (y/y_b)^{n(y)} dy}$$
- Similar to calculation in Hanawa & Fujimoto 1984

F_{crust} = flux escaping from crust ~ 0.1 MeV/nucleon
 $F_{\text{loss}}(y_c)$ = flux lost to overlying radiative region
 (get from mixing length theory)





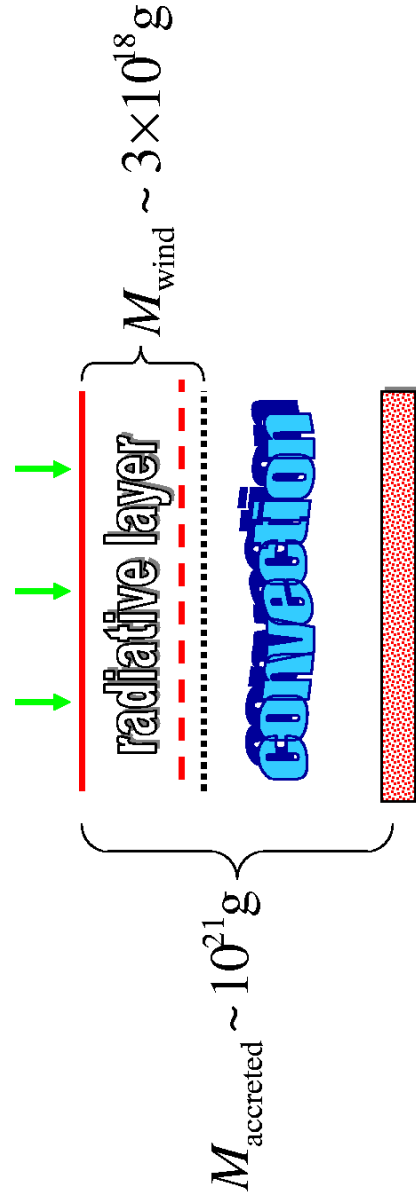
The BIG picture



$Y_{c, \min}/Y_b$ = fraction of mass above convective zone
 = 0.01% at 0.01 Edd
 = 1% at 0.1 Edd

Winds During Radius Expansion Bursts

- The **ashes** mixed throughout convective region **do not reach photosphere** (entropy barrier).
- But do the ashes reach **low enough pressures** at an **early enough times** (before $L > L_{\text{Edd}}$) to get ejected by the wind? $M_{\text{wind}}/M_{\text{accreted}} \sim E_{\text{nuclear}}/E_{\text{gravity}} \sim 1\%$



Estimate of $M_{\text{wind}} = M_{\text{acc}} Y_w / Y_b$

Suppose $L > L_{\text{edd}} = 4\pi GMc/\kappa$

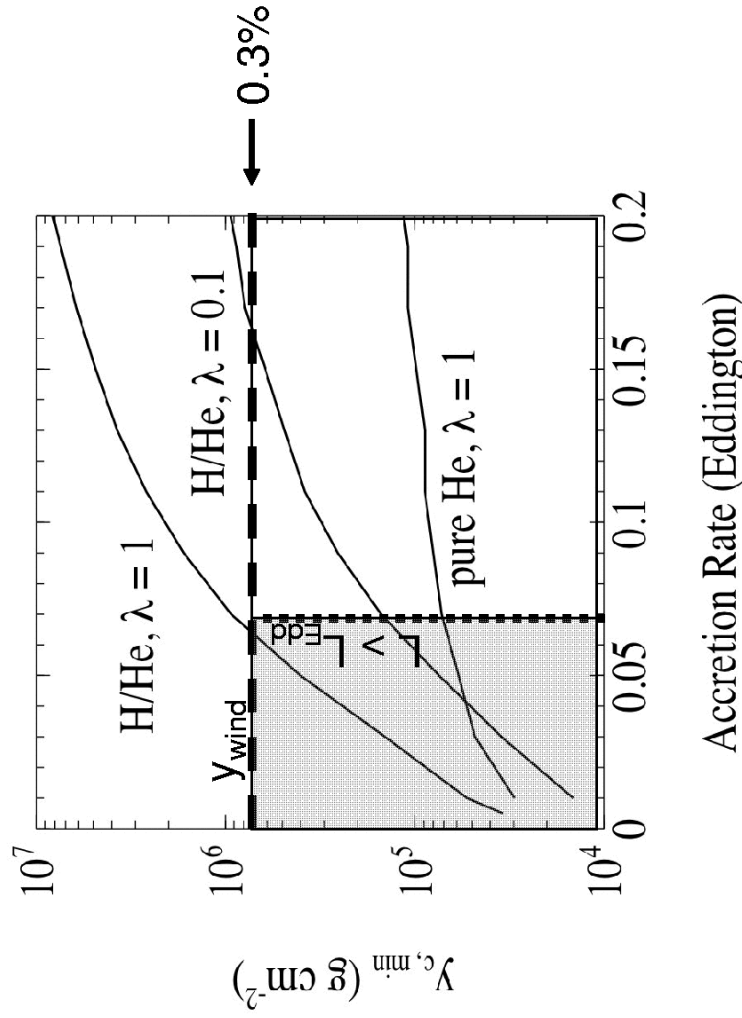
Force = $(L - L_{\text{edd}})/c \approx v_{\text{esc}} dM_{\text{wind}}/dt$

$$dM_{\text{wind}}/dt \approx 10^{18} \text{ g/s } (0.2 / \kappa) (L/L_{\text{edd}} - 1)$$

Burst duration: $\Delta t \sim M_{\text{acc}} Q_{\text{nuc}}/L$,

where $Q_{\text{nuc}} = 1.6 + 4 X_{\text{H}} \text{ MeV/nucleon}$

$$\begin{aligned} \Delta M_{\text{wind}}/M_{\text{acc}} &\approx \Delta t (dM_{\text{wind}}/dt) / M_{\text{acc}} \\ &= (1 - L_{\text{edd}}/L) Q_{\text{nuc}} / v_{\text{esc}} c \\ &= 0.003 (1 - L_{\text{edd}}/L) \end{aligned}$$



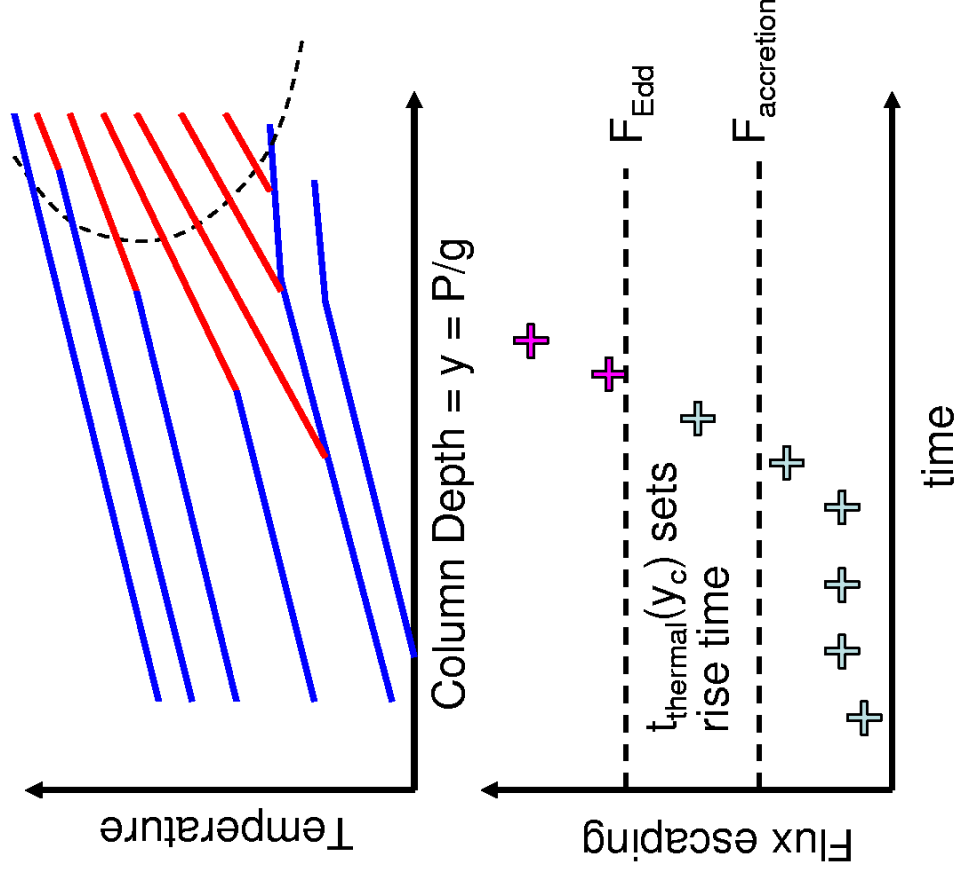
Timescale Question:

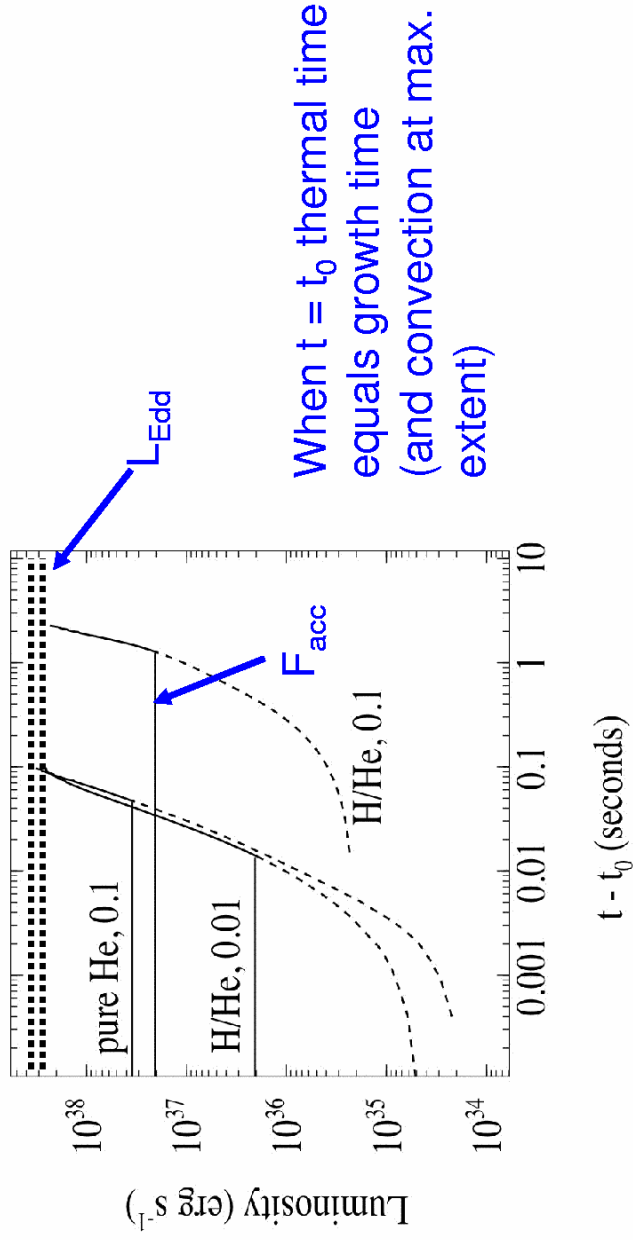
Does wind turn on after $y_c < y_{\text{wind}}$?

I.e., When $y_c < y_{\text{wind}}$ is $L < L_{\text{edd}}$?



Must calculate rise times of burst light curves



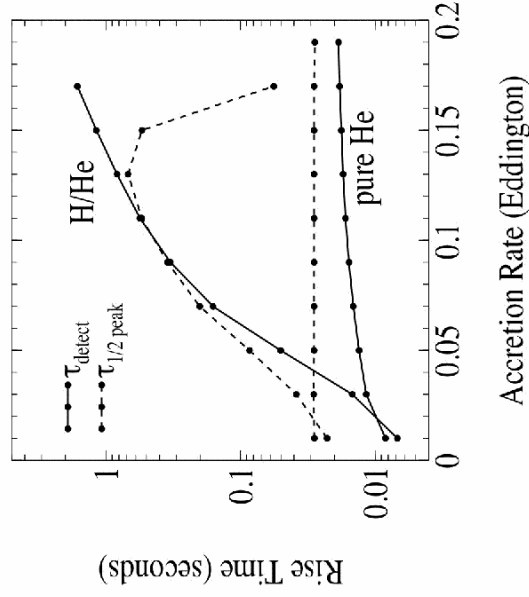


Reach $L > L_{\text{edd}}$ well after $y_c < y_{\text{wind}}$... so ashes ejected.

burst rise times

Since photons diffuse out on the thermal timescale $t_{\text{thermal}}(y_{c,\text{min}})$ sets the timescale of the burst rise

↑ The extent of convection sets the burst rise time



Ejected Ash Composition & Detectability

- Our current calculations of $\epsilon_{\text{nuclear}}$ **only account for α -captures**...reasonable if only have helium around
- Do not burn to very heavy metals
- Can get spectral lines during cooling in radius expansion
- Difficult to get lines at surface...no redshift (?)
- **New calculations** with Hendrik Schatz's full nuclear reaction network suggest there may be protons around even in initially pure He layer...changes $\epsilon_{\text{nuclear}}$ and ash composition

Ejected Ash Composition & Detectability

Where might protons come from?

The reaction $^{12}\text{C}(\text{a,g})^{16}\text{O}$ is a bottleneck

If have protons around can overcome bottleneck with the reaction $^{12}\text{C}(\text{p,g})^{13}\text{N}(\text{a,p})^{16}\text{O}$

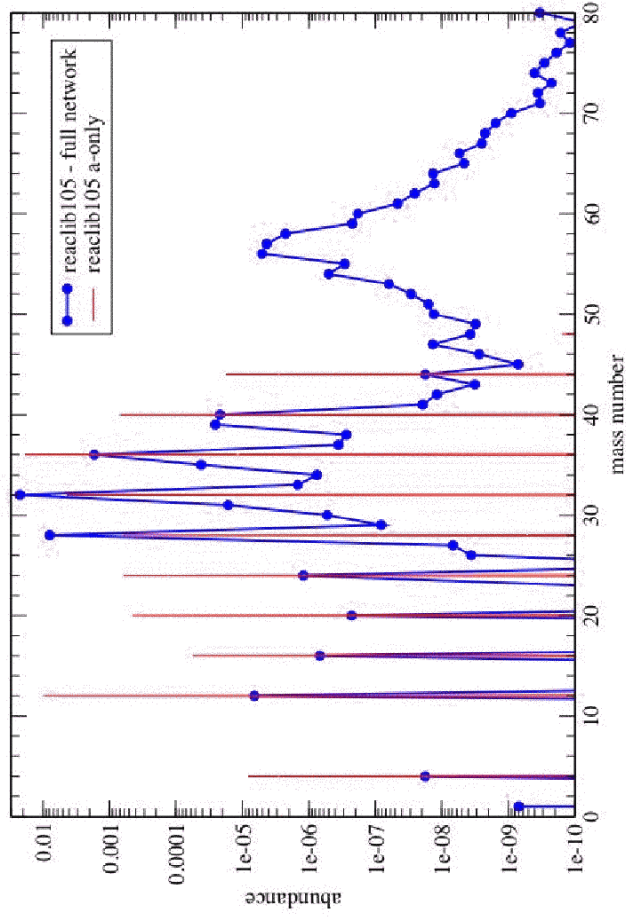
Where might protons come from if have initially pure He layer?

E.g., $^{24}\text{Mg}(\text{a,p})^{27}\text{Al}(\text{a,p})$ occurs if burning hot enough to reach ^{24}Mg but not so hot that overcome $^{12}\text{C}(\text{a,g})$ bottleneck

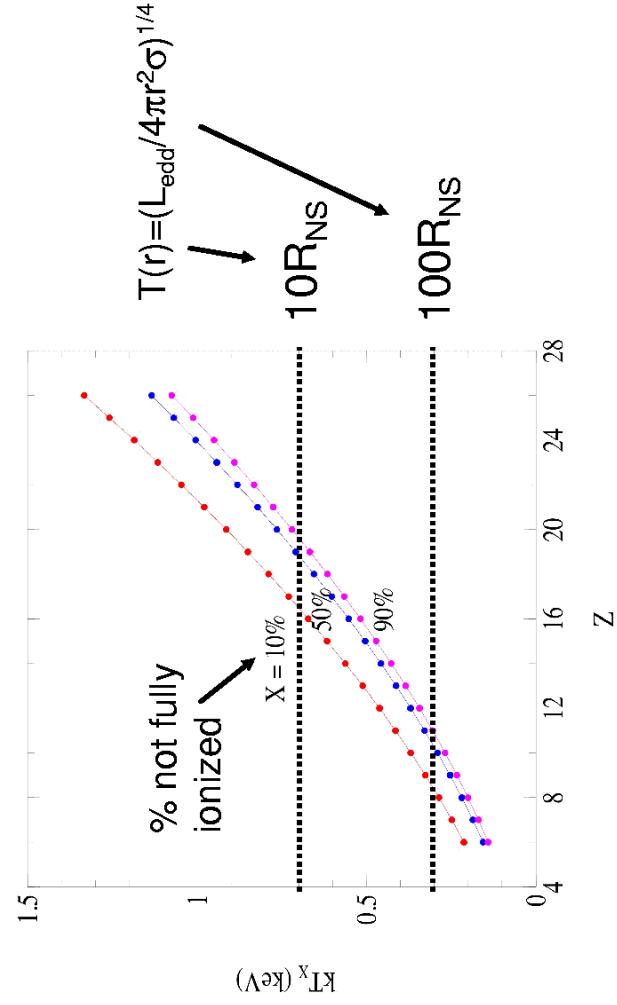


Will there be heavy metals even in low accretion rate, radius expansion bursts?
Redshift measurement?

Ejected Ash Composition & Detectability



Saha Ionization Equation:



Conclusions & Future Work

- Ashes of nuclear burning ejected during radius expansion bursts
- Detection of atomic spectral features would probe the neutron star equation of state and the nuclear burning
- Hope to soon have a better idea of the ash composition and address the detectability of lines with e.g., XMM, Astro-E2, Chandra, Constellation-X
- Putting together a proposal with D. Galloway for Chandra time