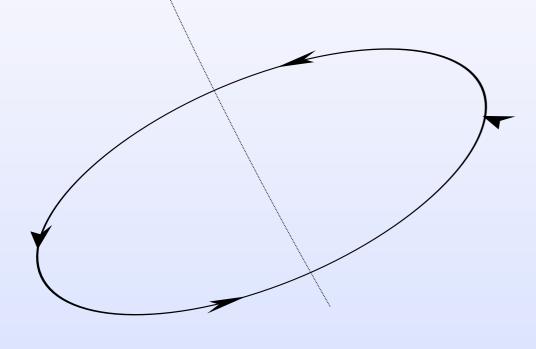


Global Simulations of Magnetised Accretion Disk Turbulence

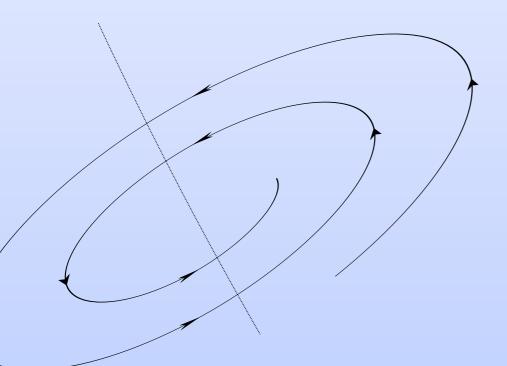
Ross Parkin & Geoff Bicknell

Rule number one



No angular momentum loss

= test particle orbits forever

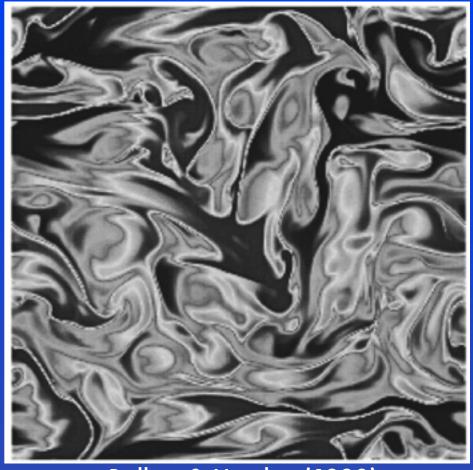


Angular momentum loss

= test particle falls on to central object

Magnetorotational turbulence

- Angular momentum transport can be facilitated by magnetorotational instability driven turbulence
- Sustained turbulence demonstrated in unstratified shearing boxes (Hawley+ 1995)
- However, stress decreases with increasing resolution in zero-net flux simulations (Fromang & Papaloizou 2007)

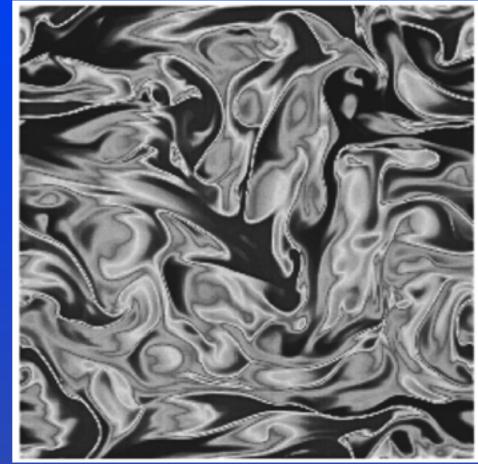


Balbus & Hawley (1998)



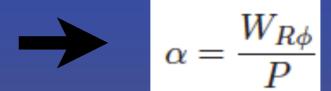
Magnetorotational turbulence

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Balbus & Hawley (1998)

$$W_{R\phi} = \rho \delta v_{\rm R} \delta v_{\phi} - B_{\rm R} B_{\phi}$$



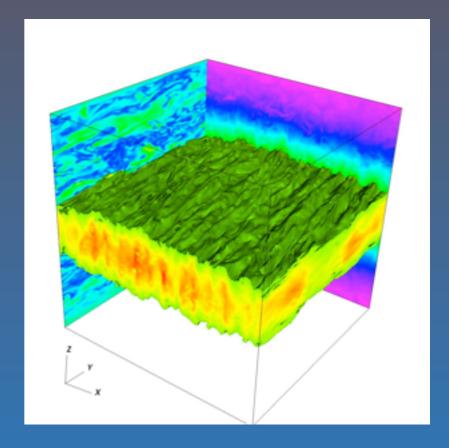
Model	Box size	Resolution	Run time (in orbits)	$\alpha_{ m Rey}$	$\alpha_{ m Max}$	α
FS64	$(H, 2\pi H, H)$	(64, 200, 64)	300	1.8×10^{-3}	4.2×10^{-3}	5.9×10^{-3}
STD64	$(H, \pi H, H)$	(64, 100, 64)	1000	9.4×10^{-4}	3.2×10^{-3}	4.1×10^{-3}
STD128	$(H, \pi H, H)$	(128, 200, 128)	250	5.0×10^{-4}	1.7×10^{-3}	2.2×10^{-3}
STD256	$(H, \pi H, H)$	(256, 400, 256)	105	2.4×10^{-4}	8.1×10^{-4}	1.1×10^{-3}



Stratified models show convergence

 Sustained turbulence demonstrated in stratified shearing boxes (e.g. Brandenburg+1995, Stone+1996, Davis+ 2010, Shi+2010, Simon +2011, Oishi & MacLow 2011)

Results converge with increasing resolution when the vertical component of gravity is included (ie stratification).

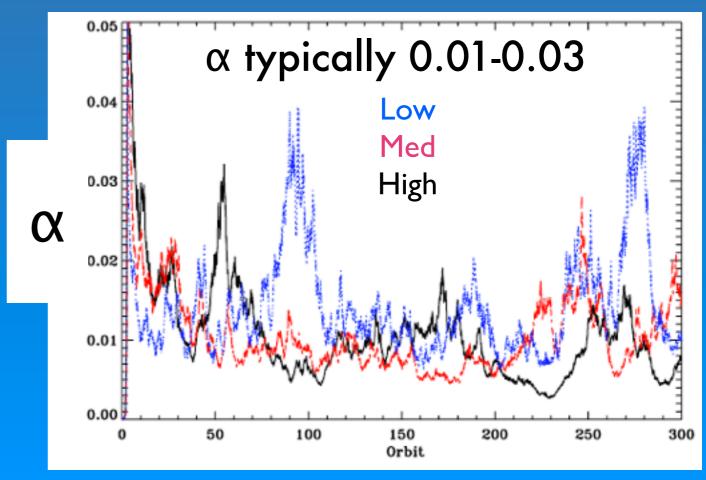


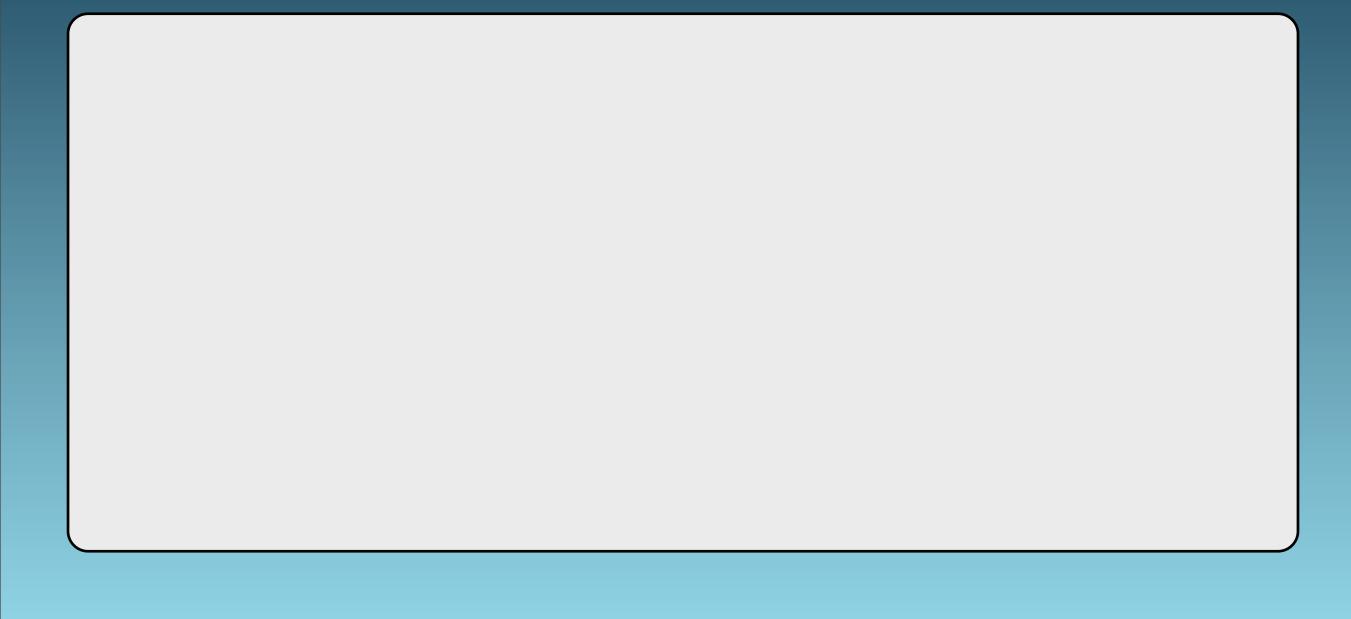
Davis+ (2010)

$$W_{R\phi} = \rho \delta v_{\rm R} \delta v_{\phi} - B_{\rm R} B_{\phi}$$

$$\alpha = \frac{W_{R\phi}}{P}$$









Will global stratified disk models converge?



- Will global stratified disk models converge?
- If so, at what resolution, and with what value for the turbulent stresses (α) ?

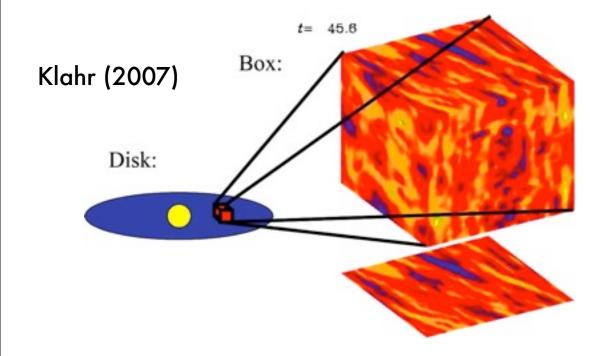


- Will global stratified disk models converge?
- If so, at what resolution, and with what value for the turbulent stresses (α)?
- How do global models and shearing-boxes differ in terms of magnetic energy production and large scale field development?





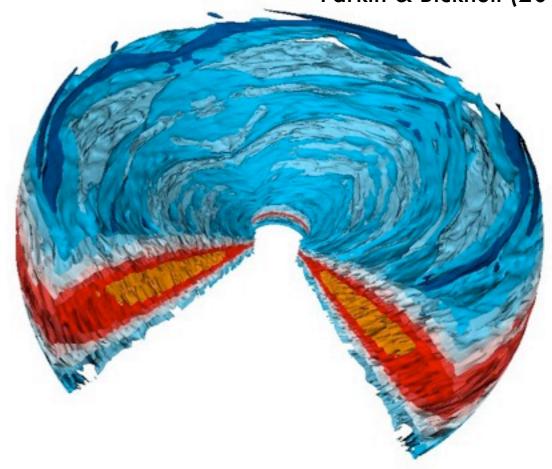
Shearing-box



See also: Brandenburg et al. 1995; Stone et al. 1996; Miller & Stone 2000; Fleming et al. 2000; Brandenburg 2005; Johansen et al. 2009; Gressel 2010; Shi et al. 2010; Davis et al. 2010; Simon et al. 2011; Guan & Gammie 2011; Oishi & Mac Low 2011; Simon et al. 2012

Global simulations

Parkin & Bicknell (2013a,b)

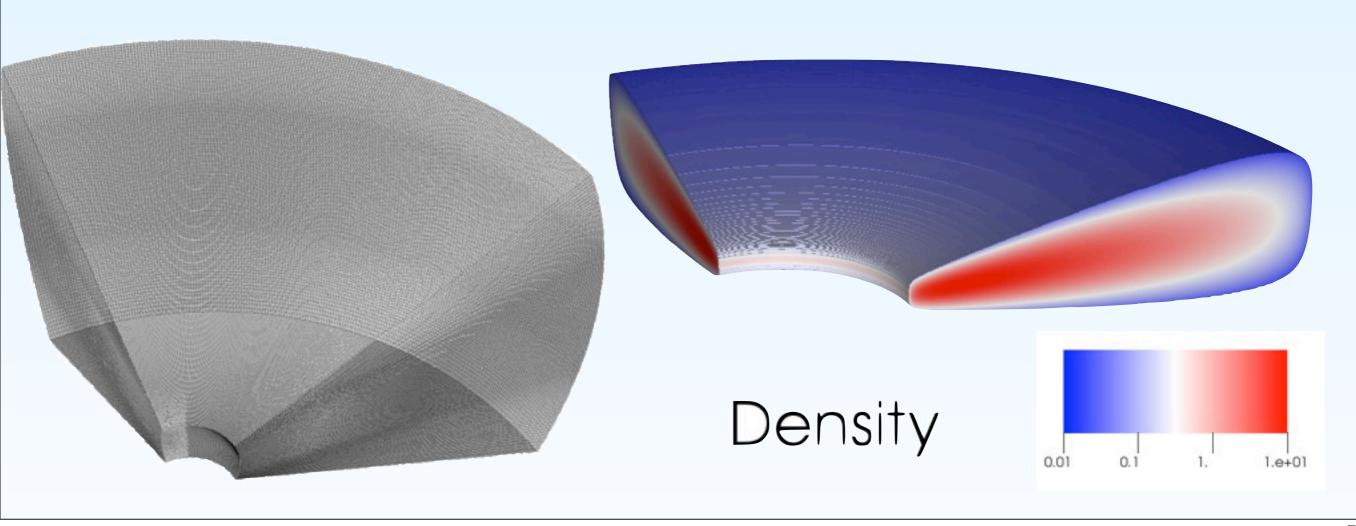


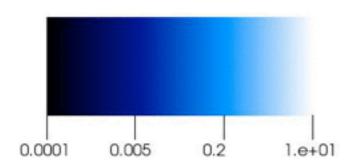
See also: Hawley 2000; Hawley & Krolik 2001; Arlt & Rudiger 2001; Fromang & Nelson 2006, 2009, Fragile et al. 2007, 2009; Beckwith et al. 2008, 2011; Lyra et al. 2008; Reynolds & Fabian 2008; Sorathia et al. 2010; O'Neill et al. 2011; Flock et al. 2011, 2012a,b; Noble et al. 2009, 2010, 2011; Hawley et al. 2011, 2013; McKinney et al. 2012, Dexter et al. 2013

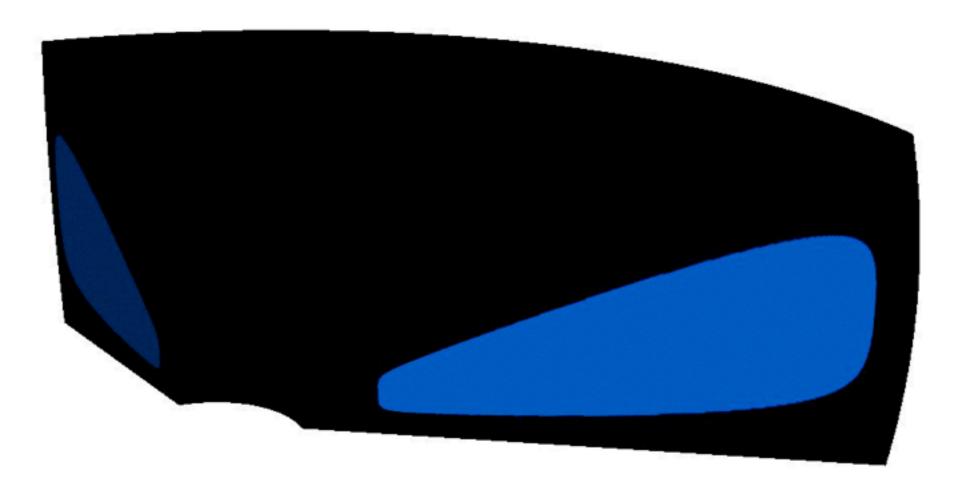
Global Simulations



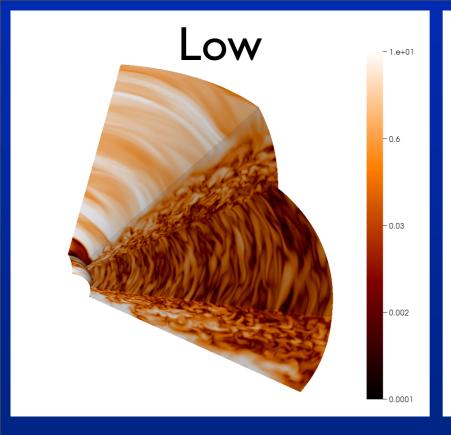
- · Equilibrium magnetized disk with purely toroidal field (Parkin & Bicknell 2013, Apj, 763, 99)
- No unphysical periodic boundaries
- · Performed on a 3D spherical grid using the PLUTO MHD code (Mignone+2007)

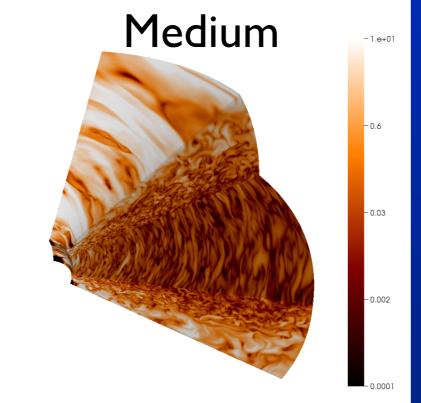


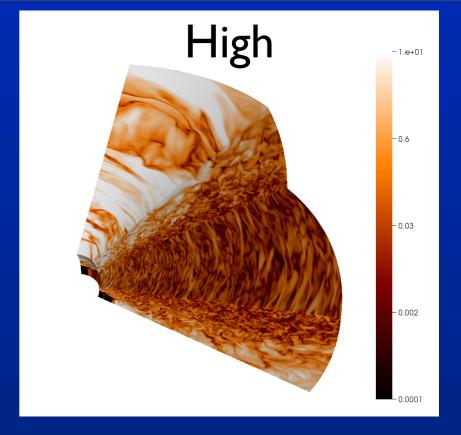




(Magnetic pressure)/(Gas pressure)







340×112×128 cells

512×170×196 cells

768×256×256 cells

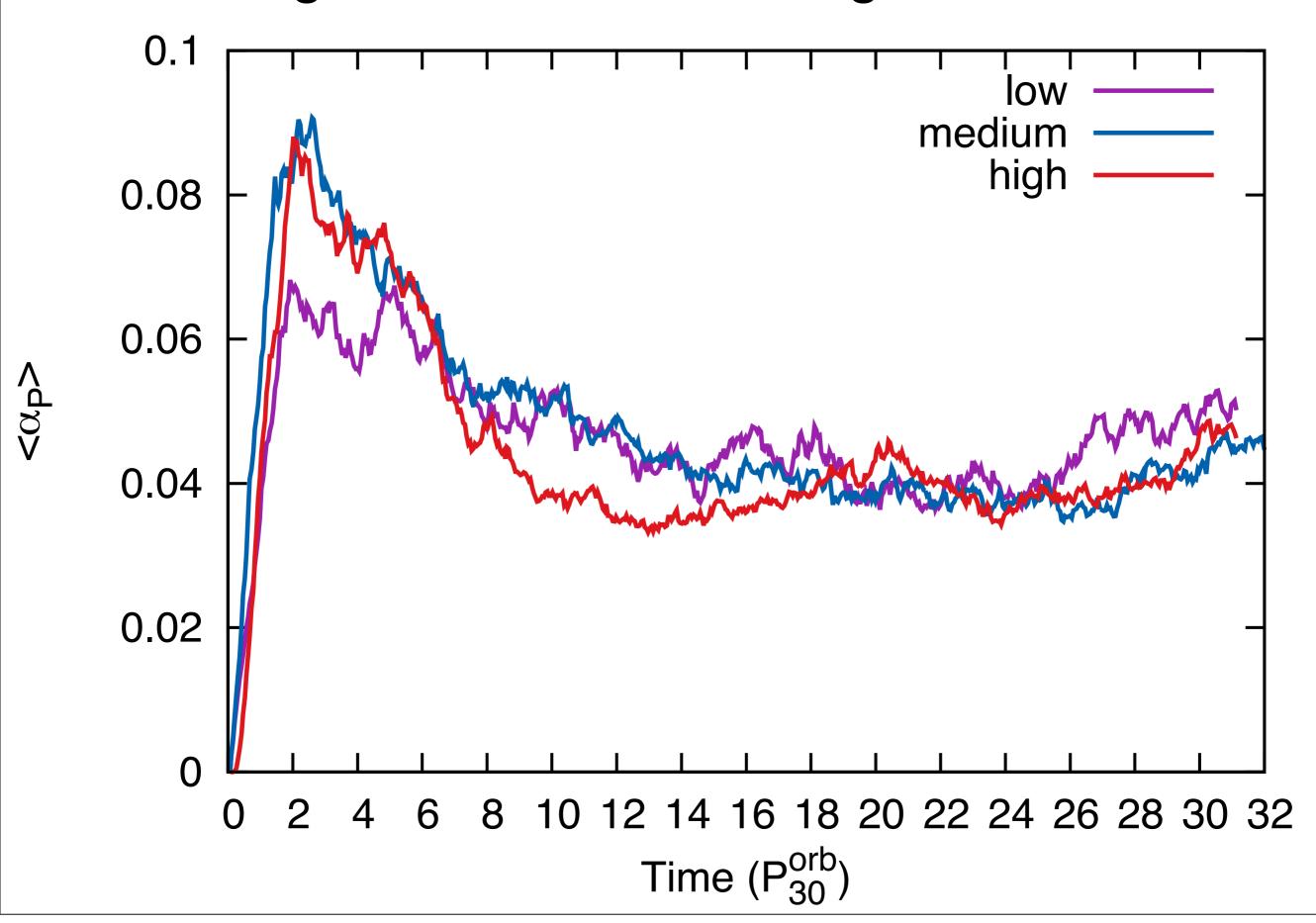
 $(n_r \times n_\theta \times n_\phi)$

Plots show (magnetic pressure)/
(gas pressure)

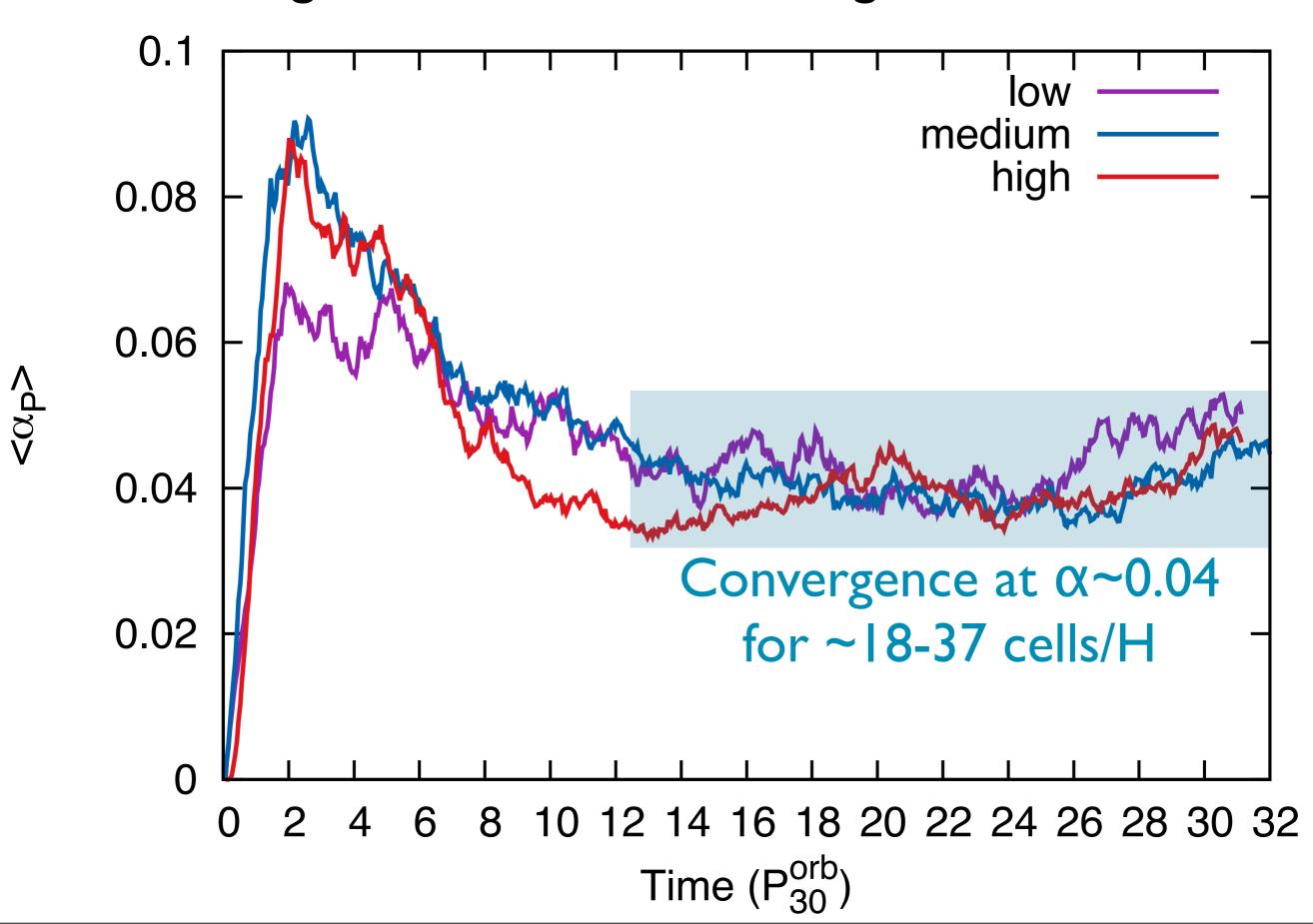
Model	Resolution (n/H) $(r \times \theta \times \phi)$			
low	$(9-36) \times 18 \times 8$			
medium	$(12-51) \times 27 \times 13$			
high	$(18-77) \times 37 \times 16$			



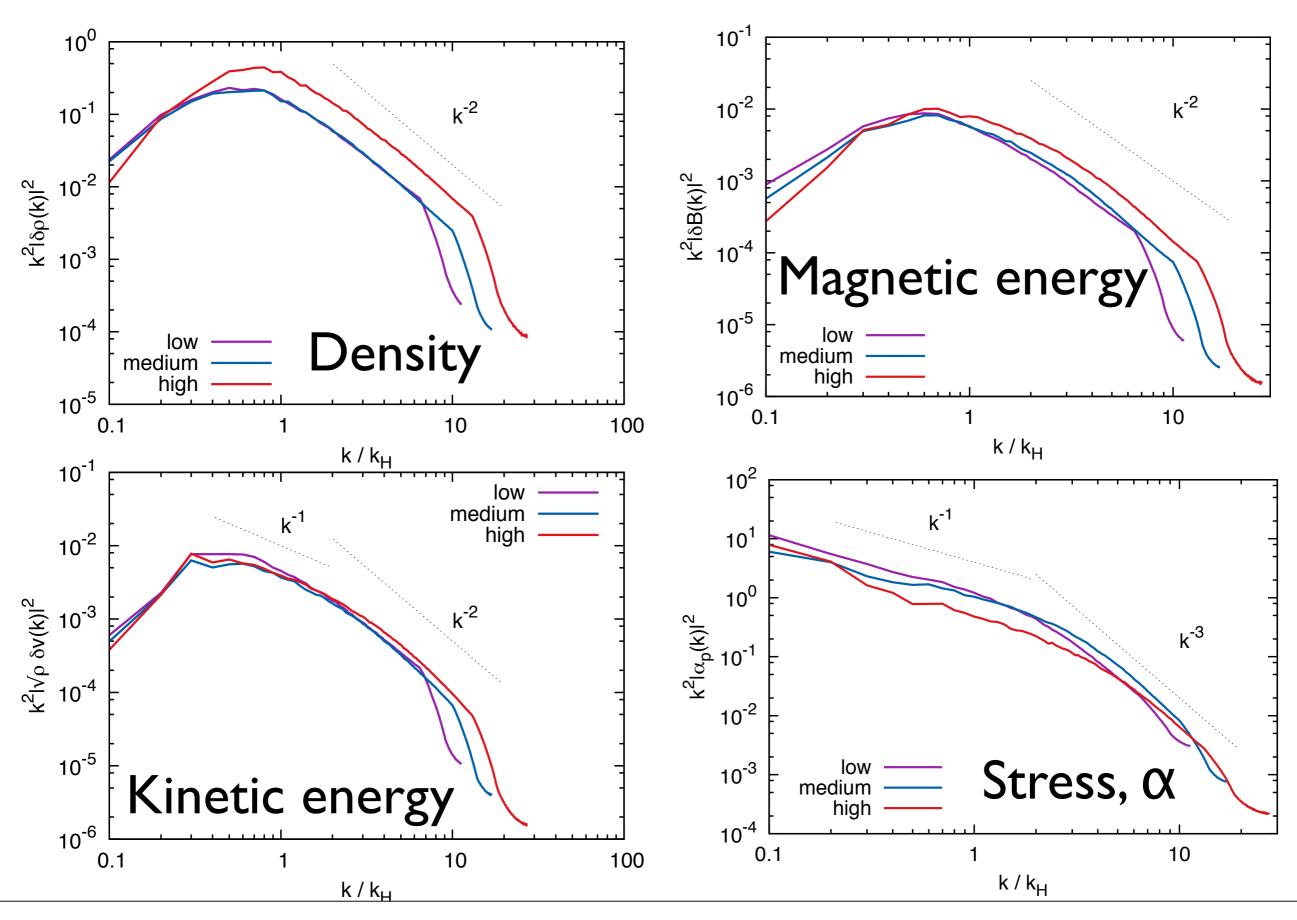
Convergence with resolution in global simulations



Convergence with resolution in global simulations

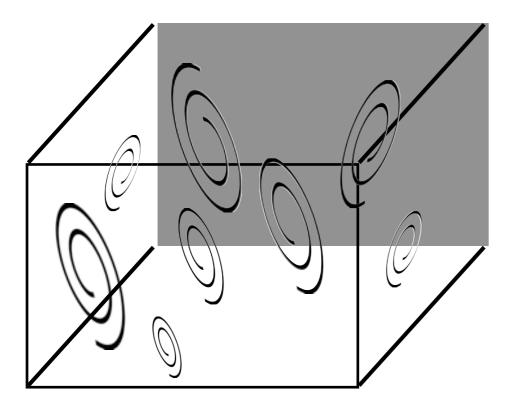


Power spectra



Magnetic energy maintenance

(Switch brain from real space to Fourier space)



Magnetic energy production must replenish energy on the largest scales.

Otherwise, the reservoir of energy at the largest scales will be drained by cascading of energy to smaller scales.







Convergence with increasing simulation resolution requires:



Convergence with increasing simulation resolution requires:

 Convergence in magnetic energy input/injection on large scales



Convergence with increasing simulation resolution requires:

- Convergence in magnetic energy input/injection on large scales
- Directly relates to the presence, or lack there-of, of magnetic field on large scales

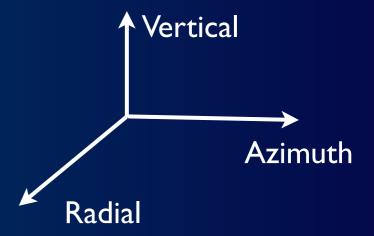


Convergence with increasing simulation resolution requires:

- Convergence in magnetic energy input/injection on large scales
- Directly relates to the presence, or lack there-of, of magnetic field on large scales
- Can assess these factors using a volume integrated analysis of the magnetic energy and induction equations with appropriate boundary conditions - see Parkin & Bicknell, 2013, arXiv:1306.1084

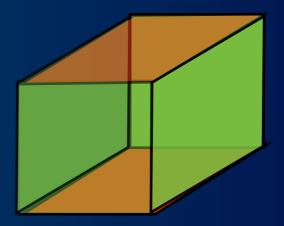
Boundary conditions for models

Periodic boundary Open boundary



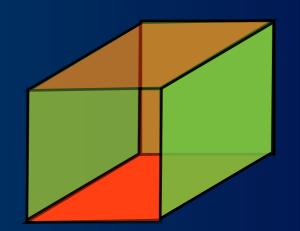
Stratified shearing box

(vertical surfaces are the disk-corona interface)



Global disk





Magnetic energy balance

The situation simplifies to,

$$\oint v_{
m i} M_{
m ij}^{
m B} n_{
m j} dS - \int v_{
m i} rac{\partial M_{
m ij}^{
m B}}{\partial x_{
m j}} \, dV pprox \eta_{
m num} \int |j|^2 dV.$$
 , where $M_{
m ij}^{
m B} = B_{
m i} B_{
m j} - \delta_{
m ij} u_{
m B}$

$$M_{
m ij}^{
m B}=B_{
m i}B_{
m j}-\delta_{
m ij}u_{
m B}$$
 (Maxwell stress tensor)

Applying appropriate boundary conditions and preserving dominant terms, we have,

Shearing box
$$-q\Omega L_{\rm x} \int_{{\rm x}_1} B_{\rm x} B_{\rm y} dS_{\rm x} - \int v_{\rm y} B_{\rm x} \frac{\partial B_{\rm y}}{\partial x} dV \approx \eta_{\rm num} \int |j|^2 dV.$$

$$-\int_{r_1} B_{\mathbf{r}} B_{\phi} v_{\phi} dS_{\mathbf{r}} - \int v_{\phi} B_{\mathbf{r}} \frac{\partial B_{\phi}}{\partial r} dV \simeq \eta_{\text{num}} \int |j|^2 dV.$$

Magnetic energy balance

The situation simplifies to,

$$\oint v_{\rm i} M_{\rm ij}^{\rm B} n_{\rm j} dS - \int v_{\rm i} \frac{\partial M_{\rm ij}^{\rm B}}{\partial x_{\rm j}} \, dV \approx \eta_{\rm num} \int |j|^2 dV. \qquad \text{, where} \qquad \frac{M_{\rm ij}^{\rm B} = B_{\rm i} B_{\rm j} - \delta_{\rm ij} u_{\rm B}}{({\sf Maxwell stress tensor})}$$

Applying appropriate boundary conditions and preserving dominant terms, we have,

Shearing box
$$\left(-q\Omega L_{\rm x}\int_{\rm x_1}B_{\rm x}B_{\rm y}dS_{\rm x}\right)-\int v_{\rm y}B_{\rm x}\frac{\partial B_{\rm y}}{\partial x}dV\right)\approx \eta_{\rm num}\int |j|^2dV.$$

Terms on LHS are larger in a global model due to open radial boundaries and large scale radial gradients

Global disk
$$-\int_{r_1} B_{\mathbf{r}} B_{\phi} v_{\phi} dS_{\mathbf{r}} - \int v_{\phi} B_{\mathbf{r}} \frac{\partial B_{\phi}}{\partial r} dV \simeq \eta_{\text{num}} \int |j|^2 dV.$$



• Our initially purely toroidal field global disk models converge with $\alpha \sim 0.04$, in agreement with Beckwith+(2011) and Hawley +(2013)



- Our initially purely toroidal field global disk models converge with $\alpha \sim 0.04$, in agreement with Beckwith+(2011) and Hawley +(2013)
- Convergence at lower resolution in a global model (~30 cells/H) than in a local model (stratified shearing-box ~ 64-128 cells/H) (- see Davis+2010; Hawley+2011 and refs there-in)



- Our initially purely toroidal field global disk models converge with $\alpha \sim 0.04$, in agreement with Beckwith+(2011) and Hawley +(2013)
- Convergence at lower resolution in a global model (~30 cells/H) than in a local model (stratified shearing-box ~ 64-128 cells/H) (- see Davis+2010; Hawley+2011 and refs there-in)
- Arises because boundary conditions in global models allow more effective energy generation than in a shearing-box (Parkin & Bicknell, 2013, arXiv:1306.1084)

