Stars, Disks and Massive Black Holes

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Gravitational vs hydrodynamical star-disk interactions



Gravitational torques are typically much stronger than ram pressure torques.

$$\left.\frac{\Delta J_{\rm ram}}{\Delta J_{\rm grav}}\right|_{\rm per \, period} \sim \left(\frac{M_{\bullet}}{M_{\star}}\right) \left(\frac{R_{\star}}{R}\right)^2$$

• Compact disks (luminous active galactic nuclei / Quasars, $R \sim 2000 r_g$):

$$\left.\frac{\Delta J_{\rm ram}}{\Delta J_{\rm grav}}\right|_{\rm per\,period} \sim 0.02 \left(\frac{M_{\bullet}}{10^6 M_{\star}}\right)^{-1}$$

• Gravitational torques significant compared to internal viscous torques ($\alpha < 1$):

$$\frac{\Delta J_{\text{grav}}}{\Delta J_{\text{visc}}}\Big|_{\text{per period}} \sim \left\{ \begin{array}{c} \alpha \\ 1/\alpha \end{array} \right\} \sqrt{N_{\star}} \left(\frac{H}{R} \right)^{-2} \left(\frac{M_{\bullet}}{M_{\star}} \right)^{-1} \sim \left\{ \begin{array}{c} 0.1 - 1 & \perp \text{ (unwarping)} \\ 1 - 10 & \parallel \text{ (inflow)} \end{array} \right\}$$

Outline

- How do stellar torques affect the disk?
 - Warped accretion disks
 - Resonant relaxation
 - Proof of concept: The maser disk of NGC 4258
 - Disk warping
 - Covering factor
 - Accretion rate
 - MBH spin jitter
- What are the consequences for MBH evolution? (work in progress, open questions)
 - Statistical properties of resonant torques.

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- Stability of thin disks
- Scaling with MBH





The randomization of stellar systems: Resonant Relaxation



N-body simulations (Eilon, Kupi & Alexander 2009)



GC dynamics (Hopman & Alexander 2006)



Maser disk warp (Bregman & Alexander 2009, 2012)

The warped maser disk in NGC 4258



(Herrnstein et al 1996)

Radio measurements of x, v, \dot{v}

 $M_{\bullet} = 3.7 \times 10^7 M_{\odot}, \quad D = 7.2 \text{ Mpc}, \quad R_{\text{in}} = 0.14, R_{\text{out}} = 0.28 \text{ pc}, \quad \omega \sim 8^{\circ}$

 $\left(\Delta J_{\rm ram}/\Delta J_{\rm RR} \sim 10^{-7}, \ \Delta J_{\rm grav}/\Delta J_{\rm visc,\perp} \sim 0.25, \ \Delta J_{\rm grav}/\Delta J_{\rm visc,\parallel} \sim 4\right)$

RR Disk warping



(Bregman & Alexander 2012)

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RR-induced warping of NGC 4258: the covering factor



External heating and ionization thought crucial for explaining disk properties.

Image: Image:

RR-warping exposes disk to the ionizing central source.

RR-induced warps and the mass accretion rate



RR-warping increases \dot{M} by up to imes4, $\left<\dot{M}\right>$ by up to imes1.5.

- Eddington-limited growth is exponential.
- RR can accelerate MBH mass and spin growth by ×100 over disk's lifetime. Alternatively, enhanced mass flow can lead to gravitational instability, fragmentation and star formation.

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Disk and MBH spin jitter



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- MBH-stars angular momentum coupling via the disk (The Bardeen-Petterson frame-dragging effect)
- Implications for radio jet opening angle and orientation?

RR warping of disks Open questions: Stability and scaling relations





The goal: Full 3D statistical description of RR (spherical harmonics in 4D) Coherent and diffusive limits, amplitudes, correlation lengths and timescales

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- When is a disk stable under RR torquing? M_•, M_{*}, N_{*}, M_{disk}, T_{disk}, v_{1,2},...
- Warp angle rms: $\omega_{RR} \propto \sqrt{t_{\rm coh}/t_{\nu_2}} \propto M_{\bullet}^{-3/7}$
- RR warping stronger for growing MBHs in early universe.
- What came first, the MBH or the stars?

Summary

- Dense stellar clusters around an accreting MBH can strongly affect the disk by fast stochastic gravitational torques.
- Demonstration of concept: RR-induced warping can explain the observed warp of NGC 4258 maser disk.
- RR-induced warping increases covering factor, accelerates the accretion of mass and angular momentum, and changes the MBH spin direction.

Work in progress:

- Disk stability under RR.
- Scaling properties with MBH mass, disk properties: RR stronger for lower-mass MBHs.

Implications for cosmic evolution of MBHs.

The disk evolution equation

Angular momentum density : $L(R) = \Sigma(R) \sqrt{GM_{\bullet}R}\ell(R)$

$$\frac{\partial \mathbf{L}}{\partial t} = \underbrace{\frac{1}{R} \frac{\partial}{\partial R} \left[3R \frac{\partial}{\partial R} (v_1 L) \ell + \frac{1}{2} v_2 R L \frac{\partial \ell}{\partial R} \right]}_{\text{Diffusive accretion of ang. mom. into the MBH}} \\ + \underbrace{\frac{1}{R} \frac{\partial}{\partial R} \left[\left(v_2 R^2 \left| \frac{\partial \ell}{\partial R} \right|^2 - \frac{3}{2} v_1 \right) L \right]}_{\text{Advective current of ang. mom. outward}} \\ + \underbrace{\mathbf{T}_{\text{BP}}(R) + \mathbf{T}_{\text{RR}}(R) + \mathbf{T}_{\text{src}}(R)}_{\text{External torques}}$$

Stochastic torques : $T_{RR}(R, t; M_{\bullet}, N)$