

# $H_{c2}$ Limiting and the Superconducting Double Transitions in $\text{Sr}_2\text{RuO}_4$



Solving the remaining puzzles

Yoshi Maeno  
*Kyoto University*

Collaborators on this issue:

Kittaka, Nakamura, Ishida, Aono, Yaguchi,  
Deguchi, Tanatar

# Symmetry of the superconducting state

Experimental identification:

1. Const. inplane spin susceptibility  
(inplane equal-spin pairing:  $d // z$ )  
There should be no Pauli limiting  
for in-plane field.

2. Intrinsic magnetism (T-violation)  
(out-of-plane orbital moment)

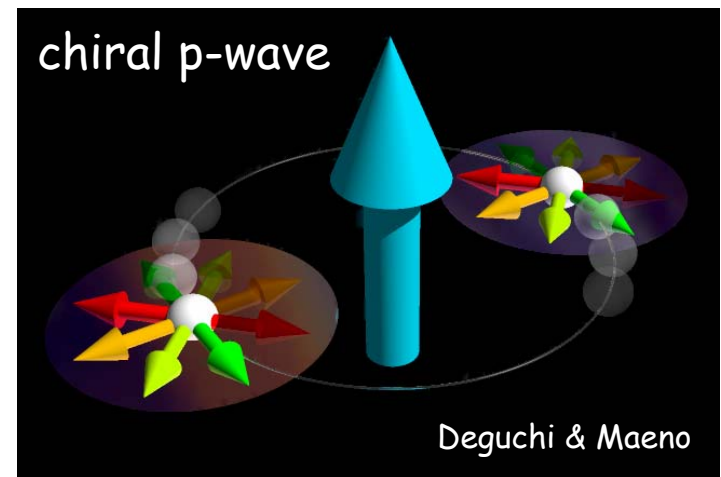
$$\vec{d} = \hat{z} (k_x \pm i k_y)$$

$\downarrow$   $\searrow$   
 $\frac{1}{\sqrt{2}} [|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle]$   $L_z = +1, -1$

- chiral: orbital magnetic moment
- degeneracy: 2

**Chiral p-wave state**

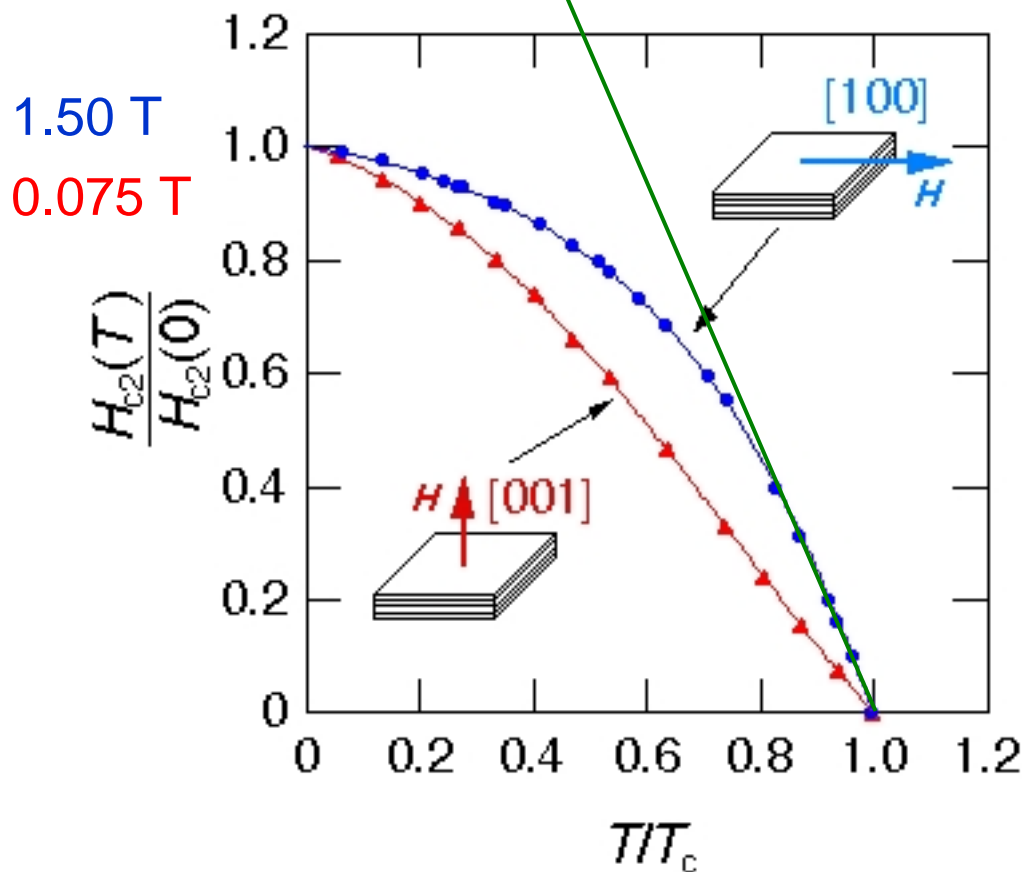
$\Gamma$	$\vec{d}$	
$A_{1u}$	$\vec{d} = \hat{x}k_x + \hat{y}k_y$	} B-phase
$A_{2u}$	$\vec{d} = \hat{x}k_y - \hat{y}k_x$	
$B_{1u}$	$\vec{d} = \hat{x}k_x - \hat{y}k_y$	
$B_{2u}$	$\vec{d} = \hat{x}k_y + \hat{y}k_x$	
$E_u$	$\vec{d} = \hat{z}(k_x \pm i k_y)$	} A-phase



# Experimental Fact 1:

## $H_{c2}$ Limiting Behavior for $H // ab$

WHH curve would predict  
much larger  $H_{c2} // ab$   
(Kittaka's presentation)

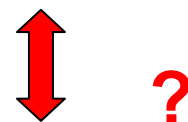


1.50 T  
0.075 T

$H // c$ :

ordinary behavior  
(conventional orbital limiting only)

$H // ab$ : unusual  $H_{c2}$  suppression at low  $T$



Pauli limiting should be irrelevant for  $H // ab$ .

(NMR Knight shift shows no change.)

$$\frac{1}{2} \chi H_p^2 = \frac{1}{8\pi} H_c^2 \quad \mu_0 H_p = 1.6 T$$

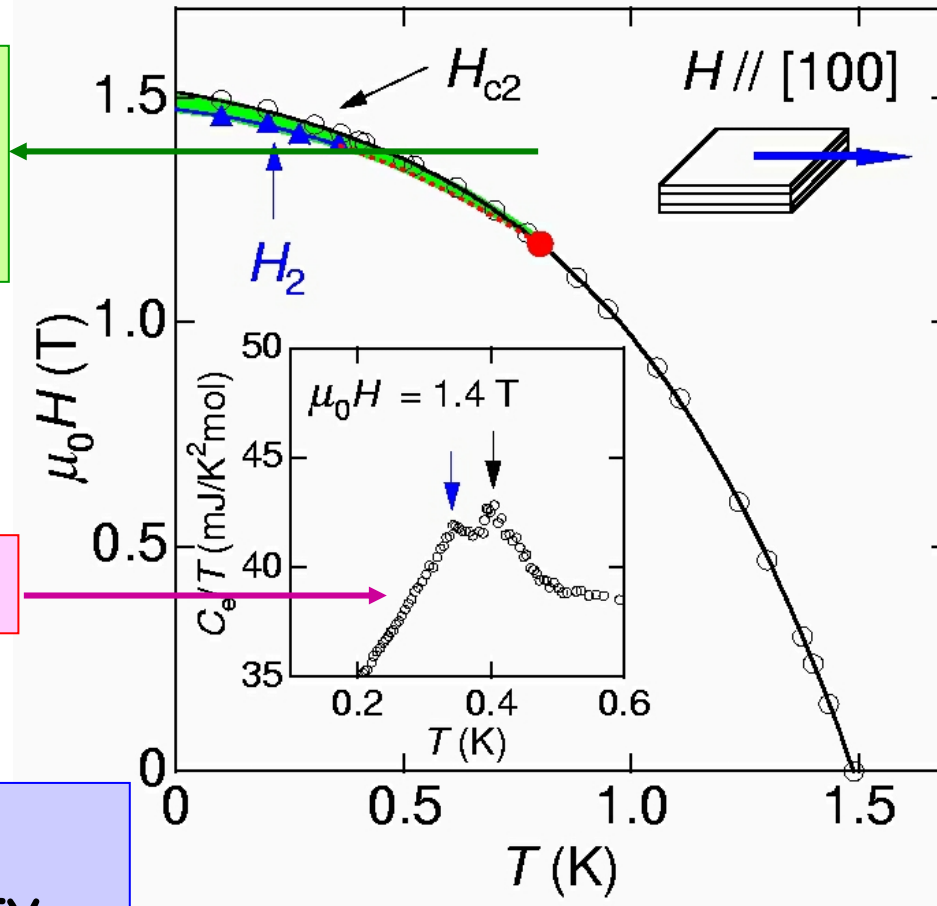
# Experimental Fact 2:

## Emergence of a Second Transition: The "Deguchi Phase"

Ru-NMR Knight  
does not change.  
Ishida et al. (2007).

Specific Heat

Observed also in  
thermal conductivity  
and AC susceptibility.

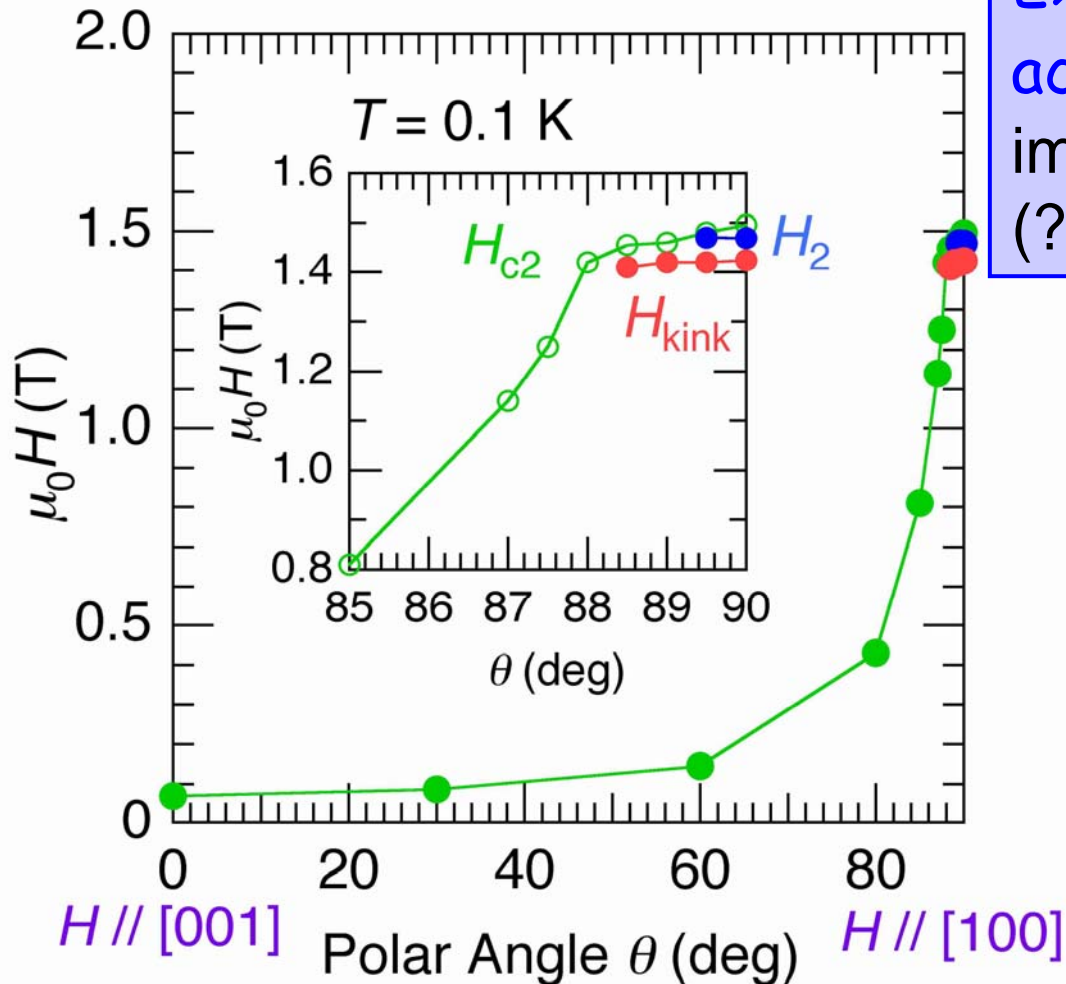


Double transitions

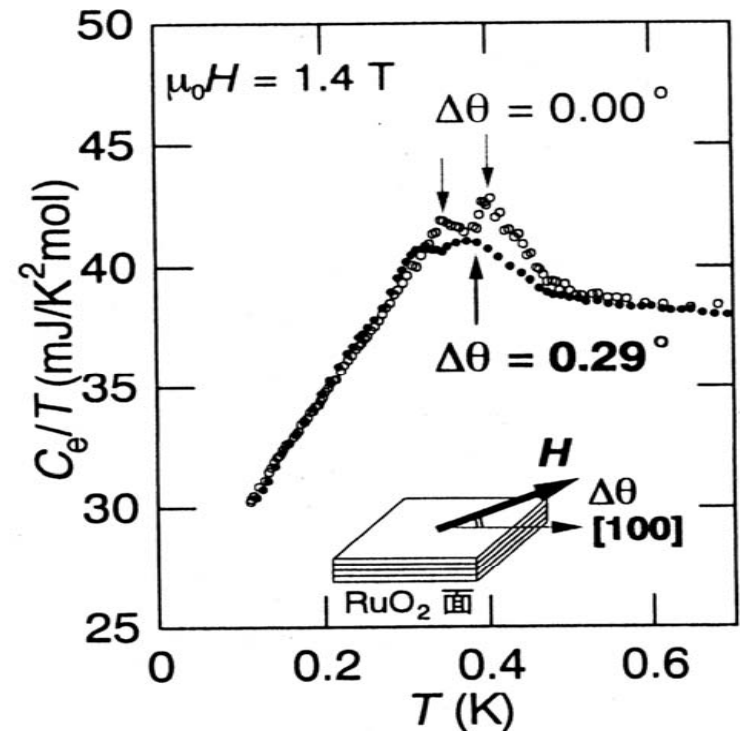
Deguchi *et al.*

# Facts 1 and 2:

## $H_{c2}$ Suppression and Double Transitions are Linked



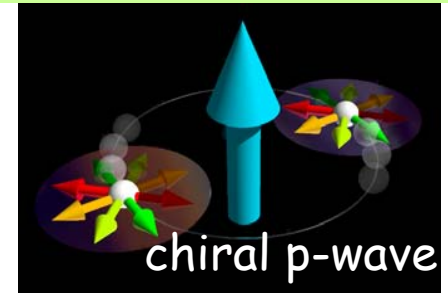
Extreme sensitivity to the accurate field alignment implies **orbital** mechanism (??).



# Proposed models for the second phase

State at zero  $H$ :  $d = z(k_x + ik_y)$

States under  $H // x$ :



Model	Spin	orb. (low $H_x$ )	orb. (high $H_x$ )	boundary	$H_{c2}$ limiting?
Agterberg et al.	$z$	$k_x + i\epsilon k_y$	$k_x$	2 <sup>nd</sup> -order	No
Udagawa et al.	$z - i\alpha y$	$k_x + i\epsilon k_y$	$k_x$	Crossover (Orb. 2 <sup>nd</sup> -order transition at lower $H$ )	No
$^3\text{He-A1}$	$z - iy$	$k_z - ik_y$ $\ell // d$ by dipole int.		2 <sup>nd</sup> -order	No

# Double Transitions: Theoretical Expectation

Orbital Scenario: D. Agterberg, PRL 80, 5184 (1998);

Kaur, Agterberg, and Kusunose, PR B72, 144528 (2005).

Some extension to simulate the experimental results

$$d = \hat{z} \Delta_0 k_{x'}$$

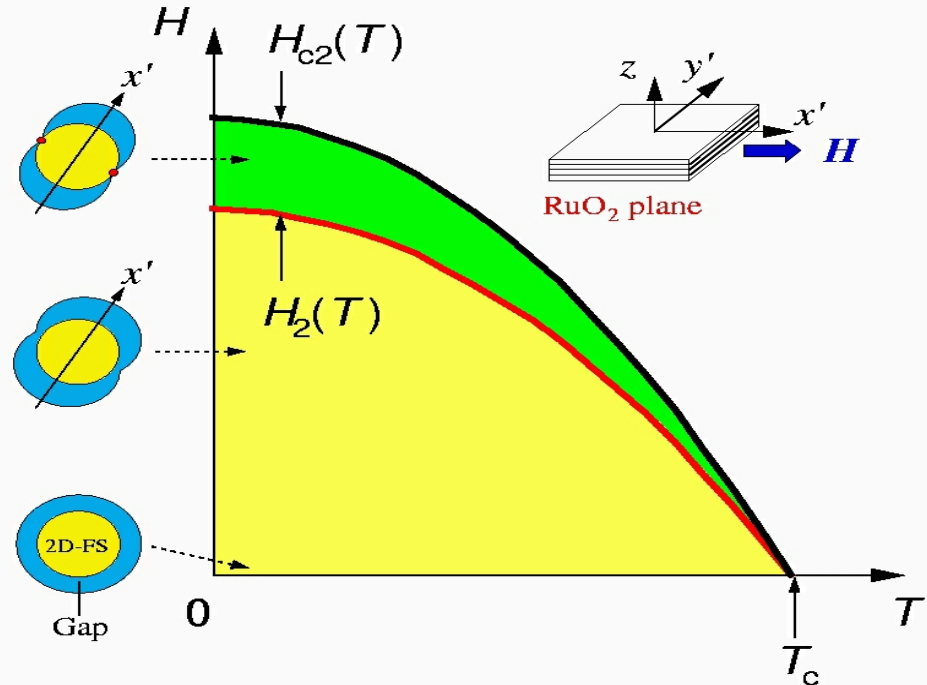
**line-node**  
( $\varepsilon = 0$ )

$$d = \hat{z} \Delta_0 (k_{x'} + i\varepsilon k_{y'})$$

**anisotropic gap**  
( $0 < \varepsilon < 1$ )

$$d = \hat{z} \Delta_0 (k_{x'} + i k_{y'})$$

**isotropic gap**  
( $\varepsilon = 1$ )



Spin + orbital scenario: M. Udagawa *et al.*, JPSJ 74, 2905 (2005).

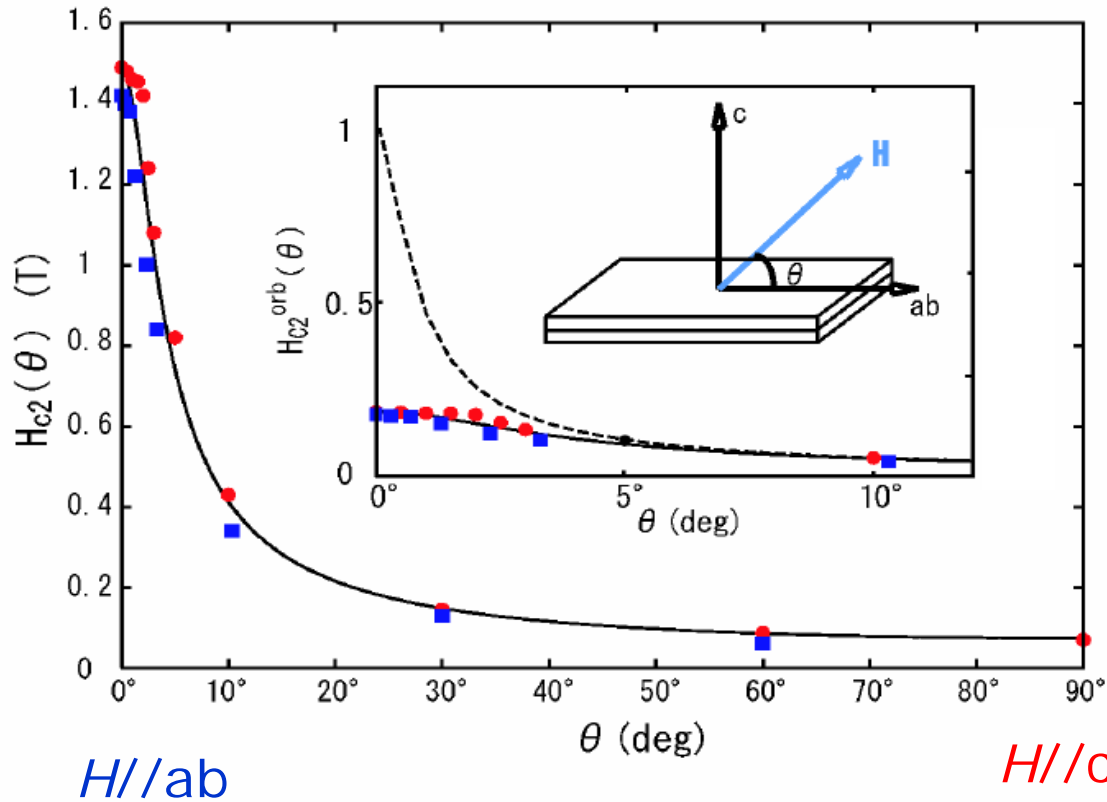
considers a non-unitary state  $(z - i\alpha y) k_x$  just below  $H_{c2}$ .

i.e.,  $|\uparrow\uparrow\rangle_x$  component.

In both scenarios,  $k_z$  is ignored because of Q2D.

# The $H_{c2}$ Suppression:

Can it be explained by Pauli limiting?



Machida and  
Ichioka (2007)

Need to use

$\Gamma$  (anisotropy of  $\xi$ ) = 107

$\Gamma$  may be estimated from  
the  $H_{c2}$  ratio near  $T_c$ .

Experimentally,

$\Gamma(T \rightarrow T_c)$  is not so large.

(Kittaka's presentation).

A single-band model is  
not appropriate, either.

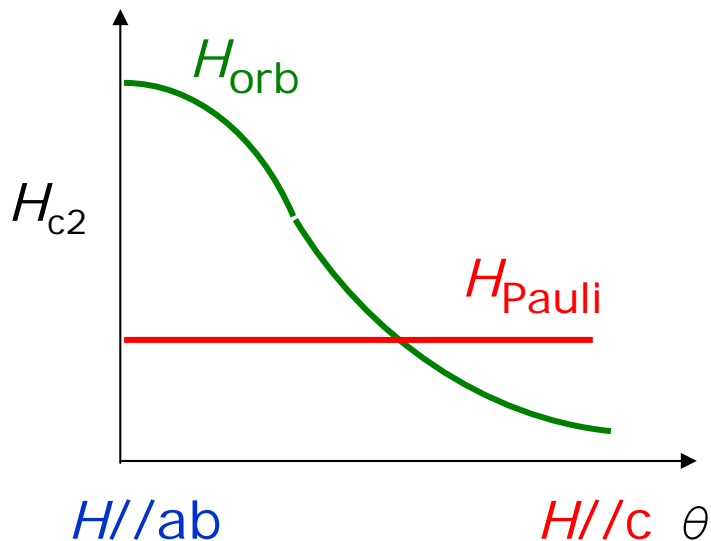
$$\frac{H_{c2}(\theta)}{H_{c2}^{orb}(\theta=0)} = \frac{1}{\sqrt{\Gamma^2 \sin^2 \theta + \cos^2 \theta + 2.4 \mu_0^2}}$$



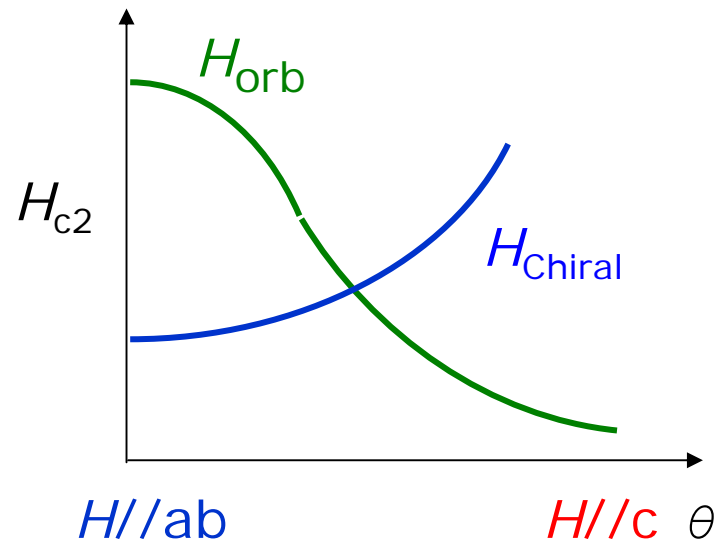
# Origin of $H_{c2}$ ( $//ab$ ) Limiting in $\text{Sr}_2\text{RuO}_4$

1. **Pauli Limiting** theory is interesting.

But it is **NOT consistent** with thermodynamic and other observations.

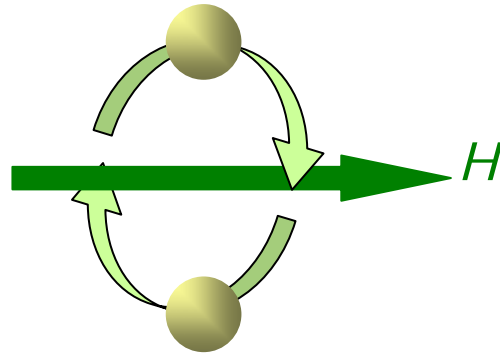
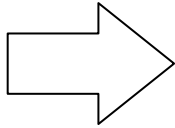


2. Novel **ORBITAL** limiting mechanism **specific to the chiral triplet pairs** seems to be operative for  $H//ab$ .

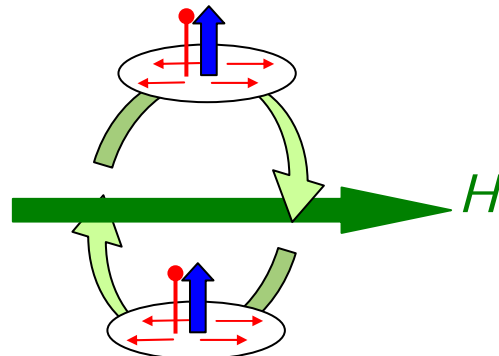
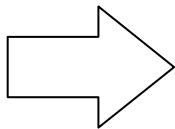
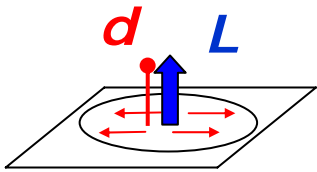


# Additional Orbital Depairing Mechanism?

Spin singlet



Chiral spin triplet

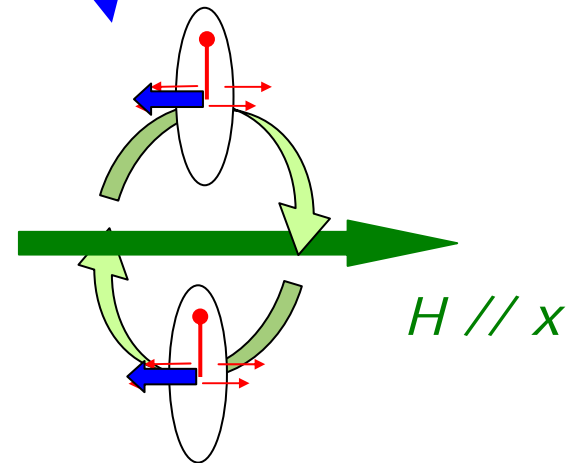


Conventional orb.  
limiting.

This state must

This effect should be too small and negligible.  
I.A.Luk'yanchuk and V.P.Mineev,  
JETP **66**, 1168 (1987).

$$d = z(k_z + ik_y)$$



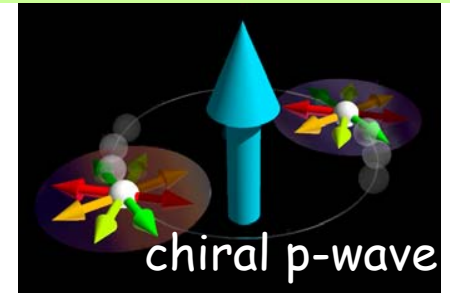
Additional effect due  
to the L-vector of a  
Cooper pair.

$$d = z(k_x + ik_y)$$

# Proposed models for the second phase

State at zero  $H$ :  $d = \mathbf{z}(k_x + ik_y)$

States under  $H // x$ :



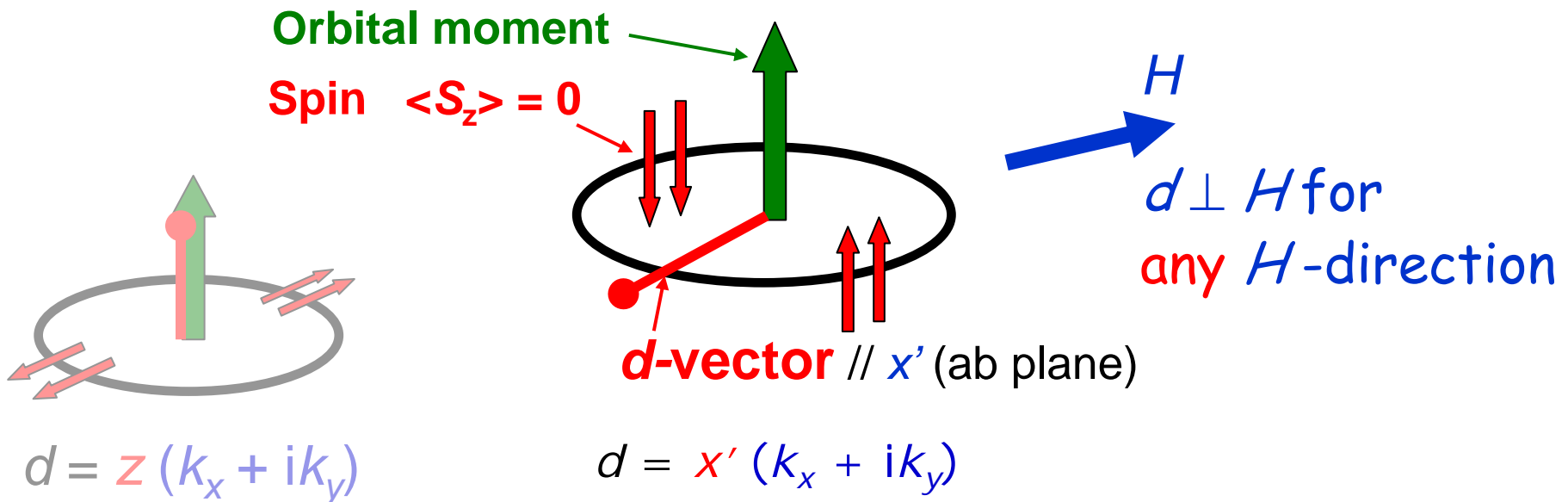
Model	Spin	orb. (low $H_x$ )	orb. (high $H_x$ )	boundary	$H_{c2}$ limiting?
Agterberg et al.	$\mathbf{z}$	$k_x + i\epsilon k_y$	$k_x$	2 <sup>nd</sup> -order	No
Udagawa et al.	$\mathbf{z} - i\alpha\mathbf{y}$	$k_x + i\epsilon k_y$	$k_x$	Crossover	No
New	$\mathbf{z} - i\alpha\mathbf{y}$	$k_x + i\epsilon k_y$ $\ell // z$	$k_z - ik_y$ $\ell // x$	?	Yes? negligible
<sup>3</sup> He-A1	$\mathbf{z} - i\mathbf{y}$	$k_z - ik_y$		2 <sup>nd</sup> -order	No

# Alternative Interpretation

At  $H=0$ ,  $d = x'(k_x + ik_y)$ :  
 The spin component pointing  
 in any direction within the  $ab$  plane.

Yoshioka, Hoshihara, Miyake

d-p model with Upp (O-2p on-site  
 Coulomb) suggests  $d \parallel ab$ .



END