

Entanglement Entropy of Critical Spin Liquids

Tarun Grover

with Frank Zhang, Ashvin Vishwanath (UC Berkeley)

Our aim: To calculate Renyi Entropy for **critical spin liquids**.

Why bother?

Demotivators:

- Gapless spin-liquids correspond to **strongly interacting gauge-matter theories** \Rightarrow very difficult to analyze.
- Entanglement entropy is a highly **non-local operator** and hence difficult to calculate.
- Experimental relevance?

Why bother?

Demotivators:

- Gapless spin-liquids correspond to **strongly interacting gauge-matter theories** \Rightarrow very difficult to analyze.
- Entanglement entropy is a highly **non-local operator** and hence difficult to calculate.
- Experimental relevance?

Why bother?

Demotivators:

- Gapless spin-liquids correspond to **strongly interacting gauge-matter theories** \Rightarrow very difficult to analyze.
- Entanglement entropy is a highly **non-local operator** and hence difficult to calculate.
- Experimental relevance?

Why bother?

Motivators:

- We have candidate **variational wave-functions** for ground state of gapless spin-liquids that may provide *non-perturbative* access to interesting properties.
- Precisely because entanglement entropy is a highly non-local operator and hence may capture non-local physics of 'quantum order'.
- Experiments \Rightarrow we study candidate spin-liquid wavefunctions for triangular lattice organic spin-liquids. Conjectured to have emergent Fermi surface.

Why bother?

Motivators:

- We have candidate **variational wave-functions** for ground state of gapless spin-liquids that may provide *non-perturbative* access to interesting properties.
- Precisely because entanglement entropy is a highly non-local operator and hence may capture non-local physics of 'quantum order'.
- Experiments \Rightarrow we study candidate spin-liquid wavefunctions for triangular lattice organic spin-liquids. Conjectured to have emergent Fermi surface.

Why bother?

Motivators:

- We have candidate **variational wave-functions** for ground state of gapless spin-liquids that may provide *non-perturbative* access to interesting properties.
- Precisely because entanglement entropy is a highly non-local operator and hence may capture non-local physics of 'quantum order'.
- Experiments \Rightarrow we study candidate spin-liquid wavefunctions for triangular lattice organic spin-liquids. Conjectured to have **emergent Fermi surface**.

- 1 Introduction to Critical Spin-liquids
 - Slave-particle construction
 - Experimental motivation for critical spin-liquids
- 2 Entanglement and Renyi entropy
- 3 Renyi entropy of critical spin-liquids
 - Benchmarking
 - Results for Critical spin-liquids
- 4 Discussion
 - Ongoing work

Oshikawa-Hastings argument:

Mott insulators with odd number of electrons per unit cell that do not break any symmetry have low lying excitations → **topological degeneracy** OR **critical phase**.

Slave particles and deconfinement



$$H = J \sum_{\langle ij \rangle} S_i \cdot S_j + \dots \quad (1)$$



$$\vec{S} = \frac{f^\dagger \vec{\sigma} f}{2} \quad (2)$$

with the constraint $f^\dagger f = 1$.

- $U(1)$ redundancy \Rightarrow gauge fields coupled to spinons f .

Slave particles and deconfinement



$$H = J \sum_{\langle ij \rangle} S_i \cdot S_j + \dots \quad (1)$$



$$\vec{S} = \frac{f^\dagger \vec{\sigma} f}{2} \quad (2)$$

with the constraint $f^\dagger f = 1$.

- $U(1)$ redundancy \Rightarrow gauge fields coupled to spinons f .

Slave particles and deconfinement



$$H = J \sum_{\langle ij \rangle} S_i \cdot S_j + \dots \quad (1)$$



$$\vec{S} = \frac{f^\dagger \vec{\sigma} f}{2} \quad (2)$$

with the constraint $f^\dagger f = 1$.

- $U(1)$ redundancy \Rightarrow gauge fields coupled to spinons f .

Slave particle matter-gauge theory

- Full gauge-matter action

$$S = \sum_{\langle xx' \rangle} (\bar{f}_{\vec{x}\sigma} f_{\vec{x}'\sigma} e^{ia_{xx'}} + c.c.) + \sum_{\square} \cos(\vec{\nabla} \times \vec{a}) \quad (3)$$

- If deconfinement \Rightarrow Non-compact gauge field

$$S = \int_{k,\omega} \bar{f}_{k\sigma} (-i\omega + \epsilon_k + \mu) f_{k\sigma} + k^2/e^2 |a(k,\omega)|^2 + j(k,\omega) a(k,\omega)$$

- Theory difficult to analyze directly for $D = 2 + 1$ and $SU(2)$ symmetry. Lee, Nagaosa, Halperin, Read, Senthil, Mross, Sung-sik Lee, Reizer ...

Slave particle matter-gauge theory

- Full gauge-matter action

$$S = \sum_{\langle xx' \rangle} (\bar{f}_{\vec{x}\sigma} f_{\vec{x}'\sigma} e^{ia_{xx'}} + c.c.) + \sum_{\square} \cos(\vec{\nabla} \times \vec{a}) \quad (3)$$

- If deconfinement \Rightarrow Non-compact gauge field

$$S = \int_{k,\omega} \bar{f}_{k\sigma} (-i\omega + \epsilon_k + \mu) f_{k\sigma} + k^2/e^2 |a(k,\omega)|^2 + j(k,\omega) a(k,\omega)$$

- Theory difficult to analyze directly for $D = 2 + 1$ and $SU(2)$ symmetry. Lee, Nagaosa, Halperin, Read, Senthil, Mross, Sung-sik Lee, Reizer ...

Slave particle matter-gauge theory

- Full gauge-matter action

$$S = \sum_{\langle xx' \rangle} (\bar{f}_{\vec{x}\sigma} f_{\vec{x}'\sigma} e^{i a_{xx'}} + c.c.) + \sum_{\square} \cos(\vec{\nabla} \times \vec{a}) \quad (3)$$

- If deconfinement \Rightarrow Non-compact gauge field

$$S = \int_{k,\omega} \bar{F}_{k\sigma} (-i\omega + \epsilon_k + \mu) f_{k\sigma} + k^2/e^2 |a(k,\omega)|^2 + j(k,\omega) a(k,\omega)$$

- Theory difficult to analyze directly for $D = 2 + 1$ and $SU(2)$ symmetry. Lee, Nagaosa, Halperin, Read, Senthil, Mross, Sung-sik Lee, Reizer ...

From slave particles to projected wavefunctions

- Consider the 'mean-field' Hamiltonian in the weak-coupling (\equiv deconfined) limit:

$$H = \sum_{\langle ij \rangle} f_i^\dagger f_j + h.c. \quad (4)$$

with $f^\dagger f = 1$.

- Ground state wavefunction

$$|\Psi\rangle_{var} \equiv |\Psi\rangle_{PFL} = \prod_i (1 - n_{i\uparrow} n_{i\downarrow}) \left(\prod_{\vec{k}\sigma} f_{\vec{k}\sigma}^\dagger \right) |0\rangle \quad (5)$$

From slave particles to projected wavefunctions

- Consider the 'mean-field' Hamiltonian in the weak-coupling (\equiv deconfined) limit:

$$H = \sum_{\langle ij \rangle} f_i^\dagger f_j + h.c. \quad (4)$$

with $f^\dagger f = 1$.

- Ground state wavefunction

$$|\Psi\rangle_{var} \equiv |\Psi\rangle_{PFL} = \prod_i (1 - n_{i\uparrow} n_{i\downarrow}) \left(\prod_{\vec{k}\sigma} f_{\vec{k}\sigma}^\dagger \right) |0\rangle \quad (5)$$

Physics of projected wavefunctions?

- Do projected wave-functions have the same symmetry as the unprojected ones?
- Do they have a well-defined Fermi surface?
- How to tell?

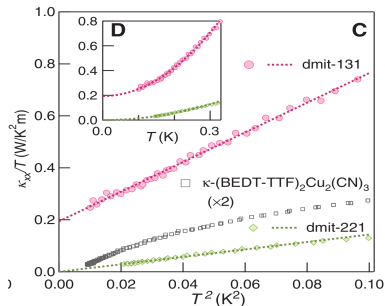
Physics of projected wavefunctions?

- Do projected wave-functions have the same symmetry as the unprojected ones?
- Do they have a well-defined Fermi surface?
- How to tell?

Physics of projected wavefunctions?

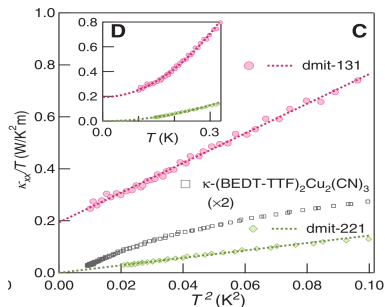
- Do projected wave-functions have the same symmetry as the unprojected ones?
- Do they have a well-defined Fermi surface?
- **How to tell?**

Insulator with metallic thermal transport and specific heat



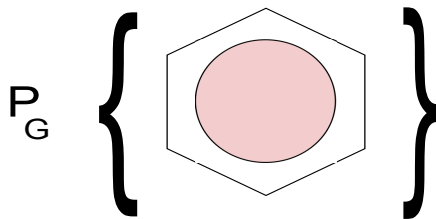
- Material: $\text{EtMe}_3\text{Sb}[\text{Pd}(\text{dmit})_2]_2$ Yamashita et al, Science 2010.
- κ/T ($T \rightarrow 0$) extrapolates to non-zero value as $T \rightarrow 0$.

Insulator with metallic thermal transport and specific heat



- Material: $\text{EtMe}_3\text{Sb}[\text{Pd}(\text{dmit})_2]_2$ Yamashita et al, Science 2010.
- κ/T ($T \rightarrow 0$) extrapolates to non-zero value as $T \rightarrow 0$.

Projected wave-functions and triangular lattice spin-liquids



where $P_G = \prod_i (1 - n_{i\uparrow} n_{i\downarrow})$ Motrunich 2005

Do projected wave-functions have the correct **entanglement properties** to faithfully capture physics of critical spin-liquids?

Basic definitions

- Wavefunction, $|\phi\rangle$: trace out B to get a density matrix on A:
 $\rho_A = \text{Tr}_B |\phi\rangle\langle\phi|$.
- Renyi entropies:

$$S_n = \frac{1}{1-n} \log(\text{Tr} \rho_A^n) \quad (6)$$

- $S_1 = S_{vN} = \text{Tr} \rho_A \log \rho_A$
- This talk: S_2

Basic definitions

- Wavefunction, $|\phi\rangle$: trace out B to get a density matrix on A:
 $\rho_A = \text{Tr}_B |\phi\rangle\langle\phi|$.
- Renyi entropies:

$$S_n = \frac{1}{1-n} \log(\text{Tr} \rho_A^n) \quad (6)$$

- $S_1 = S_{vN} = -\text{Tr} \rho_A \log \rho_A$
- This talk: S_2

Basic definitions

- Wavefunction, $|\phi\rangle$: trace out B to get a density matrix on A:
 $\rho_A = \text{Tr}_B |\phi\rangle\langle\phi|$.
- Renyi entropies:

$$S_n = \frac{1}{1-n} \log(\text{Tr} \rho_A^n) \quad (6)$$

- $S_1 = S_{vN} = \text{Tr} \rho_A \log \rho_A$
- This talk: S_2

Basic definitions

- Wavefunction, $|\phi\rangle$: trace out B to get a density matrix on A:
 $\rho_A = \text{Tr}_B |\phi\rangle\langle\phi|$.
- Renyi entropies:

$$S_n = \frac{1}{1-n} \log(\text{Tr} \rho_A^n) \quad (6)$$

- $S_1 = S_{vN} = -\text{Tr} \rho_A \log \rho_A$
- This talk: S_2

Usefulness of S_2

- Numerically easier to calculate than S_{vN} .
 - $S \geq S_2$
 - S_2 and S share:
 - Violation of area law for CFT in 1D and for free fermions in any D.
 - Non-zero topological entanglement/Renyi entropy for gapped 2D topologically ordered phases.
- Wilczek, Holzhey, Larsen, Cardy, Calabrese, Casini, Huerta, Swingle, Flammia, Hamma, Hughes, Wen ...

Usefulness of S_2

- Numerically easier to calculate than S_{vN} .
 - $S \geq S_2$
 - S_2 and S share:
 - Violation of area law for CFT in 1D and for free fermions in any D.
 - Non-zero topological entanglement/Renyi entropy for gapped 2D topologically ordered phases.
- Wilczek, Holzhey, Larsen, Cardy, Calabrese, Casini, Huerta, Swingle, Flammia, Hamma, Hughes, Wen ...

Usefulness of S_2

- Numerically easier to calculate than S_{vN} .
 - $S \geq S_2$
 - S_2 and S share:
 - Violation of area law for CFT in 1D and for free fermions in any D.
 - Non-zero topological entanglement/Renyi entropy for gapped 2D topologically ordered phases.
- Wilczek, Holzhey, Larsen, Cardy, Calabrese, Casini, Huerta, Swingle, Flammia, Hamma, Hughes, Wen ...

Entanglement entropy as a non-local probe

- For topological ordered phases,

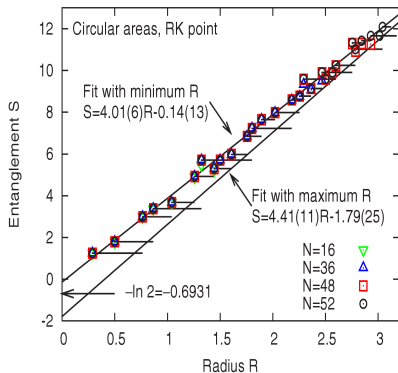
$$S_n = \alpha L - \log(D) \quad (7)$$

where

$$D = \sqrt{\sum_i d_i^2} \quad (8)$$

d_i = quantum dimension of i 'th quasiparticle.

Example: Entanglement entropy as a non-local probe



Furukawa, Misguich, More recent work: Roger Melko et al.
(Unpublished)

Calculation of Renyi entropy in Monte Carlo

$$\begin{aligned} e^{-S_2} &= \text{tr} \rho_A^2 = \sum_{a,a'} \rho_A(a, a') \rho_A(a', a) \\ &= \sum_{a,a',b,b'} \phi^*(a, b) \phi(a', b) \phi^*(a', b') \phi(a, b') \end{aligned}$$

Calculation of Renyi entropy in Monte Carlo

\Rightarrow

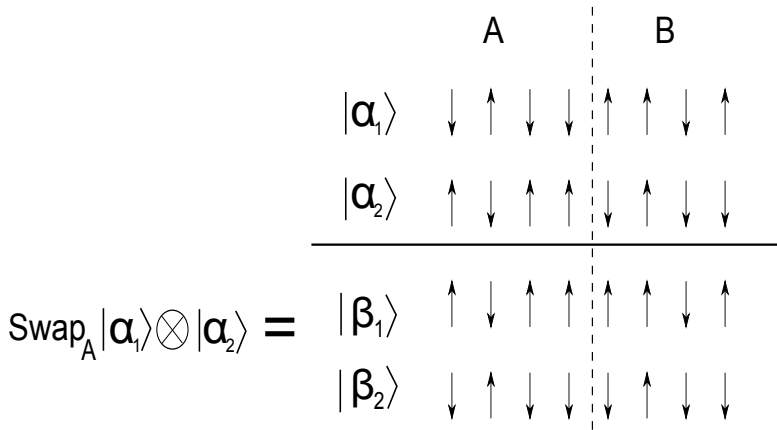
$$e^{-S_2} = \frac{\langle \Phi | \text{Swap}_A | \Phi \rangle}{\langle \Phi | \Phi \rangle} \quad (9)$$

where

$$\text{Swap}_A |a, b\rangle |a', b'\rangle = |a', b\rangle |a, b'\rangle$$

$$|\Phi\rangle = |\phi\rangle \otimes |\phi\rangle$$

Swap operator

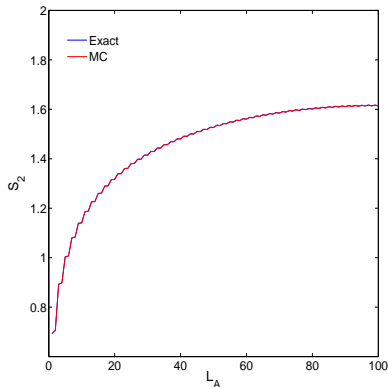


Calculation of Renyi entropy in Monte Carlo

$$\begin{aligned}\langle \text{Swap}_A \rangle &= \frac{\sum_{\alpha_1, \alpha_2} \langle \phi | \alpha_1 \rangle \langle \beta_1 | \phi \rangle \langle \phi | \alpha_2 \rangle \langle \beta_2 | \phi \rangle}{|\langle \phi | \phi \rangle|^2} \\ &= \sum_{\alpha_1 \alpha_2} \frac{|\langle \phi | \alpha_1 \rangle|^2 |\langle \phi | \alpha_2 \rangle|^2}{\langle \phi | \phi \rangle \langle \phi | \phi \rangle} \cdot \frac{\langle \beta_1 | \phi \rangle \langle \beta_2 | \phi \rangle}{\langle \alpha_1 | \phi \rangle \langle \alpha_2 | \phi \rangle} \\ &= \sum_{\alpha_1 \alpha_2} \rho_{\alpha_1} \rho_{\alpha_2} f(\alpha_1, \alpha_2)\end{aligned}$$

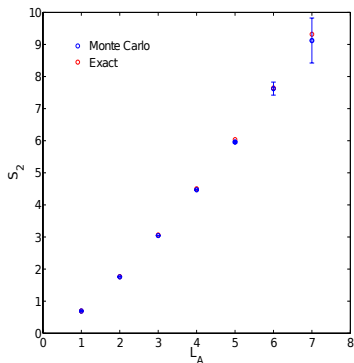
Hastings et al (2010)

Free fermions in 1D



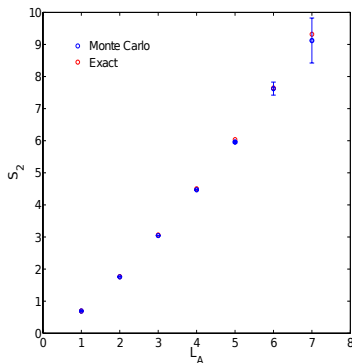
Almost exact answer!

Free fermions in 2D



- $L_{total} = 18, L_A \leq 7$.
- Very good agreement, can detect area-law violation!

Free fermions in 2D



- $L_{total} = 18, L_A \leq 7$.
- Very good agreement, **can detect area-law violation!**

Projected Fermi sea

- Does projection retain the **Fermi surface**?
- Does projected state has the same **symmetries**?

Projected Fermi sea

- Does projection retain the **Fermi surface**?
- Does projected state has the same **symmetries**?

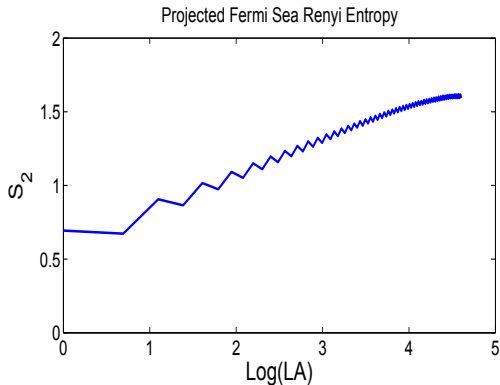
Warm-up: Projected Fermi Sea in One dimension

- Before projection: total **central charge** = 2.
- After projection: exact ground state of Haldane-Shastry model
Luttinger liquid with **central charge** = 1.

Warm-up: Projected Fermi Sea in One dimension

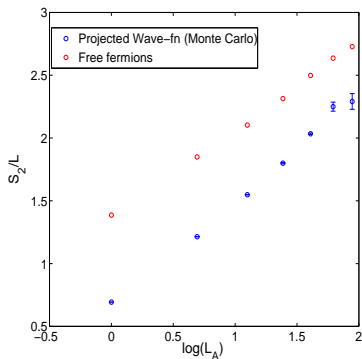
- Before projection: total **central charge** = 2.
- After projection: exact ground state of Haldane-Shastry model
Luttinger liquid with **central charge** = 1.

Projected Fermi Sea in One dimension



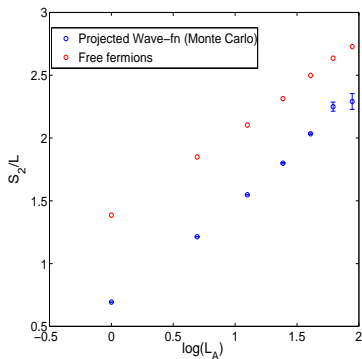
Consistent with $c = 1$.

Projected Fermi sea on triangular lattice



- $L_{total} = 18, L_A \leq 7$.
- *Area-law violation in a bosonic wavefunction!* \Rightarrow fits $L_A \log L_A$ scaling.

Projected Fermi sea on triangular lattice



- $L_{total} = 18, L_A \leq 7$.
- *Area-law violation in a bosonic wavefunction!* \Rightarrow fits $L_A \log L_A$ scaling.

Projected Fermi sea on triangular lattice

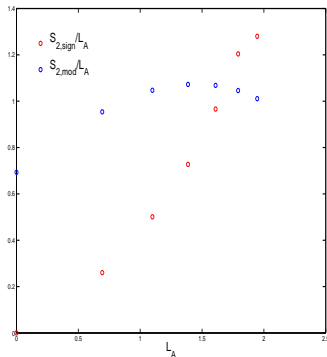
- Can write $S_2 = S_{2,sign} + S_{2,mod}$ where

$$e^{-S_{2,mod}} = \sum_{\alpha_1 \alpha_2} \rho_{\alpha_1} \rho_{\alpha_2} |f(\alpha_1, \alpha_2)| \quad (10)$$

and

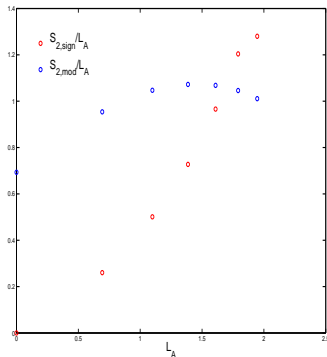
$$S_{2,sign} = S_2 - S_{2,mod} \quad (11)$$

Projected triangular: $S_{2,sign}$ Vs $S_{2,mod}$



- $S_{2,sign}$ violates area-law, $S_{2,mod}$ doesn't.
- Conjecture: holds for fermi surfaces in general.

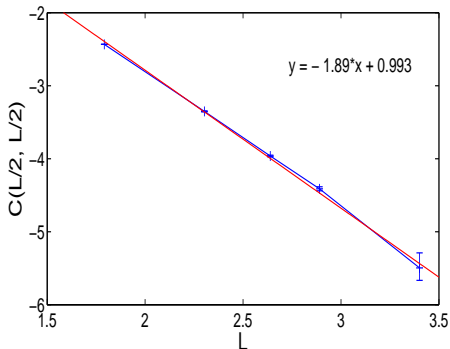
Projected triangular: $S_{2,sign}$ Vs $S_{2,mod}$



- $S_{2,sign}$ violates area-law, $S_{2,mod}$ doesn't.
- **Conjecture:** holds for fermi surfaces in general.

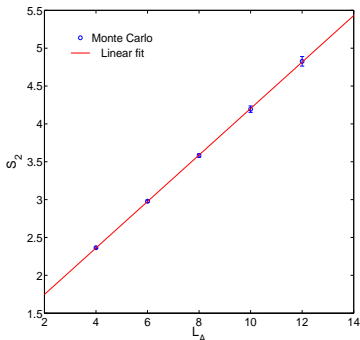
Projected Fermi sea on π -flux square lattice

Spin-spin correlations:



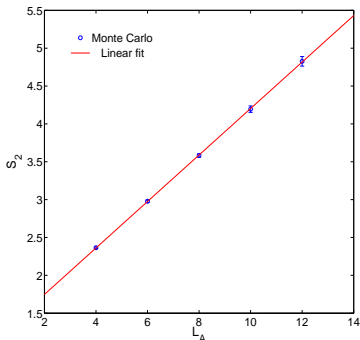
⇒ Algebraically decaying correlations!

Projected Fermi sea on π -flux square lattice



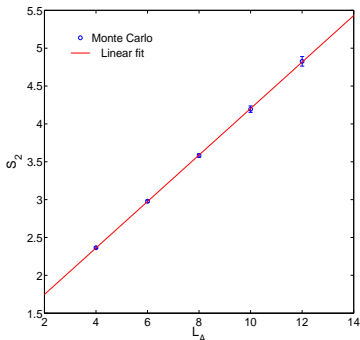
- Near perfect area law \Rightarrow consistent with 'nodal Dirac liquid'.
- Can extract 'universal constant'.
- Prediction for QED-3(?).

Projected Fermi sea on π -flux square lattice



- Near perfect area law \Rightarrow consistent with 'nodal Dirac liquid'.
- Can extract 'universal constant'.
- Prediction for QED-3(?).

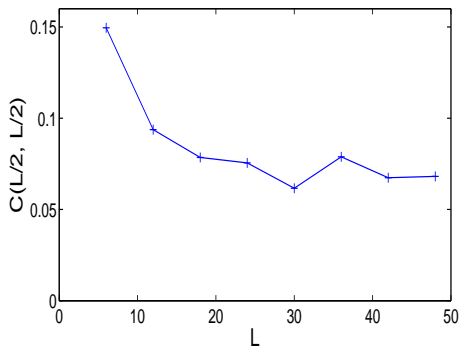
Projected Fermi sea on π -flux square lattice



- Near perfect area law \Rightarrow consistent with 'nodal Dirac liquid'.
- Can extract 'universal constant'.
- Prediction for QED-3(?).

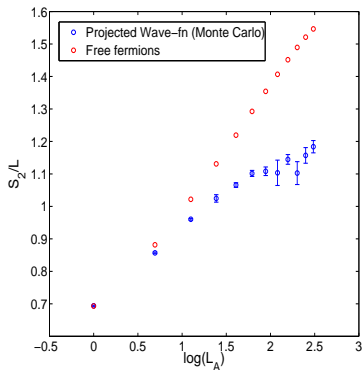
Projected Fermi sea on square lattice

Spin-spin correlations:



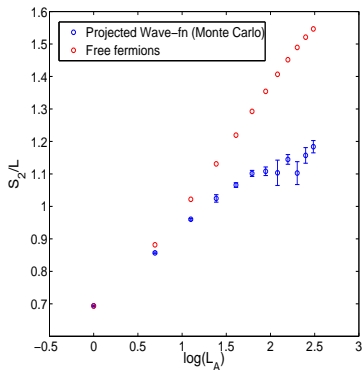
⇒ Magnetically ordered!

Projected Fermi sea on Square lattice



- $L_A \log L_A$ scaling with significant reduction.
- Partially gapped Fermi surface of ordinary fermions or FL^* ?

Projected Fermi sea on Square lattice



- $L_A \log L_A$ scaling with significant reduction.
- Partially gapped Fermi surface of ordinary fermions or FL^* ?

Ongoing work

- Topological Renyi entropy of **projected BCS** and **Quantum Hall** states.
- Renyi entropy of partially projected Fermi sea (\equiv correlated Fermi liquid).

Ongoing work

- Topological Renyi entropy of **projected BCS** and **Quantum Hall** states.
- Renyi entropy of partially projected Fermi sea (\equiv correlated Fermi liquid).

Summary

- Renyi entropy calculations can serve as a **diagnostic for critical spin-liquids**.
- First example of a **area law violation in a fully 2D bosonic wave-function**, the projected Fermi sea state on the triangular lattice.
- **Area law** and **algebraically decaying correlations** for projected Dirac metal \Rightarrow **algebraic spin-liquid**.
- Determinantal Monte Carlo opens up a very **wide range of possibilities** for calculating entanglement entropy of correlated fermions.

Summary

- Renyi entropy calculations can serve as a **diagnostic for critical spin-liquids**.
- First example of a **area law violation in a fully 2D bosonic wave-function**, the projected Fermi sea state on the triangular lattice.
- **Area law** and **algebraically decaying correlations** for projected Dirac metal \Rightarrow **algebraic spin-liquid**.
- Determinantal Monte Carlo opens up a very **wide range of possibilities** for calculating entanglement entropy of correlated fermions.

Summary

- Renyi entropy calculations can serve as a **diagnostic for critical spin-liquids**.
- First example of a **area law violation in a fully 2D bosonic wave-function**, the projected Fermi sea state on the triangular lattice.
- **Area law** and **algebraically decaying correlations** for projected Dirac metal \Rightarrow **algebraic spin-liquid**.
- Determinantal Monte Carlo opens up a very **wide range of possibilities** for calculating entanglement entropy of correlated fermions.

Summary

- Renyi entropy calculations can serve as a **diagnostic for critical spin-liquids**.
- First example of a **area law violation in a fully 2D bosonic wave-function**, the projected Fermi sea state on the triangular lattice.
- **Area law** and **algebraically decaying correlations** for projected Dirac metal \Rightarrow **algebraic spin-liquid**.
- Determinantal Monte Carlo opens up a very **wide range of possibilities** for calculating entanglement entropy of correlated fermions.

Acknowledgements

Thank you to Matthew Fisher, Ann Kallin, Michael Levin, Roger Melko, Masaki Oshikawa, Brian Swingle for interesting and useful discussions!

Thank you to KITP and all the participants for the workshop!