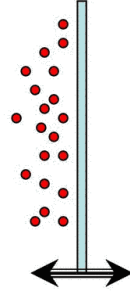


Outline

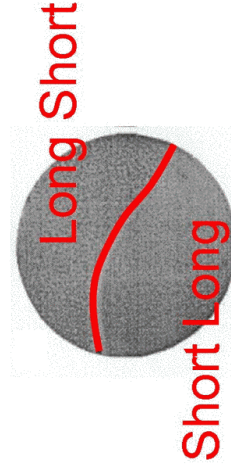
- Watanabe et al., JCE, 2001
- Echebarria and Karma, PRL, 2002;
Condmnat 2006 (50 pages)
- Echebarria and Karma, Chaos 2002
- Christini et al., PRL 2006
- Shiferaw et al., Biophys J, 2003
- Shiferaw et al., Phys. Rev. E, 2005
- Shiferaw and Karma, PNAS, 2006
- Sato et al., Circ. Res. 2006 (submitted)

Discordant Sand Alternans
Swinney et al., 1995



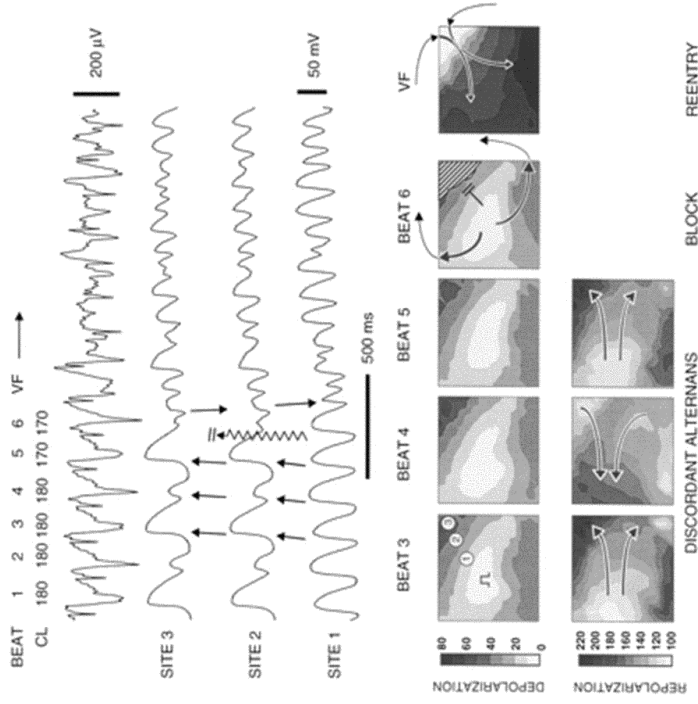
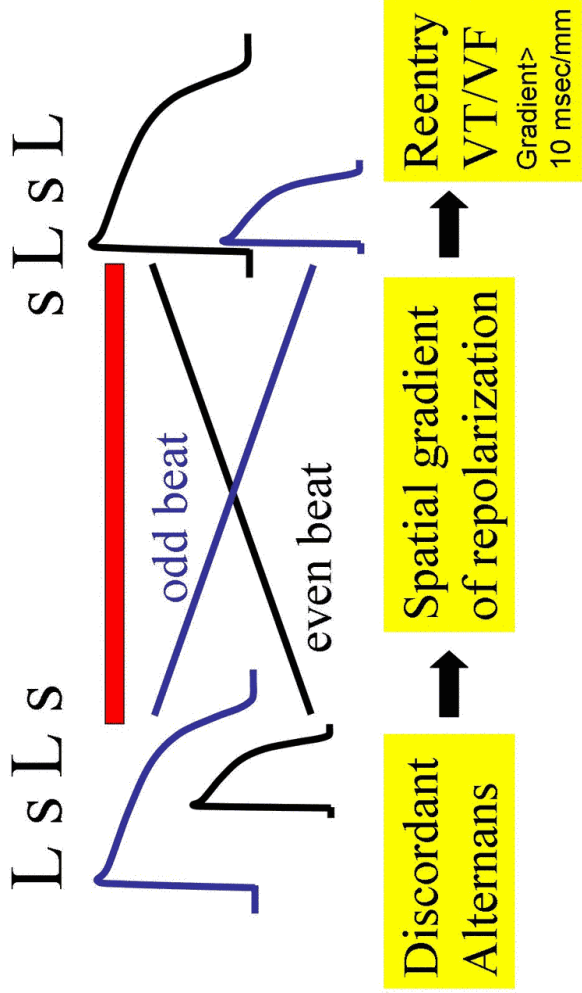
$$z = A \cos 2\pi f t$$

A single sand grain
can be driven into
alternans by a
vibrating platform



Many sand grains can
exhibit discordant alternans

Discordant Alternans

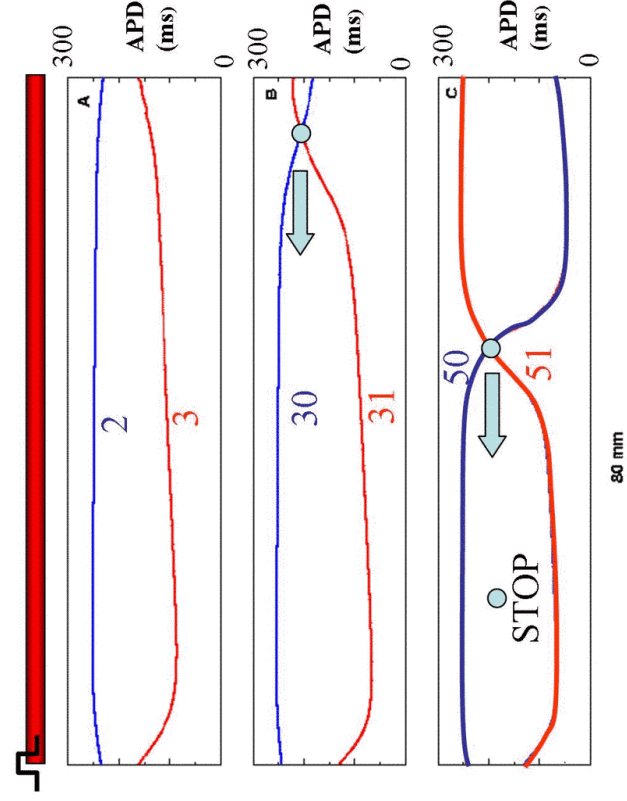


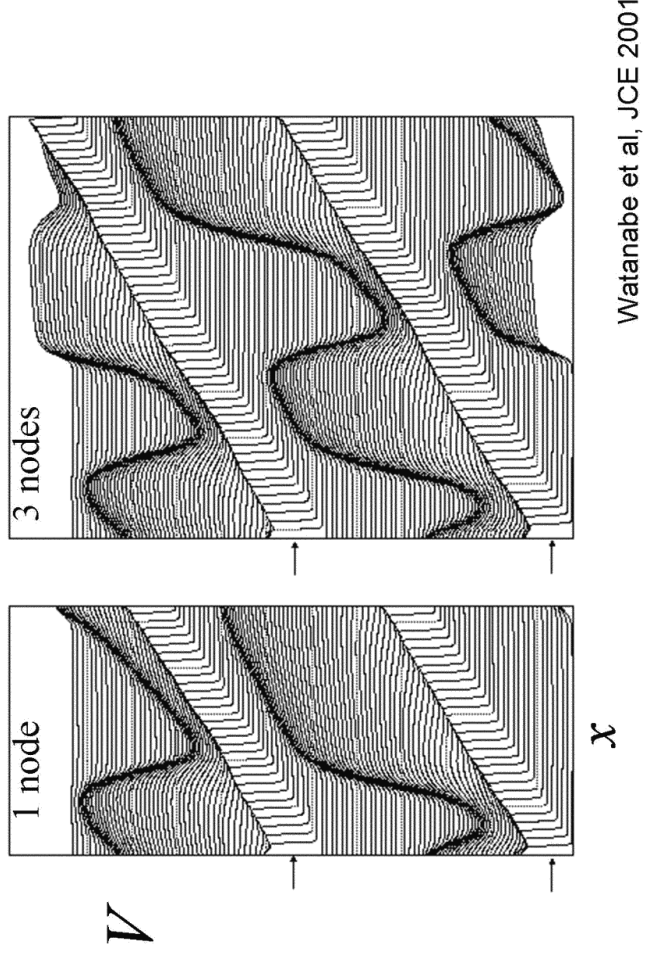
Discordant Alternans → VF

Pastore et al., 1999

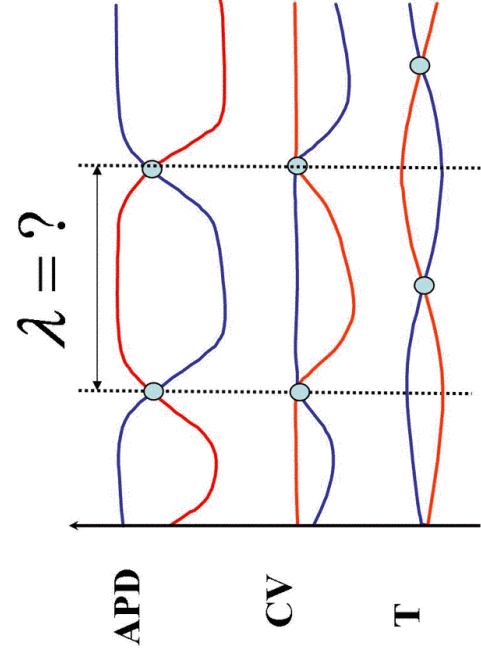
Are spatially discordant tissue
scale electrical alternans of static
or dynamic origin ?

Discordant alternans =
spatially out-of-phase period
doubling oscillations



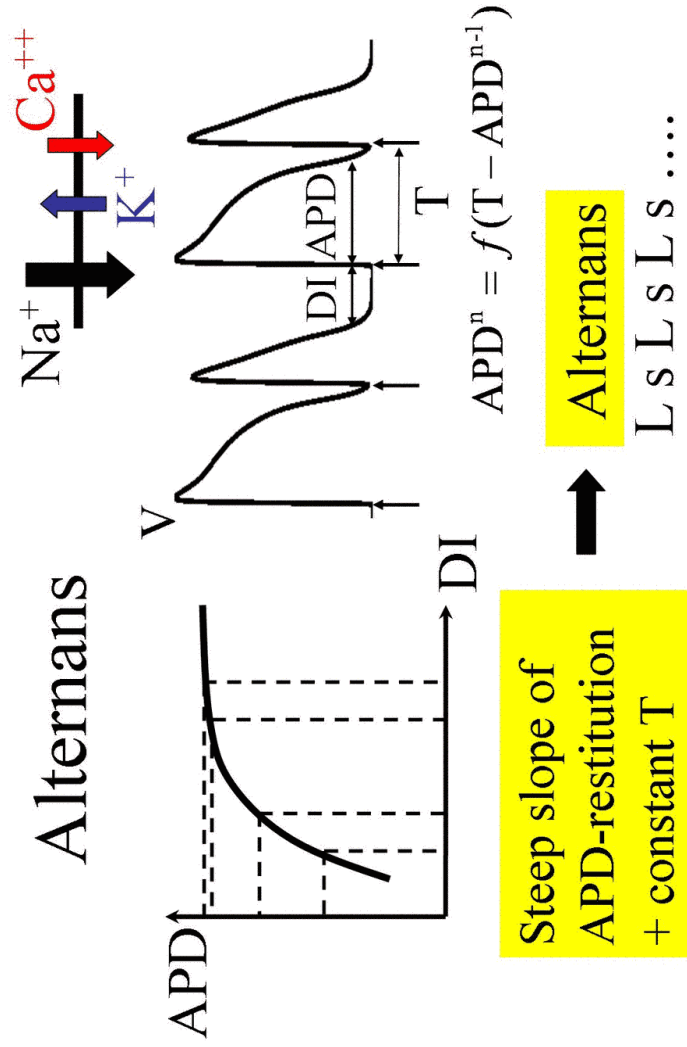


Surprise: wave patterns ?



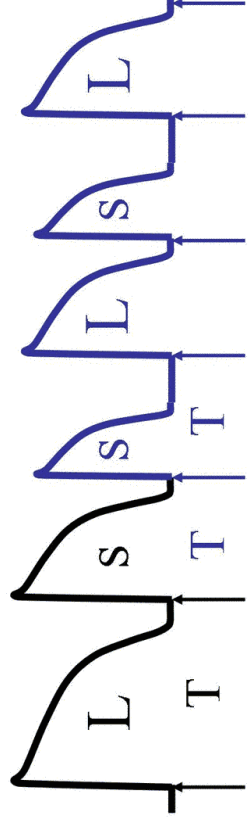
Both standing and travelling waves are possible.

How do these wave patterns form
and what determines their
wavelength ?



Dephasing

$LsLs \rightarrow sLsL$

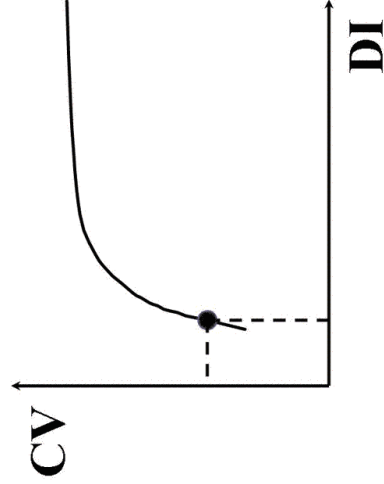
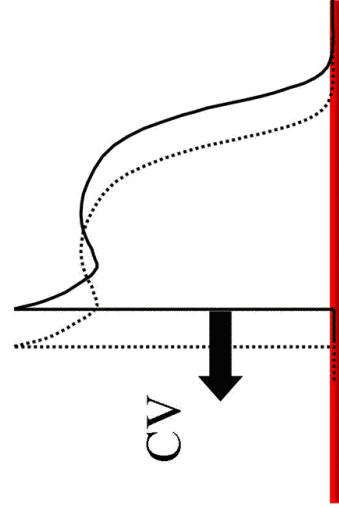


Change of T

Change of phase

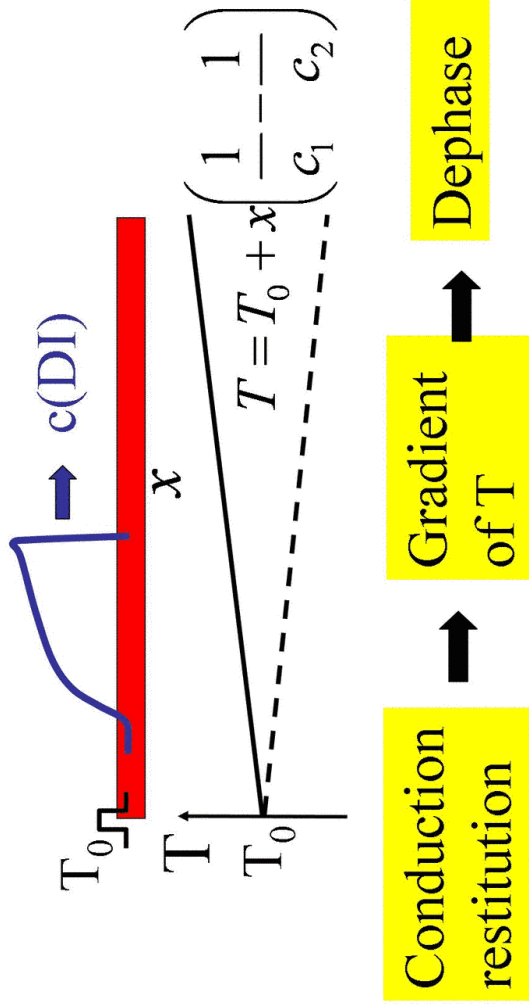
$$T < T^*$$

CV-Restitution

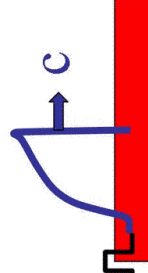



The instantaneous speed of the activation wavefront, or conduction velocity (CV), depends on the local diastolic DI, and thus on position along the cable.


Nonlocal Dephasing



Mechanisms of Discordance

- Conduction restitution 

(Qu et al, 2000; Watanabe et al, 2001; Echebarria-Karma, 2002)
- Ectopic beat 

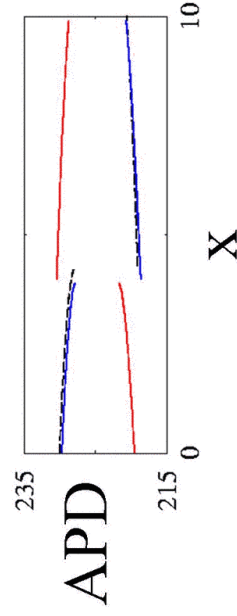
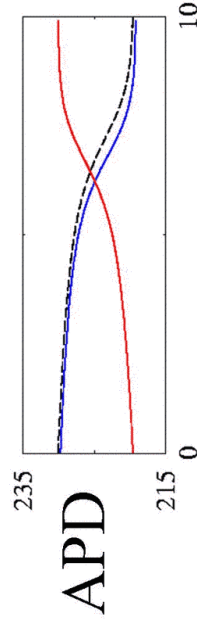
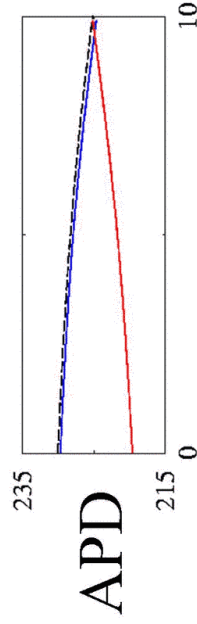
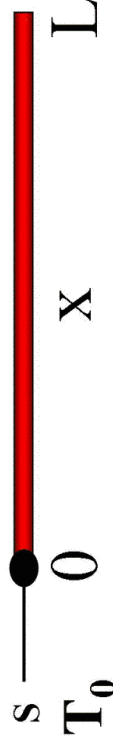
(Watanabe et al, 2000)
- Heterogeneous APD-restitution 

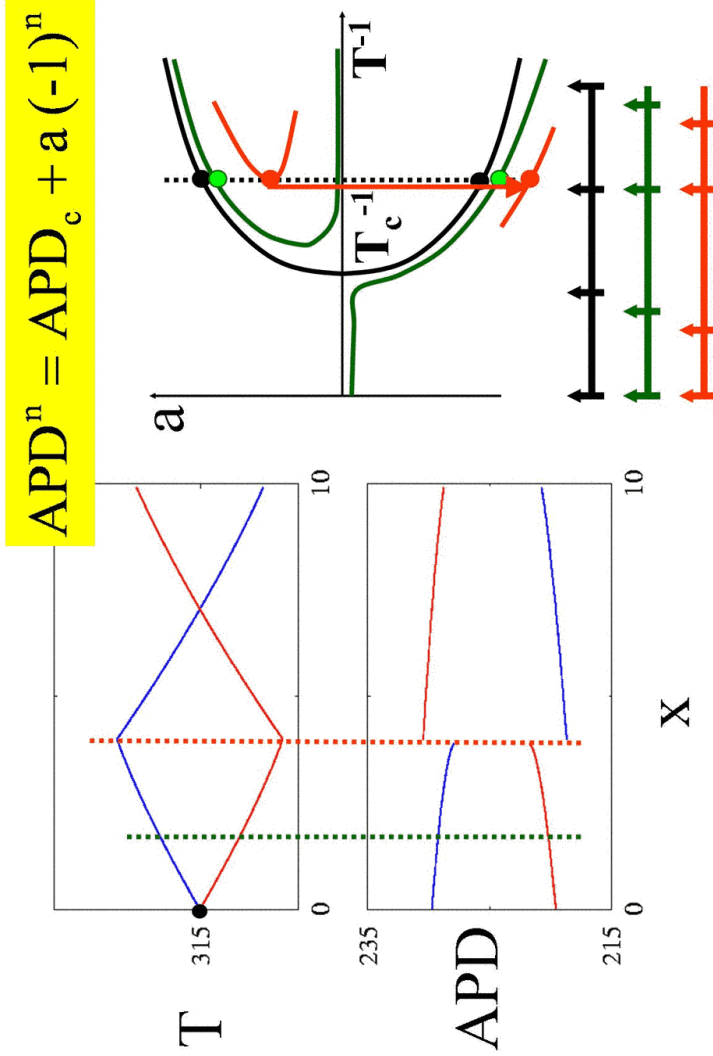
(Echebarria-Karma, 2002)

Extension to tissue of spatially coupled maps of Courtemanche, Glass and Keener (1993)

$$D^n(\mathbf{x}) = T^n(\mathbf{x}) - f(D^{n-1}(\mathbf{x}))$$

$$T^n(\mathbf{x}) = t_{n+1} - t_n = T_0 + \int_0^x \frac{dx'}{c(D^n(\mathbf{x}'))} - \int_0^x \frac{dx'}{c(D^{n-1}(\mathbf{x}'))}$$

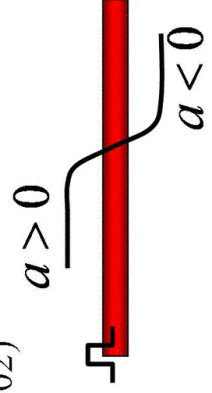




Echebarria-Karma, PRL **88**, 208101 (2002)

$$APD^n \approx APD_c + a e^{im}$$

$$T^n \approx T_c + b e^{im}$$



$\partial_t a = \sigma a - \chi a^3 + \xi^2 \partial_x^2 a - w \partial_x a - b$

$b(x) = b(0) + \frac{1}{\Lambda} \int_b^x dx' a(x')$

$$\xi \sim (D \times APD_c)^{1/2} \quad w \sim D/c \quad \Lambda = \frac{c^2}{2c'}$$

Linear stability

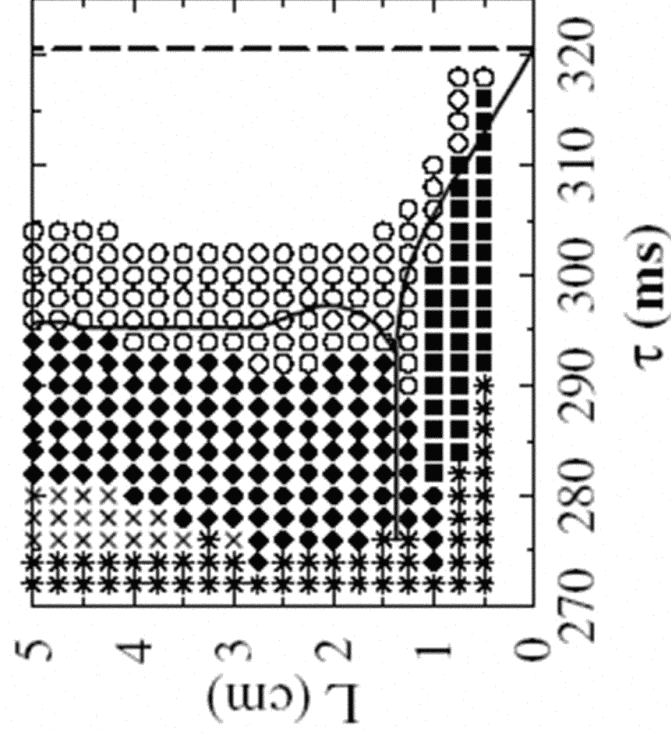
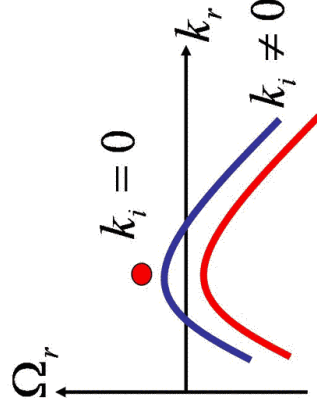
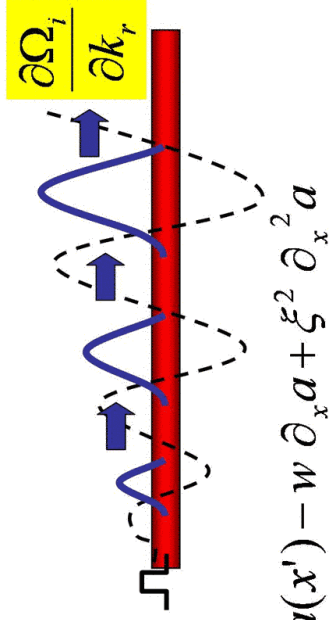
$$a \sim \exp(ikx + \Omega t)$$

$$\partial_t a = \sigma a - \frac{1}{\Lambda} \int dx' a(x') - w \partial_x a + \xi^2 \partial_x^2 a$$

$$\Omega = \sigma - \xi^2 k^2 - i \left(wk - \frac{1}{\Lambda k} \right)$$

$$\lambda_{stationary} \sim (w\Lambda)^{1/2}$$

$$\lambda_{traveling} \sim (\xi^2 \Lambda)^{1/3}$$



$$\lambda_{stationary} = 2\pi(w\Lambda)^{1/2}$$

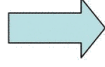
TABLE I. Values of various lengths in cm with λ_{theor} [Eqs. (12) or (13)] and λ_{sim} [from simulations of Eq. (7)].

Model	Λ	w	ξ	$\lambda_{theor}/4$	$\lambda_{sim}/4$	L_{min}
Noble	49.1	0.045	0.18	2.33	2.6	2.75
Two-variable	3.55	0.031	0.235	1.33	1.1	1.15

$$w \sim D/c \quad \Lambda = \frac{c^2}{2c'}$$

Conclusion

Period doubling
 + nonlocal dephasing
 + advection and diffusion

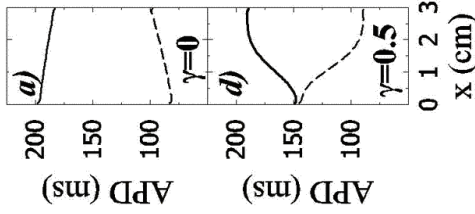


Spatial patterns via finite k instability
 stationary (advection dominated)
 convective (diffusion dominated)

Control

$$T^n = T_0 + \frac{\gamma}{2}(APD^n - APD^{n-1})$$

$$a(x, t) = \psi(x)e^{\Omega t}$$



Forced Helmholtz equation

$$\xi^2 \frac{d^2 \psi}{dx^2} + (\sigma - \Omega) \psi = \gamma \psi(0)$$

$$\Omega_0 = \sigma - \gamma$$

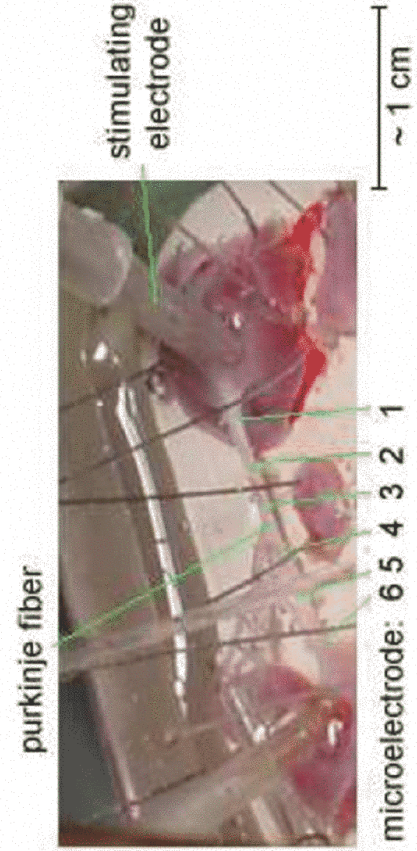
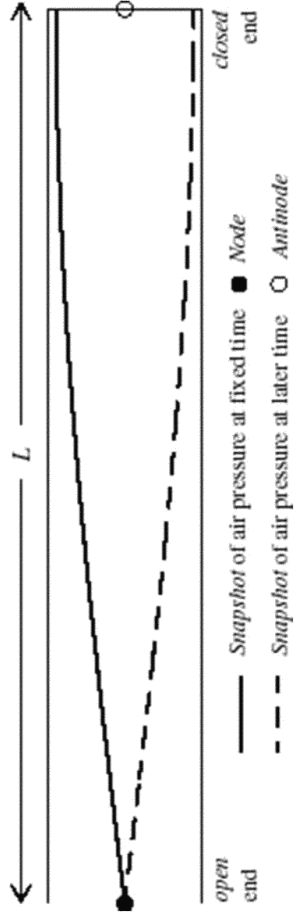
$$\Omega_n = \sigma - \frac{n^2 \pi^2 \xi^2}{L^2}$$

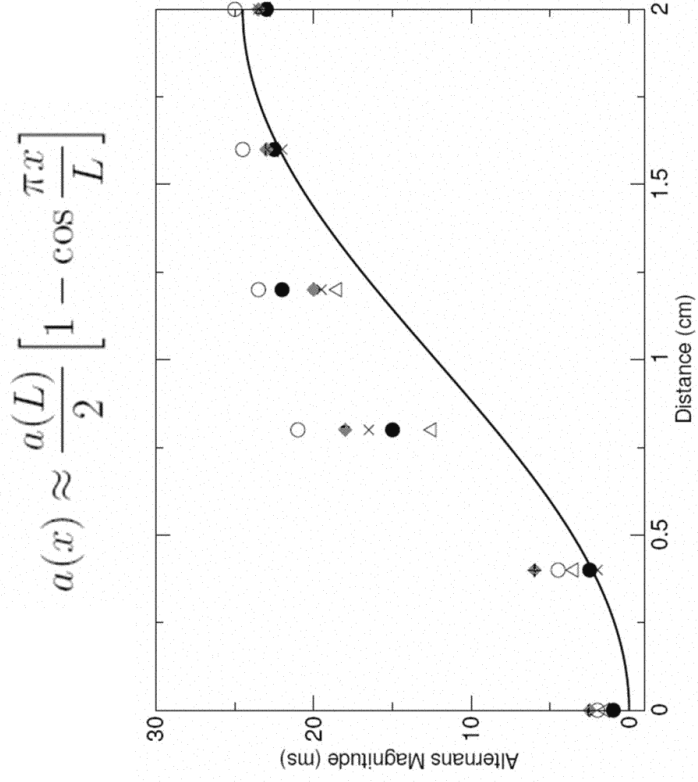
$$\psi_n \sim 1 + \left(\frac{n^2 \pi^2 \xi^2}{\gamma L^2} - 1 \right) \cos \frac{n\pi x}{L}$$

Echebarria-Karma (2002)

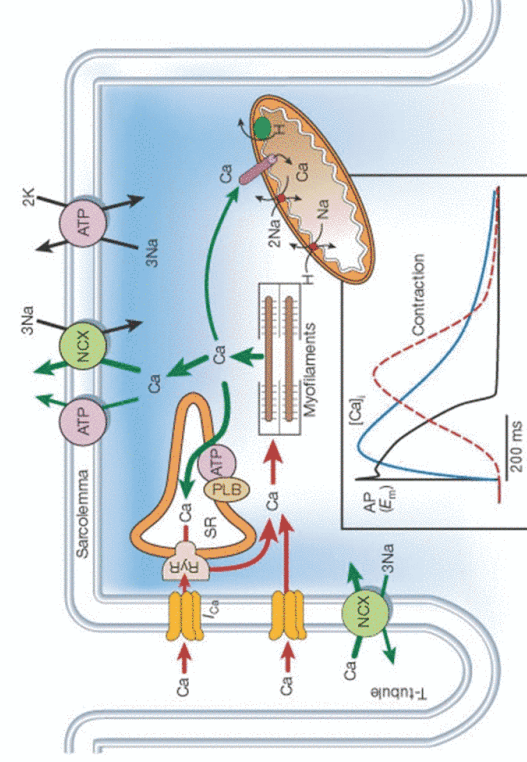


Fundamental Mode: Standing Sound Wave in a Pipe



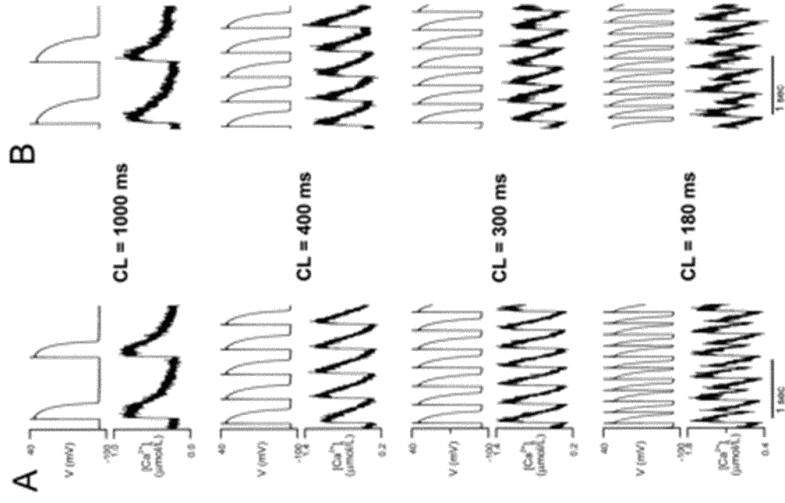


EC coupling



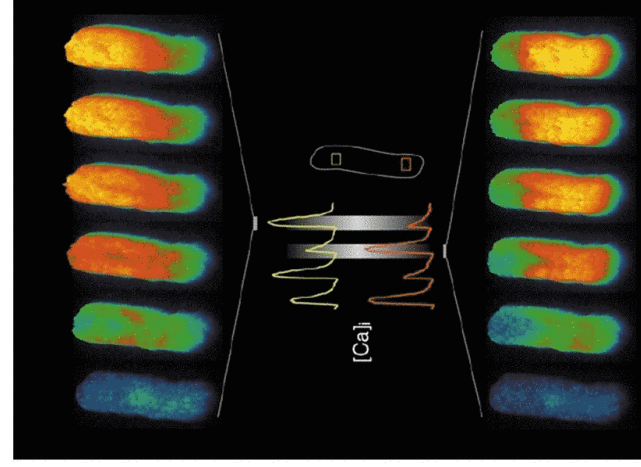
Who is the master and who is the slave???

Chudin *et al.* (2003)



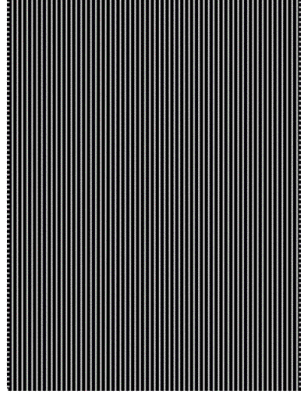
Subcellular discordant alternans

Blatter *et al.* (2003)

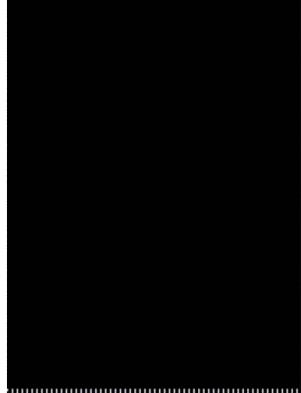


Complex APD-Ca_i Behavior in the Single Myocyte

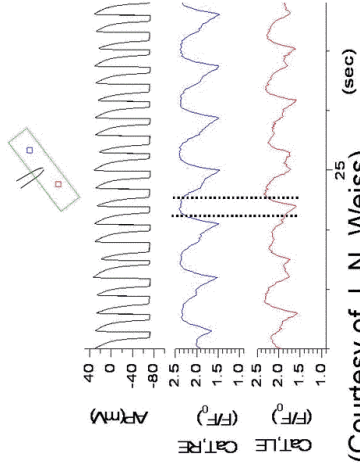
Ca_i Fluorescence



Ca_i alternans phase map

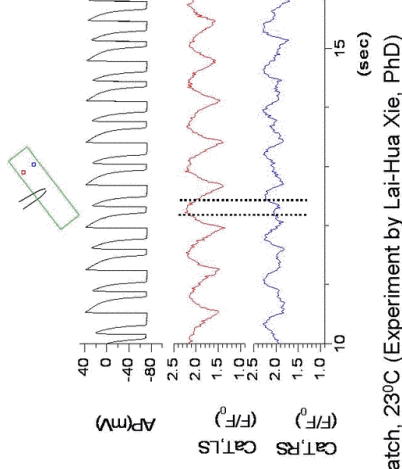


Blue = high-low
Red = low-high

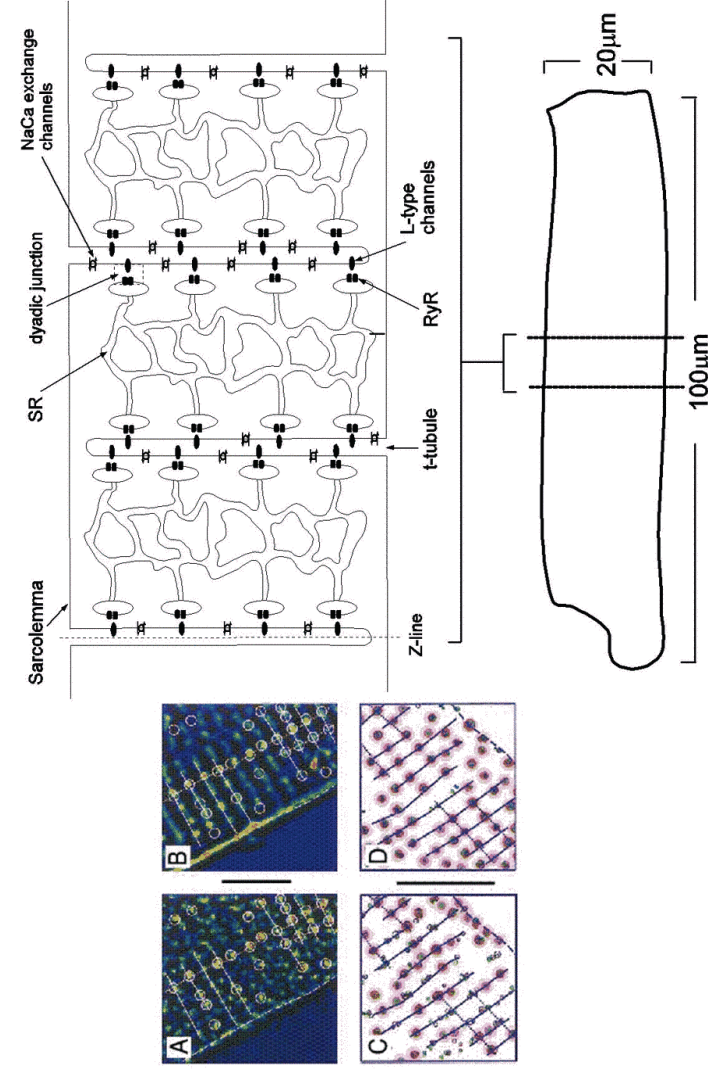
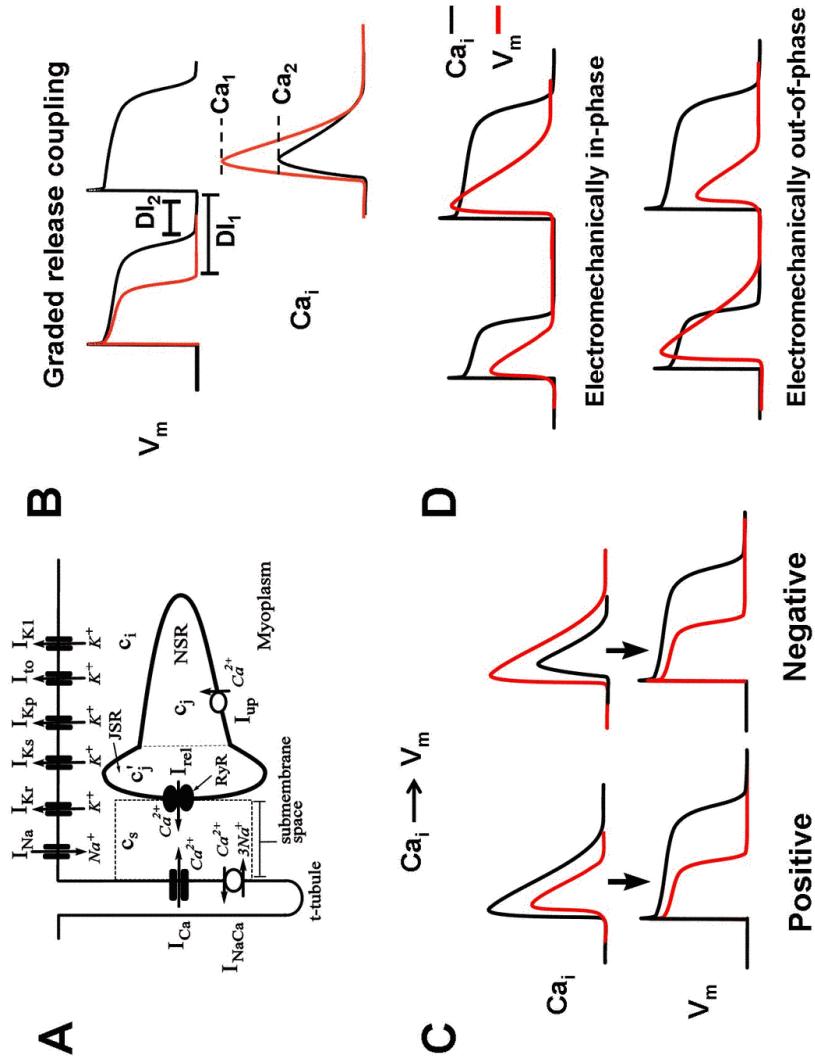


(Courtesy of J. N. Weiss)

Rabbit ventricular myocyte, perforated patch, 23°C (Experiment by Lai-Hua Xie, PhD)

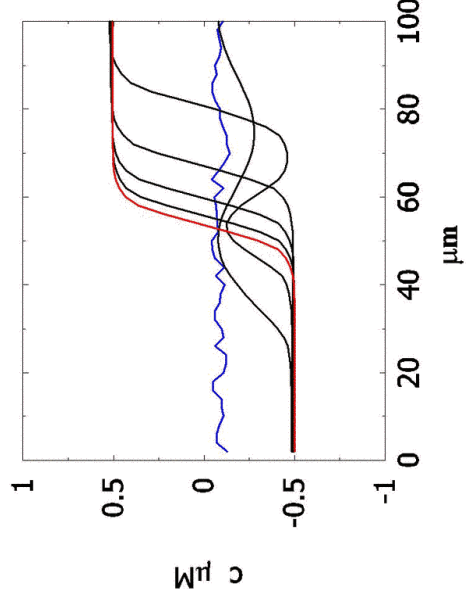


Are spatially discordant subcellular scale calcium alternans of static or dynamic origin ?



Results & interpretation

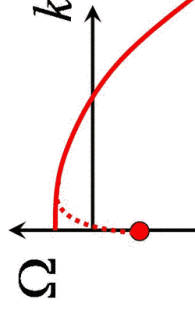
Shiferaw & Karma, PNAS (2006)



$$APD^n = APD^* - a (-1)^n$$

$$Ca_i^n = Ca_i^* + c (-1)^n$$

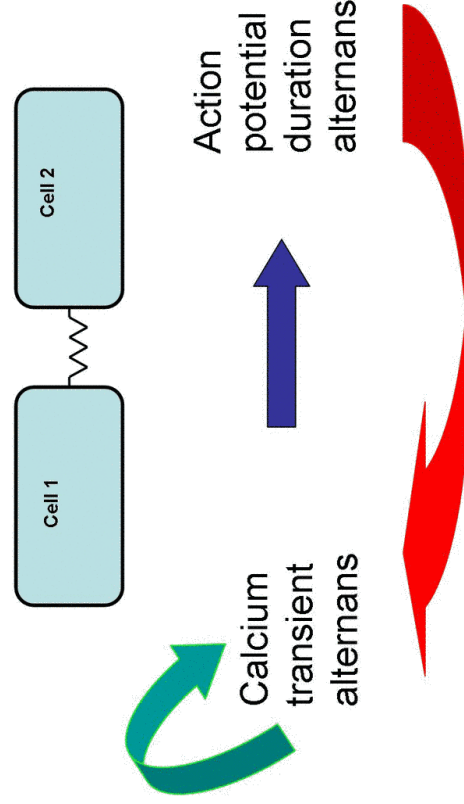
Turing-like instability



$$\partial_t a = \alpha a + \beta \bar{c}$$

$$\partial_t c = \gamma a + \delta c + \mu \bar{c} + D \nabla^2 c$$

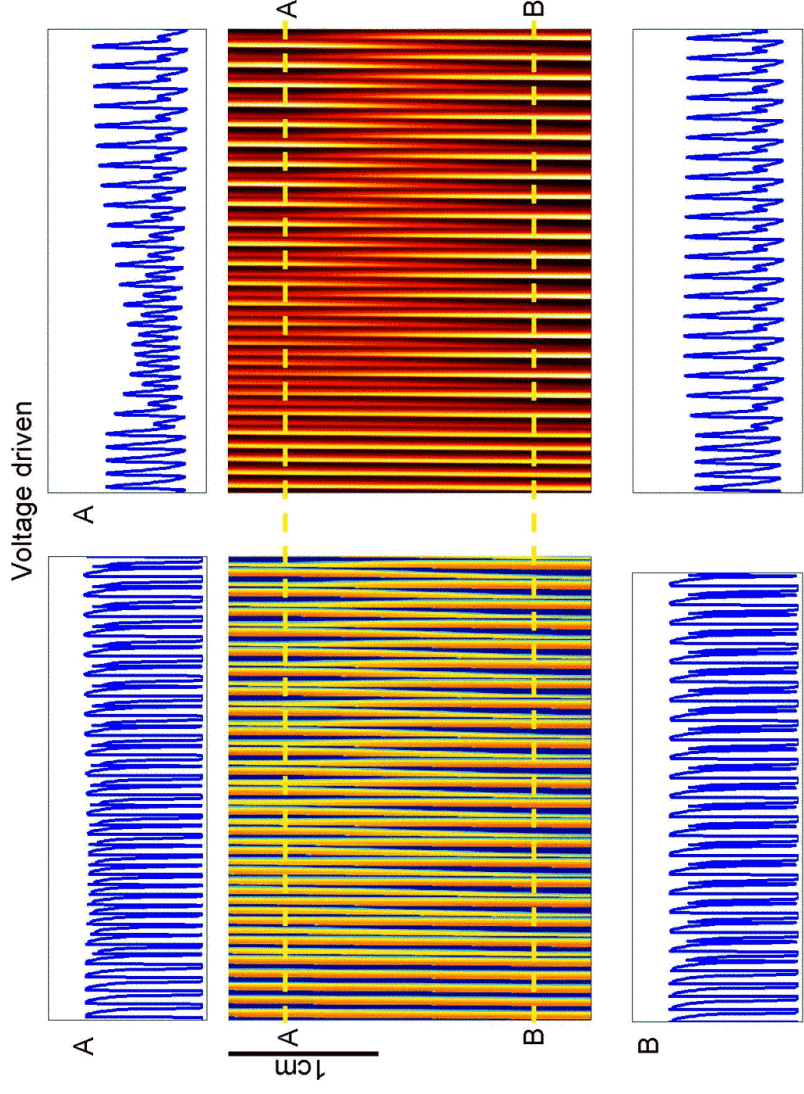
Desynchronization mechanism



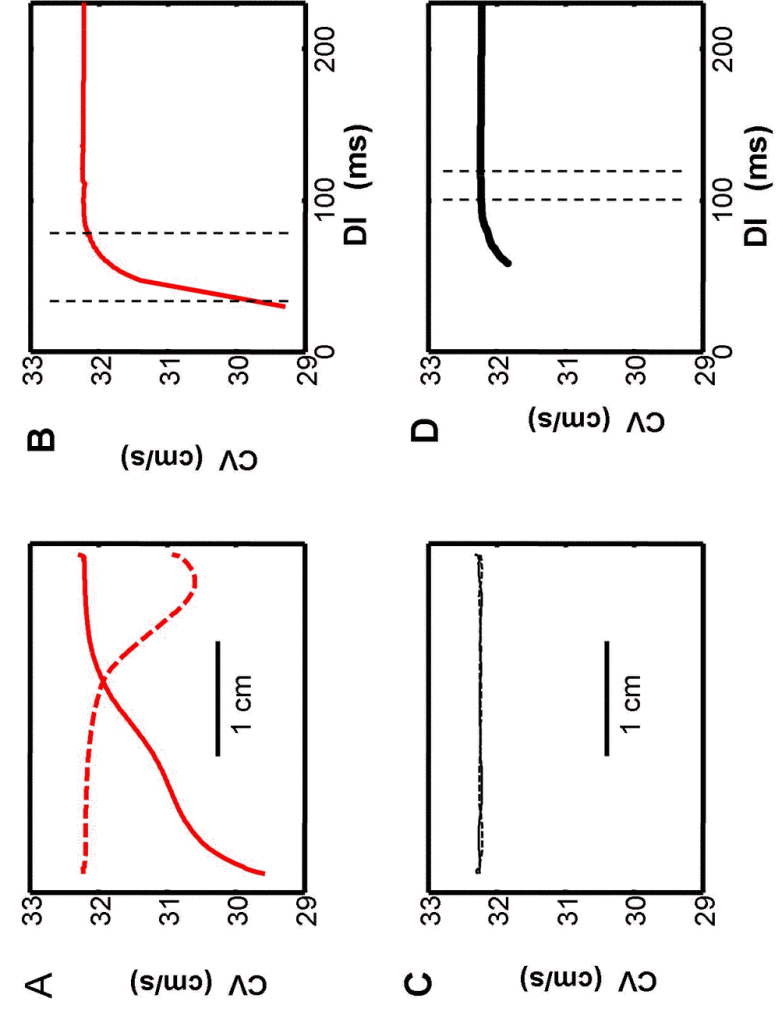
Conclusion

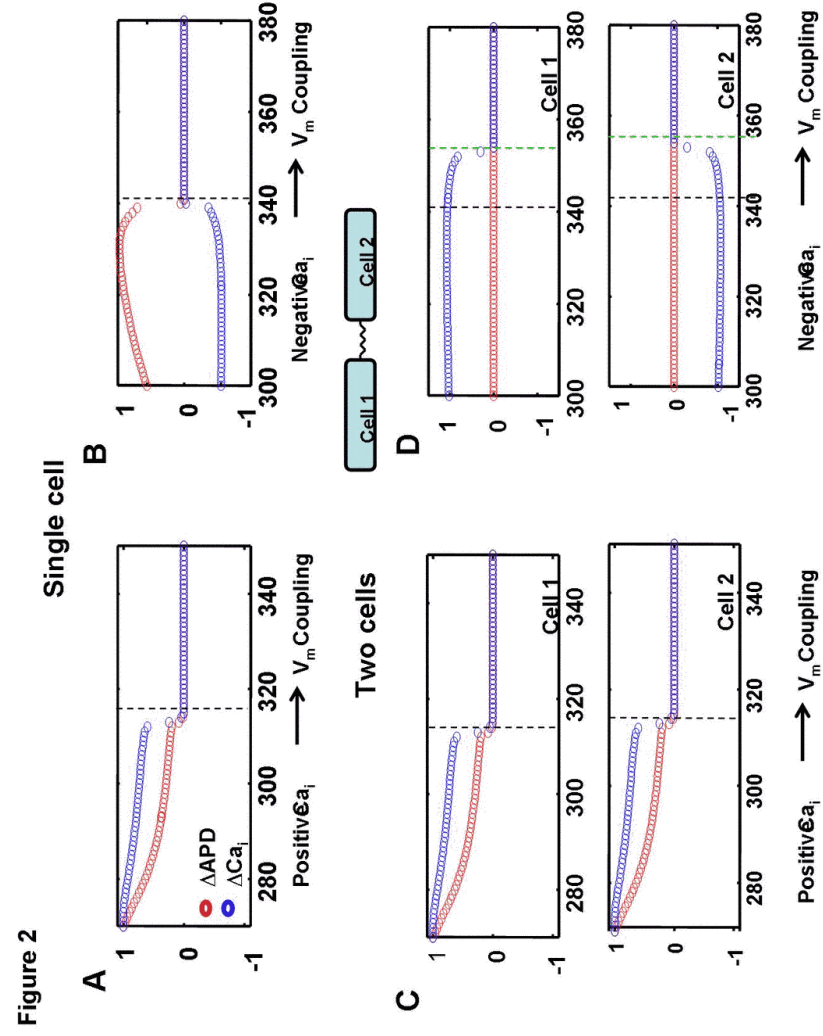
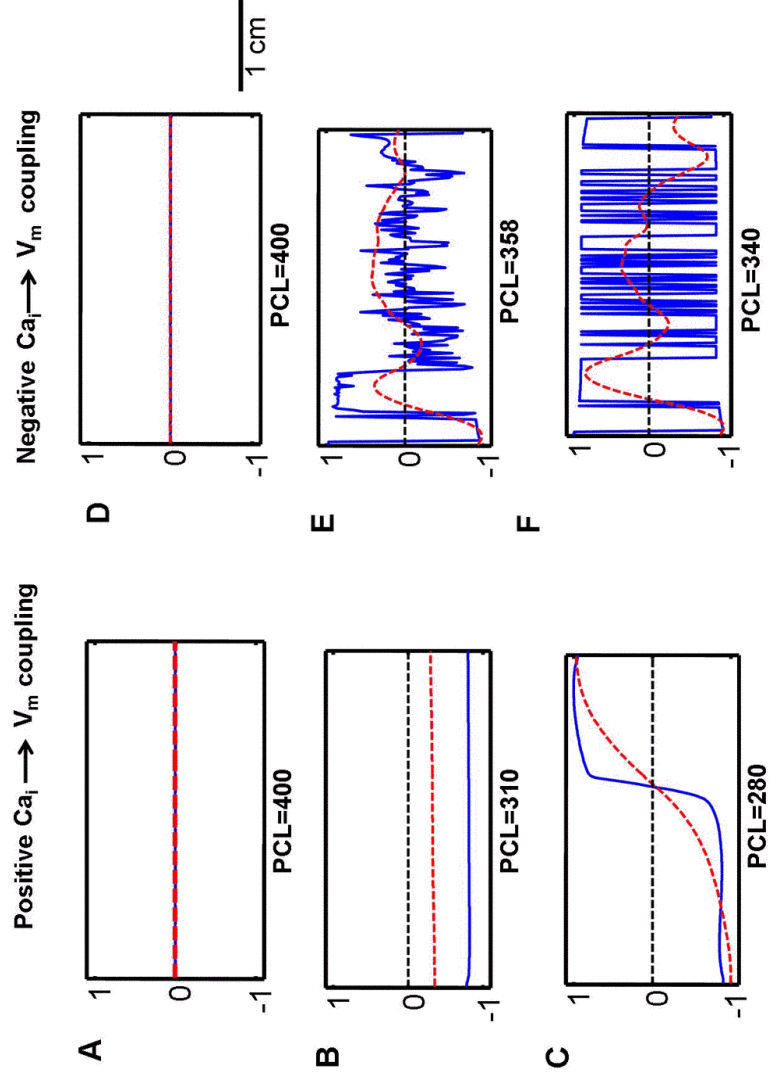
Turing-like instability mediated by calcium (activator) and voltage (inhibitor) diffusion leads to subcellular patterns of spatially discordant calcium alternans.

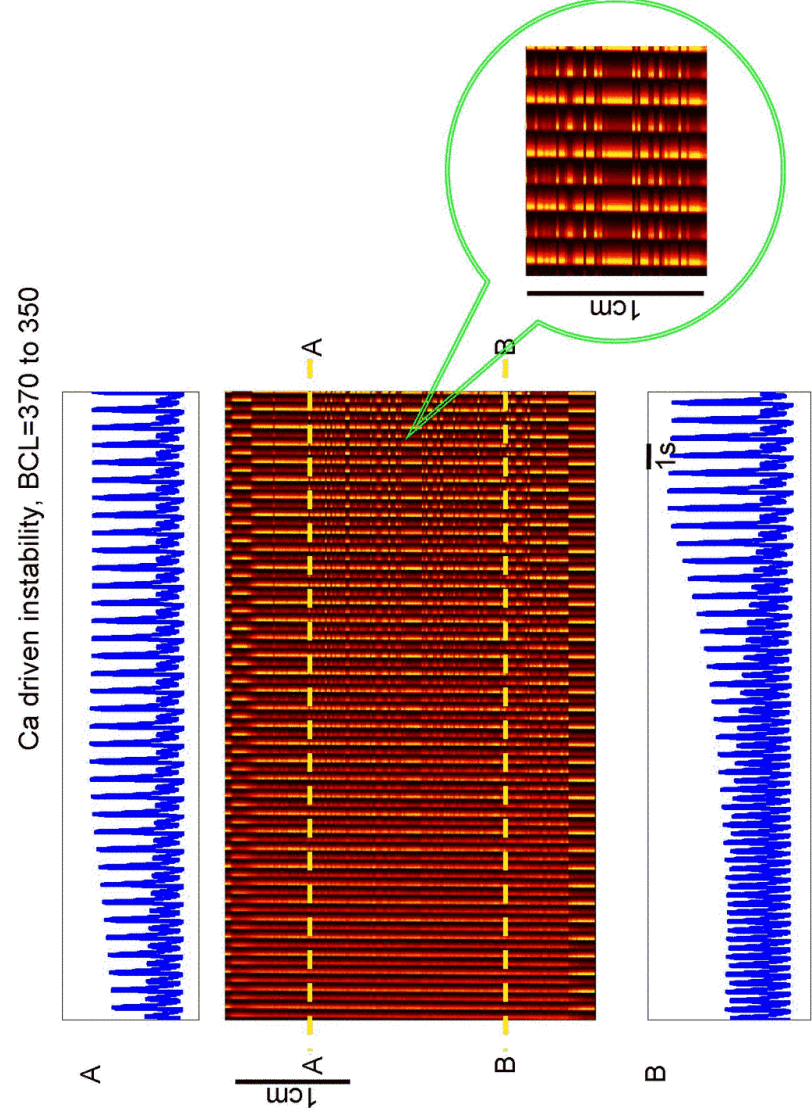
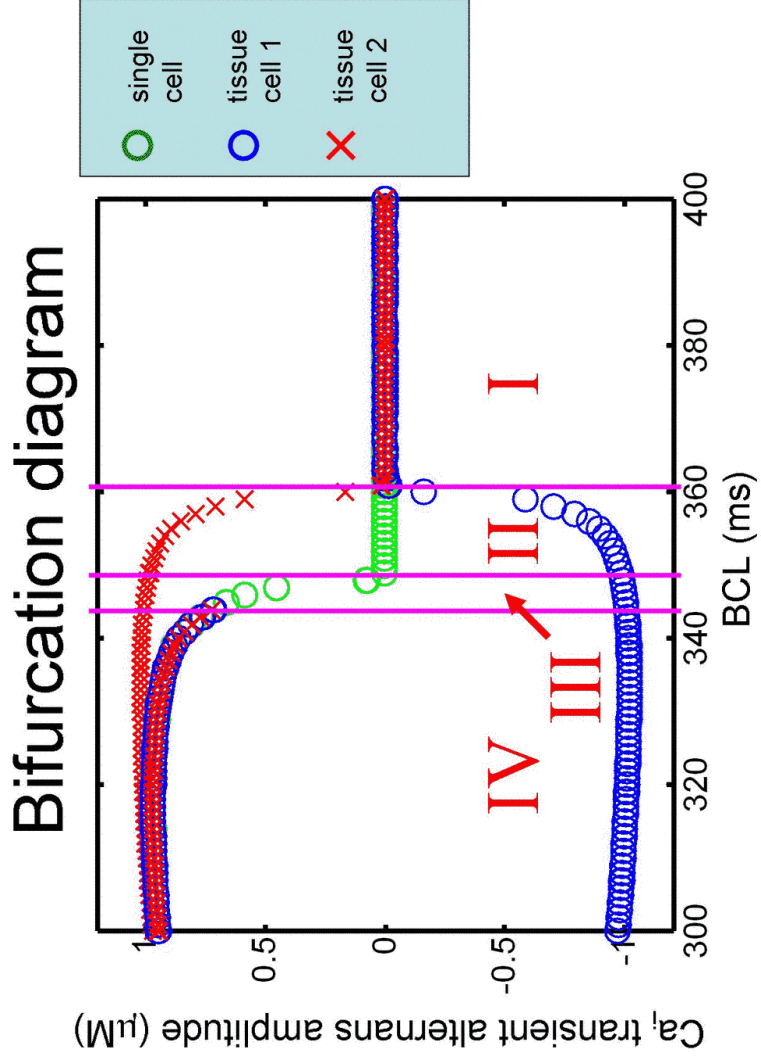
What are the tissue scale manifestations of desynchronization ?



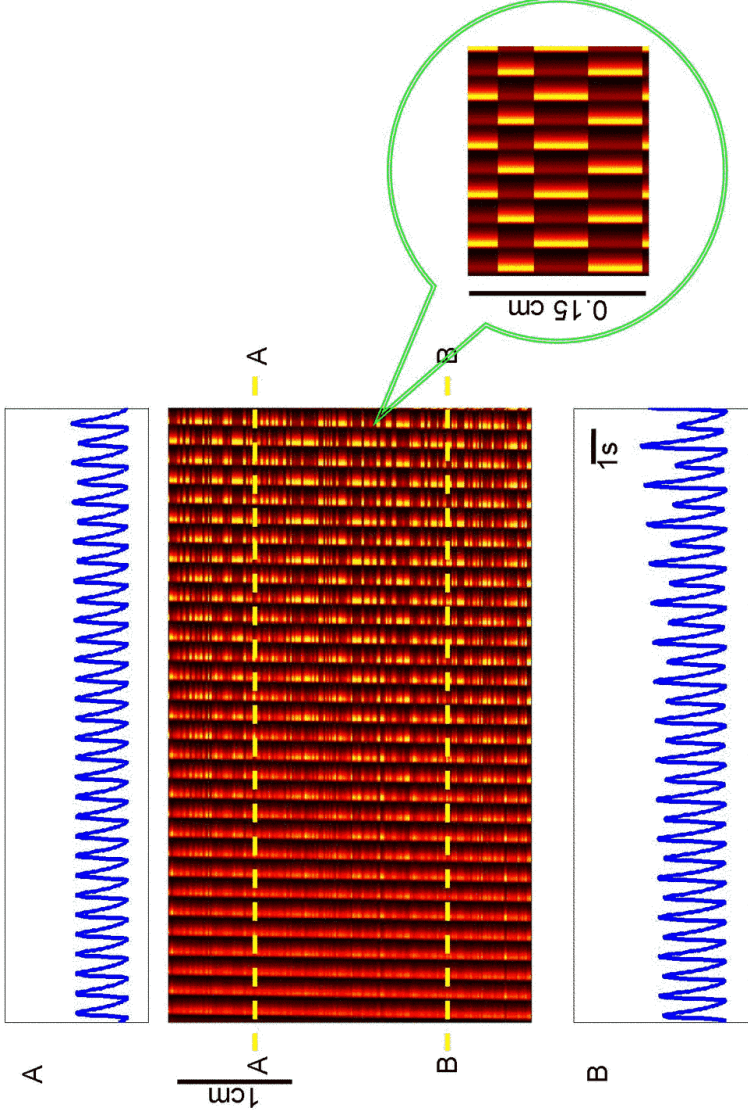
CV and CV restitution during discordant alternans



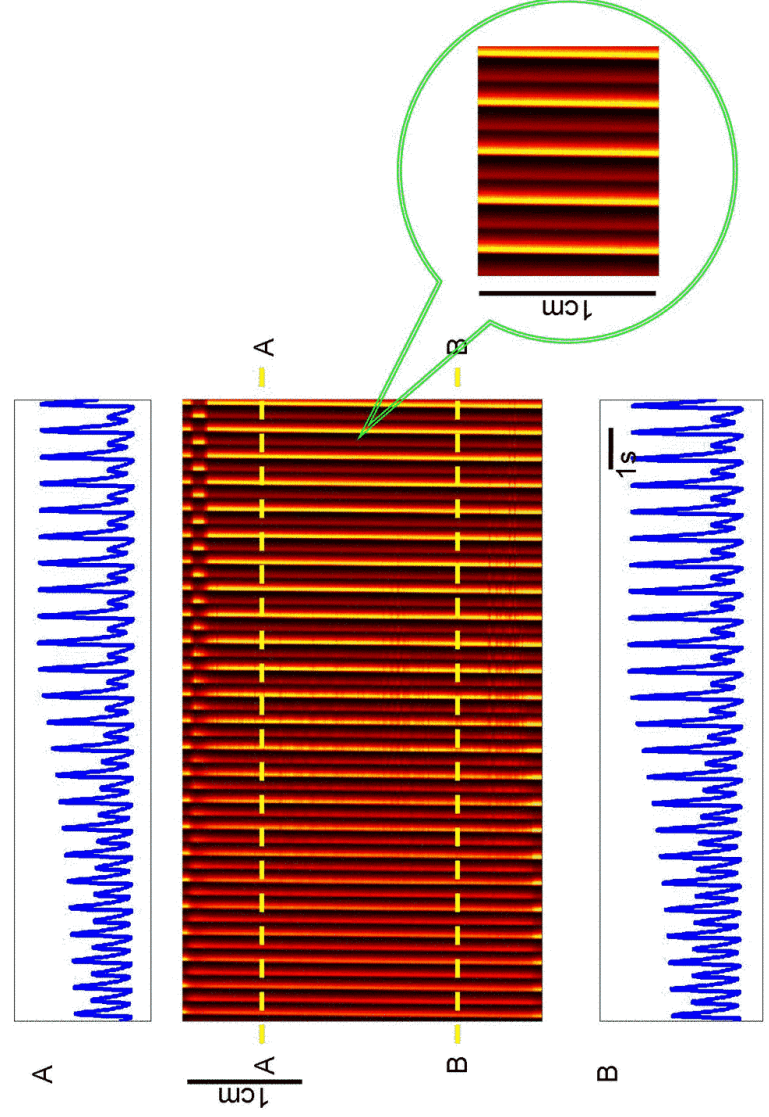




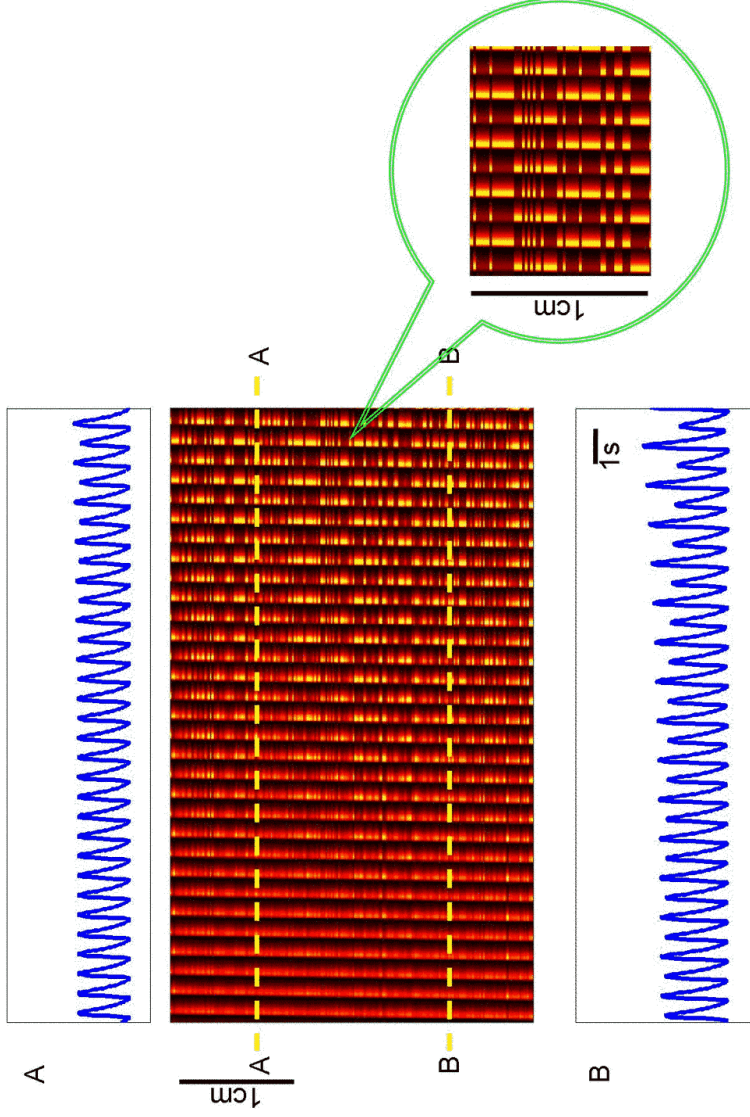
Ca driven instability, BCL=500 to 320



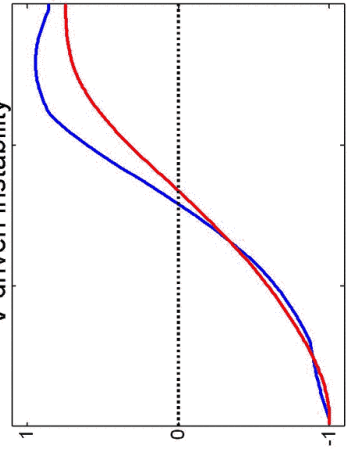
Ca driven instability, BCL=370 to 320



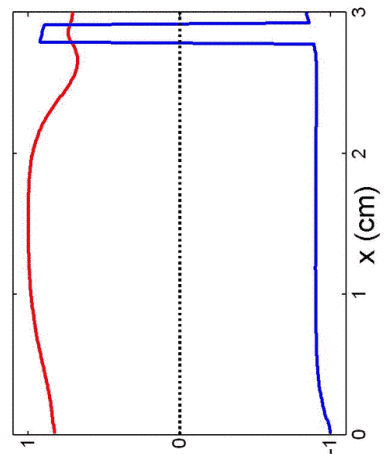
Ca driven instability, BCL=500 to 320



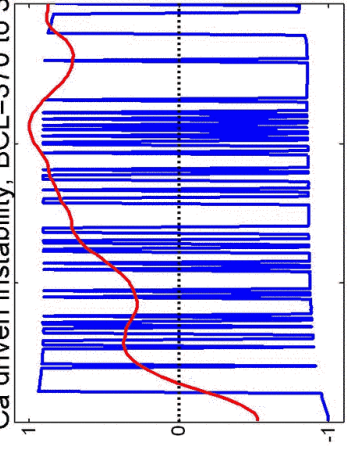
V driven instability



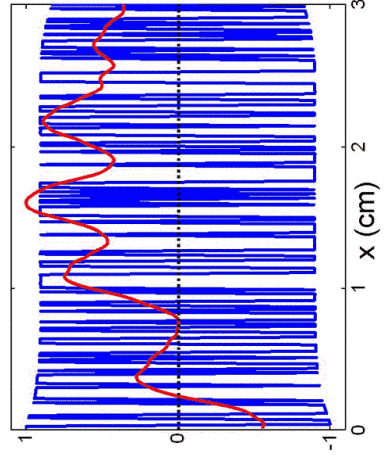
Ca driven instability, BCL=370 to 320

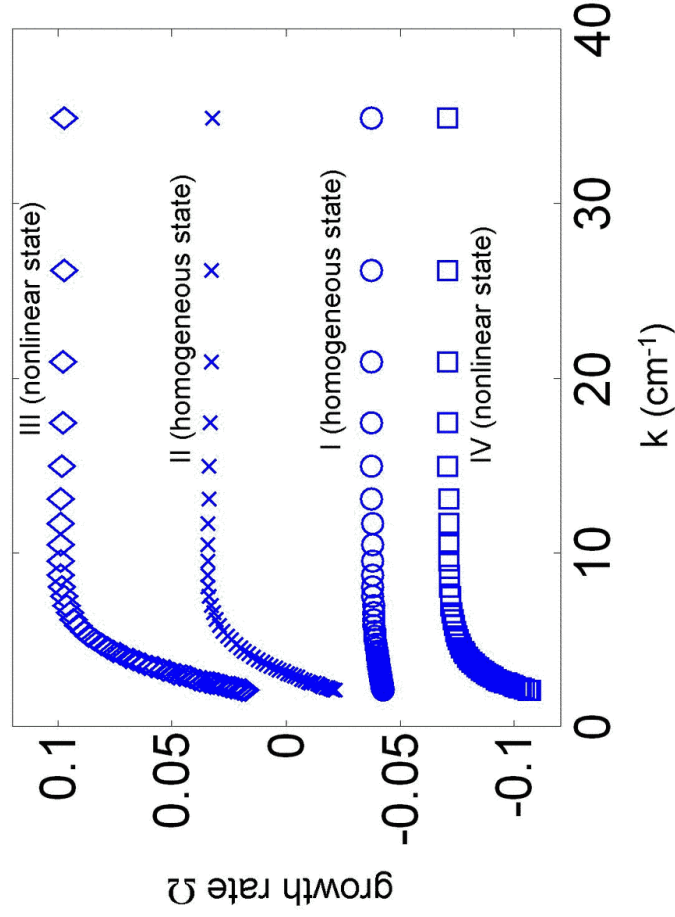


Ca driven instability, BCL=370 to 350



Ca driven instability, BCL=500 to 320





Conclusion

- Patterns of alternans influenced by both sign of coupling and initial conditions
- Spatial scale of nodal region related to cellular mechanism of alternans

