

Higgs coupling measurements at LHC:  
challenges for QCD

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Intro: LHC goals in Higgs physics  
 Overview of channels  
 Accuracy of coupling extraction  
 Requirements for signal and  
 background predictions

Introduction

Higgs search = search for  
 dynamics of  $SU(2) \times U(1)$  breaking

→ discover Higgs boson  
 → measure Higgs couplings

How to identify H as remnant of  
 $SU(2) \times U(1)$  breaking

$$\phi \rightarrow \begin{pmatrix} \chi^+ \\ \frac{v+H+i\chi^0}{\sqrt{2}} \end{pmatrix}$$

A tree level HWW or HZZ coupling  
 is the smoking gun: requires v.e.v.

$$(D_\mu \phi)^\dagger (D^\mu \phi) \rightarrow \frac{g^2}{2} \frac{(v+H)^2}{2} W_\mu^+ W^\mu$$

$$\left(\frac{g v}{2}\right)^2 W_\mu^+ W^\mu \leftrightarrow W \text{ mass}$$

$$\frac{g^2 v}{2} H W_\mu^+ W^\mu \leftrightarrow \text{HWW coupling}$$

Gauge interactions of non-vev scalar

$$\sim \phi^\dagger \phi W, \quad \phi^\dagger \phi WW$$

are bilinear in  $\phi$

- Probe fermion mass generation

$$\lambda \bar{L} \phi \tau_R \rightarrow \frac{m_\tau}{v} (v+H) \bar{\tau}_L \tau_R$$

measure relation between

fermion mass  $\leftrightarrow$   $H\bar{f}f$  coupling

$\Rightarrow$  measure Higgs' Yukawa coupl.

$H\tau\tau$

$Hbb$

$Htt$  etc.

measure couplings to gauge bosons

$HWW$

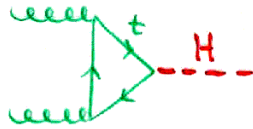
$HZZ$

$H\gamma\gamma$

$Hgg$

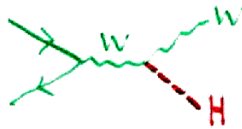
Principal production modes at hadron colliders

gluon fusion



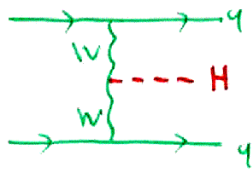
Tevatron LHC

WH/ZH production



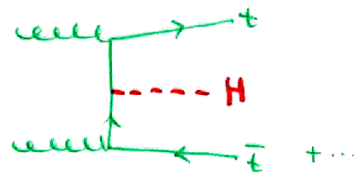
Tevatron

weak boson fusion



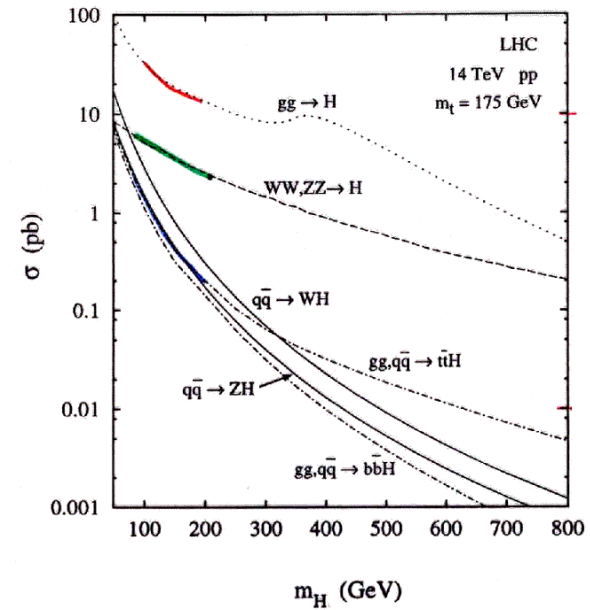
LHC

t $\bar{t}$ H production



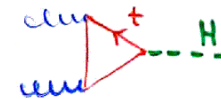
LHC

LHC cross sections



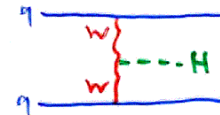
Dominant production processes

gluon fusion



10-30 pb

weak boson fusion



3-5 pb

t $\bar{t}$ H production

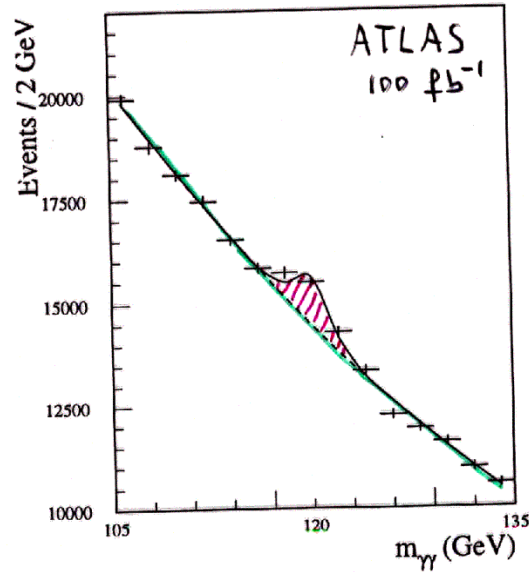
0.2-2 pb

$H \rightarrow \gamma\gamma$

CMS and ATLAS will have excellent photon energy resolution (order 1%)

- Look for narrow  $\gamma\gamma$  invariant mass peak
- Large backgrounds from

$q\bar{q} \rightarrow \gamma\gamma$   
 $gg \rightarrow \gamma\gamma$   
 isolated bremsstr.

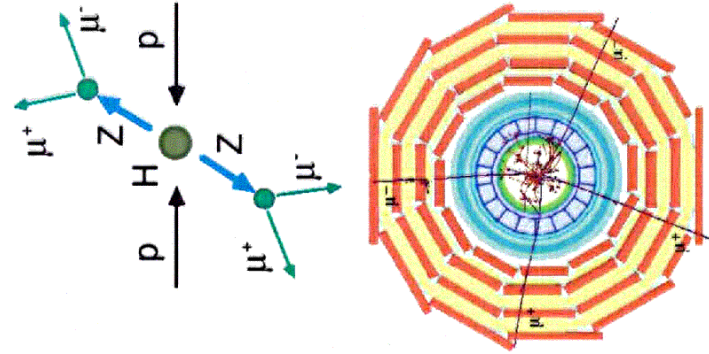


Expected spectrum for 100 fb<sup>-1</sup> of data with the ATLAS detector

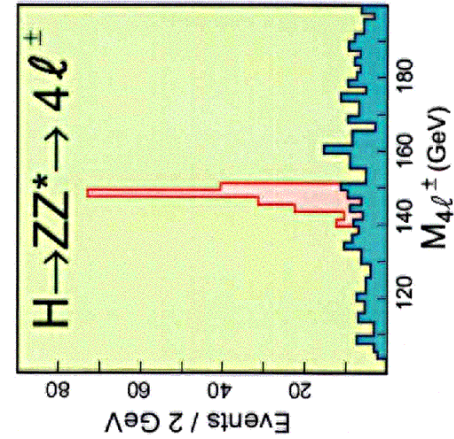
Resolution of CMS is somewhat better

100 fb<sup>-1</sup> expected after ~ 4 years

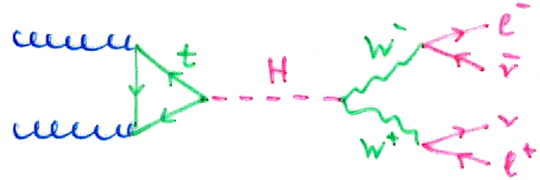
## Intermediate Mass Higgs



- $H \rightarrow ZZ \rightarrow \ell^+ \ell^- \ell^+ \ell^-$  ( $\ell = e, \mu$ )
  - Very clean
  - Resolution: better than 1 GeV
  - Valid for the mass range  $130 < M_H < 500 \text{ GeV}/c^2$

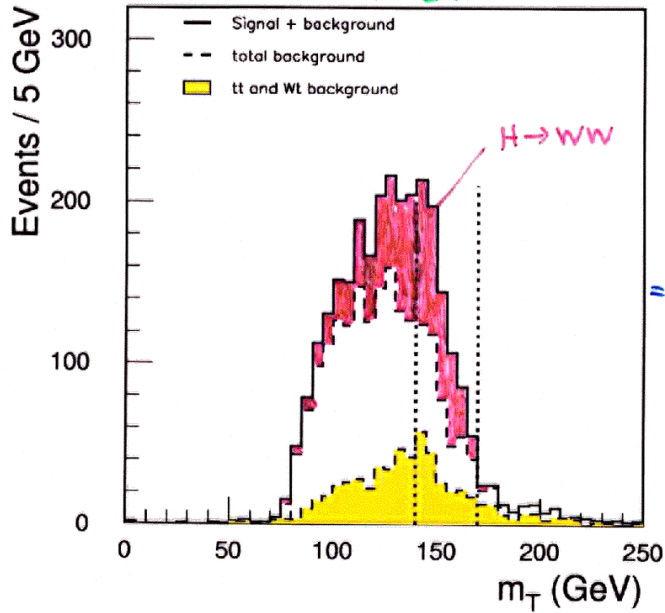


$H \rightarrow WW \rightarrow e^+e^- \nu\bar{\nu}$  (inclusive search)



Exploit  $e^+e^-$  angular correlations [Dittmar Dreiner]

ATLAS TDR



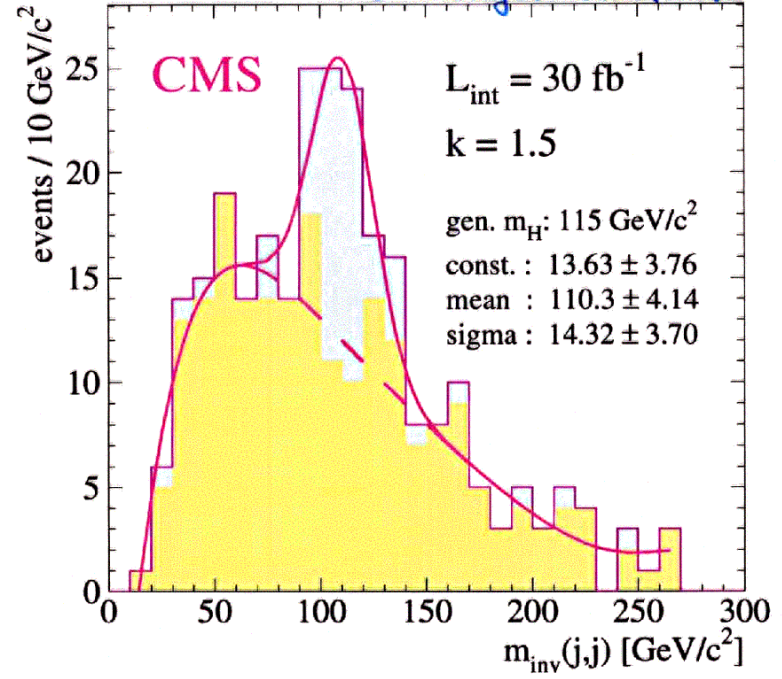
$m_H = 170 \text{ GeV}$

transverse mass  $m_T$   
 $= \sqrt{2p_T^{e^+} p_T^{e^-} (1 - \cos\theta)}$

Background and signal have similar shape  
 $\Rightarrow$  must know bkgd normalization precisely

Important for  $m_H \lesssim 120-130 \text{ GeV}$   
 $gg \rightarrow t\bar{t}H, H \rightarrow b\bar{b}$

Drollinger et al., hep-ph/0111312

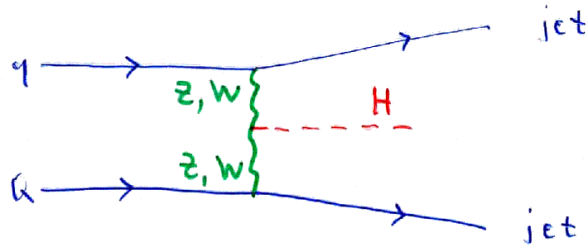


$\gamma_t = t\bar{t}H$  Yukawa coupling  
 $\Rightarrow$  measure

$\gamma_t^2 B(H \rightarrow b\bar{b})$

Weak boson fusion (WBF) has emerged as a powerful tool for

- Higgs search
  - Higgs analysis
- } at LHC



- sizable rate ( $\sim \frac{1}{5}$  of gluon fusion)
- 2 forward tagging jets for efficient background rejection
- color singlet exchange: no central jets
- well known SM cross section: small NLO correction of order 10%.

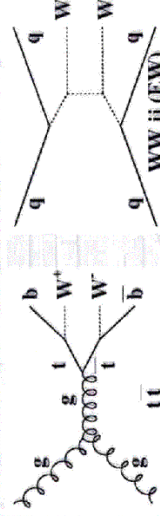
WBF is crucial for Higgs coupling measurements:  $H\tau\tau$  &  $HWW/HZZ$

### Weak Boson Fusion: $H \rightarrow WW$

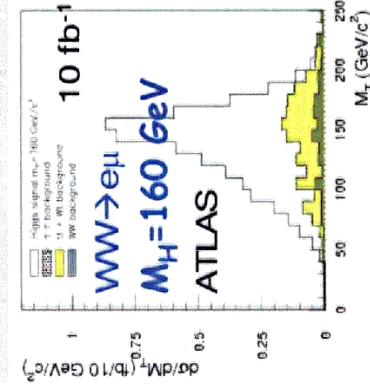
Cross section: 500 to 2000 fb for  $M_H = 120$  to 190 GeV

Dominant backgrounds

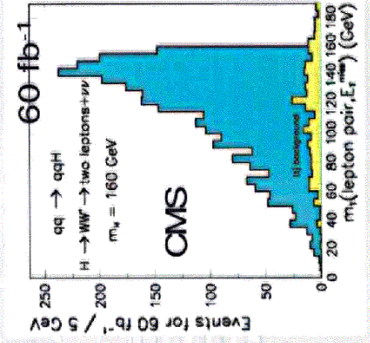
- $t\bar{t}$
- $WWjj$ ,
- $W + 4$  jets



Selection: tag jets with rapidity gap, central jet veto, b-jet veto,  $m_{jj}$ , lepton angles (Spin  $0 \leftrightarrow 1$ ), transverse mass ( $|E_{T,miss}$ )

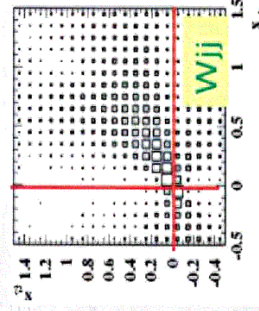
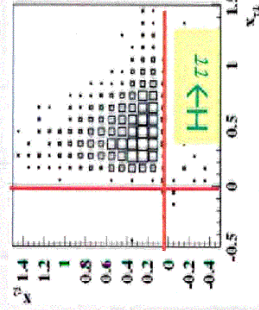
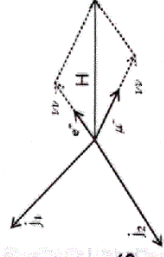


significance > 5  
for 10 fb<sup>-1</sup> and  
 $M_H = 135$  to 190 GeV  
( $WW \rightarrow ll\nu\nu$  and  $lvjj$ ,  
incl.  $\Delta BG = 10\%$ )



## Weak Boson Fusion: $H \rightarrow \tau\tau$

Mass can be reconstructed in collinear approximation  
 $X_\tau$  = momentum fraction carried by tau decay products



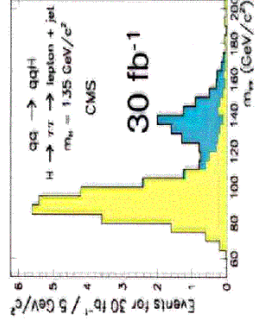
★ significance > 5 for 30 fb<sup>-1</sup> and

$M_H = 110$  to 140 GeV ( $\tau\tau \rightarrow e\mu, \tau\tau \rightarrow \mu\mu, \tau\tau \rightarrow \mu\tau \rightarrow lhad$ )

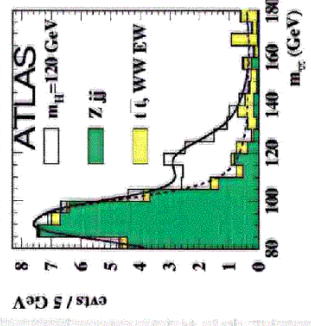
★ background estimate: ~10%

for  $M_H > 125$  GeV from side bands

for  $M_H > 125$  GeV from normalisation of  $Z \rightarrow \tau\tau$  peak



$\sigma_M = 11$  to 12 GeV

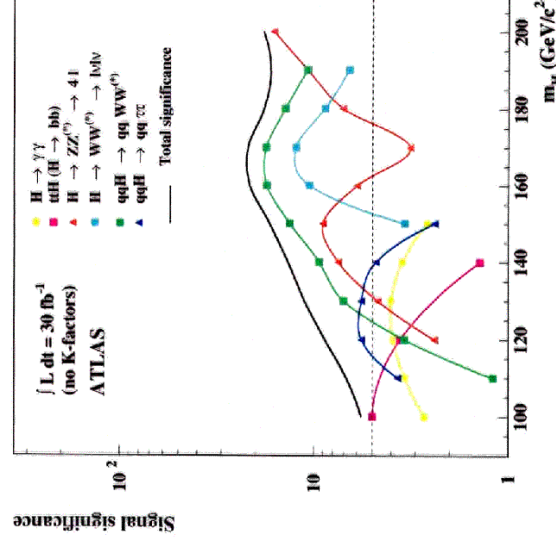
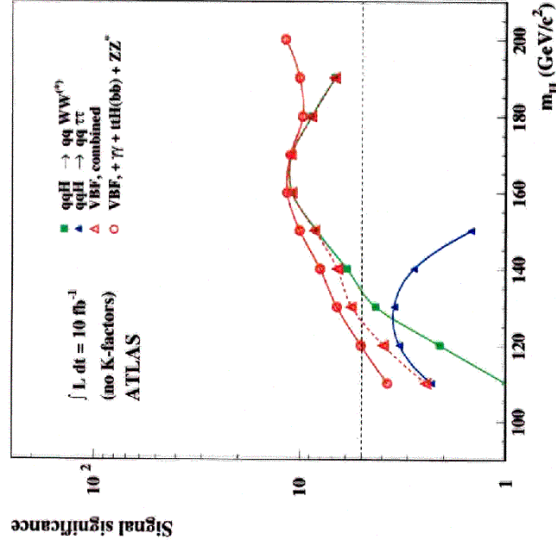


$H \rightarrow \tau\tau \rightarrow e\mu$  30 fb<sup>-1</sup>

Markus Schumacher, Bonn University Higgs Physics at LHC WIN03 Lake Geneva, Wisconsin 12

## Results from VBF Cut Analyses

J. Asai et al. SN-ATLAS-2003-024

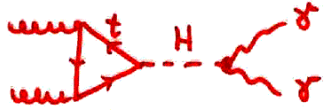


Bruce Mellado, Les Houches 2003, 29/05/03

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Summary of main SM Higgs channels

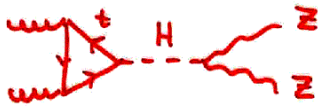
$gg \rightarrow H \rightarrow \gamma\gamma$



$m_H \lesssim 150 \text{ GeV}$

$\sim \Gamma_g \frac{\Gamma_\gamma}{\Gamma} = \gamma_\gamma$

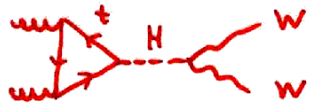
$gg \rightarrow H \rightarrow ZZ \rightarrow 4e^\pm$



$m_H \gtrsim 120 \text{ GeV}$

$\sim \Gamma_g \frac{\Gamma_{ZZ}}{\Gamma} = \gamma_{ZZ}$

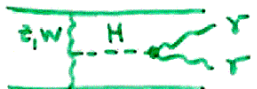
$gg \rightarrow H \rightarrow WW \rightarrow e^\pm e^\mp \mu^\pm$



$m_H \gtrsim 130 \text{ GeV}$

$\sim \Gamma_g \frac{\Gamma_W}{\Gamma} = \gamma_W$

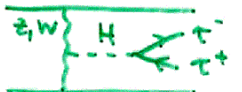
$qq \rightarrow qqH, H \rightarrow \gamma\gamma$



$m_H \lesssim 150 \text{ GeV}$

$\sim \Gamma_W \frac{\Gamma_\gamma}{\Gamma} = X_\gamma$

$qq \rightarrow qqH, H \rightarrow \tau\tau$

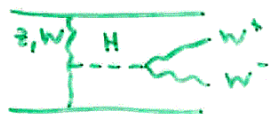


$100 \text{ GeV} \leq m_H < 150 \text{ GeV}$

$\sim \Gamma_W \frac{\Gamma_\tau}{\Gamma} = X_\tau$

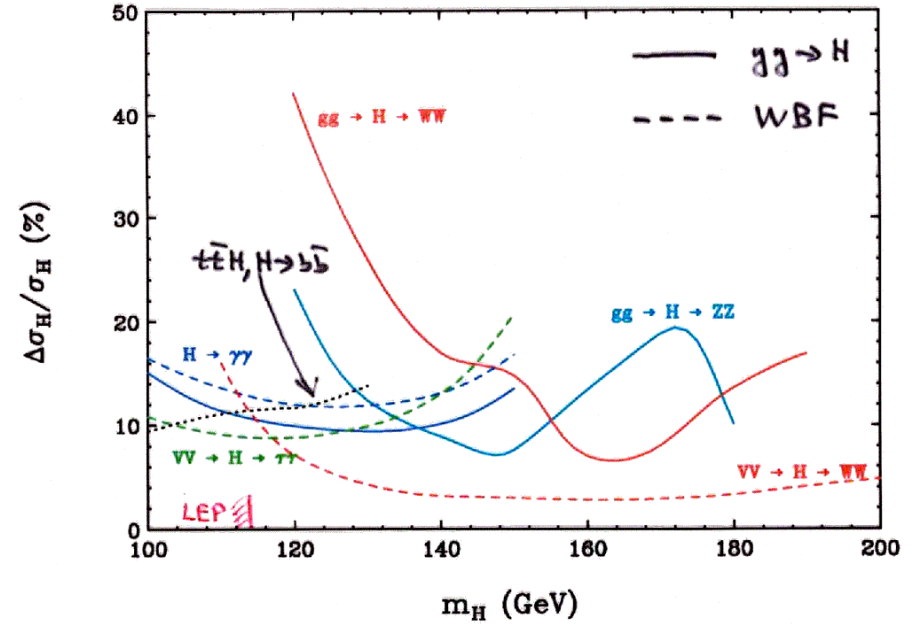
$qq \rightarrow qqH, H \rightarrow WW \rightarrow e^\pm e^\mp \mu^\pm$

$m_H \gtrsim 115 \text{ GeV}$



$\sim \frac{\Gamma_W^2}{\Gamma} = X_W$

Statistical errors with  $200 \text{ fb}^{-1}$ :



Systematic errors:

QCD/pdf uncertainties

$\pm 5\%$  for WBF

$\pm 20\%$  for gluon fusion

Luminosity/acceptance error

$\pm 5\%$

Largely cancel in cross section ratios

$\rightarrow \frac{\Gamma_\tau}{\Gamma_W}, \frac{\Gamma_\gamma}{\Gamma_W}, \frac{\Gamma_{ZZ}}{\Gamma_W}, \frac{\Gamma_\gamma}{\Gamma_W}, \frac{\Gamma_t \Gamma_b}{\Gamma_g \Gamma_\gamma}$



Generic problem for model-independent analysis at LHC:

$$\text{observed} \equiv \hat{\Gamma} = \frac{\Gamma_p \Gamma_d}{\Gamma} = \frac{x \Gamma_p \times \Gamma_d}{x^2 \Gamma}$$

Limits on rescaling factor  $x$

a) total width = sum of partial widths

$$x^2 \Gamma = \sum_i x \Gamma_i \geq \sum_{i \in \{\text{p,d|observed}\}} x \Gamma_i$$

$$\Rightarrow x \geq \sum_{i \in \text{observed}} \frac{\Gamma_i}{\Gamma} = \text{order 1}$$

observation of production puts lower bound on  $x \Gamma_p$ .

b) total width < experimental resolution (or is measured directly!)

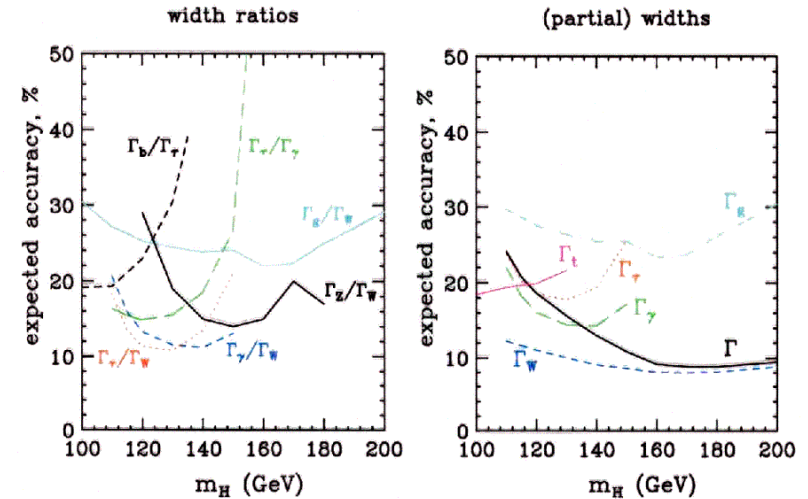
see  $H \rightarrow \gamma\gamma$  or  $H \rightarrow 4\ell$

$$\Rightarrow x \lesssim \sqrt{\frac{1 \text{ GeV}}{\Gamma}}$$

Fit LHC data within constrained models:

- $\frac{g_{H\tau\tau}}{g_{Hbb}} = \text{SM value}$
- $\frac{g_{HWW}}{g_{HZZ}} = \text{SM value}$
- no exotic channels

Assume  $100 \text{ fb}^{-1}$  of data in each of two detectors



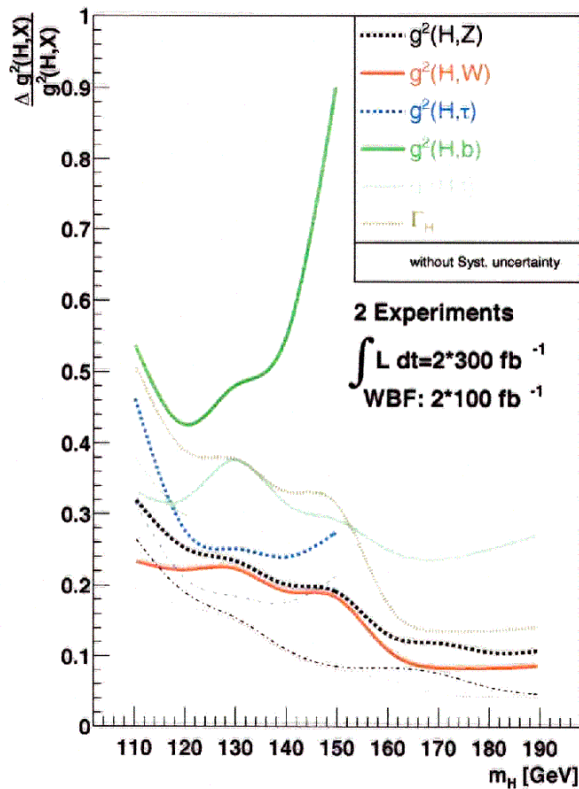
Coupling ratios may differ from SM values in generic models

Example: Large bottom squark corrections to the  $hbb$  vertex in SUSY models

Assume Higgs doublets/singlets only

$$\Rightarrow g_{HWW}^2 = (g_{HWW}^2)^{SM} \sin^2(\alpha - \beta) < (g_{HWW}^2)^{SM}$$

Perform global fit to couplings for expected LHC data



Les Houches 2003

fit allows new particles in  $H \rightarrow \gamma\gamma, H \rightarrow gg$  loops and open exotic channels

## QCD requirements

Signal: would like to have theory errors on production cross section at 5-10% level

$\sigma(gg \rightarrow H)$   $\pm 10-20\%$  @NNLL  $\rightarrow$  Harlander  $\rightarrow$  Grazzini

$\sigma(WBF)$   $\pm 4\%$  @NLO

$\sigma(t\bar{t}H)$   $\pm 10-15\%$  @NLO  $\rightarrow$  Reina

QCD corrections are available for all relevant production processes

- Would like to have NLO MC for  $t\bar{t}H$  including  $t \rightarrow bW$  in narrow width approx.
- interface with parton shower MC
- separation of Hjj events: " $gg \rightarrow Hgg$ " vs. WBF

Backgrounds

(1) Narrow resonances:  $H \rightarrow \gamma\gamma$ ,  $H \rightarrow ZZ \rightarrow 4e$

obtain background experimentally from sideband analysis

(2)  $H \rightarrow \tau\tau$  in WBF

dominant backgrounds:  $(Z \rightarrow \tau\tau) + 2 \text{ jets}$

NLO MC's are available

Measure bkgd in  $Z \rightarrow \mu^+\mu^-$  events

(3)  $H \rightarrow b\bar{b}$  in  $t\bar{t}H$  production

dominant background:  $t\bar{t}b\bar{b}$

need background shape at 10% level

→ sideband analysis with LO MC

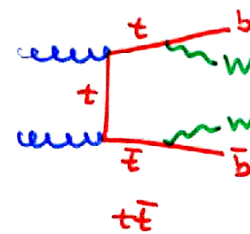
NLO calculation ???

(4)  $H \rightarrow WW \rightarrow e^+\nu e^-\bar{\nu}$

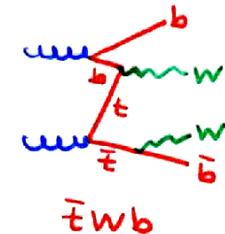
dominant background:  $t\bar{t} \rightarrow W^+W^-b\bar{b}$

(a) inclusive search (Grazzini)

$\sigma(t\bar{t}) < \sigma(tbW)$  after severe jet veto



vs.



LO:  
Kauer

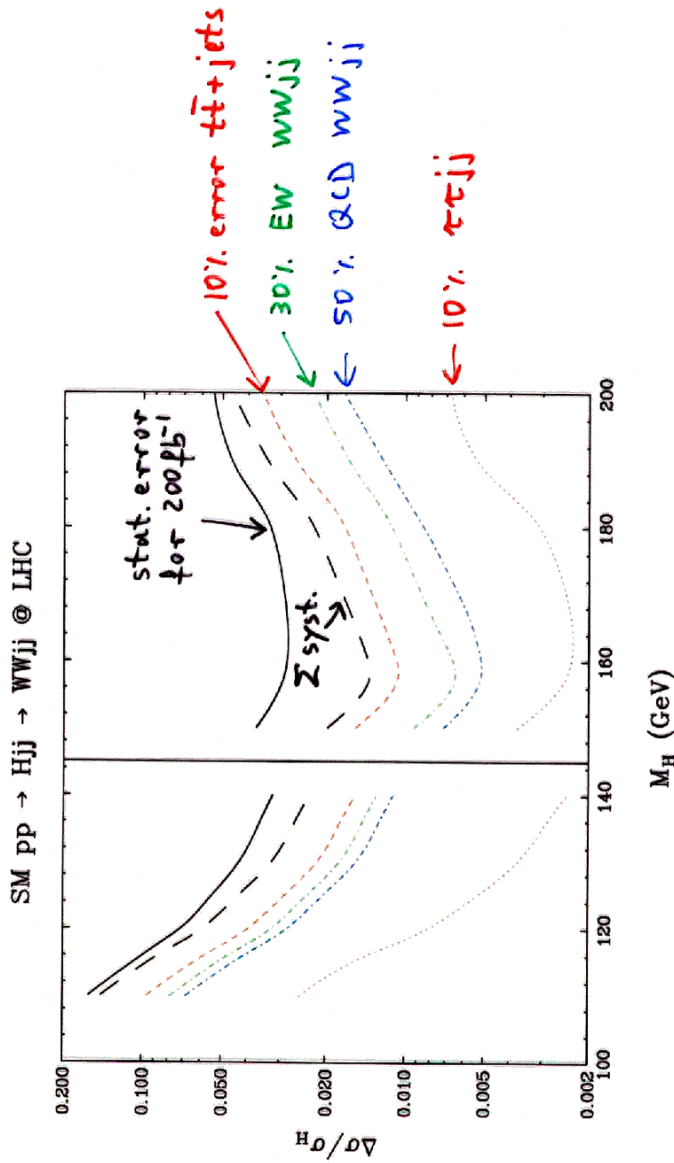
5% normalization uncertainty probably requires NLO calculation including off-shell effects

(b) WBF:  $H+jj \rightarrow WWjj$

need  $t\bar{t}j$  @ NLO

off-shell effects give ~15% correction at LO

Contributions to Higgs measurement errors for  $\sigma_{\text{WBF}}$



Other areas where improvement is needed

- Use of central jet veto for
  - WBF
  - $t\bar{t} \rightarrow WWb\bar{b}$  suppression
- disentangling WBF and  $gg \rightarrow Hgg$

Conclusion:

Higgs discovery is "easy"

Higgs measurements are the true challenge