

# Entanglement and coherence in many-body dipolar systems

Susanne Yelin

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Efi Shahmoon, Elie Wolfe

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# Entanglement and coherence in many-body dipolar systems superradiant

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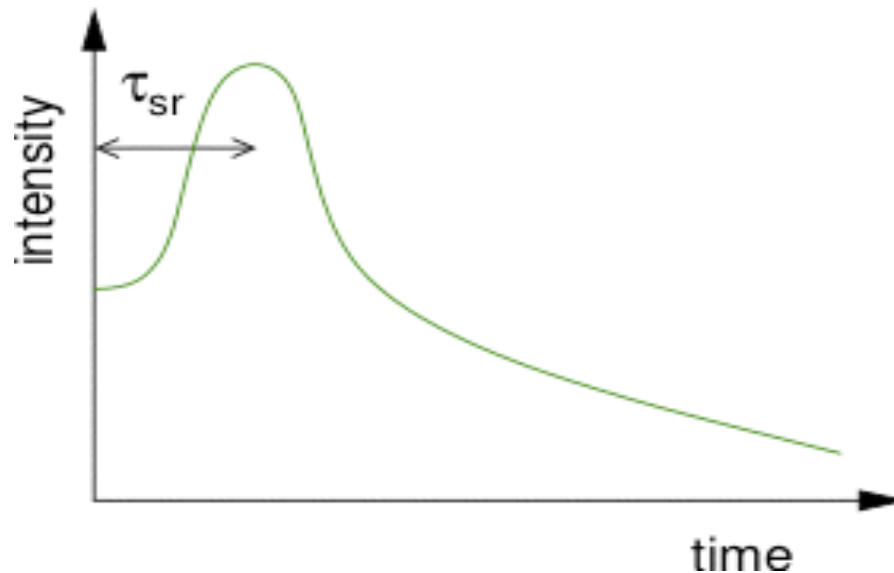
\$\$: ARO, NSF

# Atom-atom correlations in superradiance: Classic example



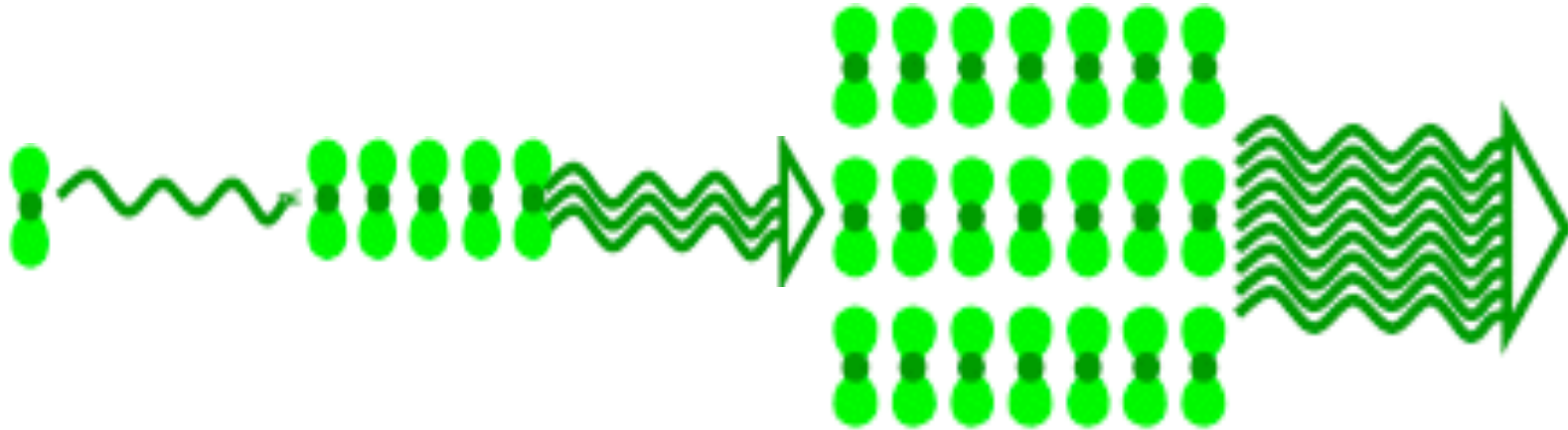
$\lambda$

- Superradiance



# What is **super** in superradiance?

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- Superradiance (superfluorescence)  $\propto N^2$   
(particle number)
- Build-up of collective dipoles

# What is “superradiance”?

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1. Everything that involves Dicke states
  - (e.g., collective  $\sqrt{N}$  effects,
  - bad-cavity limit,
  - ...)
2. Only systems involving cooperative (and nonlinear) effects
  - i.e., effect of exchange interaction
  - more than single excitation

only for purists

# Questions - guideline

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- Superradiance - What? Why?
- How do we calculate it (better)?
- Is there a collective (Lamb) shift?
- How does entanglement come into the picture?

# Questions - guideline

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- Superradiance - What? Why?

# so far...

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- ideal (Dicke)

PHYSICAL REVIEW

VOLUME 93, NUMBER 1

JANUARY 1, 1954

## Coherence in Spontaneous Radiation Processes

R. H. DICKE

*Palmer Physical Laboratory, Princeton University, Princeton, New Jersey*

(Received August 25, 1953)

By considering a radiating gas as a single quantum-mechanical system, energy levels corresponding to certain correlations between individual molecules are described. Spontaneous emission of radiation in a transition between two such levels leads to the emission of coherent radiation. The discussion is limited first to a gas of dimension small compared with a wavelength. Spontaneous radiation rates and natural line breadths are calculated. For a gas of large extent the effect of photon recoil momentum on coherence is calculated. The effect of a radiation pulse in exciting "super-radiant" states is discussed. The angular correlation between successive photons spontaneously emitted by a gas initially in thermal equilibrium is calculated.



# Dicke states

Fully symmetric state of  $n$  excitations in  $N$  particles, for example

$$|2\rangle_4 = \frac{1}{\sqrt{6}} (|1100\rangle + |1010\rangle + |1001\rangle + |0110\rangle + |0101\rangle + |0011\rangle)$$

$N$ -particle Dicke states decay with up to  $N^2$  speedup

# so far...

- ideal (Dicke)

- classical

PHYSICAL REVIEW A

VOLUME 2, NUMBER 3

SEPTEMBER 1970

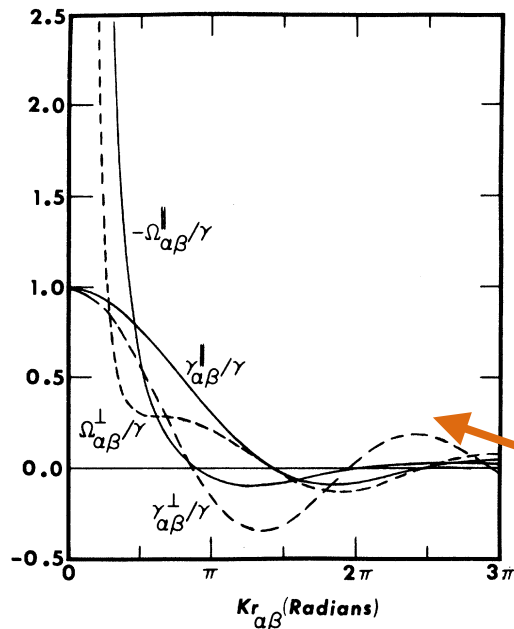
## Radiation from an $N$ -Atom System. I. General Formalism

R. H. Lehmburg

*U. S. Naval Air Development Center, Warminster, Pennsylvania*

(Received 19 November 1969)

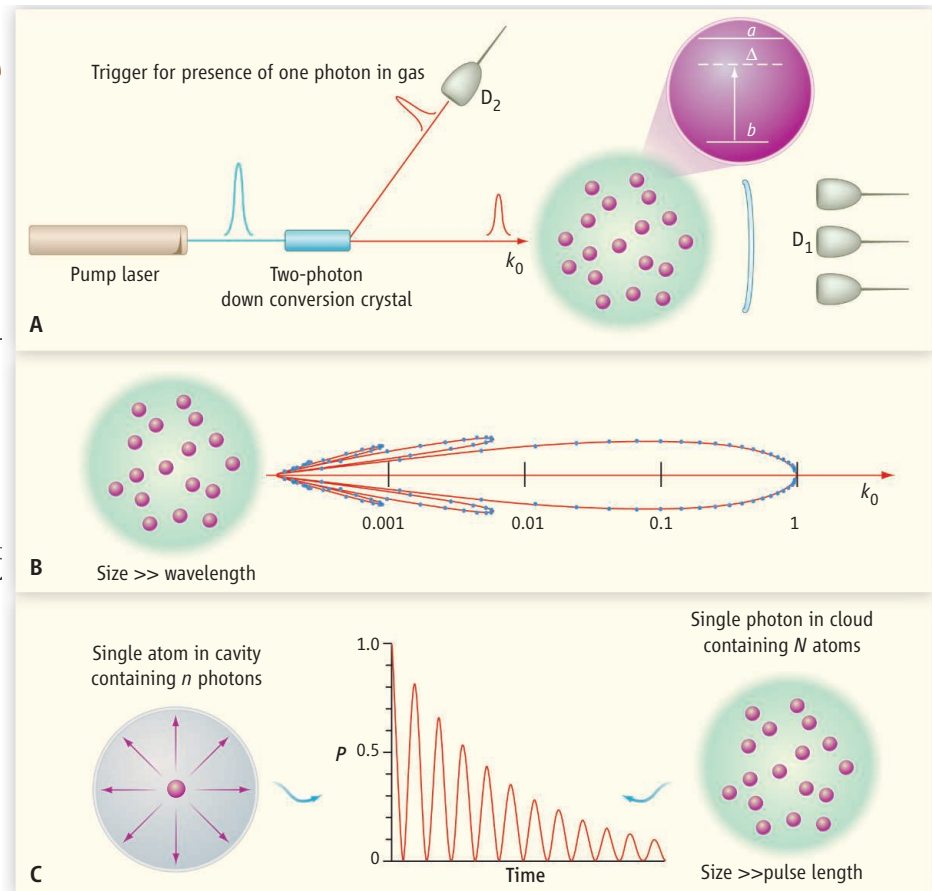
We consider the radiation from a system of  $N$  identical two-level atoms coupled to a continuum of quantized em modes, and possibly, to an external driving field near resonance. The atoms can be distributed over a region large in comparison to the resonant wavelength, but smaller than the spontaneous pulse length. Radiation rates and correlation functions are expressed in terms of expectation values of time-dependent atomic operators, which are shown to satisfy coupled first-order differential equations involving similar atomic operators and the initial radiation operators. The corresponding equations for the expectation values simplify considerably if no driving field is present. Similar results are derived for a model in which each atom is replaced by a harmonic oscillator.



collective shift

# so far...

- ideal (Dicke)
- classical (Friedberg et al.)
- single-photon (Svidzinsky et al.)



## PHYSICS

# The Super of Superradiance

Marlan O. Scully<sup>1,2</sup> and Anatoly A. Svidzinsky<sup>1</sup>

In 1954, Robert Dicke introduced the concept of superradiance in describing the emission of radiation from a collection of atoms. An even more interesting kind of superradiance can occur when a single photon is emitted from a cloud of atoms.

# so far...

- ideal (Dicke)
- classical (Friedberg, Sjoquist, Lamborg,...)
- single-photon
- linear (Javanainen)

in correlation order

Shifts of a Resonance Line in a Dense Atomic Sample

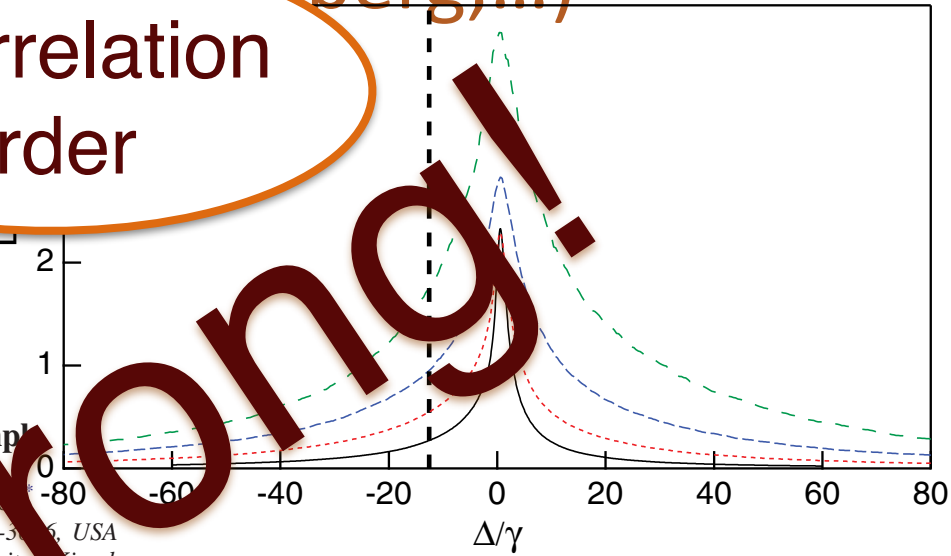
Juha Javanainen,<sup>1</sup> Janne Ruostekoski,<sup>2</sup> Yi Li,<sup>1</sup> and Sung-Min Yoon<sup>\*</sup>

<sup>1</sup>Department of Physics, University of Connecticut, Storrs, Connecticut 06269-3046, USA

<sup>2</sup>Mathematical Sciences, University of Southampton, Southampton SO9 7 1BJ, United Kingdom

(Received 28 August 2013; published 21 March 2014)

We study the collective response of a dense atomic sample to light, generally exactly using classical-electrodynamics simulations. In a homogeneously broadened atomic sample there is no overt Lorentz-Lorentz local field shift of the resonance, nor a collective Lamb shift. However, the addition of inhomogeneous broadening restores the usual mean-field phenomenology.



# so far...

- ideal (Dicke)
- classical (Friedberg et al., Lemberg,...)
- single-photon (Svidzinsky et al)
- linear (Javanainen)

## Shifts of a Resonance Line in a Dense Atomic Sample

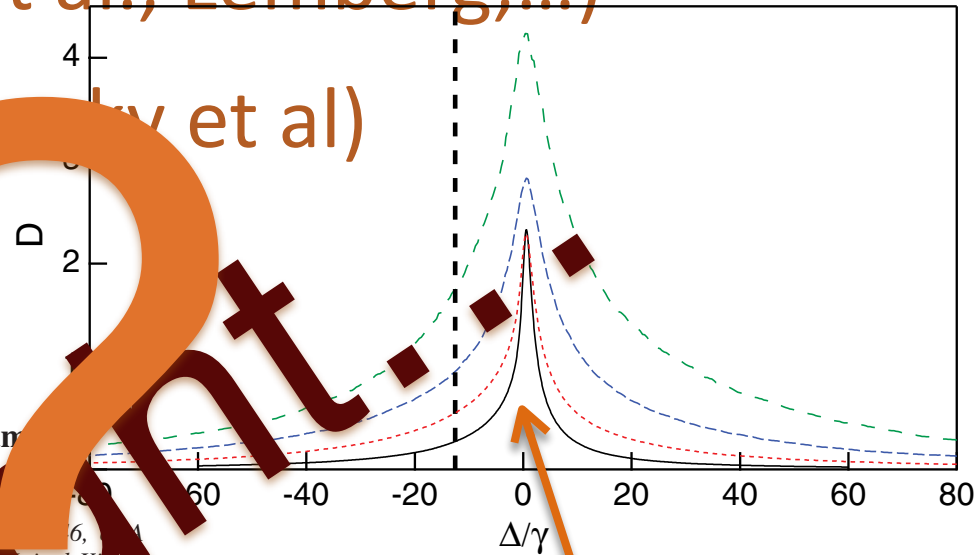
Juha Javanainen,<sup>1</sup> Janne Ruostekoski,<sup>2</sup> Yi Li,<sup>1</sup> and Wang-Ming

<sup>1</sup>Department of Physics, University of Connecticut, Storrs, Connecticut 06269, USA

<sup>2</sup>Mathematical Sciences, University of Southampton, Southampton SO1 8BJ, United Kingdom

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We study the collective response of a dense atomic sample to light using essentially classical-electrodynamics simulations. In a homogeneously broadened atomic sample there is no Lorentz-Lorentz local field shift of the resonance, nor a collective Lamb shift. However, the addition of inhomogeneous broadening restores the usual mean-field phenomenology.



no shift!

# Cooperative Lamb Shift in an Atomic Vapor Layer of Nanometer Thickness

J. Keaveney,<sup>1</sup> A. Sargsyan,<sup>2</sup> U. Krohn,<sup>1</sup> I. G. Hughes,<sup>1</sup> D. Sarkisyan,<sup>2</sup> and C. S. Adams<sup>1,\*</sup>

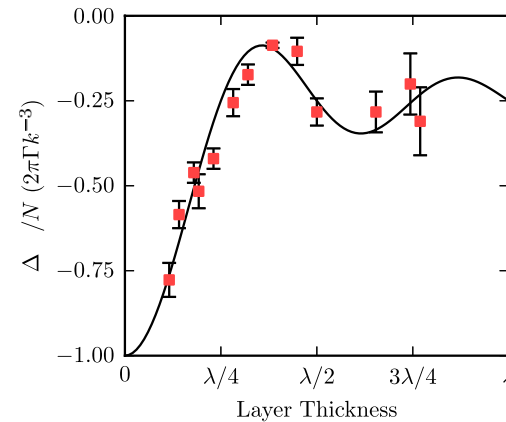
<sup>1</sup>Department of Physics, Rochester Building, Durham University, South Road, Durham DH1 3LE, United Kingdom

<sup>2</sup>Institute for Physical Research, National Academy of Sciences—Ashtarak 2, 0203, Armenia

(Received 25 January 2012; published 23 April 2012)

We present an experimental measurement of the cooperative Lamb shift and the Lorentz shift using a nanothickness atomic vapor layer with tunable thickness and atomic density. The cooperative Lamb shift

- ideal (Dicke)
- classical
- single-photon
- linear (Jaynes-Cummings)
- dilute (Adams, Ye,...)



# so far...

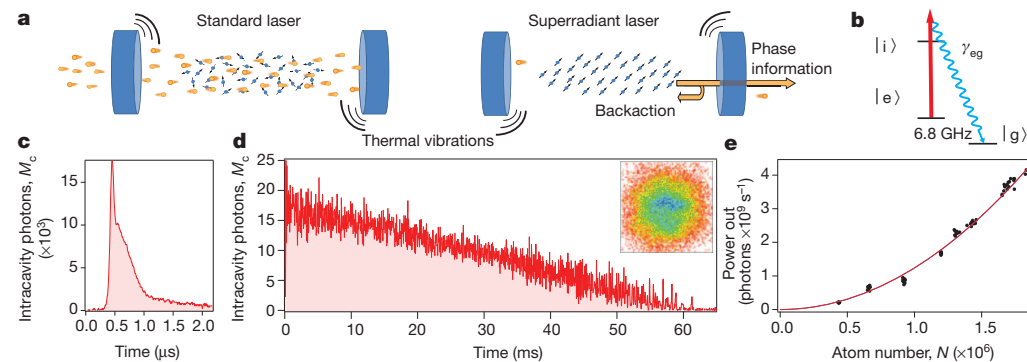
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- ideal (Dicke)
- classical (Friedberg et al., Lemberg,...)
- single-photon (Svidzinsky et al)
- linear (Javanainen)
- dilute **Collective atomic emission and motional effects in a dense coherent medium**

S. L. Bromley<sup>1</sup>, B. Zhu<sup>1</sup>, M. Bishof<sup>1</sup>, X. Zhang<sup>1</sup>, T. Bothwell<sup>1</sup>, J. Schachenmayer<sup>1</sup>, T. L. Nicholson<sup>1</sup>,  
R. Kaiser<sup>2</sup>, S. F. Yelin<sup>3</sup>, M. D. Lukin<sup>4</sup>, A.M. Rey<sup>1</sup>, & J. Ye<sup>1</sup>

# A steady-state superradiant laser with less than one intracavity photon

Justin G. Bohnet<sup>1</sup>, Zilong Chen<sup>1</sup>, Joshua M. Weiner<sup>1</sup>, Dominic Meiser<sup>1†</sup>, Murray J. Holland<sup>1</sup> & James K. Thompson<sup>1</sup>



- ideal (
- classic
- single.
- linear
- dilute (Adams, Ye,...)
- very low excitation (Holland/Thompson...)



# What is “superradiance”?

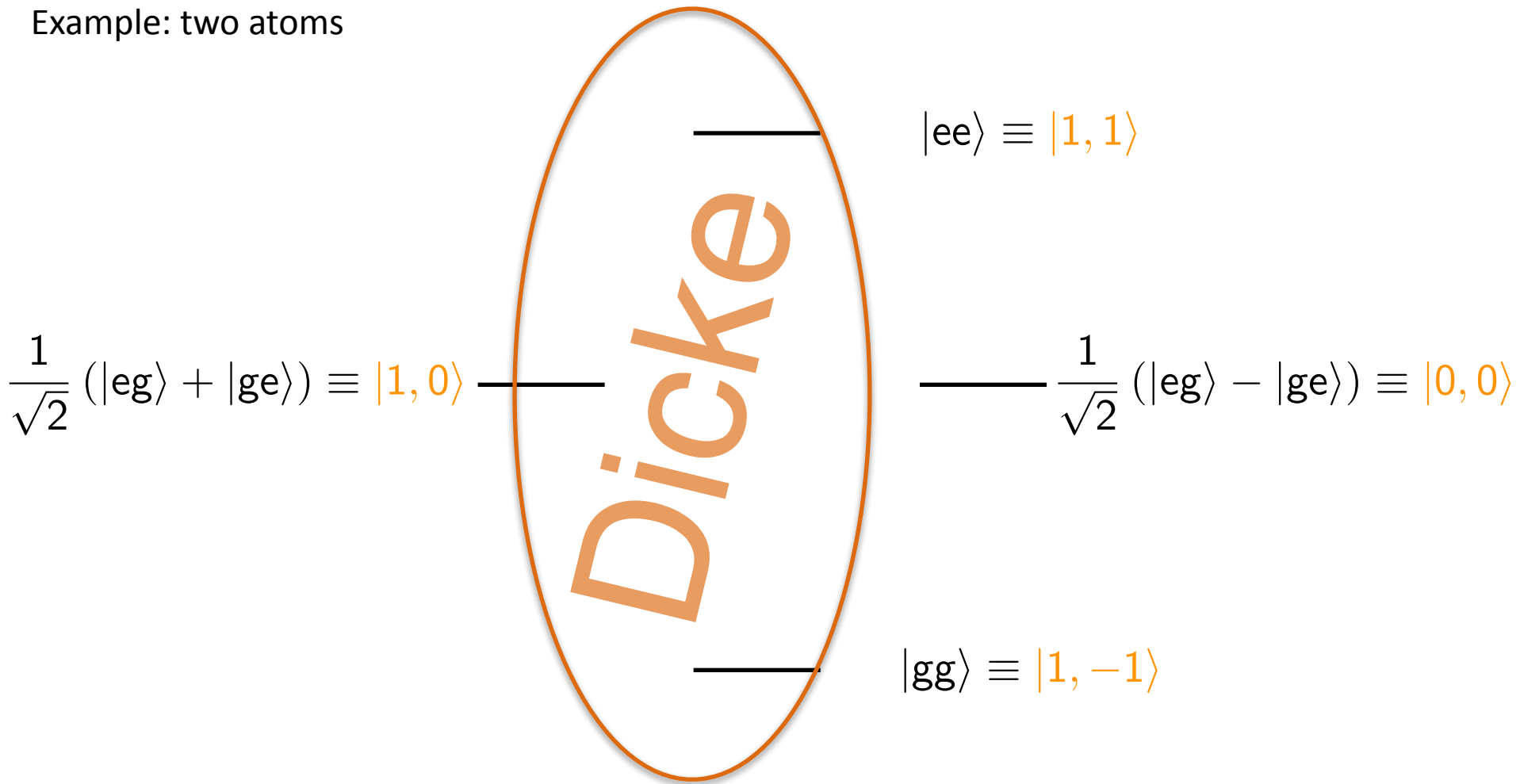
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  - bad-cavity limit,
  - ...)
2. Only systems involving cooperative (and nonlinear) effects
  - i.e., effect of exchange interaction
  - more than single excitation

only for purists

# Cooperativity

Example: two atoms



# Cooperativity

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Example: two atoms

$$\text{————} \quad |ee\rangle \equiv |1, 1\rangle$$

$$\frac{1}{\sqrt{2}} (|eg\rangle + |ge\rangle) \equiv |1, 0\rangle \text{————}$$

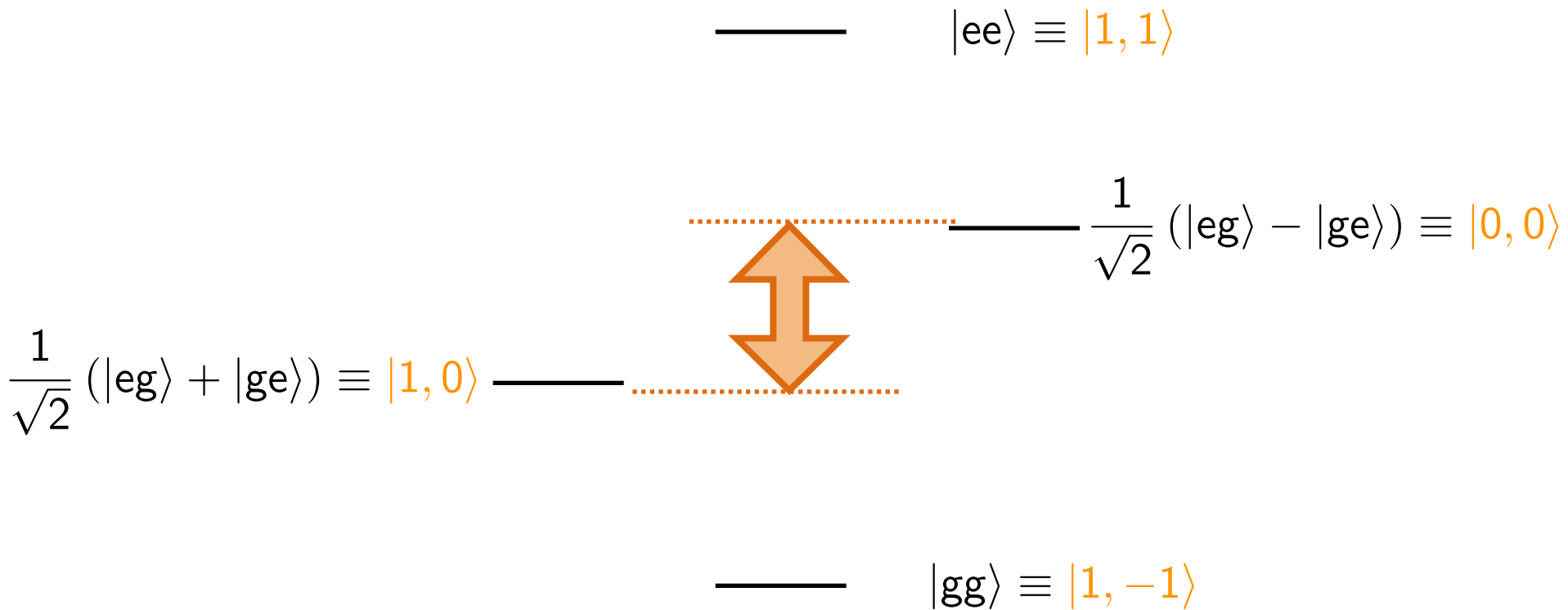
$$\text{————} \frac{1}{\sqrt{2}} (|eg\rangle - |ge\rangle) \equiv |0, 0\rangle$$

$$\text{————} \quad |gg\rangle \equiv |1, -1\rangle$$

(anti)symmetry costs energy: “exchange interaction”

# Cooperativity

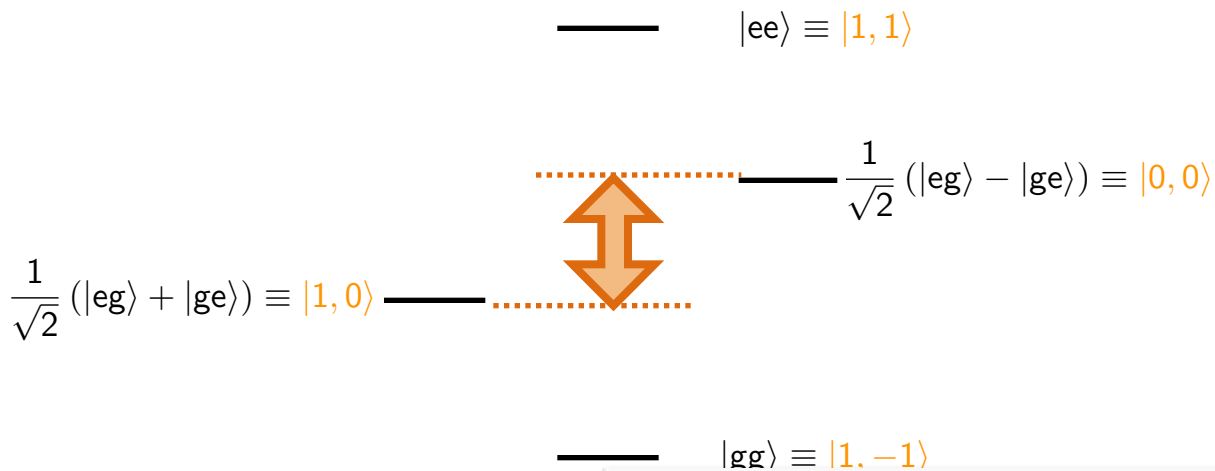
Example: two atoms



(anti)symmetry costs energy: “exchange interaction”

# Cooperativity

Example: two atoms

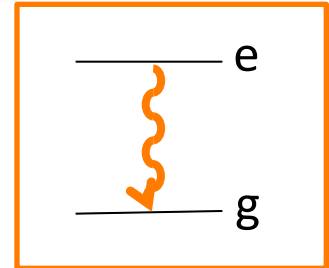


exchange interaction:

- usually dipole-dipole mediated
- creates shift and broadening (Kramers-Kronig)

# Dicke model

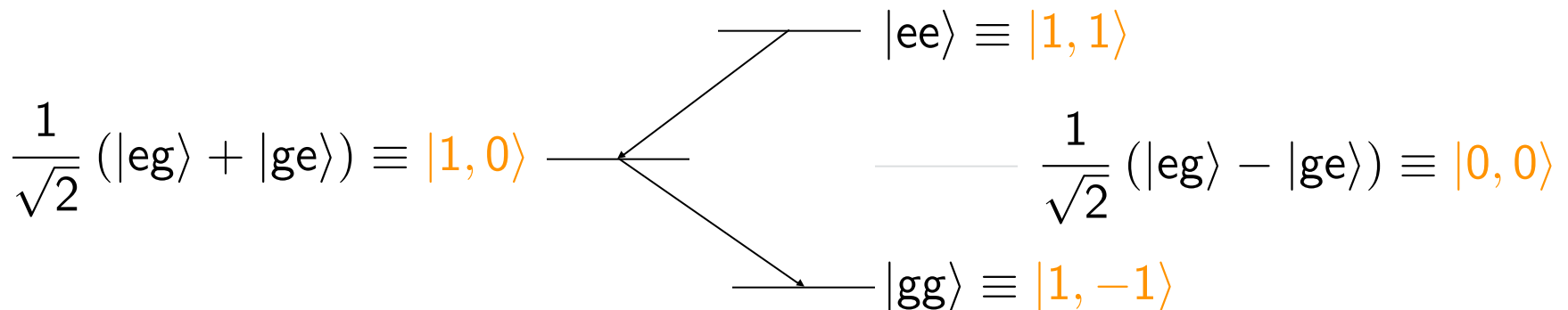
Dicke states:  $|J, M_J\rangle = \text{Sym} \left| \underbrace{e \dots e}_{J+M_J} \underbrace{g \dots g}_{J-M_J} \right\rangle$



angular momentum formulation  $\Rightarrow J \leq \frac{N}{2}$  (= max. excitation)  
 $M_J = -J \dots J$  (= actual excitation)

Radiation couples only states with equal  $J$

Example: two atoms



# Questions - guideline

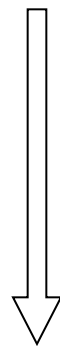
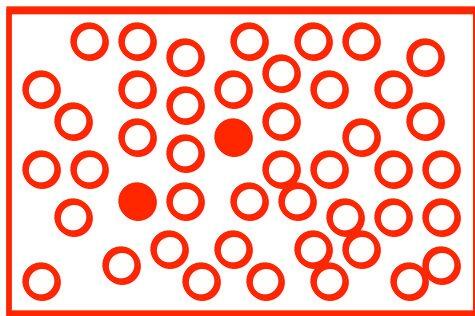
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- Superradiance - What? Why?
- How do we calculate it (better)?

# Dynamics of atoms in dense media - Schwinger-Keldysh & Dyson Eq.

Full dynamics (all degrees of freedom of atoms, fields)

$$H = H_{\text{atoms}} + H_{\text{field}} - \sum \mathbf{p}_i \mathbf{E}_i$$



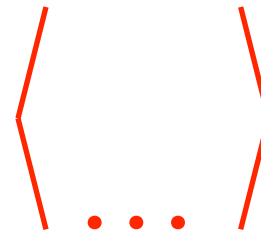
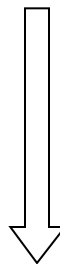
two probe atoms  
+  
surrounding atoms

Two atoms + field

$$V_{\text{probe}} = \sum_{i=1,2} \mathbf{p}_i \mathbf{E}_i \Rightarrow S = T e^{-\frac{i}{\hbar} \int d\tau V_{\text{probe}}(\tau)}$$

- $\langle e^{sX} \rangle = e^{\sum \frac{s^m}{m!} \langle\langle X^m \rangle\rangle}$

- Gauss:  $m \leq 2$



field degrees  
of freedom

effective two-atom description

Two atom Master equation



# Equations of motion

---

$$\dot{a} = \Gamma - (\gamma + 2\Gamma) a - \bar{\gamma} x - i \Omega (\rho_{eg} - \rho_{ge})$$

$$\dot{n} = 2\gamma - 2(\gamma + 2\Gamma) n - 4\gamma a + 4(\bar{\gamma} + 2\bar{\Gamma}) x - 4i \Omega (m_{eg} - m_{ge}) - 4i \mathcal{C}\gamma (\rho_{ge} m_{eg} - \rho_{eg} m_{ge})$$

$$\dot{x} = -(\gamma + 2\Gamma) x + \frac{\bar{\gamma} + 2\bar{\Gamma}}{2} n + \bar{\gamma} a - \frac{\bar{\gamma}}{2} + i \Omega (m_{eg} - m_{ge}) + i \mathcal{C}\gamma (\rho_{ge} m_{eg} - \rho_{eg} m_{ge})$$

$$\dot{\rho}_{eg} = - \left( \frac{\gamma + 2\Gamma}{2} + i(\delta + 2\Delta - \Delta_\Omega) \right) \rho_{eg} + \frac{\bar{\gamma} - 2i\bar{\delta}}{2} m_{eg} - i(\Omega + \mathcal{C}\gamma \rho_{eg}) (2a - 1)$$

$$\dot{m}_{eg} = - \left( 3 \frac{\gamma + 2\Gamma}{2} + \bar{\gamma} + 2\bar{\Gamma} + i(\delta + 2\Delta - \Delta_\Omega) \right) m_{eg} - \left( \gamma + \frac{\bar{\gamma}}{2} + i\bar{\delta} \right) \rho_{eg} \\ - i(\Omega + \mathcal{C}\gamma \rho_{eg}) (n - 2x) - 2i(\Omega + \mathcal{C}\gamma \rho_{ge}) \rho_{eg,eg}$$

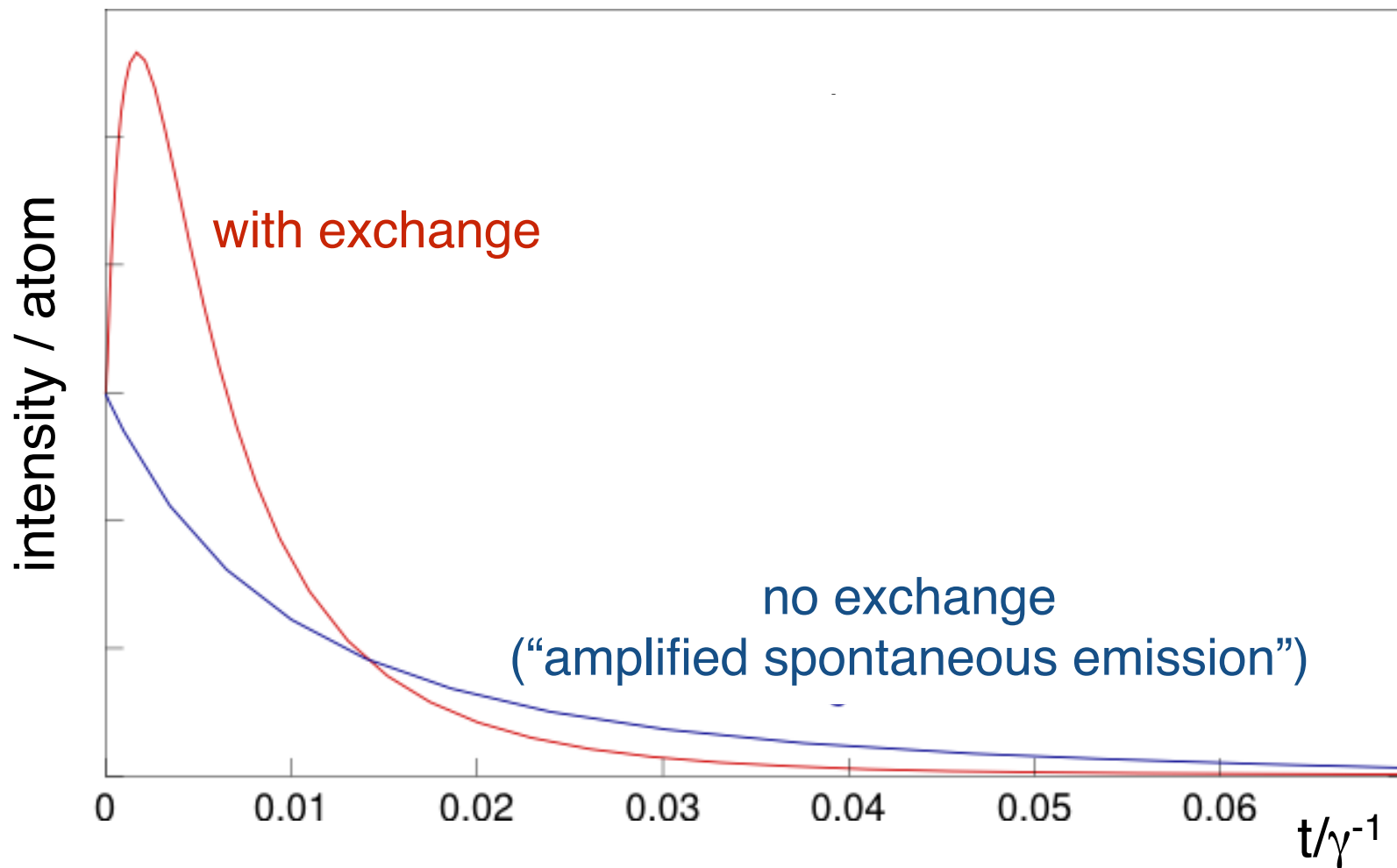
$$\dot{\rho}_{eg,eg} = - ((\gamma + 2\Gamma) + 2i(\delta + 2\Delta - \Delta_\Omega)) \rho_{eg,eg} - 2i(\Omega + \mathcal{C}\gamma \rho_{eg}) m_{eg}$$

$$\begin{aligned}
 \dot{a} &= \Gamma - (\gamma + 2\Gamma) a - \bar{\gamma} x - i \Omega (\rho_{eg} - \\
 \dot{n} &= \Gamma - (\gamma + 2\Gamma) n - 4\gamma a + 4(\bar{\gamma} + \\
 \dot{x} &= \Gamma - \frac{\bar{\gamma} + 2\bar{\Gamma}}{2} n + \bar{\gamma} a - \\
 \dot{\rho}_{eg} &= \Gamma - \frac{\gamma + 2\Gamma}{2} + i (\delta + 2\Delta - \Delta_{\Omega}) \\
 \dot{m}_{eg} &= - \left( 3 \frac{\gamma + 2\Gamma}{2} + \bar{\gamma} + 2\bar{\Gamma} + i (\delta + 2\Delta - \Delta_{\Omega}) \right) \\
 &\quad - i (\Omega + \mathcal{C} \gamma \rho_{eg}) (n - 2x) - 2i (\Omega \\
 \dot{\rho}_{eg, eg} &= - ((\gamma + 2\Gamma) + 2i (\delta + 2\Delta - \Delta_{\Omega}))
 \end{aligned}$$

average  
excited state

effective  
inversion

exchange  
term



$$\dot{a} = \Gamma - (\gamma + 2\Gamma) a - \bar{\gamma} x - i \Omega (\rho_{eg} -$$

$$\dot{n} = 2\gamma - 2(\gamma + 2\Gamma) n - 4\gamma a + 4(\bar{\gamma} +$$

$$\dot{x} = (\gamma + 2\Gamma) x + \frac{\bar{\gamma} + 2\bar{\Gamma}}{2} n + \bar{\gamma} a -$$

$$\dot{\rho}_{eg} = - \left( \frac{\gamma + 2\Gamma}{2} + i (\delta + 2\Delta - \Delta_{\Omega}) \right) \rho_{eg}$$

$$\dot{m}_{eg} = - \left( 3 \frac{\gamma + 2\Gamma}{2} + \bar{\gamma} + 2\bar{\Gamma} + i (\delta + 2\Delta - \Delta_{\Omega}) \right) m_{eg}$$

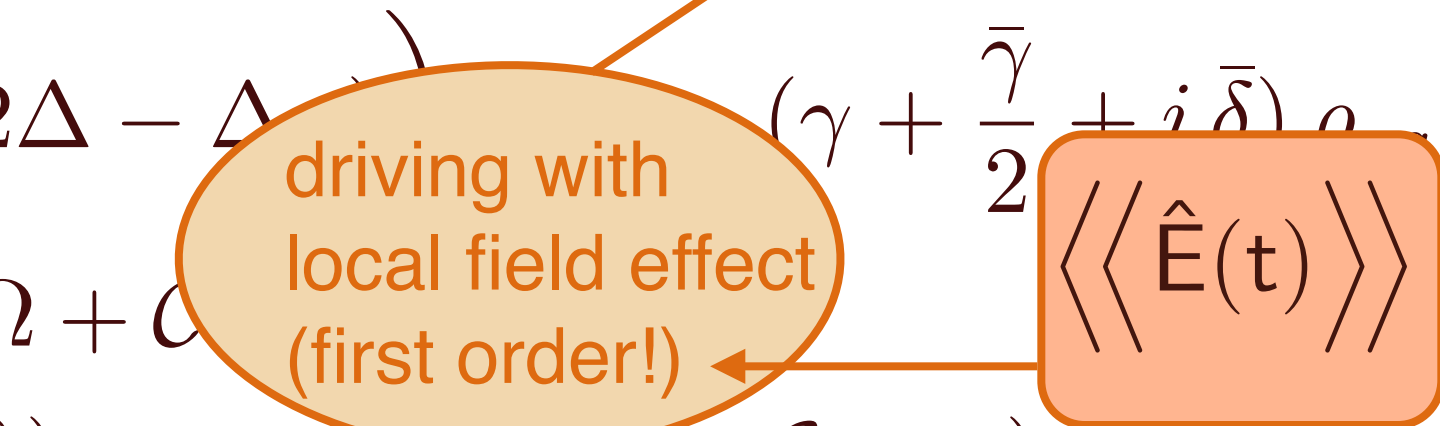
$$- i (\Omega + \mathcal{C} \gamma \rho_{eg}) (n - 2x) - 2i (\Omega + \mathcal{C} \gamma \rho_{eg}) x$$

$$\dot{\rho}_{eg, eg} = - ((\gamma + 2\Gamma) + 2i (\delta + 2\Delta - \Delta_{\Omega})) \rho_{eg, eg}$$

$$- 2\bar{\Gamma}) x - 4i\Omega(m_{eg} - m_{ge}) - 4i\mathcal{C}\gamma(\rho_{ge}m_{eg} - \rho_{eg}m_{ge})$$

$$- \frac{\bar{\gamma}}{2} + i\Omega(m_{eg} - m_{ge}) + i\mathcal{C}\gamma(\rho_{ge}m_{eg} - \rho_{eg}m_{ge})$$

$$\rho_{eg} + \frac{\bar{\gamma} - 2i\bar{\delta}}{2} m_{eg} - i(\Omega + \mathcal{C}\gamma\rho_{eg})(2a - 1)$$



$$\rho_{eg, eg} - 2i(\Omega + \mathcal{C}\gamma\rho_{eg})m_{eg}$$

$$\begin{aligned}
& \gamma - 2(\gamma + 2\Gamma) n - 4\gamma a + 4(\bar{\gamma} + 2\bar{\Gamma}) x - 4i\Omega (m_{eg} - n) \\
& (\gamma + 2\Gamma) x + \frac{\bar{\gamma} + 2\bar{\Gamma}}{2} n + \bar{\gamma} a - \frac{\bar{\gamma}}{2} + i\Omega (m_{eg} - n) \\
& \left( \frac{\gamma + 2\Gamma}{2} + i(\delta + 2\Delta - \Delta_\Omega) \right) \rho_{eg} + \frac{\bar{\gamma} - 2i\bar{\delta}}{2} m_{eg} \\
& \left( 3\frac{\gamma + 2\Gamma}{2} + 2\bar{\Gamma} + i(\delta + 2\Delta - \Delta_\Omega) \right) m_{eg} - \\
& i(\Omega - \dots) \left( \frac{\gamma + 2\Gamma}{2} + i(\delta + 2\Delta - \Delta_\Omega) \right) \rho_{eg, eg} - 2i(\Omega - \dots)
\end{aligned}$$

superradiance rate and shifts (second order!)

$$\langle\langle \hat{E}_1(t_1) \hat{E}_2(t_2) \rangle\rangle_{,eg}$$


# Can one expect superradiance?

B

... and Chirping

**Dynamics of atoms in dense media – Schwinger-Keldysh & Dyson Eq.**

Full dynamics (all degrees of freedom of atoms, field)

$$H = H_{atoms} + H_{field} = \sum p_i E_i$$


Two probe atoms + surroundings

Two atoms + field

$$V_{probe} = \sum p_i E_i \Rightarrow S = S_0^{-1} T^{\dagger} V_{probe} S_0^{-1}$$

effective two-atom description  
Two-atom Master equation

**Symmetric small sample approximation**

$$\dot{\rho} = - (2\Gamma + \gamma) \rho + \Gamma$$

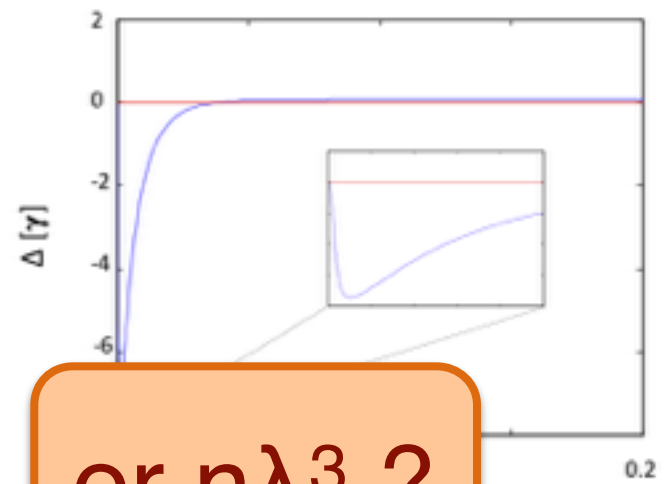
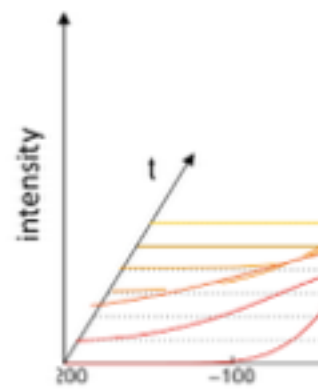
$$\dot{\rho} = -2(2\Gamma + \gamma) \rho - 2\gamma(2n - 1) + 8\Gamma \rho$$

$$\dot{\rho} = - (2\Gamma + \gamma) \rho + \Gamma \rho$$

where

- $\rho$ : average excited state population
- $\rho$ : two-atom "coherence"
- $\rho$ : correlation

$$S(\omega) = FT$$



or  $n\lambda^3$  ?

The impor

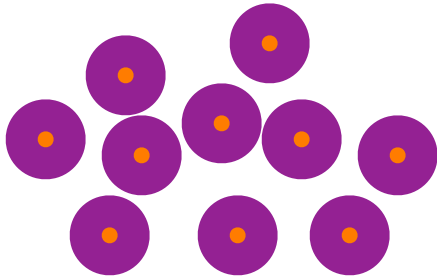
$$n\lambda^3 r$$

optical depth

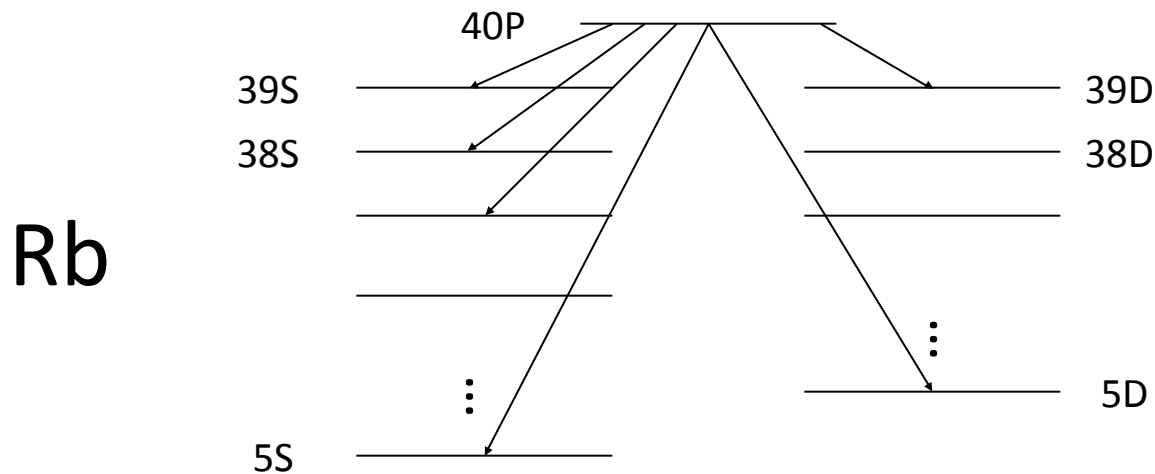
$n$ : density,  $\lambda$ : wavelength,  $r$ : system size

# New experimental systems: example

- Ultracold Rydberg atoms

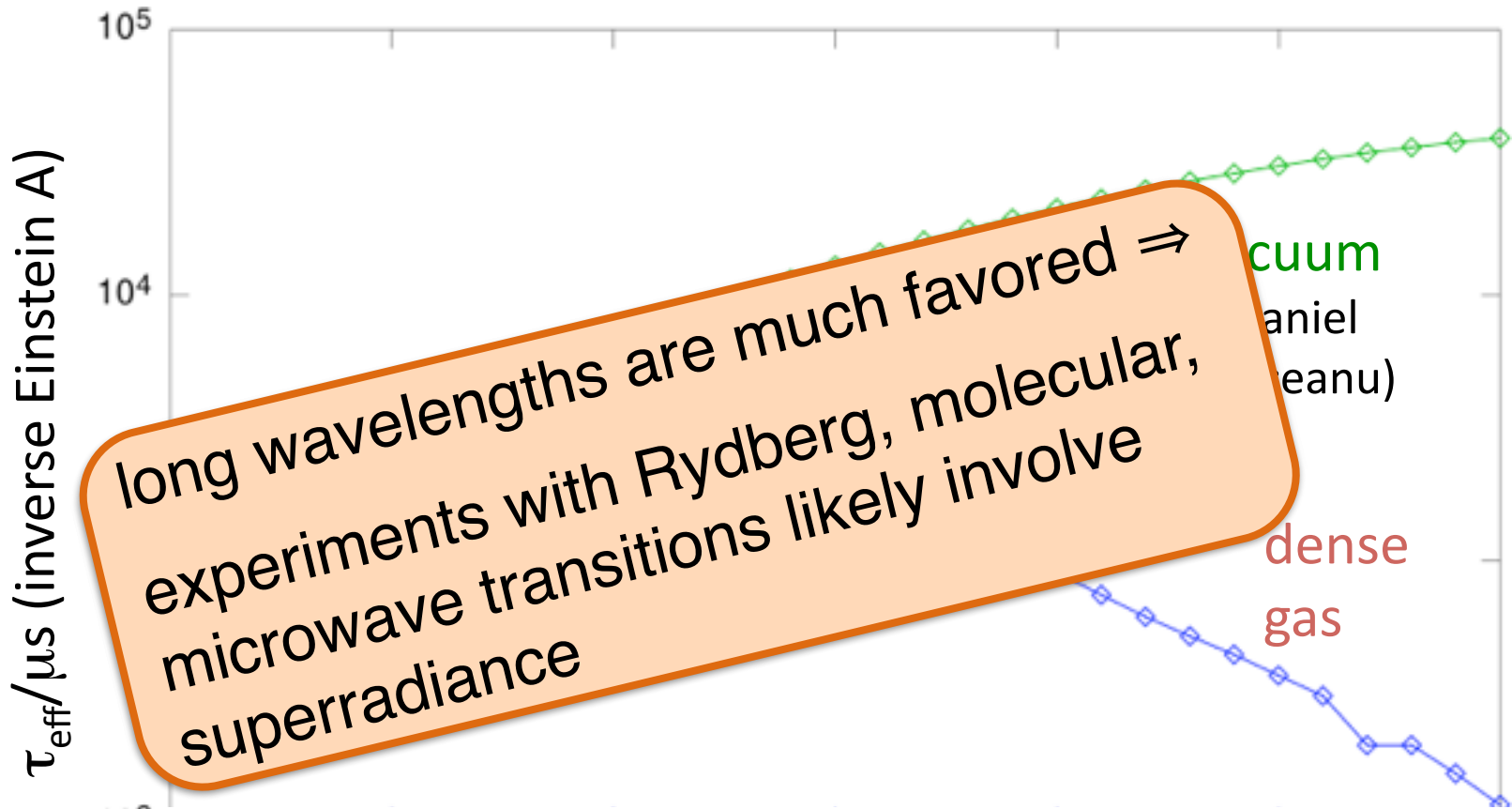


(Phil Gould, Ed Eyler, Uconn)





# Effective decay times from 40P into nS

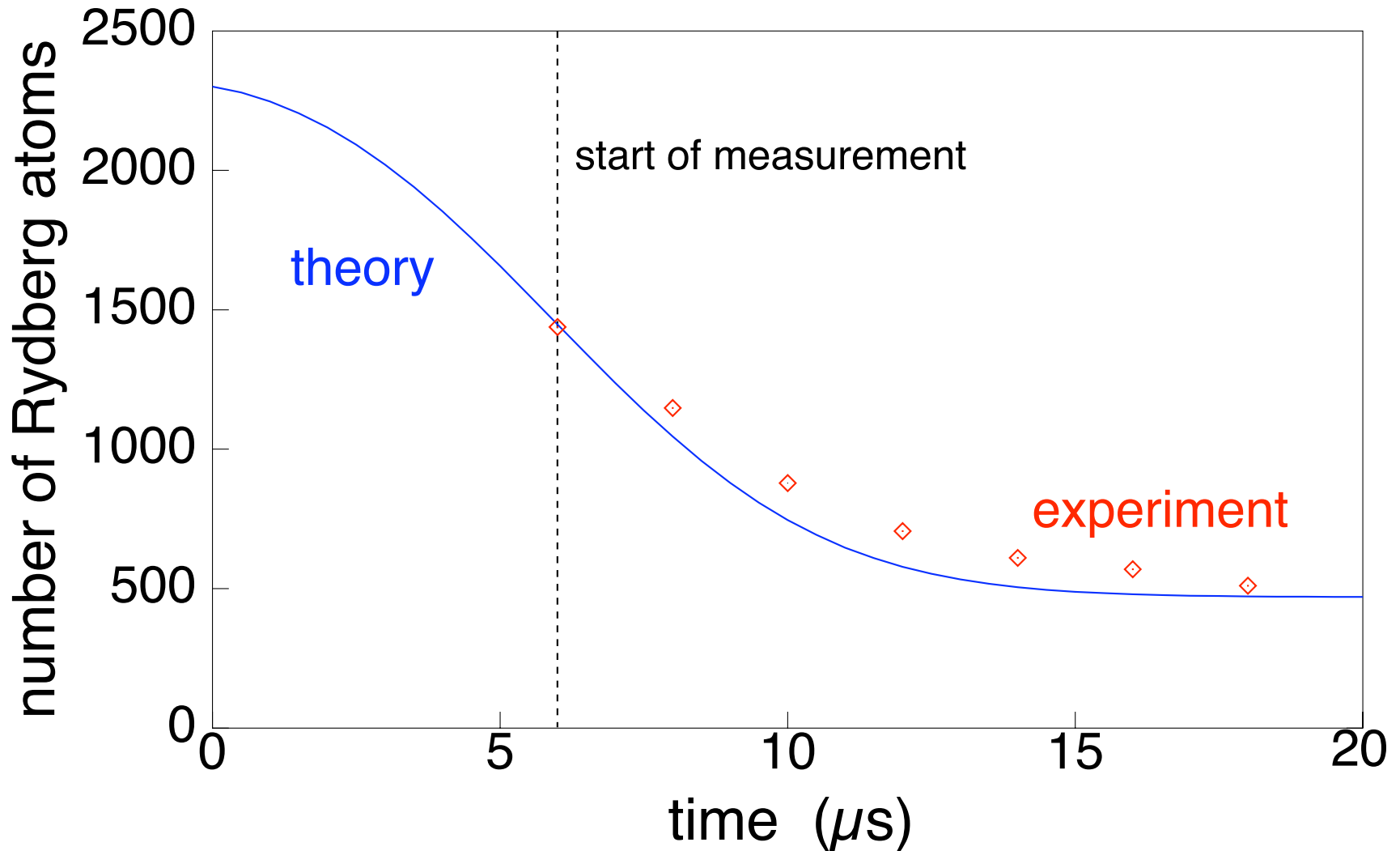


In vacuum: decay into low n is favored

In dense gas: decay into high n is favored  $\Rightarrow \lambda$  large,  $n \lambda^2 r$  large!

**superradiant decay!**

# Superradiance in Rydberg systems



# Questions - guideline

---

- Superradiance - What? Why?
- How do we calculate it (better)?
- Is there a collective (Lamb) shift?

# Collective Lamb shift

---

- “Lamb shift” is the result of interaction with the vacuum fluctuations
- In the case of altered density of states of the “vacuum” (i.e., the surrounding space), the value of the shift changes
- With a high (superradiant) density of radiators, the density of states inside the medium can be considerably altered



“Collective Lamb shift”

# Collective Shift

has **spontaneous** part....

$$\langle\langle E^- E^+ \rangle\rangle \approx \langle\langle a^\dagger a \rangle\rangle = n$$

dependent on number of photons

$$\langle [E^-, E^+] \rangle \propto \langle [a, a^\dagger] \rangle = 1$$

independent on number of photons

$$\Delta_{\text{stim}}^{(ij)} = \frac{1}{2\pi} \mathcal{P} \int d\omega' \frac{\Gamma_{ij}(\omega')}{\omega - \omega'}$$

# Collective Shift

has **spontaneous** part....

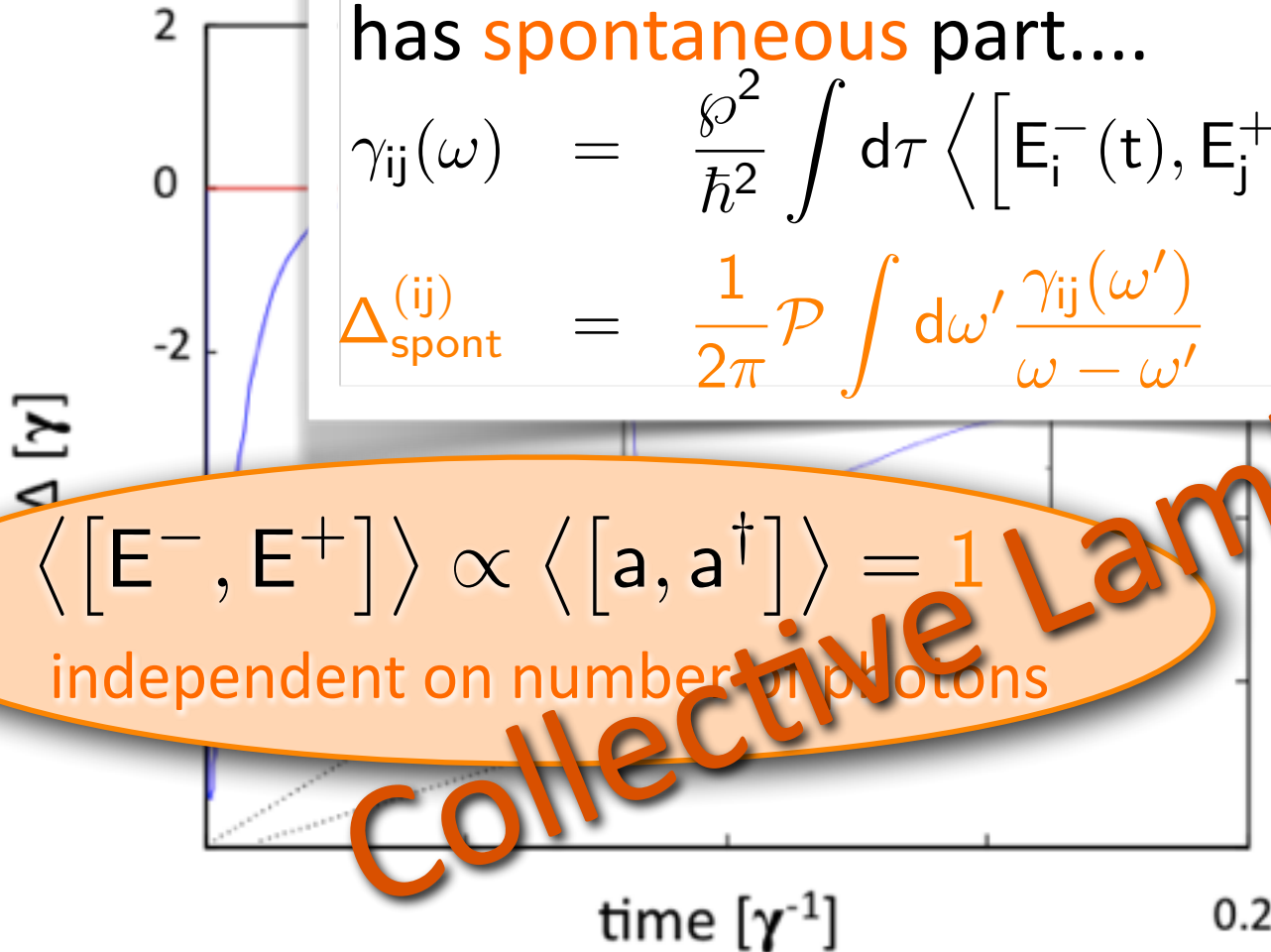
$$\gamma_{ij}(\omega) = \frac{\rho^2}{\hbar^2} \int d\tau \langle [E_i^-(t), E_j^+(t + \tau)] \rangle e^{i\omega\tau}$$

$$\Delta_{\text{spont}}^{(ij)} = \frac{1}{2\pi} \mathcal{P} \int d\omega' \frac{\gamma_{ij}(\omega')}{\omega - \omega'}$$

$$\langle [E^-, E^+] \rangle \propto \langle [a, a^\dagger] \rangle = 1$$

independent on number of photons

collective Lamb Shift



# Collective Lamb Shift in the low-excitation limit ( $\propto \Omega^2$ )

$$\delta_{\text{coll. Lamb}} = \frac{\cos k_0 r - e^{-\tilde{\gamma} k_0 r} \cos(k_0 r + \tilde{\delta} k_0 r)}{r}$$

with

$$k_0 = \frac{\omega}{c}$$

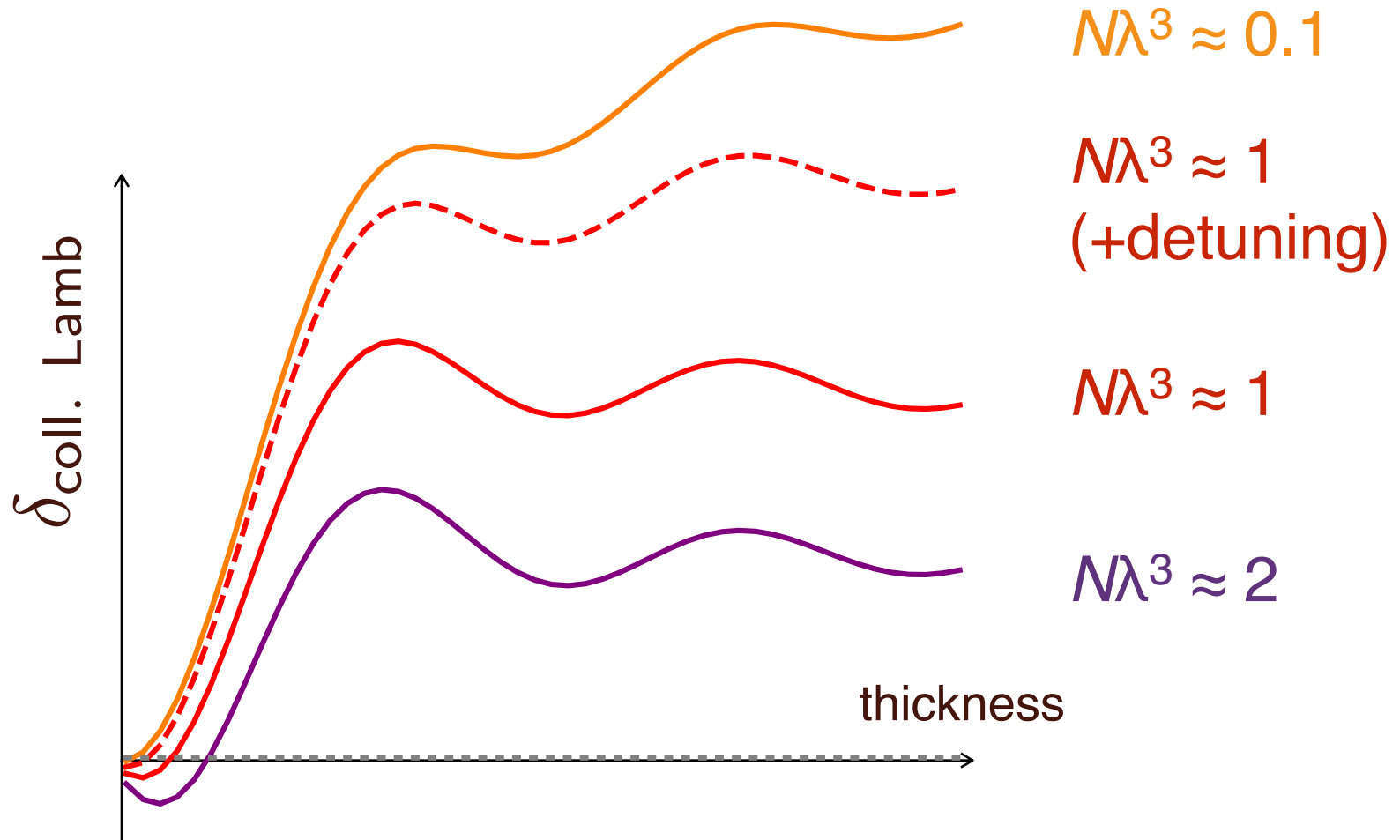
$$\tilde{\gamma} = \mathcal{C} \frac{\frac{\gamma}{2}}{\left(\frac{\gamma}{2}\right)^2 + (\Delta + \mathcal{C}\gamma - \delta_{\text{coll. Lamb}})^2}$$

$$\tilde{\delta} = \mathcal{C} \frac{\Delta + \mathcal{C}\gamma - \delta_{\text{coll. Lamb}}}{\left(\frac{\gamma}{2}\right)^2 + (\Delta + \mathcal{C}\gamma - \delta_{\text{coll. Lamb}})^2}$$

$\mathcal{C}$  = volume per  
cubic  
wavelength

solve self-consistently, normalized vs. (vacuum) Lamb shift

# Collective Lamb Shift

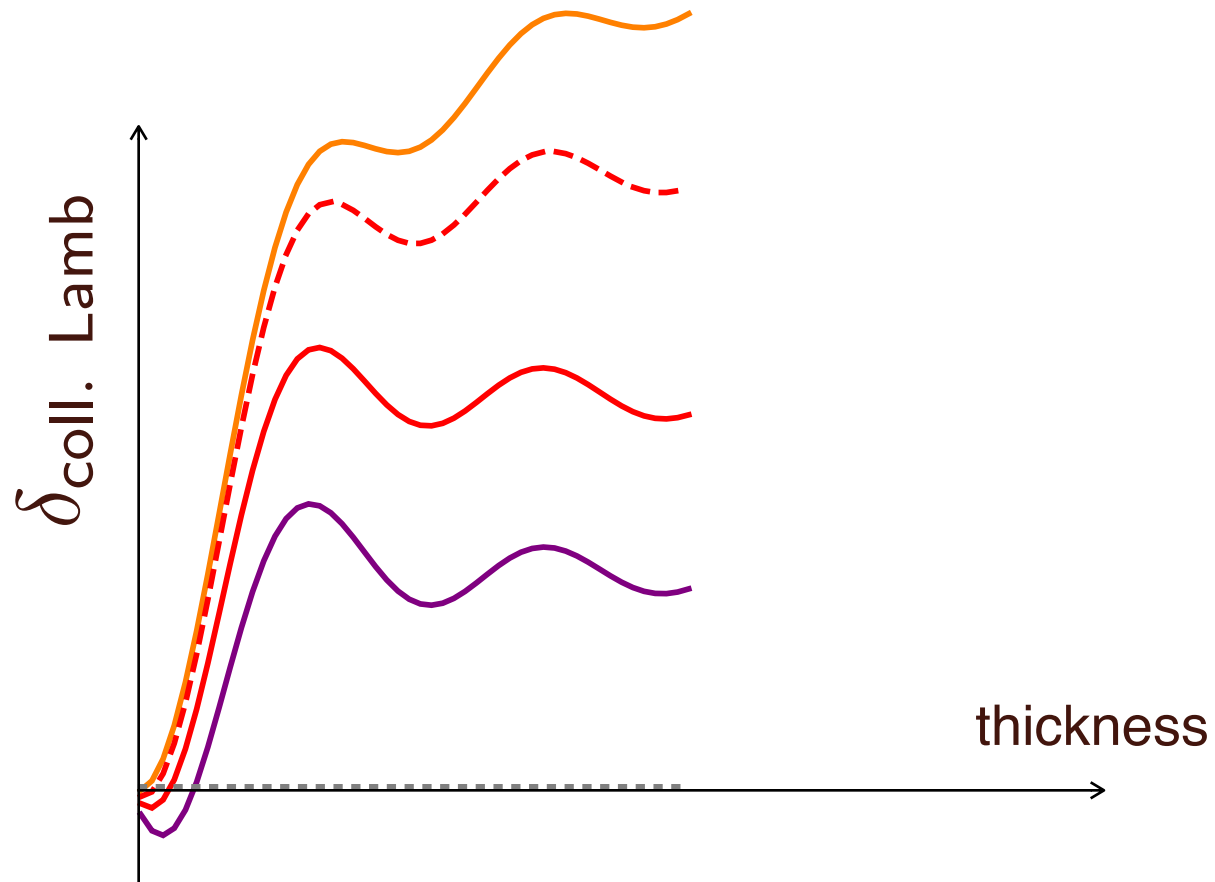


(integrated over thickness)

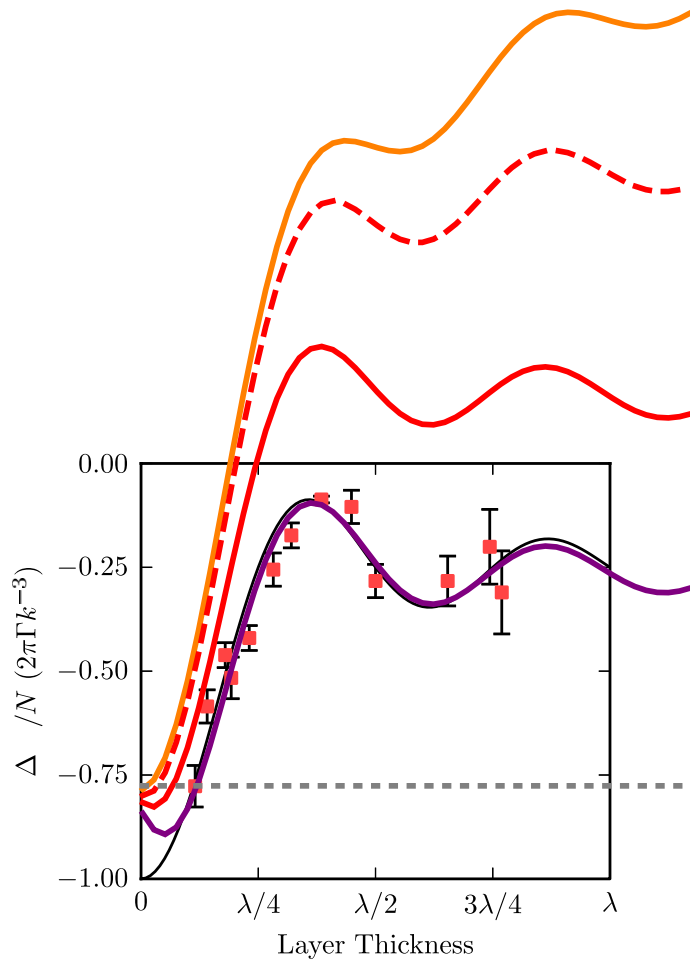


# Collective Lamb Shift

---



# Collective Lamb Shift



# Collective Shift

has **spontaneous** part....

$$\langle\langle E^- E^+ \rangle\rangle \approx \langle\langle a^\dagger a \rangle\rangle = n$$

dependent on number of photons

... and **"stimulated"** part

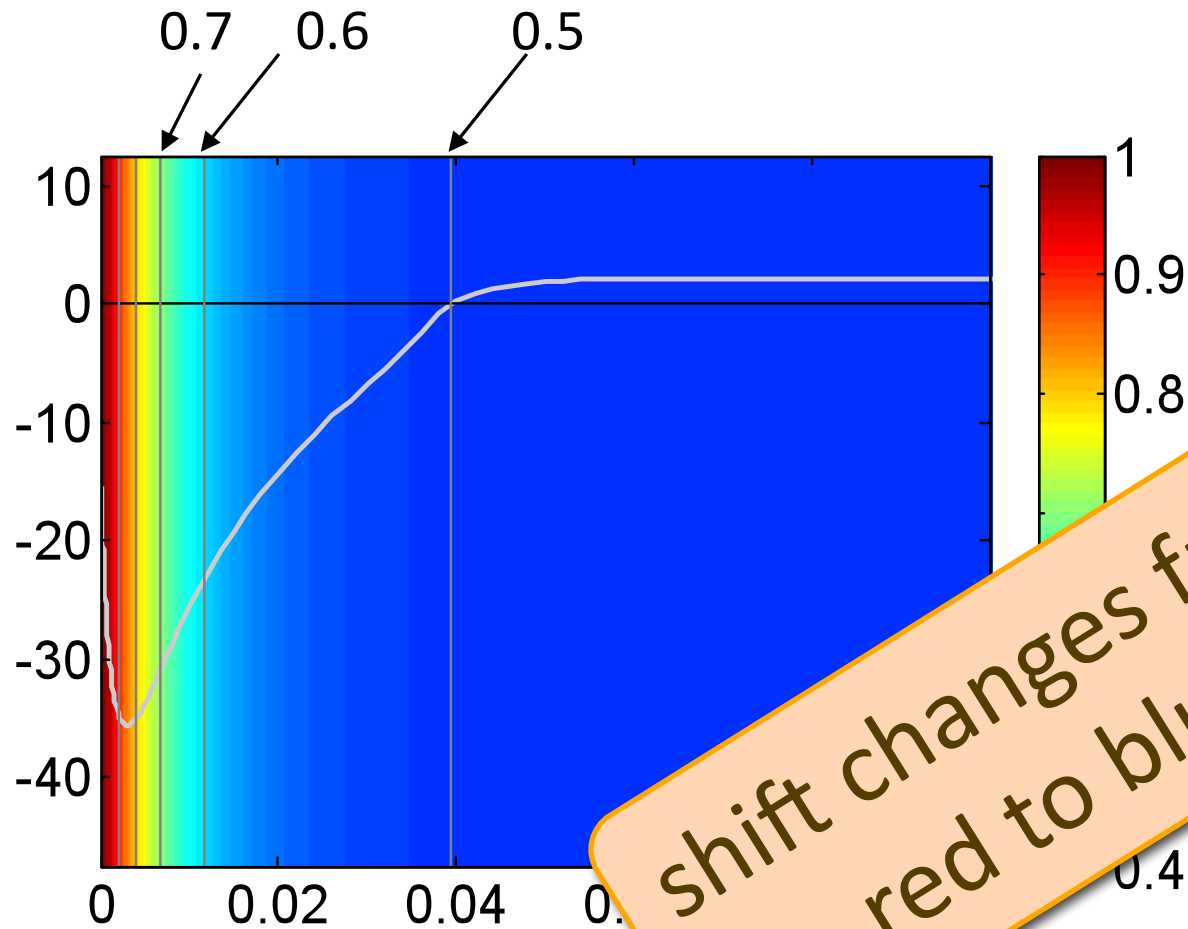
$$\Gamma_{ij}(\omega) = \frac{\wp^2}{\hbar^2} \int d\tau \langle\langle E_i^-(t) E_j^+(t + \tau) \rangle\rangle e^{i\omega\tau}$$

$$\Delta_{\text{stim}}^{(ij)} = \frac{1}{2\pi} \mathcal{P} \int d\omega' \frac{\Gamma_{ij}(\omega')}{\omega - \omega'}$$

Induced shift

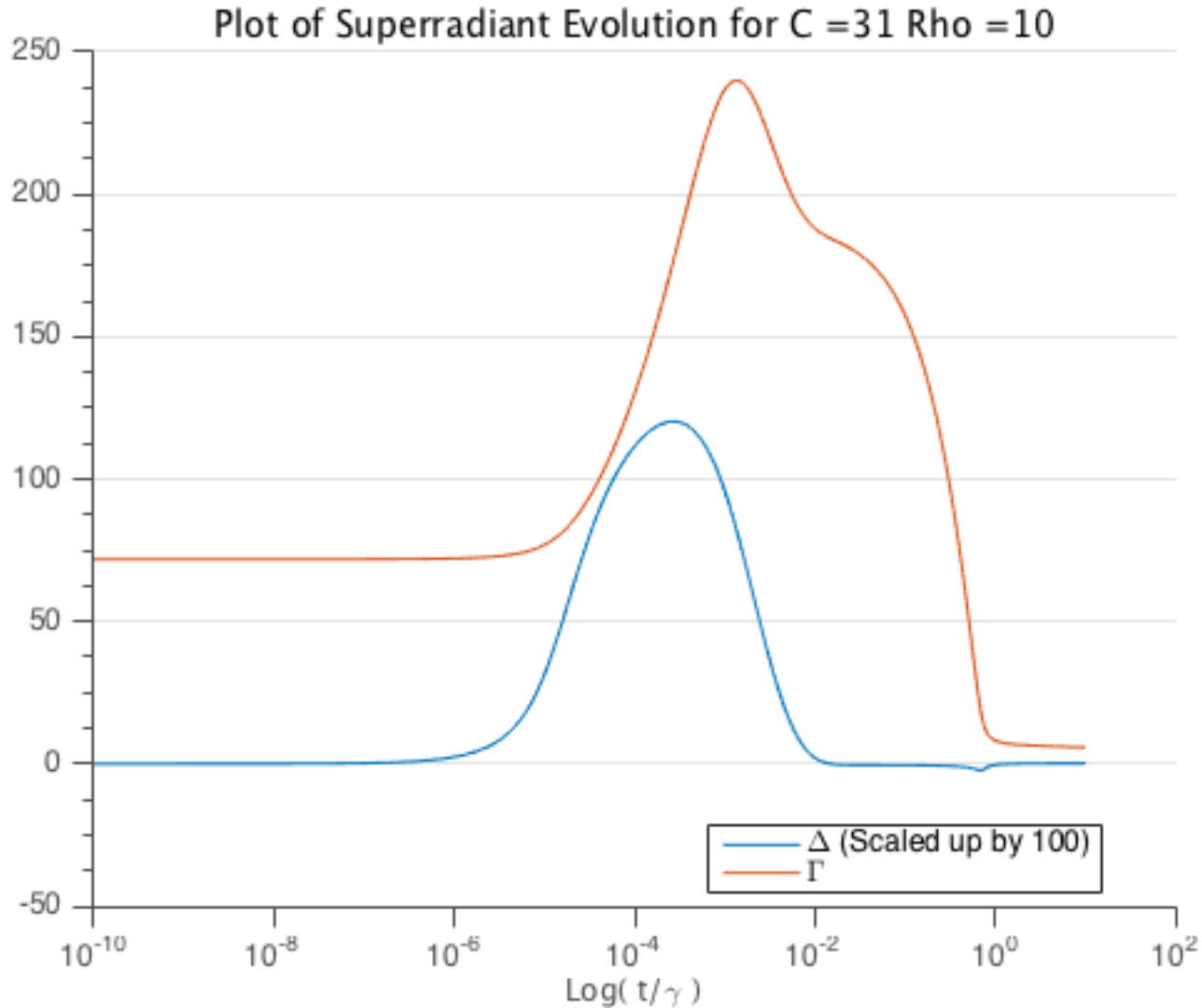


# Induced Shift: Decay of inverted system



shift changes from red to blue

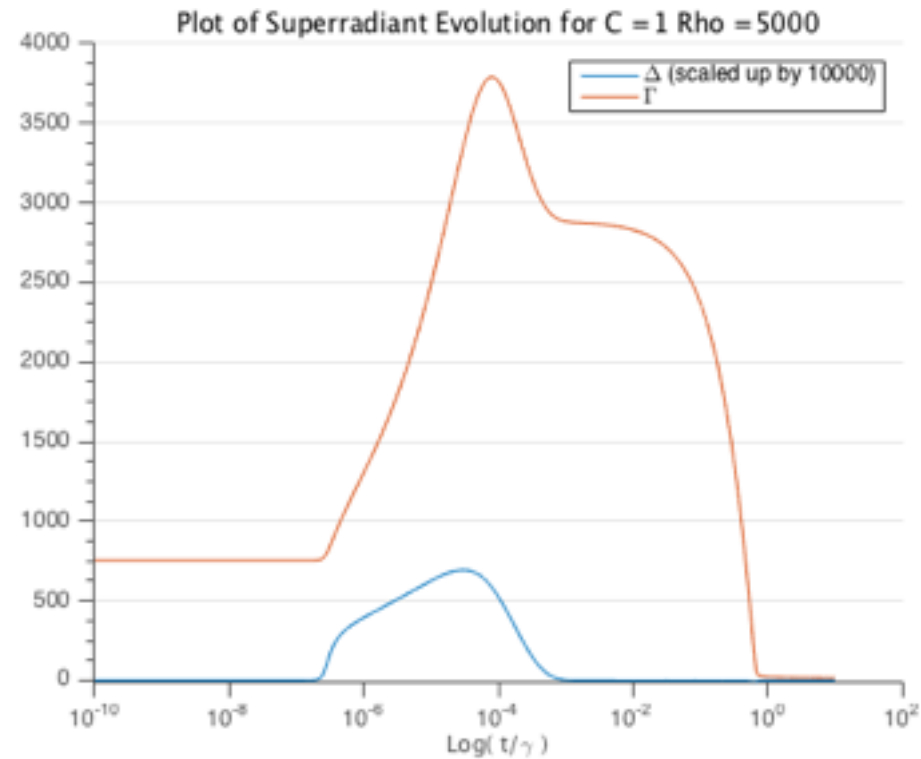
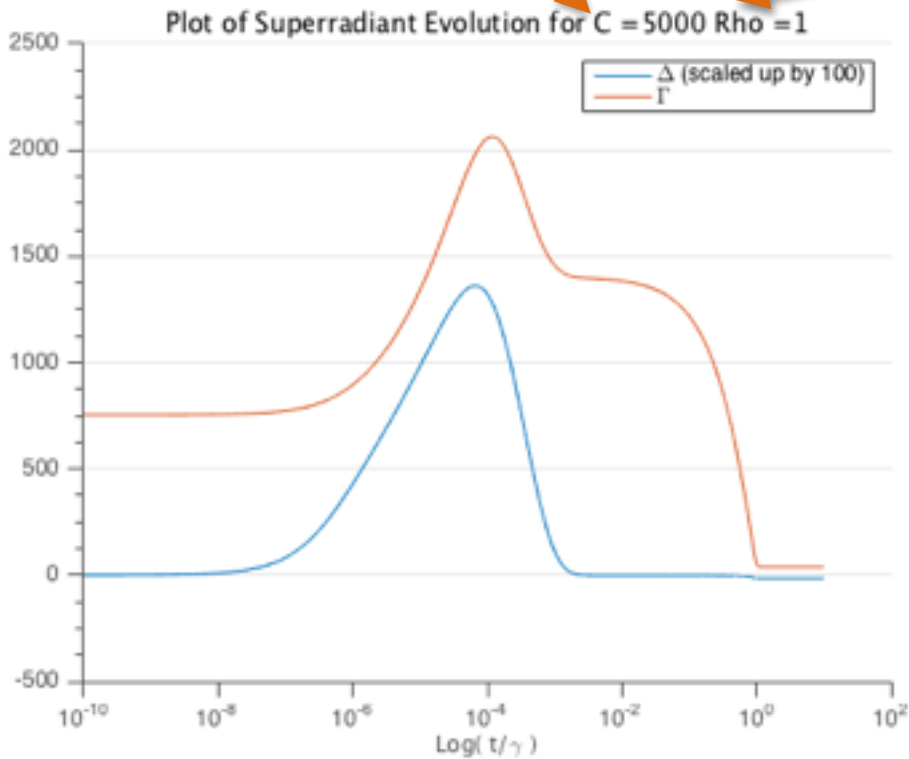
# Collective Shift: decay of inverted TLS



# Collective Shift: decay of inverted TLS

$$C \propto N\lambda^3$$

$$C\rho \propto N\lambda^2L$$

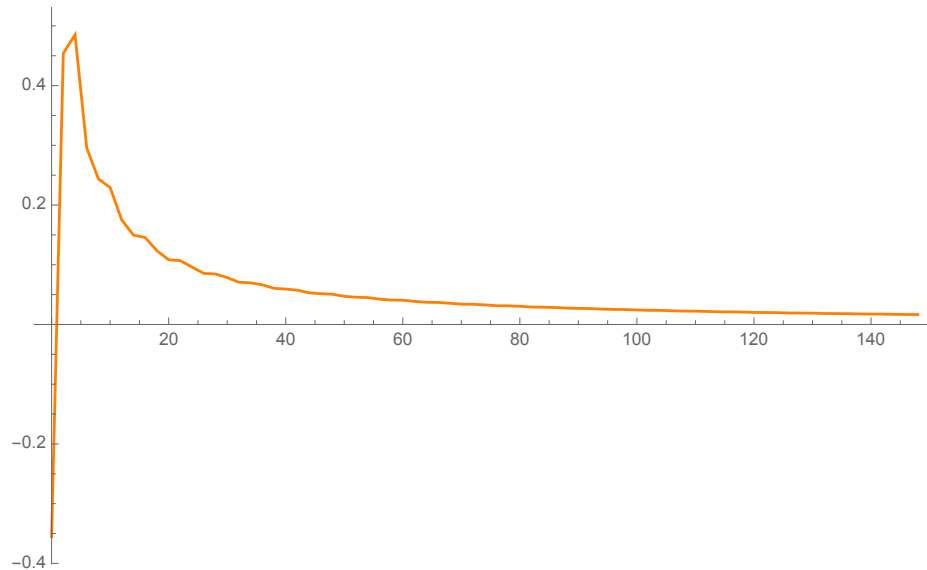


# Collective Shift: low-excitation limit ( $\propto \Omega^2$ )

---

0

# Measurable shift?



## Shifts of a Resonance Line in a Dense Atomic Sample

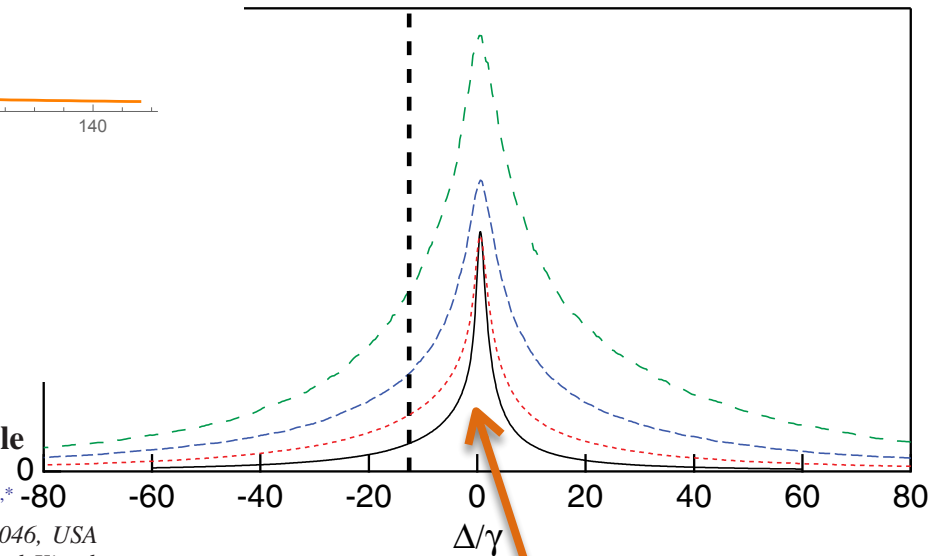
Juha Javanainen,<sup>1</sup> Janne Ruostekoski,<sup>2</sup> Yi Li,<sup>1</sup> and Sung-Mi Yoo<sup>1,\*</sup>

<sup>1</sup>Department of Physics, University of Connecticut, Storrs, Connecticut 06269-3046, USA

<sup>2</sup>Mathematical Sciences, University of Southampton, Southampton SO17 1BJ, United Kingdom

(Received 28 August 2013; published 21 March 2014)

We study the collective response of a dense atomic sample to light essentially exactly using classical-electrodynamics simulations. In a homogeneously broadened atomic sample there is no overt Lorentz-Lorentz local field shift of the resonance, nor a collective Lamb shift. However, the addition of inhomogeneous broadening restores the usual mean-field phenomenology.



no shift!



# Questions - guideline

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- Superradiance - What? Why?
- How do we calculate it (better)?
- Is there a collective (Lamb) shift?
- How does entanglement come into the picture?

# Superradiance and Entanglement

---

Does (Dicke) superradiance need/create entanglement?

NO

# Superradiance and Entanglement

---

How to define/calculate many-particle entanglement?

Spin Squeezing Inequalities and Entanglement of N Qubit States J. K. Korbicz, J. I. Cirac, M. Lewenstein

Separability in  $2 \times N$  composite quantum systems B. Kraus, J. I. Cirac, S. Karnas, M. Lewenstein

Entangled symmetric states of N qubits with all positive partial transpositions R. Augusiak, J. Tura, J. Samsonowicz, M. Lewenstein

Four-qubit entangled symmetric states with positive partial transpositions J. Tura, R. Augusiak, P. Hyllus, M. Kuś, J. Samsonowicz, M. Lewenstein

Separability criteria and entanglement witnesses for symmetric

# Superradiance and Entanglement

---

Does (Dicke) superradiance need/create entanglement? (Initial state: no entanglement)

Dicke superradiant  
time evolution

=

separable states

**constructive proof**

# Superradiance and Entanglement

---

Does (Dicke) superradiance need/create entanglement? (Initial state: no entanglement)

our system: mixed state of  
N-atom Dicke states with  
N+1 known independent  
coefficients  $p_i$

=

compare to mixture of  
symmetric product states of  
N (two-level) atoms (needs  
N+1 coefficients  $y_i$ )

(N+1) - dim.  
equation  
system

# Superradiance and Entanglement

Does (Dicke) superradiance need/create entanglement? (Initial state: no entanglement)

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N (two-level) atoms (needs  
N+1 coefficients  $y_i$ )

condition:  
all coefficients  
 $0 \leq p_i \leq 1$

(N+1) - dim.  
equation  
system

# Superradiance and Entanglement

Does (Dicke) superradiance need/create entanglement? (Initial state: no entanglement)

our system: mixed state of  
N-atom Dicke states with  
N+1 known independent  
coefficients  $p_i$

=

complete mixture of  
symmetric Dicke states of  
N (terms  $y_i$ ) (needs  
N+1 parameters  $y_i$ )

condition:  
all coefficients  
 $0 \leq p_i \leq 1$

(N+1) - dim.  
equation  
system



# Superradiance and Entanglement

---

Driven superradiant system:



# Fuzzy Bunny?

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# Spin Squeezing

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- Correlated (“squeezed”) spins could improve resolution in one direction (“quadrature”).

# (Spin) Squeezing

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- How to measure squeezing/measurement improvement?

$$\xi^2 \equiv \frac{\text{optimal variance}}{\text{unsqueezed optimal variance}}$$

# Spin squeezing

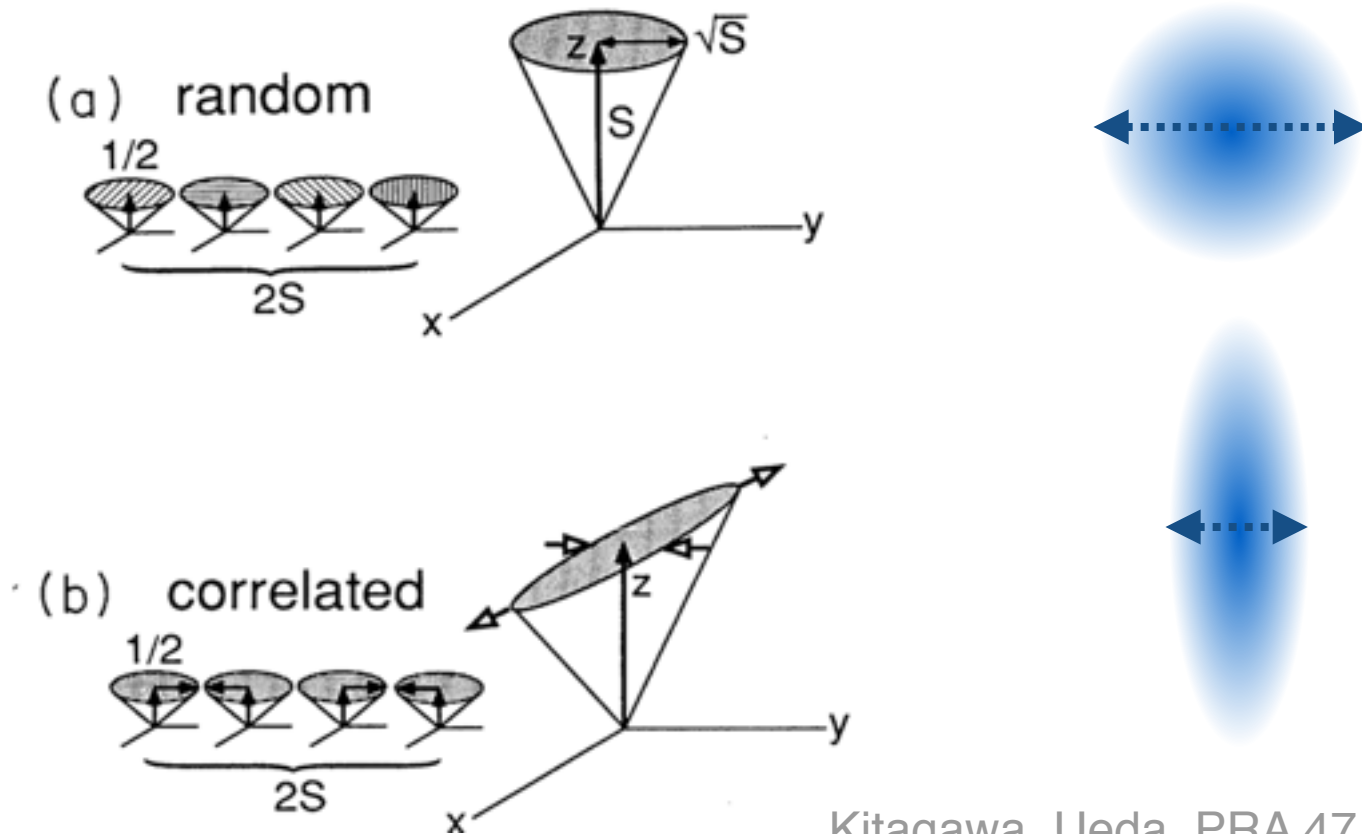
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Old problem: How to improve metrology by spin squeezing ensembles

- ➔ Groups of Bigelow, Kuzmich, Lewenstein, Mølmer, Polzik, Sanders, Sørensen, Vuletic, Wineland,...

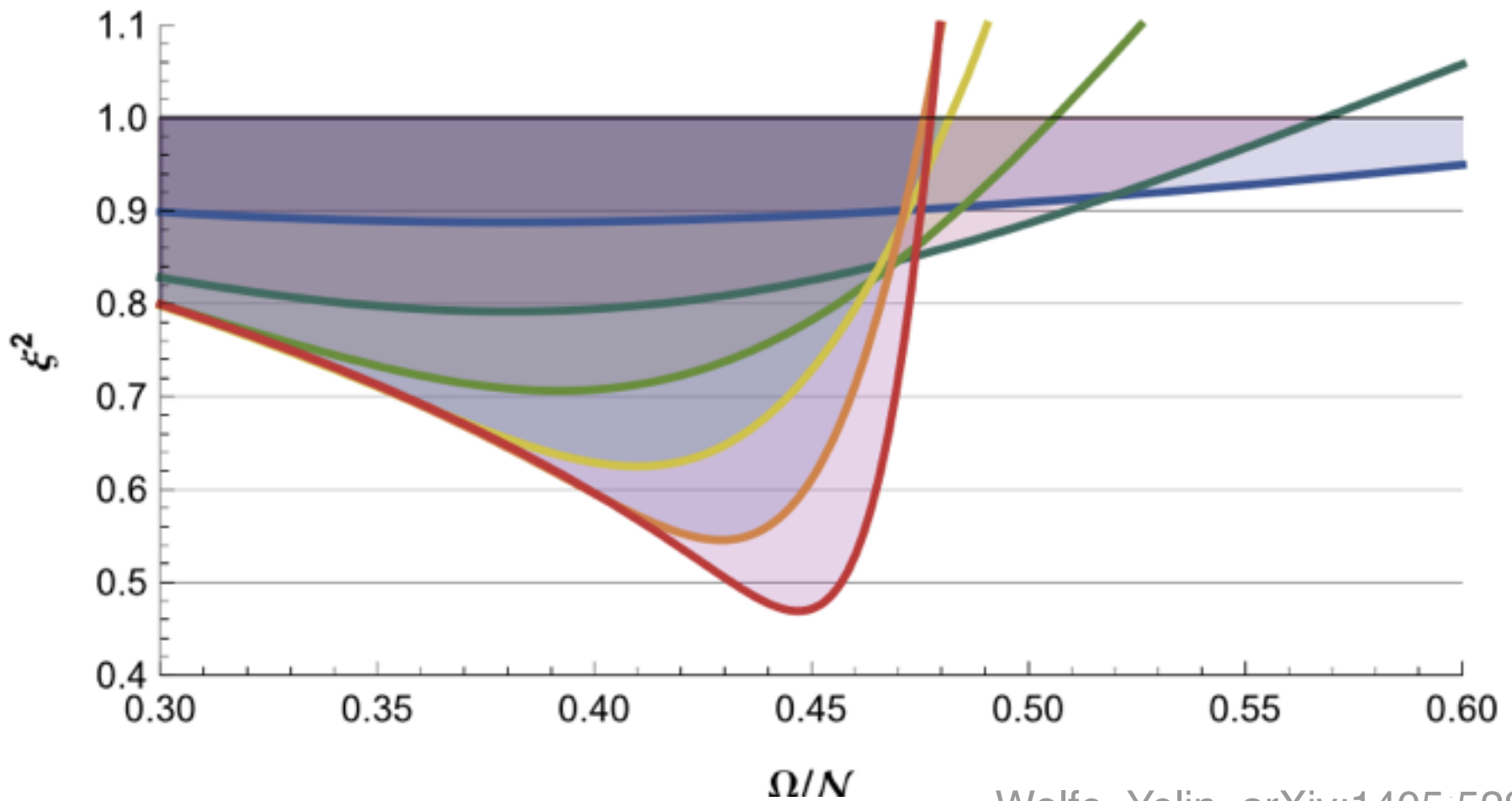
# Spin Squeezing

- Correlated (“squeezed”) spins could improve resolution in one direction (“quadrature”).



# Superradiant Spin Squeezing

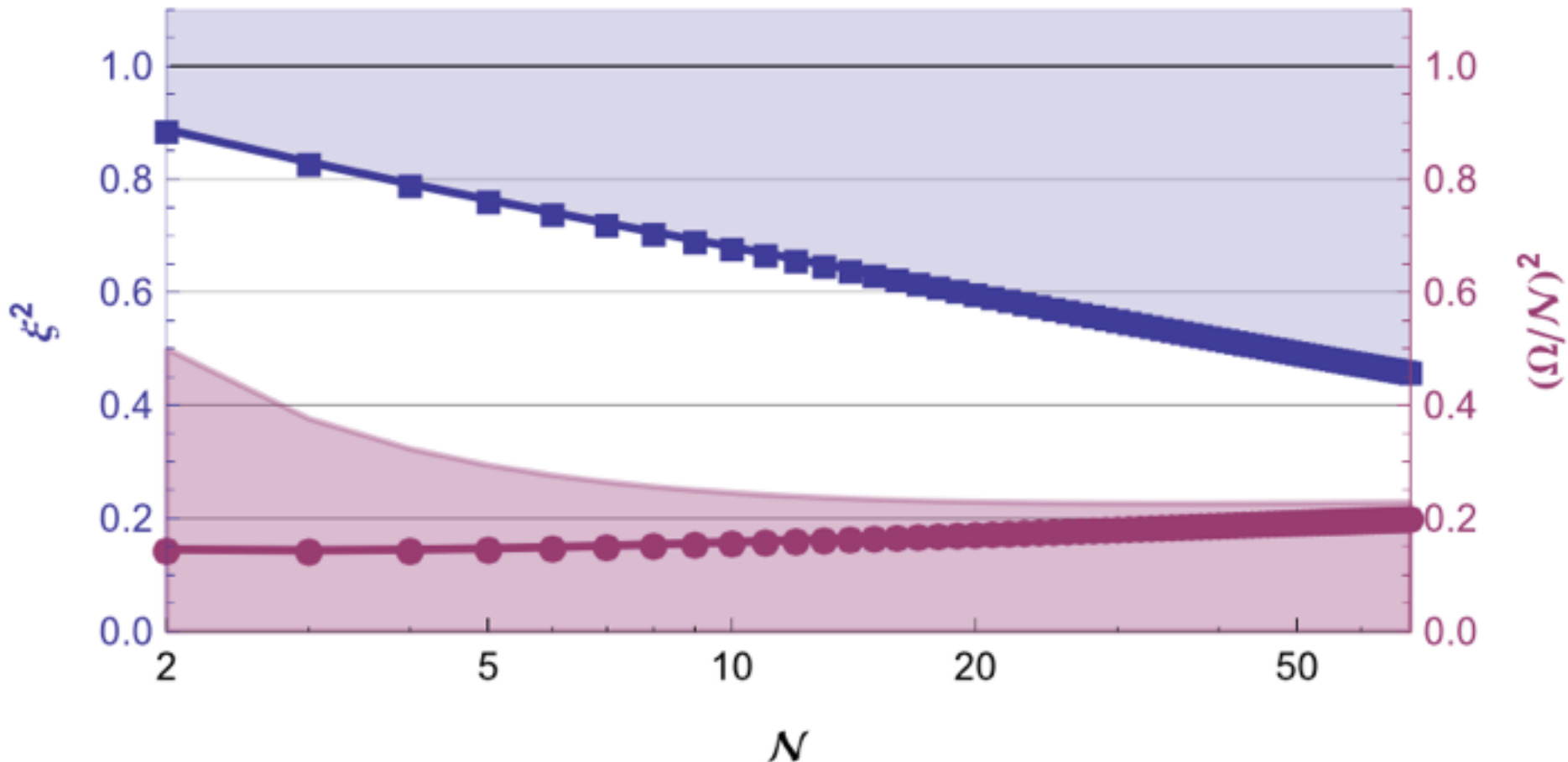
$\xi^2$  vs.  $\Omega/N$ ,  $N =$  ■ 2 ■ 4 ■ 8 ■ 16 ■ 32 ■ 64



# Best case for Dicke ensemble

■ minimal possible  $\xi^2$

● most optimal  $(\Omega/\mathcal{N})^2$



# Conclusions, Applications and Outlook

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- Superradiance - What? Why?
  - \* Collective effect + exchange
- How do we calculate it (better)?
  - \* large, homogeneous, self-consistently
  - ➔ small, ordered? higher correlations?
- Is there a collective (Lamb) shift?
  - \* Yes, and yes.
  - ➔ Find schemes to measure!
- How does entanglement come into the picture?
  - \* Cooperativity alone does not create entanglement, but cooperative + driving interaction squeeze
  - ➔ Find in more realistic systems + squeeze THz light fields



Full dynamics (all degrees of freedom of atoms, fields)

---

$$H = H_{\text{atoms}} + H_{\text{field}} - \sum \mathbf{p}_i \mathbf{E}_i$$



two possible atoms  
some other atoms

# Thank

# you!



