BITS & BRANES, MAY 24, 2012

ON THE EMERGENCE OF SPACE-TIME

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Emergent Space-time

- Gravity and space-time physics, as described by GR
 + QFT, are likely to be emergent.
- String theory gives important hints, but itself still needs to be developed into a general framework.

Questions

- Is there a universal scenario for the emergence of space-time and gravity?
- What are the basic principles and mechanisms?

Holography & Cosmology

- A proposed scenario should explain the value of the Bekenstein-Hawking entropy.
- And ideally also give a hint towards the origin of the de Sitter entropy?

- What are the implications of the emergence of space-time and gravity in a cosmological setting?
- Does it tell us something about the origin of dark energy and possibly even dark matter?

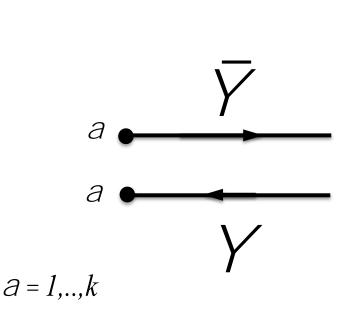
Open String Field Theory

 $k \operatorname{Tr} \left(AQ^* A + \frac{2}{3} A^* A^* A \right)$

Can open string field theory be derived from the anomaly of an

"Half-String" Field Theory

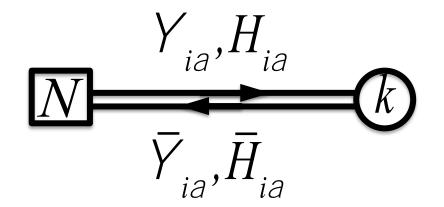
 $\operatorname{Tr}(\overline{Y}_{a}(Q_{L} + A) \times Y_{a})$



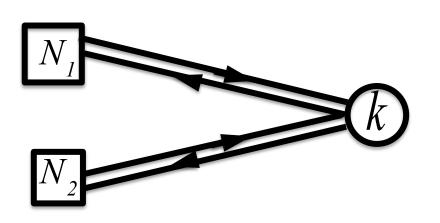
 $\leftrightarrow X$

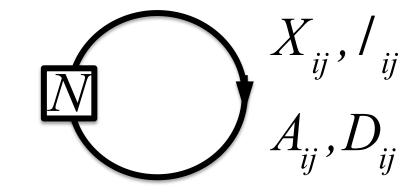
'ADHM' Quiver QM = "atoms of space-time"

N

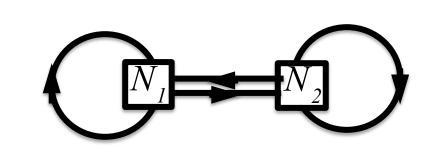


Bifundamental Hypermultiplets





Adjoint Vectormultiplets



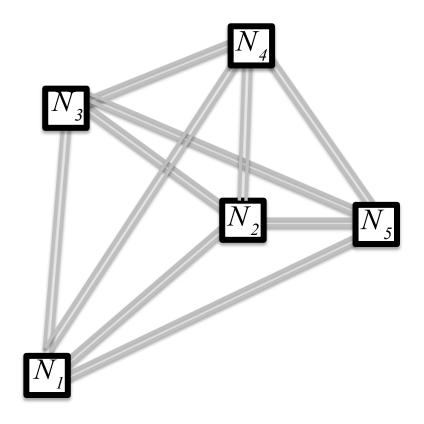
 $\left|X_{ij}H_{a}^{j}\right|^{2} + \overline{Y}_{a}^{i}X_{ij} \times g Y_{a}^{j}$ $\operatorname{Tr}_{\check{\mathsf{C}}\check{\mathsf{e}}}^{\check{\mathfrak{e}}}X^{I}, X^{J}\check{\mathsf{l}}^{2}$ + $\overline{I} g_{I\check{\mathsf{e}}}^{\check{\mathsf{e}}}X^{I}, /\check{\mathsf{l}}^{0}_{\check{\mathsf{e}}}$

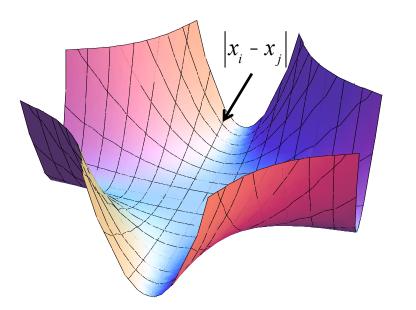
Either *H* or *X* has large but slow variations while the other has fast but small variations. The same holds for the eigenvalues and "off-diagonal" modes of *X*. Space time corresponds to the moduli space of these eigenvalues.

M(atrix) theory

Matter is described as bound states of N_i 'eigenvalues'.

Gravity arises due to integrating out the offdiagonal modes.



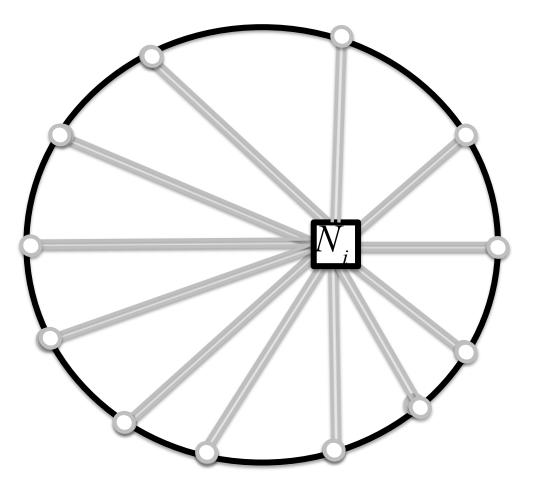


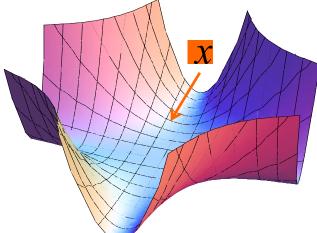
 $\operatorname{Tr}_{\check{\mathsf{Q}}\check{\mathsf{e}}}^{\check{\mathsf{d}}}X^{I}, X^{J}\check{\mathsf{U}}^{2} + \overline{I} g_{I}\check{\mathsf{e}}X^{I}, /\check{\mathsf{U}}_{\check{\mathsf{d}}}^{\circ}$ » + $|x_i - x_j|^2 |X_{ij}|^2 + |x_i - x_j| |I_{ij}|^2$...

In the coulomb phase the fermions "condense:

 $\bar{Y}^i_a Y^j_a \ ^1 \theta$

The eigenvalues of X get drawn towards the fermi surface.





 $\left|X_{ij}H_{a}^{j}\right|^{2} + \overline{Y}_{a}^{i}X_{ij} \times \mathcal{G} Y_{a}^{j}$

Counting of states:



The phase space is characterized by *c=6k* and *N*. and roughly can be thought of as a symmetric orbifold

$$M_N \gg M_N^N / S^N$$
 with $C(M) = C$

The dimension of the Hilbert space is typically

$$\dim (H_N) = C(M_N)$$

and obeys Cardy's formula

$$\log \dim (H_N) = 2\rho \sqrt{\frac{c}{6}} \frac{c}{e} N - \frac{c}{24} \frac{\ddot{o}}{\phi}$$

What values do N and c=6k have?

AdS-Schwarschild Black Holes

In general we expect

 $\ell \gg R$

Schwarschild Black Holes

Various authors have derived a Virasoro algebra with

$$\frac{c}{12} = \frac{A}{8\rho G}$$

$$N = MR$$

For black holes these are equal so that

(Carlip, Padmanabhan)

'de Sitter space'

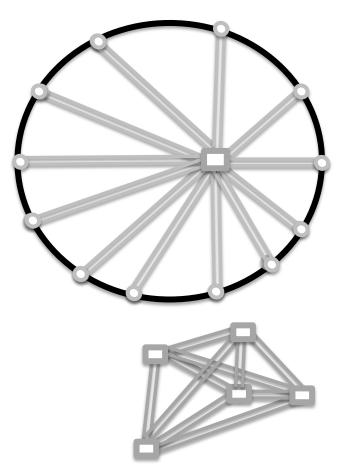
$$ds^{2} = -(1 - H_{0}^{2}R^{2})dt^{2} + \frac{dR^{2}}{1 - H_{0}^{2}R^{2}} + R^{2}dW^{2}$$

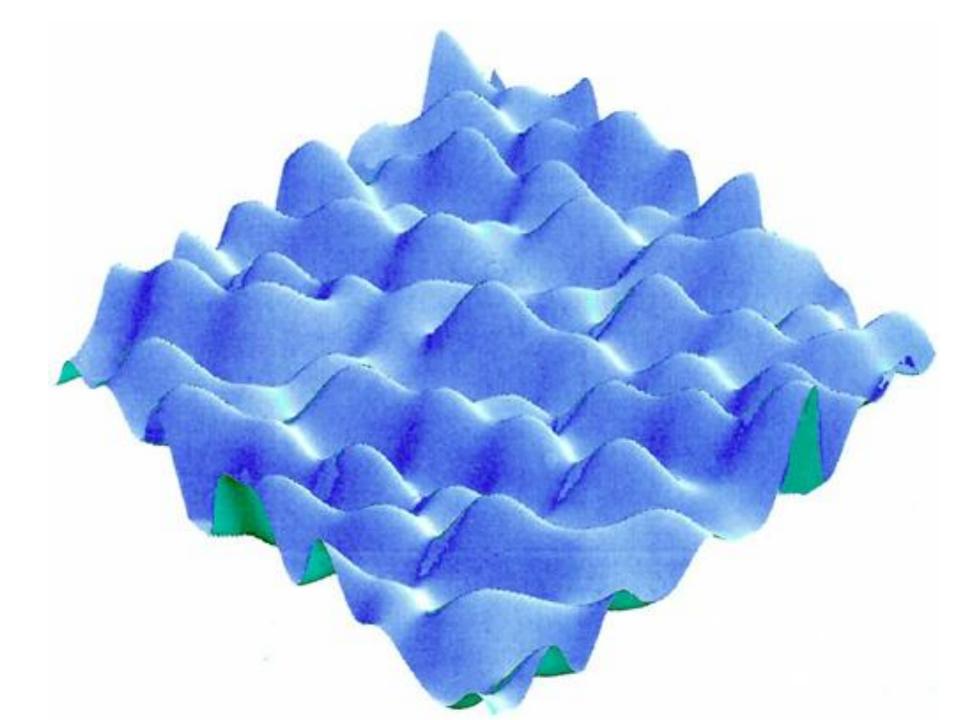
Represents a dynamical "quiver system" again with

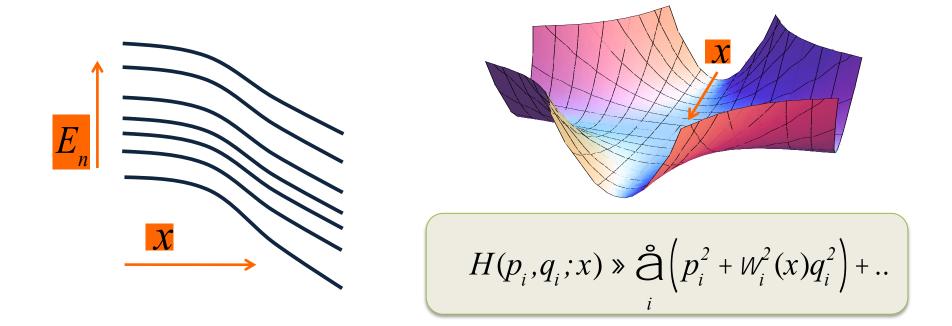
$$\frac{c}{12} = \frac{A_{hor}}{8\rho G} \qquad \qquad N = ER$$

Leading to the correct temperature and entropy (density)

$$T = \frac{H_0}{2\rho} \qquad \qquad s \gg \frac{H_0}{G}$$



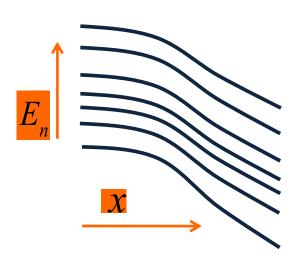




The gravitational (= inertial) force arises as an adiabatic reaction (=Born-Oppenheimer/Casimir) force.

Corrections are given by: Magnetic forces due to (non-abelian) Berry phases. Dynamics of "off-diagonal" modes.

$$F = - \overset{\mathfrak{A}}{\varsigma} \overset{\P E_n \ddot{0}}{\underset{e}{\overset{\circ}{\P}} x \overset{\circ}{\vartheta}}$$



When the separation of time scales between the fast and slow variables is large the phase space volume

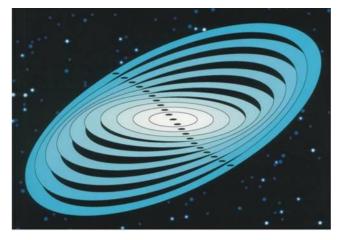
$$\Omega(E,x) = \int d^N p \ d^N q \Big|_{H(p,q;x) \le E}$$

is preserved

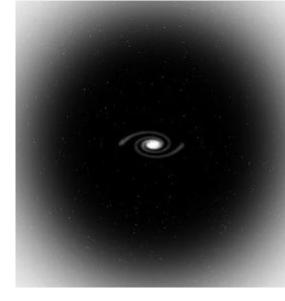
The force can then be written in the form:

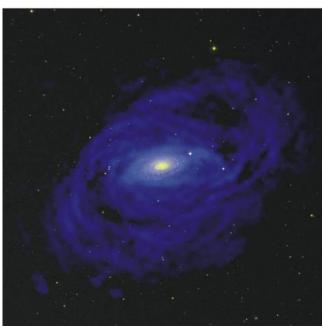
$$F = - \mathop{\mathbb{C}}\limits^{\mathfrak{A}} \underbrace{\P E \, \stackrel{\mathbf{\ddot{0}}}{\stackrel{\div}{\P}}}_{\check{\mathbb{C}}} = T \mathop{\mathbb{C}}\limits^{\mathfrak{A}} \underbrace{\P S \, \stackrel{\mathbf{\ddot{0}}}{\stackrel{\div}{\P}}}_{\check{\mathbb{C}}} \underbrace{\P S \, \stackrel{\mathbf{\ddot{0}}}{\stackrel{\div}{\P}}}_{\check{\mathbb{C}}}$$

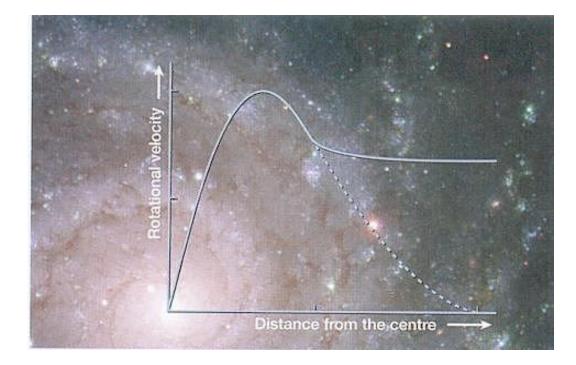
$$S = \log W$$
$$\frac{1}{T} = \frac{\P_E W}{W}$$

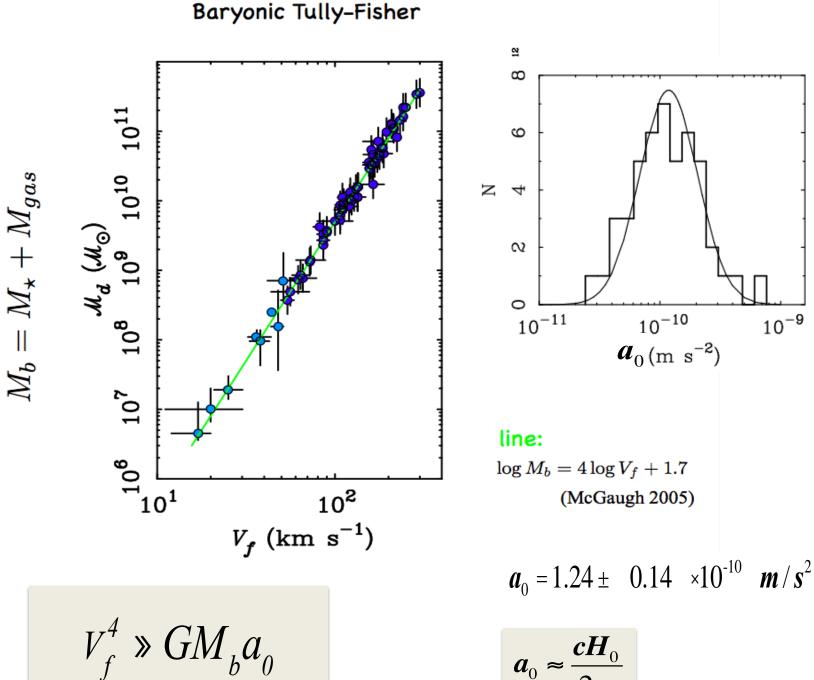


Dark matter In Galaxies









 $a_0 \approx \frac{cH_0}{2\pi}$

Suppose we may define a microscopic phase space volume W(E; r(x))

for a given matter density ρ and interpret Newton's potential Φ as the field dual to ρ in a canonical ensemble

$$Z(b; \mathsf{F}) = \int [dr] W(E; r) e^{-b(E - \int \mathsf{F}r)}$$

What would this say about the gravitational force? The saddle point equations give the relations

$$\mathsf{F} = -T\frac{d}{dr}\log W(E, r) \qquad \frac{1}{T} = \frac{\partial}{\partial E}\log W(E, r)$$

The gravitational force can thus be written as

$$F = \int r \nabla F = -T \int \left(r \nabla \frac{d}{dr} \log W(E, r) \right)$$

which after partial integration and by inserting

$$\Gamma(x) = \mathop{\mathrm{a}}\limits^{\circ} m_i \mathcal{O}(x - x_i)$$

gives the following suggestive equation

$$F_i = m_i \nabla F(x_i) = T \nabla_i S(E, \{x_i\})$$

where S(*E*, {*x*_{*i*}}) equals the microscopic entropy. => defining equation for an **adiabatic reaction force**! Let us go back to the relation

$$\langle \nabla F \rangle = -T \nabla \frac{d}{dr} \log W(E, r)$$
 "linear response"

And use the fact that we know that To compute the fluctuations

$$\langle \nabla^2 \mathsf{F} \rangle = 4\rho G \Gamma$$

$$\langle |\nabla F|^2 \rangle = \left(T \nabla \frac{d}{dr} \right)^2 \log W(E, r) = 4\rho G d^{(3)}(0)$$

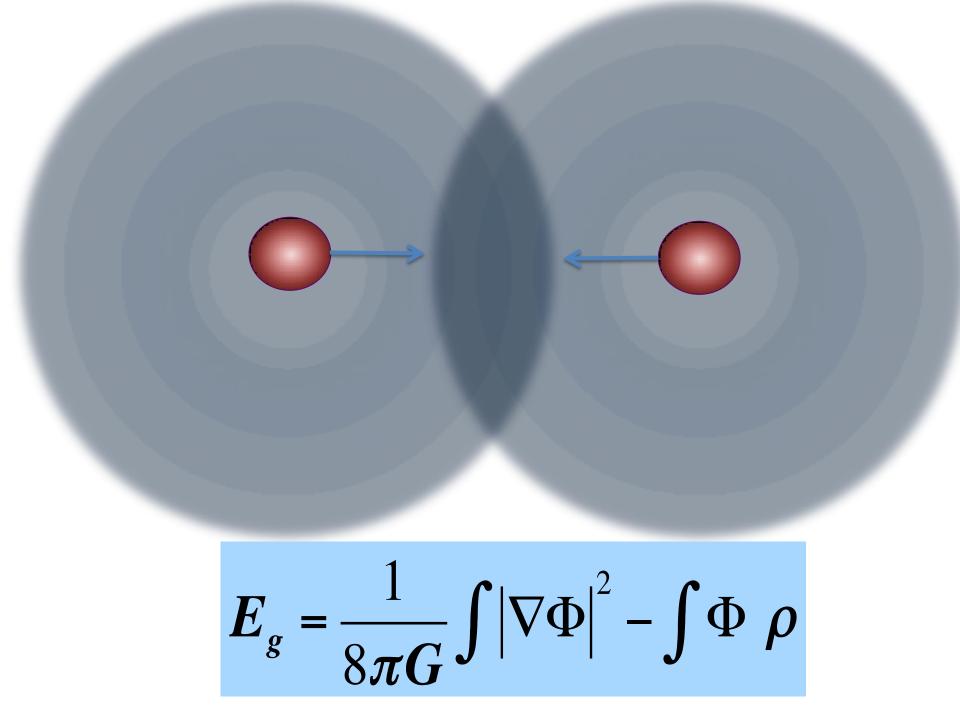
"fluctuation -dissipation"

Or integrated

$$\frac{1}{8\rho G} \int_{V} \left\langle |\nabla \mathsf{F}|^{2} \right\rangle = \frac{1}{2} N k_{B} T$$

N = number of modes contained in volume V

Note: the size of fluctuations is determined by the UV cut off.



Alternative derivation: consider

$$Z[b;r] = \int [dF]e^{-b\left(\frac{1}{8\rho G_V}\int_V |\nabla F|^2 - \int_V Fr\right)}$$

and compute the one and two point functions. $\left< \nabla {\rm F} \right>$ gives the Newtonian acceleration for ρ

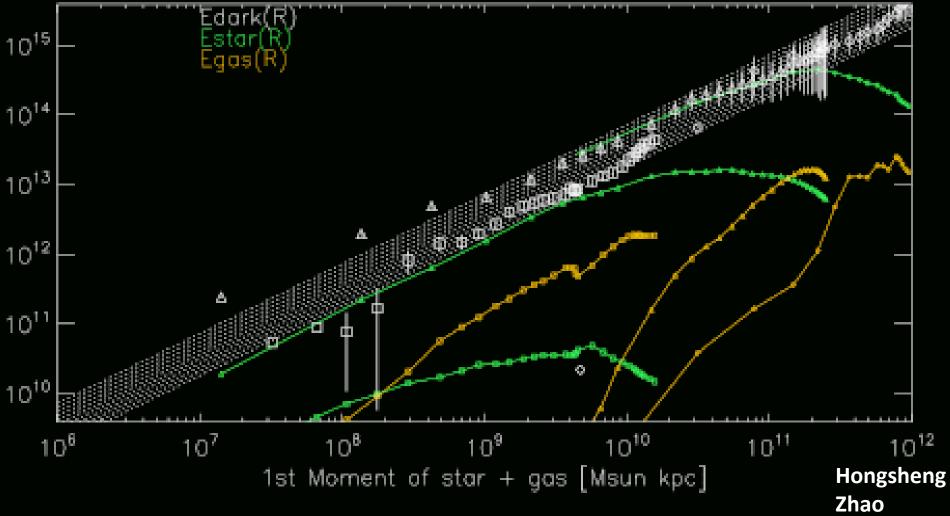
while

$$\frac{1}{8\rho G} \int_{V} \left\langle \left| \nabla F \right|^{2} \right\rangle = \frac{1}{2} N k_{B} T$$

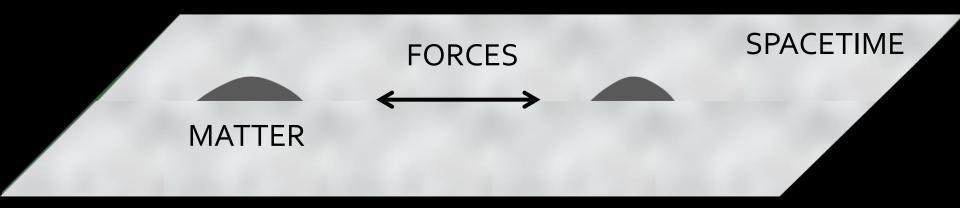
where N is the number of modes of Φ inside the volume V.

$$C_g(R) = \frac{1}{8\rho G} \int_0^R \left| \nabla F_D \right|^2 = \frac{1}{2} N k_B T \qquad k_B T = \frac{\hbar H_0}{2\rho} \qquad N = \frac{M_B c R}{\hbar}$$





EMERGENT



DARK MATTER

DARK ENERGY

