

Convection Modifies MHD Turbulence in Accretion Disks

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KITP

*Disks, Dynamos, and Data:
Confronting MHD Accretion Theory with Observations*

Collaborators:

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Iwona Kotko
Julian Krolik
Jean-Pierre Lasota

Greg Salvesen
Takayoshi Sano
Evan Yerger

Outline

1 Context

- Accretion Disk Theory
- Dwarf Novae

2 Convective Accretion Disks

- Convection Enhances α
- Convection Quenches Magnetic Field Reversals

3 Observational Implications (Lightcurves)

4 Summary

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α -Disk Model

Shakura & Sunyaev (1973)

$$\alpha \equiv \frac{\langle \text{stress} \rangle_z}{\langle \text{pressure} \rangle_z} = \frac{\int_{-\infty}^{\infty} w_{r\phi} dz}{\int_{-\infty}^{\infty} p dz} \equiv \frac{W_{r\phi}}{\int_{-\infty}^{\infty} p dz}$$

- Prescription to estimate turbulent stresses
 - Turbulent stresses can also be thought of as an effective viscosity

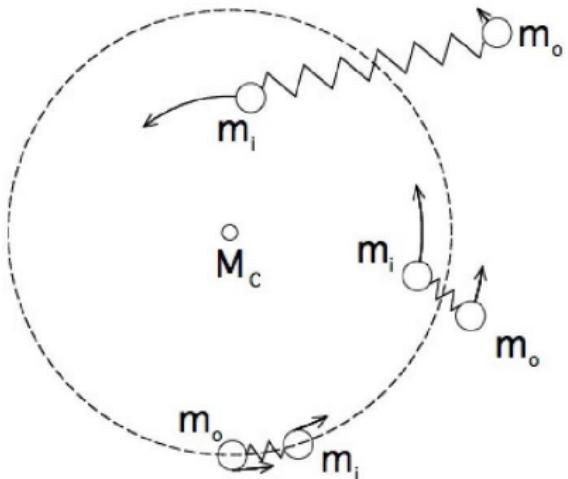
$$\bar{\nu}_{\text{eff}} = \frac{2}{3} \frac{W_{r\phi}}{\Sigma \Omega} = \frac{2}{3} \alpha \frac{\int_{-\infty}^{\infty} p dz}{\Sigma \Omega}$$

- Allows us to approximate the structure of disks

Source of Stress

$$w_{r\phi} = -\frac{B_r B_\phi}{4\pi} + \rho v_r \delta v_\phi$$

- Magnetorotational Instability (MRI)
 - Generates turbulent stress
 - Can be simulated



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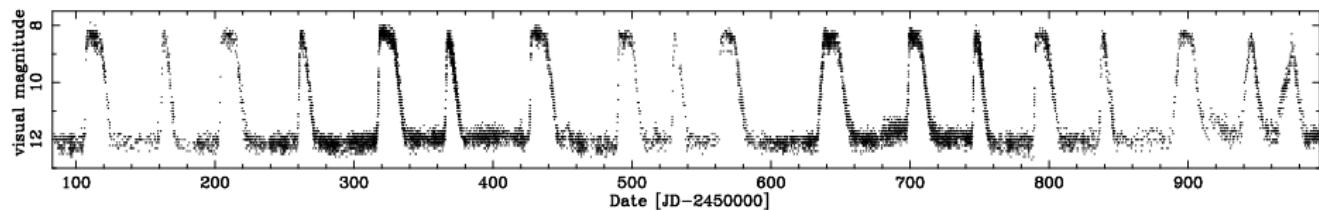
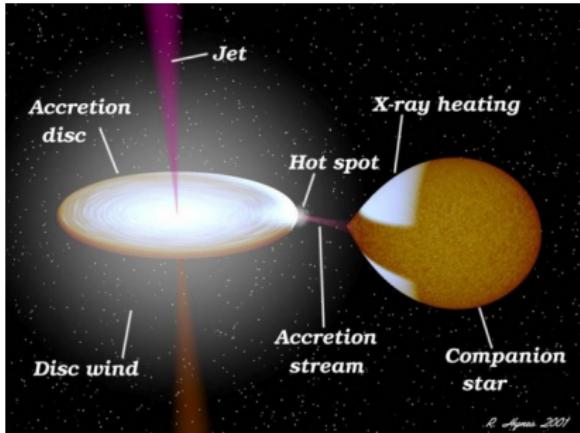
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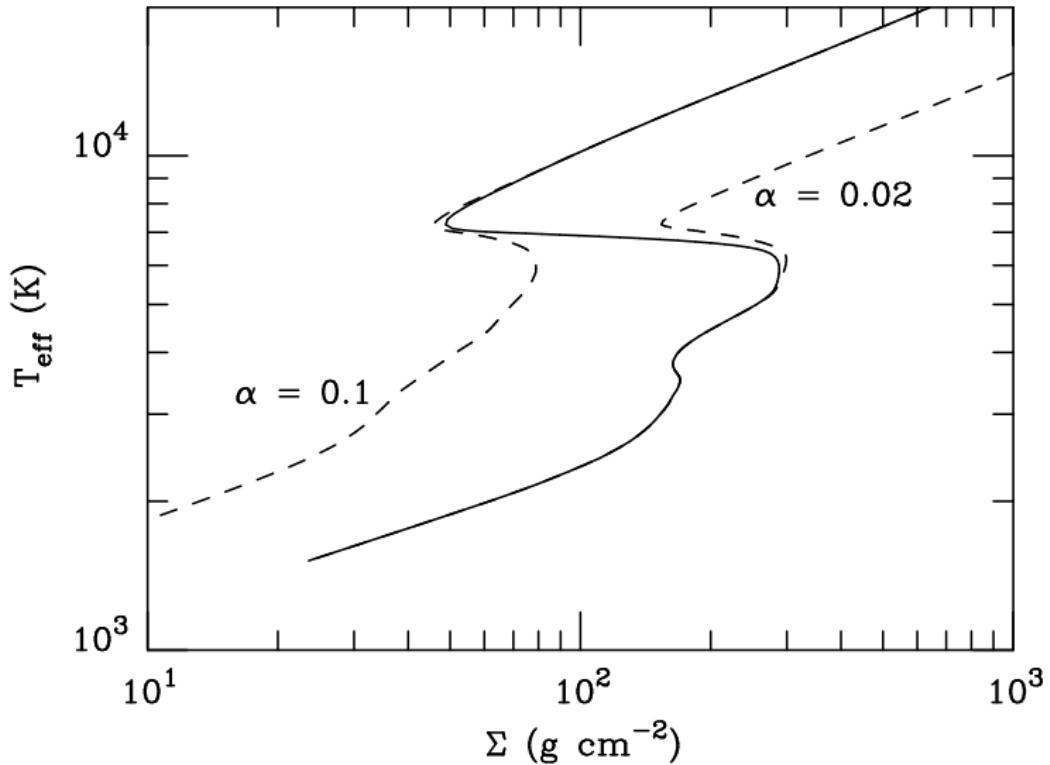
Dwarf Novae

- DNe are white dwarfs in binaries (Cataclysmic variables)
- Powered by thermal instability in the accretion disk
- Extremely well studied
 - SS Cygni was discovered in 1896



SS Cygni (Wheatley et al., 2003)

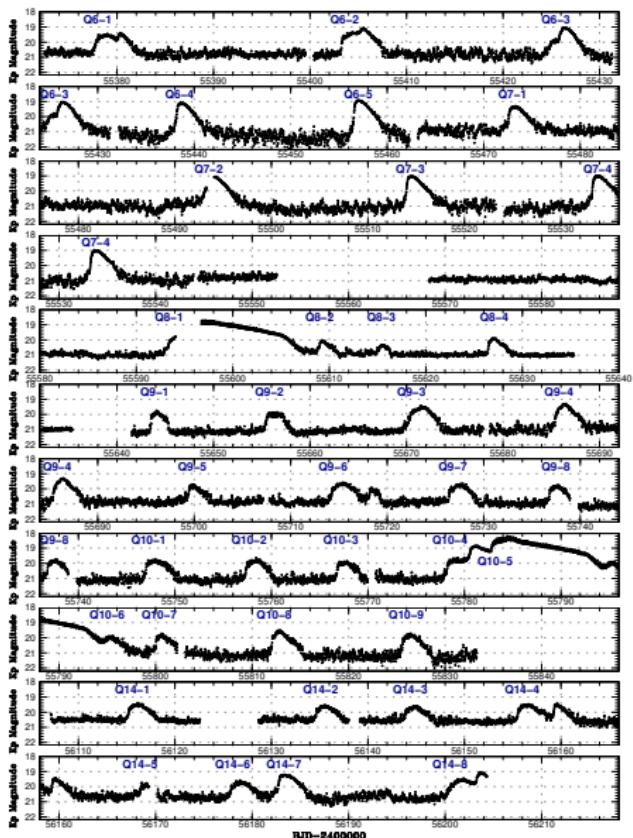
S-curve



Lasota (2001)

Dwarf Novae can Illuminate the Underlying Physics

- Wealth of data
 - Kepler data
 - Properties of steady state disks are insensitive to α
 - Time variations give α
 - The mechanism which determines α is not well understood
 - Previous accretion disk simulations give $\alpha \sim 0.01$
 - These lay outside the dwarf nova regime



Kato & Osaki (2013)

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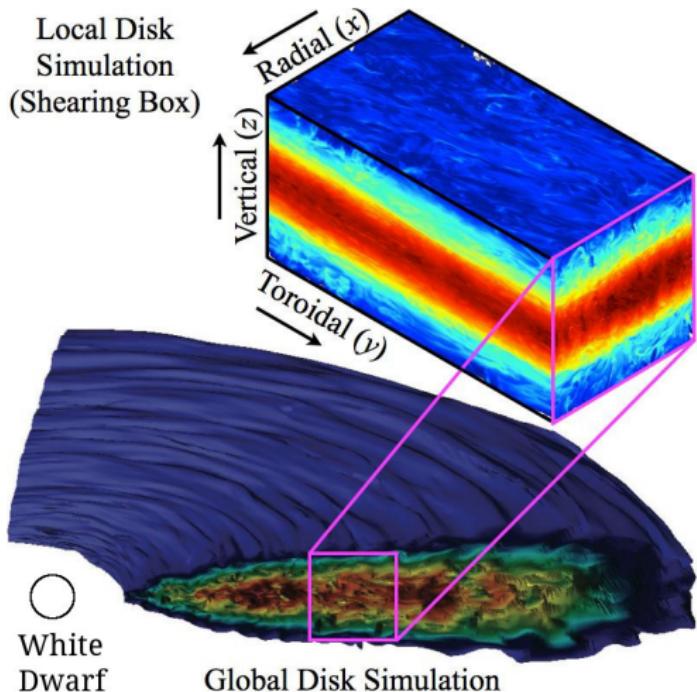
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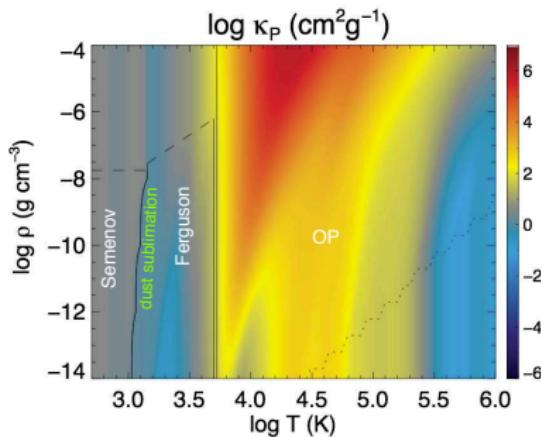
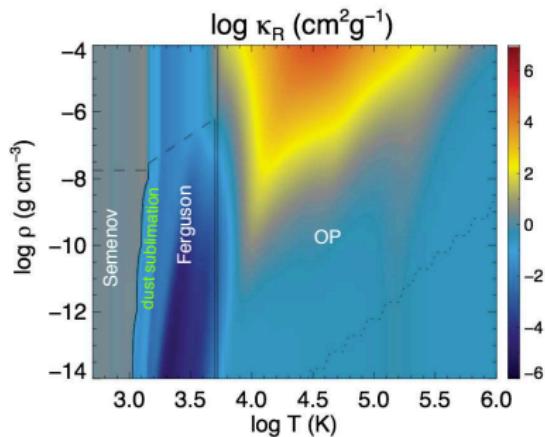
4 Summary

Simulations Methods

- Modified ZEUS code
- Ideal MHD
- Flux limited diffusion
- Stratified shearing box geometry
- Initial conditions based on 1D model
- Realistic dust free opacity
- Realistic equation of state

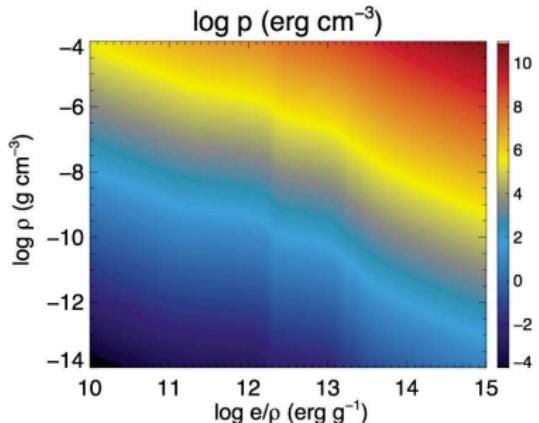
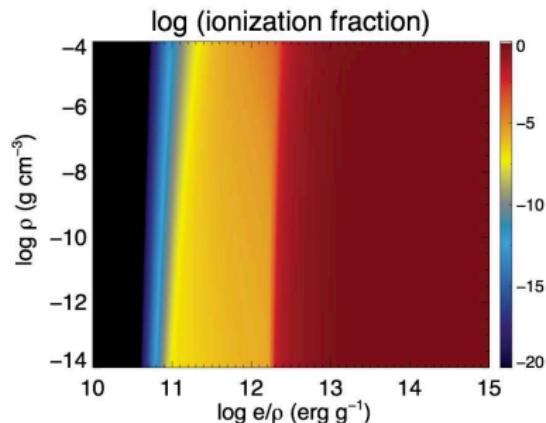
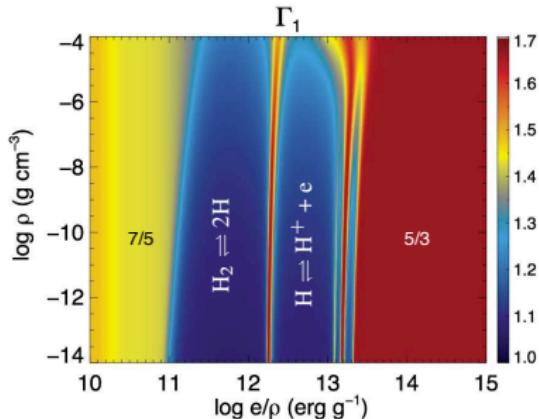
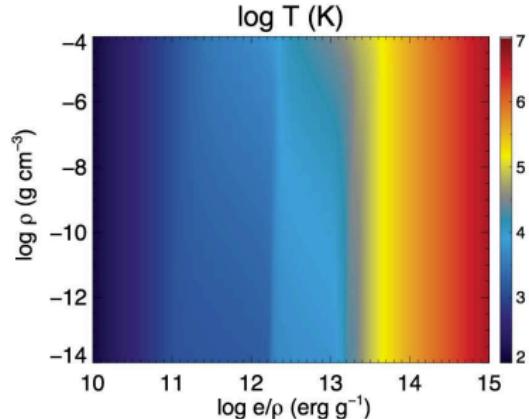


Opacity Tables

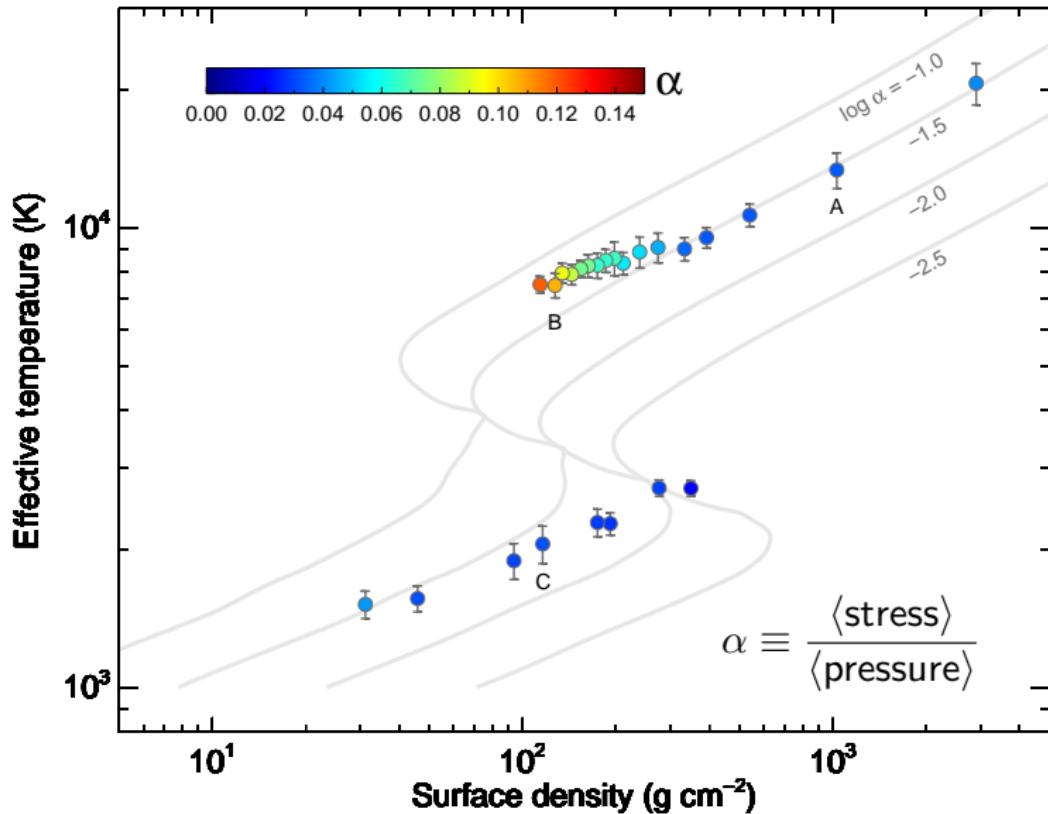


Mean opacities combining three published opacity tables by Semenov et al. (2003), Ferguson et al. (2005) and the Opacity Project (OPCD_3.3).

Equation of State

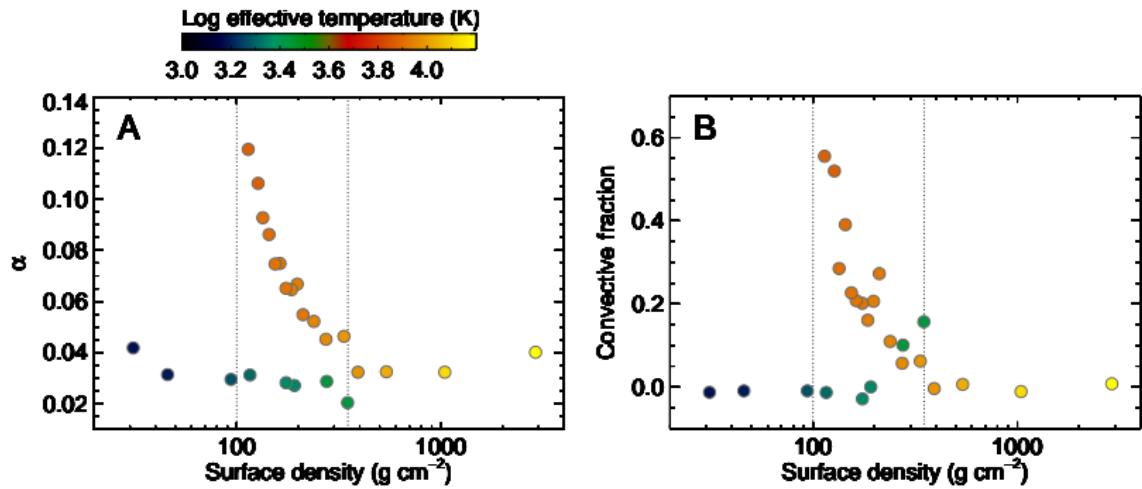


Simulation S-curve



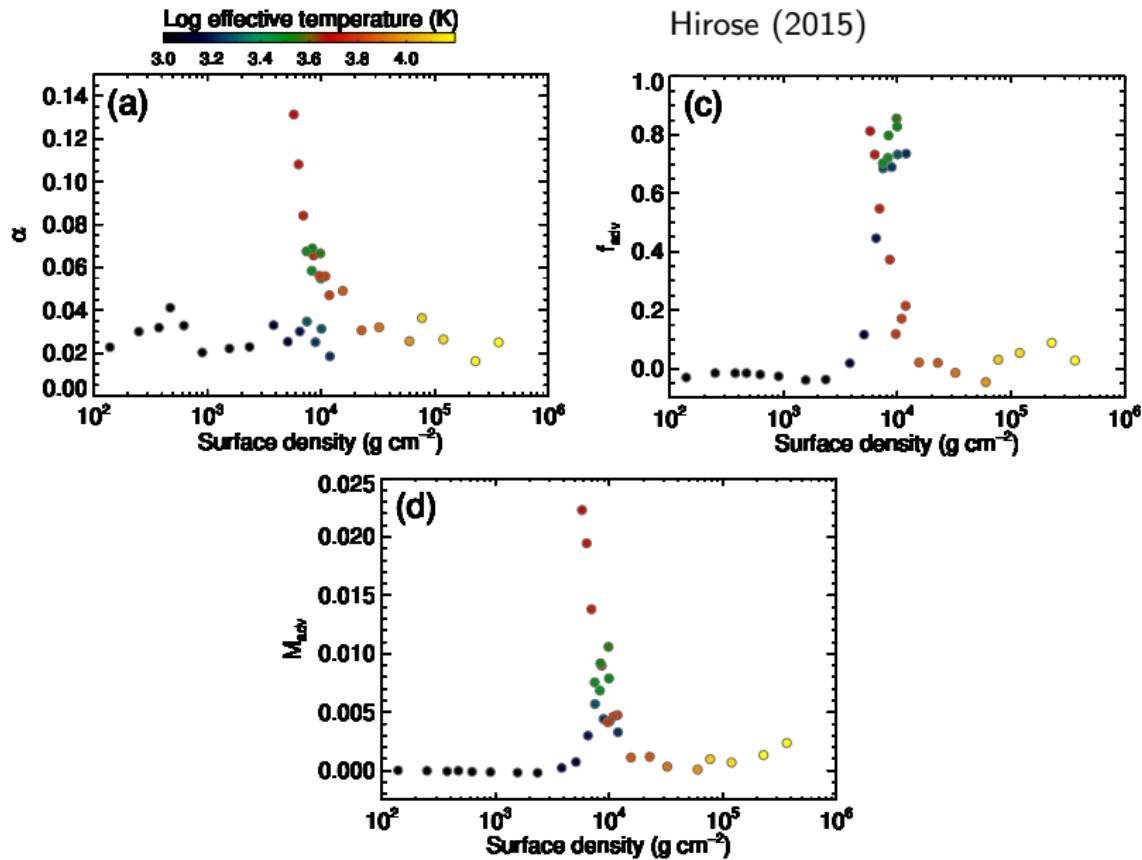
Hirose et al. (2014)

Convection Enhances α

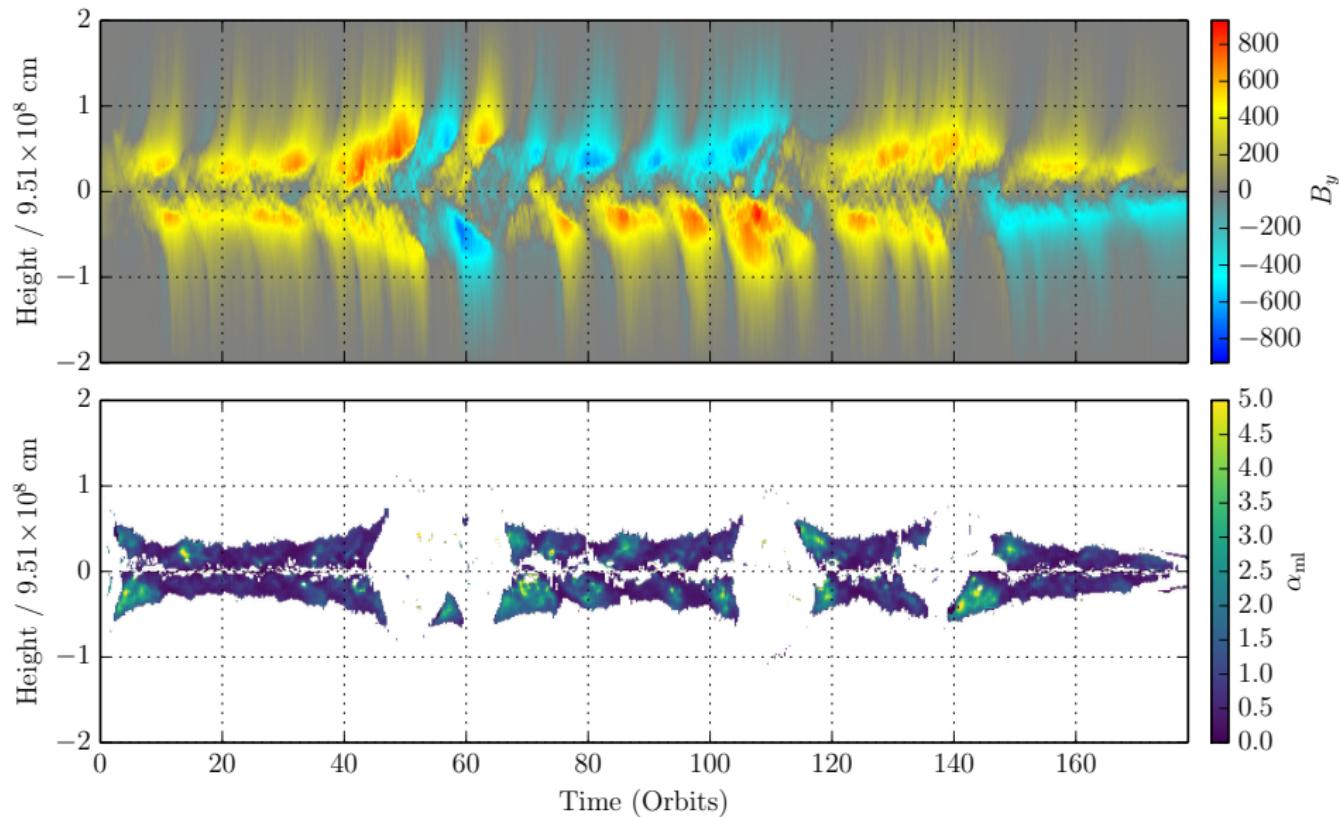


Hirose et al. (2014)

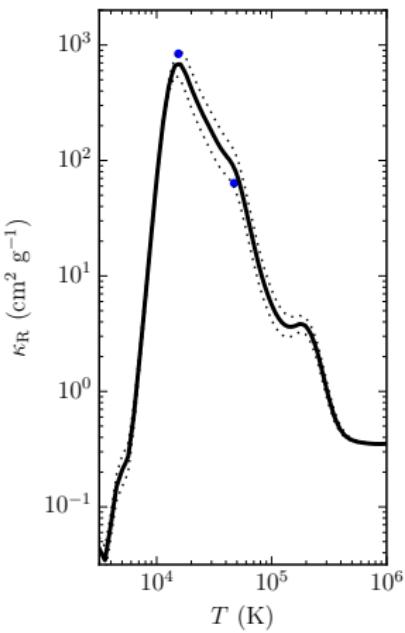
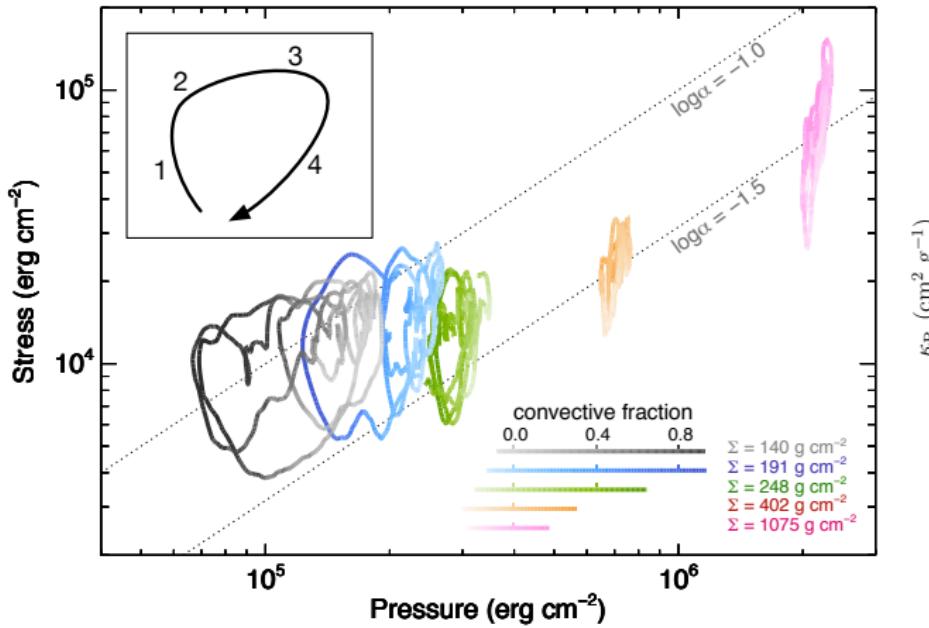
Convection Enhances α in FU Ori



Intermittent Convection



Convective Limit Cycle



Hirose et al. (2014)

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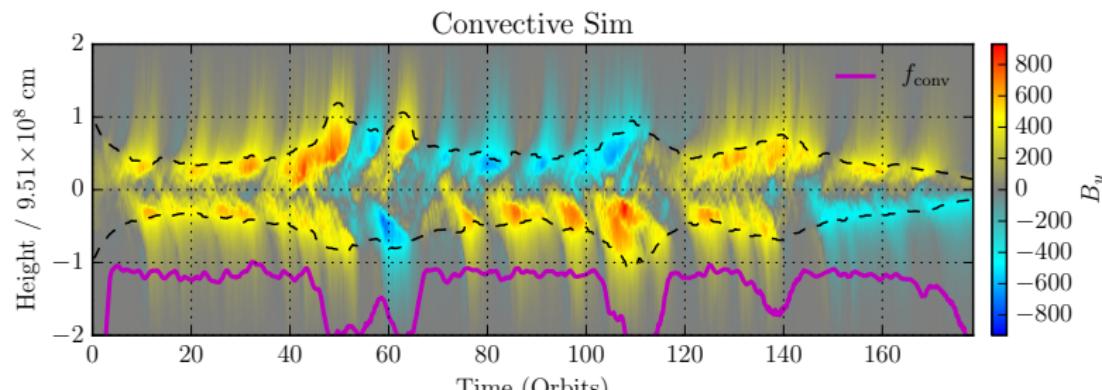
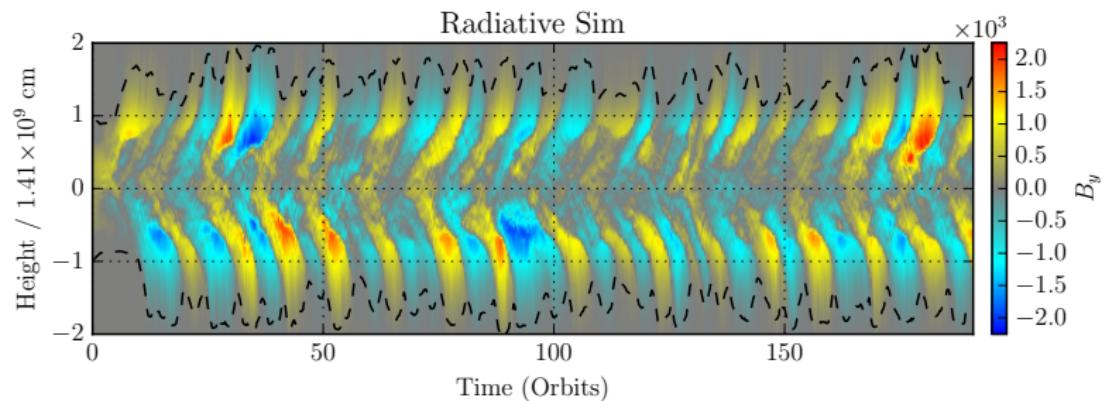
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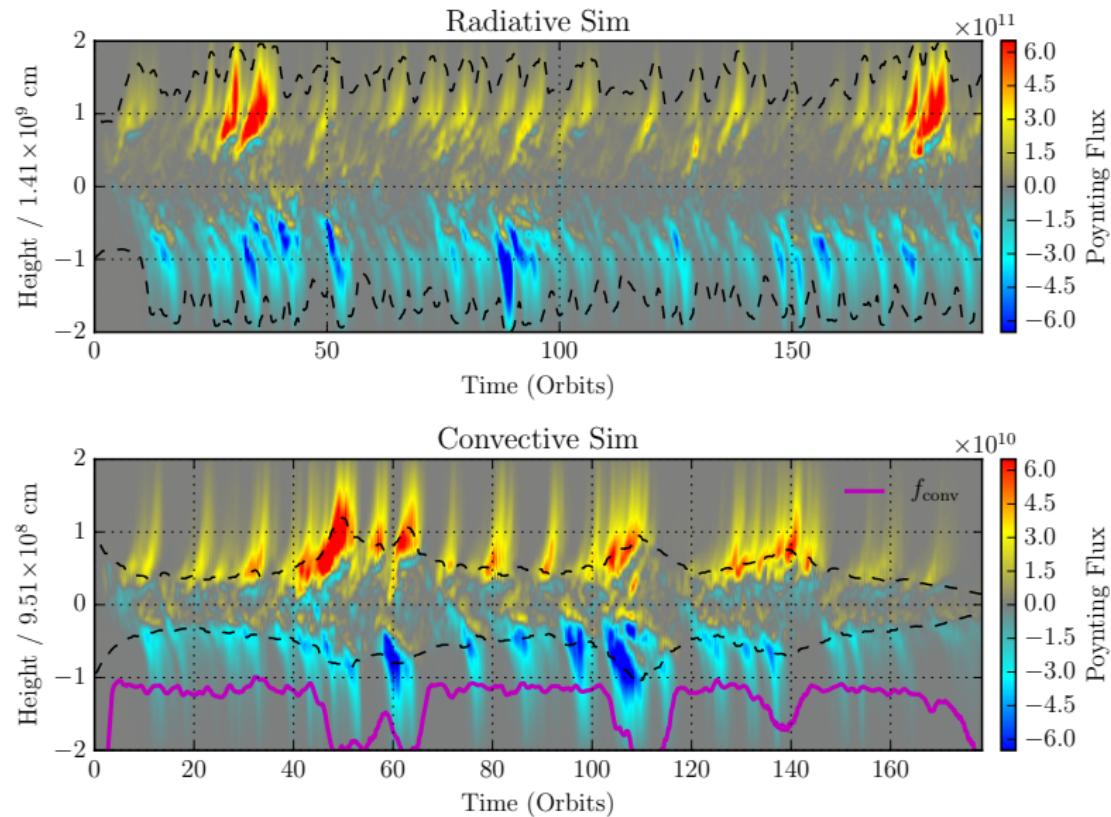
4 Summary

Magnetic Dynamos and the Butterfly Diagram

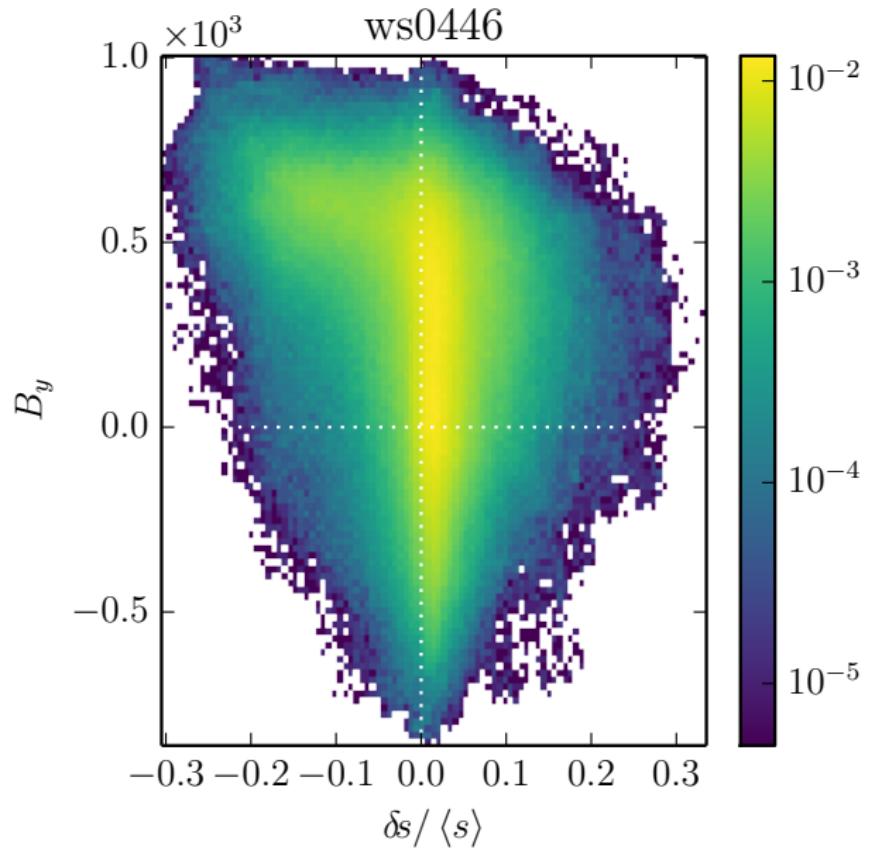


Coleman et al. (2017)

Poynting Flux



Convection Drags Magnetic Field



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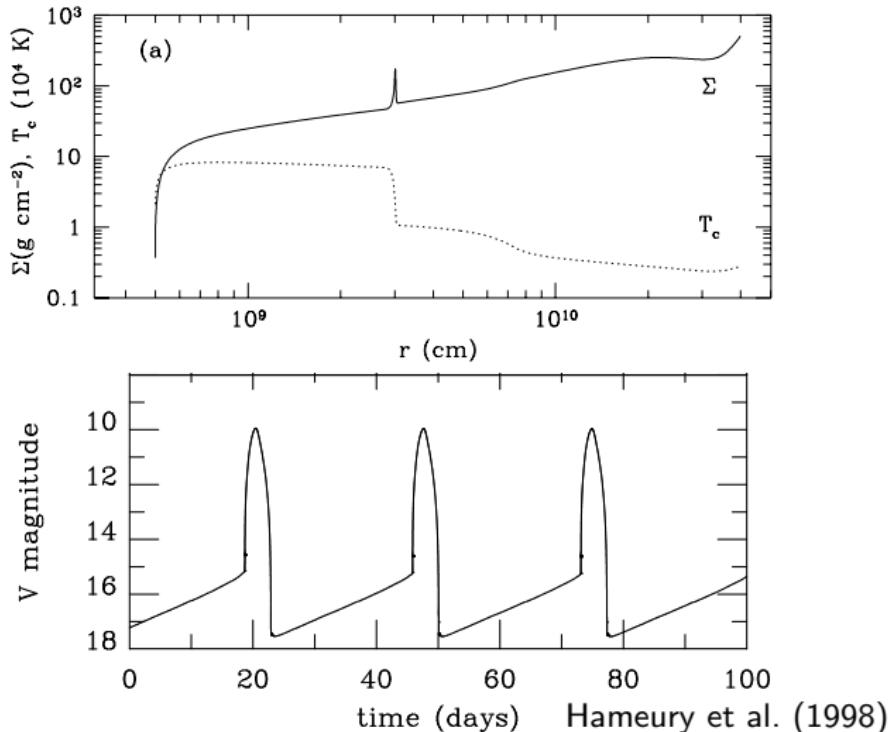
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Disk Instability Model

- Developed by Jean-Marie Hameury and Guillaume Dubus
- Uses library of equilibrium vertical structures to determine how adjacent annuli in the disk communicate.



Improving Vertical Structures

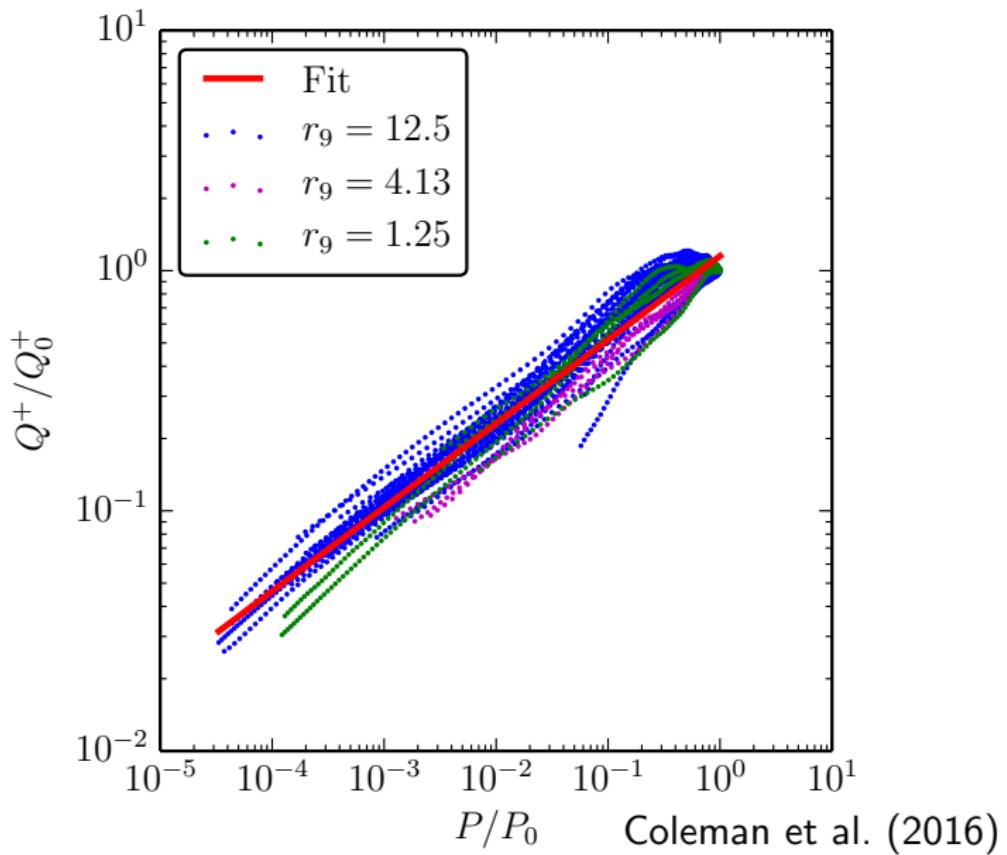
Existing Models

- $Q^+(z) = \frac{3}{2}\alpha\Omega P(z)$
- Lack description of magnetic pressure
- $\alpha(T_0) \approx \begin{cases} \alpha_{\text{hot}} & T_0 > T_c \\ \alpha_{\text{cold}} & T_0 < T_c \end{cases}$
- Mixing length based on Sun

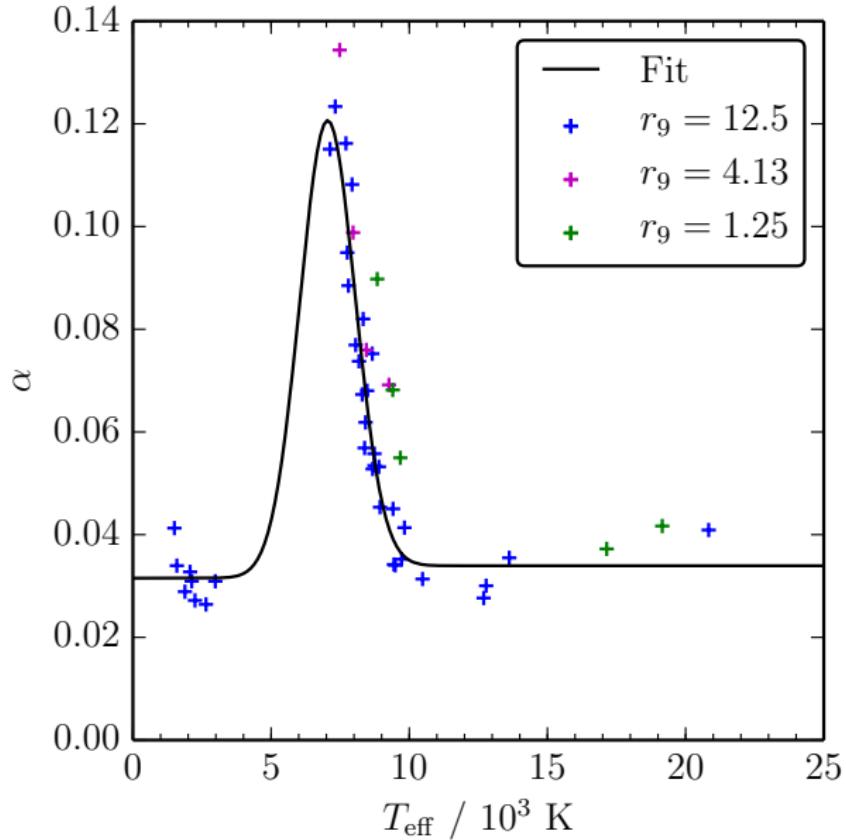
Modified Models

- $Q^+(z) \propto P(z)^\beta$
 - $\beta \approx 0.35$
- Still lacks magnetic pressure
- $\alpha(T_{\text{eff}}) = \alpha_{\text{sim}}(T_{\text{eff}})$
- Mixing length based on time-averaged DN simulations ($\alpha_{ml} = 6$)

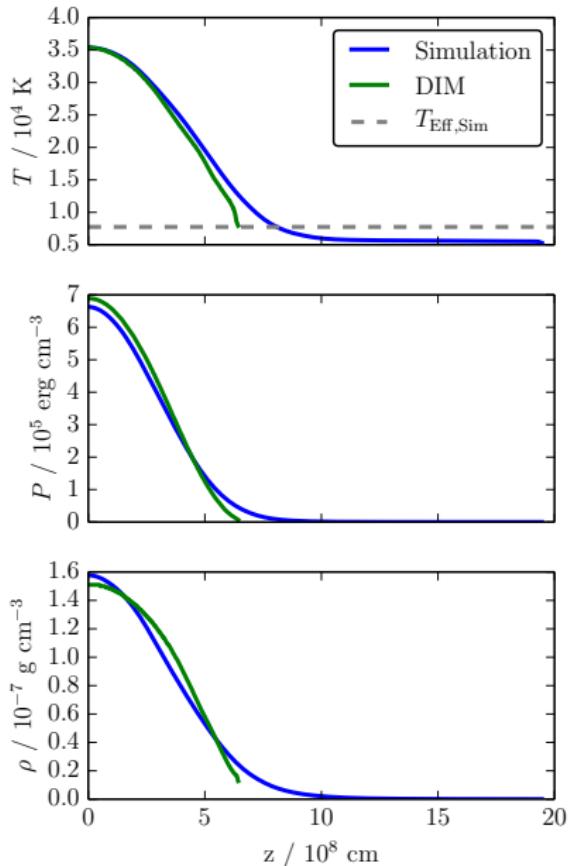
Simulation Profile data



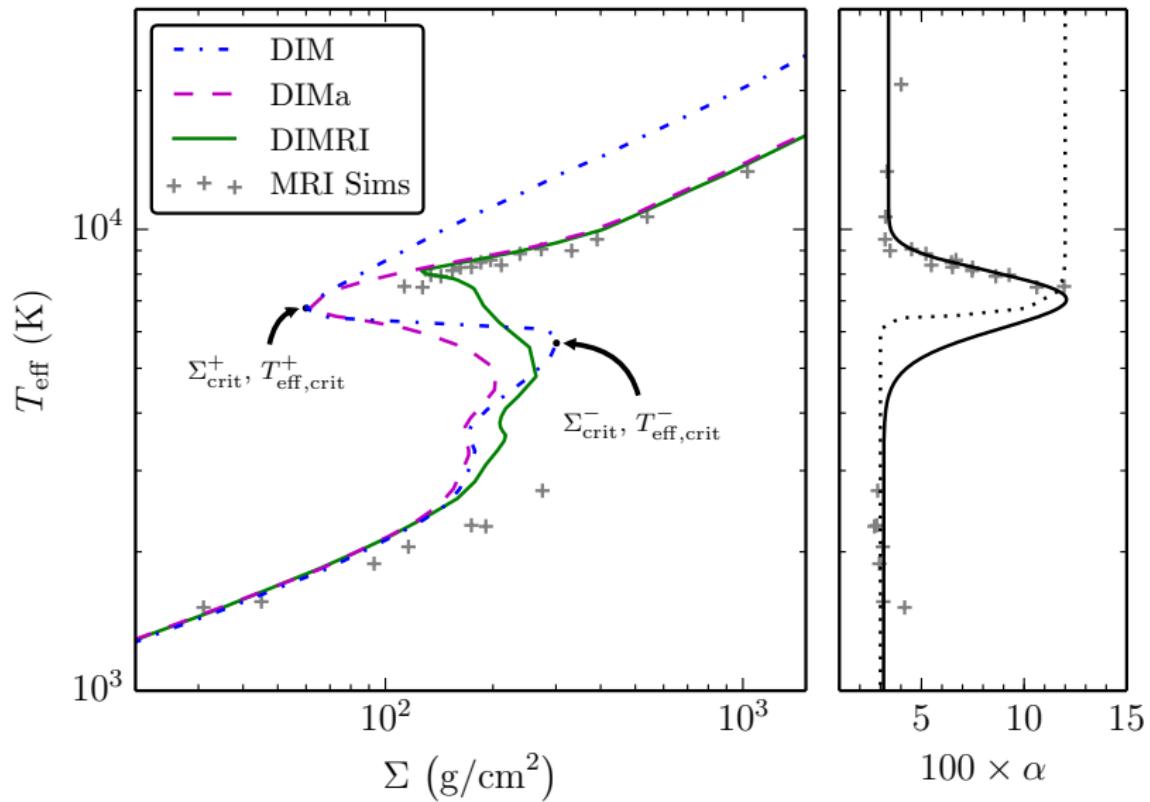
α from Simulations is Well Behaved



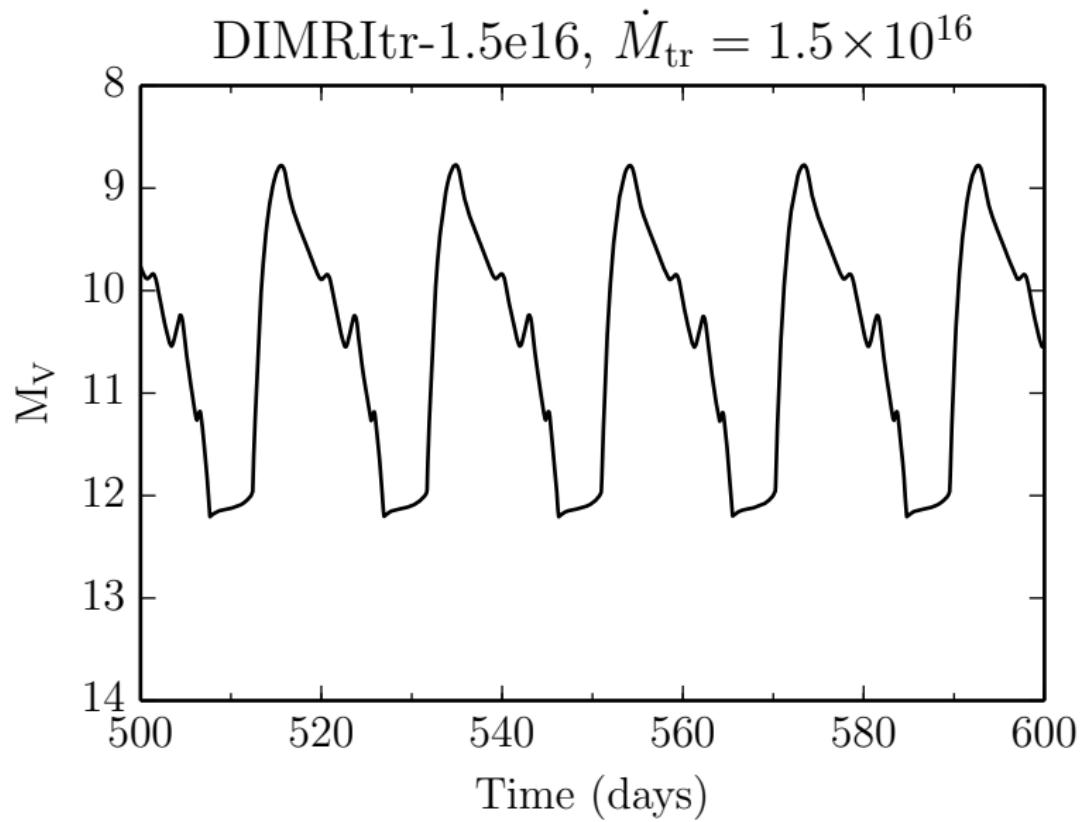
Vertical Structure Comparison



S-curves



Light Curve



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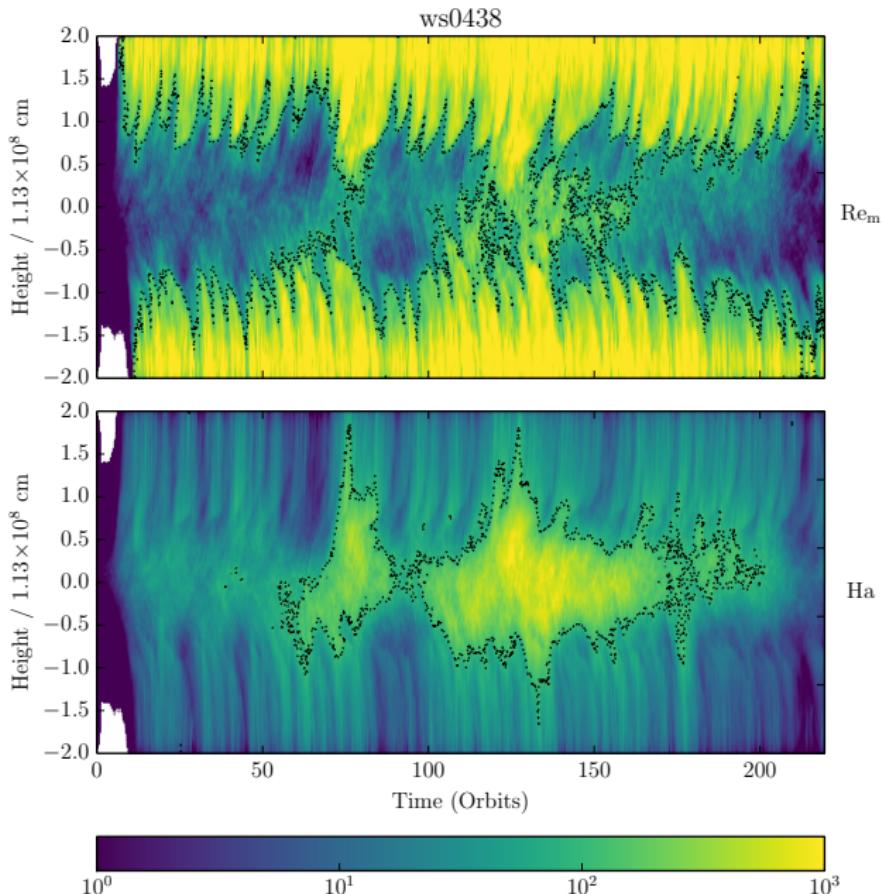
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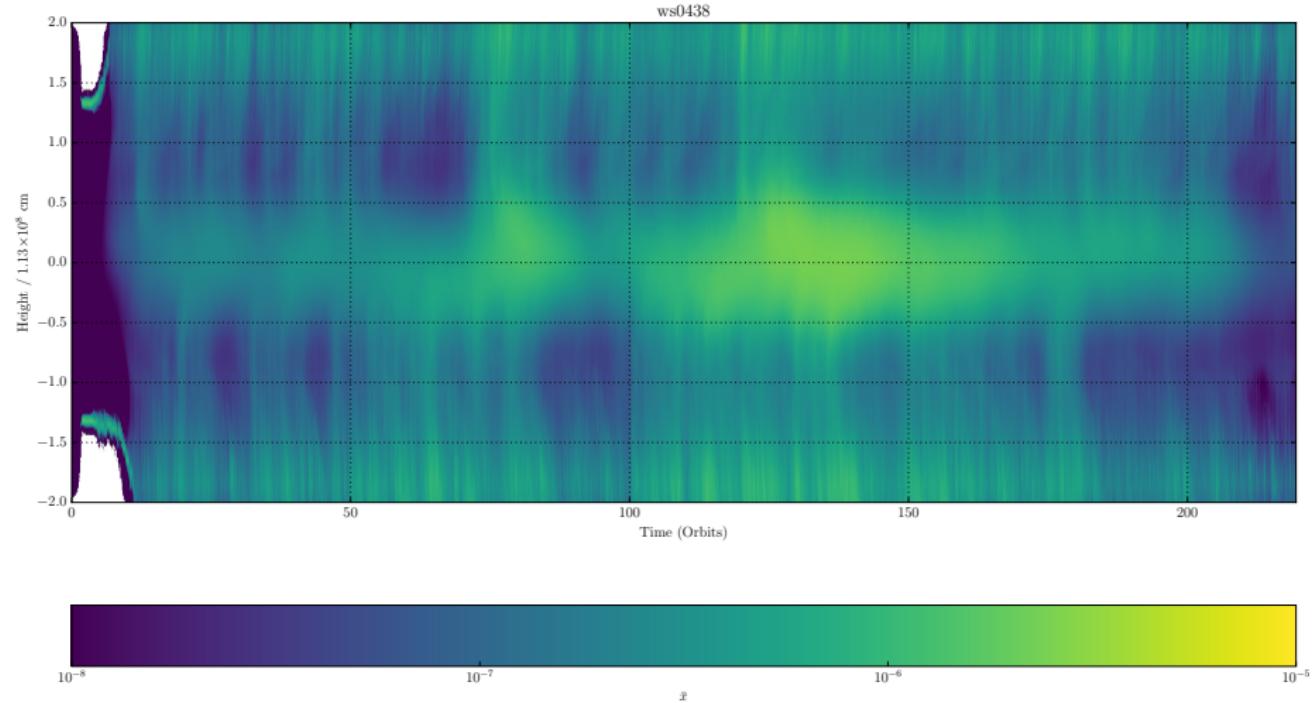
- DNe provide insight into accretion disks
 - Time variation
 - Well studied
- Convective accretion disks
 - α varies with $\alpha \sim .1$ at the tip of the upper branch
 - α highly correlated with convection
 - Convection Quenches Field Reversals
- Generated lightcurves
 - Modified DIM to incorporate local simulation data
 - Gives lightcurves (mostly) consistent with observations
 - Puts enhancement of α on stronger theoretical grounds
- Global simulations to come (Future work with Yan-Fei Jiang)
- Featured work
 - Coleman, M. S. B., Yerger, E., Blaes, O., Salvesen, G., & Hirose, S., 2017, MNRAS (in press), arXiv:1701.08177
 - Coleman, M. S. B., Kotko, I., Blaes, O., Lasota, J.-P., & Hirose, S., 2016, MNRAS, 462, 3710
 - Hirose, S. 2015, MNRAS, 448, 3105
 - Hirose, S., Blaes, O., Krolik, J. H., Coleman, M. S. B., & Sano, T., 2014, ApJ, 787, 1

Bonus Slides!

Non-ideal MHD

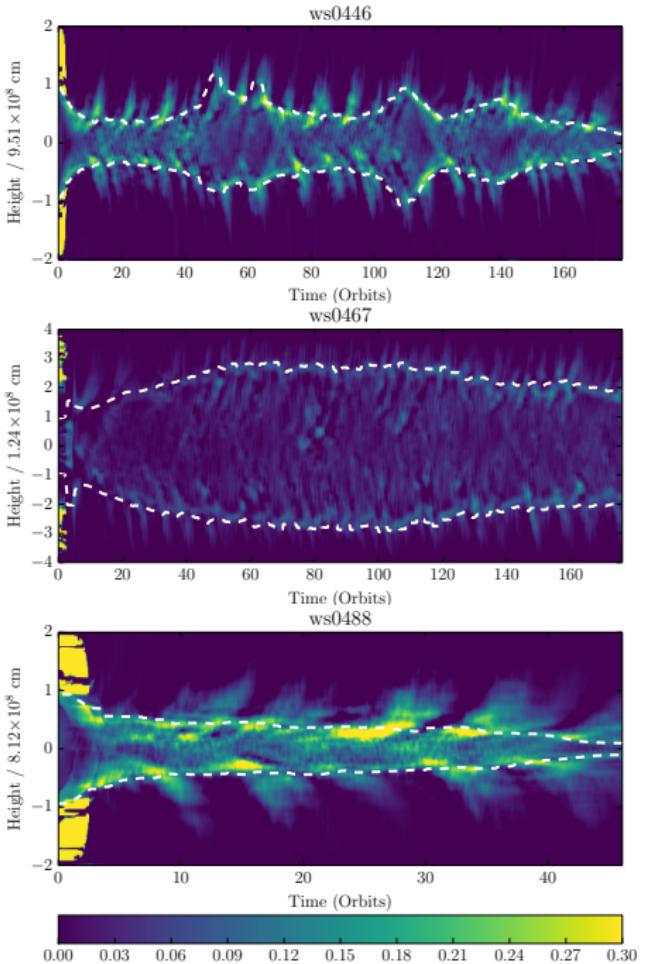


Ionization Fraction

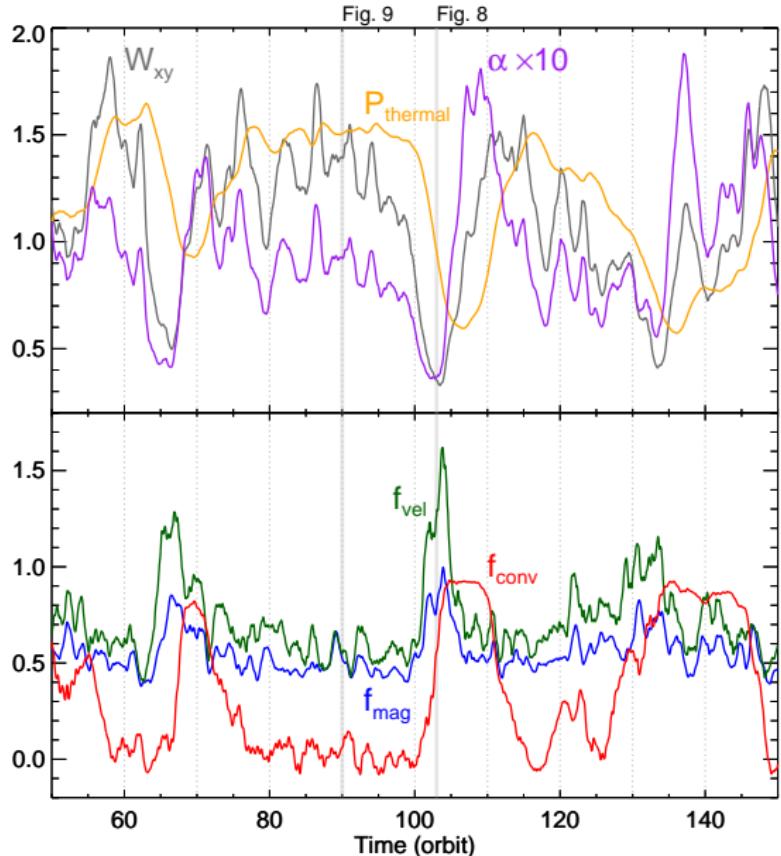


Non-Equilibrium

$$\alpha_{\text{eff}}(z, t) \equiv \frac{2}{3} \frac{1}{\Omega P_0} \frac{\partial F_z}{\partial z} \left(\frac{P_0}{P} \right)^{\delta}$$



Time Variation



$$f_{\text{mag}} \equiv \frac{\int [\langle B_z^2 \rangle] [\langle p_{\text{gas}} \rangle] dz}{\int [\langle B_r^2 \rangle] [\langle p_{\text{gas}} \rangle] dz}$$
$$f_{\text{vel}} \equiv \frac{\int [\langle v_z^2 \rangle] [\langle p_{\text{gas}} \rangle] dz}{\int [\langle v_r^2 \rangle] [\langle p_{\text{gas}} \rangle] dz}$$

Hirose et al. (2014)

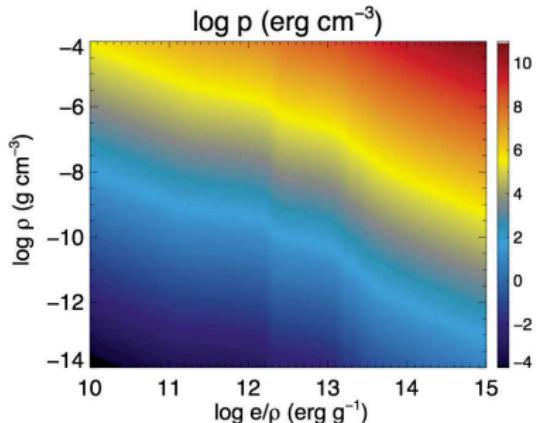
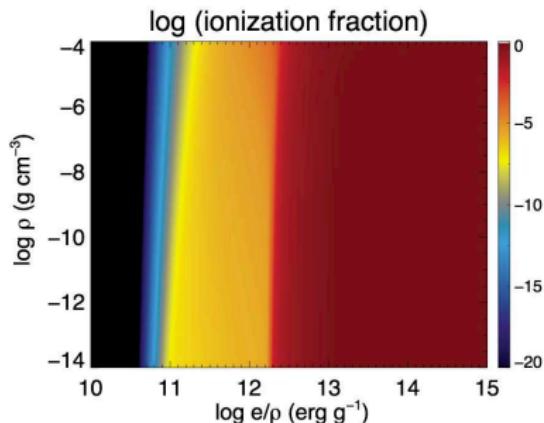
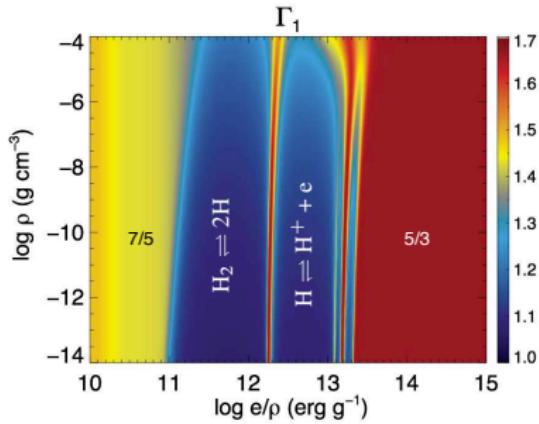
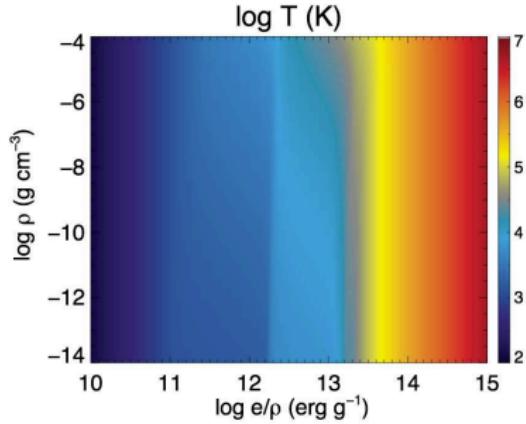
Equation of State

We wrote our own equation of state (EOS) code.
Preexisting codes lack alkali metals which ionize easily.

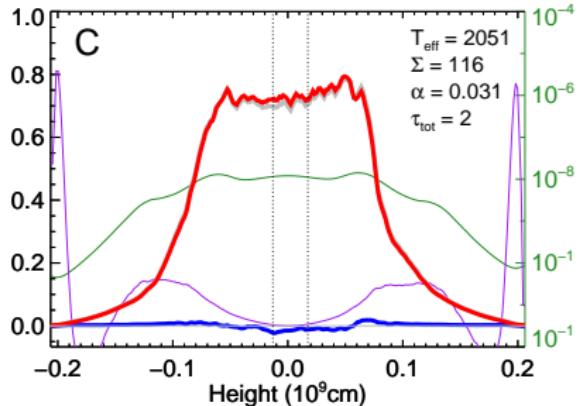
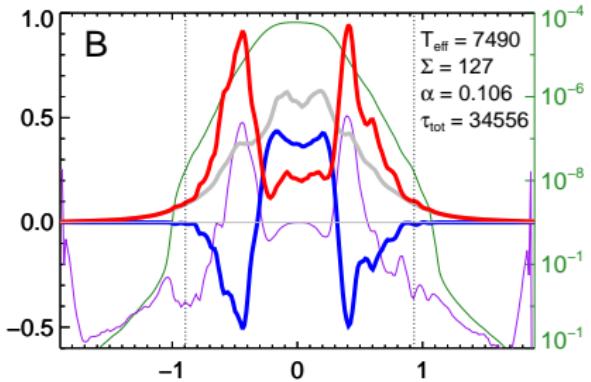
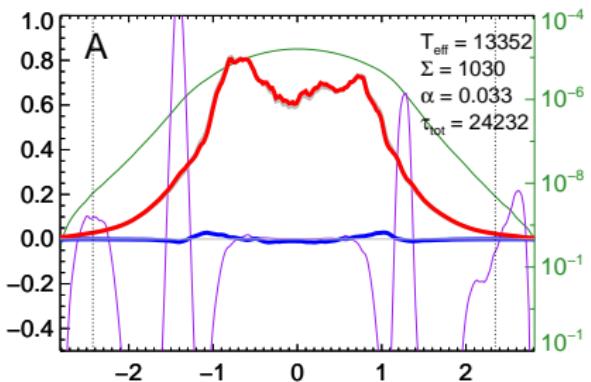
EOS properties

- Local thermal equilibrium
- Independent variables are density (ρ) and specific internal energy (e/ρ)
- Hydrogen species: H^- , H , H^+ , H_2 , H_2^+
- Helium species: He , He^+ , He^{++}
- Metals (neutral and singly ionized only): C, N, O, Ne, Na, Mg, Al, Si, P, S, Cl, Ar, K, Ca, Cr, Mn, Fe, Ni
- Solar metallicity

Equation of State Cont.

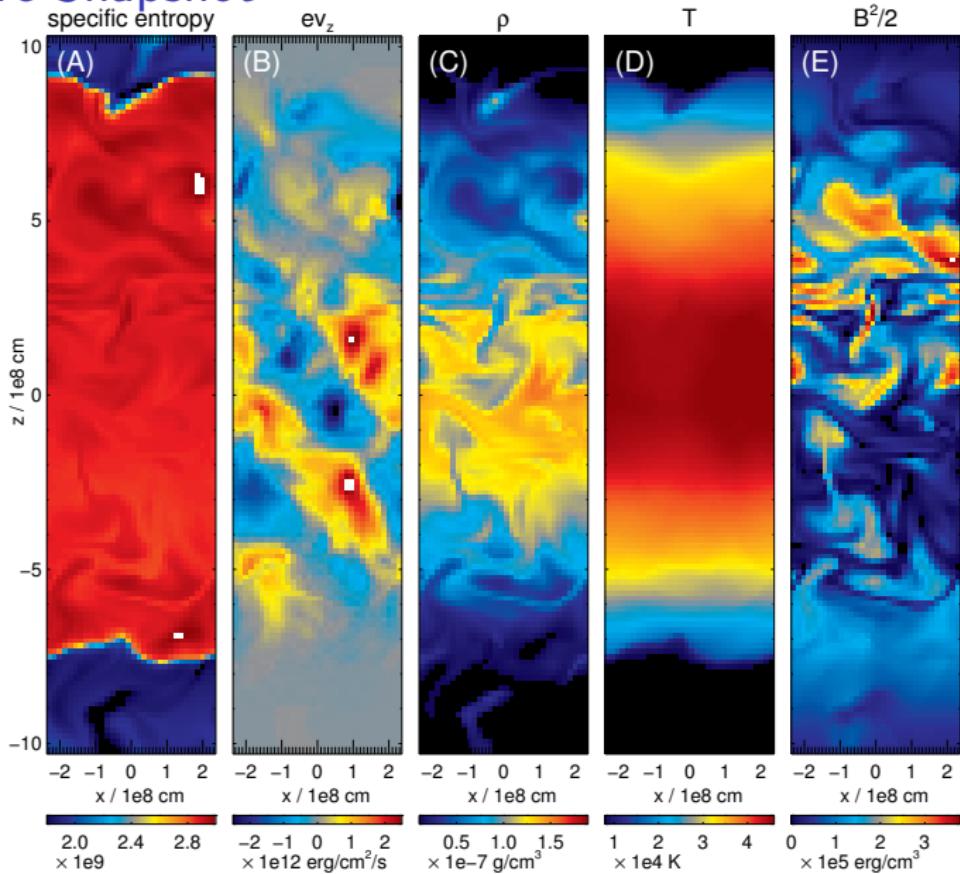


Evidence for True Thermal Convection

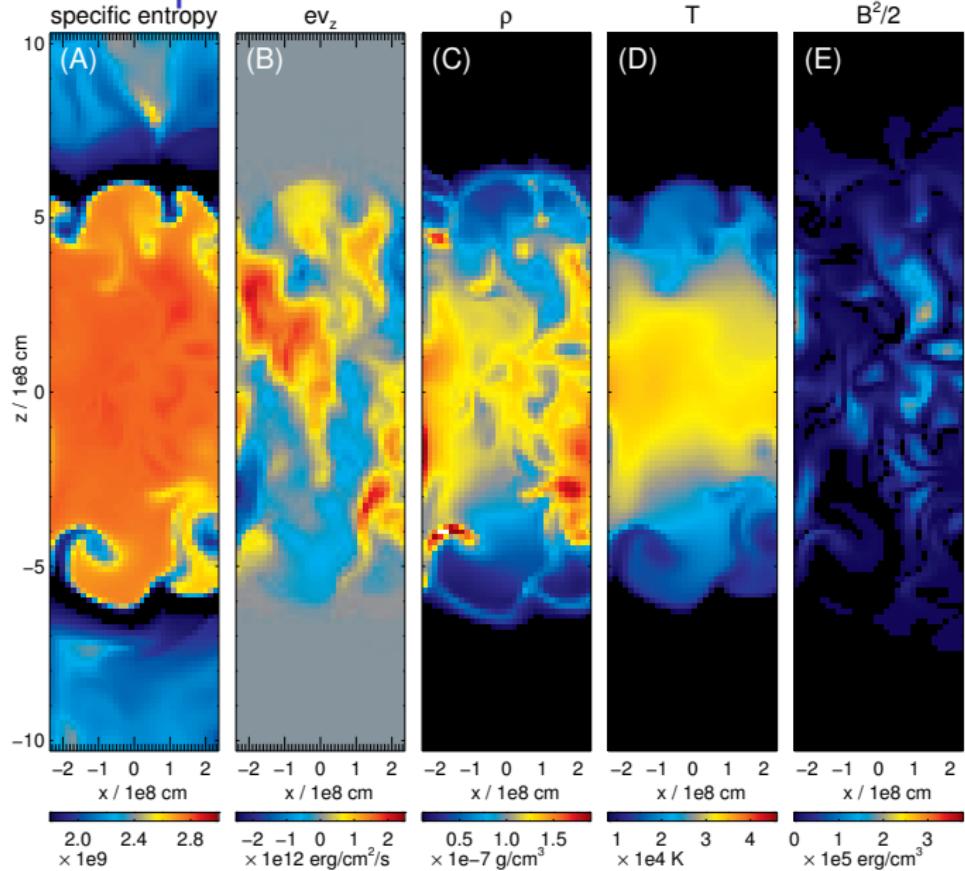


- Dissipated energy \bar{Q}^+
Normalized by 2000, 500, and 15
($\text{erg cm}^{-3} \text{s}^{-1}$)
- Radiative diffusion \bar{Q}_{rad}^-
- Gas advection \bar{Q}_{rad}^-
- Brunt-Väisälä freq. $0.1 \frac{N^2}{\Omega^2}$
- Opacity $\kappa_R \rho$ in cm^{-1}

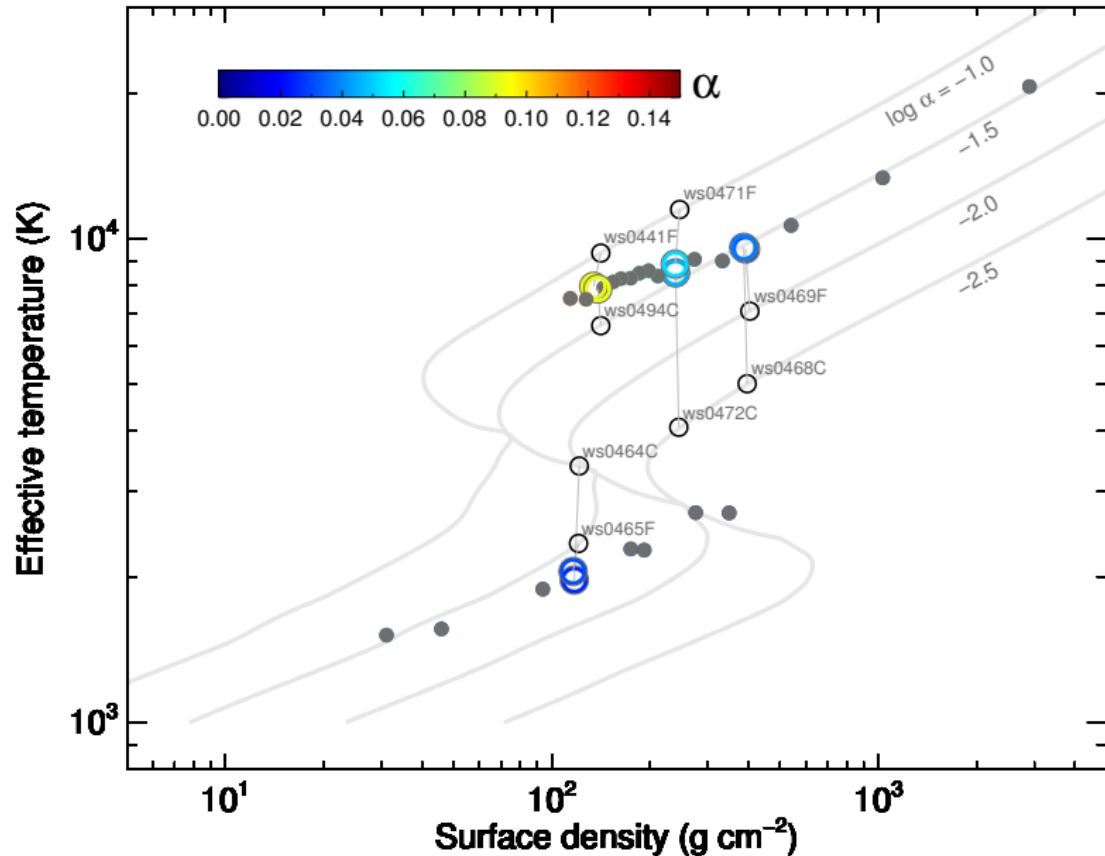
Radiative Snapshot



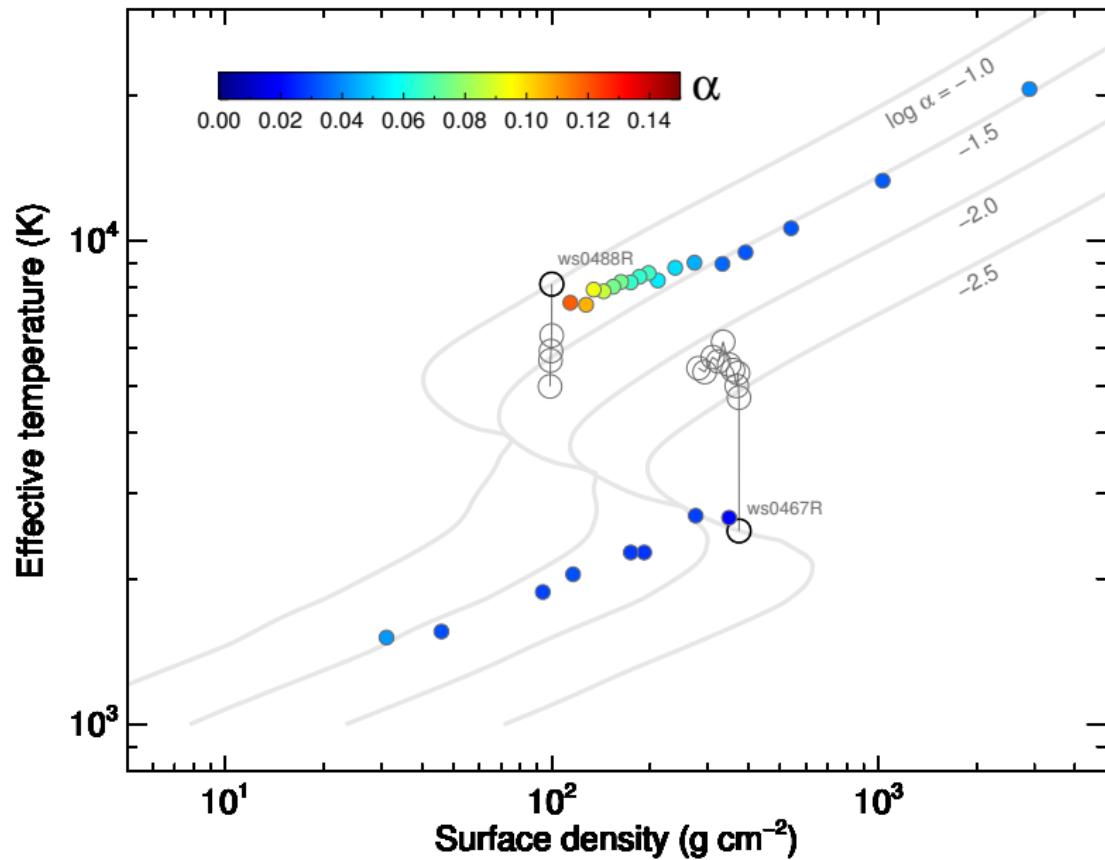
Convective Snapshot



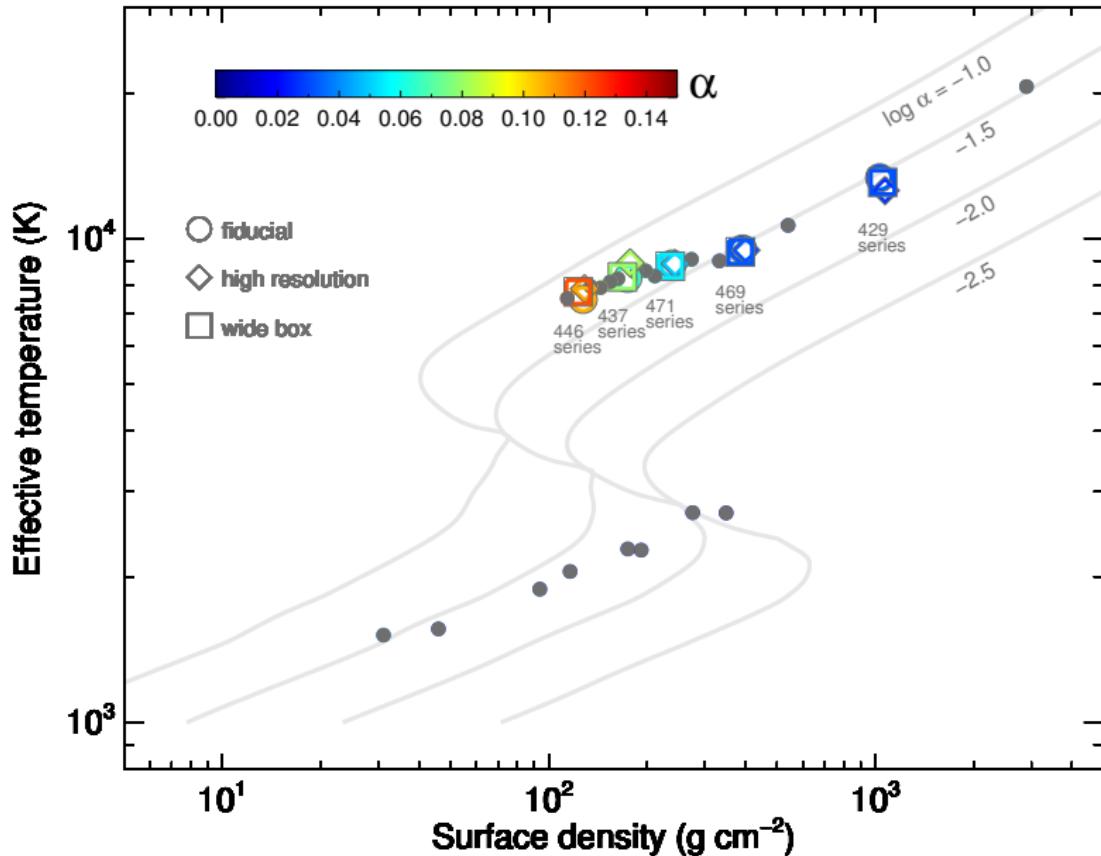
Branches are Attractors



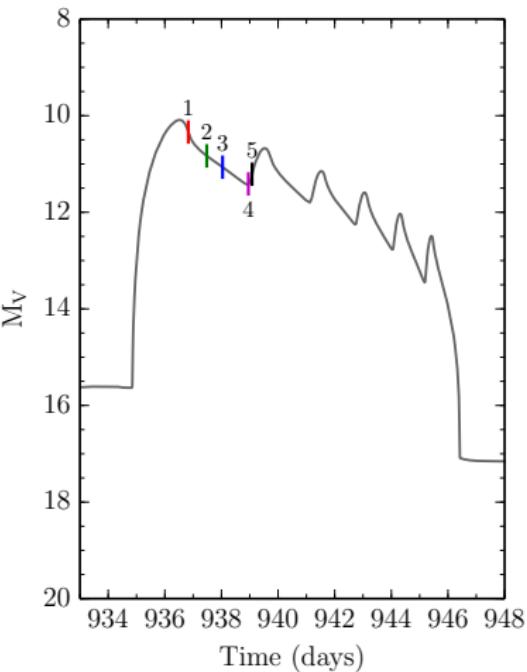
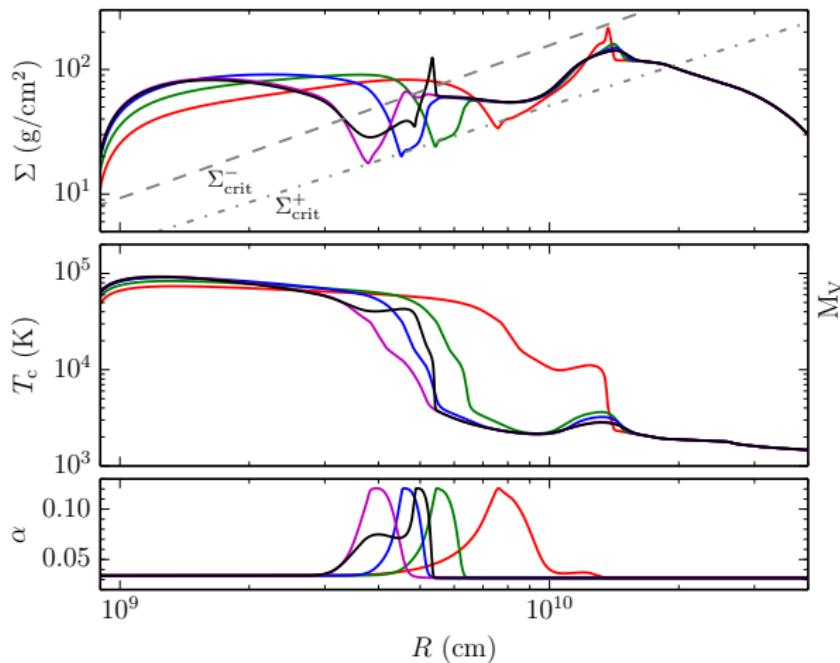
Branches have Ends



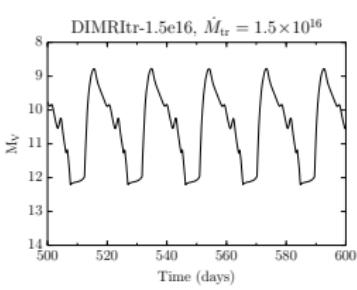
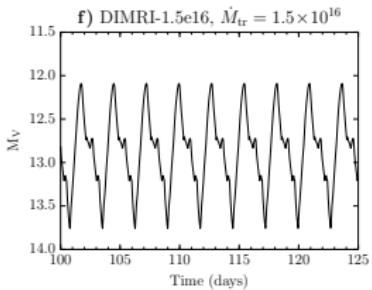
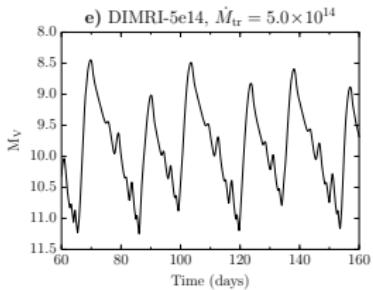
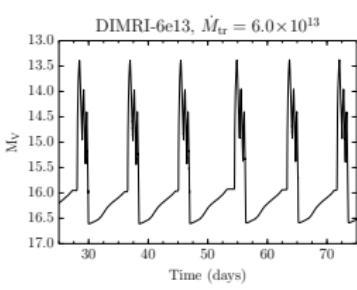
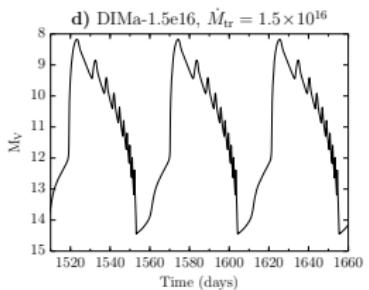
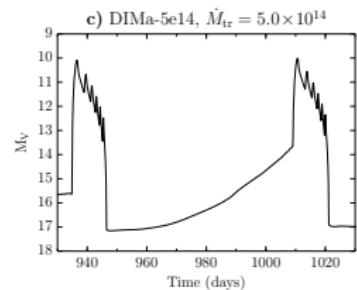
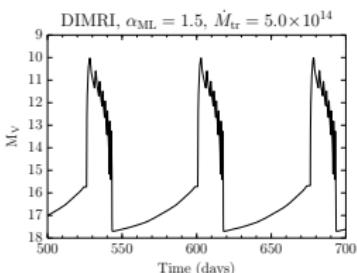
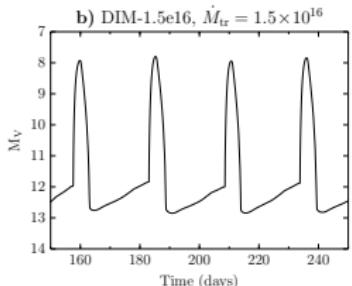
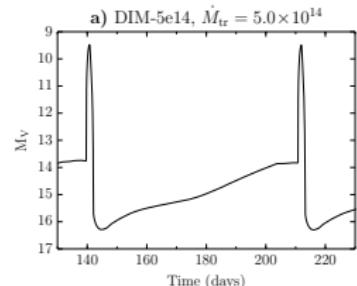
Equilibria are Robust to Numerics



Reflares



Light Curves Continued



Work Cited

- Coleman, M. S. B., Kotko, I., Blaes, O., Lasota, J.-P., & Hirose, S. 2016, MNRAS, 462, 3710
- Coleman, M. S. B., Yerger, E., Blaes, O., Salvesen, G., & Hirose, S. 2017, MNRAS (in press), arXiv:1701.08177
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- Shakura, N. I., & Sunyaev, R. A. 1973, , 24, 337
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