

# Inhomogeneous phases driven by competing orders

Ilya Vekhter

*Louisiana State U.*

Z. Nussinov and A. V. Balatsky,

*Los Alamos National Lab*

*cond-mat/0409474 and unpublished*



# Message

Inhomogeneous states and glassiness may occur spontaneously in systems *without disorder* as a result of competing interactions or competing orders.

*S. Brazovskii, 1975; J. Schmalian and P. Wolynes, 2000-2004; ZN, IV, AVB 2004.*

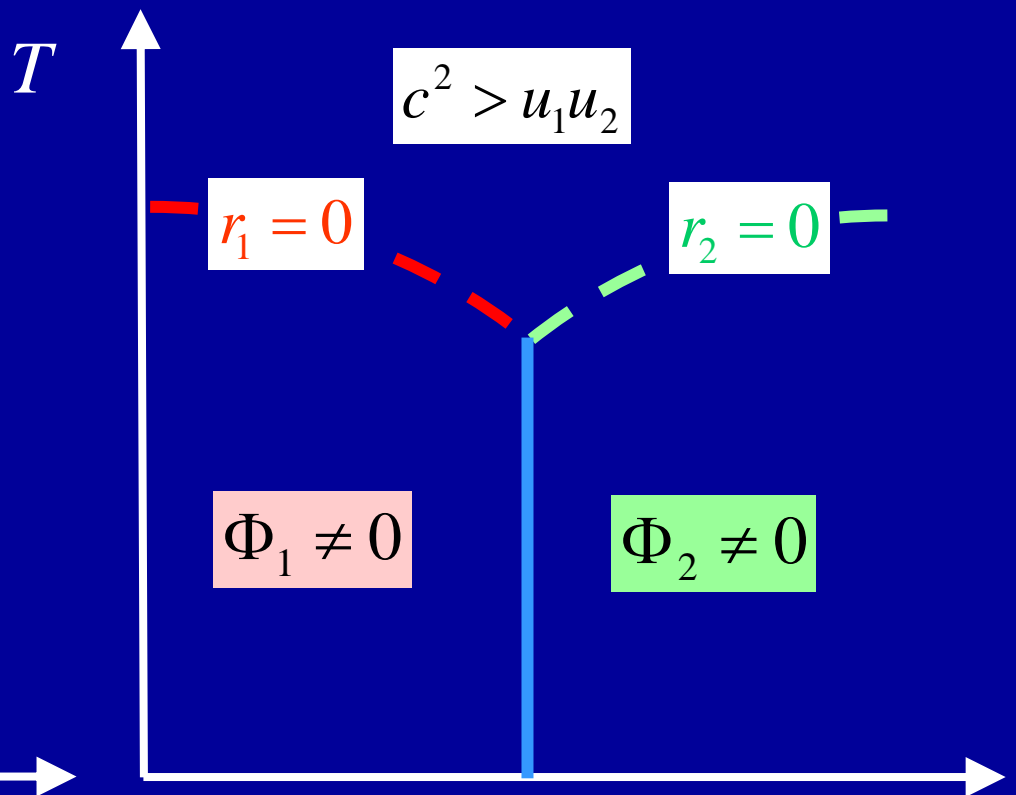
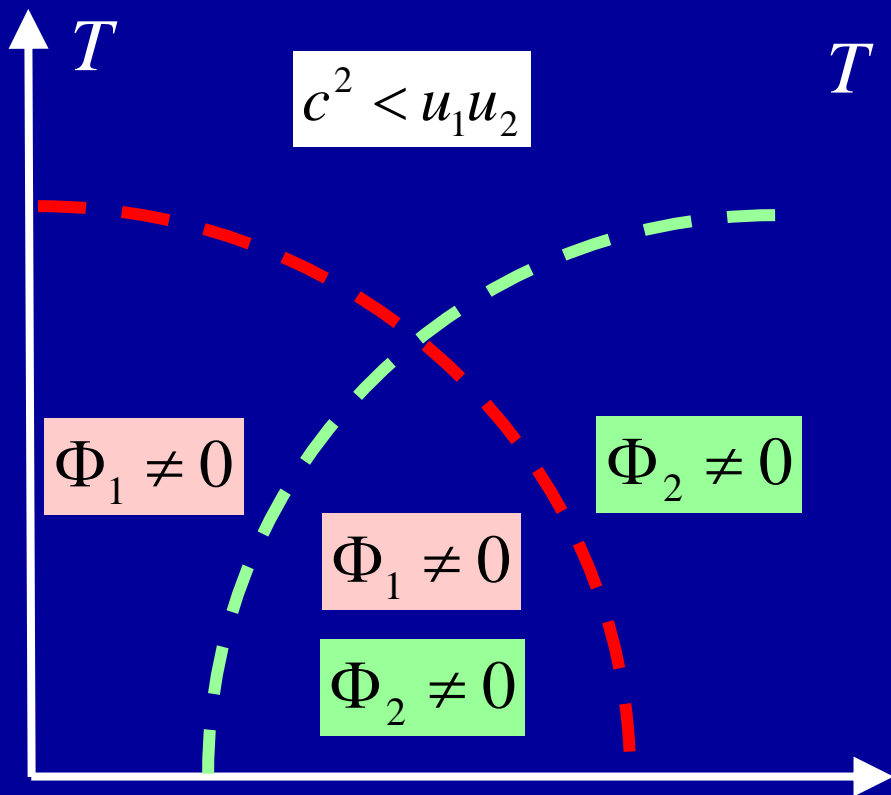
# Inhomogeneous ordered states

- **Coexistence and competition** of different orders:  
*manganites, heavy fermions, cuprates,...*
- **Inhomogeneities** on the micro-meso-nano- scale
- Microscopic origins: **disorder, competing interactions, competing orders**
- Description at the level of **effective theories** (Ginzburg-Landau)?

# Competing orders: a reminder

$$F_0 = \frac{r_1}{2} |\Phi_1|^2 + \frac{u_1}{4} |\Phi_1|^4 + \frac{r_2}{2} |\Phi_2|^2 + \frac{u_2}{4} |\Phi_2|^4 + \frac{c}{2} |\Phi_1|^2 |\Phi_2|^2 + \frac{1}{2} (\nabla \Phi_1)^2 + \frac{1}{2} (\nabla \Phi_2)^2$$

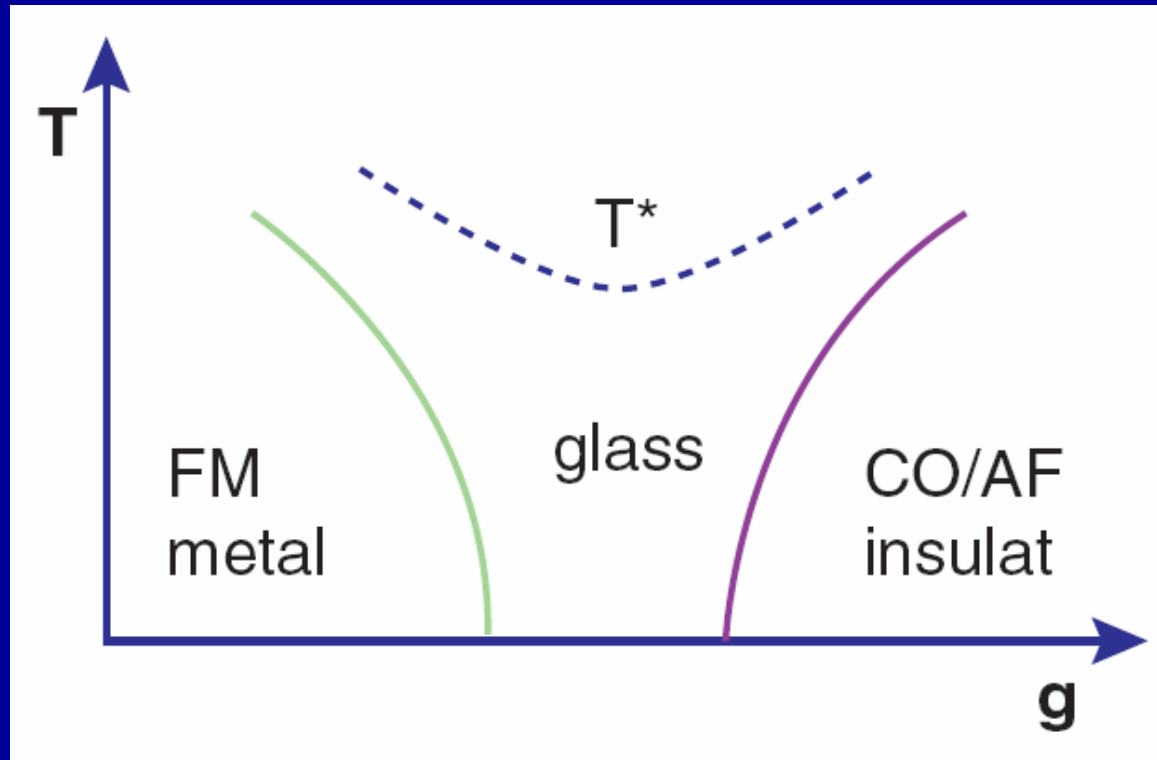
*E. M. Lifshitz (1944), K. Wilson and M. Fisher (1972),  
Liu and Fisher, 1973*



# Disorder-induced glass

Disorder + competing orders

*E. Dagotto et al. 2001-2005;*



Can we obtain inhomogeneous states in a uniform system?

# Gradient couplings

- System has a preferred wave vector:

$$F_1 = a |\Phi_1|^2 |\Phi_2|^2 (\nabla \varphi_2 - q_0), \text{ where } \Phi_2 = |\Phi_2| \exp(i\varphi_2)$$

*manganites, G. Milward, M. Calderon, P. Littlewood, 2005*

- System selects the wave vector from interactions

$$F_1 = a |\Phi_2|^2 |\nabla \Phi_1|^2 + \text{higher}, \quad a < 0 \quad \text{ZN, IV, and AVB, 2004-05}$$

e.g. stripes: charges like to sit at magnetic domain walls

- Inhomogeneity **only in the coexistence region**

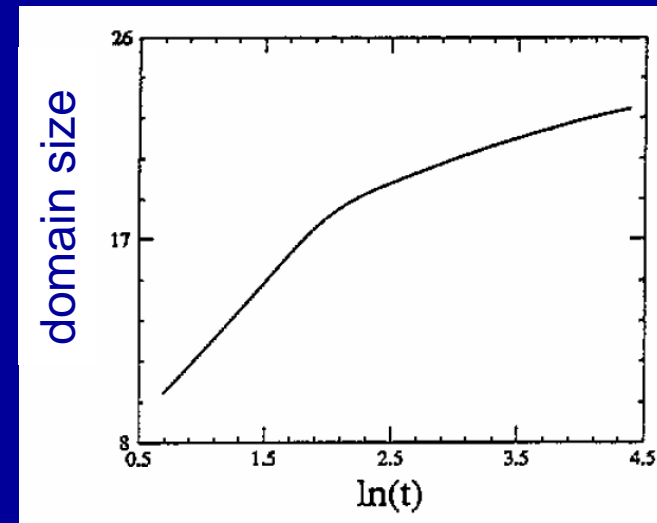
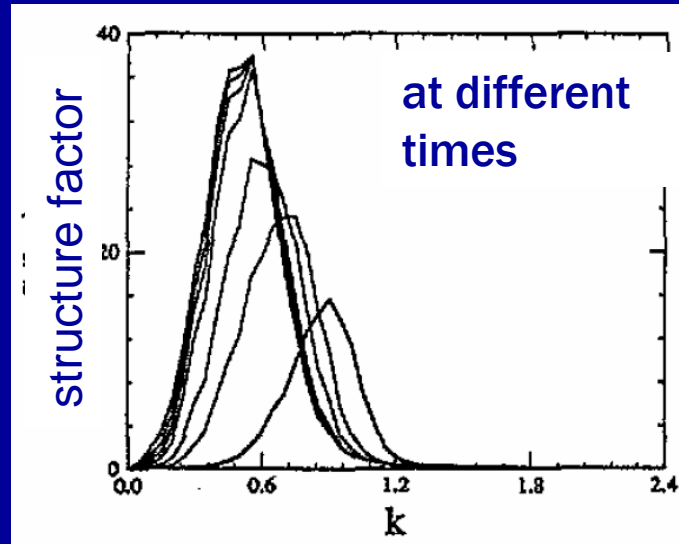
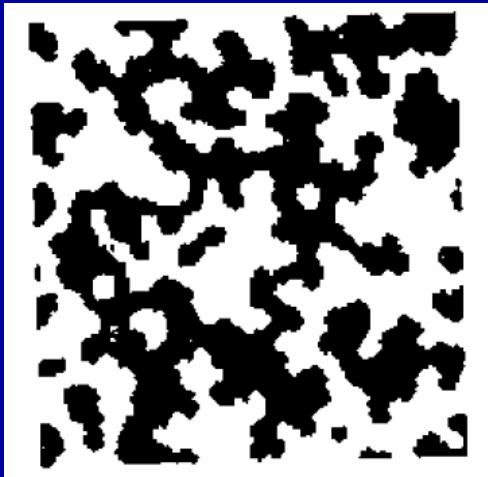
$$F_0 = \frac{r_1}{2} |\Phi_1|^2 + \frac{u_1}{4} |\Phi_1|^4 + \frac{r_2}{2} |\Phi_2|^2 + \frac{1}{4} |\Phi_2|^4 + \frac{c}{2} |\Phi_1|^2 |\Phi_2|^2$$

$$F_1 = \frac{a}{2} |\Phi_2|^2 |\nabla \Phi_1|^2 + \frac{1}{4} |\nabla \Phi_1|^4 + \frac{1}{2} |\nabla \Phi_1|^2 + \frac{1}{2} |\nabla \Phi_2|^2 + \text{higher}$$

- Example: mean field  $T_{c2} > T_{c1}$  :

$$F_{11} = -\frac{|a|}{4} |\Phi_2|^2 q^2 |\Phi_1|^2$$

- **Modulated phase** with weakly T-dependent  $q_0 \propto \Phi_2$
- **Effective model similar to surfactants in liquids**
- **Complicated inhomogeneous states** *M. Laradji et al. 1992*



**Inhomogeneous; Slow dynamics; Glassy?**

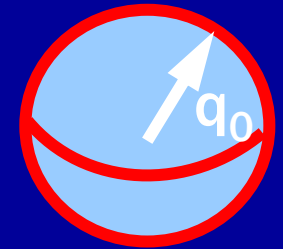
# Effective theory

*S. Brazovskii, 1975*

$$H = \frac{1}{2} \sum_{\mathbf{k}} v(\mathbf{k}) \Phi_1(\mathbf{k}) \Phi_1(-\mathbf{k}) + \frac{u}{4} \sum_{\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3 + \mathbf{k}_4 = 0} \Phi_1(\mathbf{k}_1) \Phi_1(\mathbf{k}_2) \Phi_1(\mathbf{k}_3) \Phi_1(\mathbf{k}_4)$$

$$v(\mathbf{k}) \approx r_0 + (\mathbf{k}^2 - q_0^2)^2, \quad r_0 \equiv \xi_0^{-2} = a(T - T_{c1})$$

- Isotropic model - shell of modes  $|\mathbf{k}| = q_0$ .



- Large phase space for fluctuations:

classical dynamics

or quantum dynamics

$$\langle \Phi^2 \rangle = bTq_0^2 / \sqrt{r_0}$$

$$\langle \Phi^2 \rangle = bq_0^2 \log \frac{E_0}{r_0}$$

- Drives system away from transition



# 1<sup>st</sup> order transition

*Brazovskii transition: S. Brazovskii, 1975*

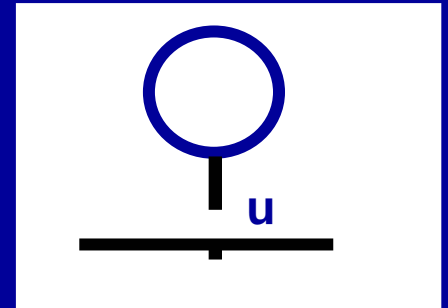
- Self-consistency (large N)

$$r = r_0 + u \langle \Phi^2 \rangle = r_0 + ubTq_0^2 / \sqrt{r}$$

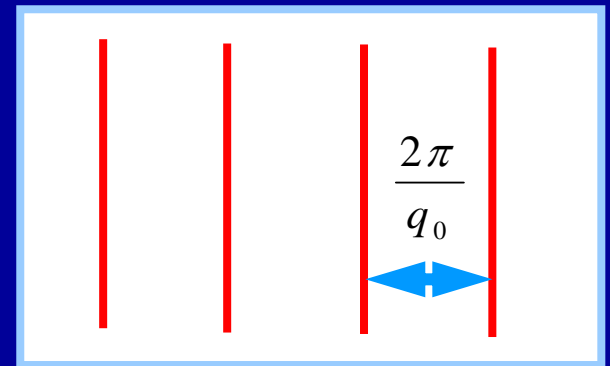
2<sup>nd</sup> order transition impossible (large N) ;  
real transition may occur at  $0 < T_c \ll T_{c1}$

- Fluctuation driven 1<sup>st</sup> order transition
- Mean field: lamellar phase

$$\Phi_1 = A_1 \cos(\mathbf{q}_0 \mathbf{r})$$



$$r^{3/2} \approx ubTq_0^2$$



# Two length scales and glassiness

- Competition between

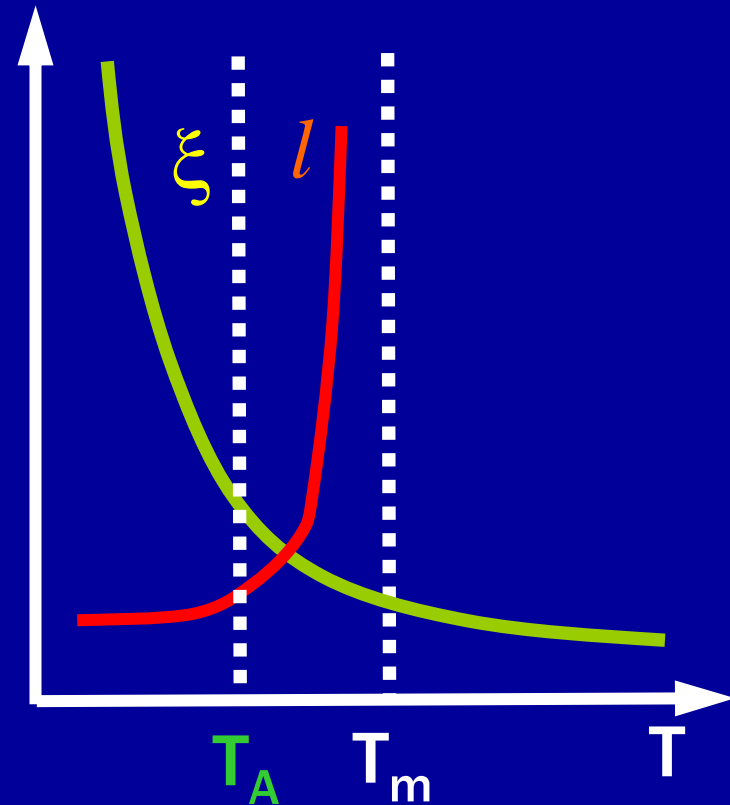
*Z. Nussinov et al., 1999*

- correlation length,  $\xi^{-2} = r(T) + q_0^2$
- modulation length  $l^{-2} = q_0^2 - r(T)$
- at  $q_0^2 = r$  short range correlations
- at  $r + q_0^2 = 0$  long range order

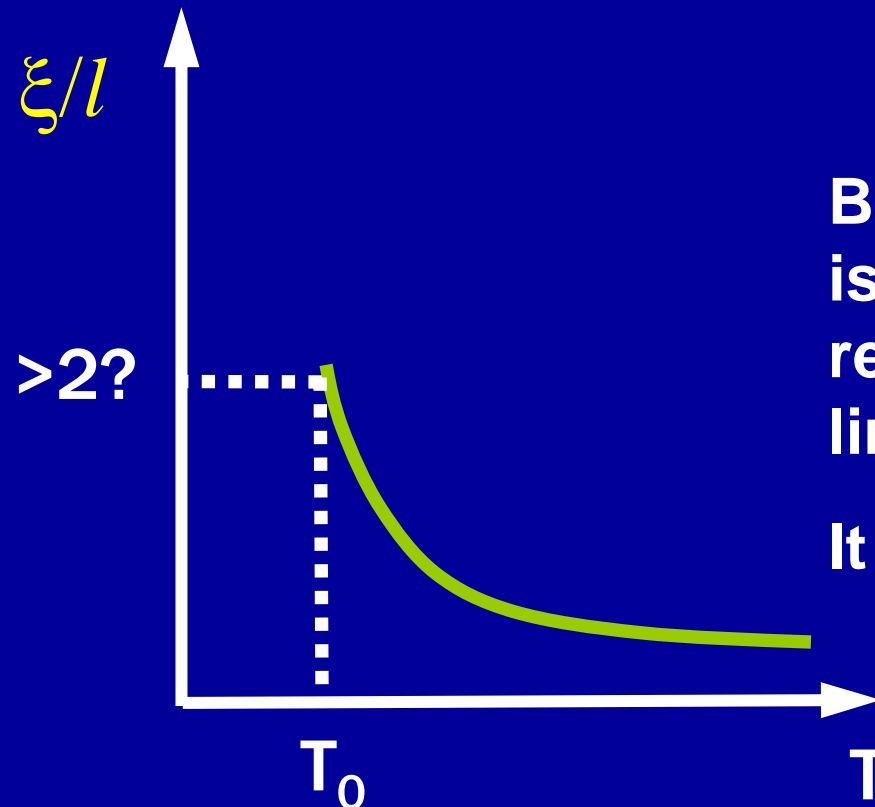
- Glass emerges when  $\xi/l > 2$

*J. Schmalian and P. Wolynes, 2000*

- $N \propto \exp(q_0^3 V)$  metastable states below  $T_A(q_0)$
- Low cost of creating regions of order parameter ( $\xi^{-2}$ ) correlated over short distance of order  $l$



- If  $T_{c2} \gg T_{c1}$  : fluctuation-induced first order transition into inhomogeneous or glassy phase at  $T_{c2} > T_1 > T_{c1}$  possibly followed by a strongly first order transition to a uniform phase at  $T^* \ll T_{c1}$ .
- For mean field  $T_{c2} \approx T_{c1}$ : 1<sup>st</sup> order into modulated phase
- Reason:  $q_0$  depends on  $T$
- Under investigation



Basic question:  
is it enough to  
reach glassy  
limit?

It depends...

# Summary

- You can get modulated and inhomogeneous states of different orders in nominally uniform systems. *Disorder need not be there.*
- These states may exhibit glassy dynamics.
- For glassiness details matter (unfortunately?)

## Open question:

- Hamiltonian from which such GL follows?
- Plan: use mean field microscopic theories (see talks by G. Alvarez and W. Atkinson), and check the relevant Ginzburg-Landau theories.