

Turbulence and waves “but how much mixing?” *and dissipation?*

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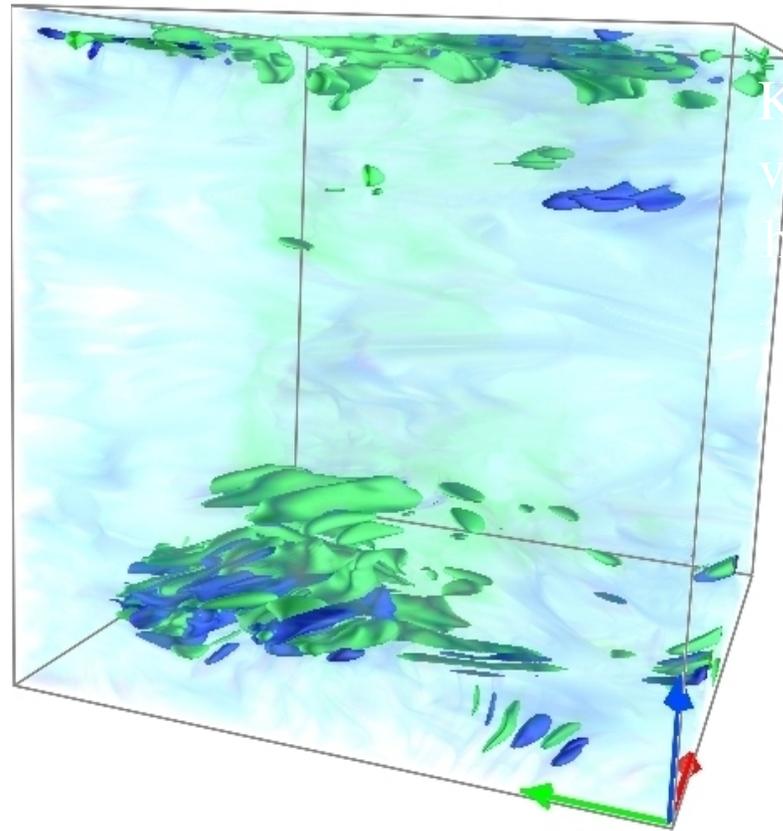
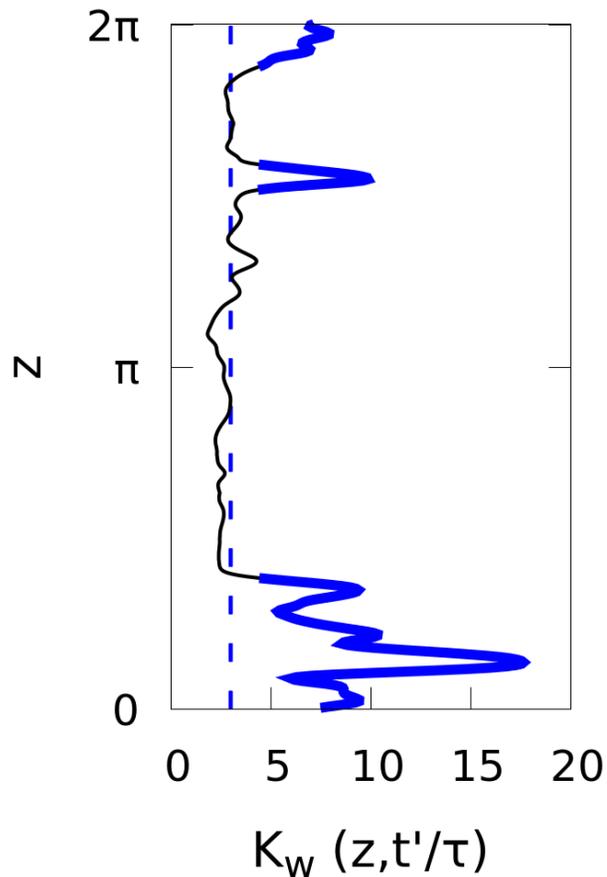
Raffaele Marino, LMFA-ECL/Lyon

Duane Rosenberg, CIRA/NOAA-Boulder

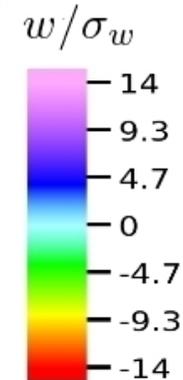
NSF/NCAR/Yellowstone: 76 decay or forced RST runs @1024³ res., 6 forced runs @2048³ ; ~ 85% background (free) time
DOE/Titan: Decay RST run @ 4096³ point resolution, a few outputs of which are on the *John Hopkins turbulence data base*

Scaling laws for mixing & dissipation in unforced rotating stratified turbulence, *J. Fluid Mech.* **844**, 519, 2018
Variations of characteristic time scales in rot. strat. turb. using a large parametric numerical study, *Eur. Phys. J-E* **39**, 8, 2016
Evidence for Bolgiano-Obukhov scaling in rotating stratified turbulence using high-resolution Direct Num. Sim., *Phys. Fluids* **27**, 055105, 2015

Forced 512^3 run,
 $Fr \sim 0.08$, $Re \sim 3800$,
No rotation (*Newton-Calabria*)



Kurtosis of vertical
velocity w ,
horizontally averaged,
& visualization
of vertical velocity w



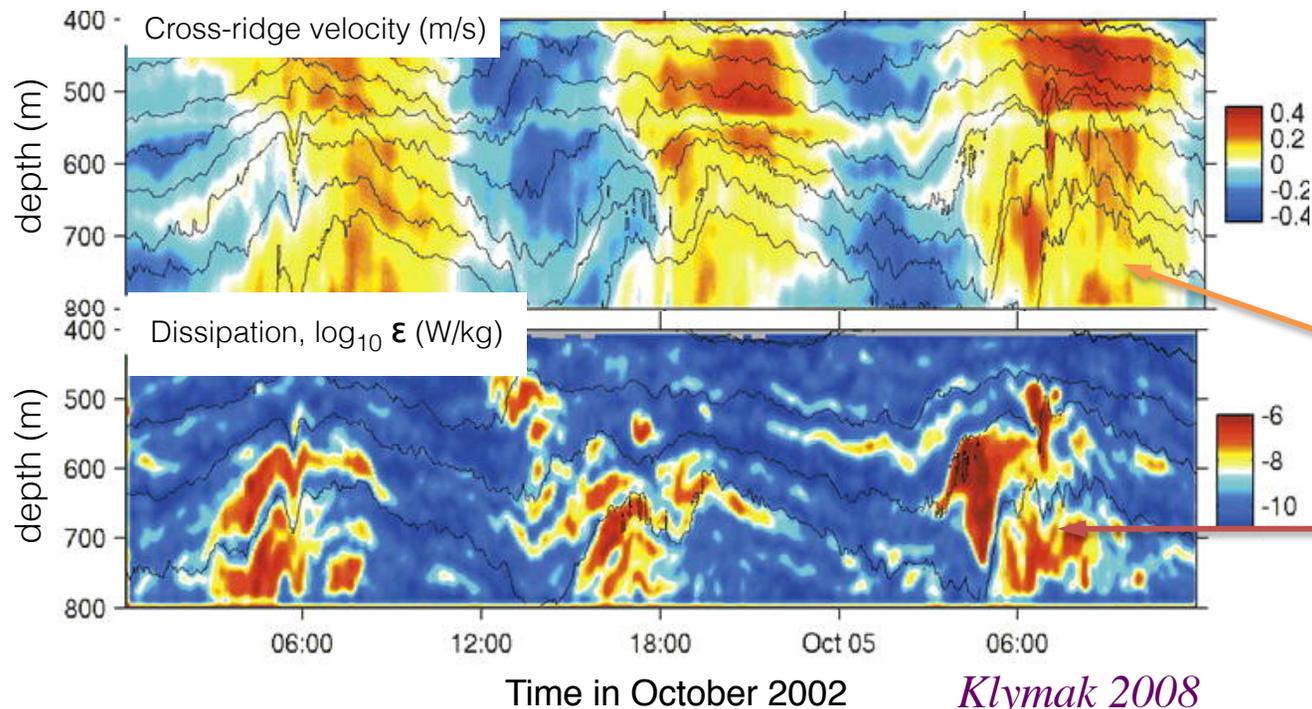
Strong intermittency of the vertical *velocity*

Model through the Vieillefosse system with gravity waves, forcing and dissipation

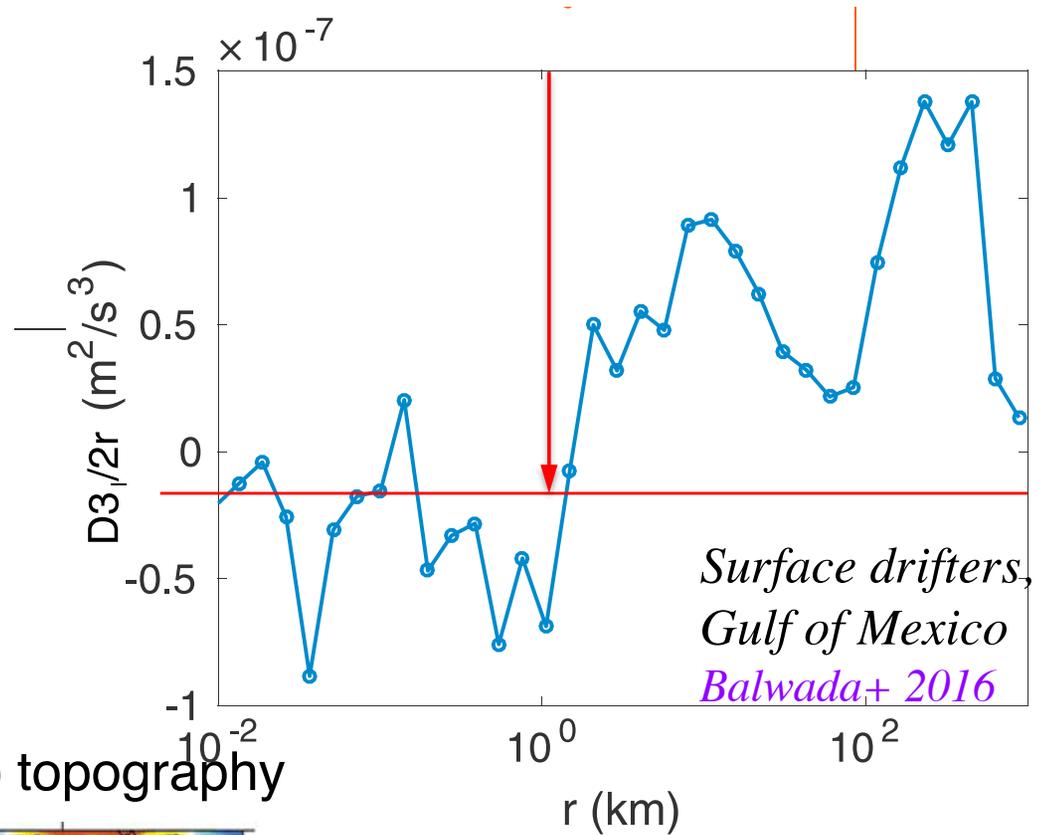
Feraco+, Vertical drafts and mixing in stratified turbulence: sharp transition with Froude number. Submitted to *Eur. J. Phys. Lett.*. ArXiv/ 1806.00342

(see also Rorai+, *Turbulence comes in bursts in stably stratified flows*, *Phys. Rev. E* **89**, 043002, 2014)

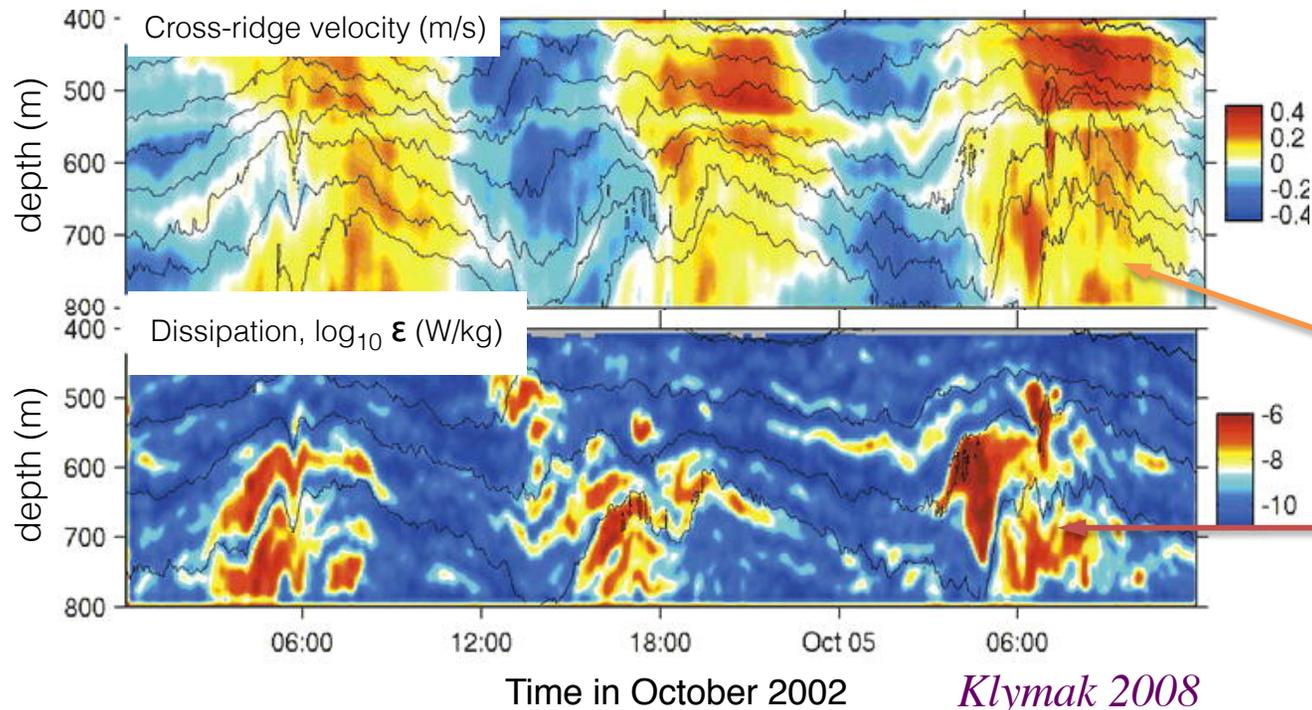
b) Breaking internal tides over tall steep topography



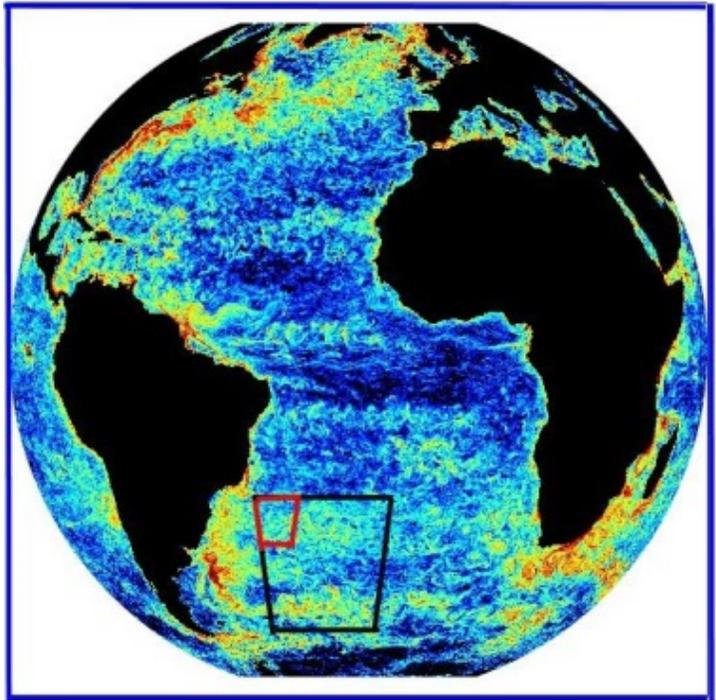
Hawaiian ridge
 $U=0.1\text{ m/s}$, $L=1000\text{ m}$
 $\tau_{NL}=L/U \sim 3\text{ hrs}$
 $N=0.001\text{ s}^{-1}$, $Fr \sim 0.1$
 $\epsilon_V \sim 10^{-6}\text{ W} \sim \epsilon_D = U^3/L$
 $Re \sim 10^8$, $R_B \sim 10^6$



b) Breaking internal tides over tall steep topography



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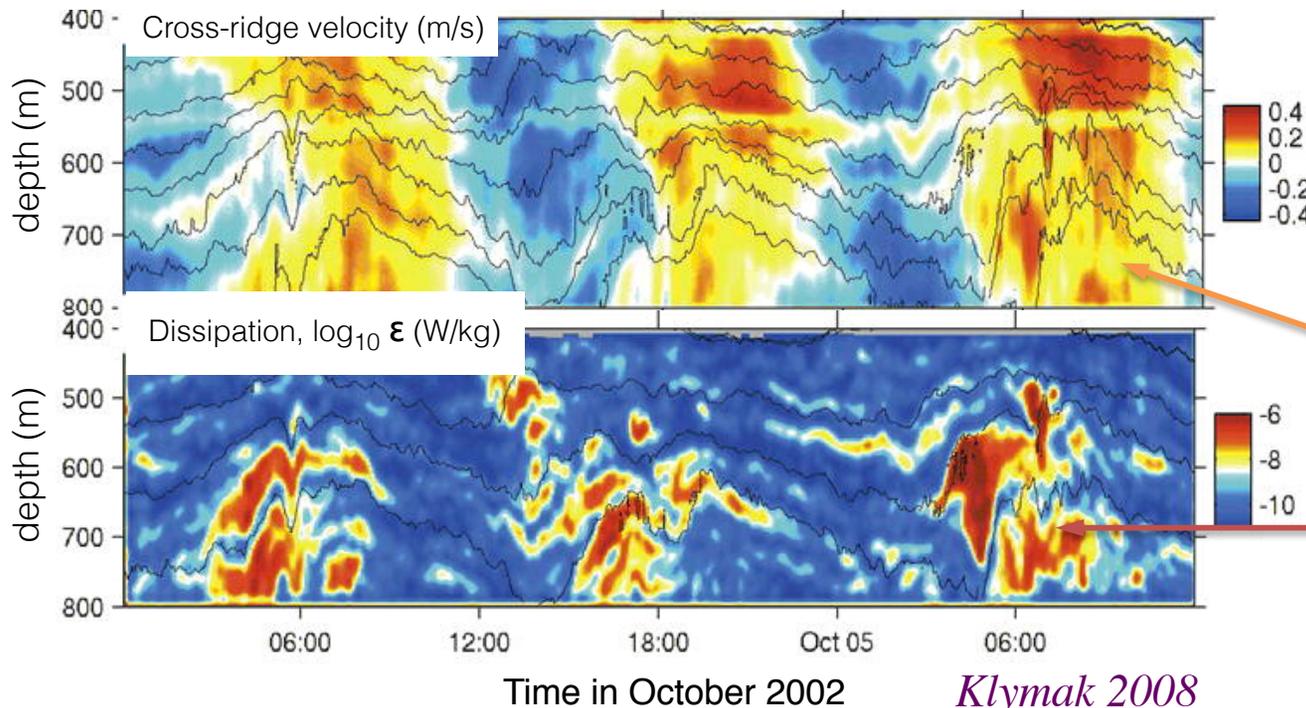


Ocean model,
smallest resolved scale: 20km

Kinetic energy dissipation

Pearson+2018

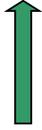
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Klymak 2008

Outline

- Introduction
- Equations and characteristic scales
- Parametric study with direct numerical simulations
- Three constitutive laws for normalized E_p , E_w , ε_v
- Consequences of  for scaling of mixing efficiency+
- Discussion, Conclusions and Perspectives

Incompressible Boussinesq equations + rotation
 3D cubic box, periodic boundary conditions

$$\frac{\partial \mathbf{u}}{\partial t} + \boldsymbol{\omega} \times \mathbf{u} + 2\boldsymbol{\Omega} \times \mathbf{u} = -N\vartheta \hat{e}_z - \nabla \mathcal{P} + \nu \nabla^2 \mathbf{u},$$

$$\frac{\partial \vartheta}{\partial t} + \mathbf{u} \cdot \nabla \vartheta = Nw + \kappa \nabla^2 \vartheta,$$

$$[\theta] = [L T^{-1}]$$

Governing dimensionless parameters:

Reynolds, Froude, Rossby, Prandtl=1

$$Re = \frac{U_0 L_0}{\nu}, \quad Fr = \frac{U_0}{L_0 N}, \quad Ro = \frac{U_0}{L_0 f}, \quad Pr = \frac{\nu}{\kappa}, \quad f = 2\Omega$$

Fr < 1 , together with **Re >> 1**

$N/f = Ro/Fr > 2.5$ (ocean, atmosphere)

SCALES: Purely stratified flow ($f=0$): $Fr = U_0/[NL_0] < 1$

Scale at which $Fr = 1$?

Purely stratified turbulence ($f=0$): $Fr = U_0/[NL_0] < 1$

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→ $L_B = U_0/N$, buoyancy scale

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l ? such that $Fr(l)=1$, for a Kolmogorov spectrum: $u(l) \sim \epsilon_v^{1/3} l^{1/3}$

→ $l = l_{Oz} = [\epsilon_v/N^3]^{1/2}$, Ozmidov scale

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Buoyancy Reynolds number: $R_B \equiv \epsilon_v / [\nu N^2] = [l_{Oz} / \eta]^{4/3}$

$R_B = 1$ for $l_{Oz} = \eta = [\epsilon_v / \nu^3]^{-1/4}$ (η : Kolmogorov dissipation scale)

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Numerical conundrum: large R_B , small Fr for geophysical flows

With rotation: Ro also small; $Ro/Fr = N/f \sim 5$ (ocean) or 100 (atmosphere)

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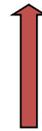
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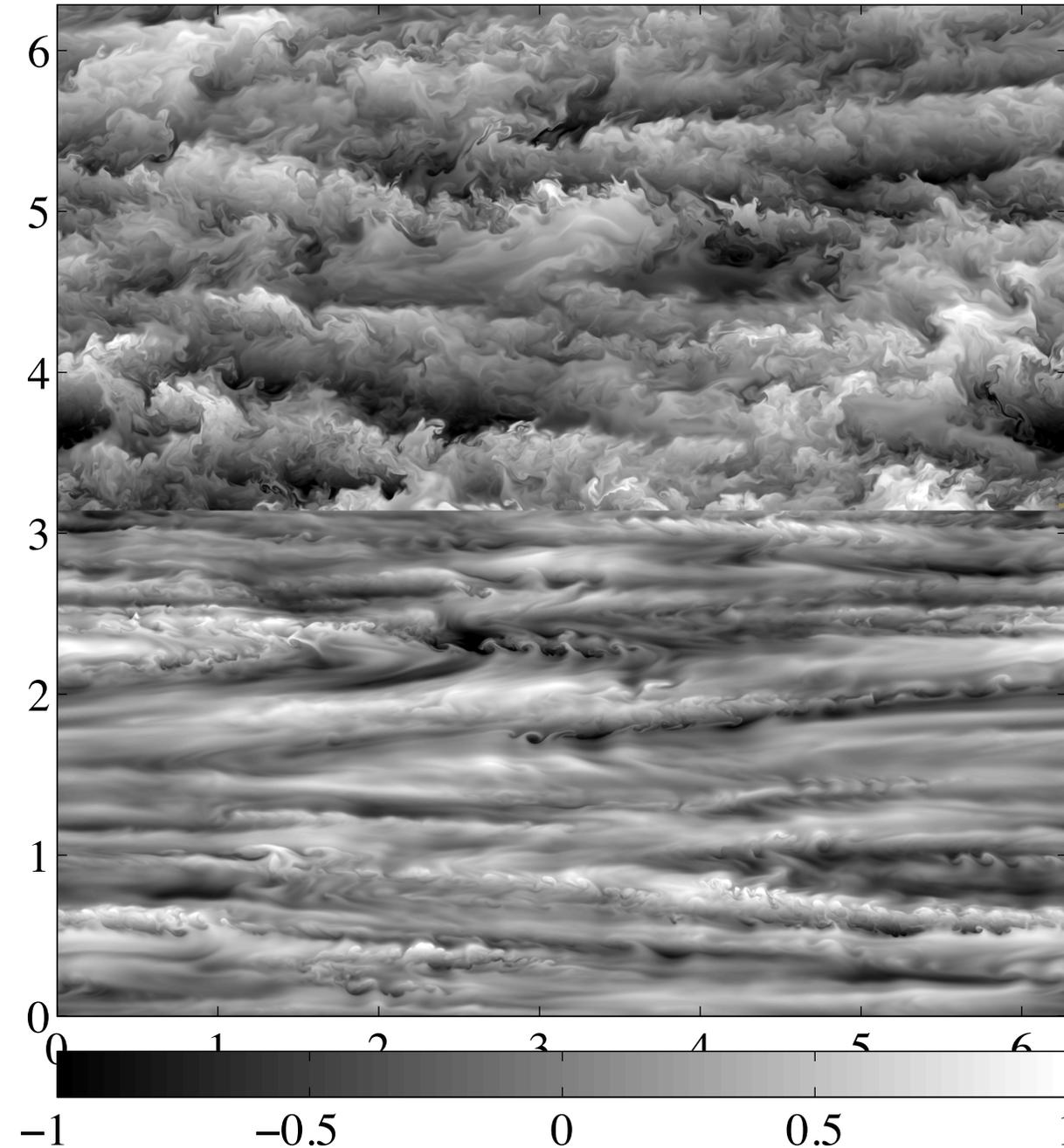
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Stratification, no rotation, large scale forcing, $Re \sim 24000$, 2048^3 grid:
Temperature fluctuations, xz slice



$N=4, \quad Fr \sim 0.11$
 $R_B = ReFr^2 \sim 300$

$N=12, Fr \sim 0.03$
 $R_B \sim 22$

Vertical vorticity
at peak of dissipation
($\omega_{z\text{-mag}}$, *horizontal cut*):

Eddies & lanes

Plot @ full 4096^2 res.
*GHOST pseudo-spectral
code (DOE/titan, 2014)*

Log scale

$f=2.7$, $\omega_{\text{rms}} \sim 17$

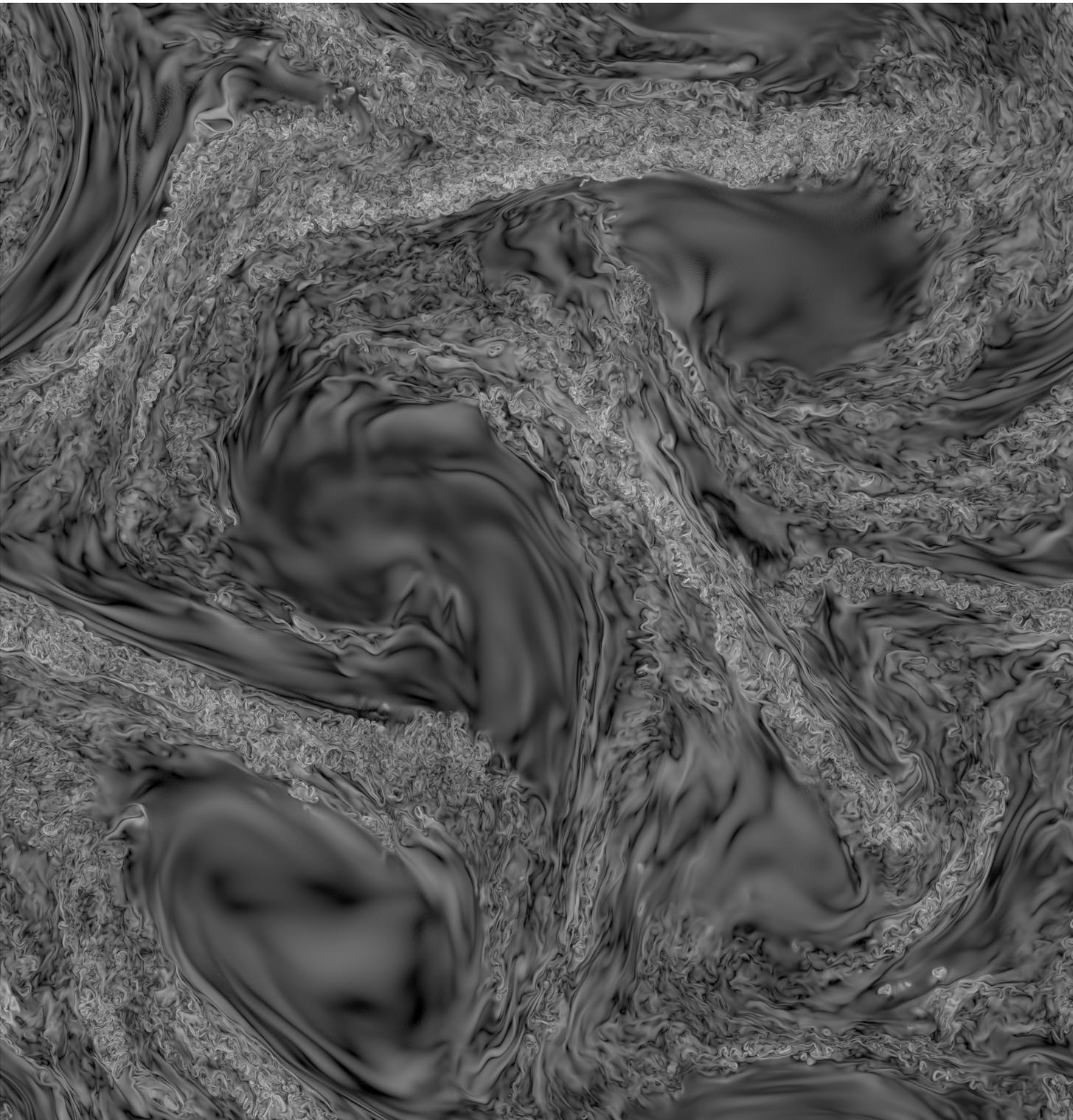
$Re=55000$ ``*ocean*''

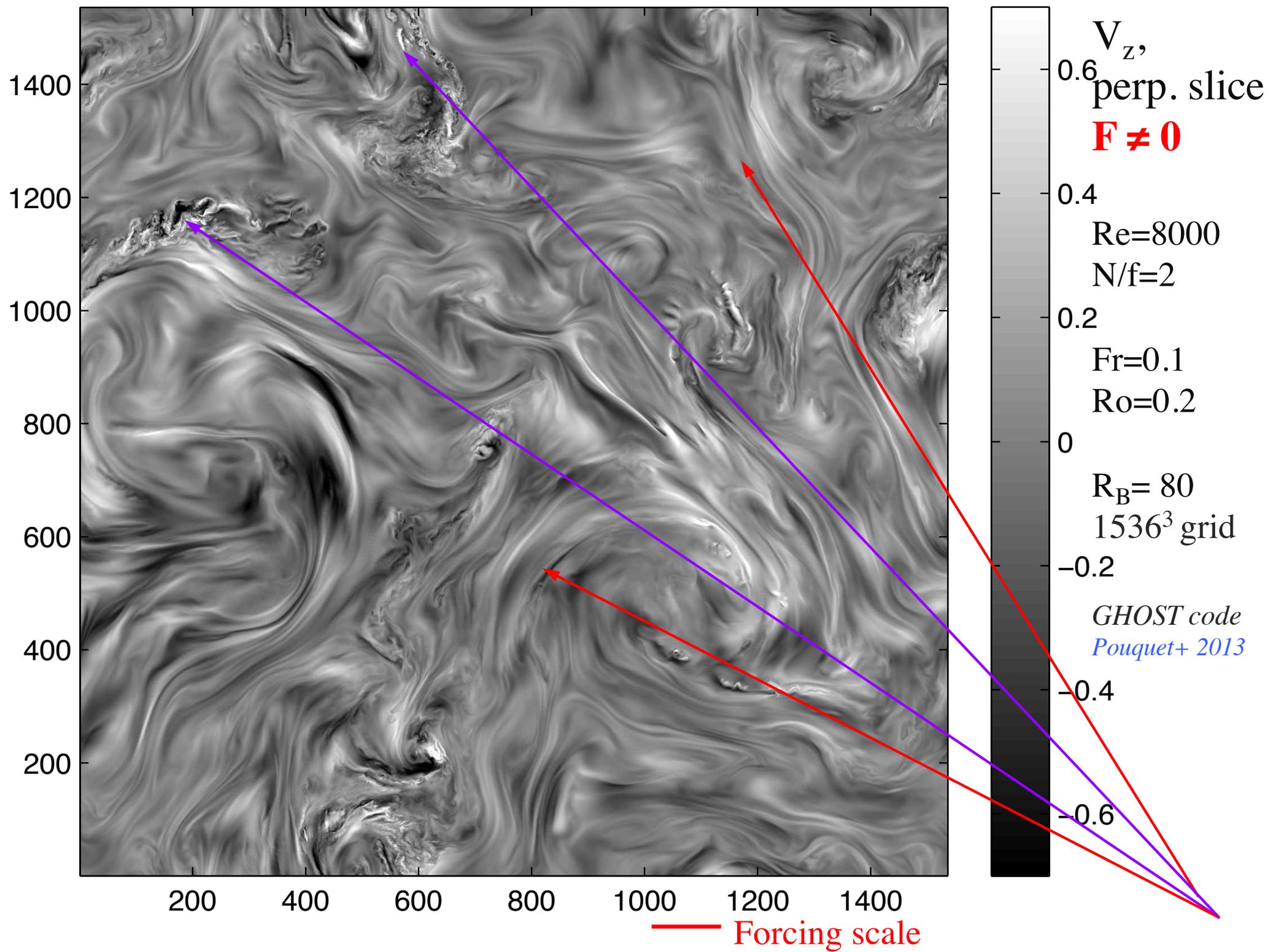
$Fr=0.024$, $N/f=5$

$R_B=32$, $k_{\text{max}}\eta \sim 2$

No forcing, $k_0 \sim 2.5$

*Bolgiano-Obukhov scaling
Rosenberg+ 2015*





THE PARAMETRIC STUDY

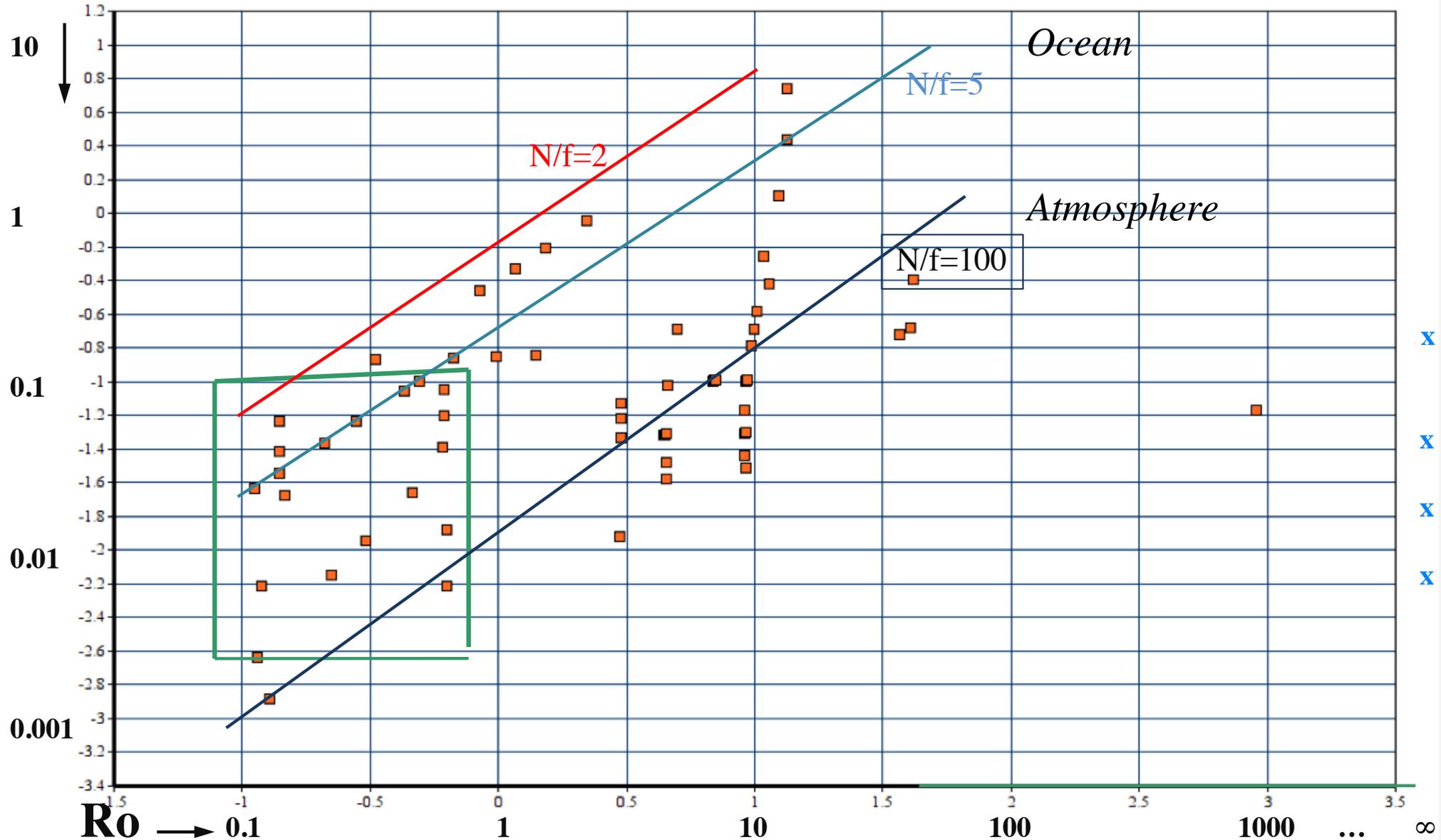
1024^3 , $Re \sim 1.2 \times 10^4$, $\theta(t=0)=0$, $F=0$,
 $k_0 \sim 2.5$, isotropic (I); GHOST code.

Fr

$2.5 \leq N/f \leq 312$

... and ...

∞



THE PARAMETRIC STUDY

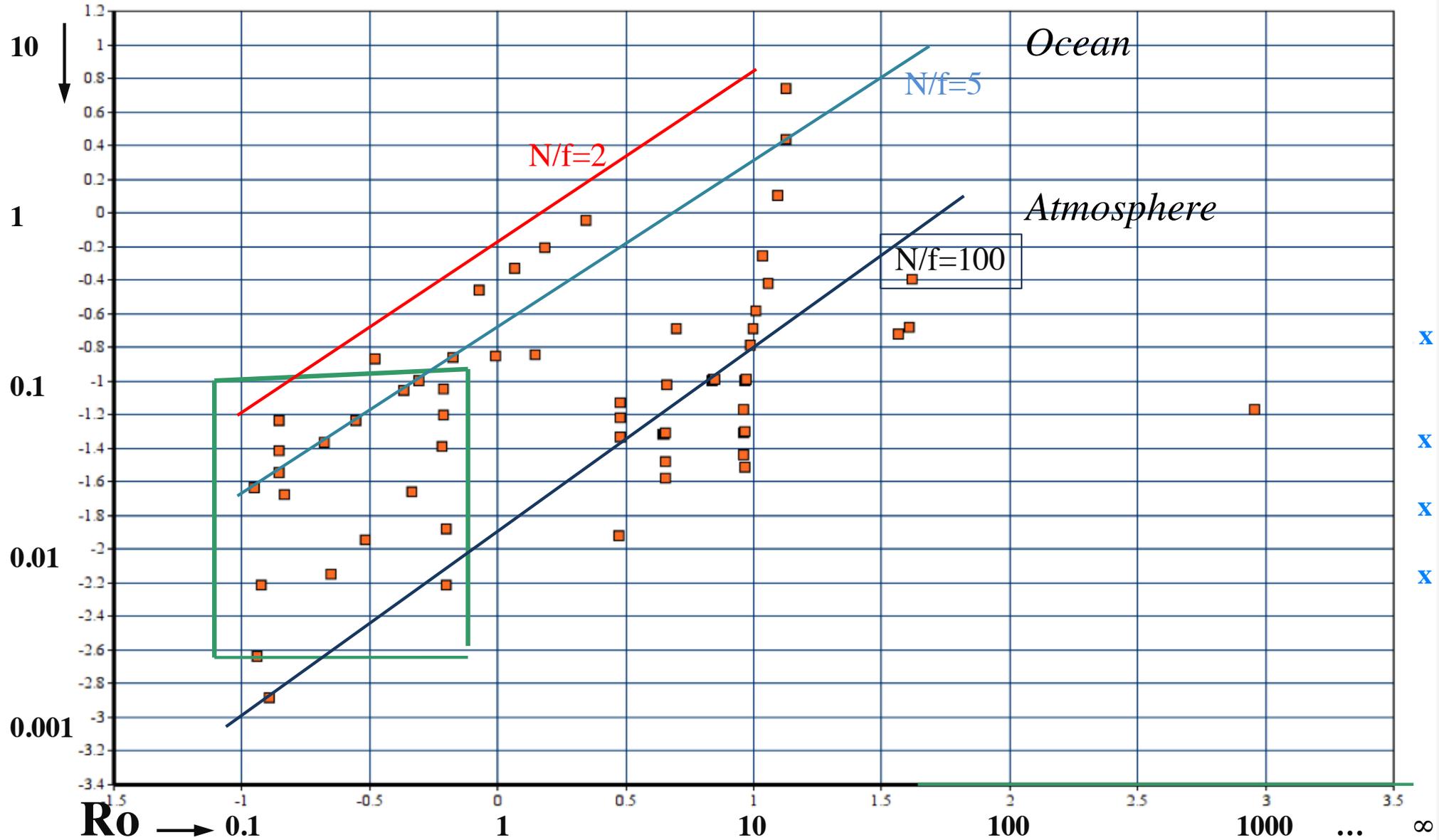
1024^3 , $Re \sim 1.2 \times 10^4$, $\theta(t=0)=0$, $F=0$,
 $k_0 \sim 2.5$, isotropic (I); GHOST code.
Some 512^3 (I or QG) runs, lower Re

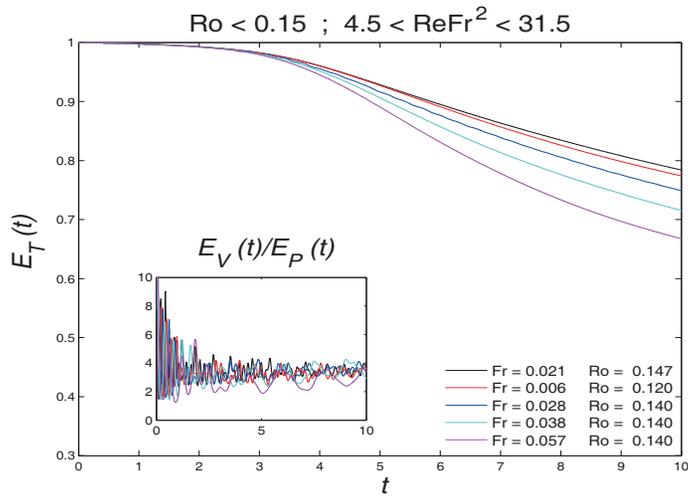
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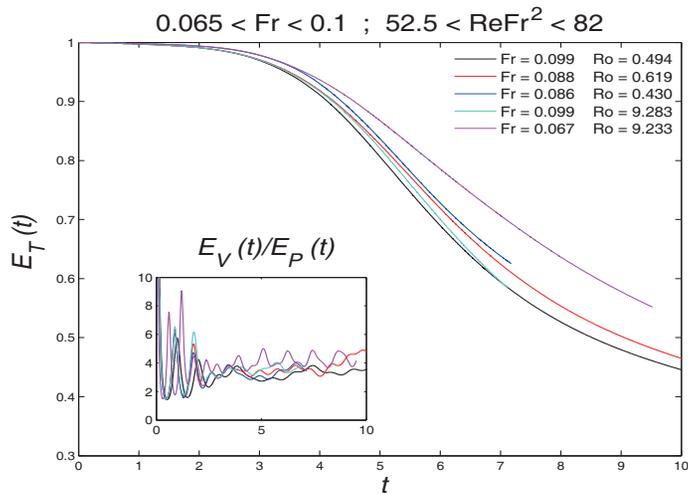




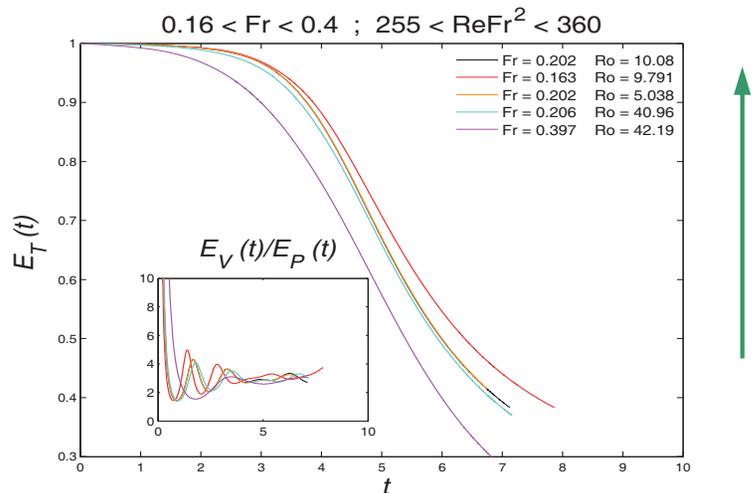
$$E_T = f(t) \text{ and}$$

$$E_V/E_P = f(t)$$

← Low Fr, low Ro
Fr~0.03

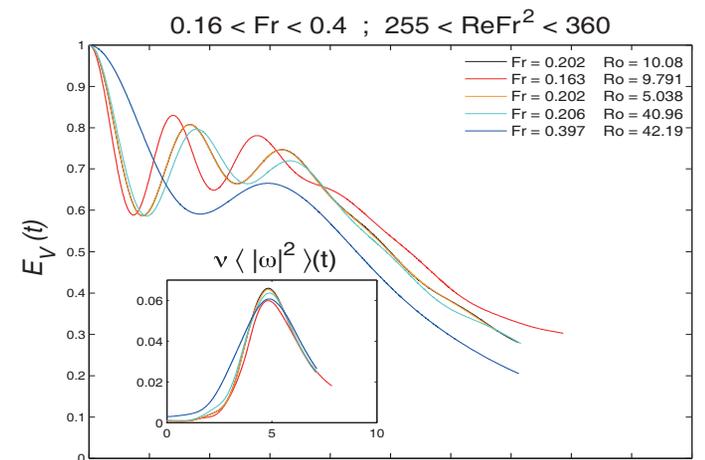
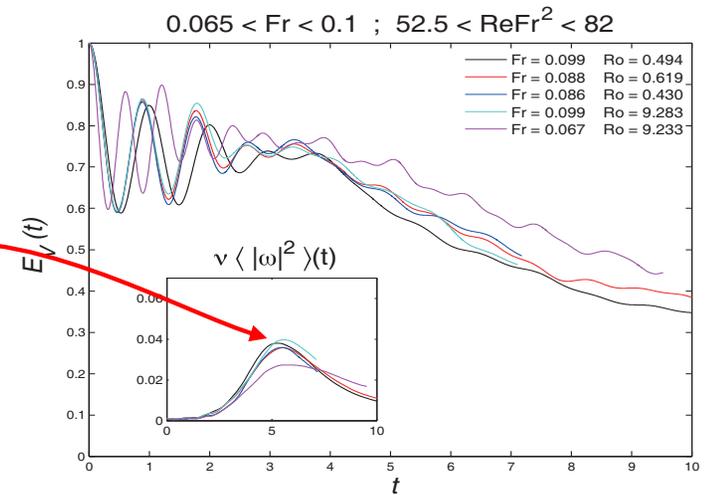
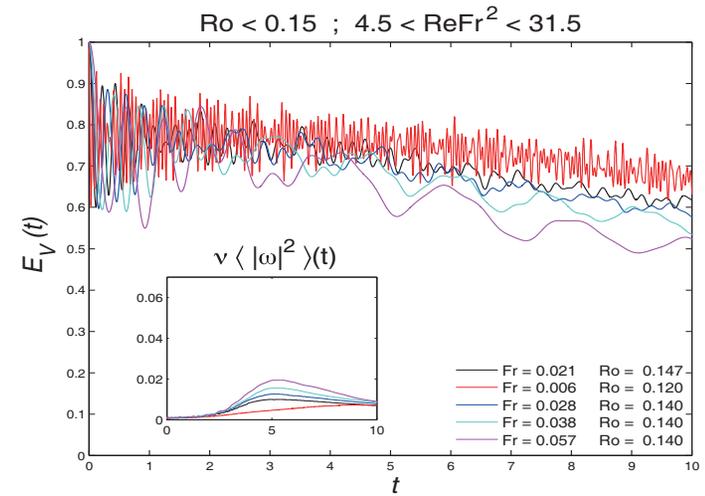


← Intermediate range
Fr~0.09



← High R_B
Fr~ 0.2

$E_V = f(t)$
 and kinetic energy dissipation $\varepsilon_V = f(t)$



Statistics taken at peak of dissipation
 with a temporal averaging within
 variations of $\pm 0.025 \varepsilon_V$

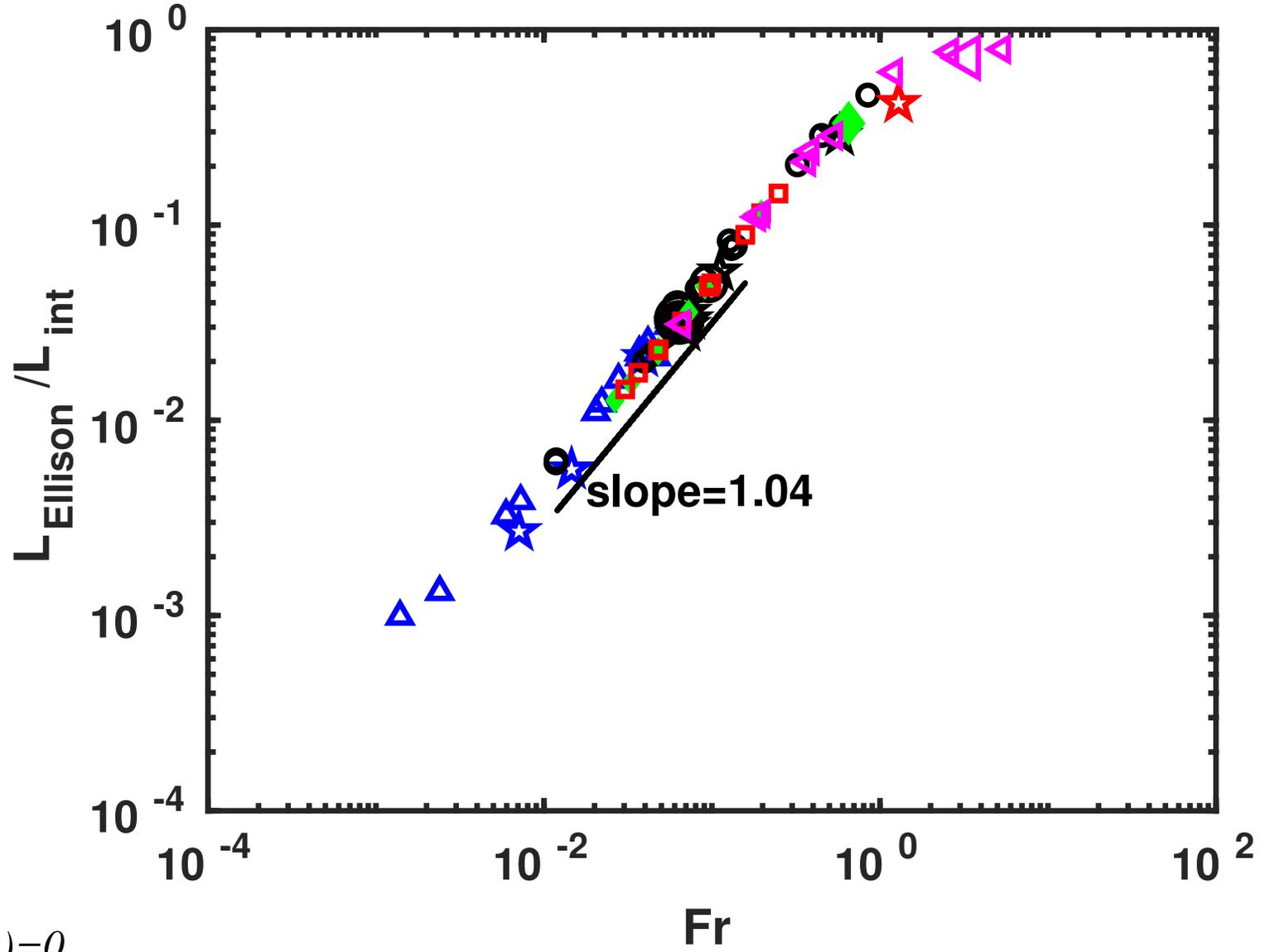
Ellison scale

$$L_E = \sqrt{E_P / N}$$

Integral scale

$$L_{\text{int}} = f(E_V)$$

Color binning in Ro: 0 → 0.3 → 2.9 → 6.0 → 10 →



Initial conditions:

Stars: QG

Otherwise, HIT, $\theta(t=0)=0$

Size of symbol: proportional to viscosity

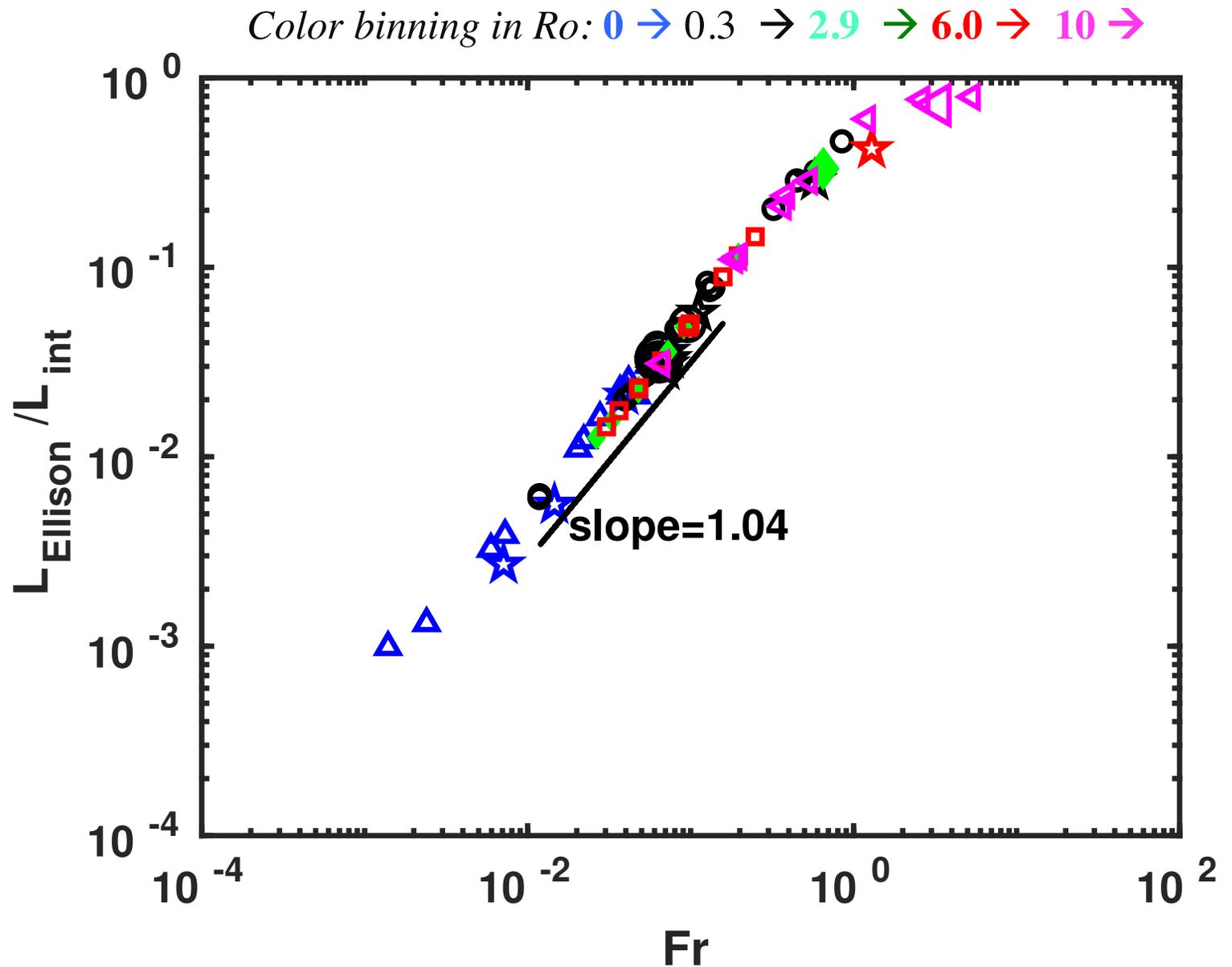
$$L_E / L_{\text{int}} \sim Fr \sim U_{\text{rms}} / [L_{\text{int}} N]$$

Ellison scale

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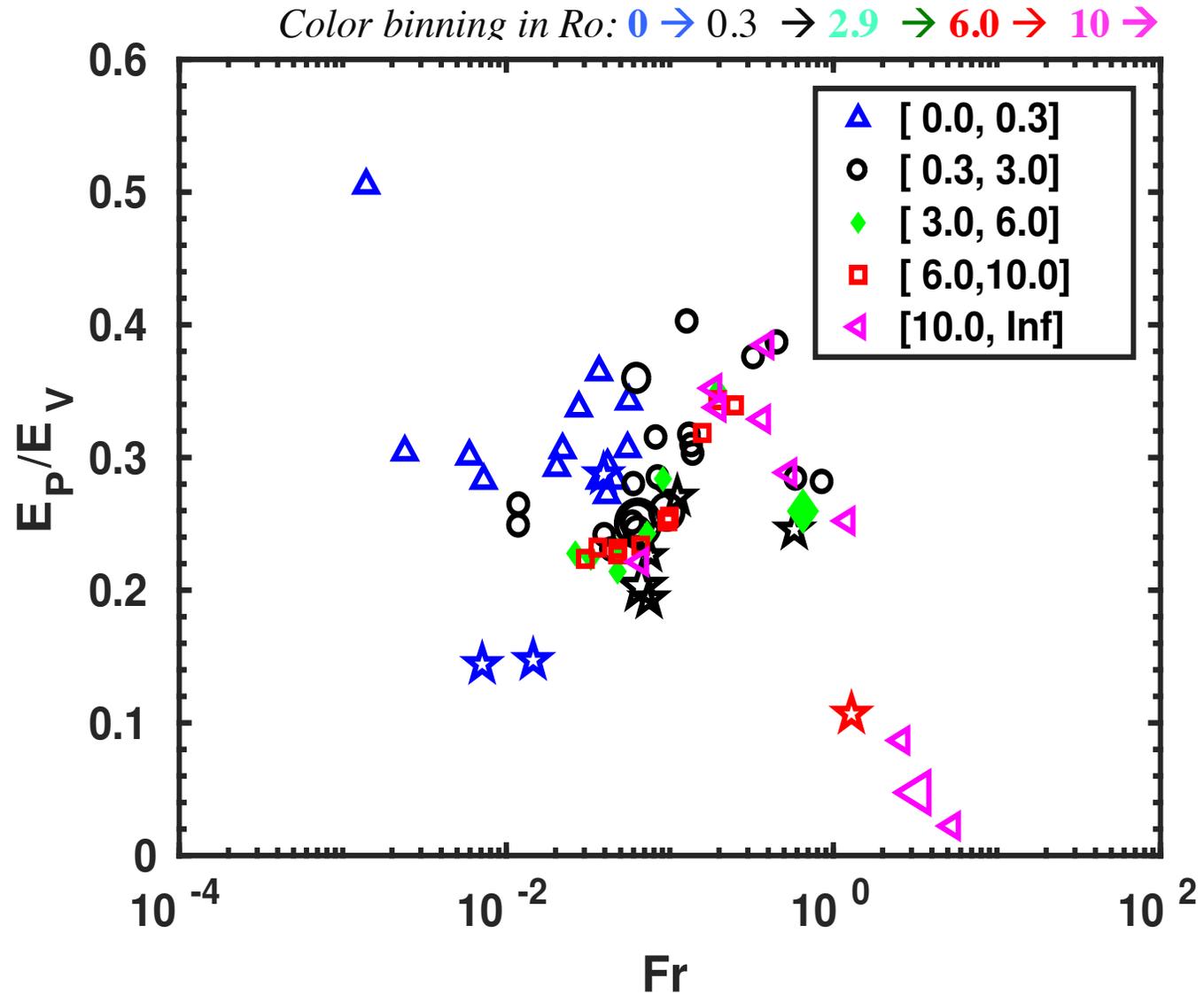
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$$\rightarrow L_E \sim L_B = U_{\text{rms}} / N$$

Ratio of potential to kinetic energy



Initial conditions:

Stars: QG

Otherwise, $u_{perp} \sim w$ and $\theta=0$

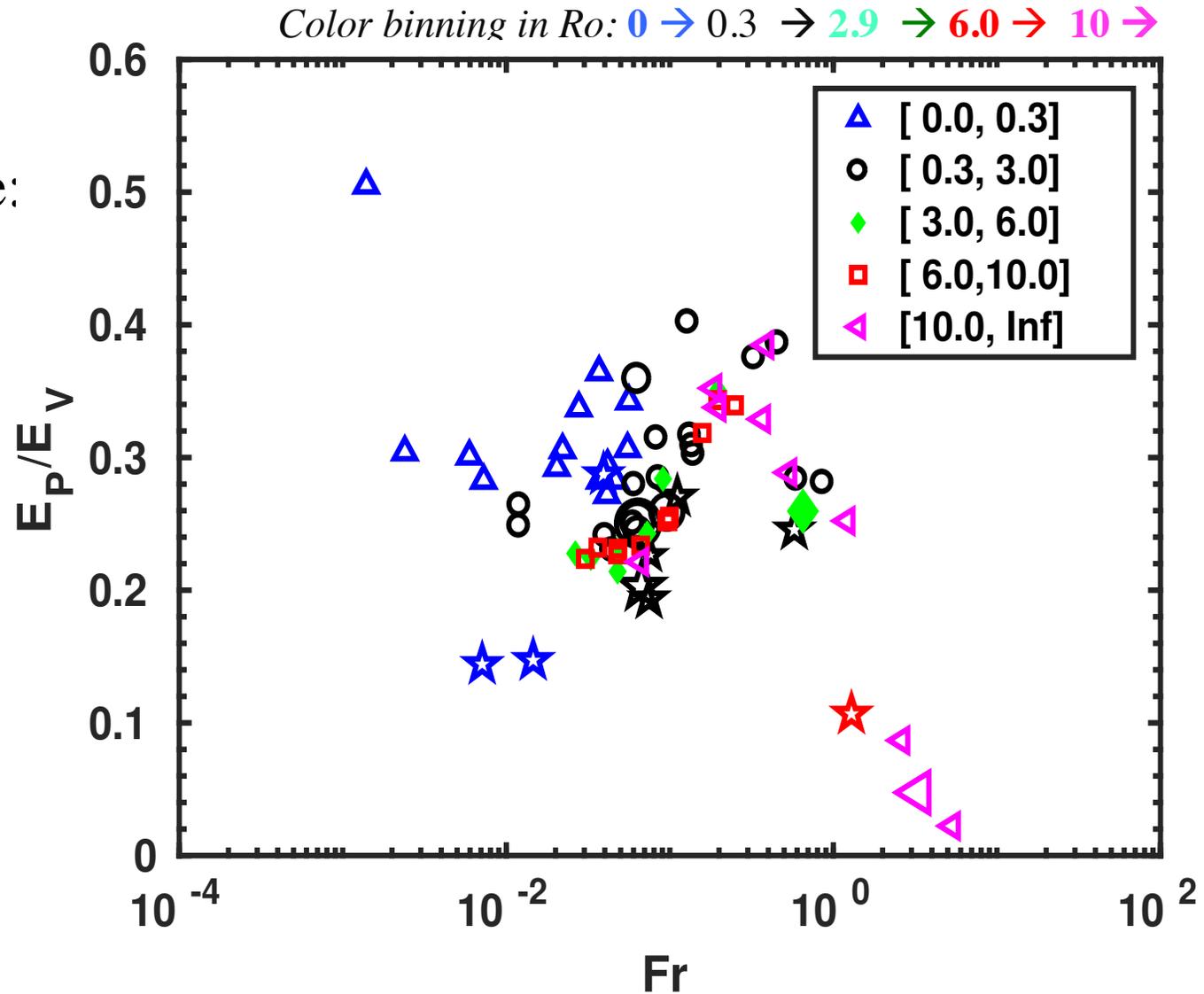
Size of symbols: Roughly inversely proportional to numerical resolution, that is, in fact, to Reynolds number

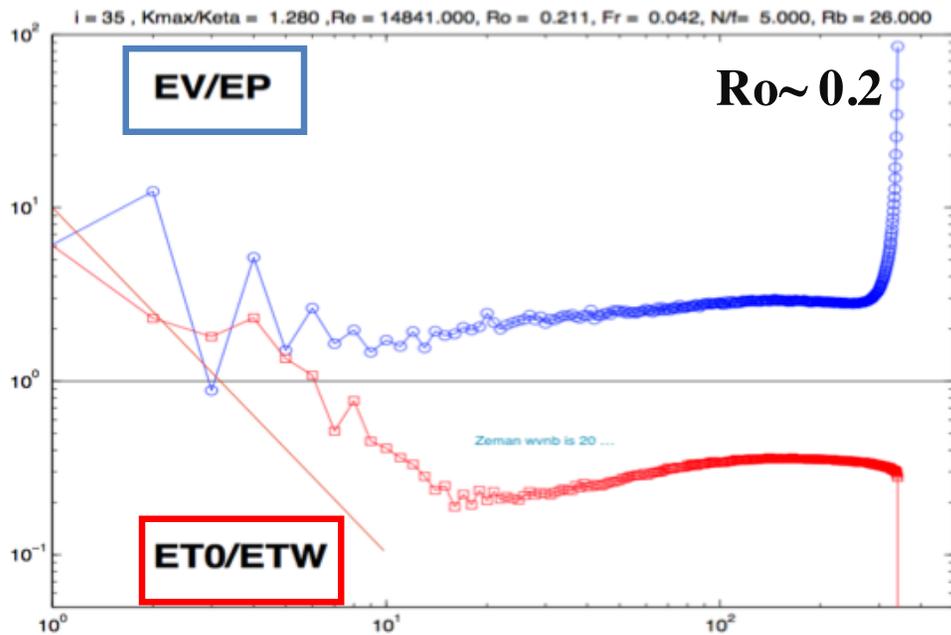
Ratio of potential to kinetic energy

Intermediate regime:
 $E_P \sim E_V$
Or

$$\Theta_{\text{rms}} \sim U_0 \quad (1)$$

Stars: Quasi-geostrophic ICs





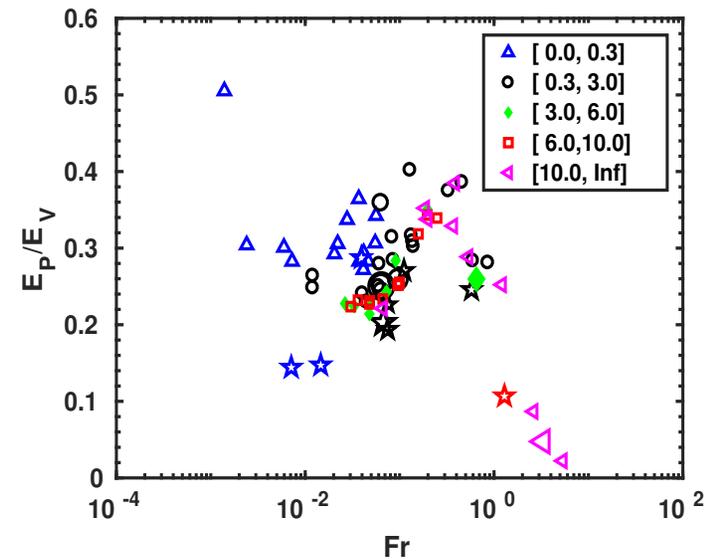
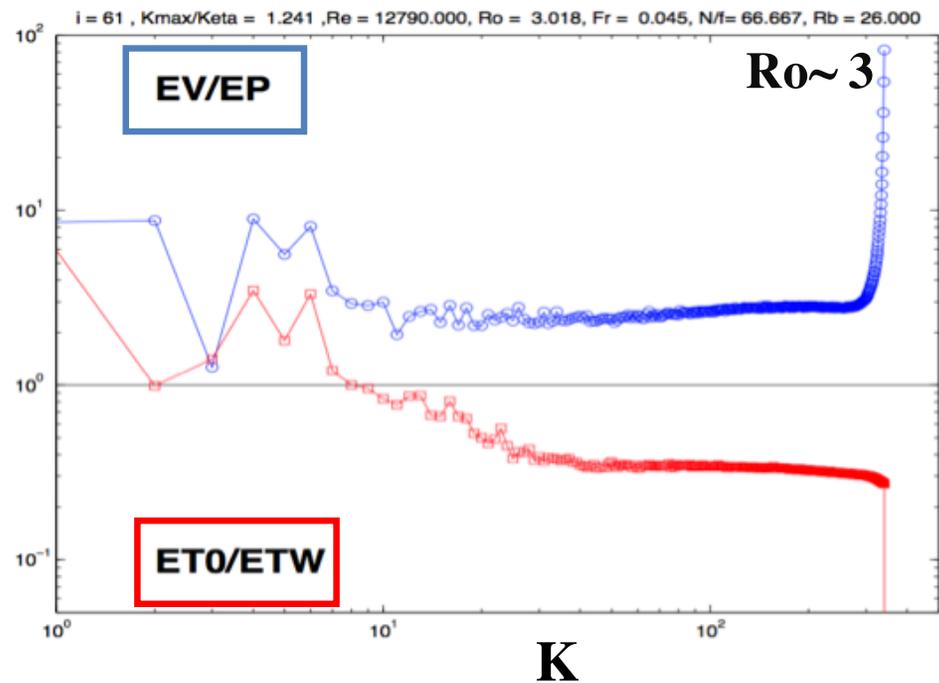
Ratios of energy spectra at peak:

Kinetic to potential

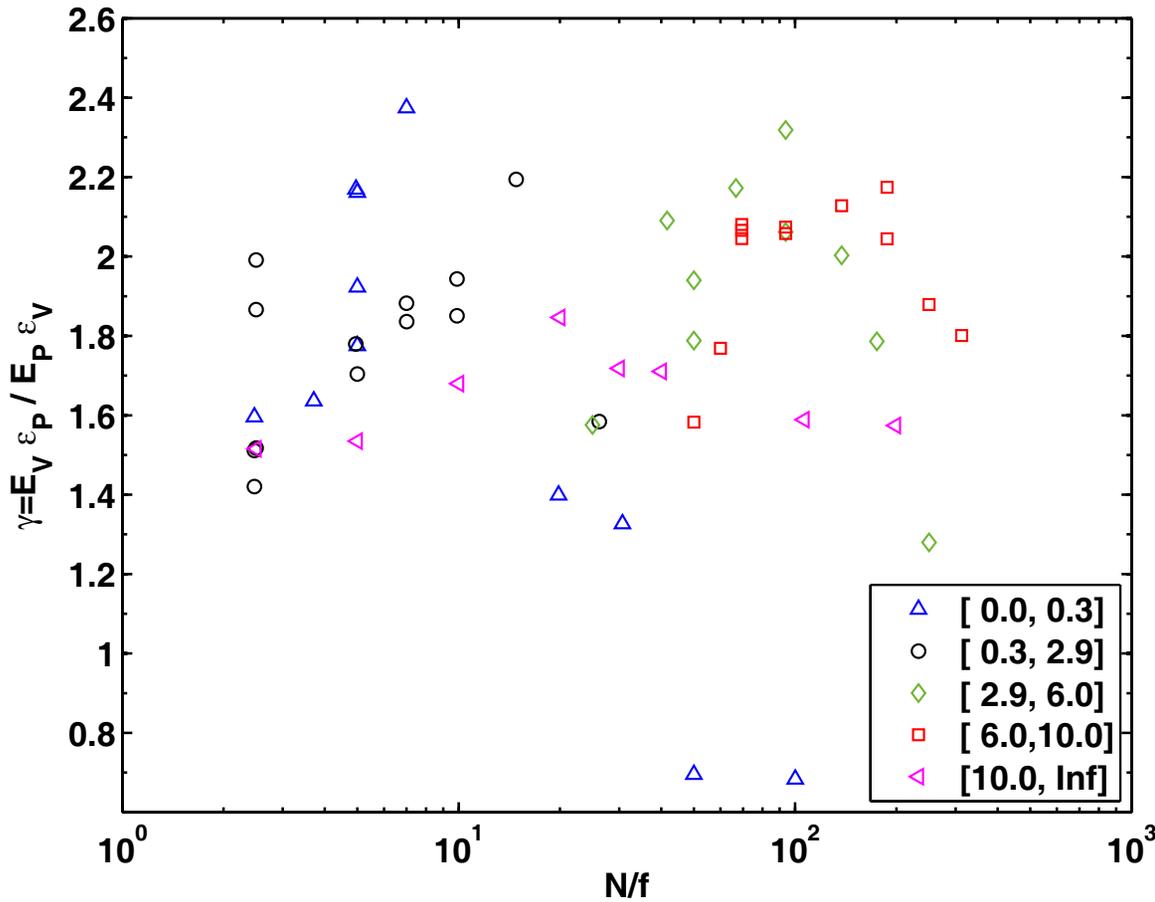
and

zero to wave mode

Both runs: $Fr = 0.04$, $R_B \sim 26$

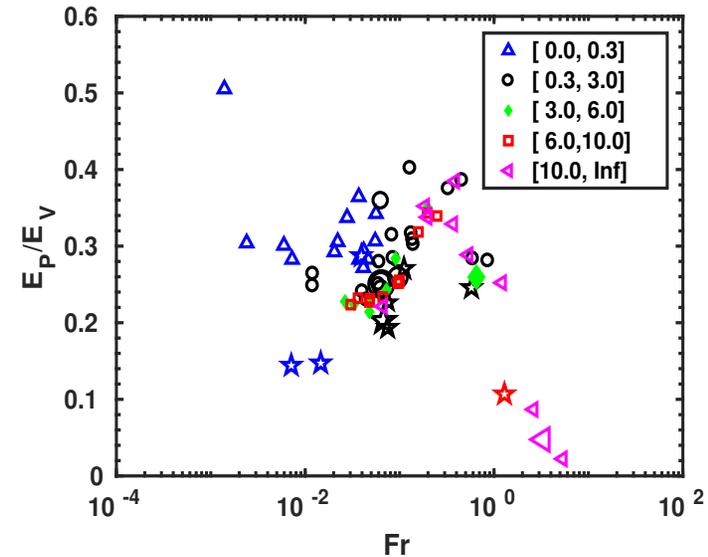


Color binning in Ro : $0 \rightarrow 0.3 \rightarrow 2.9 \rightarrow 6.0 \rightarrow 10 \rightarrow$



Ratio of kinetic to potential energy effective dissipation times

$$T_V/T_P = [E_V/E_P] * [\epsilon_P/\epsilon_V]$$



Low Fr:
influence of initial conditions

Vertical velocity
around peak of dissipation

Intermediate values:
 $w^2 / 2E_V \sim Fr^0$
 Or

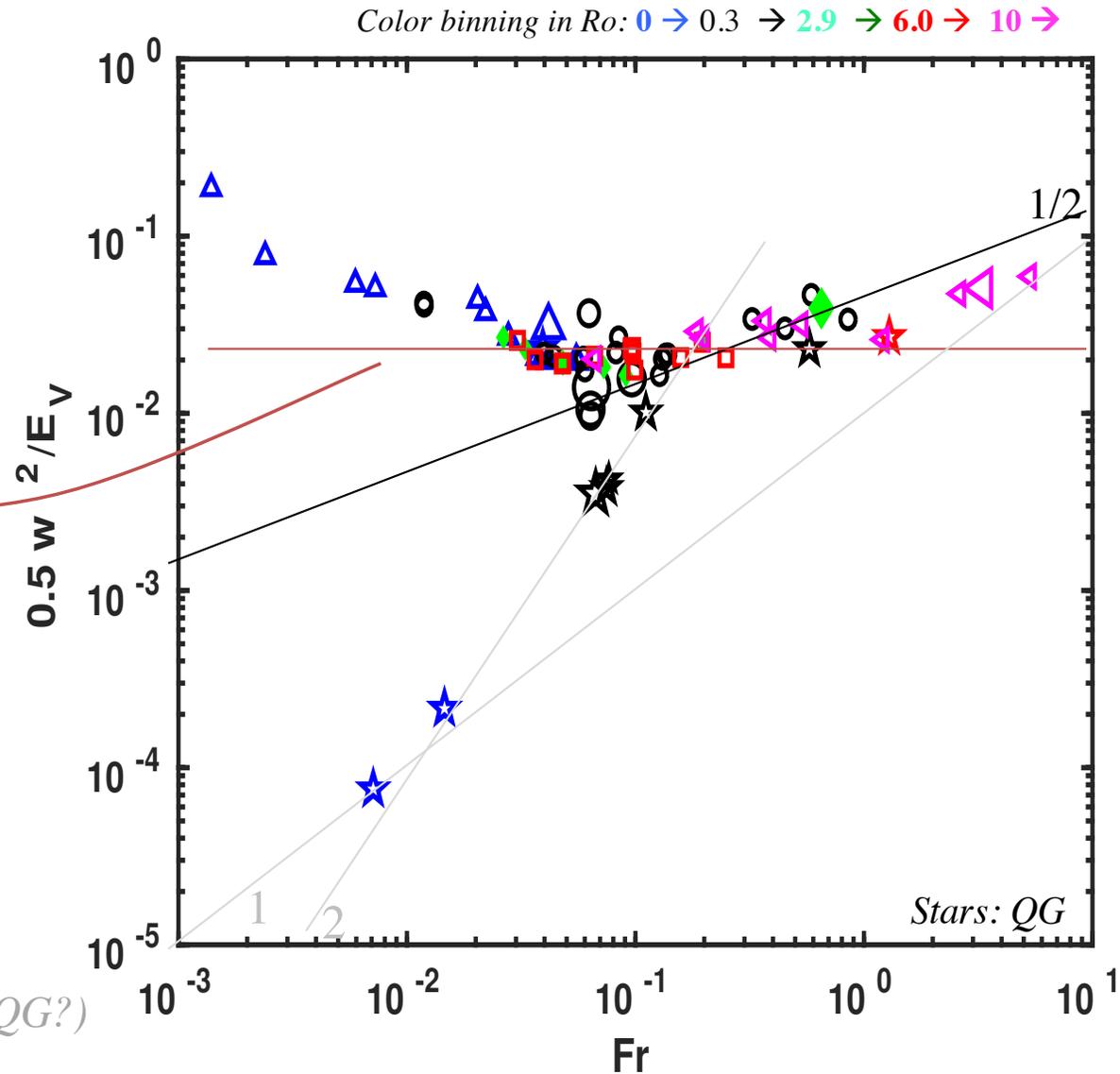
$$w/U_0 = a Fr^0 \quad (2a)$$

Other possible scaling:

$$w/U \sim Fr^{1/4} \quad (2b, \text{data?}),$$

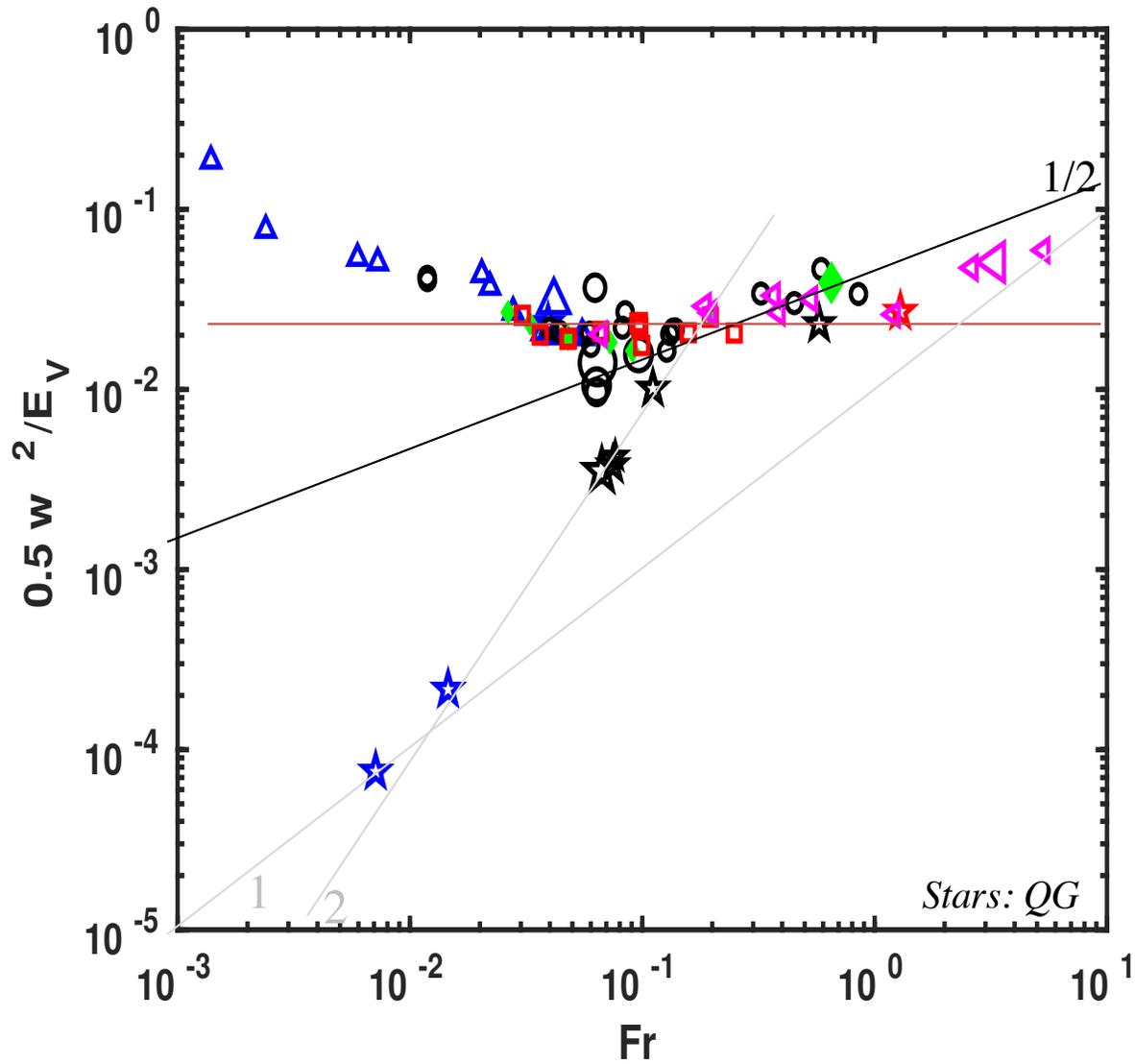
$$\text{or } w/U \sim Fr^{1/2} \quad (2c, \text{Maffioli+ 2016}),$$

$$\text{or } \sim Fr \quad (2d, \text{incompressibility; \& QG?})$$

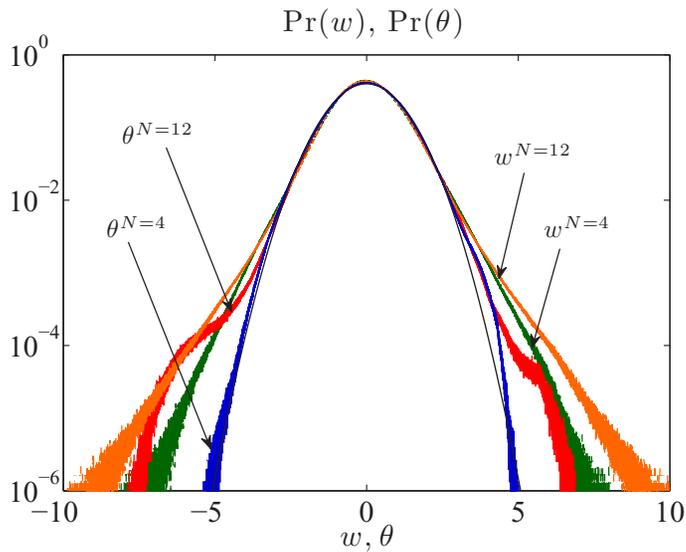


Vertical velocity

around peak of dissipation



Rorai+ PHYSICAL REVIEW E **89**, 043002 (2014)
2048³ forced



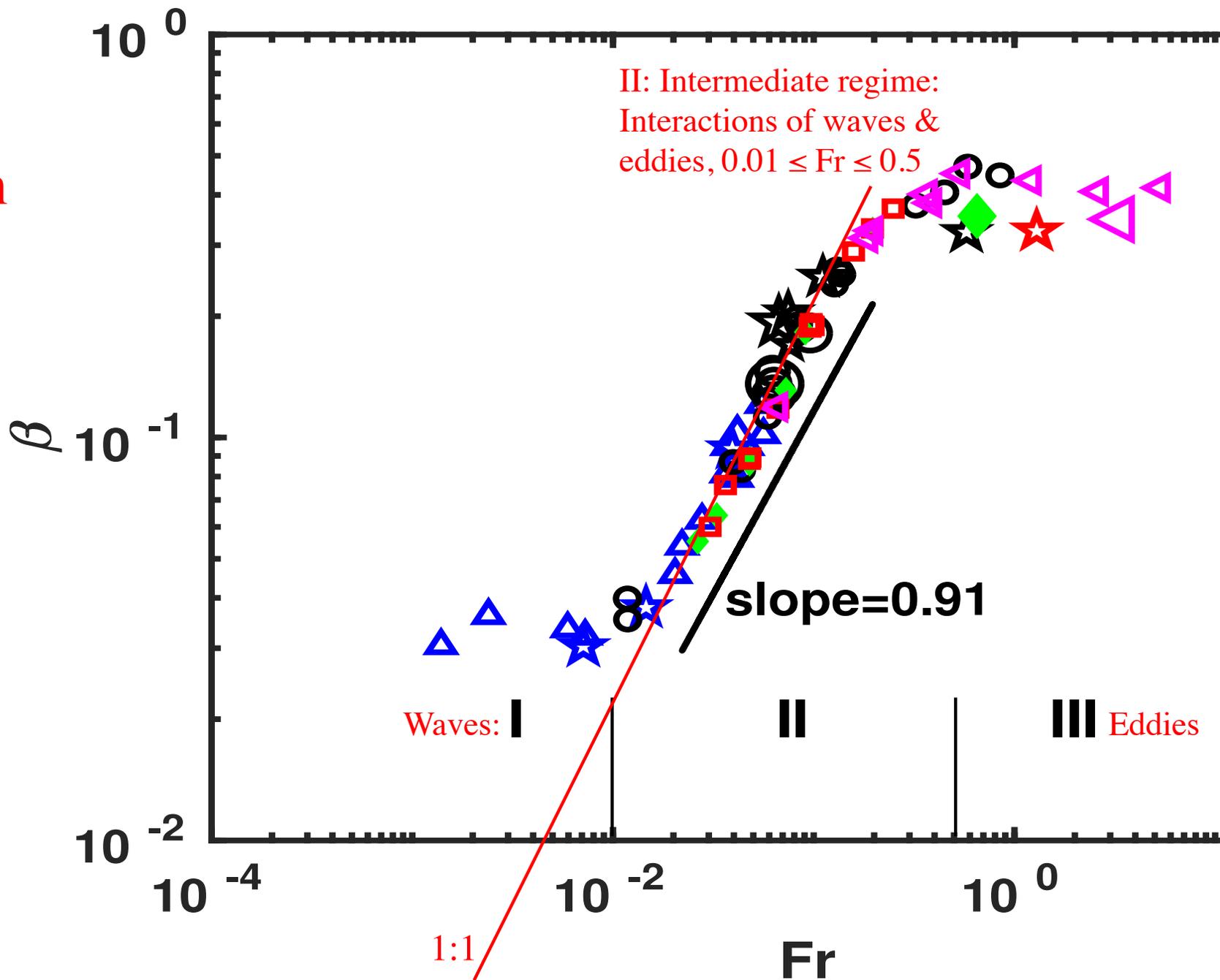
also Feraco+ 2018

Color binning in Ro : 0→0.3→2.9→6.0→10→

Normalized
kinetic
energy
dissipation
rate

$$\beta = \varepsilon_V / \varepsilon_D$$

→ 3
regimes

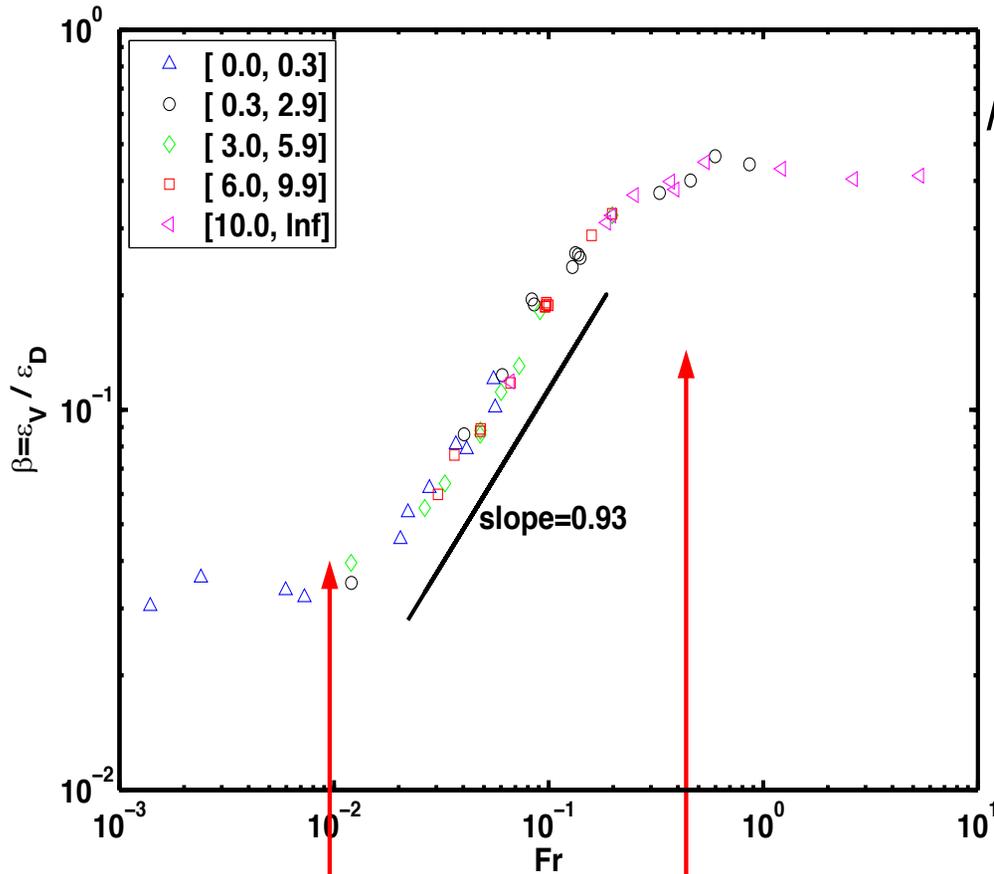


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Normalized kinetic energy dissipation rate

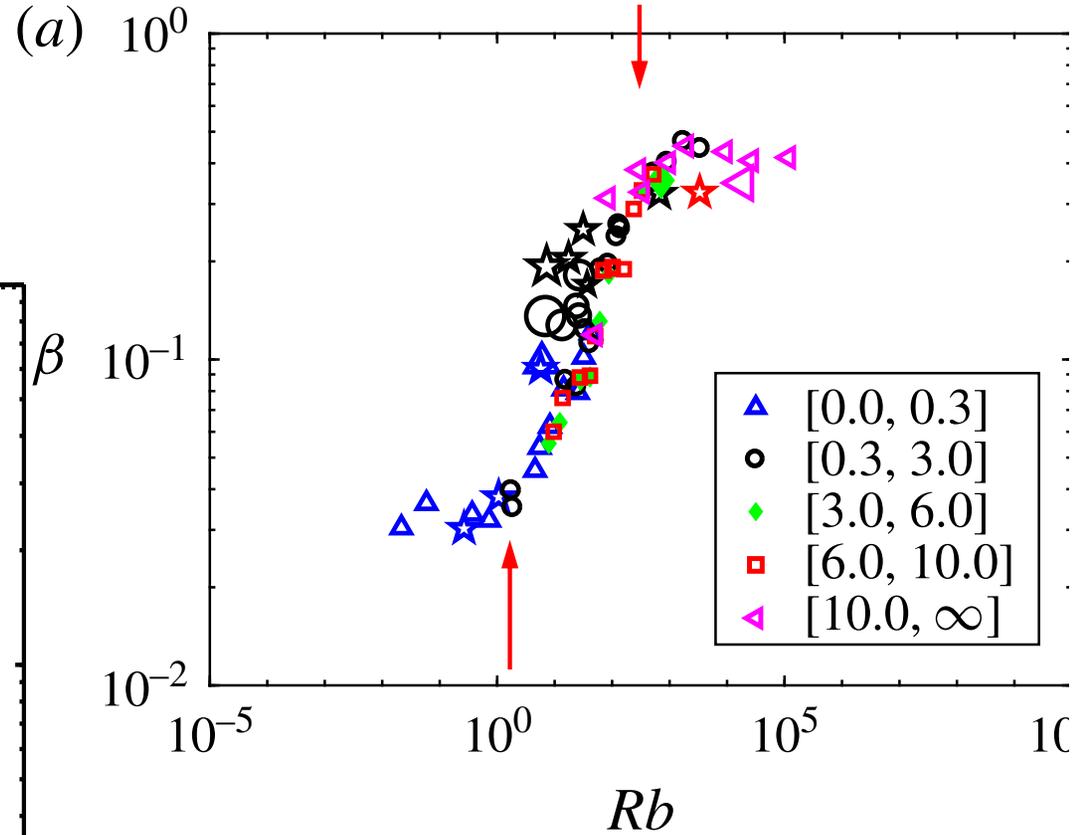
$\varepsilon_V/\varepsilon_D \rightarrow 3$ regimes

No QG, no low Re runs



II: Intermediate regime:

$$0.010 \leq Fr \leq 0.5$$



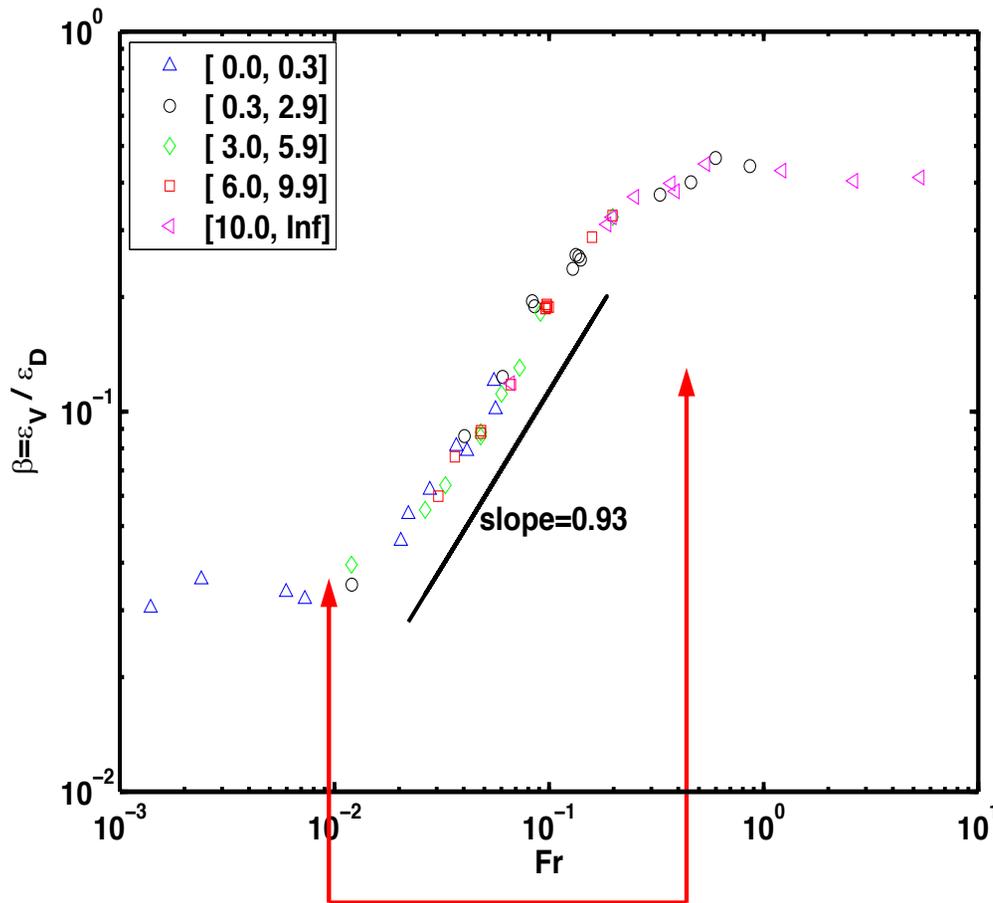
$$\text{or } 1 \leq R_B \leq 200$$

Color binning in Ro : $0 \rightarrow 0.3 \rightarrow 2.9 \rightarrow 6.0 \rightarrow 10 \rightarrow$

Normalized kinetic energy dissipation rate

$\varepsilon_V/\varepsilon_D \rightarrow 3$ regimes

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II: Intermediate regime:
 $0.010 \leq Fr \leq 0.5$

Classical model of **weak**,
wave turbulence: energy
transfer **slower** than eddy

$$\tau_{NL} = L_0 / U_0:$$

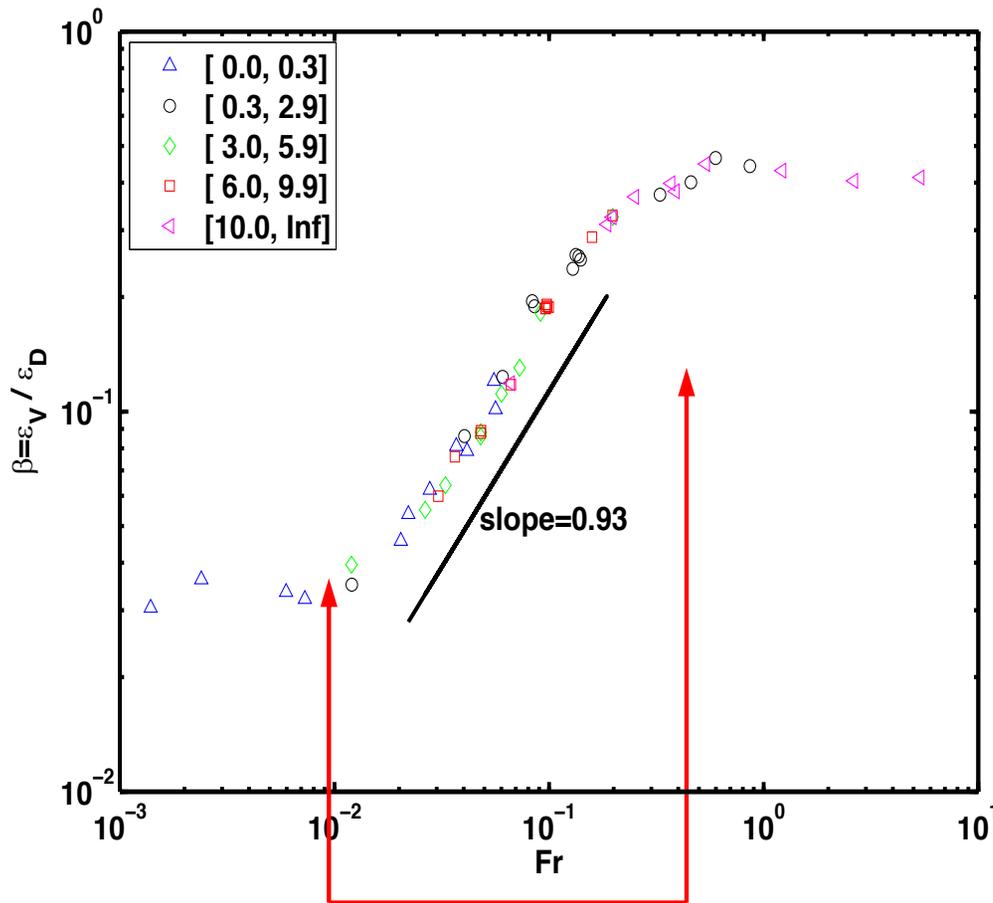
$$\tau_{\text{transfer}} \sim \tau_{NL} / Fr$$

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$$\tau_{NL} = L_0 / U_0:$$

$$\tau_{transfer} \sim \tau_{NL} / Fr$$

$$\sim \tau_{NL} * [\tau_{NL} / \tau_W]$$

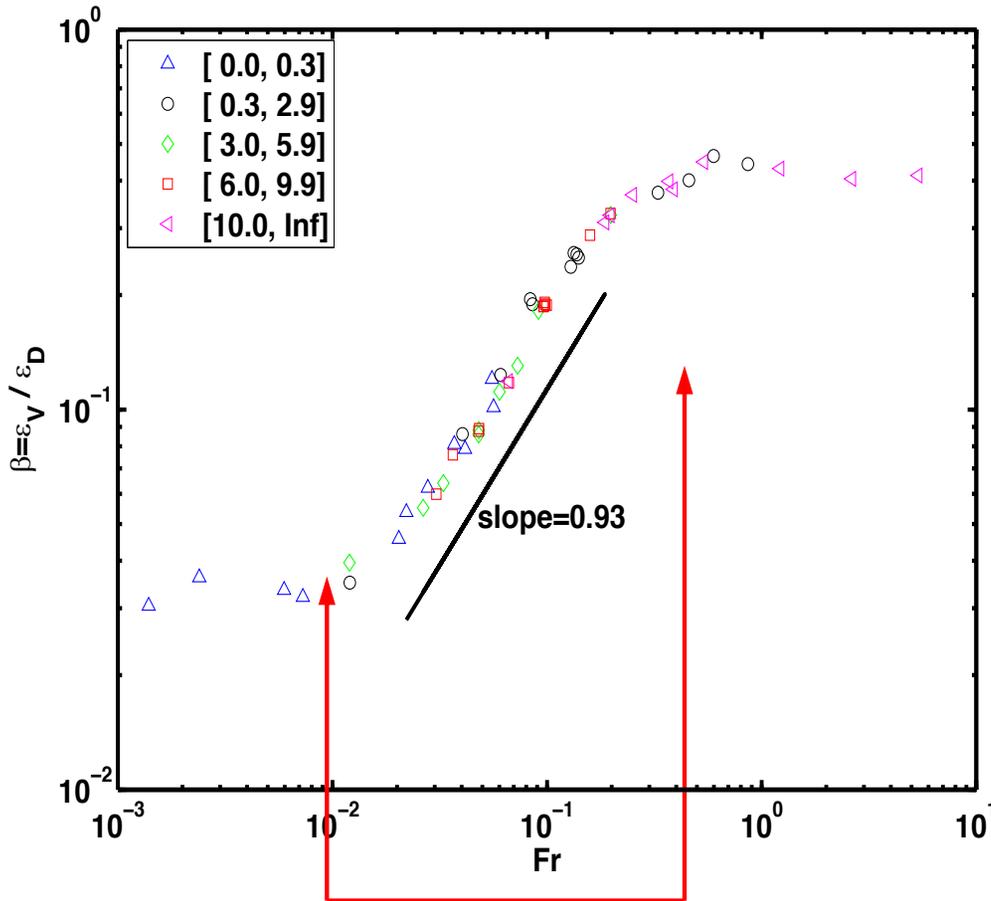
(MHD: Iroshnikov-Kraichnan, ~ '60s,
 Zakharov+ '80s, weak/wave turbulence)

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Normalized kinetic energy dissipation rate

$\varepsilon_V/\varepsilon_D \rightarrow 3$ regimes

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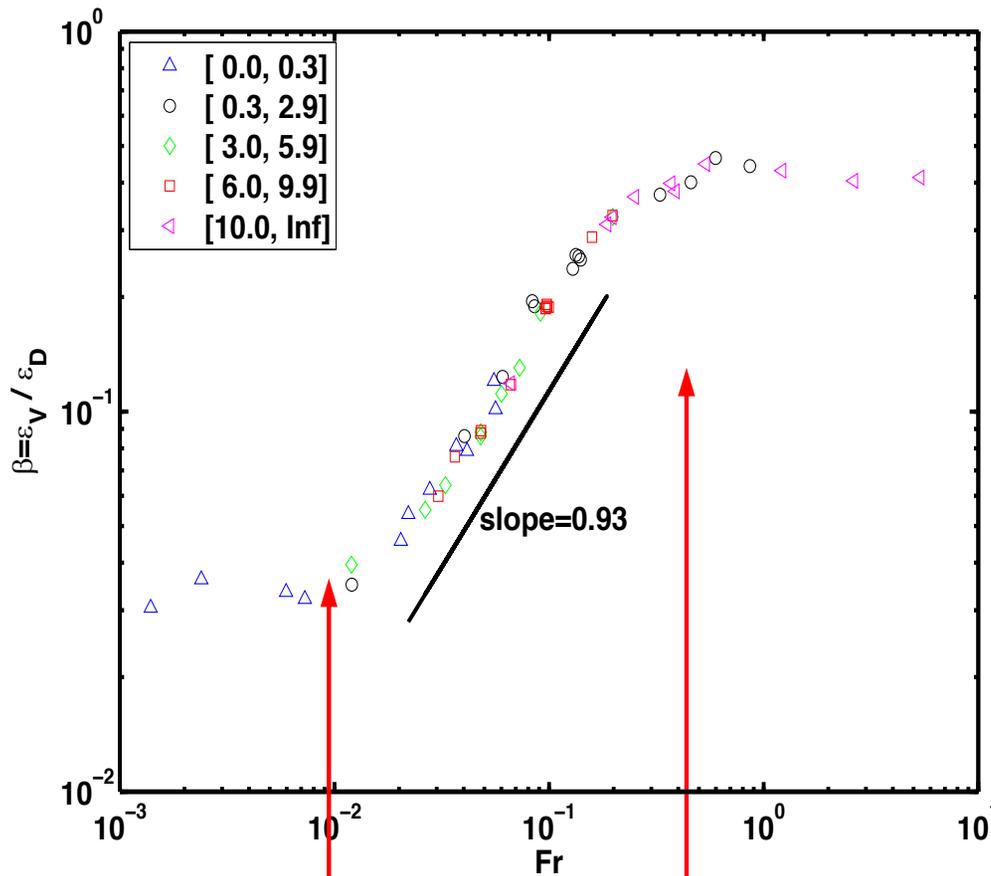
$$\varepsilon_V = dE_V/dt \sim E_V / \tau_{transfer}$$

Color binning in Ro : $0 \rightarrow 0.3 \rightarrow 2.9 \rightarrow 6.0 \rightarrow 10 \rightarrow$

Normalized kinetic energy dissipation rate

$\varepsilon_V/\varepsilon_D \rightarrow 3$ regimes

No QG, no low Re runs



II: Intermediate regime:
 $0.010 \leq Fr \leq 0.5$

Classical model of **weak**,
wave turbulence: energy
transfer **slower** than eddy

$$\tau_{NL} = L_0 / U_0:$$

$$\tau_{transfer} \sim \tau_{NL} / Fr$$

$$\sim \tau_{NL} * [\tau_{NL} / \tau_W]$$

(MHD: Iroshnikov-Kraichnan, ~ '60s,
Zakharov+ '80s, weak/wave turbulence)

$$\varepsilon_V = dE_V/dt \sim E_V / \tau_{transfer}$$

$$\varepsilon_V \sim Fr \varepsilon_D = Fr U_0^3 / L \quad (3)$$

Or: $\beta = \varepsilon_V / \varepsilon_D \sim Fr$ *as ~ observed*

Role of anisotropy?

Three laws → Energy balance & mixing: Γ_f, R_f & Γ_D, R_D

$$D_t E_v = -B_f + \varepsilon_v$$

$$D_t E_p = +B_f + \varepsilon_p$$

$$B_f = \langle N\theta w \rangle : \textit{Vertical buoyancy flux}$$

Three laws → Energy balance: Γ_f, R_f ; & Γ_D, R_D

$$D_t E_v = -B_f + \varepsilon_v$$

$$D_t E_p = +B_f + \varepsilon_p$$

$B_f = \langle N\theta w \rangle$: *Vertical buoyancy flux*

$$E_v = \frac{1}{2} \langle |u|^2 \rangle, E_p = \frac{1}{2} \langle \theta^2 \rangle$$

$$\varepsilon_v = \nu \langle \omega^2 \rangle, \varepsilon_p = \kappa \langle |\text{grad } \theta|^2 \rangle, E_T = E_v + E_p, \varepsilon_T = \varepsilon_v + \varepsilon_p$$

$B_f / \varepsilon_v = \Gamma_f = R_f / [1 - R_f]$: *Mixing efficiency* *(momentum equation)*

$R_f = B_f / [B_f + \varepsilon_v]$: *Flux Richardson number* \in [0,1]

$\varepsilon_p / \varepsilon_v = \Gamma_D = R_D / [1 - R_D]$, $R_D = \varepsilon_p / \varepsilon_T$ \in [0,1] *(coupled equations)*

Scaling?

$$D_t E_v = -B_f + \varepsilon_v$$

$$D_t E_p = +B_f + \varepsilon_p$$

$B_f = \langle N\theta w \rangle$: *Vertical buoyancy flux*

$$E_v = \frac{1}{2} \langle |u|^2 \rangle, E_p = \frac{1}{2} \langle \theta^2 \rangle$$

$$\varepsilon_v = \nu \langle \omega^2 \rangle, \varepsilon_p = \kappa \langle |\text{grad } \theta|^2 \rangle, E_T = E_v + E_p, \varepsilon_T = \varepsilon_v + \varepsilon_p$$

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$R_f = B_f / [B_f + \varepsilon_v]$: *Flux Richardson number*

$\varepsilon_p / \varepsilon_v = \Gamma_D = R_D / [1 - R_D]$, $R_D = \varepsilon_p / \varepsilon_T \in [0, 1]$ (coupled equations)

High R_B : $\Gamma_f = 0.2$ (Osborn-Cox '80) vs. $\Gamma_f \sim R_B^{-1/2}$ (Lozovatsky & Fernando 2013) vs. ?

→ Prediction of scaling for mixing efficiency, flux Richardson number and dissipation in the two regimes of wave-eddy interactions and of strong eddies

using the three constitutive scaling laws for temperature, vertical velocity and dissipation efficiency versus Froude number:

$$\Theta_{\text{rms}} \sim U_0 \quad (1)$$

$$w/U_0 = a \text{Fr}^0 \quad (2a)$$

$$\varepsilon_v \sim \text{Fr} \varepsilon_D = \text{Fr} U_0^3/L \quad (3)$$

[0 → ¼? (2b) Or → 1 (2d)?]

Intermediate regime II, $Fr \leq 1$, using the 3 scaling laws:

$$B_f = \langle N\theta w \rangle, \theta_{\text{rms}} \sim U_0, w \sim Fr^0 U_0, \varepsilon_v \sim Fr \varepsilon_D, \varepsilon_D \sim U_0^3/L$$

$R_B > 1, Re \gg 1$ but irrelevant otherwise

$$\rightarrow \Gamma_f^{\text{II}} = B_f/\varepsilon_v = N\langle w \theta \rangle/\varepsilon_v \sim 1/Fr^2 \sim [R_B]^{-1} \quad (\text{observed})$$

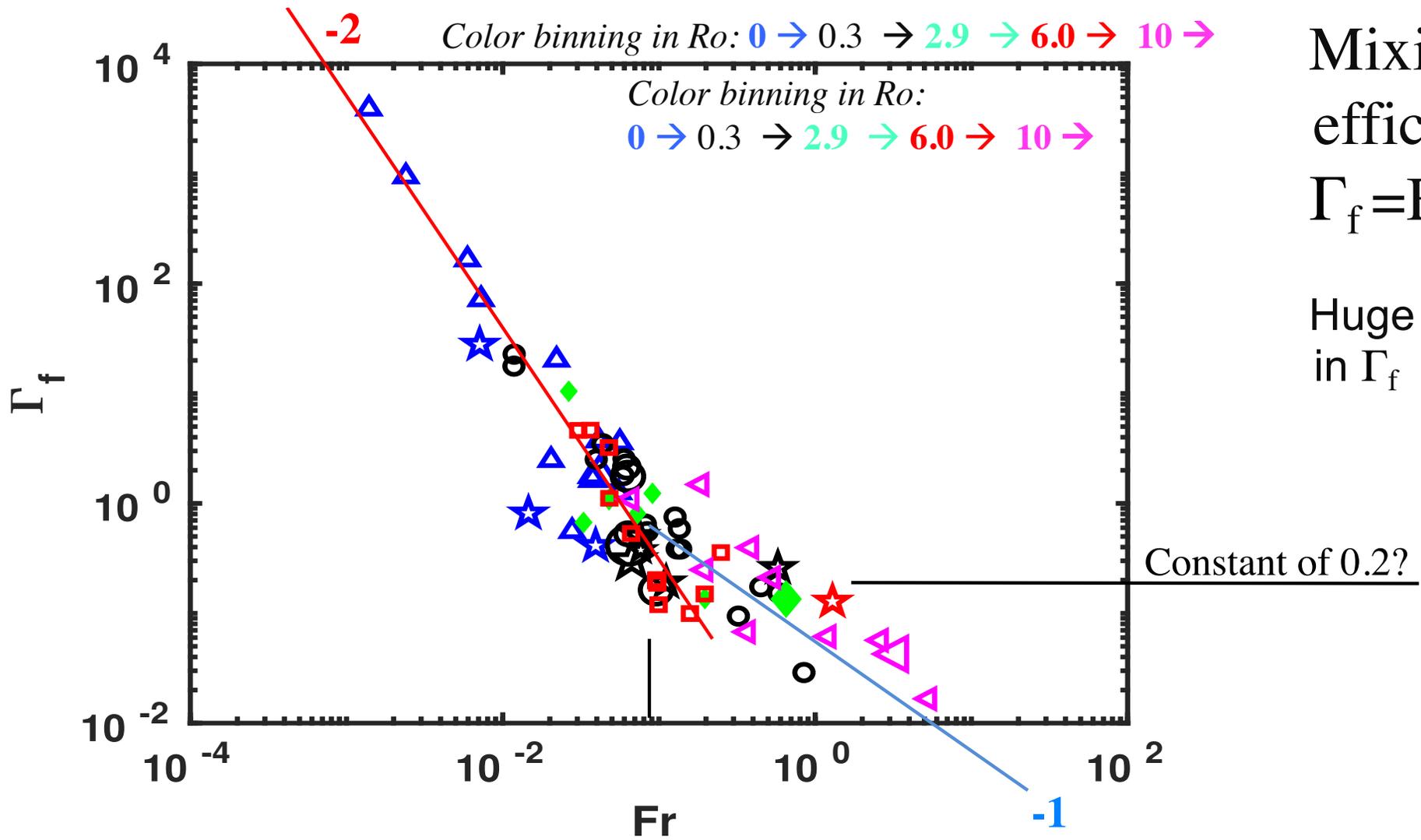
Higher regime III, $Fr \geq 1$:

$$\theta_{\text{rms}} \sim U_0, w \sim Fr^0 U_0, Fr=1 = U_0, \varepsilon_v \sim Fr \varepsilon_D \text{ at } Fr = 1, \varepsilon_v \sim \varepsilon_D$$

$R_B > 1, Re \gg 1$ but irrelevant otherwise

$$\rightarrow \Gamma_f^{\text{III}} = B_f/\varepsilon_v = N\langle w \theta \rangle/\varepsilon_v \sim 1/Fr \sim [R_B]^{-1/2} \quad (\text{observed})$$

\rightarrow *Our numerical data?*

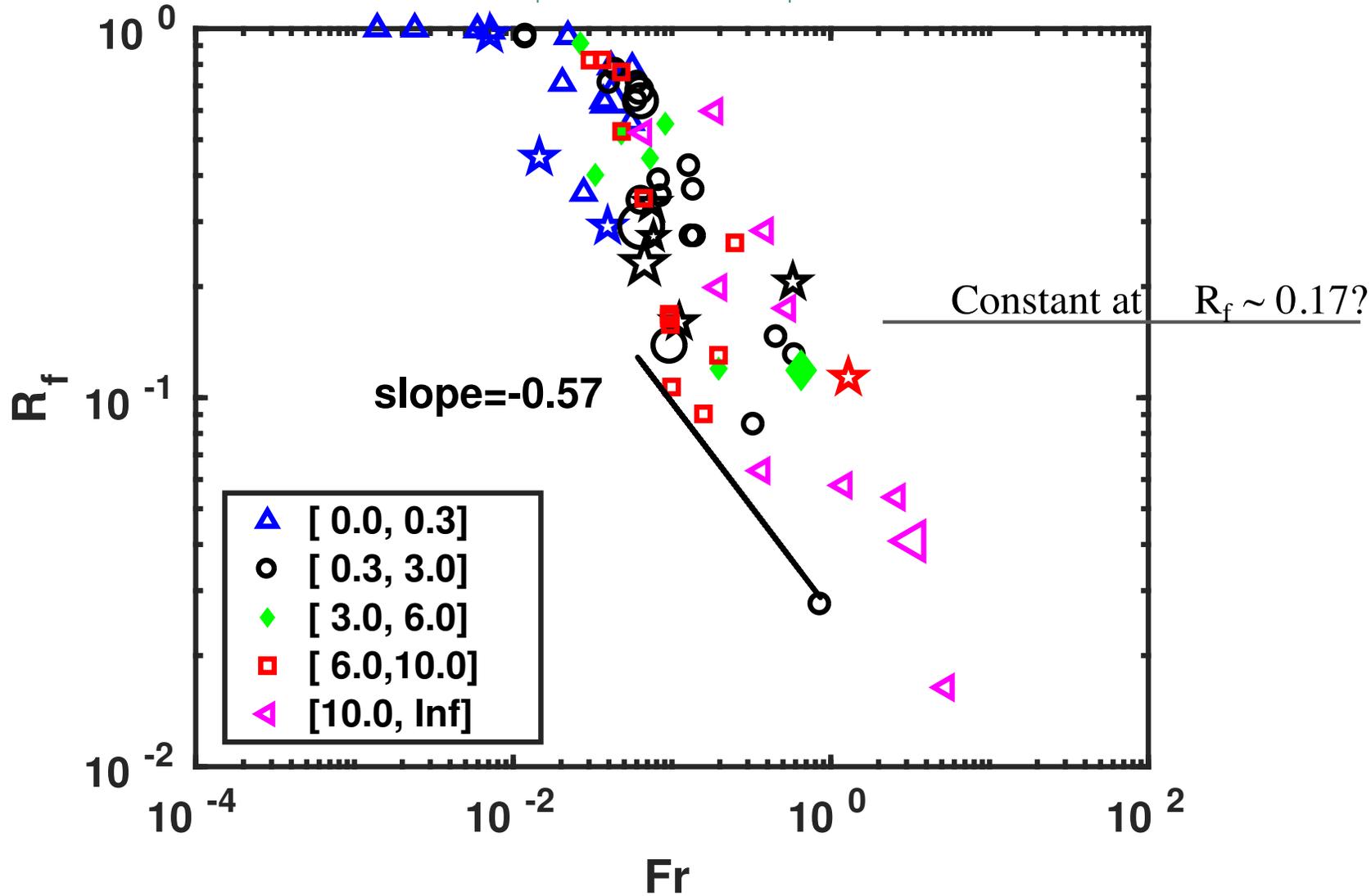


$$N \langle w \theta \rangle / \epsilon_v = \Gamma_f \sim 1 / Fr^2 \sim [R_B]^{-1} \text{ (regime I \& II?)}$$

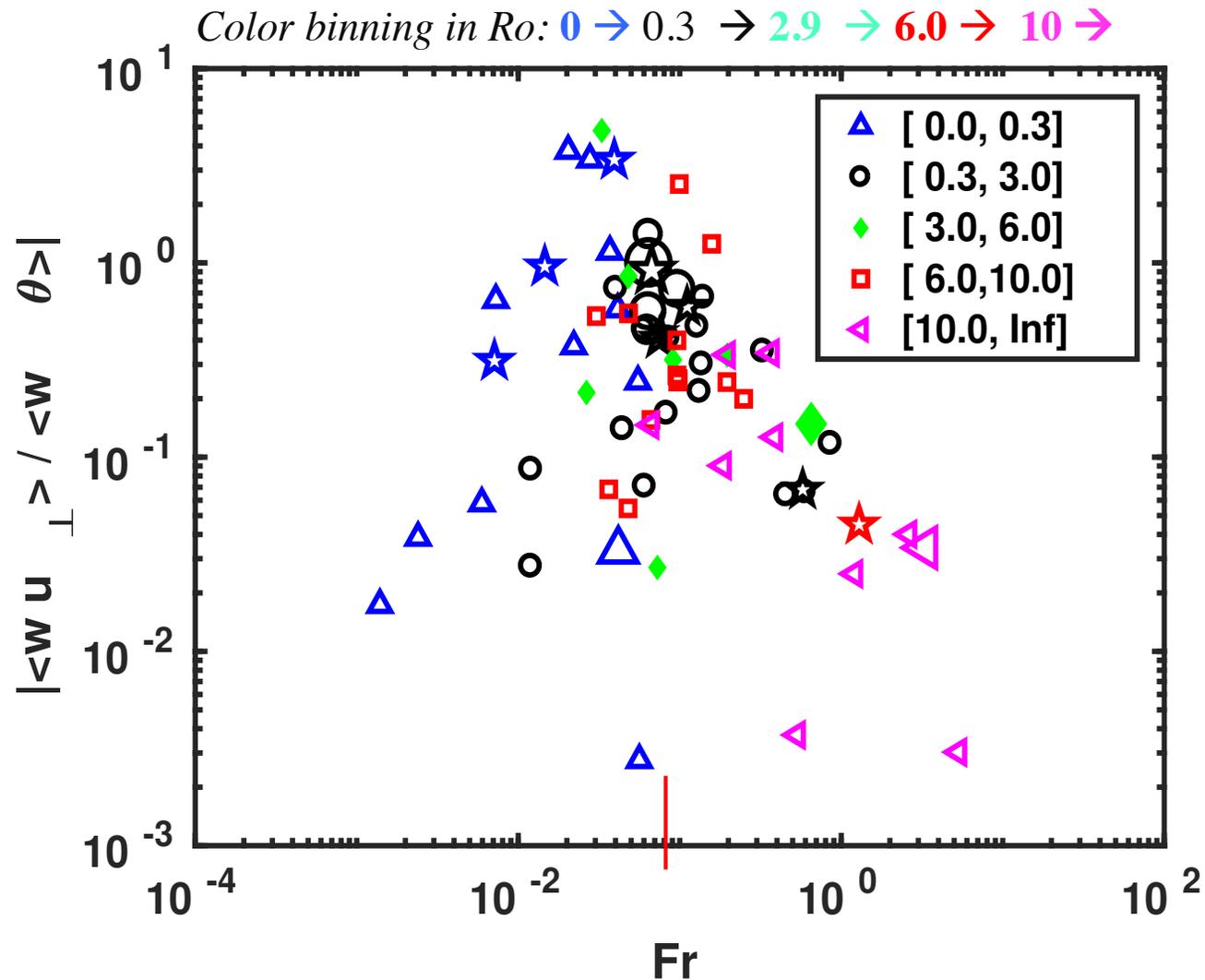
$$\text{or for } Fr \sim 1: \Gamma_f \sim 1 / Fr \sim [R_B]^{-1/2} \text{ (regime III)}$$

Flux Richardson number: $R_f = B_f / [B_f + \epsilon_v]$

Color binning in Ro : 0 \rightarrow 0.3 \rightarrow 2.9 \rightarrow 6.0 \rightarrow 10 \rightarrow



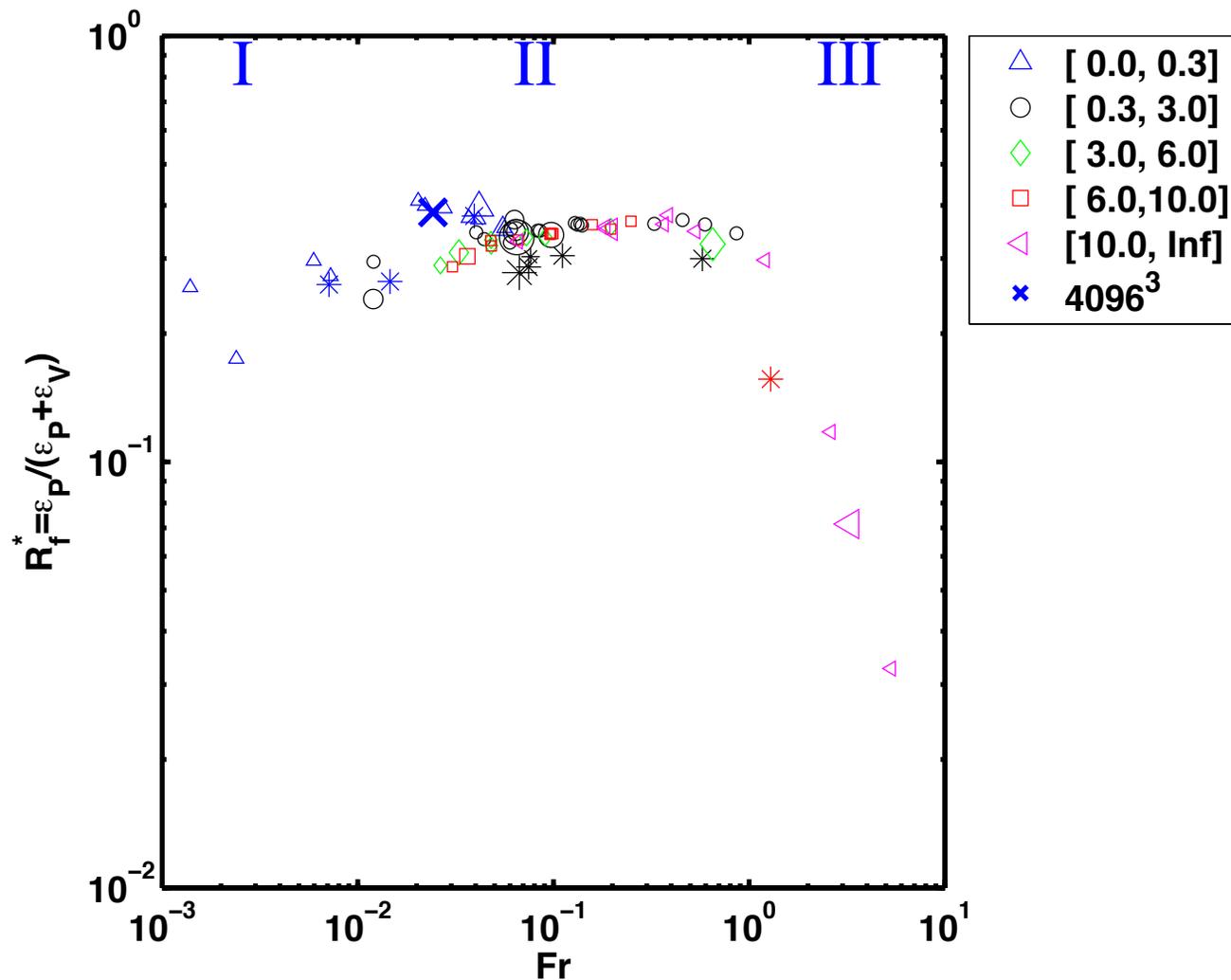
Ratio of vertical fluxes
of horizontal velocity
to that of
temperature fluctuations
versus
Froude Number



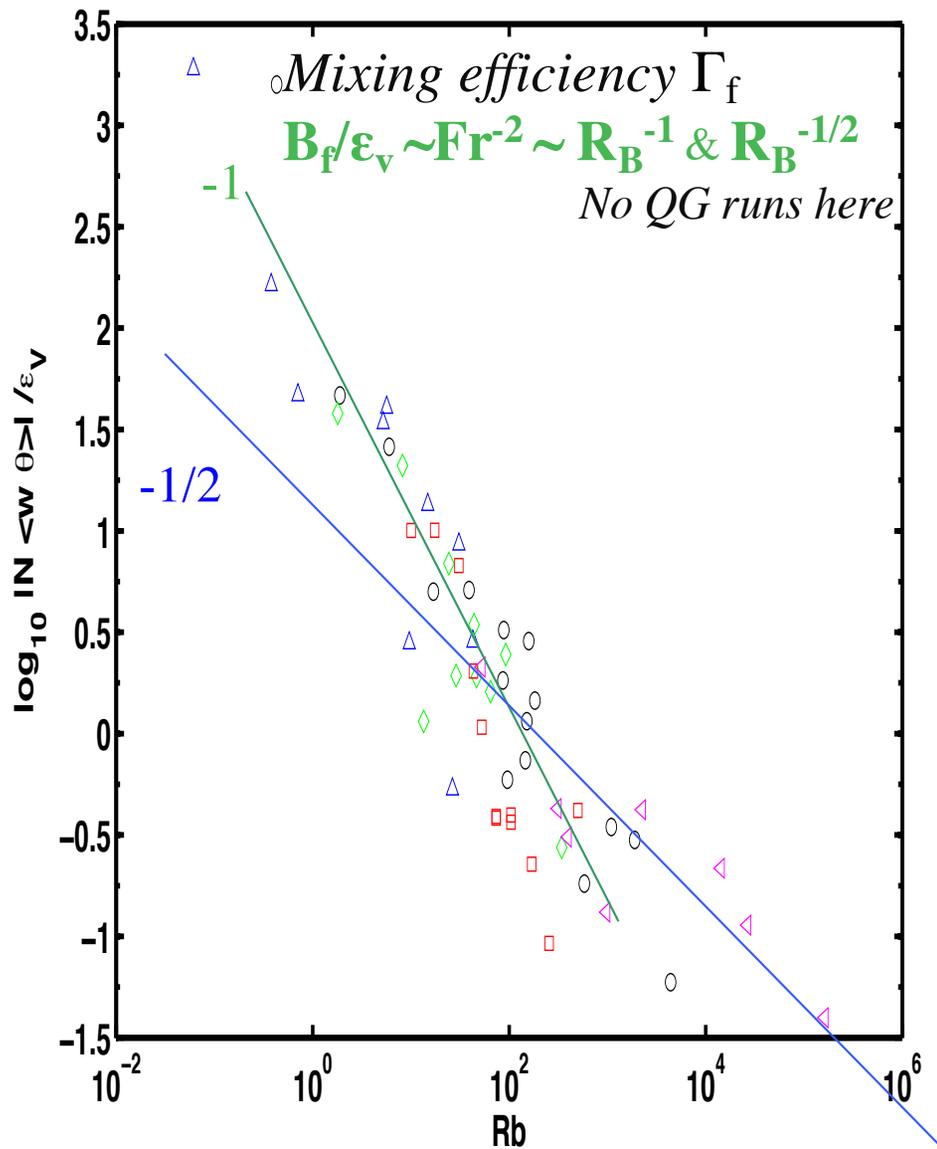
Non-monotonic, with a
plateau/peak around $Fr \sim 0.07$?

Ratio of potential to
total energy dissipation

$$\varepsilon_p / [\varepsilon_p + \varepsilon_v]$$



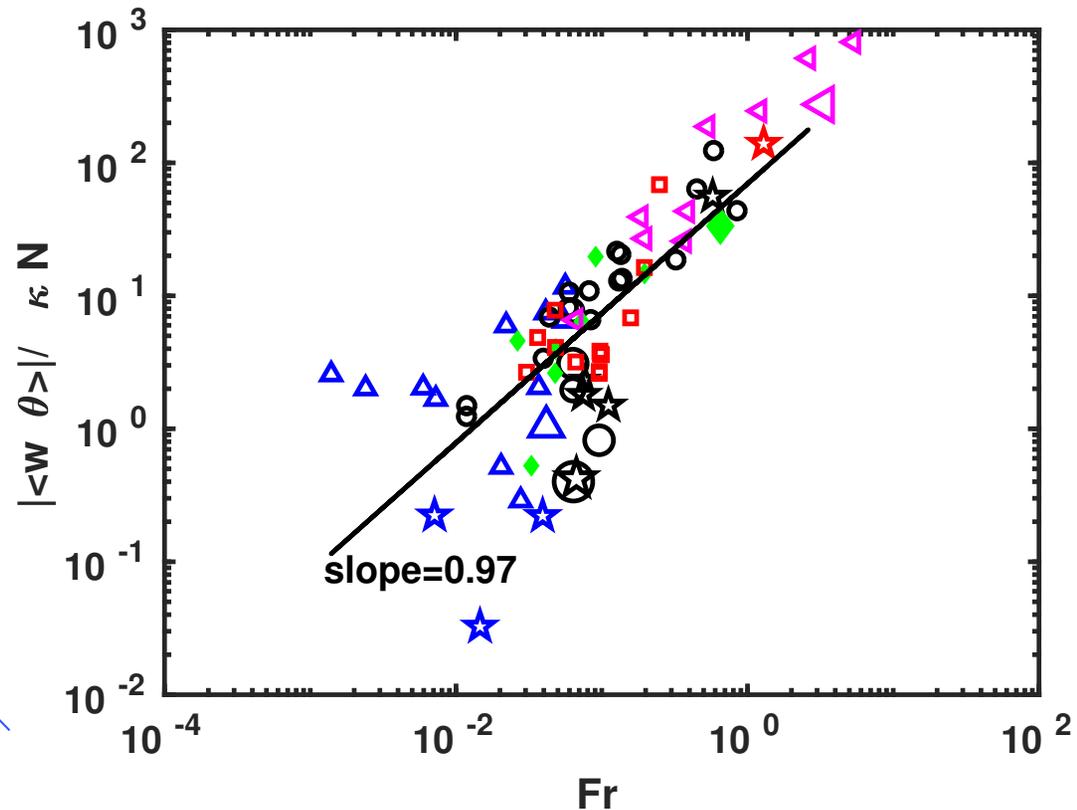
Color binning in Ro : 0 \rightarrow 0.3 \rightarrow 2.9 \rightarrow 6.0 \rightarrow 10 \rightarrow



Buoyancy flux $B_f = \langle N w \theta \rangle$
 with two different normalizations for the
 \leftarrow momentum and buoyancy equations

↓

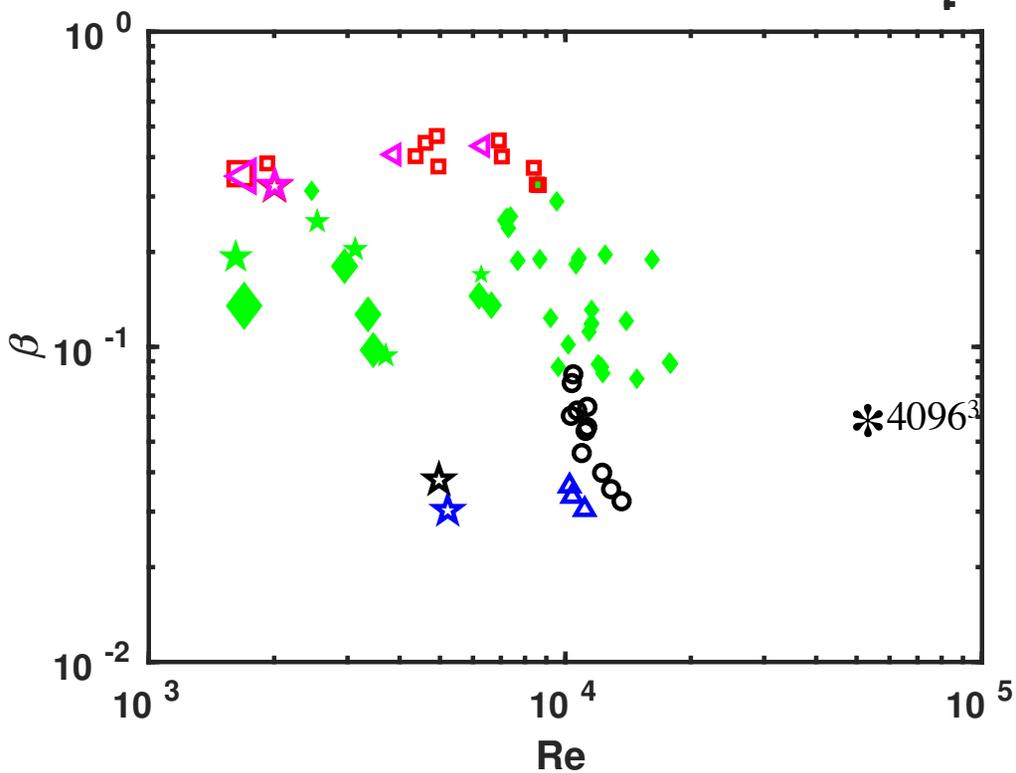
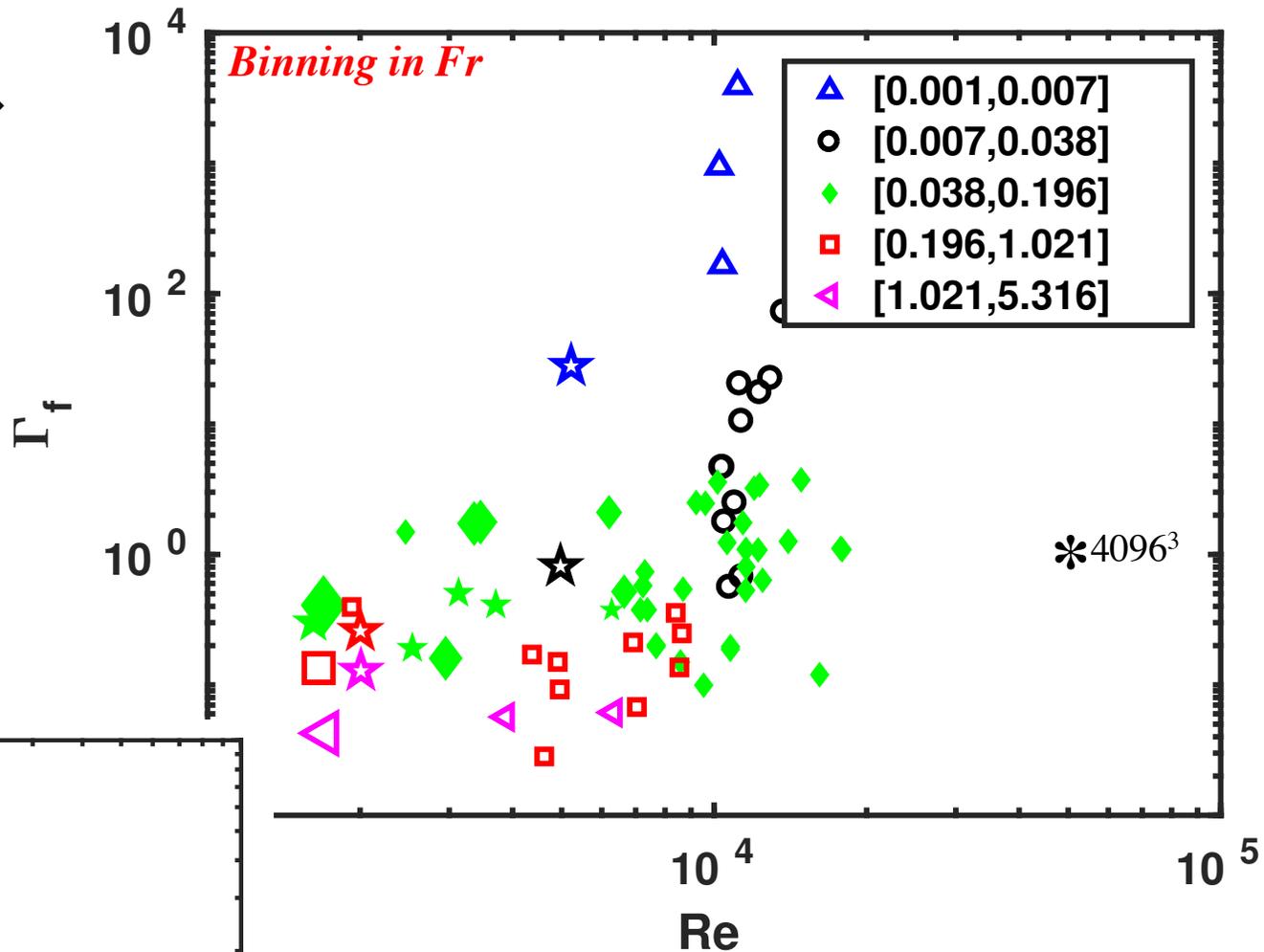
$$B_f / [\kappa_0 N^2] \sim Fr$$



Normalized anomalous diffusivity κ_θ
 $B_f / [\kappa N^2] = \kappa_\theta / \kappa \sim Fr$ in regimes II & III,
 $\kappa_\theta \sim U_{rms} L_{int} Fr$

Mixing efficiency \rightarrow
&
Dissipation efficiency \downarrow

versus
Reynolds Number



Questions and Perspectives

- Role of rotation?
- Role of $k_0 \sim 2.5$, poor large-scale statistics & weaker resonances?
- Importance of QG?
- Role of lack of stationarity due to lack of forcing?
→ Add forcing and large-scale friction → Temporal averaging
- But, what about anisotropy – dispersion relation and – forcing?
- Role of boundary conditions?
- Will small aspect ratio help, role of vertical shear?
- Approach through small-scale modeling: will *Artificial Intelligence* & *Machine Learning* help?

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- Approach through small-scale modeling: will *Artificial Intelligence & Machine Learning* help?

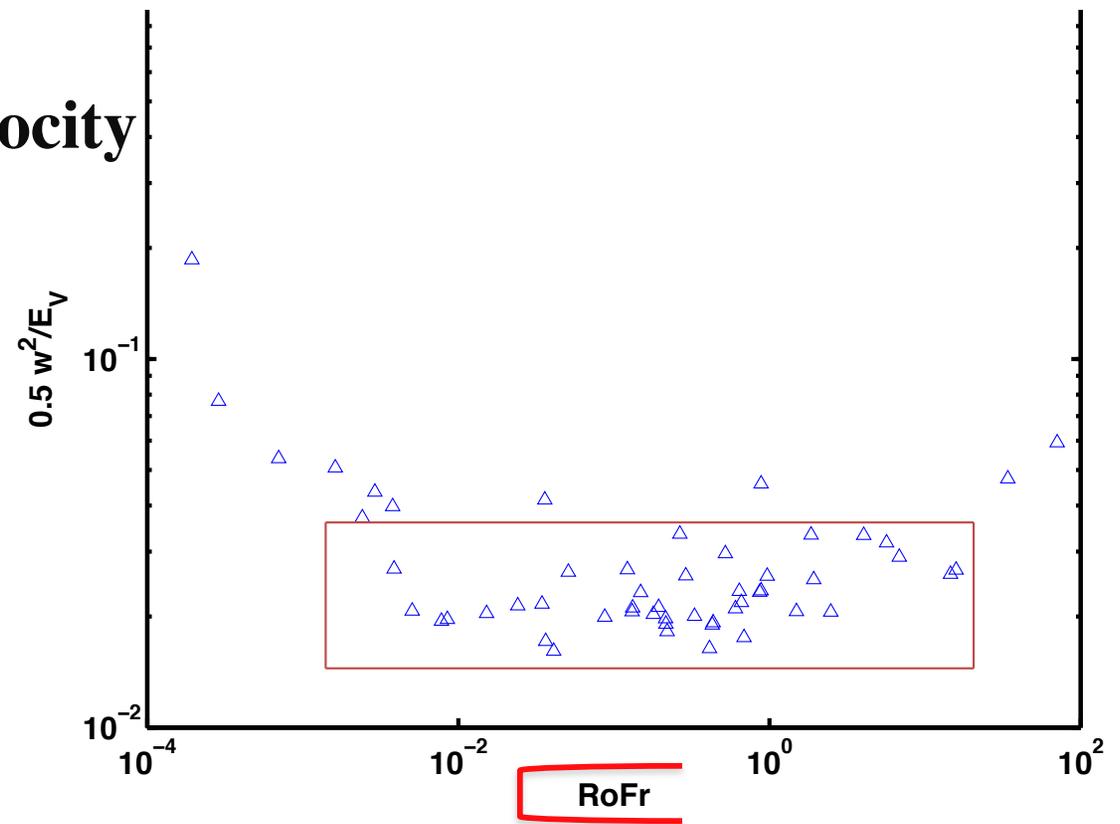
Perspective (1): Vertical velocity

in the presence of rotation

Intermediate regime:

$$w / U_0 \sim [\text{Ro.Fr}]^{-0} \quad (3b)$$

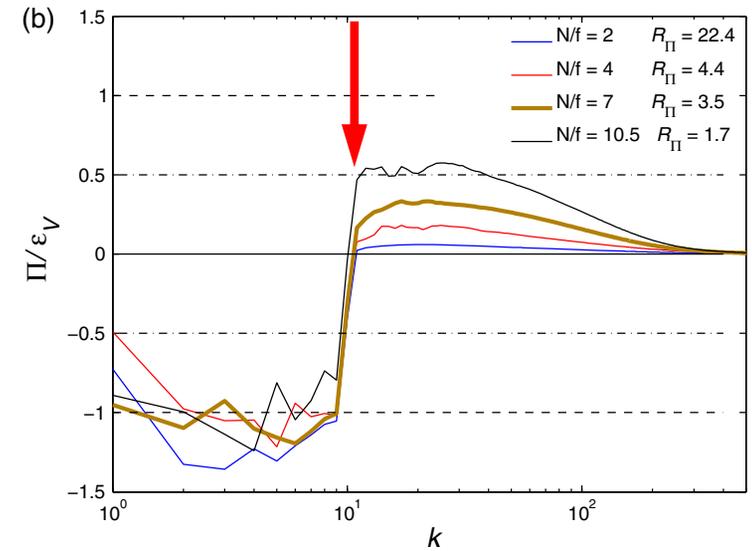
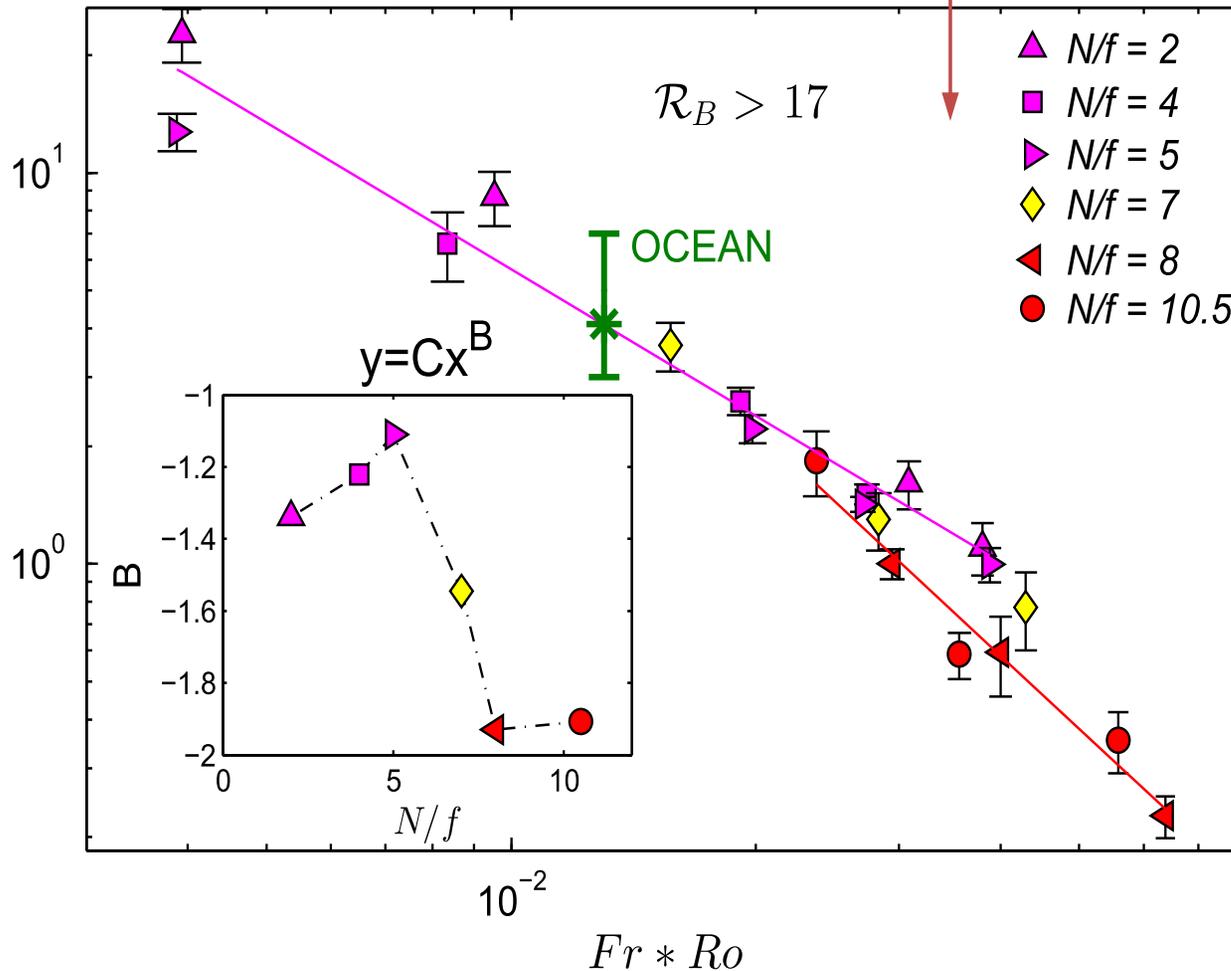
$$\rightarrow \Gamma_f = \langle N\theta w \rangle / \varepsilon_v \sim \text{Fr}^{-2}$$



Perspective (2): Forced runs

- Analyze runs of rotating stratified turbulence with **isotropic forcing** at intermediate scale ($k_f \sim 10$): what are their mixing properties?

Pouquet+ 2013 ; Marino+ 2015 (rot+strat)



$$R_{\Pi} = \varepsilon_{LS} / \varepsilon_{ss}$$

$$\varepsilon_{ss} \sim \varepsilon_D Fr, \quad \varepsilon_D \sim U^3 / L$$

$$\varepsilon_{LS} \sim [1/Ro] \varepsilon_D$$

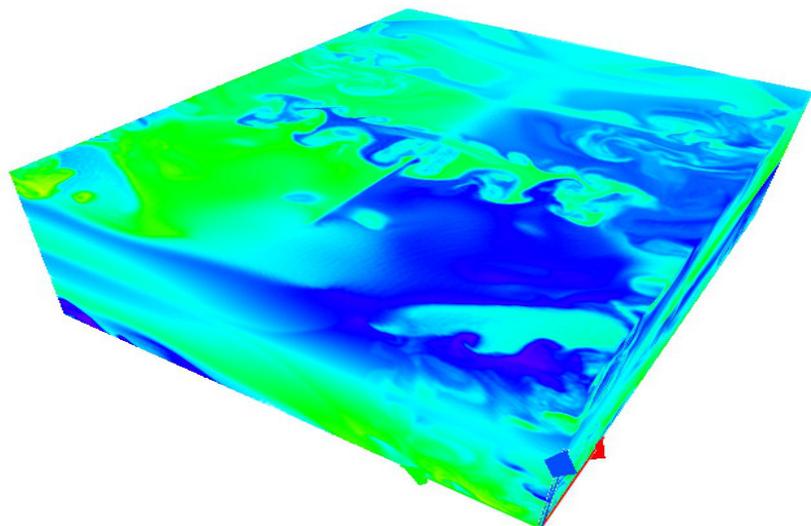
$$\rightarrow R_{\Pi} = \varepsilon_{LS} / \varepsilon_{ss} \sim [Ro \cdot Fr]^{-1}$$

with transition at $Ro \sim 0.45$

Perspective (3): Shallow fluids

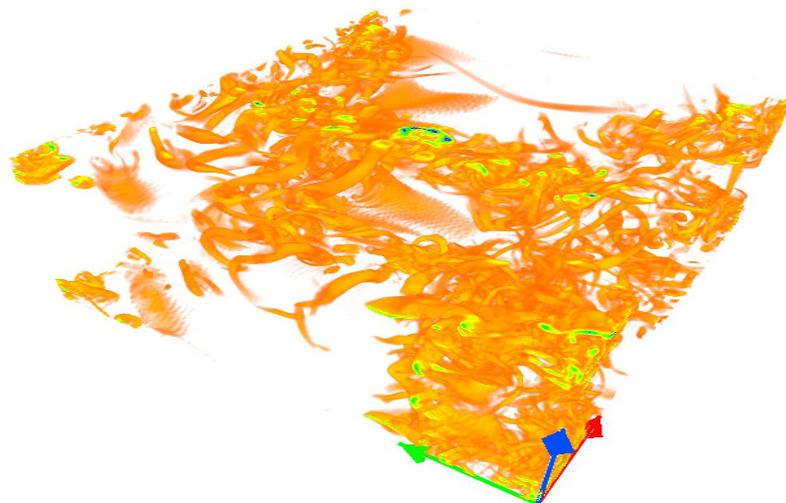
- Explore the dynamics in match boxes:
- Example; $2048^2 \times 256$ points with Taylor-Green forcing, resulting locally in strong vertical shear, and strong dissipation, with fronts and filaments

Temperature



&

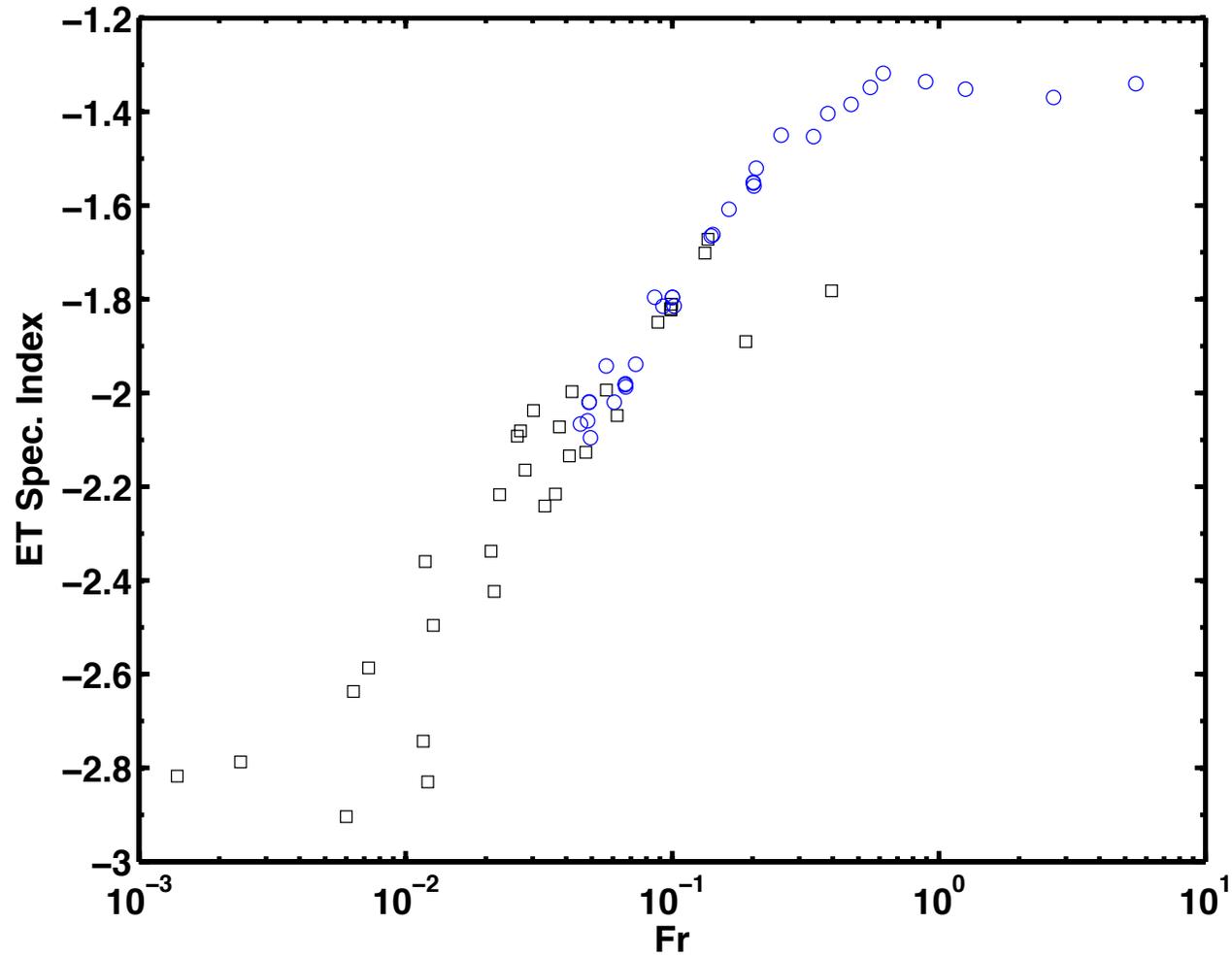
vorticity (*Oks+*, [arXiv 1706.10287](https://arxiv.org/abs/1706.10287))



Summary and perspectives:

- There are three distinct regimes in rotating stratified turbulence, determined by the Froude number for high enough Reynolds number, with somewhat different thresholds when analyzing different fields: Strong waves, eddy-wave interactions and full turbulence
- Mixing efficiency and flux Richardson number vary measurably in the intermediate regime, $0.01 \leq Fr \leq 0.5$ ($1 \leq R_B \leq 1000$), and with scaling laws
- Dissipation increases as Froude, “weak” turbulence regime, for $R_B \approx 10-10^3$
- Lower values of mixing efficiency at high R_B → *the velocity and temperature fluctuations are only weakly coupled (passive scalar regime)*
- Together, large numerical resolution and large parametric study allow for some scale separation and for understanding separately the roles of some of the players
- Local instabilities and very intense local small-scale dynamics (*more so than in FDT*)
- Role of Reynolds number in mixing, in the intermediate regime in particular?
- Roles of I.C. or forcing (3D vs. 2D, θ or not, vortices vs. waves, balanced or not ...), of non-local interactions, of large-scale instabilities & of large-scale friction?
- Intermittency high kurtosis and non-monotonicity ([Rorai+ 2014](#); [Feraco+ ArXiv 2018](#))

ET Spec. Index v. Fr, $k_{max}/k_{eta} > 1.25$



Spectral index of total energy spectra:

Again 3 regimes and going smoothly from “-3” to “-5/3”

Machine learning and modeling

- Use data to guess the functional form for the role of small scales in order to write sub-grid scale models

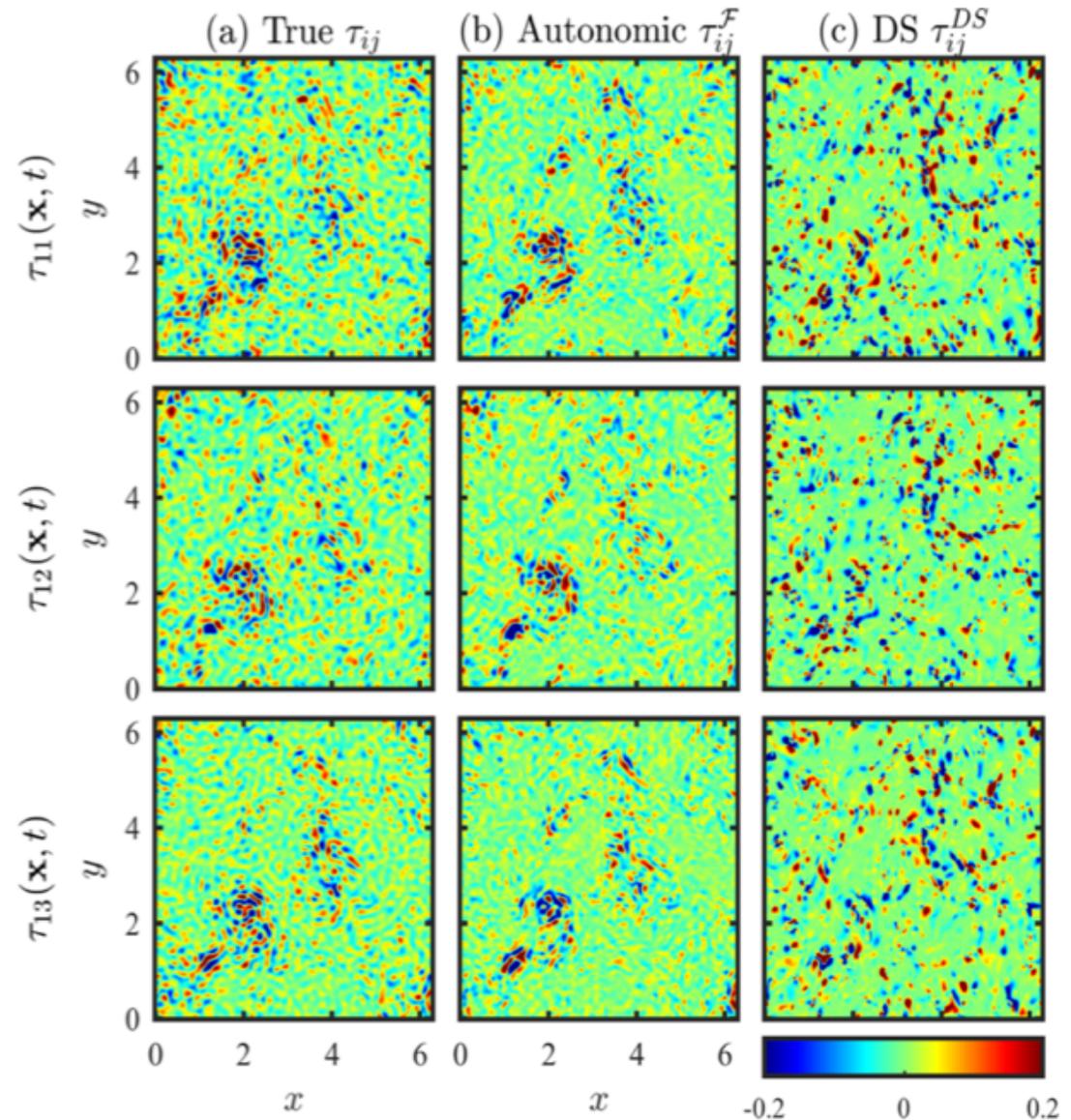


FIG. 2. Coarse-grained turbulent stress fields $\tau_{11}(\mathbf{x}, t)$ (top row), $\tau_{12}(\mathbf{x}, t)$ (middle row), and $\tau_{13}(\mathbf{x}, t)$ (bottom row), showing results for (left column) the true stress $\tau_{ij}(\mathbf{x}, t)$, (middle column) the autonomic closure $\tau_{ij}^{\mathcal{F}}(\mathbf{x}, t)$, and (right column) the dynamic Smagorinsky model $\tau_{ij}^{DS}(\mathbf{x}, t)$.

King+2016

Machine learning and modeling

- Use data to guess the functional form for the role of small scales in order to write sub-grid scale models

(f) pdf(τ_{13})

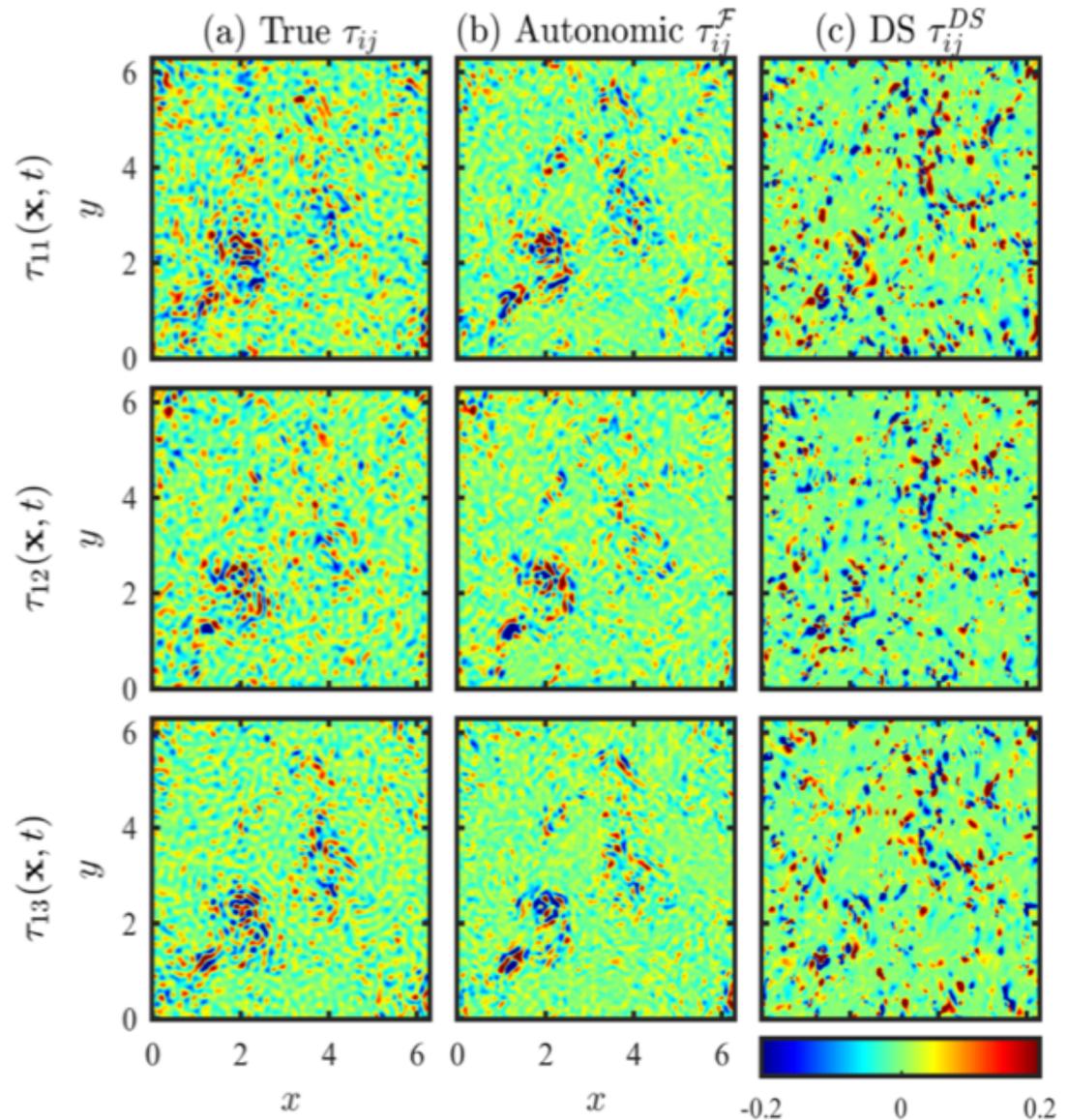
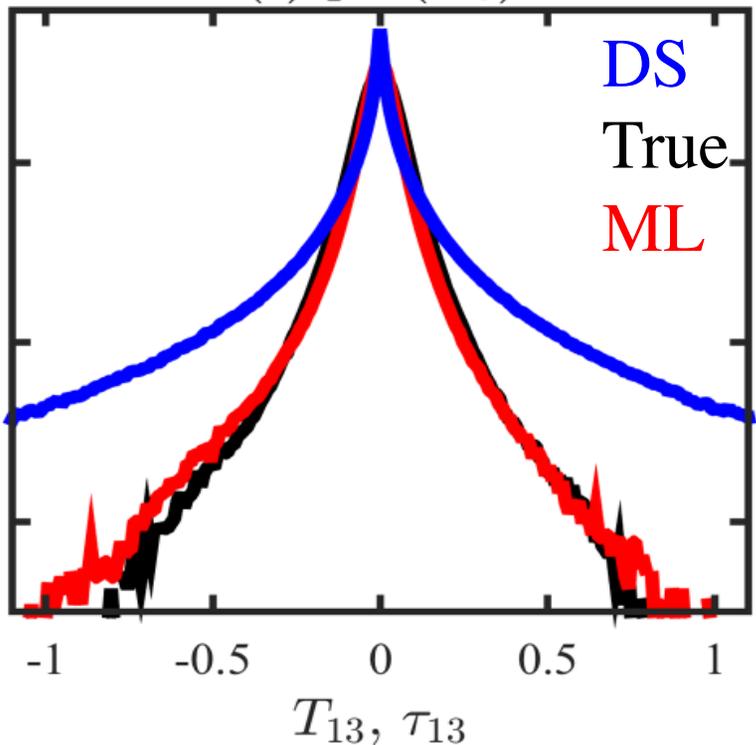
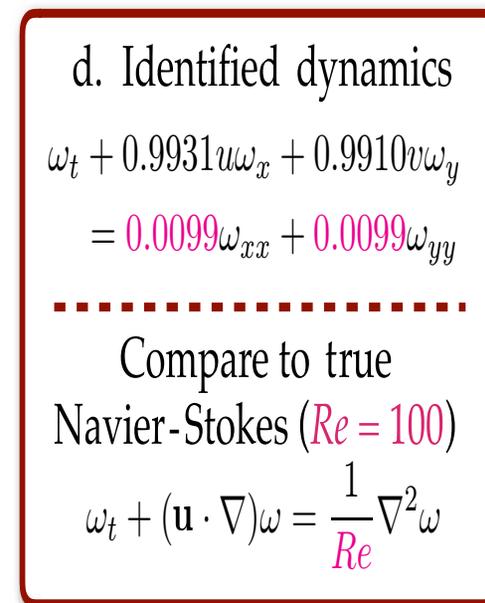
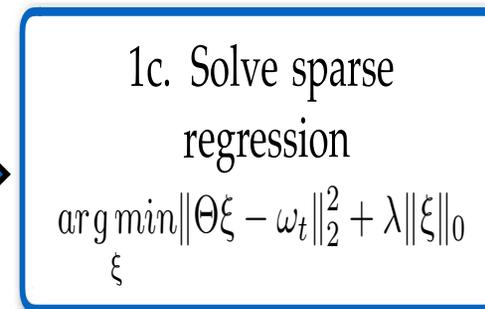
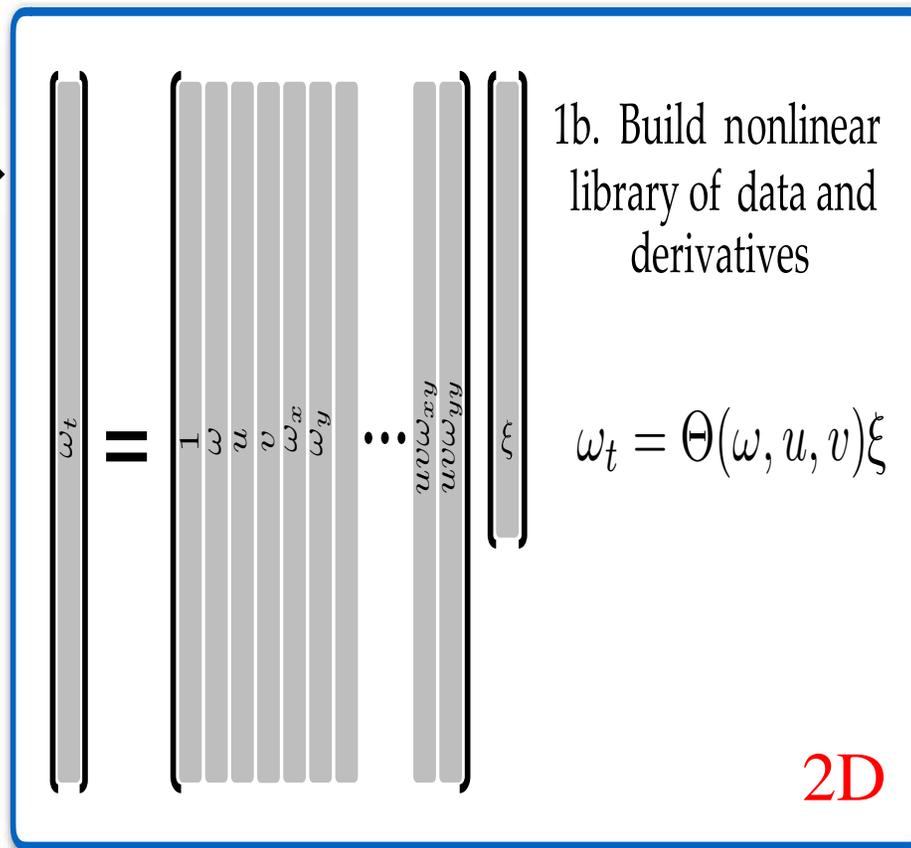
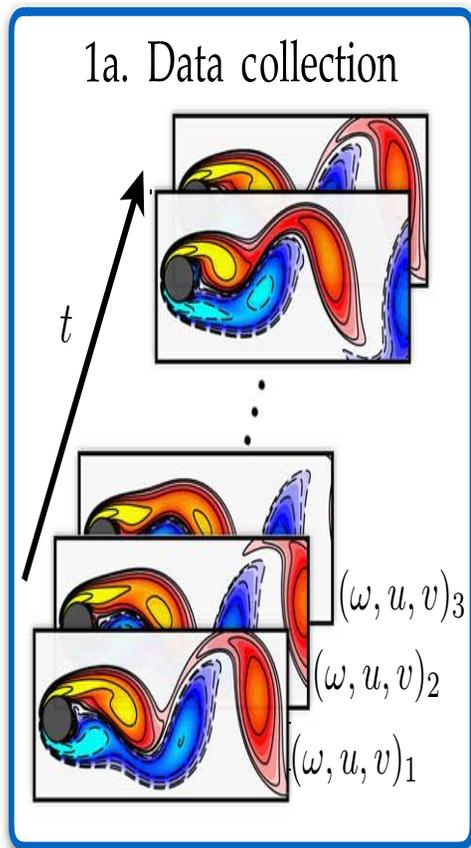


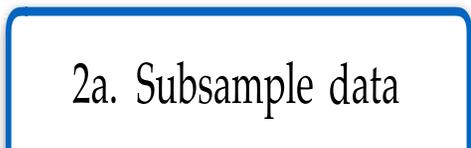
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Full data



data

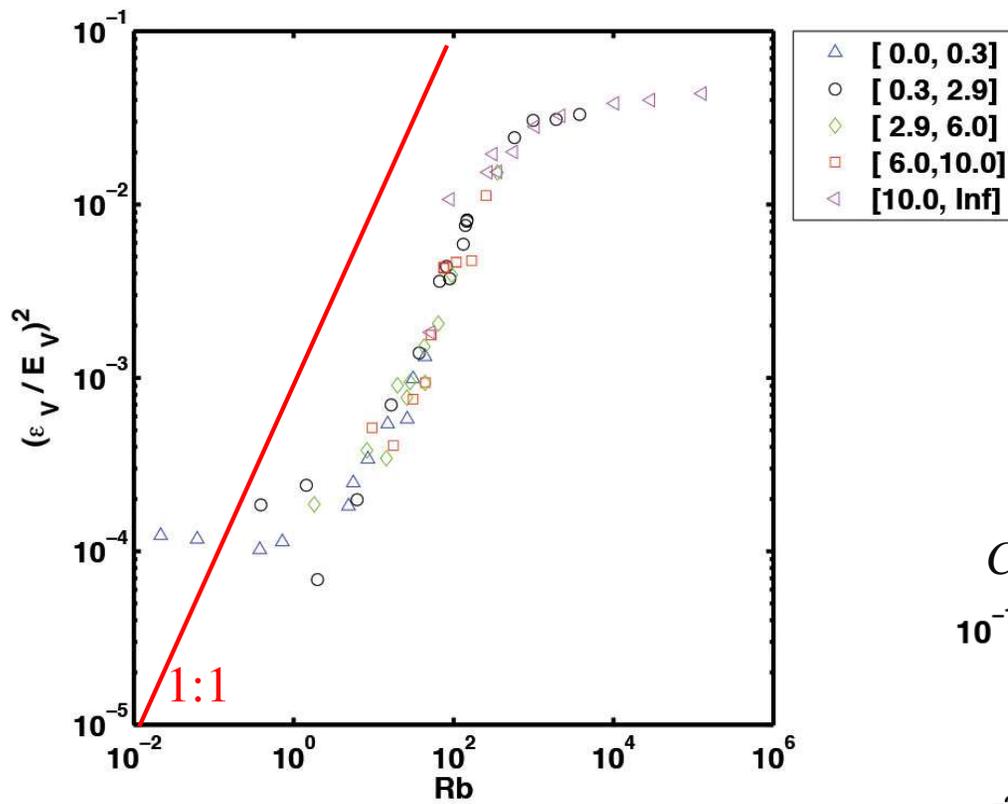


Thank you for your attention

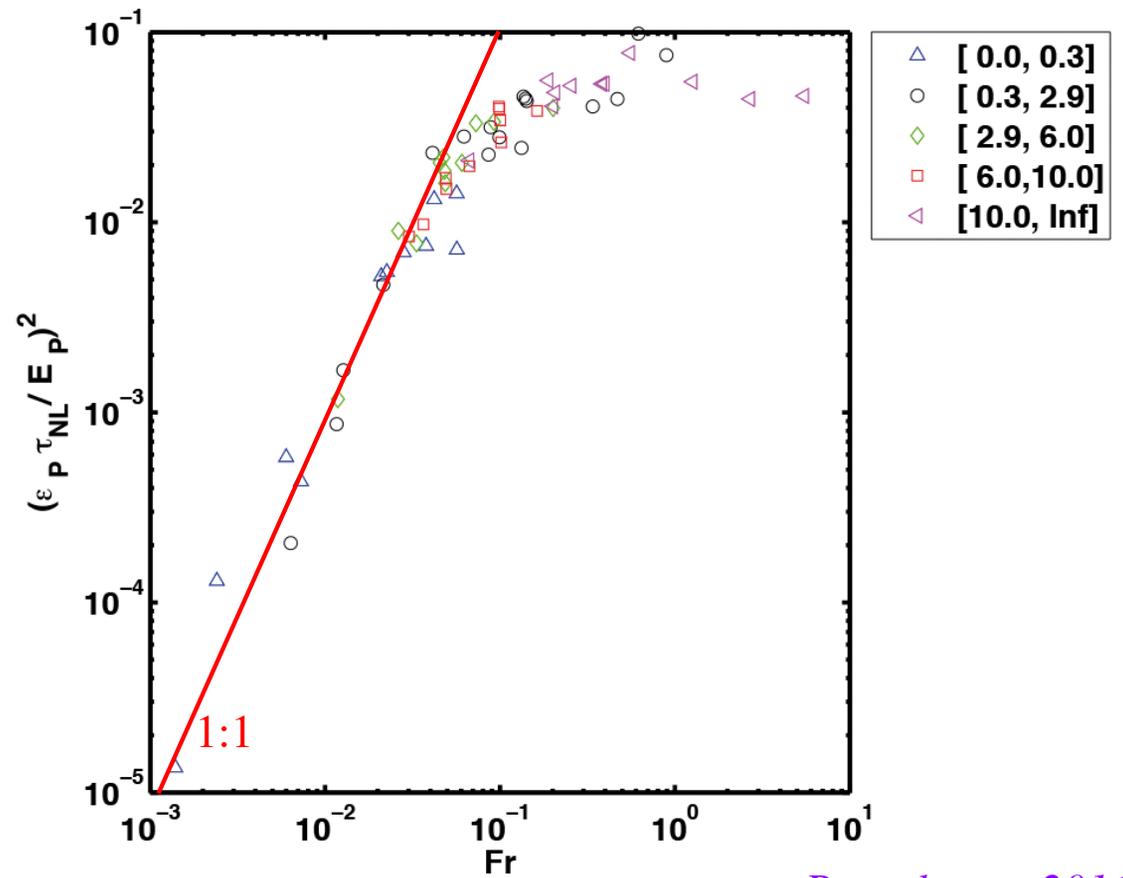
Some references

pouquet@ucar.edu

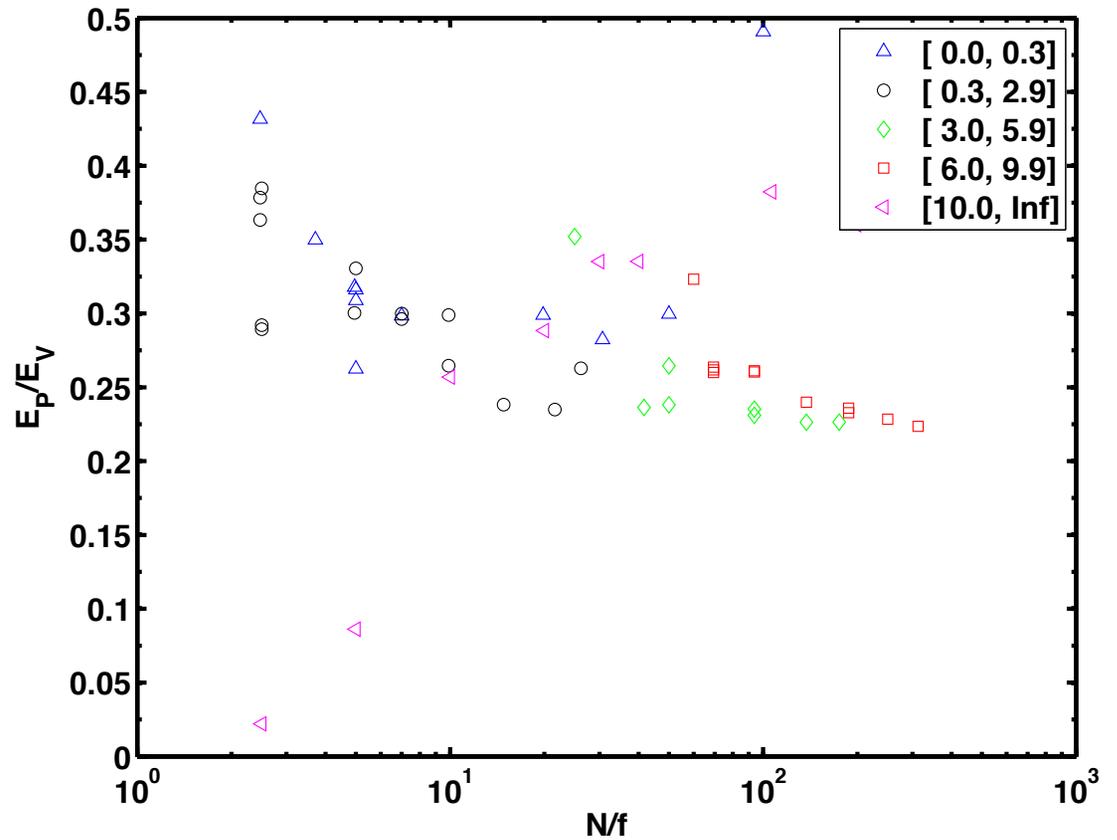
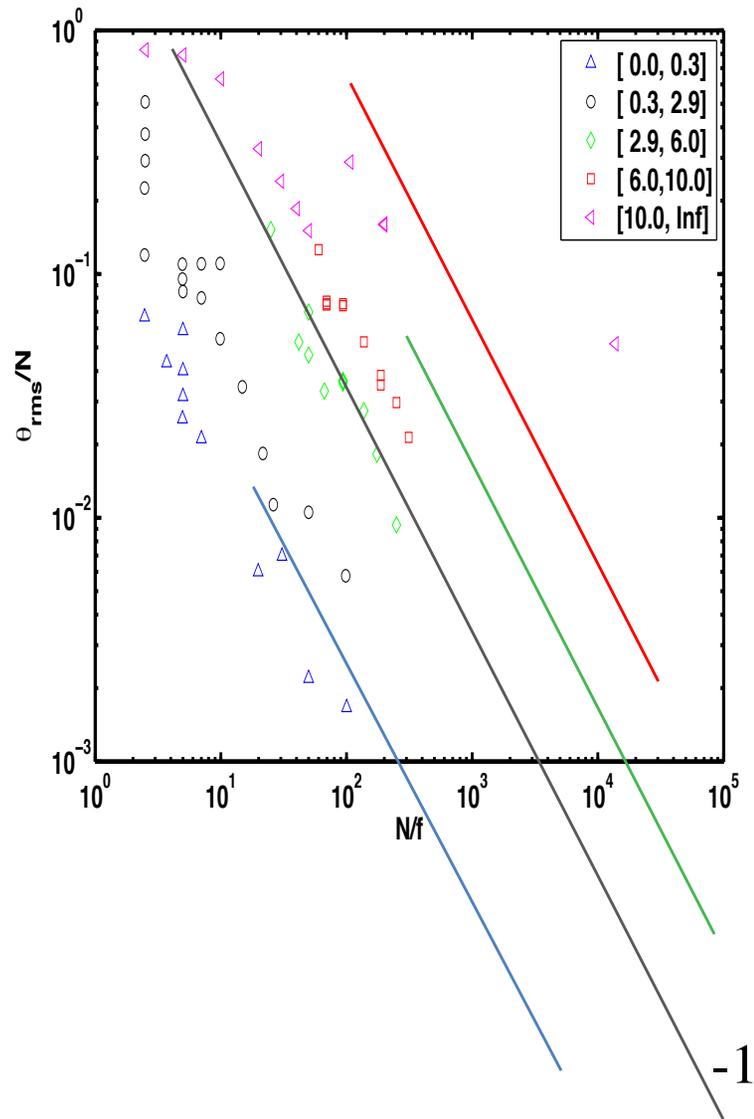
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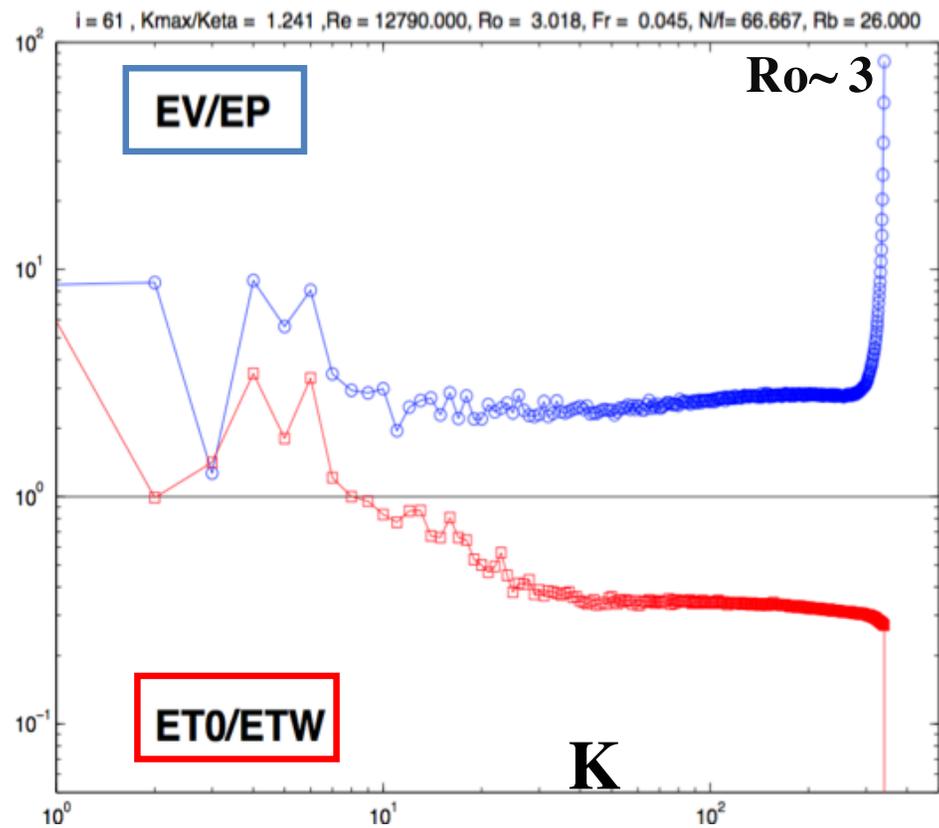
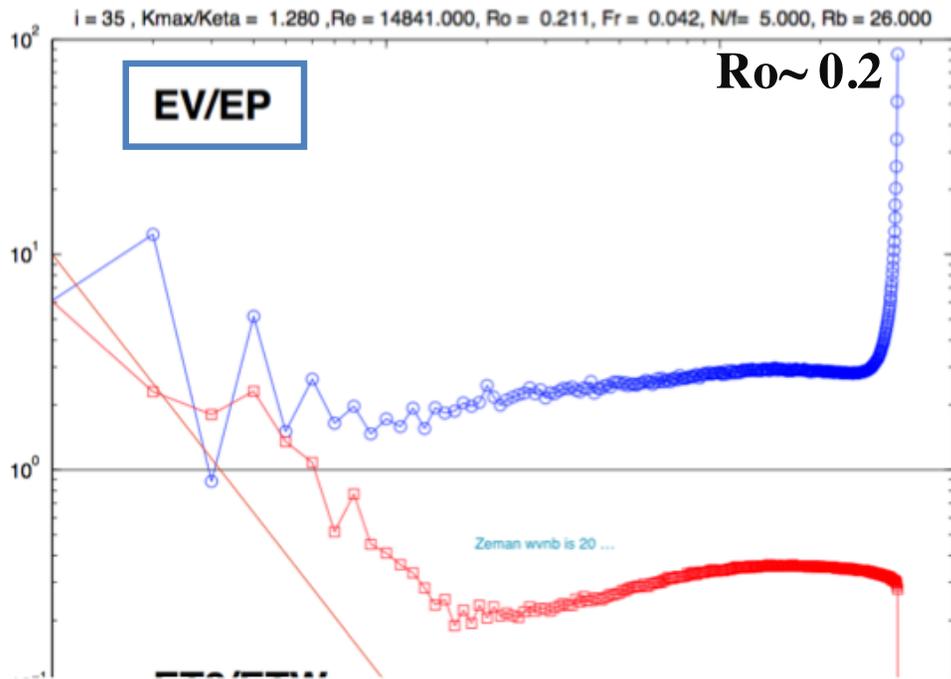


Color binning in Ro: 0 → 0.3 → 2.9 → 6.0 → 10 →



Variations with N/f of ←Ellison scale and $E_P/E_V \downarrow$





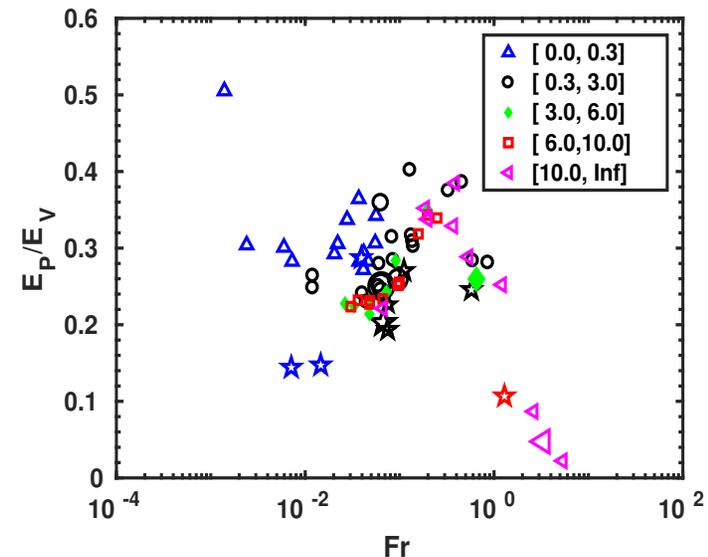
Ratios of energy spectra at peak:

Kinetic to potential

and

zero to wave mode

← $Fr=0.04, R_B \sim 26$



Color binning in Ro: 0 → 0.3 → 2.9 → 6.0 → 10 →

Ratio of potential to total energy dissipation

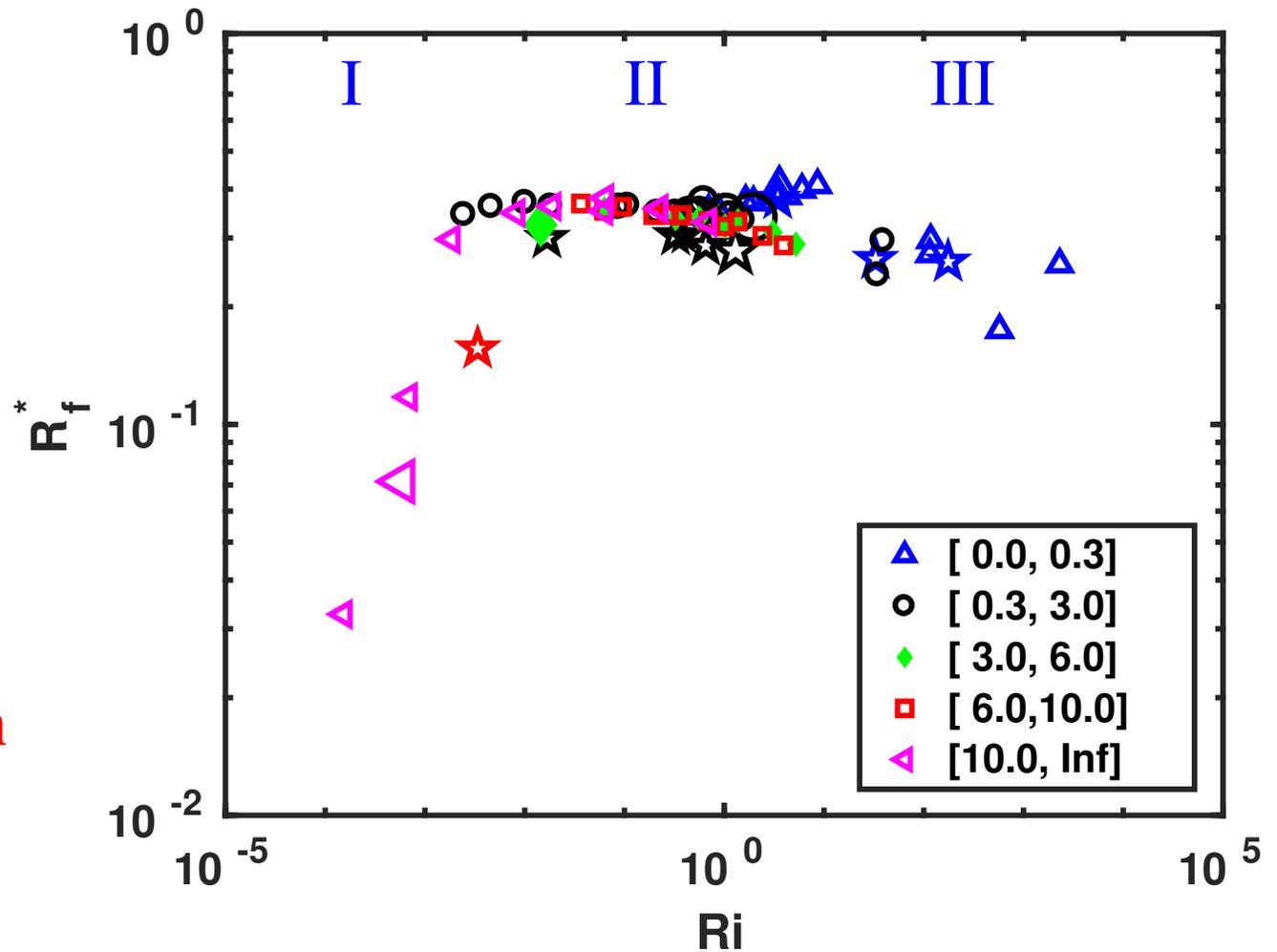
$$\frac{\varepsilon_p}{[\varepsilon_p + \varepsilon_v]}$$

versus

Richardson Number

$$Ri \equiv [N / \langle \partial_z u_{\perp} \rangle]^2$$

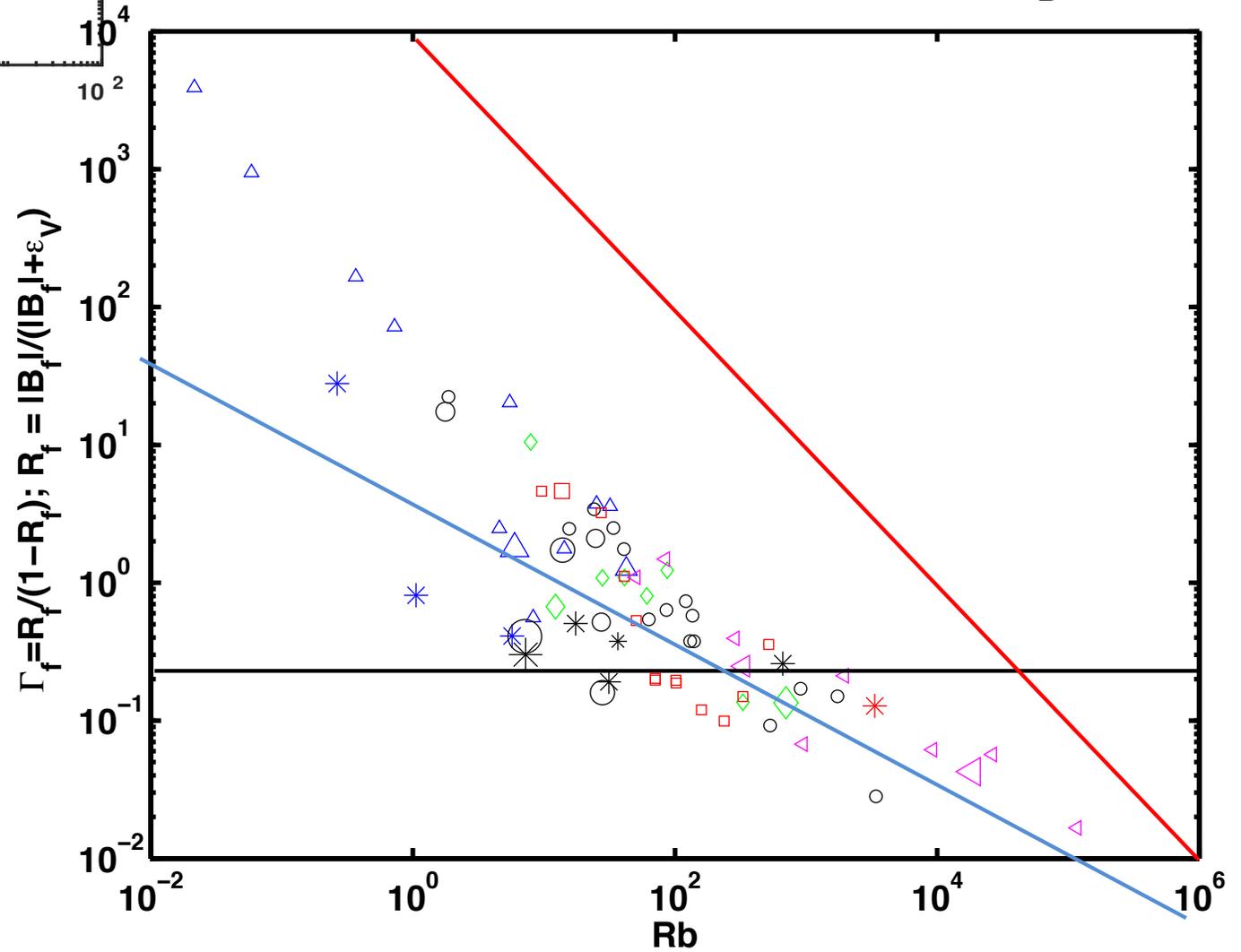
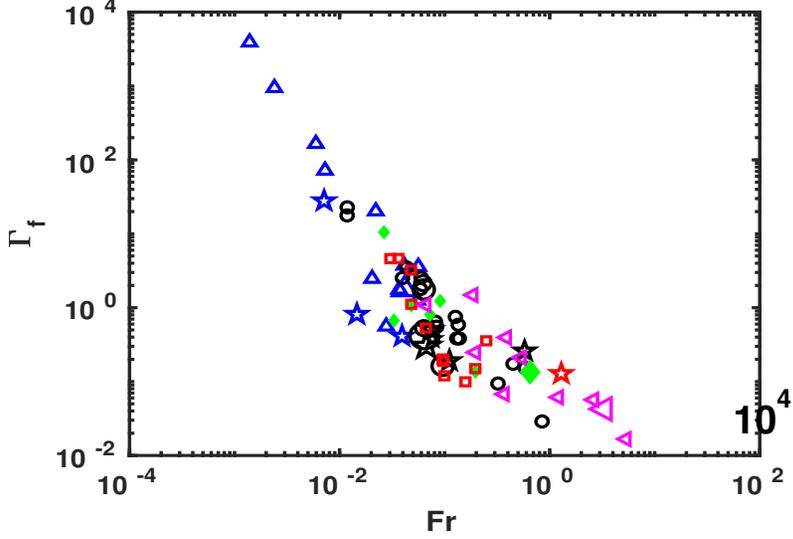
Rather constant in regime II,
but slight effect of rotation (blue triangles)



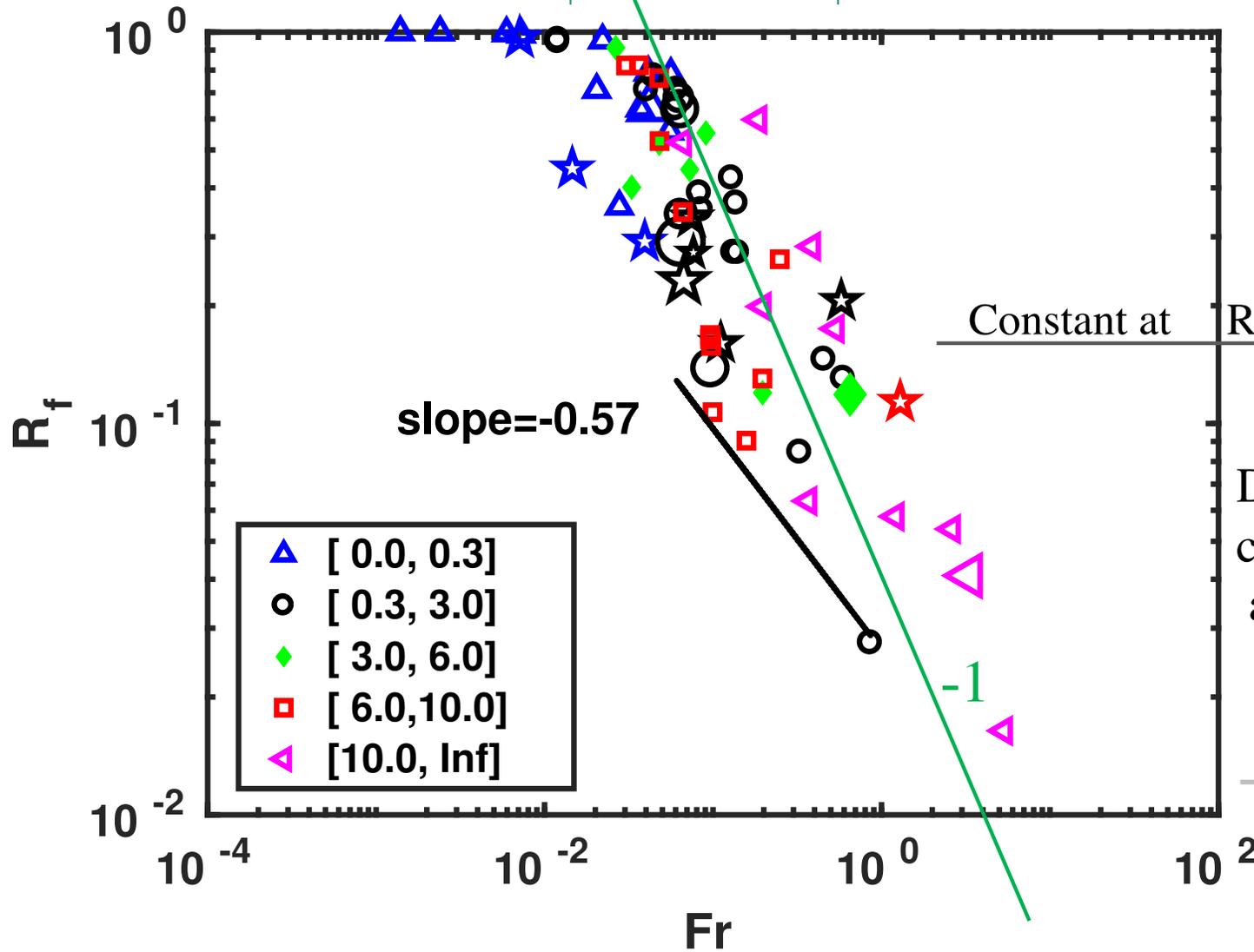
Color binning in Ro : 0 \rightarrow 0.3 \rightarrow 2.9 \rightarrow 6.0 \rightarrow 10 \rightarrow

Color binning in Ro : 0 \rightarrow 0.3 \rightarrow 2.9 \rightarrow 6.0 \rightarrow 10 \rightarrow

Mixing
efficiency
 $\Gamma_f = B_f / \epsilon_V$
in terms of
 R_B



Flux Richardson
 number: $R_f = B_f / [B_f + \epsilon_v]$

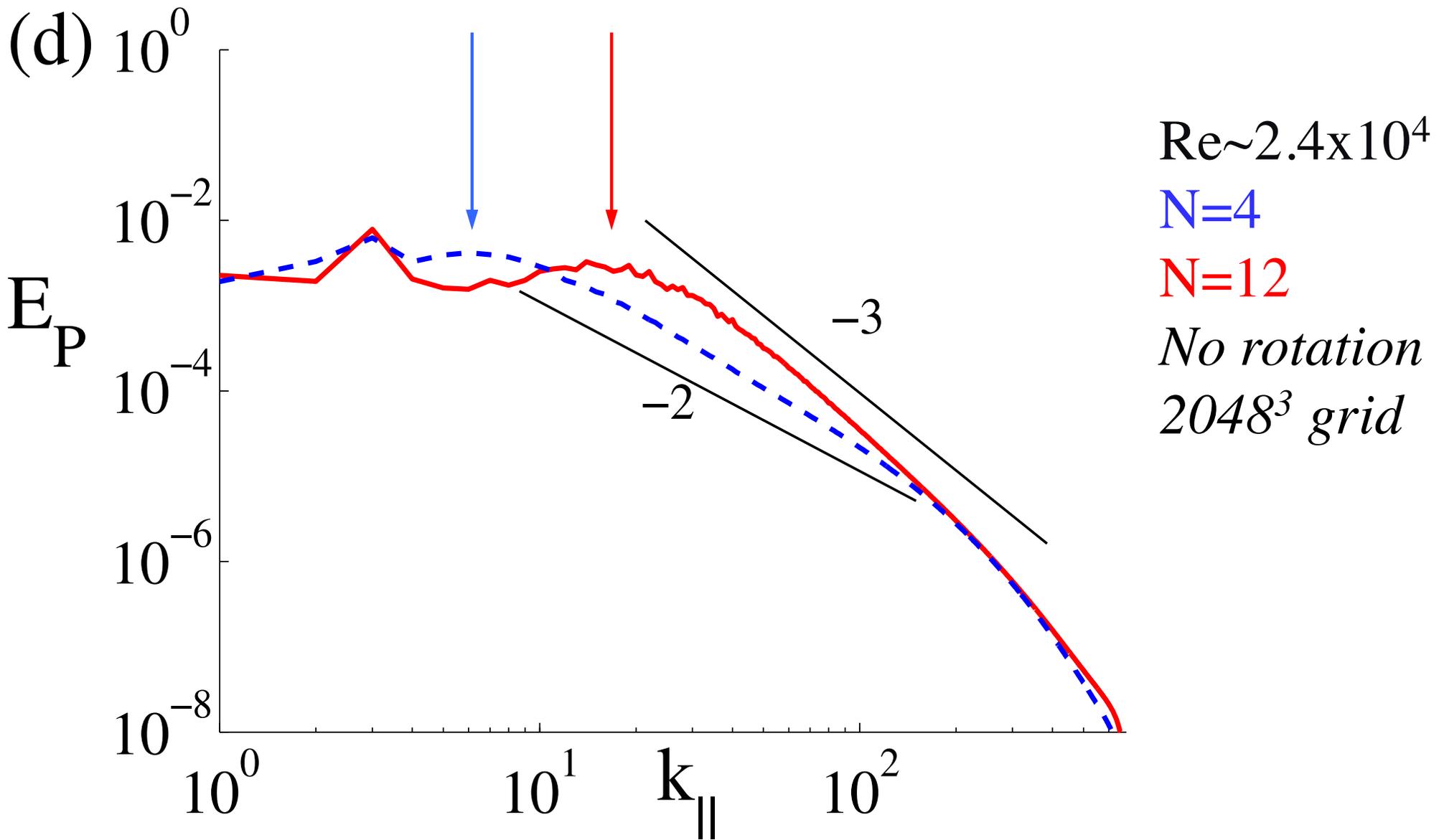


Constant at $R_f \sim 0.17?$ Perhaps for QG??

Dependency on initial conditions for θ (small), and on N/f (large)?

$v_{shw} \sim [N/f] d_{\perp} \theta$
 $\rightarrow \theta \sim U[f/N] < 1$

Color binning in Ro : $0 \rightarrow 0.3 \rightarrow 2.9 \rightarrow 6.0 \rightarrow 10 \rightarrow$

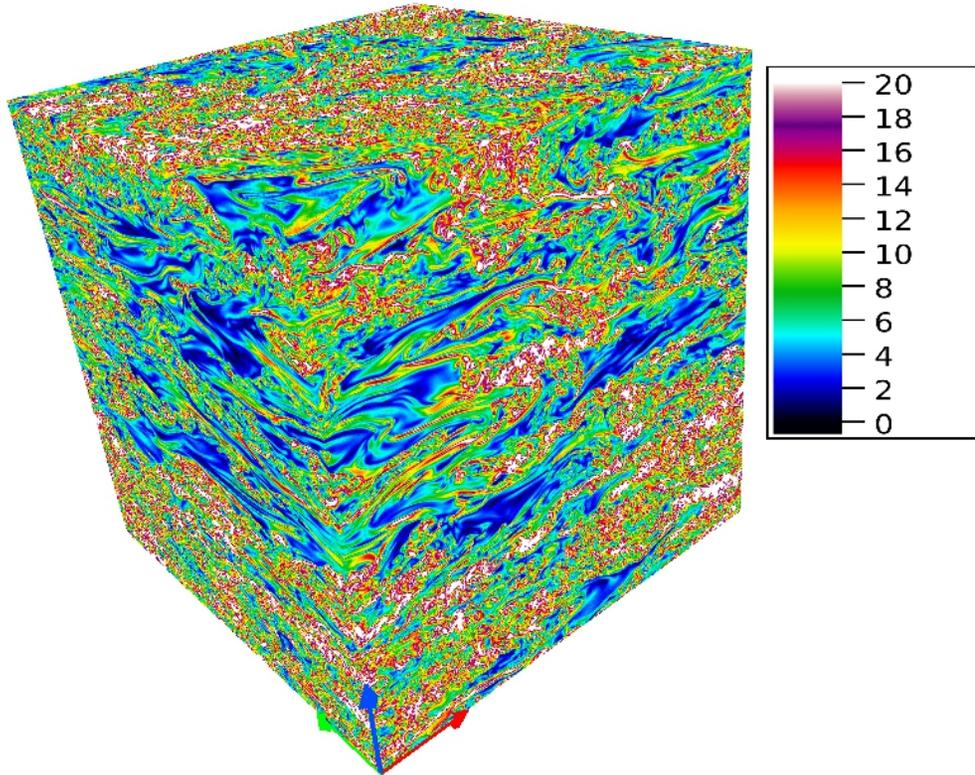


Potential energy = $f(k_z)$:
 Plateau until $k_B \sim N/U$,
 the buoyancy wvnb.

Incompressible Boussinesq equations + rotation

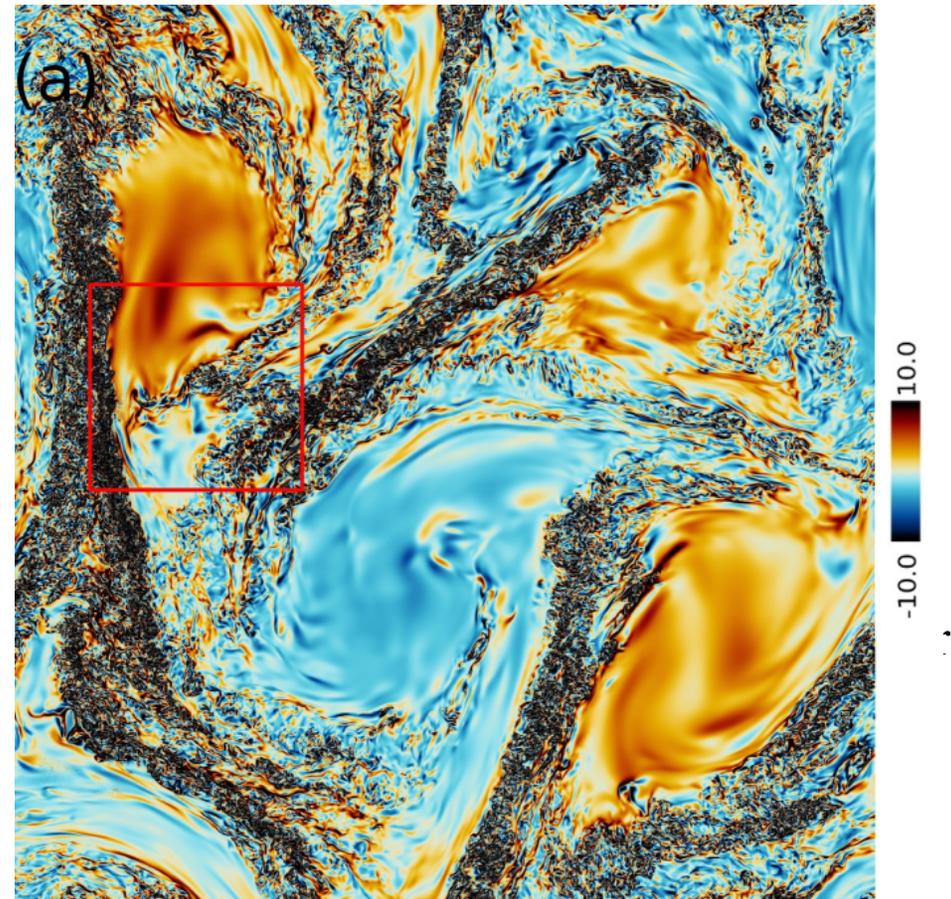
Vorticity, 3D rendering “atmosphere”

$Ro = 9.2$, $Fr = 0.067$, $Re \simeq 12000$, $R_B \sim 53$,
 $N/f = 137$, 1024^3 grid, ... (Rosenberg+ 2016)



Vorticity, 2D cut, “ocean”, $Ro \sim 0.12$,

$Fr \sim 0.024$, $Re \sim 54000$, $R_B \sim 32$, $N/f = 5$,
 4096^3 grid, decaying, resolved & strongly
intermittent (Rosenberg+ 2015)

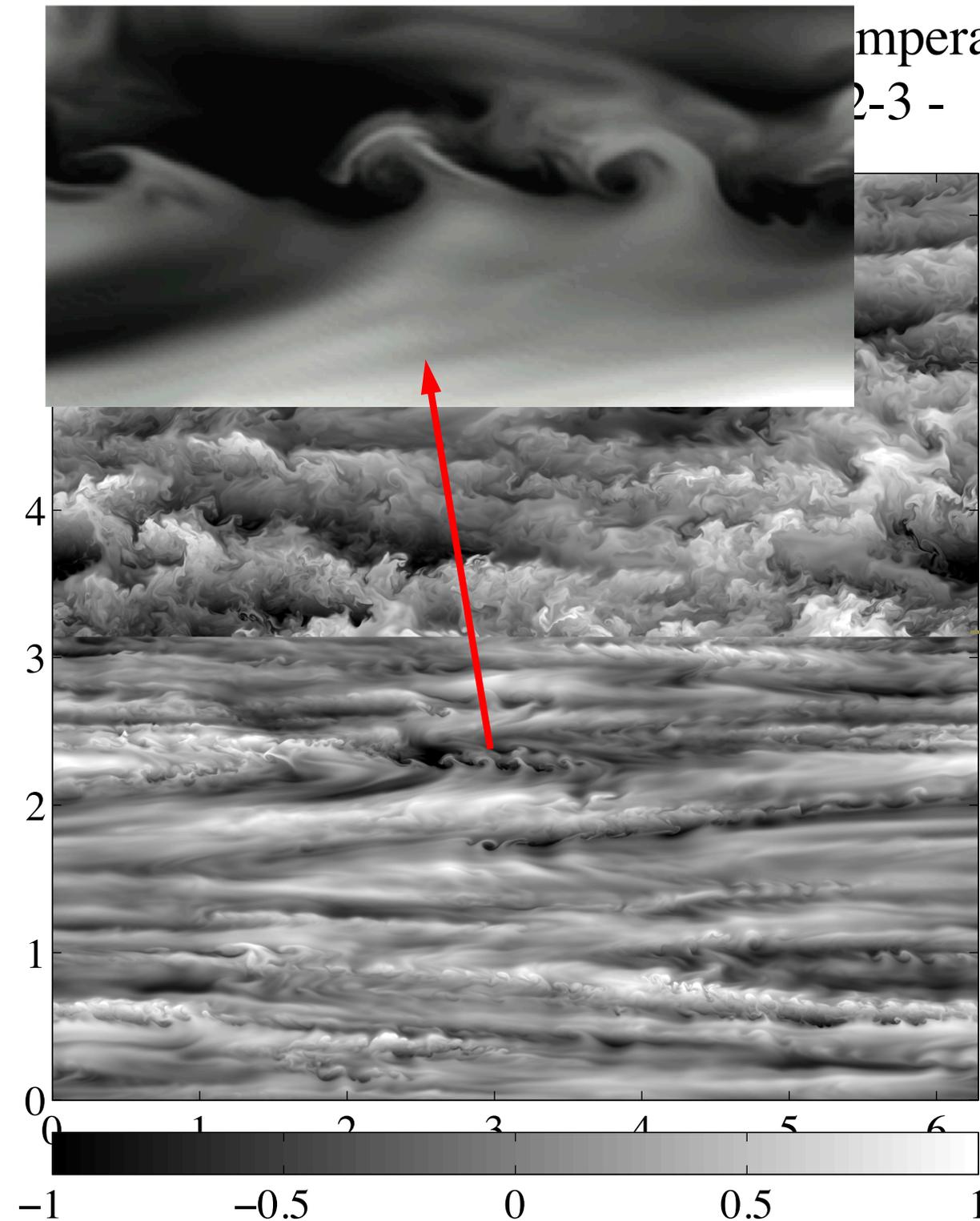


Temperature fluctuations, xz slice
2-3 - **Re ~ 24000, 2048³ grids, f=0**

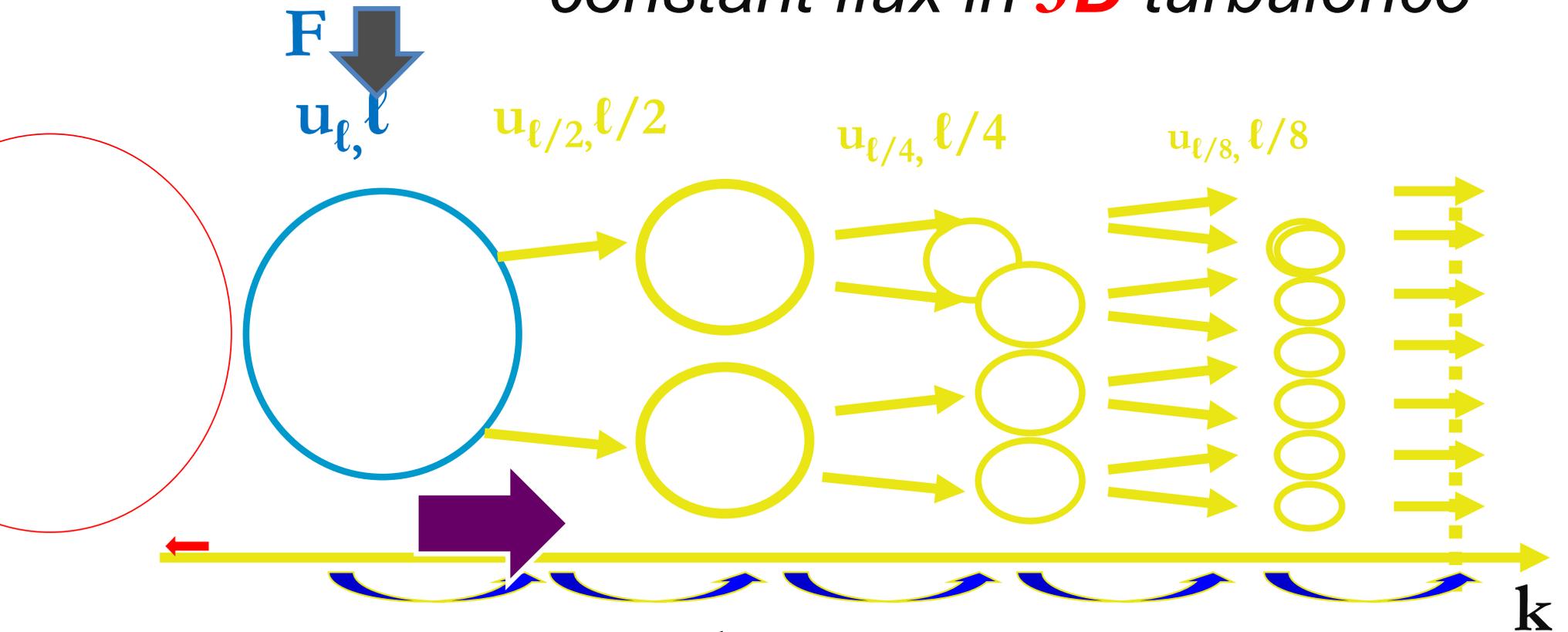
$N=4, \quad Fr \sim 0.11$
 $R_B = ReFr^2 \sim 300$

$N=12, Fr \sim 0.03$
 $R_B \sim 22$

Rorai et al., 2014



Classical energy cascade with constant flux in 3D turbulence



$$\sin a \cos b = \frac{1}{2} (\sin(a + b) + \sin(a - b))$$

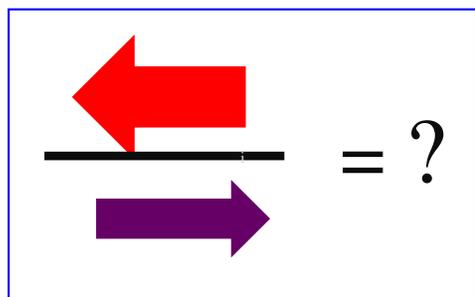
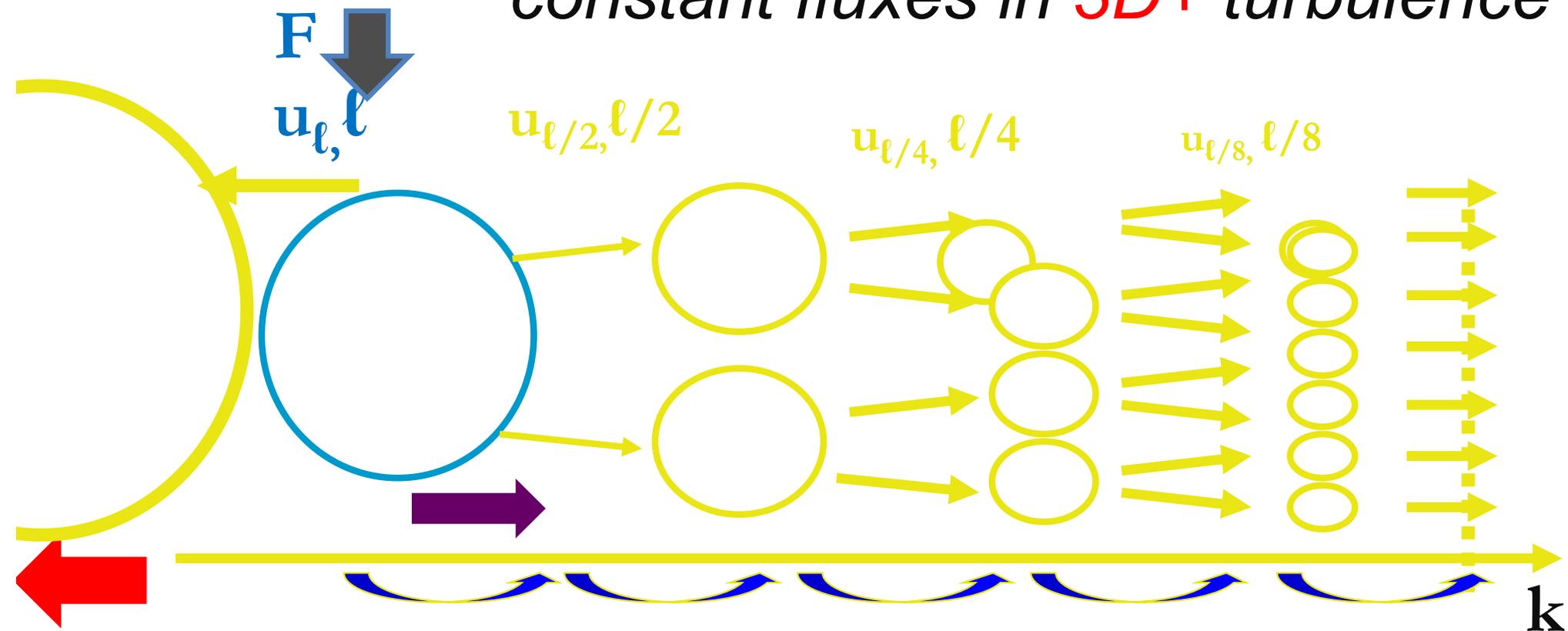
Advection

$u \cdot \text{grad}(u) \rightarrow$ Fourier : E goes to small & large scales
 \rightarrow Convolution

Eddy viscosity & eddy noise

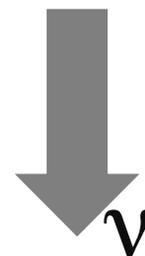


Dual energy cascades with dual constant fluxes in 3D+ turbulence



Rotation
with or without
stratification

And/or magnetic field



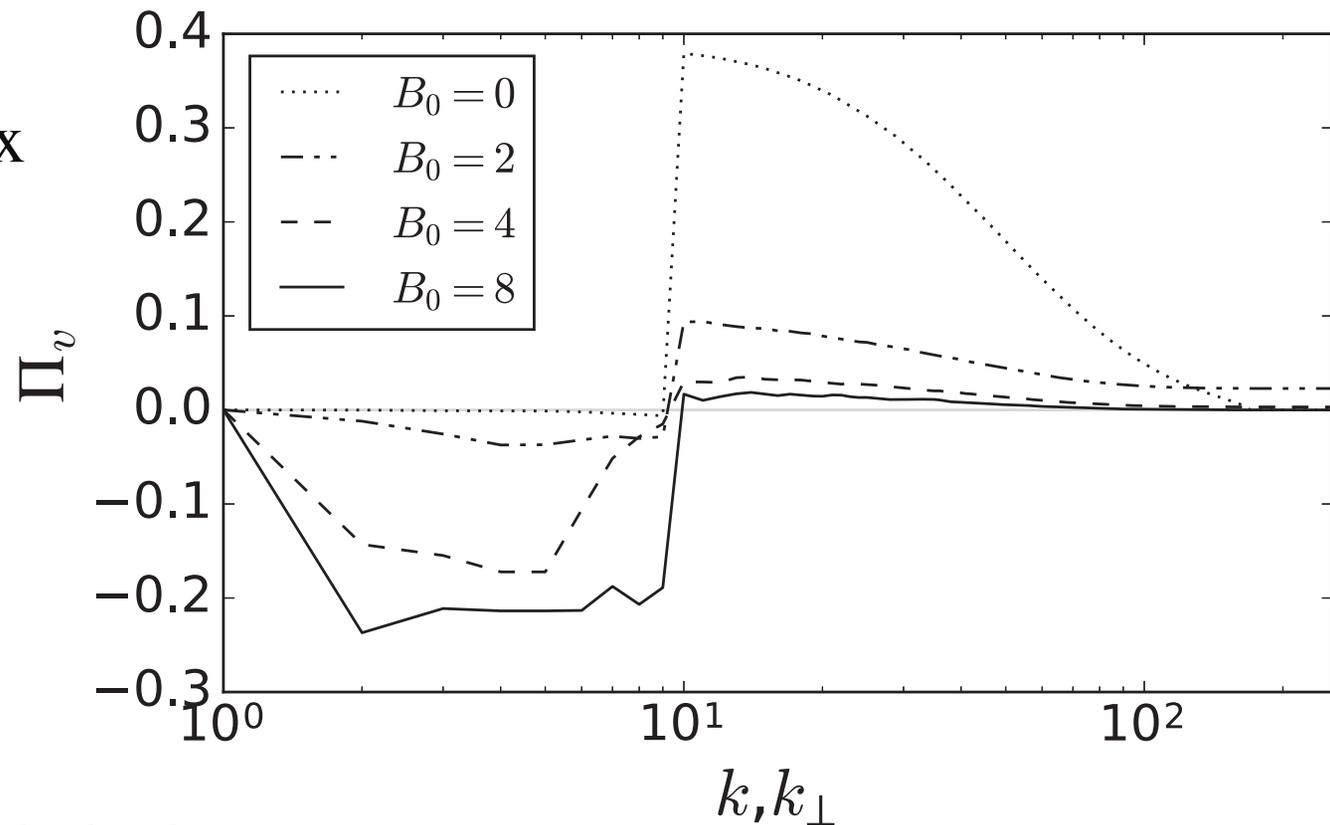
3D-MHD + imposed external field B_0

+ F_V but with $F_M = 0$

(Alexakis 2011; Sujovolsky+ 2016)

TRIDIMENSIONAL TO BIDIMENSIONAL TRANSITION IN ...

Kinetic energy flux



Also: kinetic Alfvén and whistler waves in the “Solar Wind” (Che+ 2014)

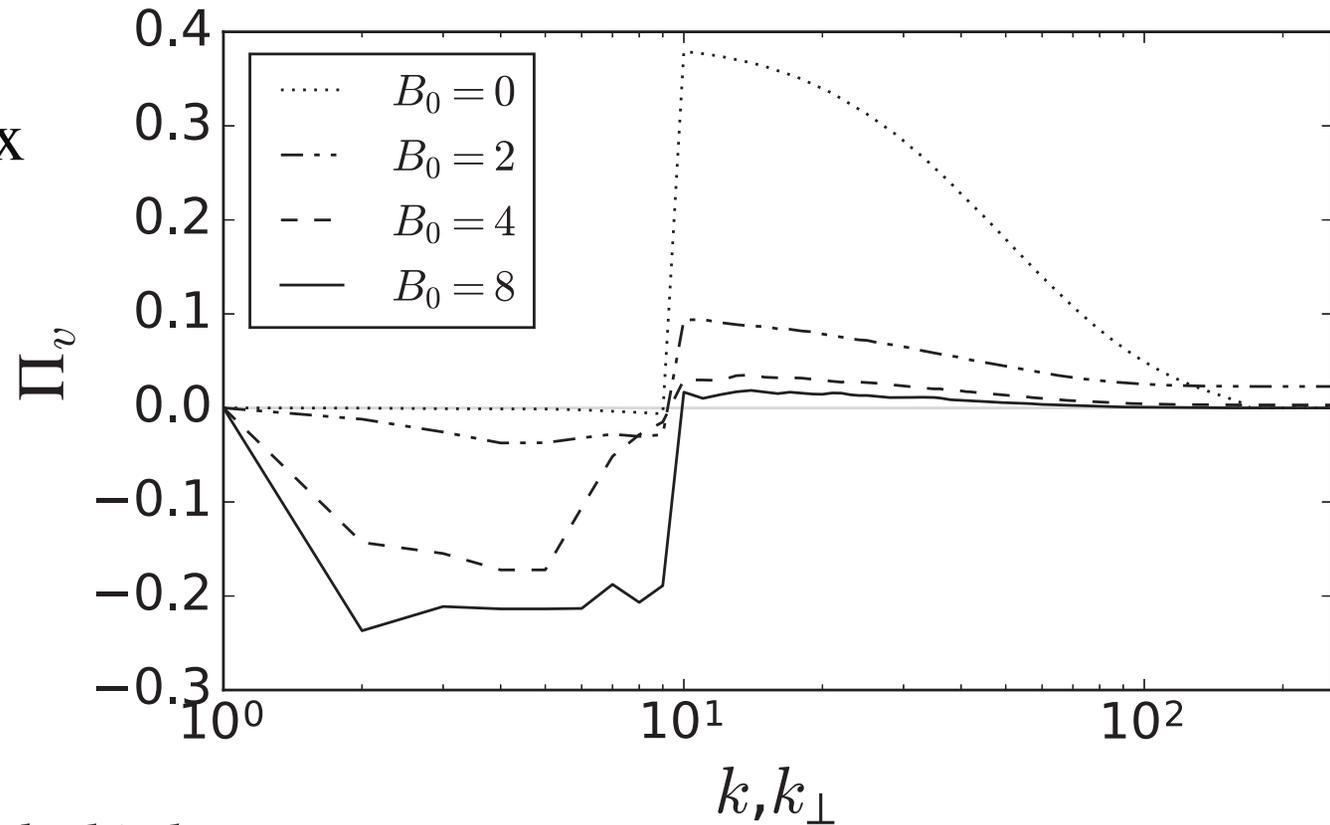
3D-MHD + imposed external field B_0

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TRIDIMENSIONAL TO BIDIMENSIONAL TRANSITION IN ...

Kinetic energy flux



Also: kinetic Alfvén and whistler waves in the “Solar Wind” (Che+ 2014)

Possibility of lab. experiment?

Strictly two-dimensional forced MHD

Control parameter: $\mu = F_M/F_V$

FLUXES

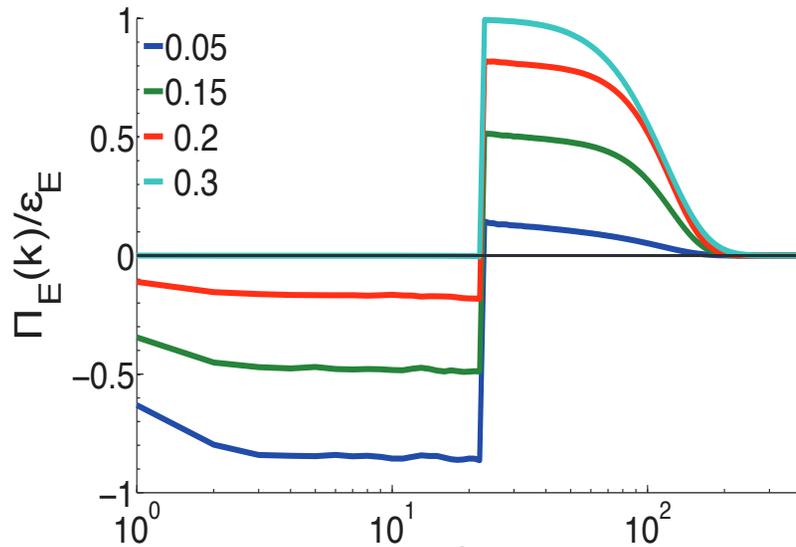
$$\epsilon_E^- \propto (\mu_c - \mu)^{\gamma_E}$$

$$\epsilon_A^- \propto (\mu - \mu_c)^{\gamma_A}$$

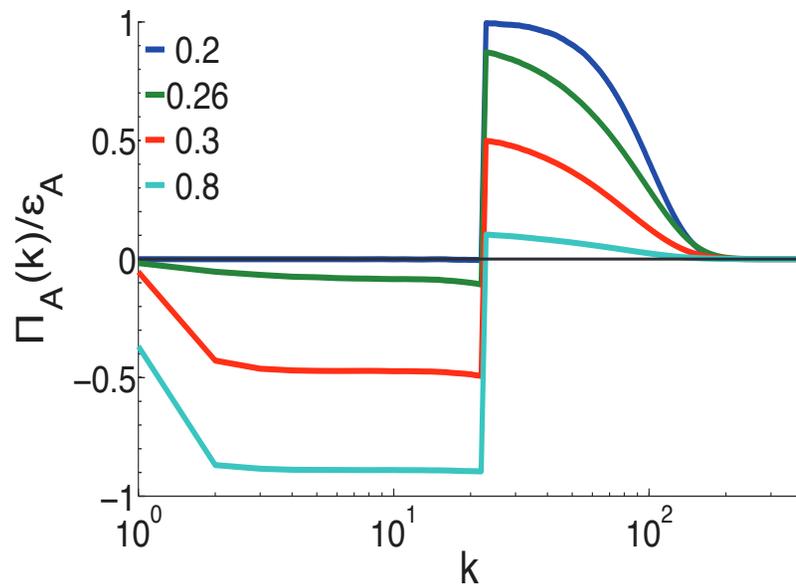
$$\gamma_E \sim 0.82$$

$$\gamma_A \sim 0.24$$

$$\mu_c \sim 0.22_{(E)} \text{ or } 0.25_{(A)}$$



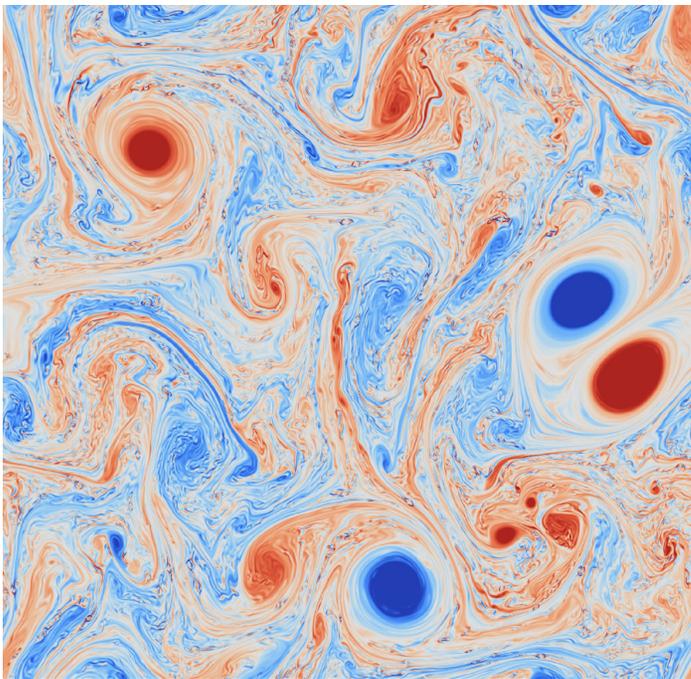
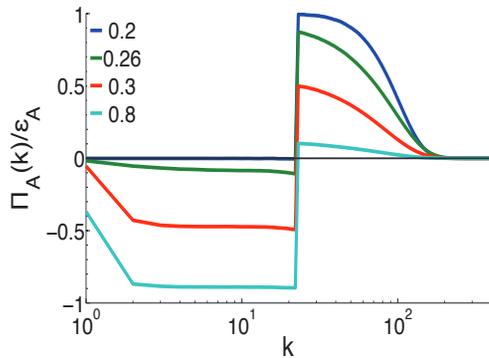
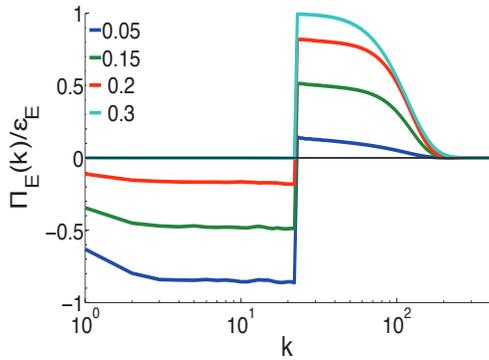
Energy \uparrow and magnetic potential \downarrow



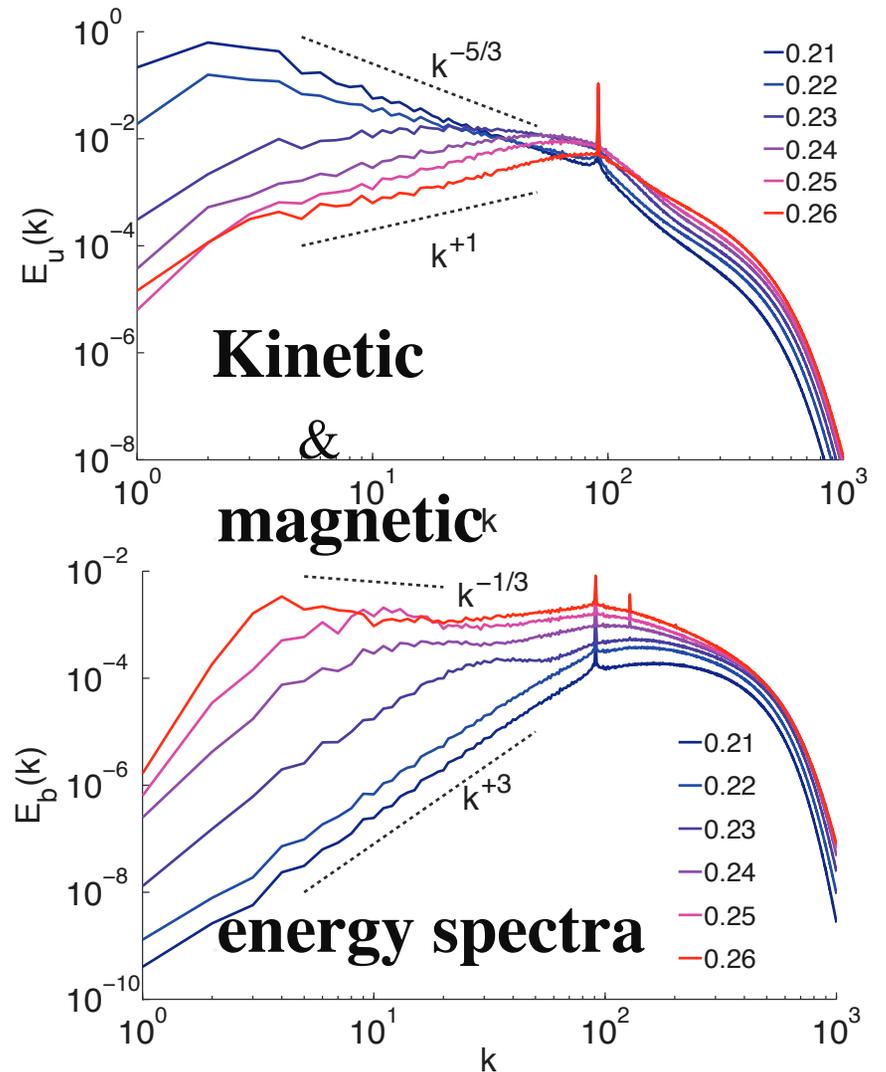
Strictly 2D forced MHD, Control parameter $\mu = F_M/F_V$

$$E_A(k) = k^{-2} E_M(k)$$

SESHASAYANAN, BENAVIDES, AND ALEXAKIS (2014, 2016)



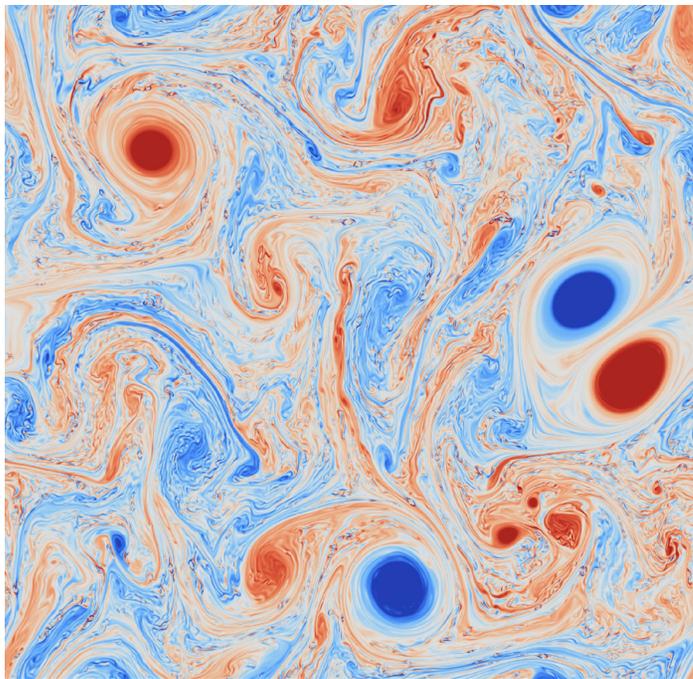
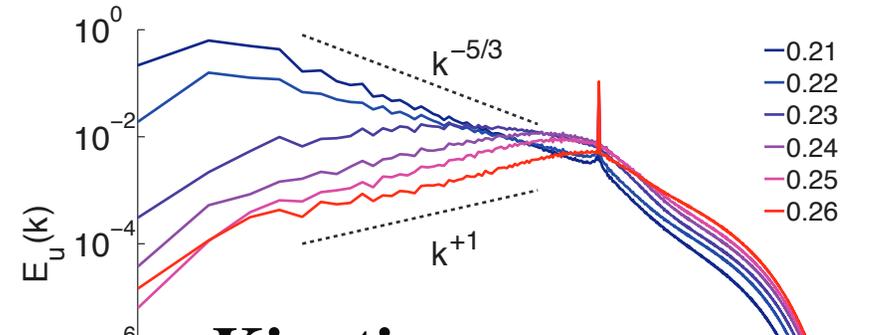
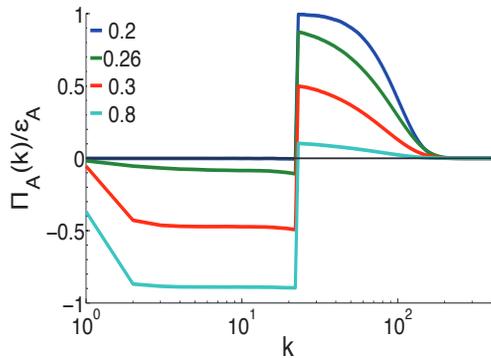
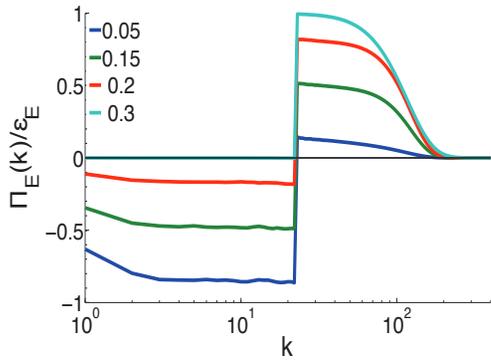
Vorticity, forcing scale 1/20: _



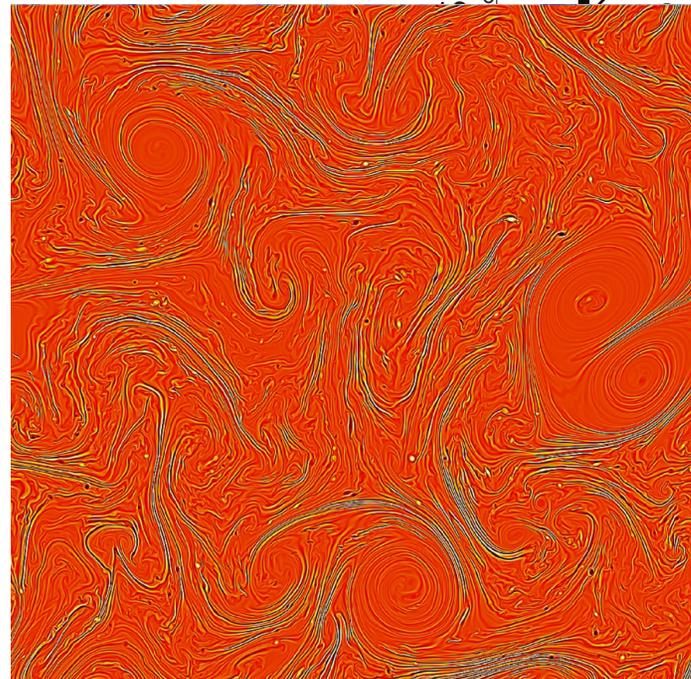
Strictly 2D forced MHD, Control parameter $\mu = F_M/F_V$

$$E_A(k) = k^{-2} E_M(k)$$

SESHASAYANAN, BENAVIDES, AND ALEXAKIS (2014, 2016)



Vorticity, forcing scale 1/20: $_$



Current

