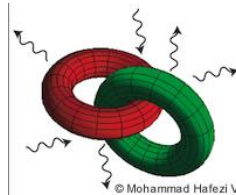


# COLLECTIVE ELECTRONIC EXCITATIONS FOR ULTRA-STRONG COUPLING AND SUPERRADIANCE

**Angela Vasanelli**

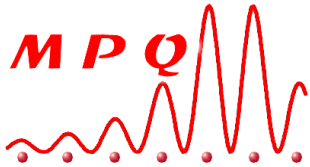
MPQ Laboratories, University Paris Diderot, France

KITP, 5-9 october 2015



# Acknowledgements

**Yanko Todorov and Carlo Sirtori**



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Clean room: Pascal Filloux, Christophe Manquest, Stéphane Suffit



Growth: G. Beaudoin, I. Sagnes

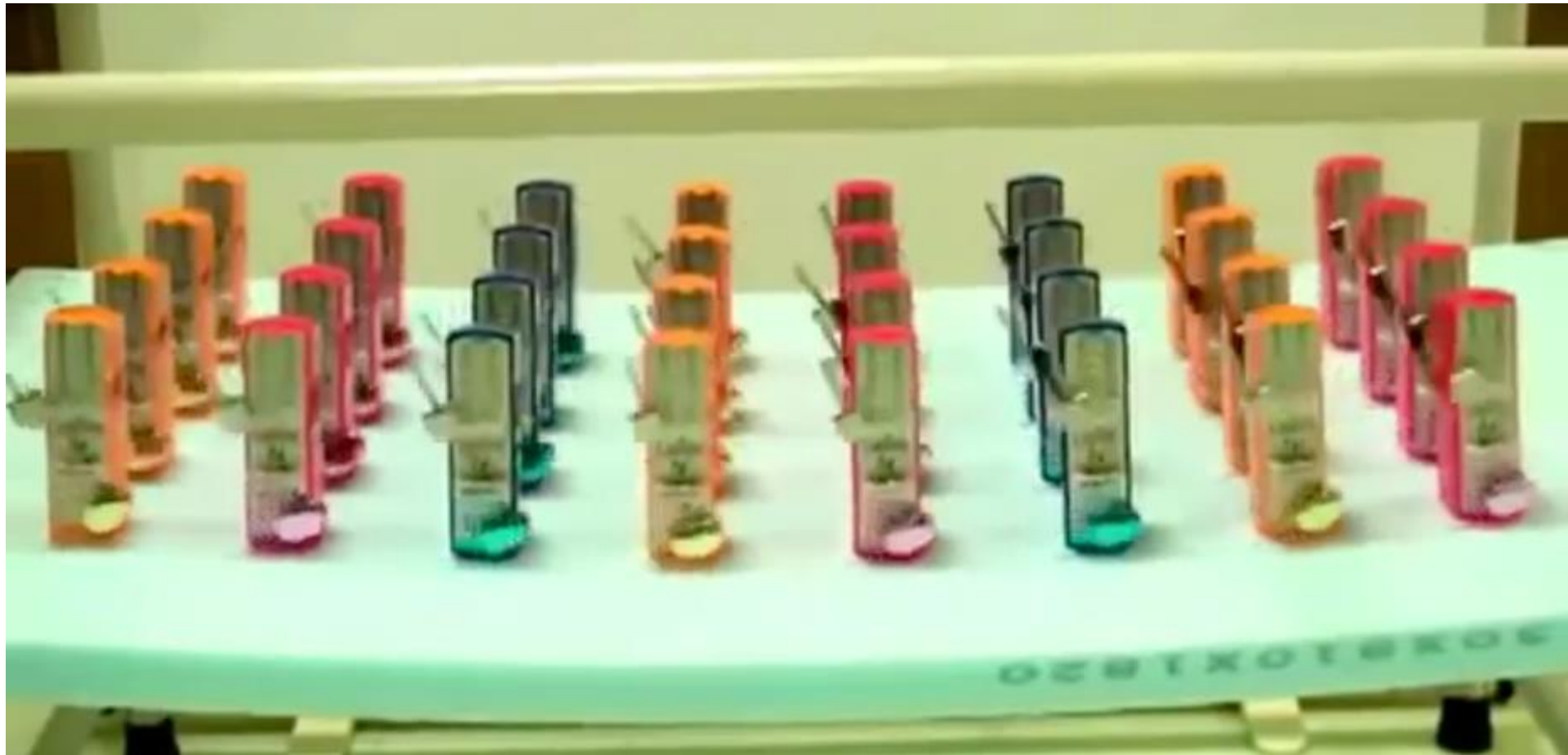
Funding:

European Research Council  
Executive Agency

ADEQUATE

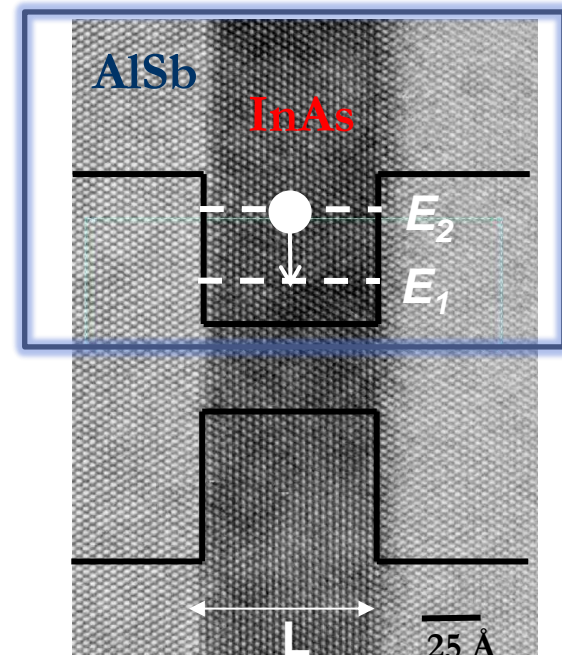
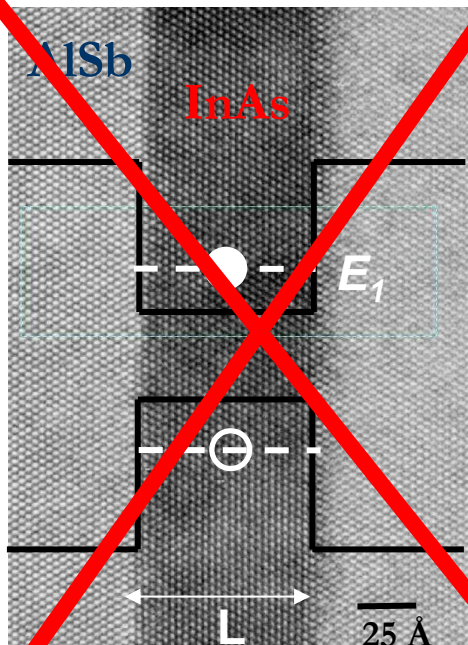


**This work is not on interacting photons...**



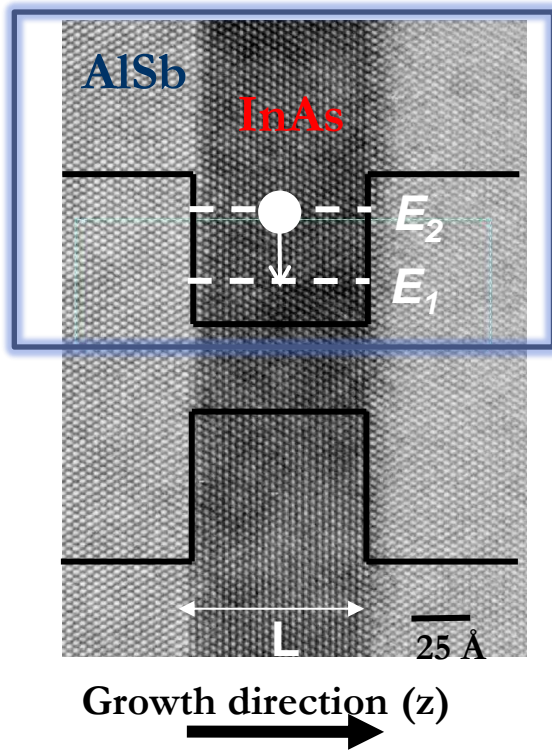
**...but on interacting electrons!**

# It is not about excitons...

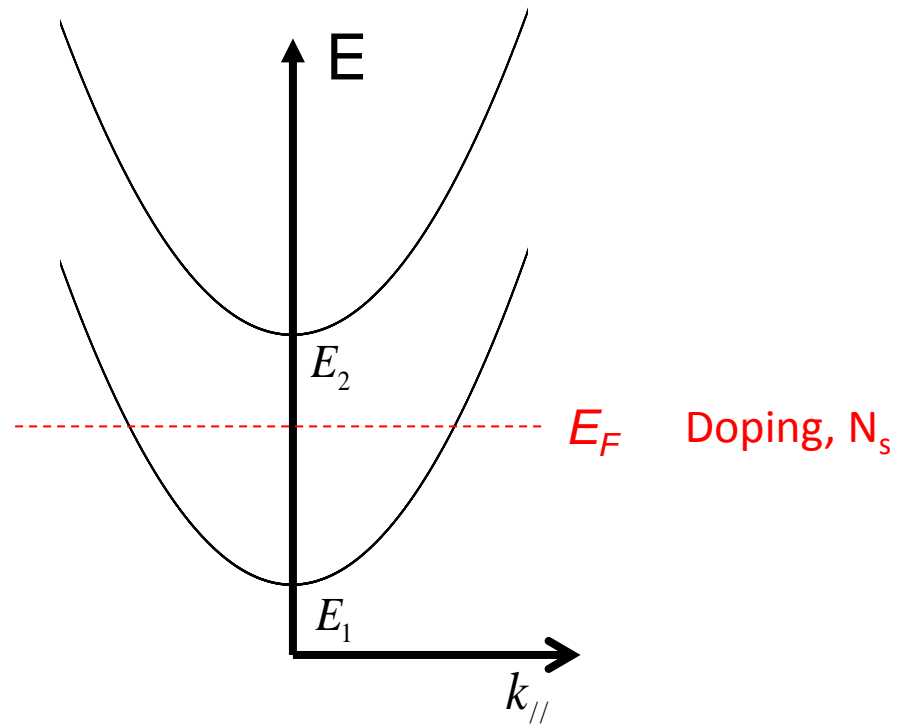


...everything happens in the conduction band of a semiconductor quantum well

# Intersubband transition in quantum wells



$$E_2 - E_1 \propto \frac{1}{L^2}$$



$$E_{n,k_{||}} = E_n + \frac{\hbar^2 k_{||}^2}{2m^*}$$

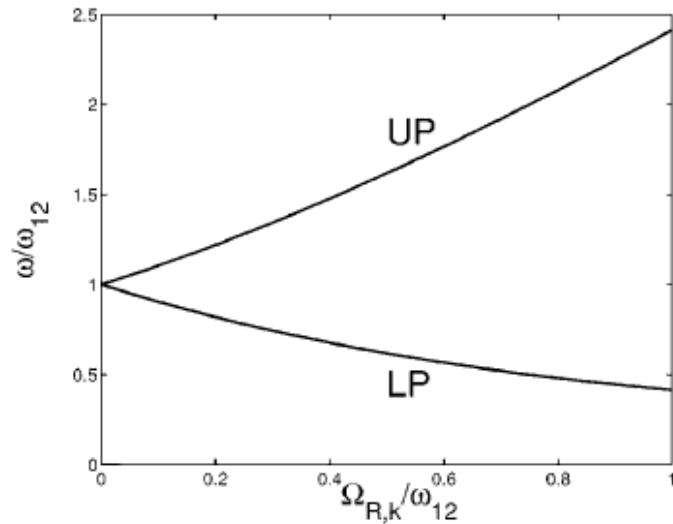
Photon wavelength : mid and far infrared  
 Spontaneous emission lifetime  $\sim 10 - 100$  ns  
 Non-radiative lifetime  $\sim 1$  ps

# Two different quantum phenomena:

- ❑ **Ultra-strong light-matter coupling**
- ❑ **Superradiance**

# Two different quantum phenomena:

- Ultra-strong light-matter coupling (USC)**

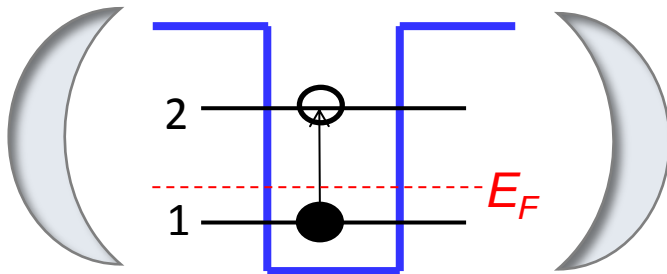


C. Ciuti *et al.*, PRB 2005

$$\frac{\text{Rabi energy}}{\text{Matter excitation energy}} = \frac{E_R}{E_{exc}} \approx 1$$

Anti-resonant and quadratic terms of the light – matter interaction are not negligible

*Günter et al Nature 2009, Todorov et al. PRL 2010, Delteil et al. PRL 2012, Geiser et al. PRL 2012; Niemczyk et al Nat. Phys. 2010; Scalari et al Science 2012, Schwartz et al. PRL 2011, Kéna-Cohen et al. Adv. Opt. Mater. 2013, Schlather et al. Nano Lett. 2013, Nagasawa et al. Phys. Chem. Lett. 2014...*

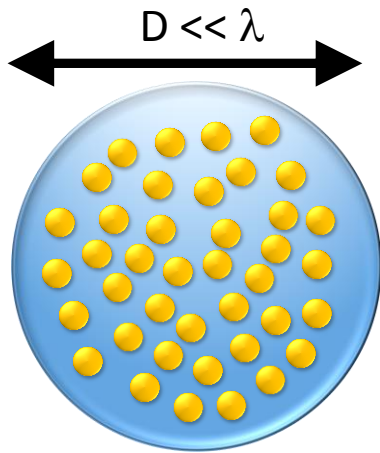


Rabi splitting:  $\Omega_R \propto \sqrt{f_{12} \frac{N_{QW}}{L_{cav}} N_1}$

D. Dini *et al.* PRL 2003

# Two different quantum phenomena:

- **Superradiance**



Phenomenon of cooperative recombination taking place in an ensemble of two-level emitters, when they are separated by a distance smaller than the wavelength.

Only one atom is excited:

$$|\psi\rangle = (|egg \dots\rangle + |geg \dots\rangle + |gge \dots\rangle) / \sqrt{N_e}$$

$$1 / (\underbrace{N_e}_{\text{circled}} \Gamma_0^{sp}) \text{ Spontaneous decay time}$$

Spontaneous emission rate for a single atom

R. H. Dicke, PR 1954

M. Gross, S. Haroche, Phys. Rep. 1982

Skribanowitz et al. PRL 1973

Scheibner et al. Nat. Phys. 2007

Van Loo et al. Science 2013

Zhang et al. PRL 2014

Goban et al. PRL 2015

Laurent et al. PRL 2015



**Common point:**  
**Many electronic excitations!**

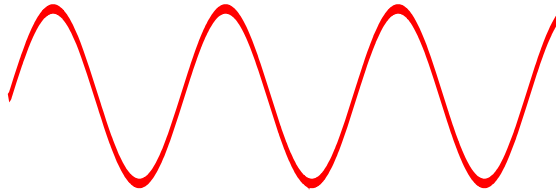
→ Dense electron gases in semiconductor QWs

# Outline

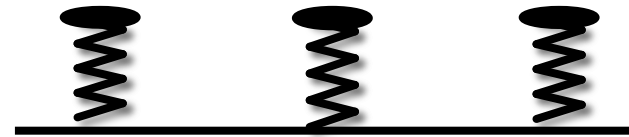
- Many-body excitations in dense electron gases
- Ultra-strong coupling regime with collective electronic excitations
- Spontaneous emission of collective states: superradiance
- Conclusion and perspectives

# **Many-body excitations in dense electron gases**

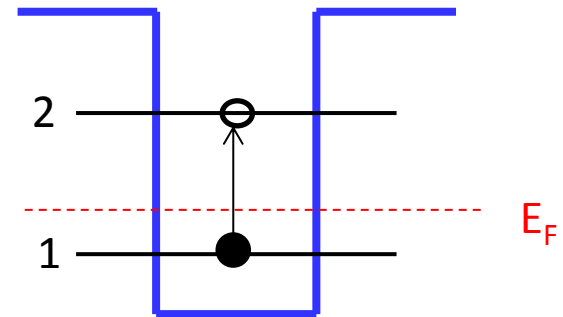
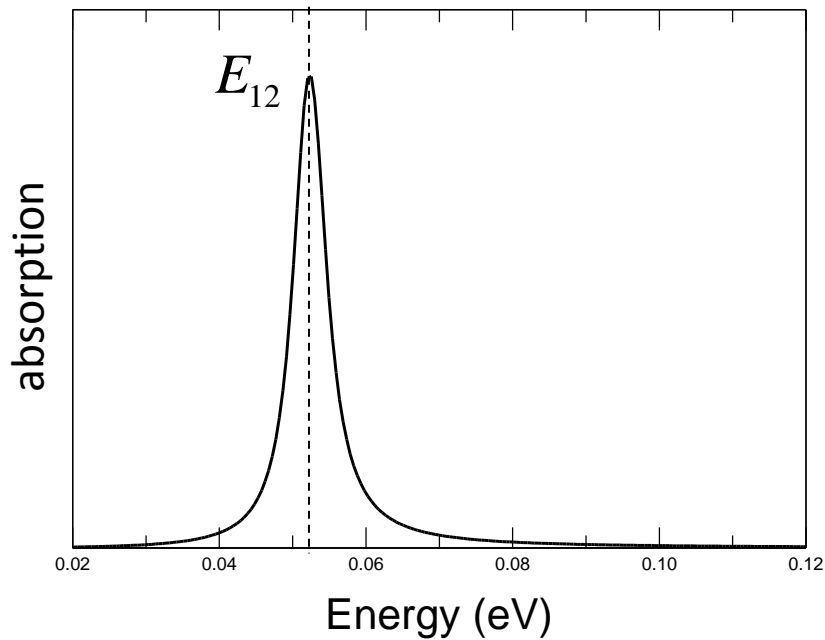
# Interaction free space radiation - 2DEG



Radiation

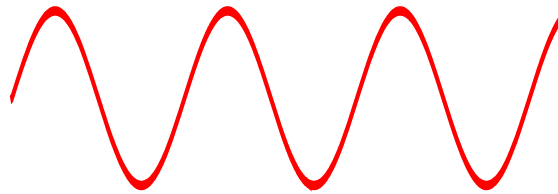


Matter polarization: excitation of an ensemble of non interacting dipolar oscillators

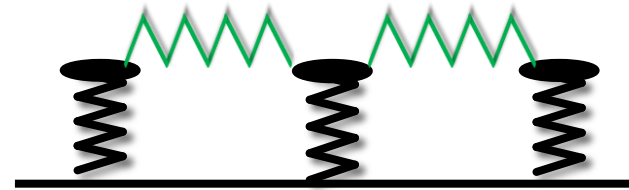


$$\alpha \propto N_s f_{12}$$

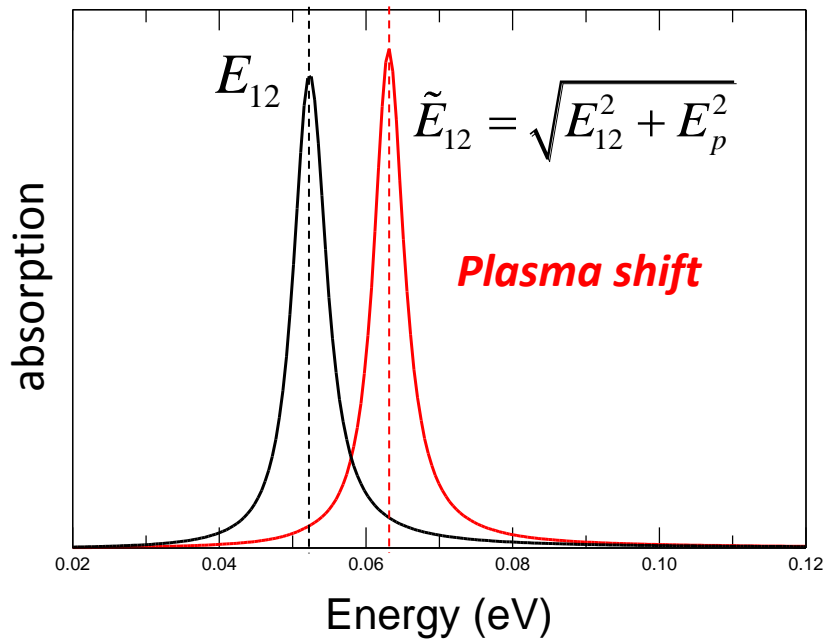
# Interaction free space radiation - 2DEG



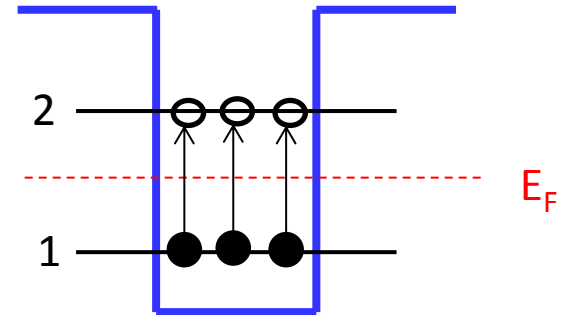
Radiation



Matter polarization: excitation of an ensemble of interacting dipolar oscillators

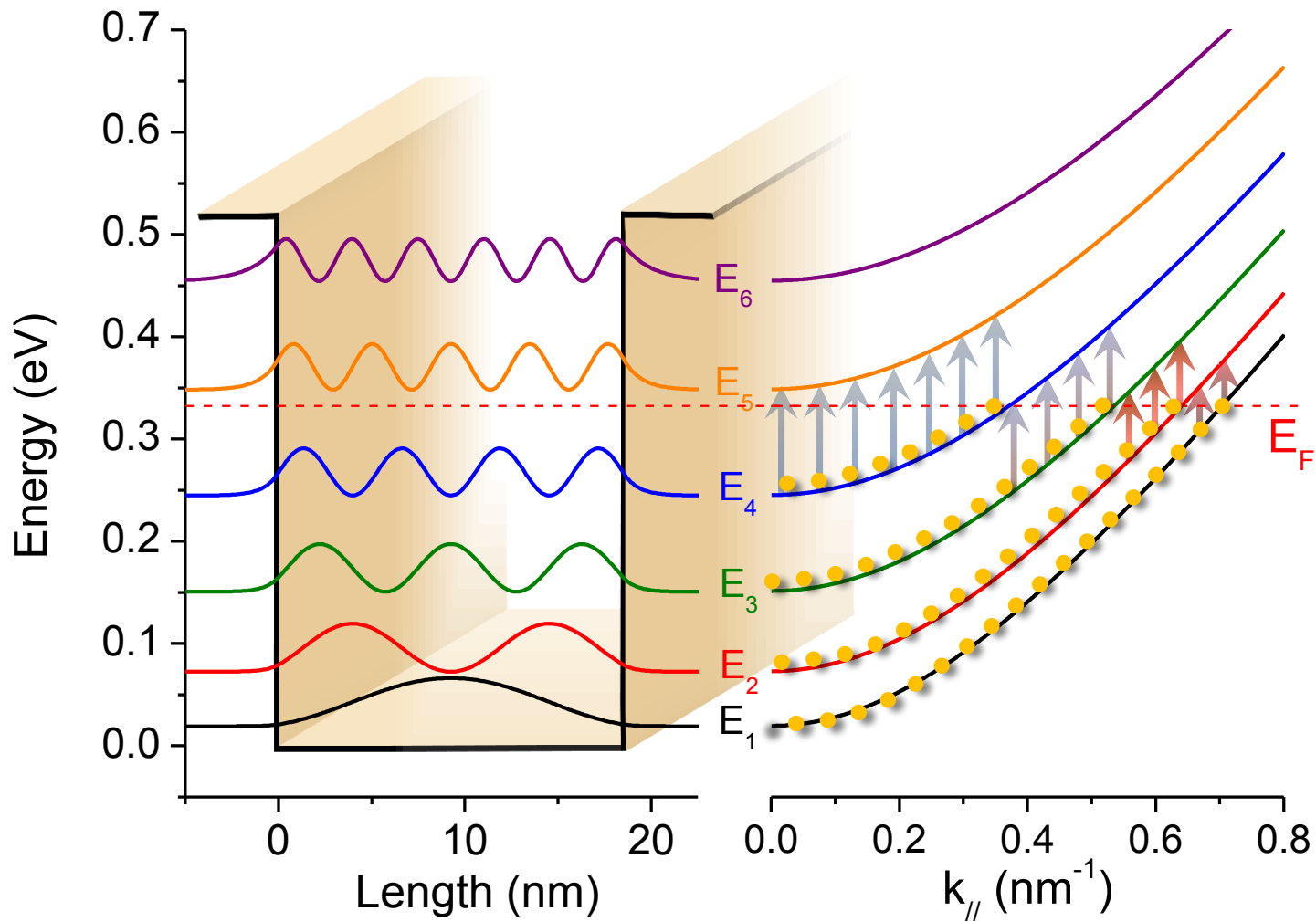


Ando, Fowler and Stern Rev. Mod. Phys. 1982



$$\alpha \propto E_p^2 = \frac{\hbar^2 e^2 f_{12}}{m^* \epsilon_0 \epsilon_{st} L_{eff}} (N_1 - N_2)$$

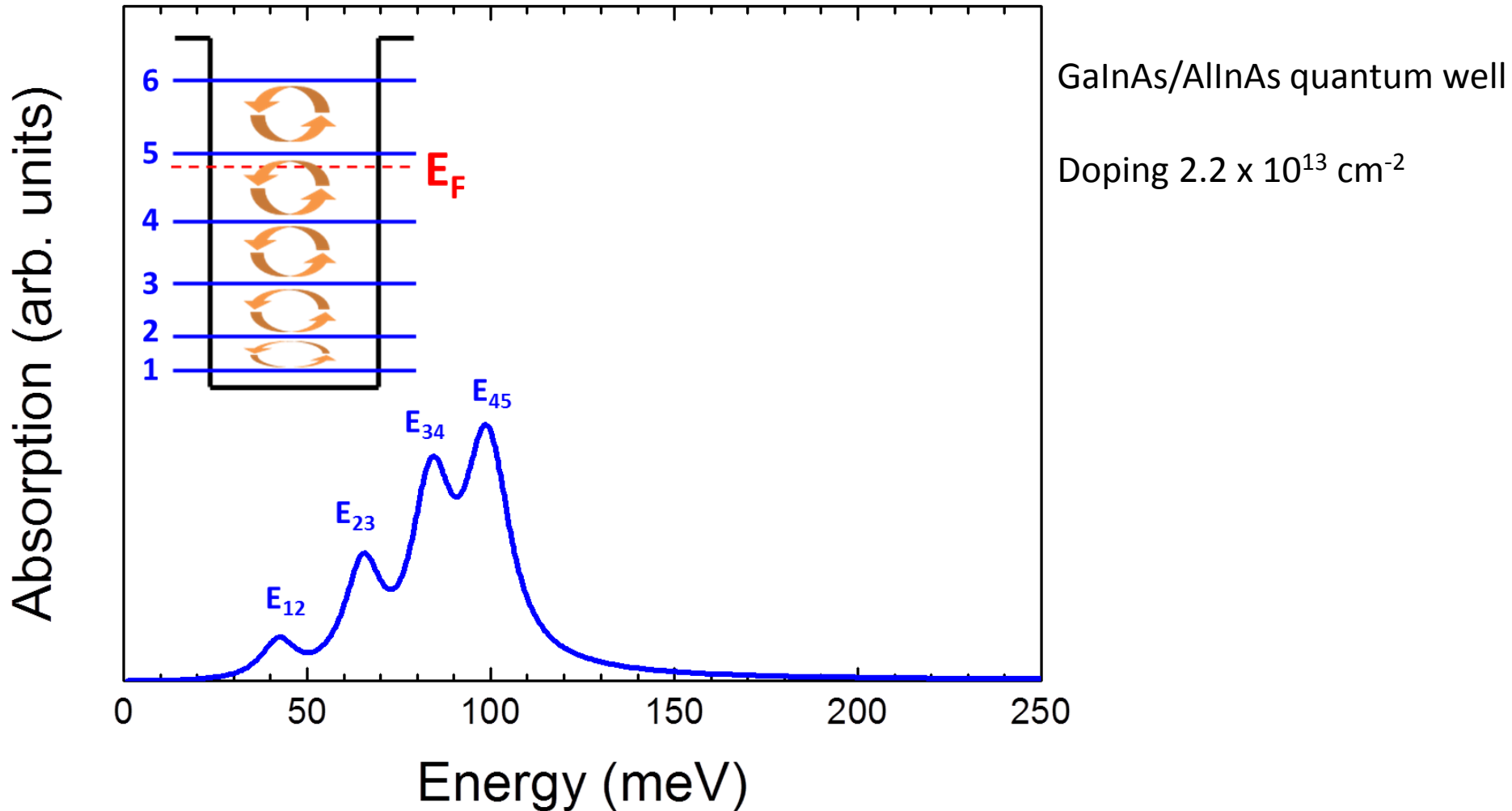
# Highly doped quantum well



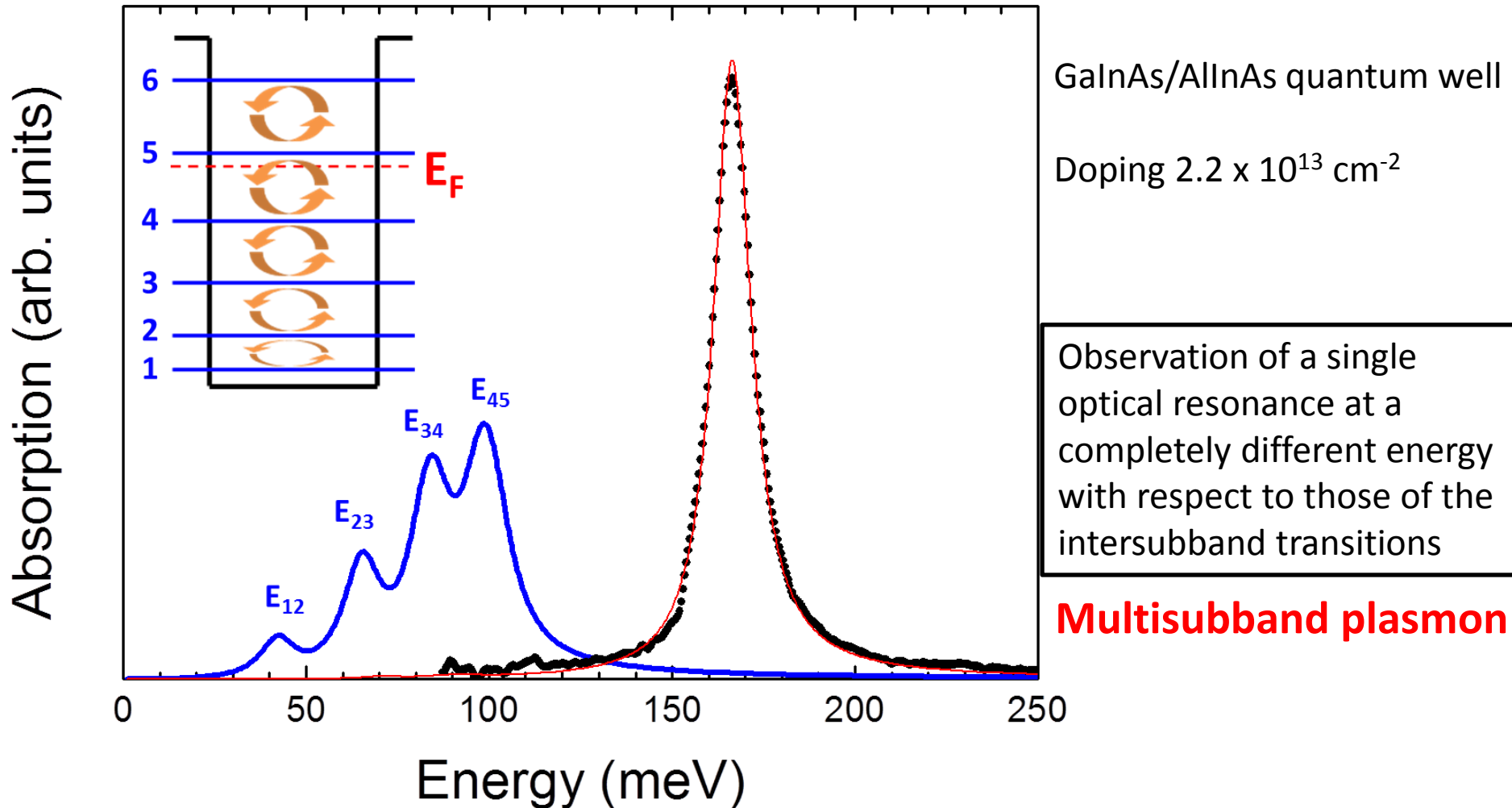
18.5 nm GaInAs/AlInAs quantum well

Uniformly doped in the well:  $2.2 \times 10^{13} \text{ cm}^{-2}$  ( $N_v = 1.2 \times 10^{19} \text{ cm}^{-3}$ )

# Single particle absorption spectrum

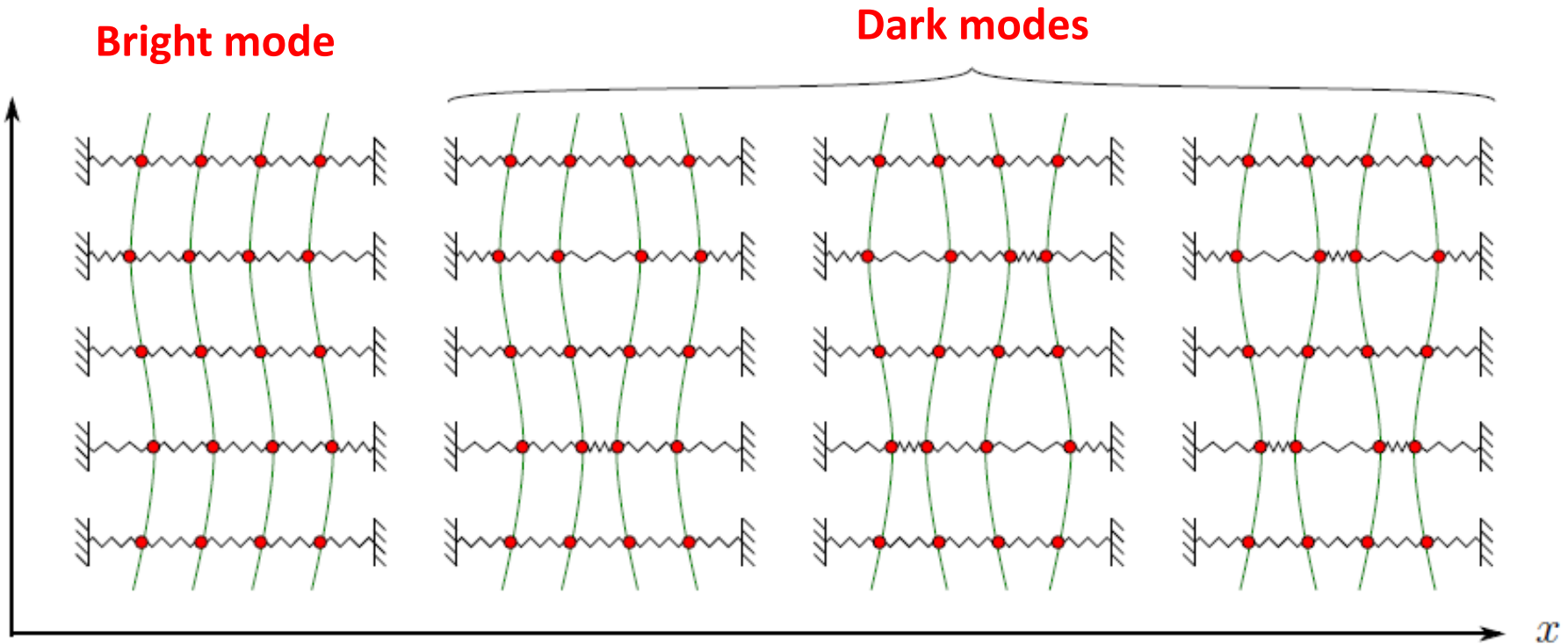


# Observation of the cooperative regime



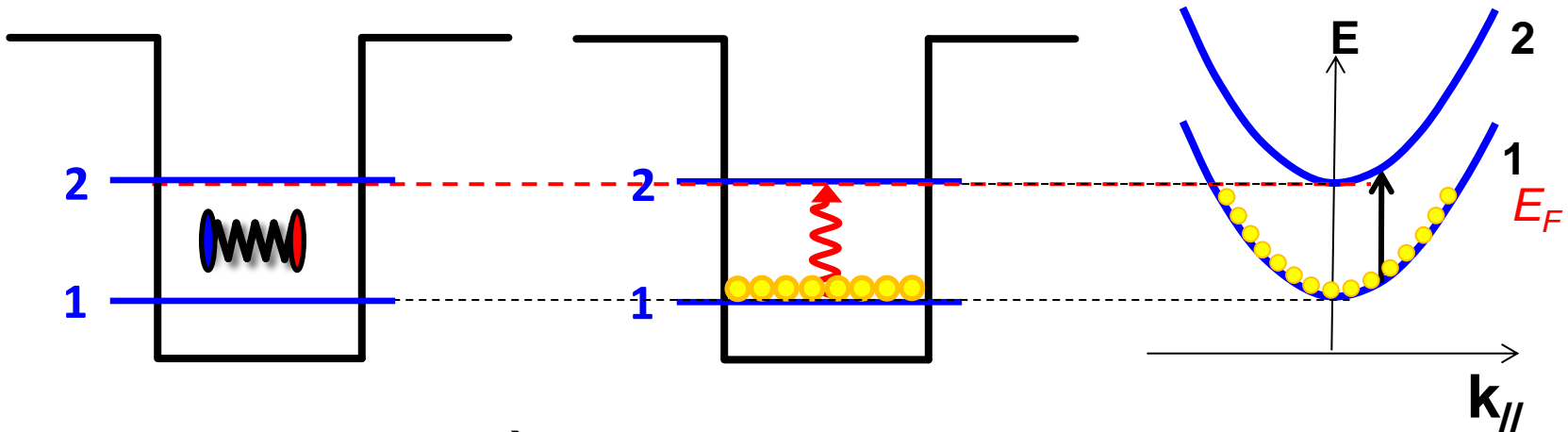


# Bright mode: intersubband dipoles locked in phase



The bright mode is the one in which the center of mass moves

# Quantum model: basic ideas



High electronic density  $\rightarrow$  each intersubband dipolar oscillator interacts not only with the external field, but also with all the other oscillators (**depolarization field**).

This can be taken into account by using the dipole representation to describe the light-matter interaction:

$$\hat{H}_{\text{int}} = \int \frac{1}{\epsilon_0 \epsilon(z)} \left[ -\hat{\mathbf{D}}(\mathbf{r}) \cdot \hat{\mathbf{P}}(\mathbf{r}) + \frac{1}{2} \hat{\mathbf{P}}^2(\mathbf{r}) \right] d^3 \mathbf{r}.$$

**Dipole – dipole interaction between intersubband excitations**

# Basic steps of the calculation

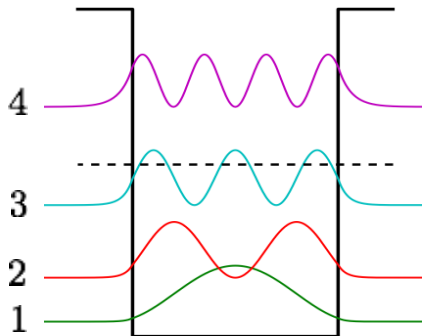
1

**Polarization** operator

$$P_z = \sum_{\alpha} \frac{j_{\alpha}(z)}{\omega_{\alpha}} [b_{\alpha}^{\dagger} + b_{\alpha}]$$

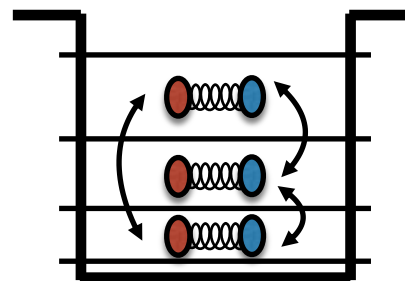
← intersubband excitation operator

$$j_{ij} \propto \psi_i \frac{\partial \psi_j}{\partial z} - \psi_j \frac{\partial \psi_i}{\partial z}$$



2

$$\hat{H}_{\text{int}} = \int \frac{1}{\epsilon_0 \epsilon(z)} \left[ -\hat{\mathbf{D}}(\mathbf{r}) \cdot \hat{\mathbf{P}}(\mathbf{r}) + \frac{1}{2} \hat{\mathbf{P}}^2(\mathbf{r}) \right] d^3 \mathbf{r}.$$

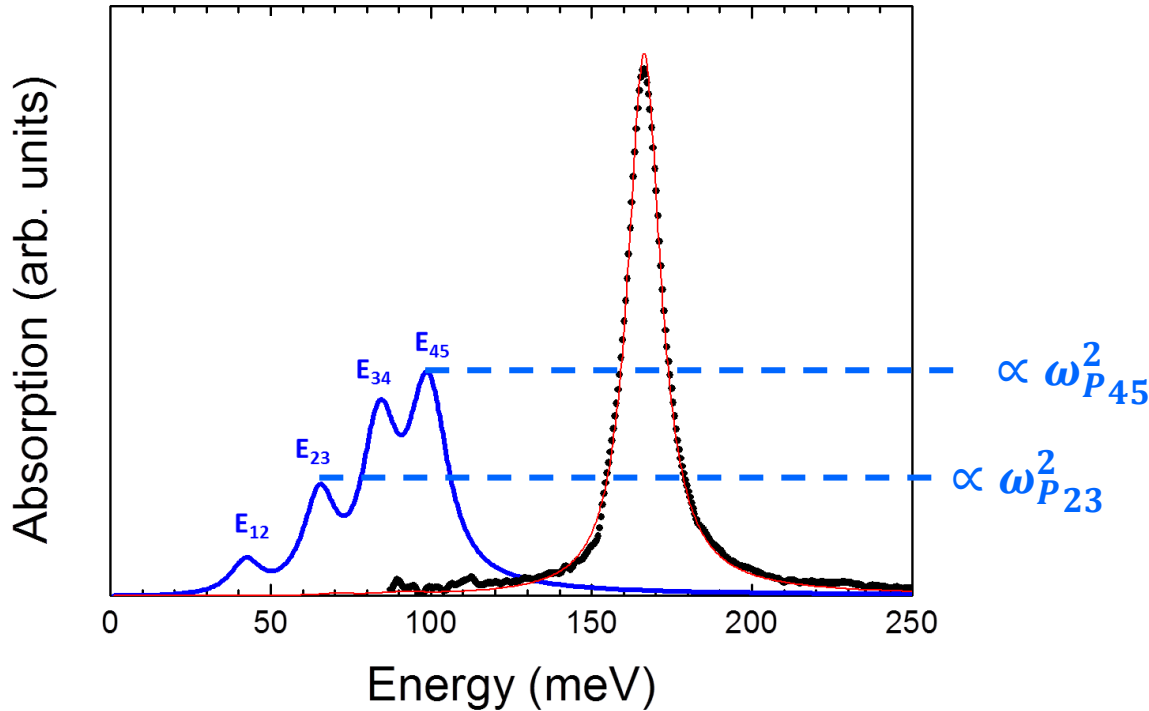


**Dipole – dipole  
interaction**

**Collective modes of the system (bright + dark)**

3

# Properties of the multisubband plasmon



The multisubband plasmon (MSP) concentrates the whole interaction with light



We can associate to the MSP an effective plasma frequency:

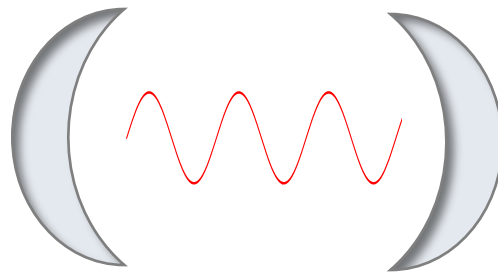
$$\Omega_P^2 = \sum_i \omega_{P_{i,i+1}}^2$$

We can also define an intersubband contribution such that:

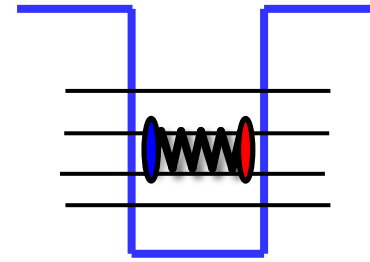
$$E_{MSP}^2 = E_{ISB}^2 + (\hbar\Omega_P)^2$$

$$E_{ISB}^2 = \frac{(\hbar\Omega_P)^2}{\sum_i \frac{\omega_{P_i}^2}{(\omega_{i+1} - \omega_i)^2}}$$

Harmonic mean of the intersubband transitions



Cavity mode



Intersubband  
polarization

# Ultra-strong light-matter coupling regime

$$H = E_{MSP} P^+ P + E_c \left( a^+ a + \frac{1}{2} \right) + i E_R (a^+ - a) (P^+ + P)$$

MSP  
plasmon

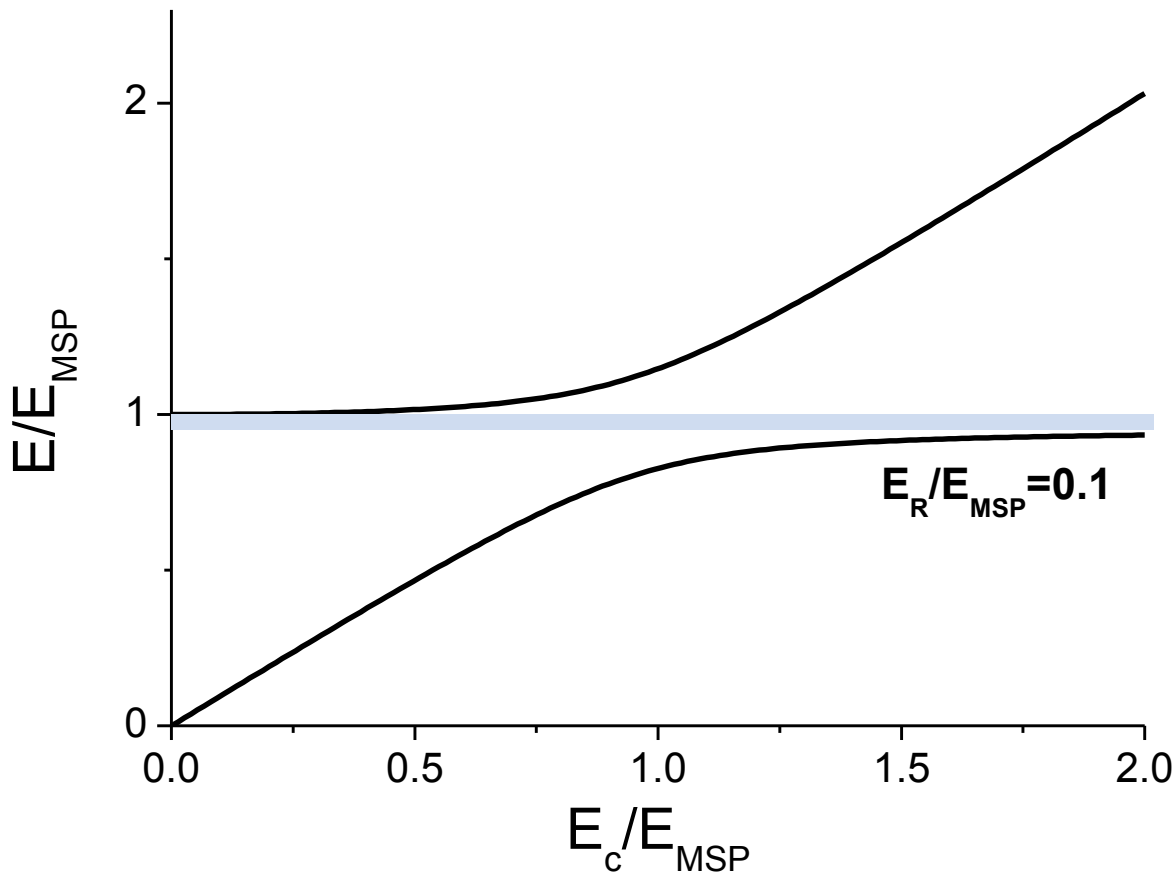
light

$$E_R \propto \hbar \Omega_P$$

Light-matter coupling

# Coupling of a MSP and a cavity mode

$$(E^2 - E_{MSP}^2)(E^2 - E_C^2) = E_R^2 E_C^2$$

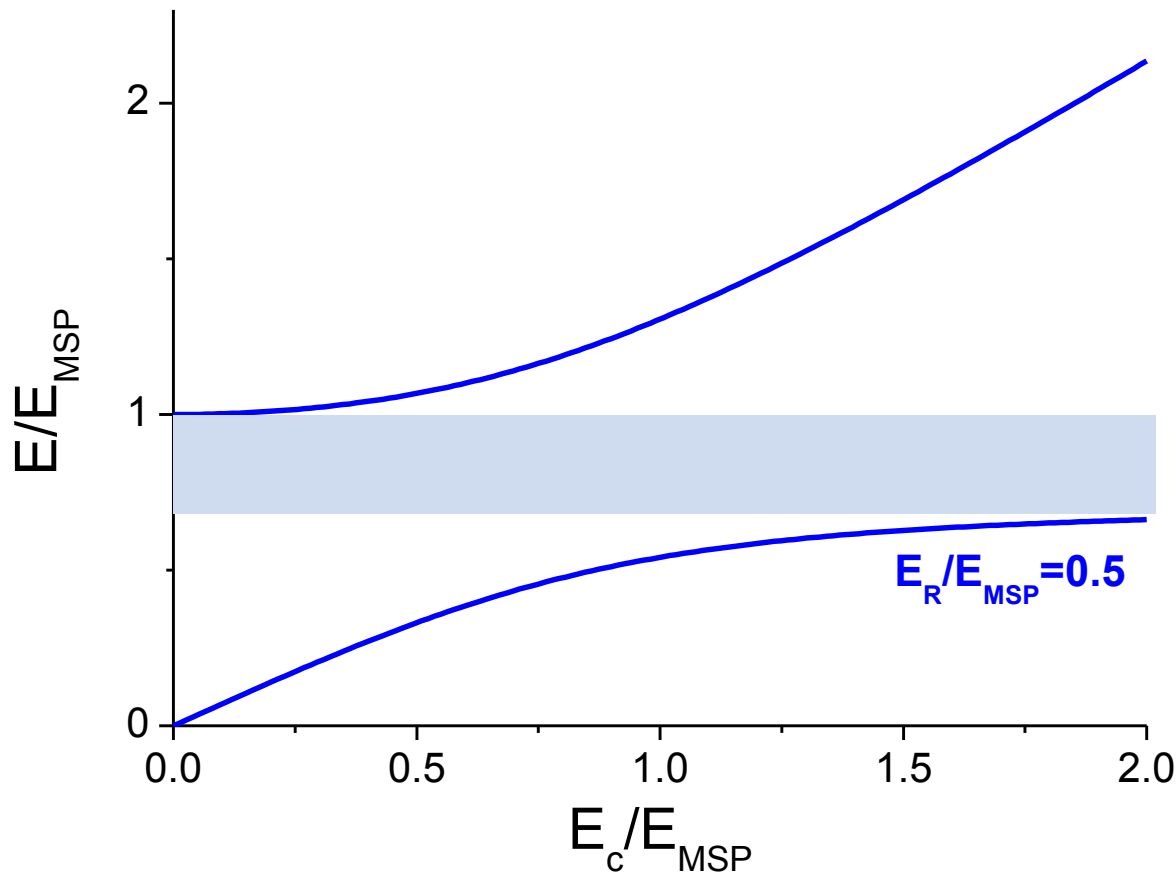


Opening of a polaritonic gap  
in the polariton dispersion

$$E_{MSP} \sqrt{1 - \left(\frac{E_R}{E_{MSP}}\right)^2}$$

# Coupling of a MSP and a cavity mode

$$(E^2 - E_{MSP}^2)(E^2 - E_c^2) = E_R^2 E_c^2$$

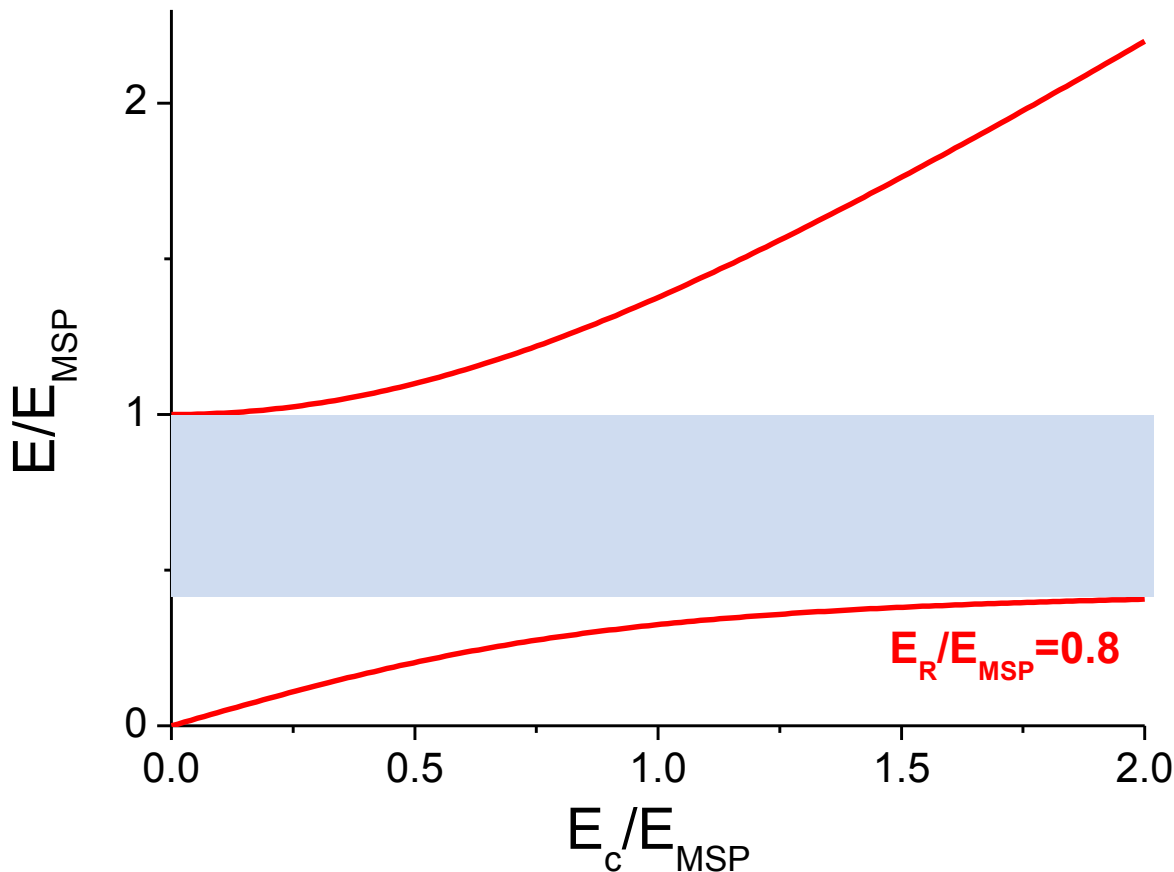


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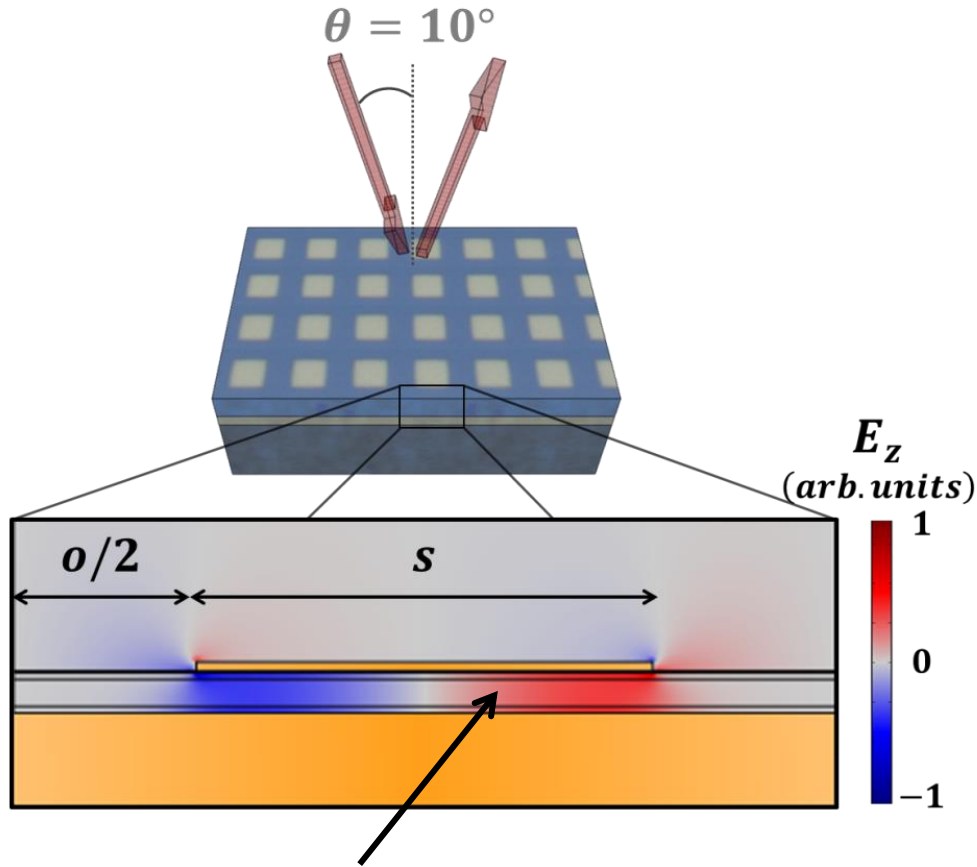


Opening of a polaritonic gap  
in the polariton dispersion

$$E_{MSP} \sqrt{1 - \left(\frac{E_R}{E_{MSP}}\right)^2}$$

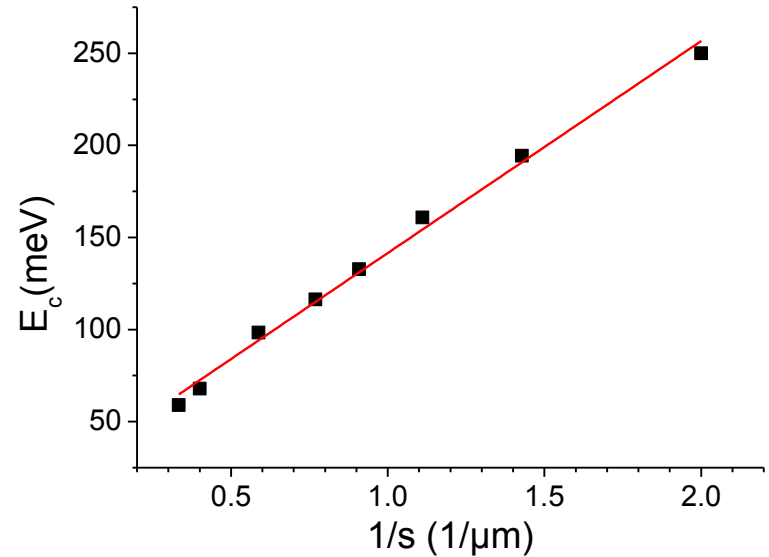


# Highly sub- $\lambda$ plasmonic cavities



Highly doped semiconductor layer displaying a Berreman mode

**Cavity thickness: 240nm**



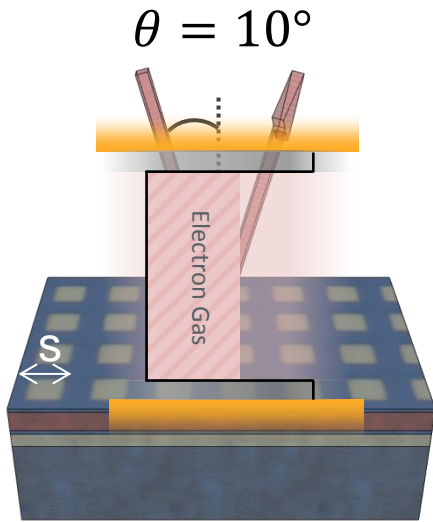
First order cavity mode

$$s = \lambda / (2n_{eff})$$

Y. Todorov *et al*, PRL **105**, 196402 (2010)

P. Jouy *et al*. APL **98**, 231114 (2011)

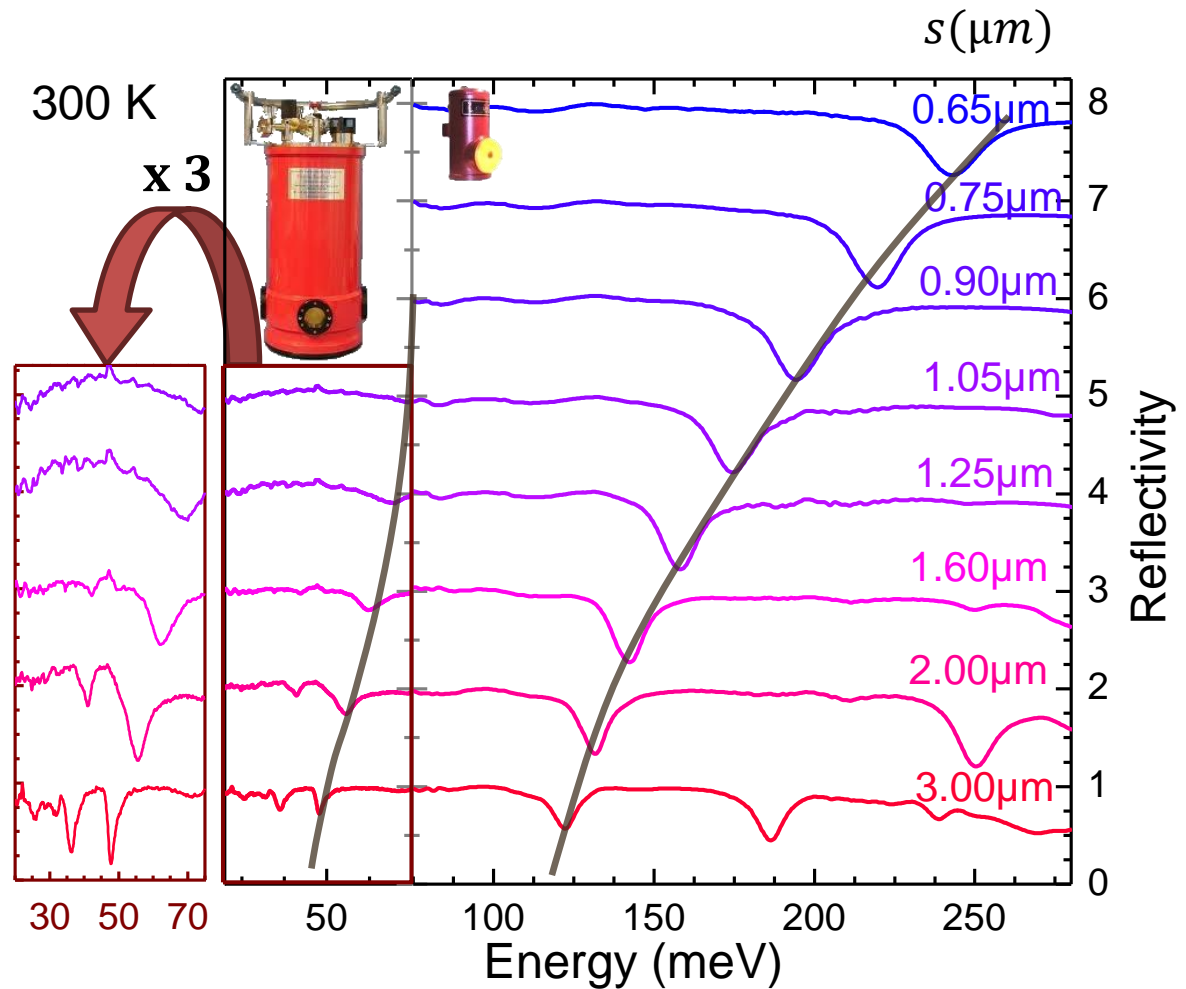
# Reflectivity Spectra



$$s = \lambda / (2n_{eff})$$

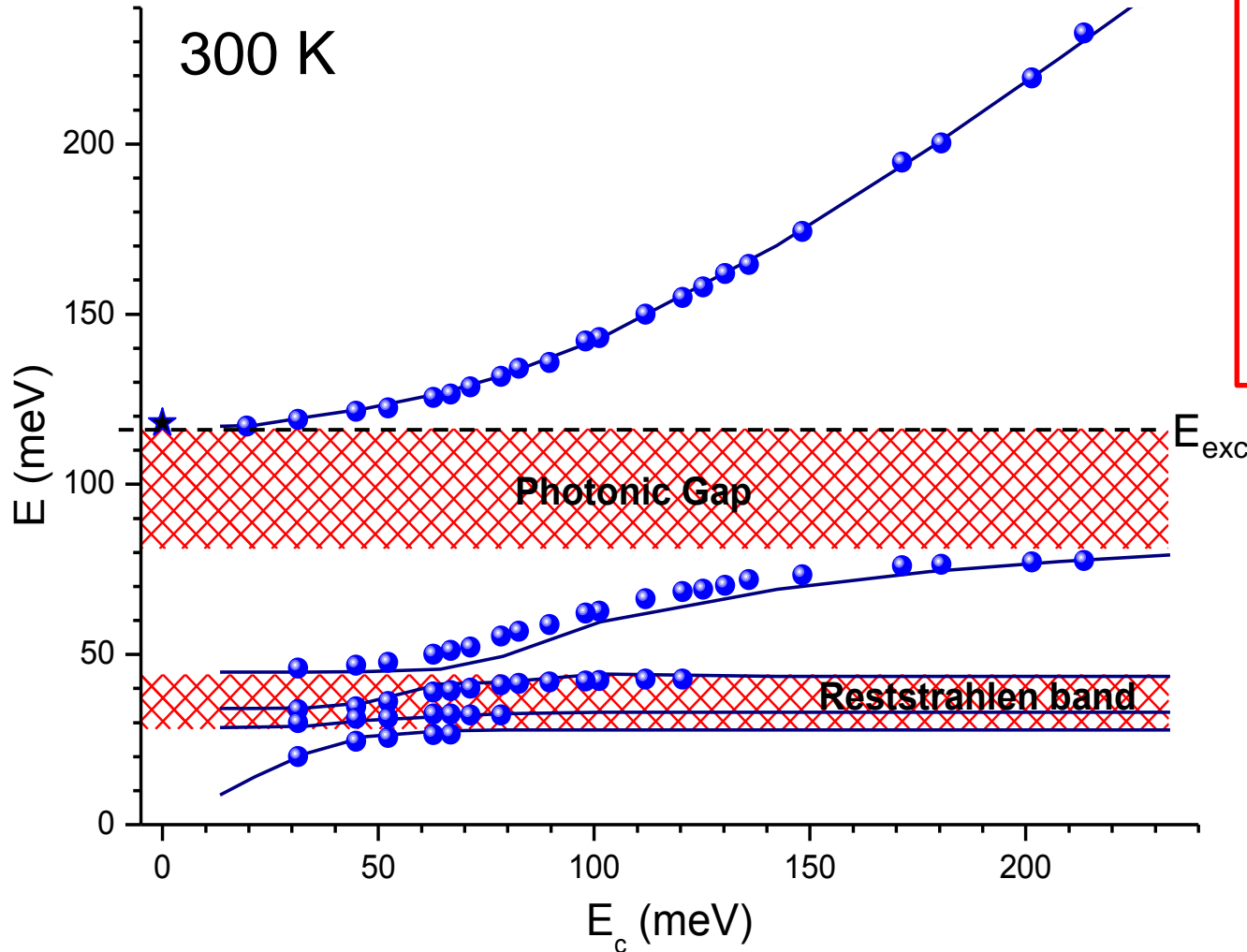
Y. Todorov *et al*, PRL (2010)

P. Jouy *et al*. APL (2011)



B. Askenazi *et al.*, New Journal of Physics **16**, 043029 (2014)

# Ultra-strong coupling dispersion



$$E_R = 86 \text{ meV}$$

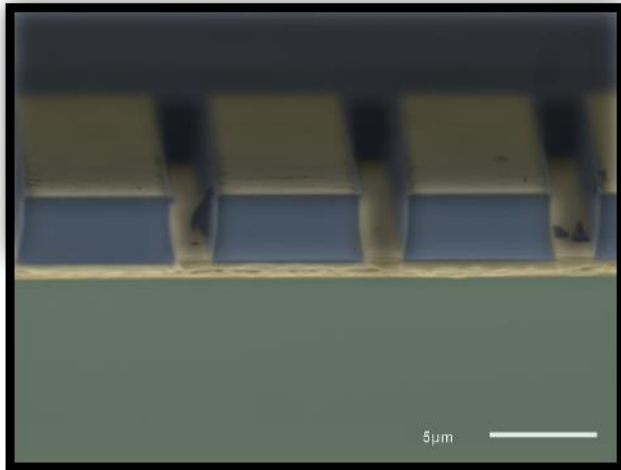
$$\frac{E_R}{E_{exc}} = 73\%$$

$$\frac{E_R}{E_{exc}} \approx 1$$

Ultra-strong coupling regime

C. Ciuti et al. PRB 2005

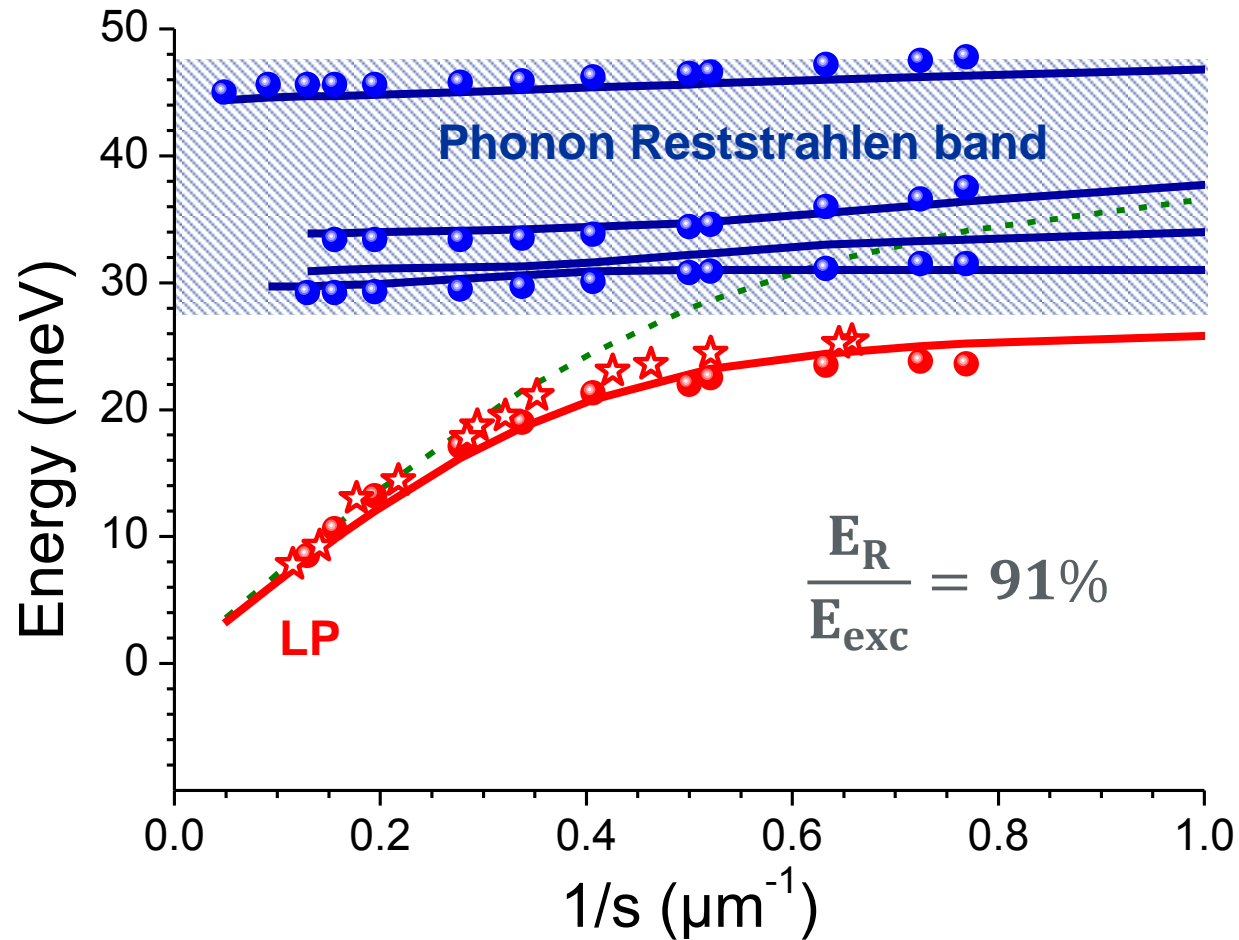
# Increasing the overlap factor



Etched cavities to improve the mode confinement

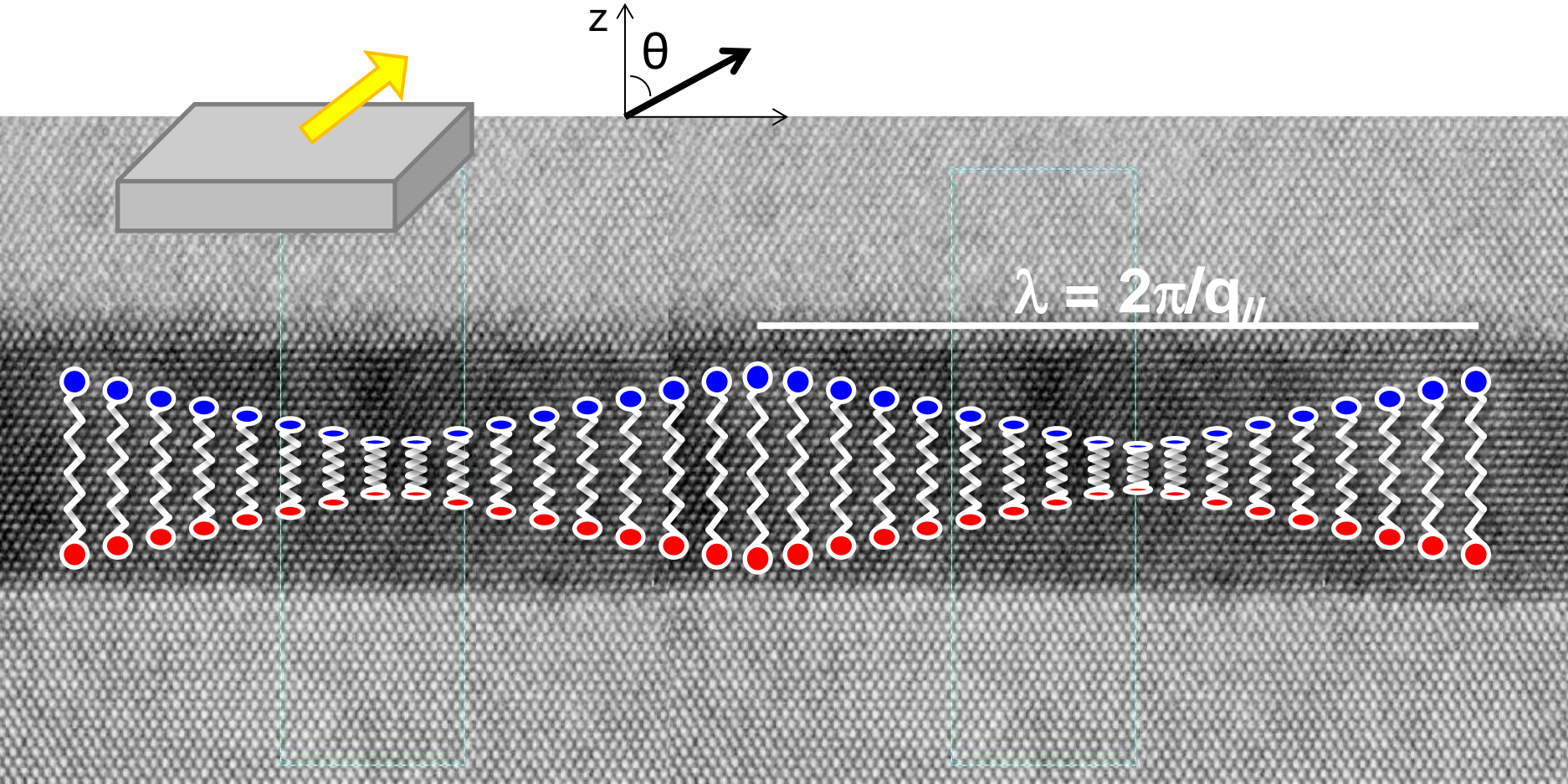
$$E_R = 90.0 \text{ meV}$$

$$E_{exc} = 99 \text{ meV}$$



# **Coupling of the plasmon with free space radiation**

# The multisubband plasmon: a transverse density wave



The dipole of the collective excitation is proportional to  $N$ , the number of particles involved

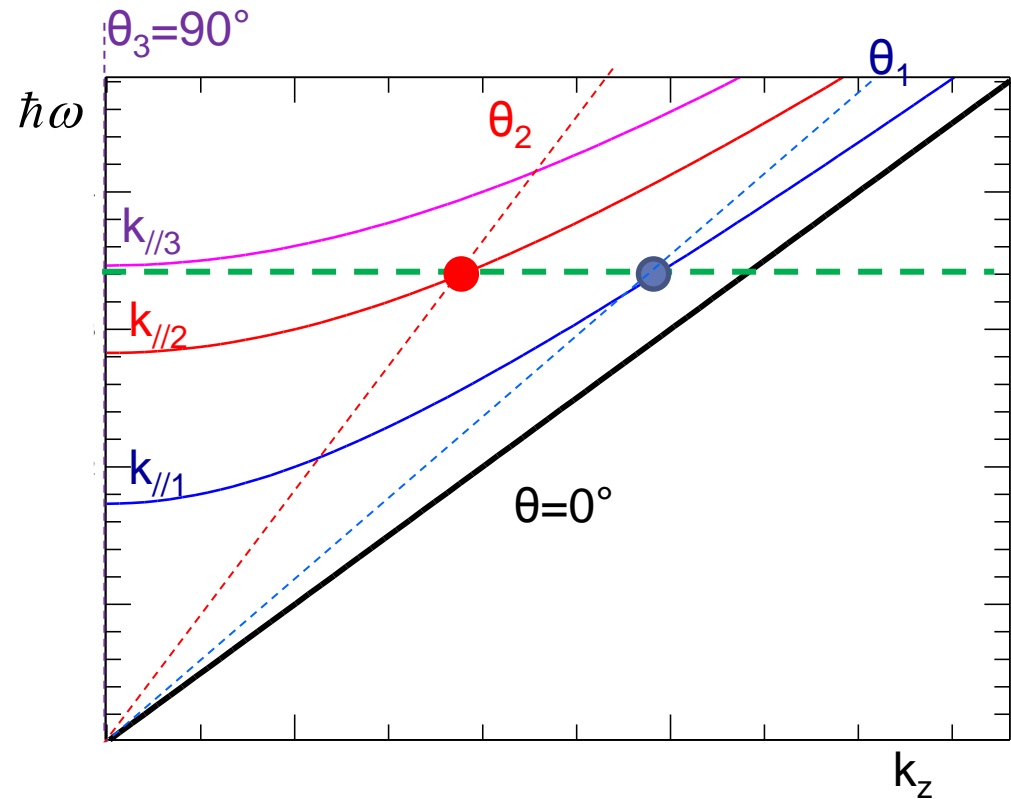
**Macroscopic polarization!**

# Spontaneous emission rate of a plasmon

Intrinsic radiative decay due to the lack of wavevector conservation along the confinement direction  $z \rightarrow$  The 2D plasmon interacts with a quasi-1D density of photon states:

$$\rho(k_{//}, \omega) = \sum_{k_z} \delta\left(\hbar\omega - \frac{\hbar c}{n} \sqrt{k_{//}^2 + k_z^2}\right)$$

$$\propto \frac{1}{\cos \theta}$$



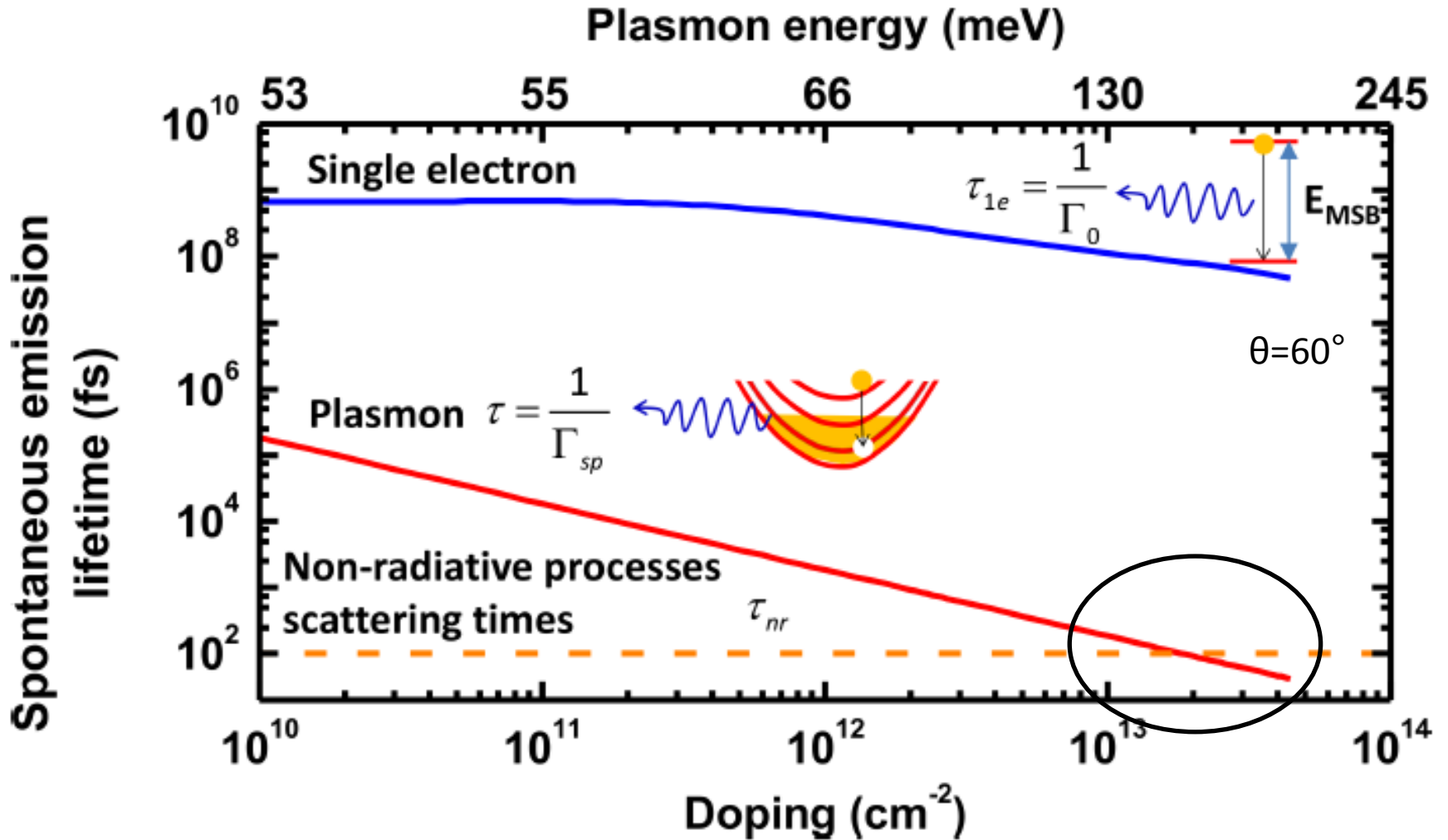
**Spontaneous emission rate by Fermi golden rule for a given plasmon state (selected by  $\theta$  angle):**

$$\Gamma_{\theta} = N_s \frac{1}{2} \frac{e^2}{m^* c n \varepsilon_0} \frac{\sin^2 \theta}{\cos \theta}$$

**Superradiant decay!**

C. Ciuti and I. Carusotto, PRA 2006  
 F. Alpegiani and L. C. Andreani, PRB 2014  
 T. Laurent et al. PRL (in press)

# Spontaneous emission rate of a plasmon

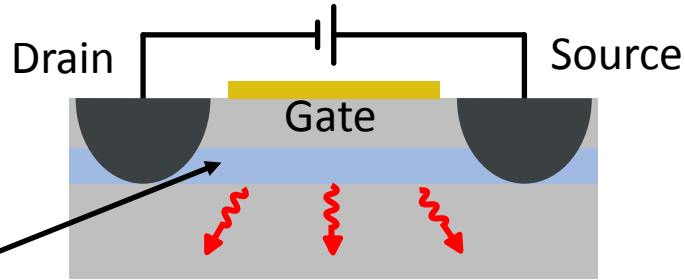


**Spontaneous emission** can be the dominant relaxation mechanism for multisubband plasmons!

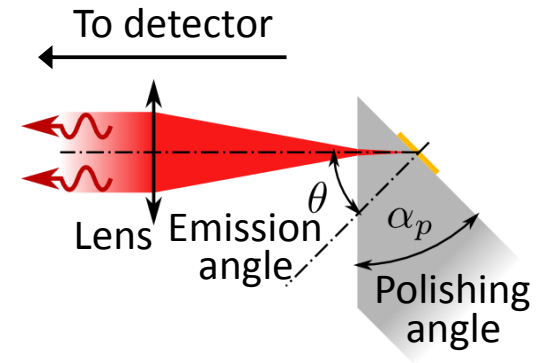
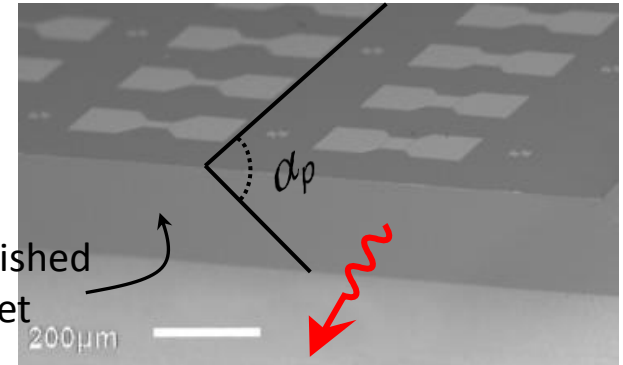
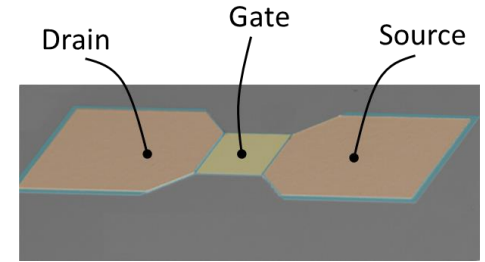
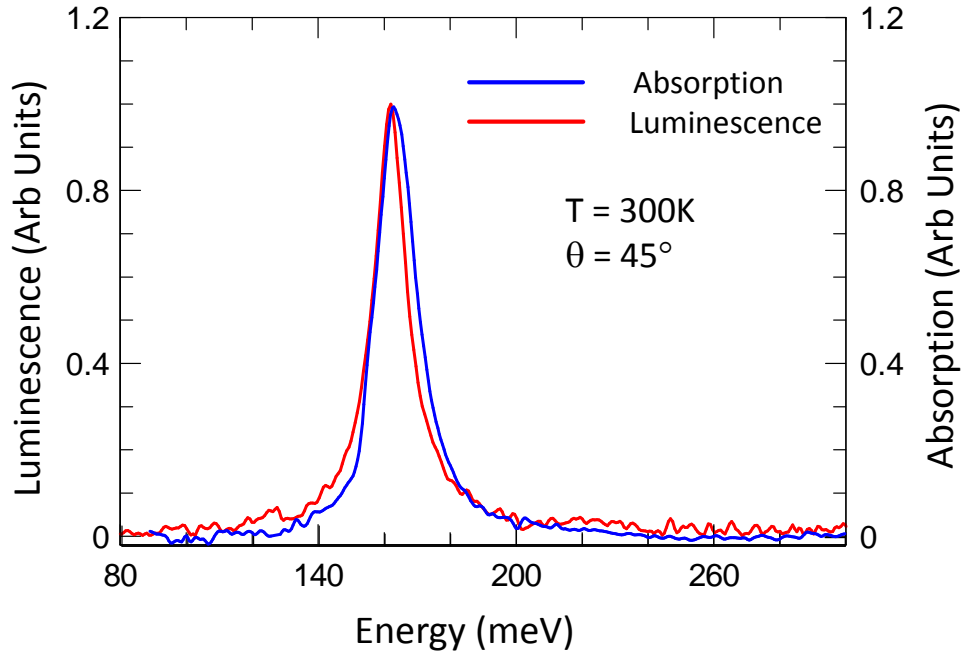


# Thermal excitation

Plasmons are heated by the current: *the plasmonic lamp*

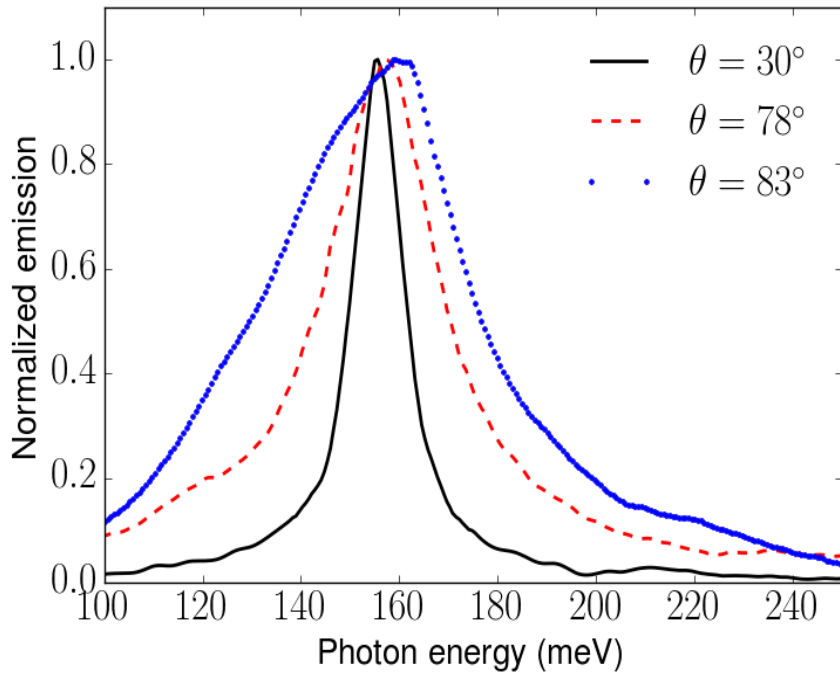


Heavily doped quantum well

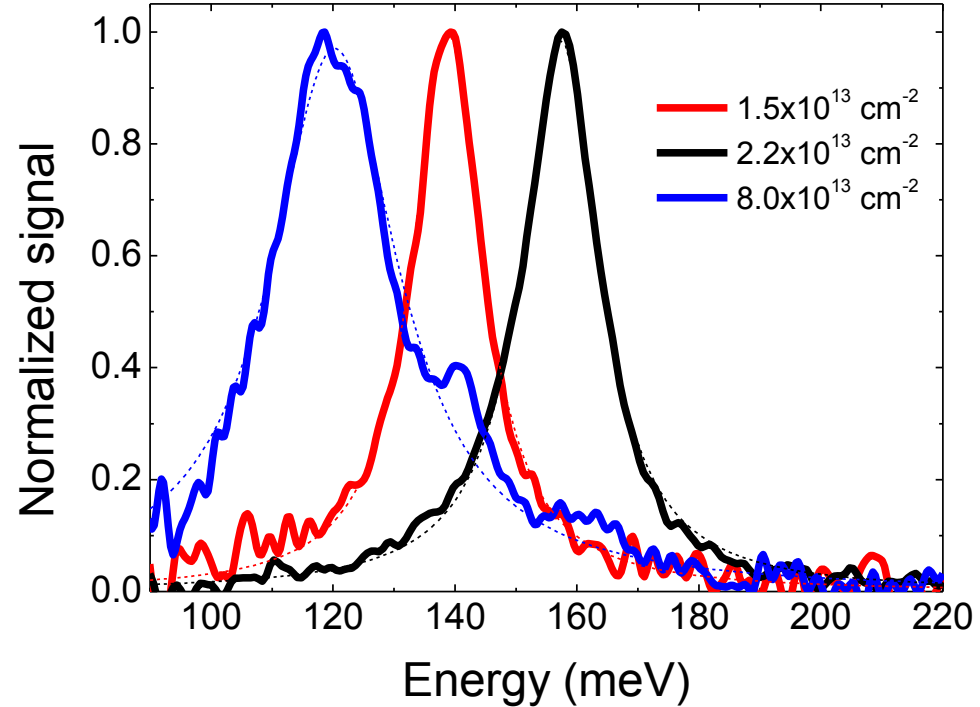


# Plasmon linewidth

The linewidth of the resonance reflects the plasmon lifetime

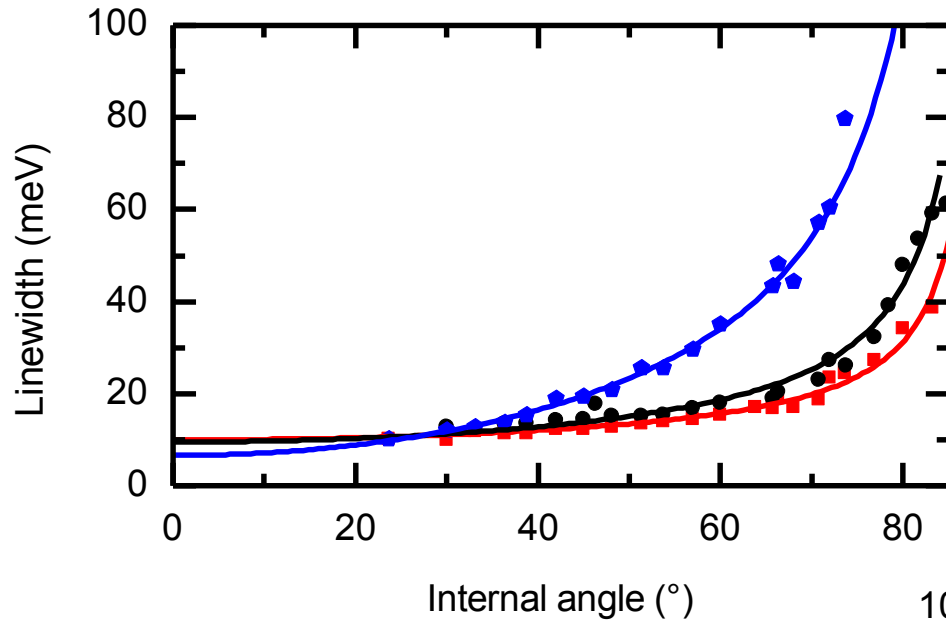


Doping  $N_s = 2.2 \times 10^{13} \text{ cm}^{-2}$



$\theta = 55^\circ$

# Radiative broadening

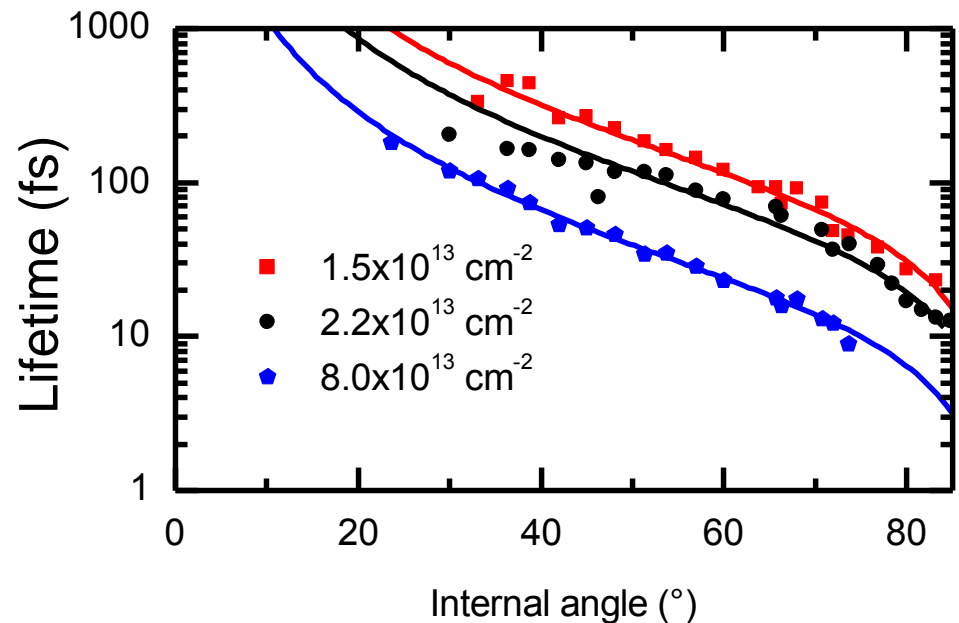


$$\Gamma = \Gamma_{sp}(\theta) + \gamma$$

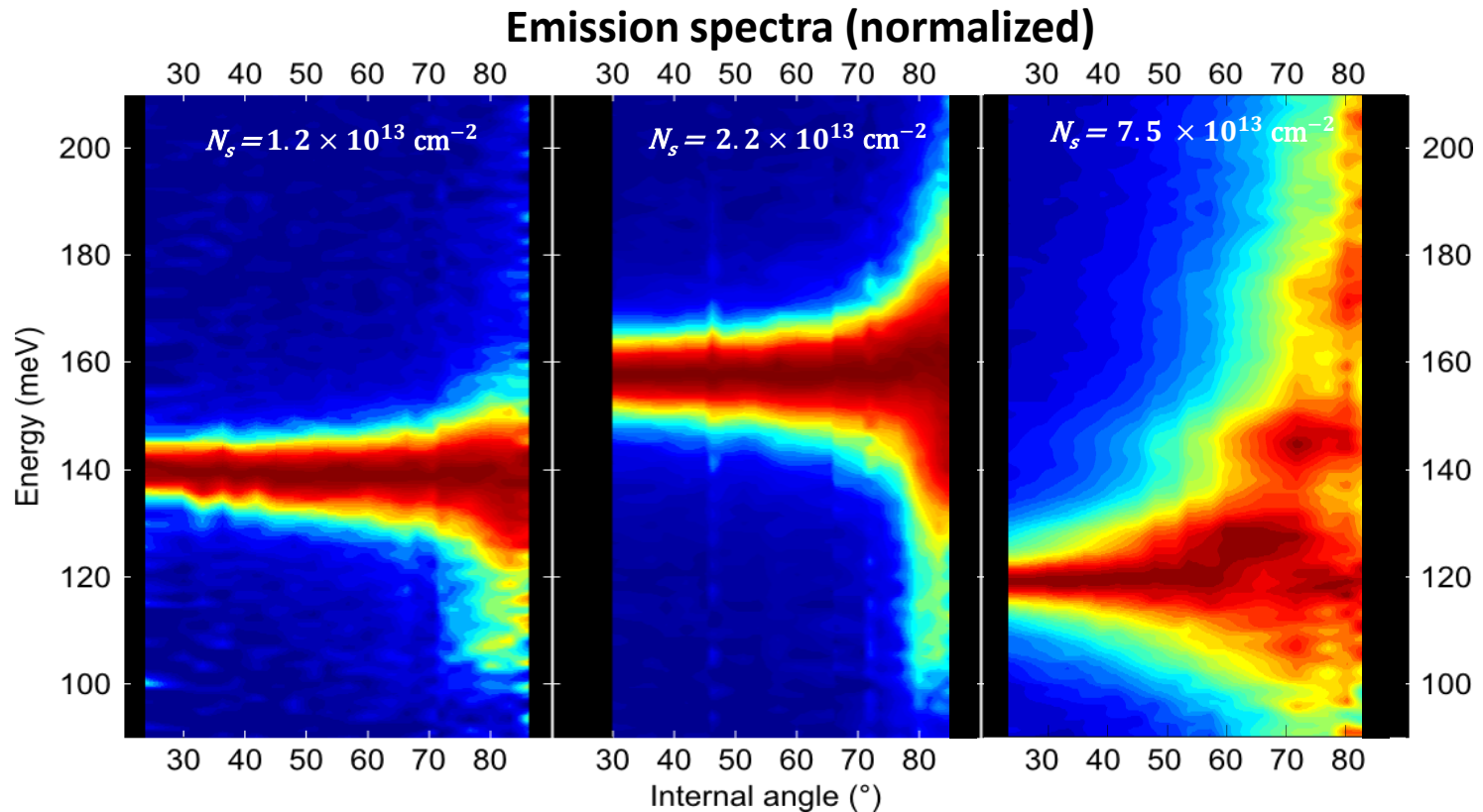
$$\Gamma_{sp}(\theta) = N_s \frac{1}{2} \frac{e^2}{m^* c n \epsilon_0} \frac{\sin^2 \theta}{\cos \theta}$$

Spontaneous emission of 10's of fs!

$$\Gamma_0 \sim 50\text{ns} \rightarrow \frac{\Gamma_{sp}}{\Gamma_0} \sim 10^6$$



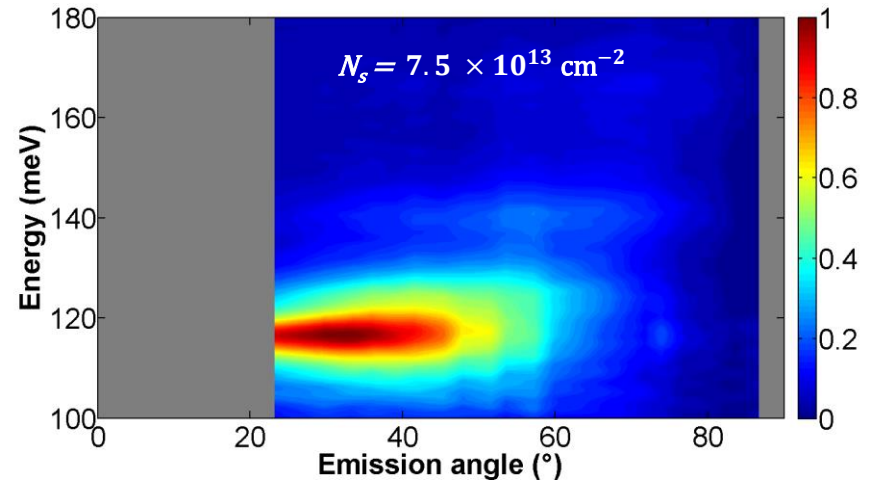
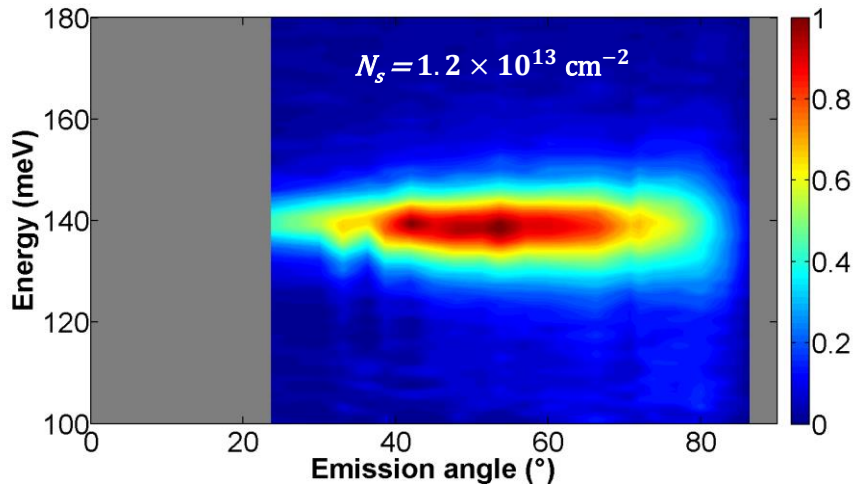
# Radiative broadening color maps



- Emission spectra are radiatively broadened
- Density dependence
- Extra features at high energy in highly doped samples

# Directional emission

## Emission spectra (non-normalized)



- Existence of a preferential emission direction
- The emission direction strongly depends on the electron density

Emitted power per unit solid angle:

$$dP = E_{MSP} \Gamma_{sp}(\theta) n(E_{MSP}, T) dN_{pl} \propto \sin^2 \theta \quad (\text{Larmor formula})$$

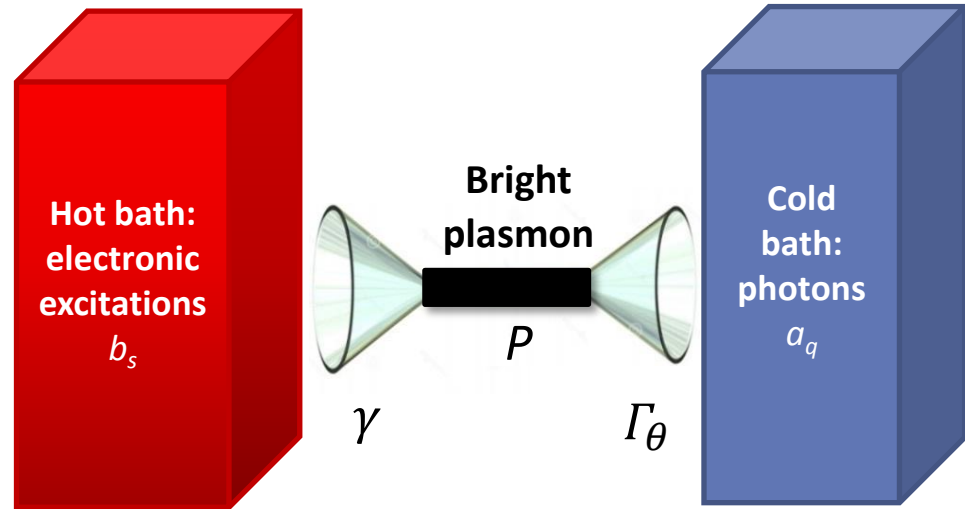
Plasmon occupancy

Number of plasmon modes in the solid angle of detection



# Quantum Langevin model

Bright plasmons are generated by thermal electronic excitation and decay by emitting a photon



Coupled system hamiltonian :

$$H = \sum_n \hbar\omega_n P_n^\dagger P_n + \sum_s \hbar\omega_s b_s^\dagger b_s + \sum_q \hbar\omega_q \left( a_q^\dagger a_q + \frac{1}{2} \right) + i\hbar \sum_{n,q} W_{n,q} [a_q^\dagger P_n - a_q P_n^\dagger] + i\hbar \sum_{n,s} K_{n,s} b_s^\dagger P_n - K_{n,s}^* b_s P_n^\dagger$$

Light-matter coupling

Coupling to the electronic bath



$\Gamma_\theta$

$\gamma$

# Kirchhoff's law of thermal emission

Analytical solutions in the one-plasmon approximation  
(including anti-resonant terms)

$$I_{\omega}^e = \epsilon_{\omega} I_{\omega}^{BB}(T)$$

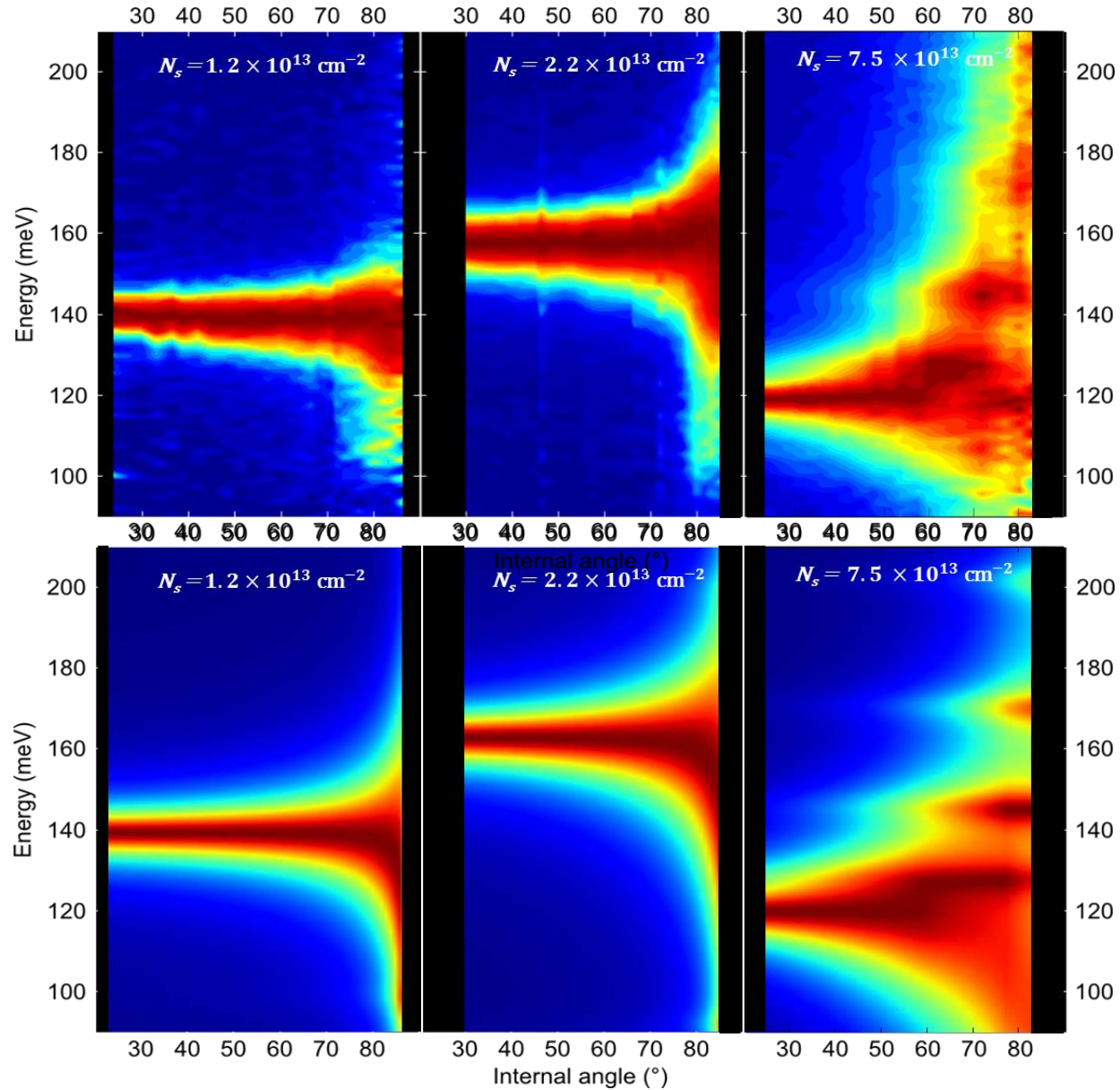
**Intensity of the plasmon emission**

**Blackbody intensity**

**Emissivity**  
=  
**Absorptivity**

$$\alpha_g(\omega) = \frac{4 \frac{4\omega_{MSP}^2}{(\omega_{MSP} + \omega)^2} \gamma \Gamma_g(\omega)}{(\omega_{MSP} - \omega)^2 + \frac{4\omega_{MSP}^2}{(\omega_{MSP} + \omega)^2} (\gamma + \Gamma_g(\omega))^2}$$

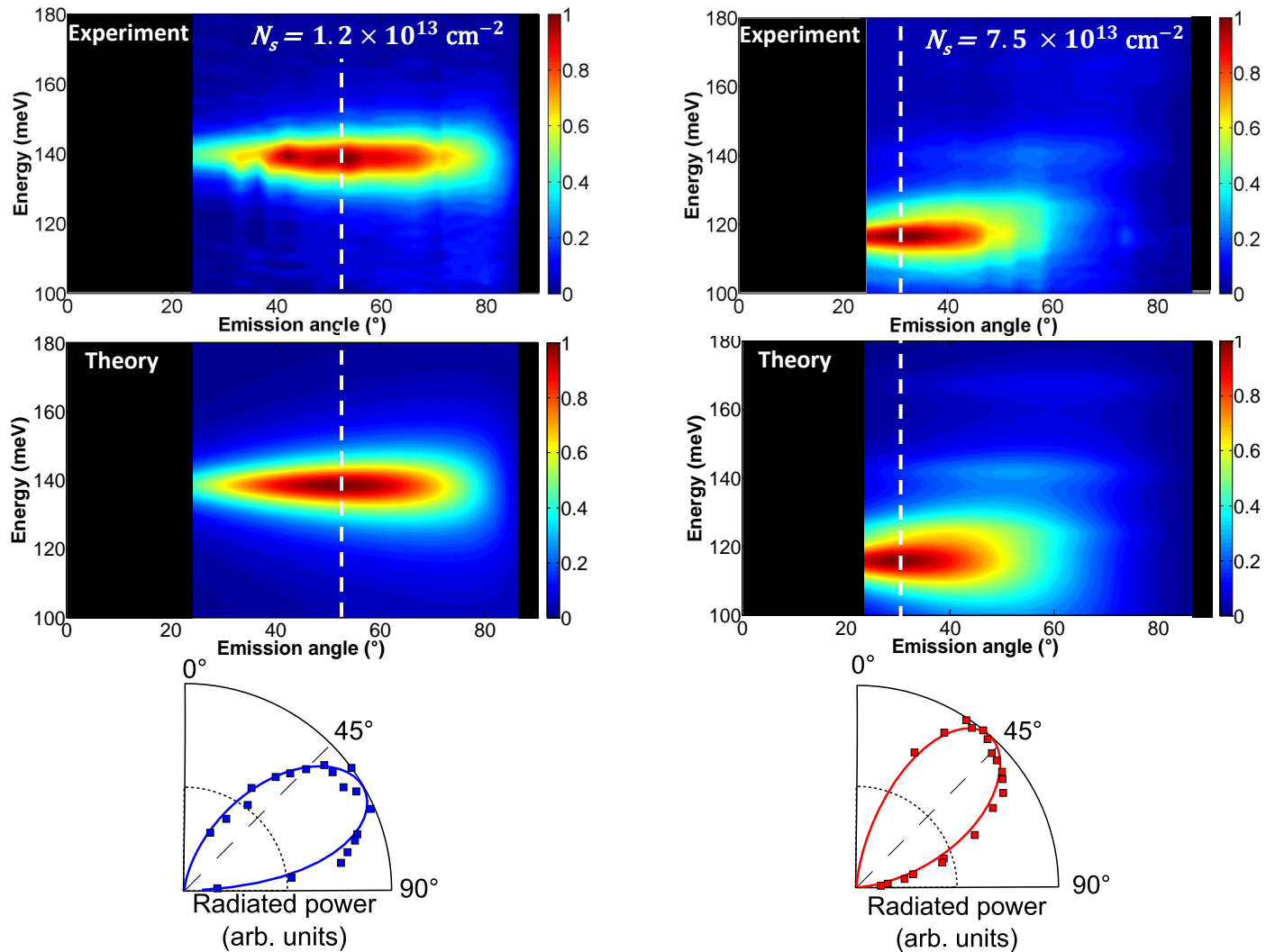
# Experiment vs theory



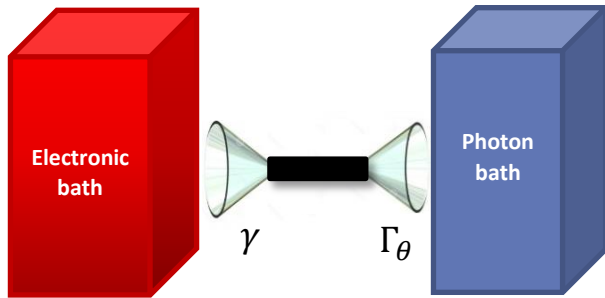


# Directional emission

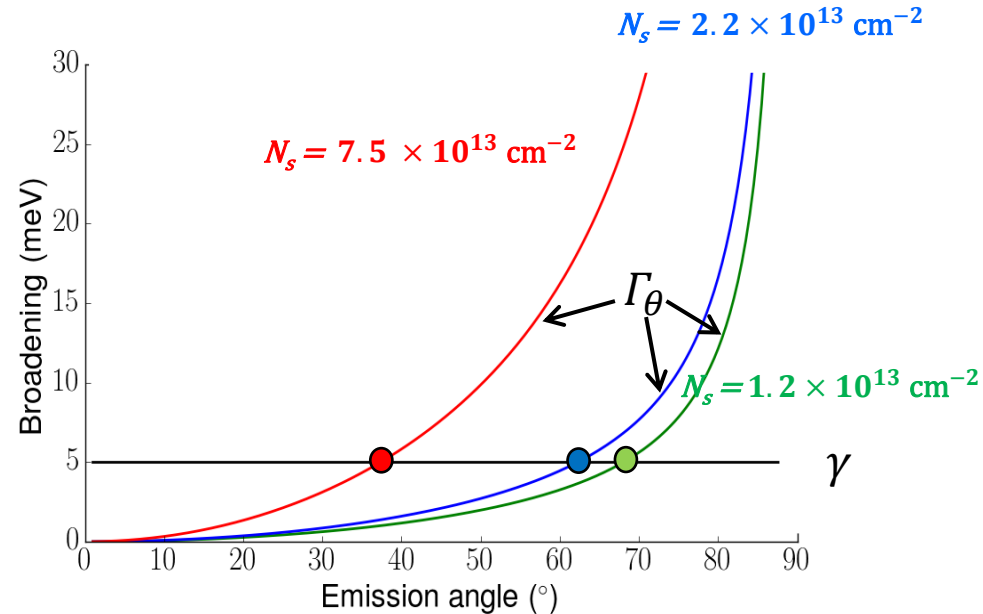
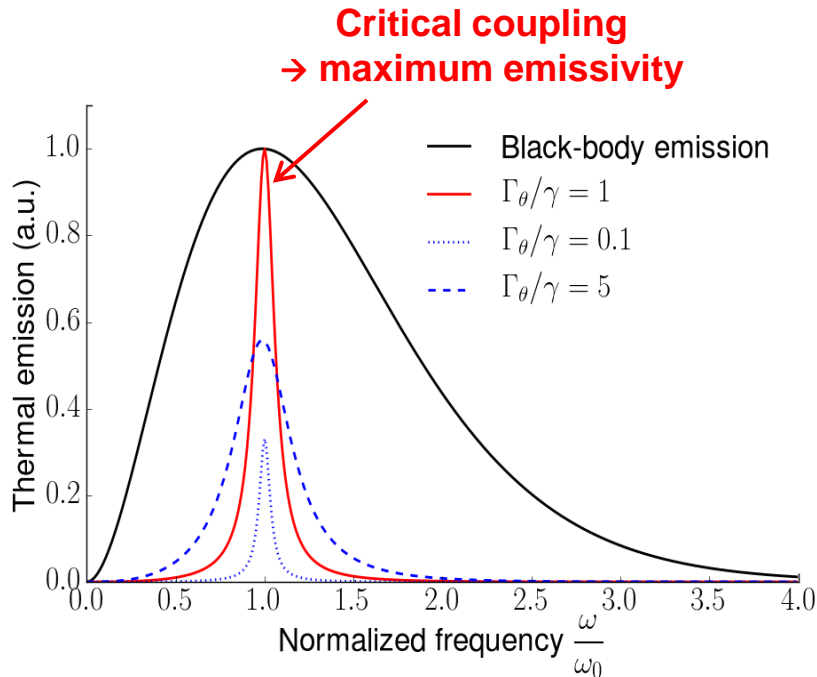
## Emission spectra (non-normalized)



# Critical coupling



Thermal emission is a twofold process:  
the limiting step depends on the ratio between  
radiative and non-radiative rates



The critical coupling condition is always  
met at a certain angle

# Conclusions



Emergence of multisubband plasmon modes in the optical spectrum of dense electron gases



Ultra-strong coupling regime with record values of the relative coupling at room temperature



Superradiant decay of the multisubband plasmon



Directional thermal emission

# Perspectives



## **Transport + quantum optics:**

- **Signatures of the light-matter coupling in the electronic transport (see G. Pupillo's talk)**
- **Engineer the resonant injection of plasmons/polaritons**



New devices based on quantum engineering of plasmons



Superradiance with spatially separated quantum wells



*Quantum optics with intersubband plasmons/polaritons?*