

Quantum Information, Tensor Networks & MBL

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D. Huse (Princeton), N. Yao (Berkeley), S. Choi, M. Lukin (Harvard)



Max Planck Institut
of Quantum Optics
(Garching)

KITP September 2018

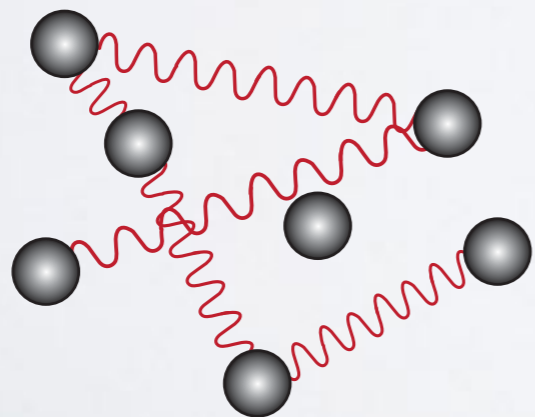
What are TNS?

- TNS = Tensor Network States

A general state of the N -body Hilbert space has exponentially many coefficients

$$|\Psi\rangle = \sum_{i_j} c_{i_1 \dots i_N} |i_1 \dots i_N\rangle$$

N

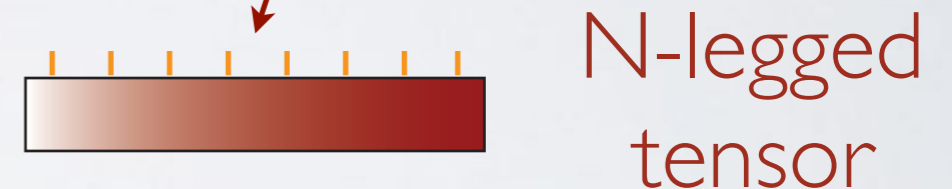


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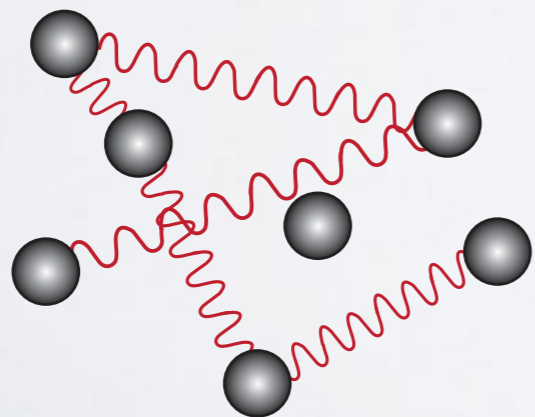
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$$d^N$$

N

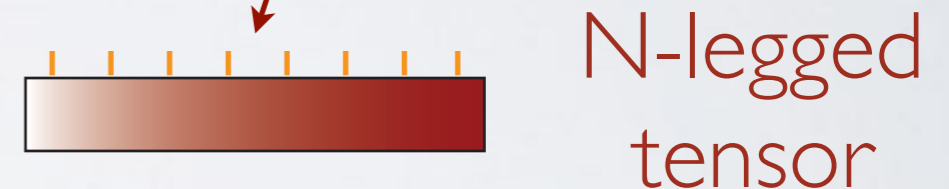


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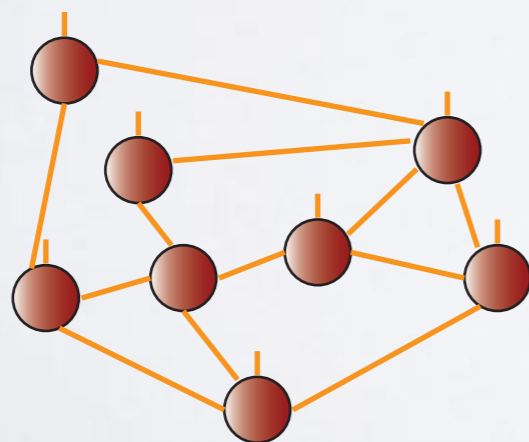
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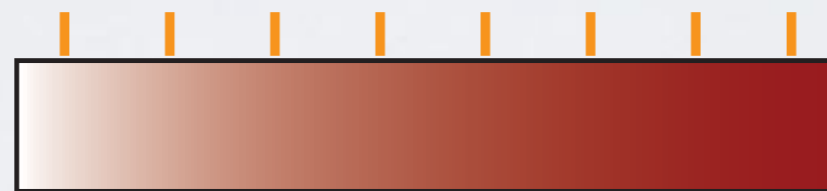
A TNS has only a polynomial number of parameters

$\text{poly}(N)$



Paradigmatic: MPS

- MPS = Matrix Product States



$$|\Psi\rangle = \sum_{i_1 \dots i_N} c_{i_1 \dots i_N} |i_1 \dots i_N\rangle$$

Affleck, Kennedy, Lieb, Tasaki, PRL 1987

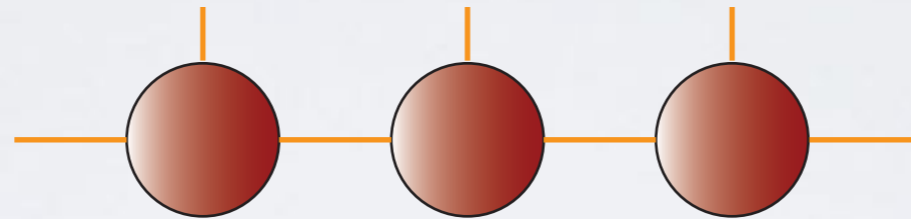
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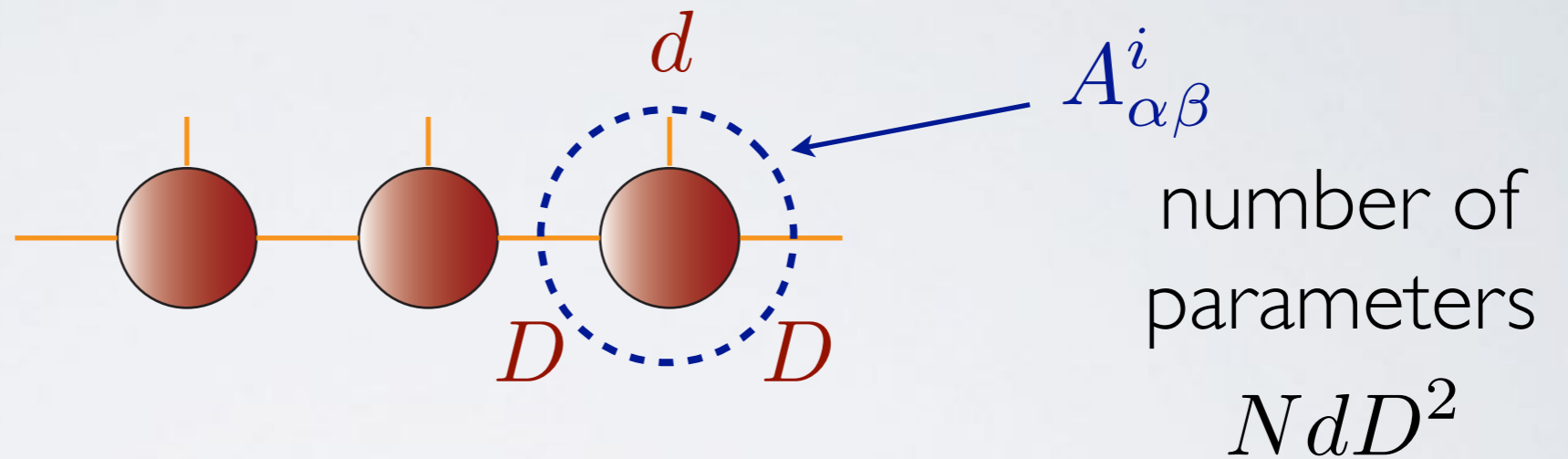
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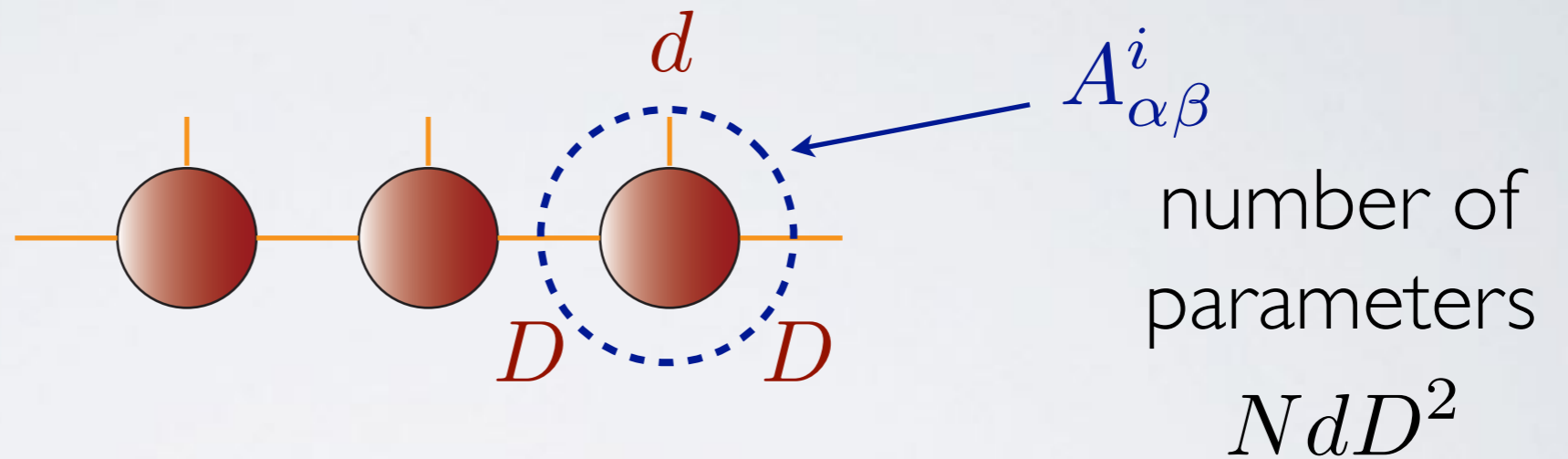
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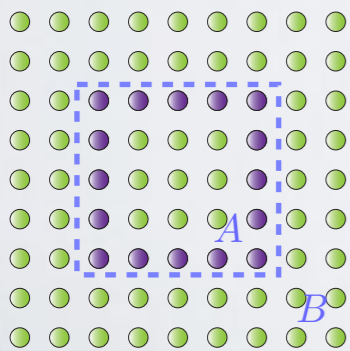
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Area law by construction

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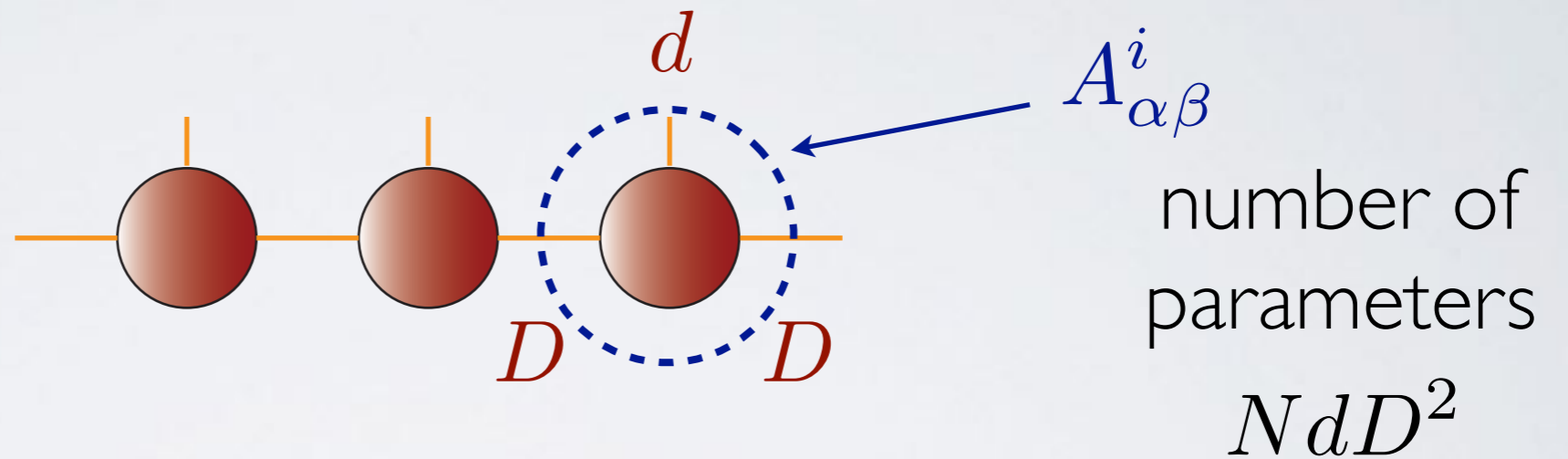
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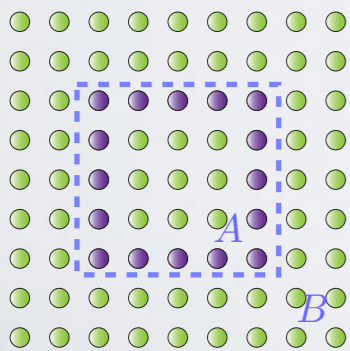
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Area law by construction

Bounded entanglement

$$S(L/2) \leq \log D$$

Affleck, Kennedy, Lieb, Tasaki, PRL 1987

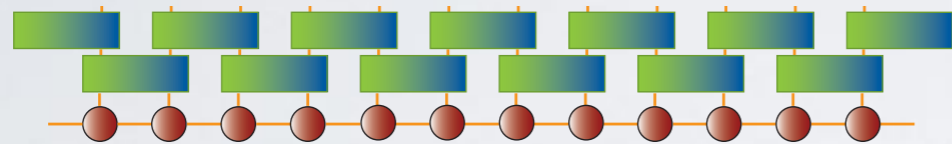
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time evolution with MPS

evolving the (pure state) ansatz



entanglement can grow fast!

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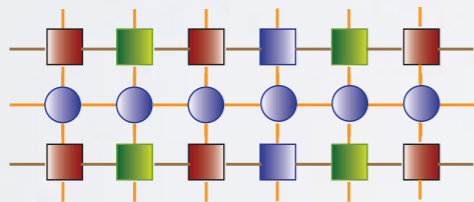


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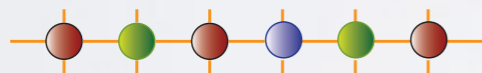


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also for mixed states

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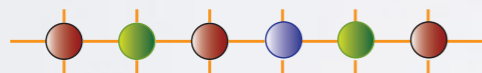


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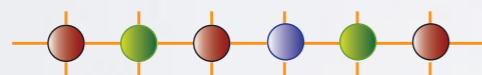


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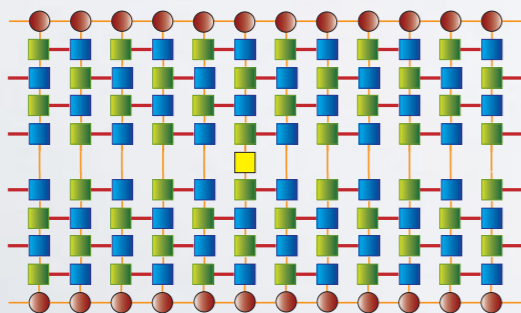


also for mixed states

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Prosen Pizorn, PRL 2008

observables as TN to contract



MCB, Hastings, Verstraete, Cirac, PRL 2009
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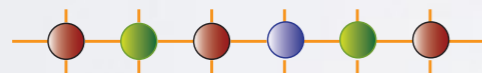


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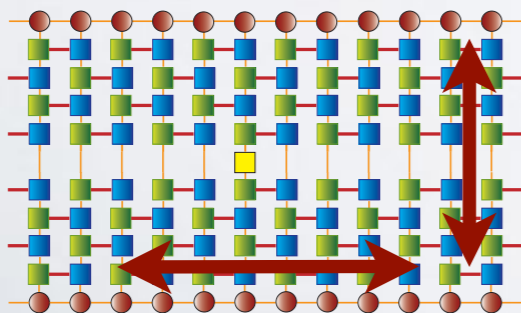
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different *entanglement* quantities



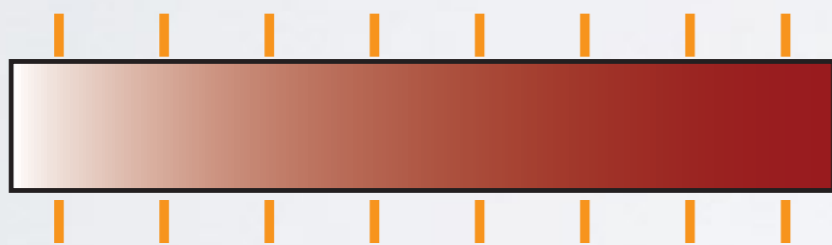
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TNS: mixed states & evolution

$$\rho = \sum_{i_1, j_1 \dots i_N, j_N} C^{i_1 j_1 \dots i_N j_N} |i_1 \dots i_N\rangle \langle j_1 \dots j_N|$$

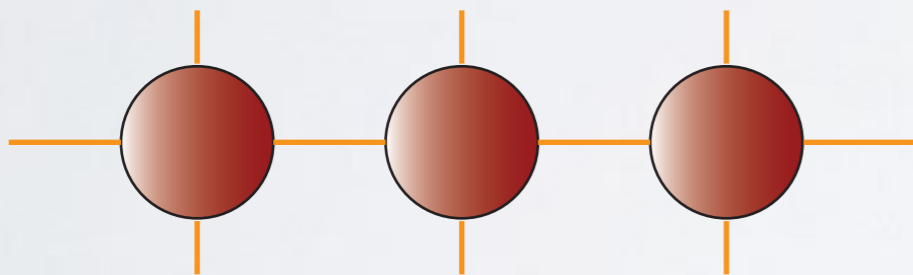
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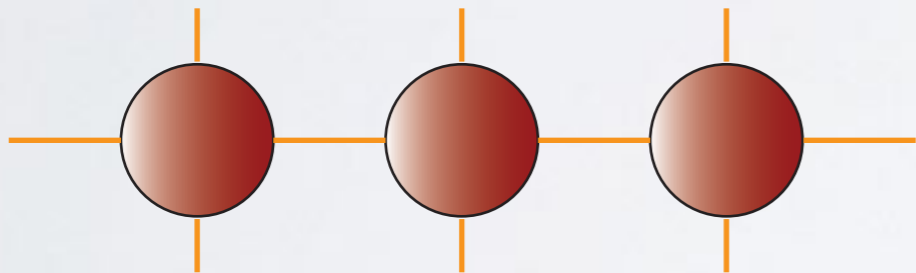
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can describe mixed states and operators
e.g. thermal equilibrium

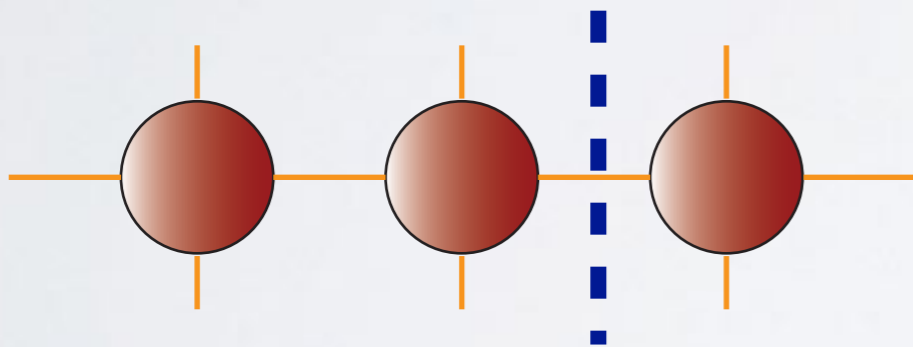


Verstraete et al., PRL 2004
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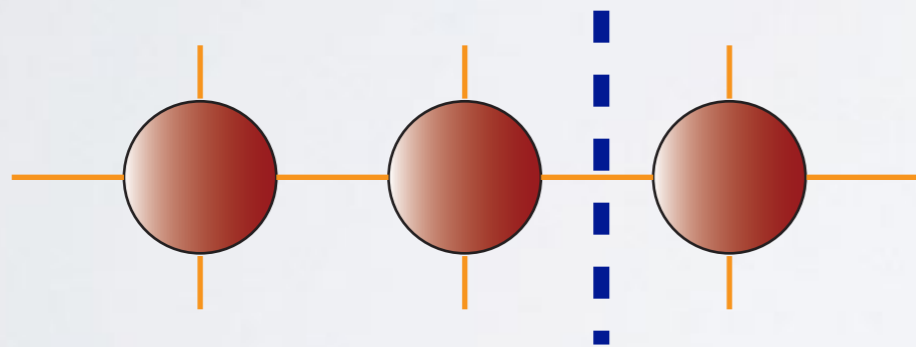
bounded operator space
entanglement entropy

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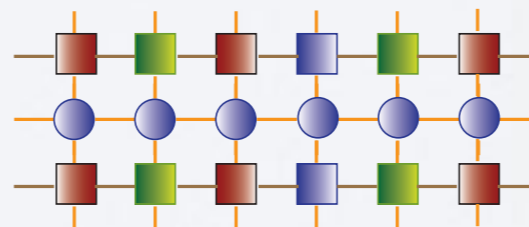
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time evolution: unitary and non-unitary

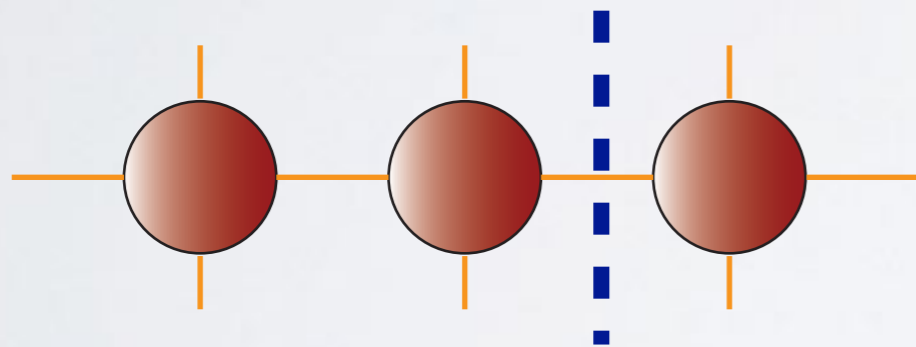


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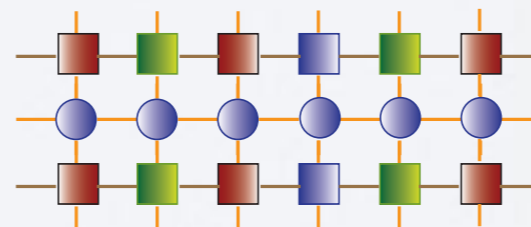
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time evolution: unitary and non-unitary



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long-time properties: slowest operators

Kim et al, PRE 2015

FOCUS: MBL dynamical scenario

interactions + strong disorder \Rightarrow localizing regime

absence of thermalization

Basko, Aleiner, Altshuler, Ann. Phys. 2006

Gornyi, Mirlin, Polyakov, PRL 2005

Altman, Vosk, Ann.Rev.CM 2015

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here: some quantum information + TN perspectives

Some questions we are asking

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propagation of correlations

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quantum memory features

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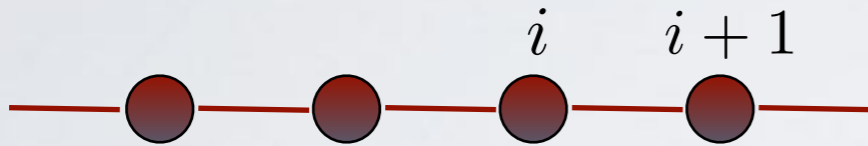
simulability with MPO

Hamiltonian
properties

local conserved quantities

the model

$$H = \sum \left(S_x^{[i]} S_x^{[i+1]} + S_y^{[i]} S_y^{[i+1]} + J S_z^{[i]} S_z^{[i+1]} + h_i S_i^z \right)$$



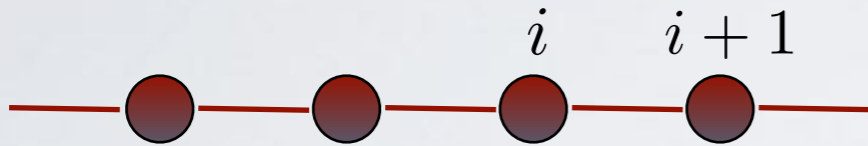
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$J=0 \Rightarrow$ non-interacting XY

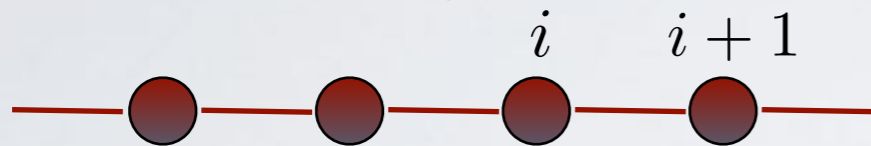
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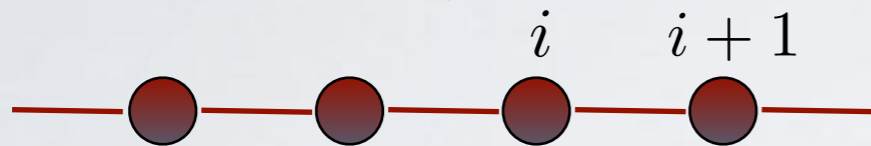
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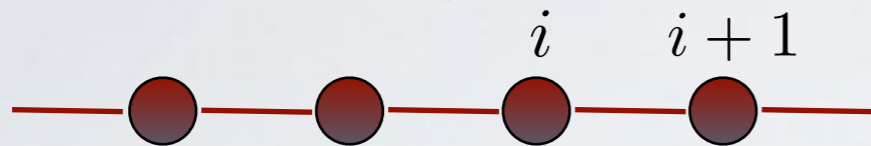
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$J=1 \Rightarrow$ shows MBL for $h \sim 3-3.5$

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initial states

mixed states at high T



$$\rho_{\Phi} \propto |\Phi\rangle\langle\Phi| \otimes Id^{\otimes \frac{L-L_0}{2}}$$

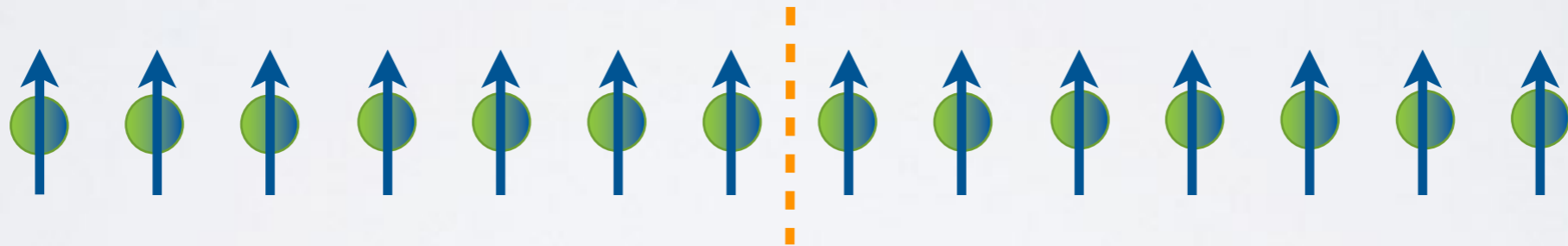
propagation of correlations

propagation of correlations

usual scenario: global quenches
for pure states

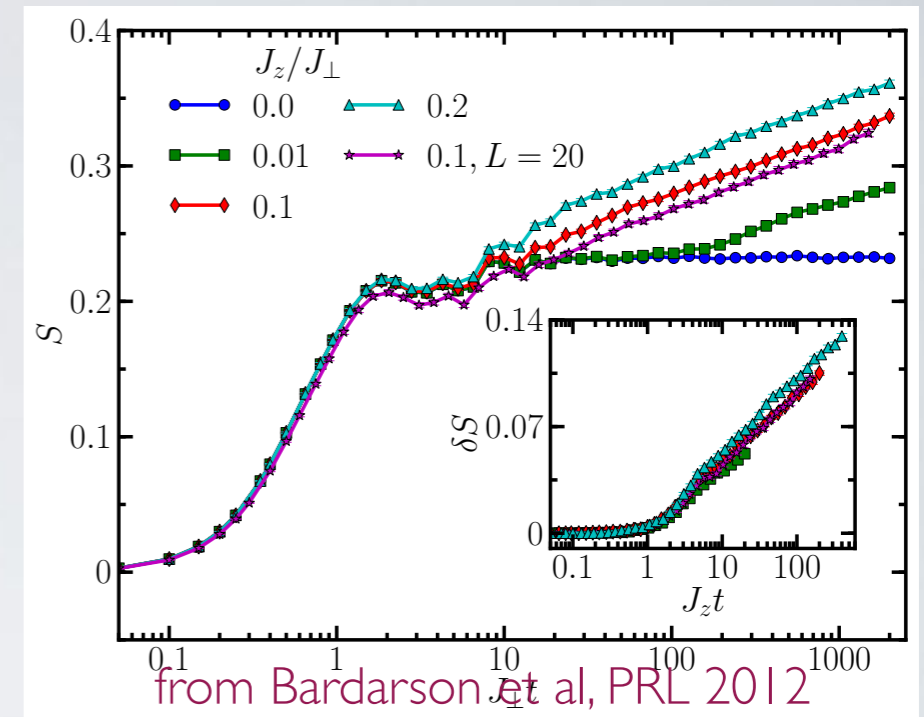
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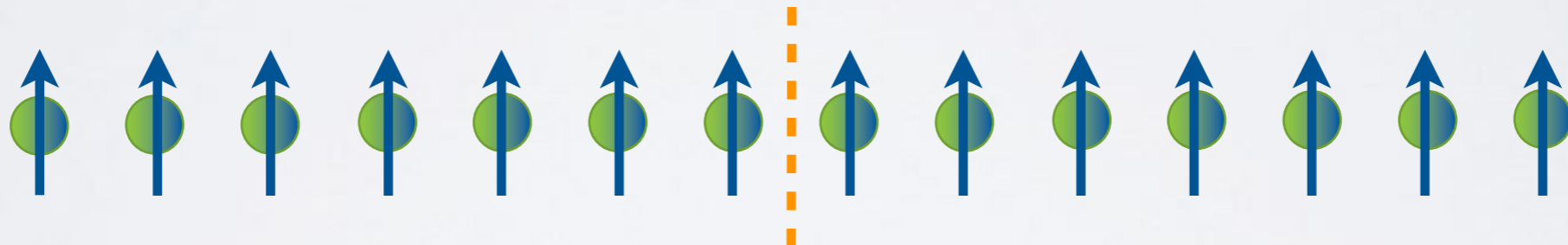
usual scenario: global quenches for pure states



single particle localization \rightarrow saturation of S

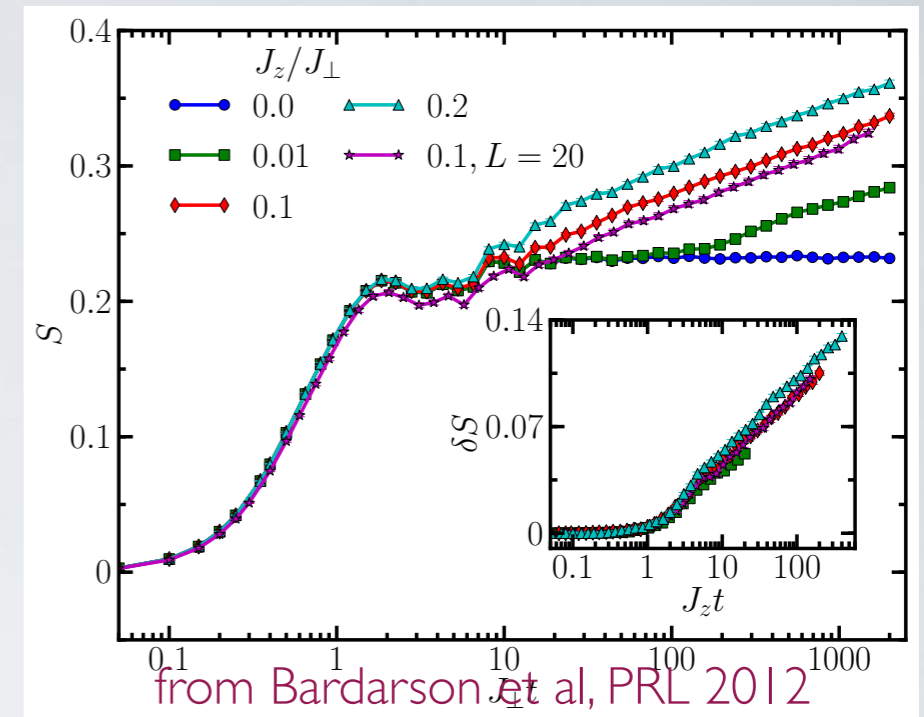
MBL \rightarrow logarithmic growth of S
explained by l-bit model

Bardarson et al, PRL 2012
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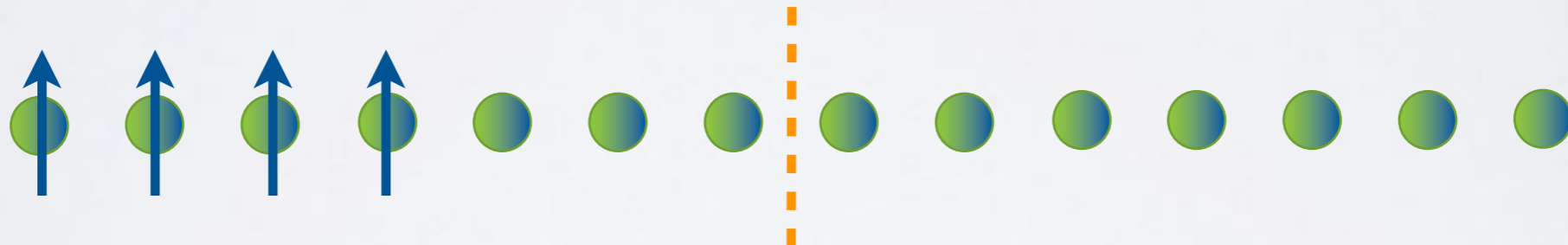
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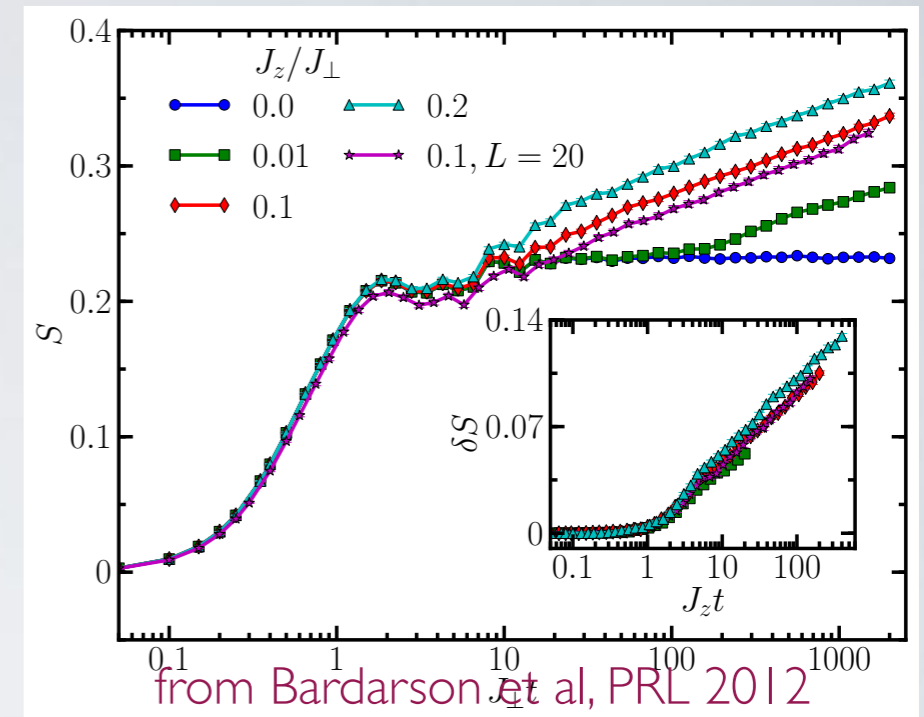
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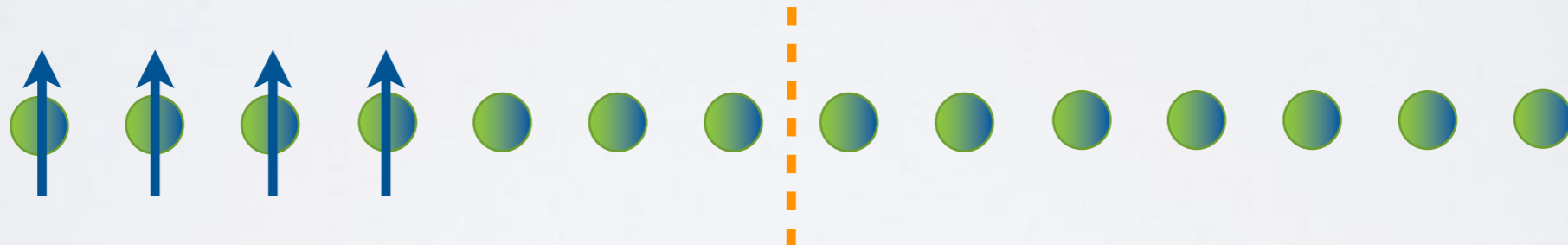
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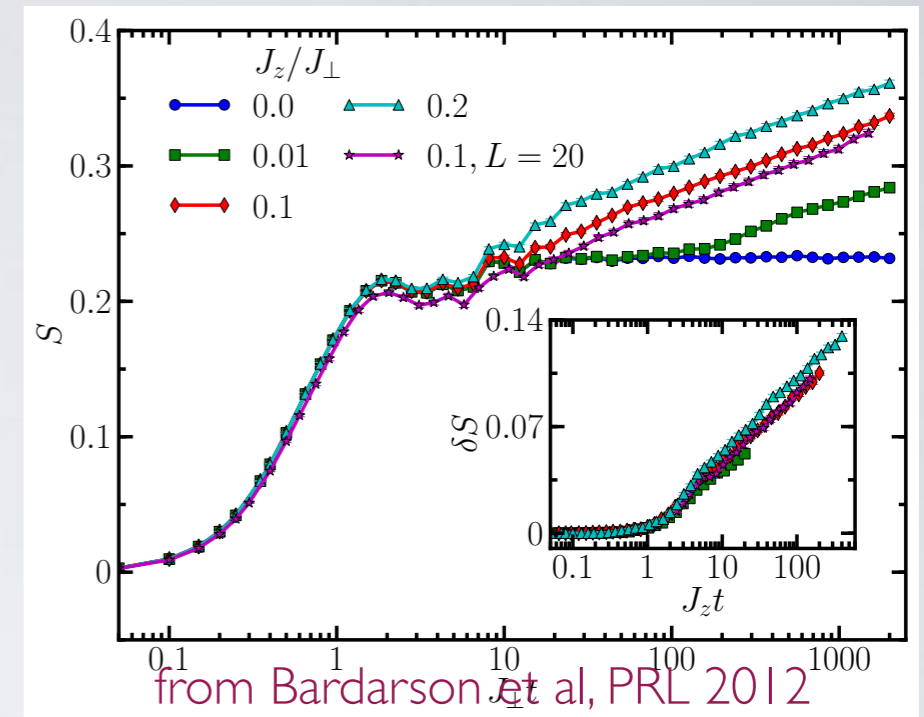
different for mixed states

$$I(\text{left } L_c \text{ sites} : \text{rest})$$

measures correlations between subsystems

propagation of correlations

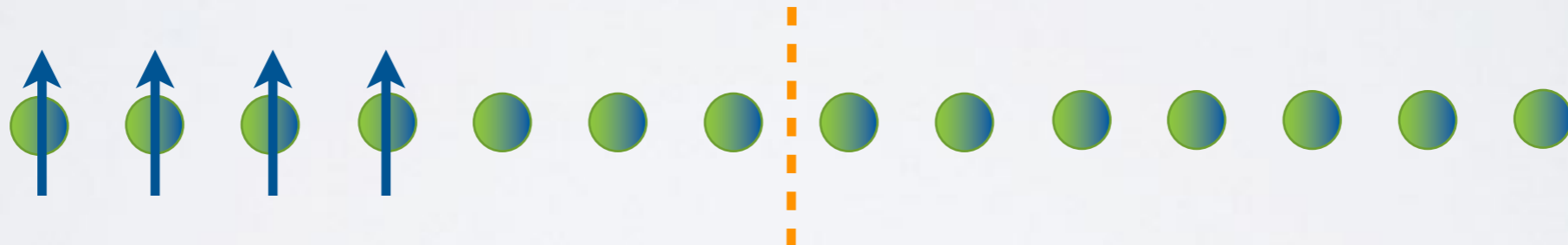
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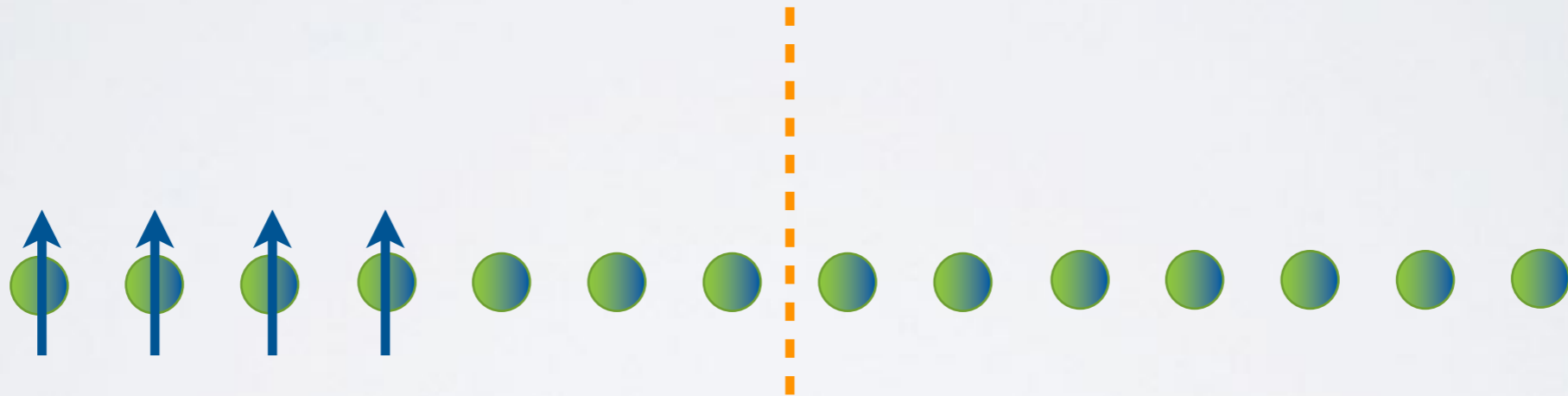
$$I(A : B) = S(A) + S(B) - S(AB)$$

measures correlations between subsystems

propagation of correlations

upper bounded

$$I(\ell : N - \ell) = S(\ell) + S(N - \ell) - S(N)$$



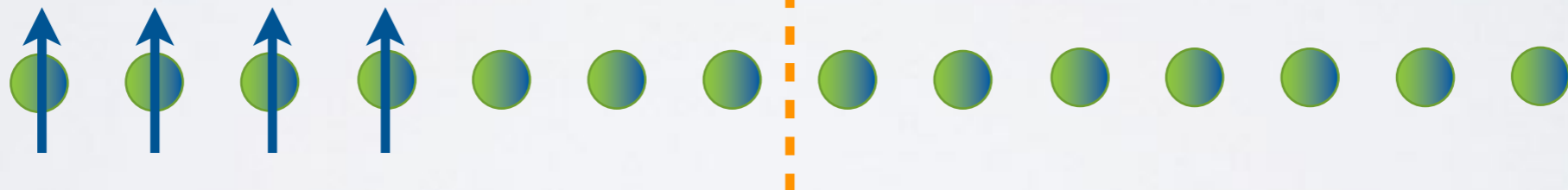
propagation of correlations

upper bounded

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$$\leq \ell$$

$$\leq N - \ell$$



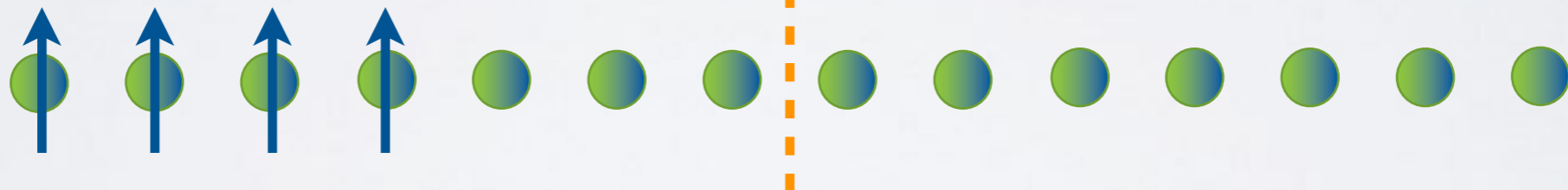
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initially

$$\rho_{\Phi} \propto |\Phi\rangle\langle\Phi| \otimes Id^{\otimes \frac{L-L_0}{2}}$$

$$L_0 = 1$$

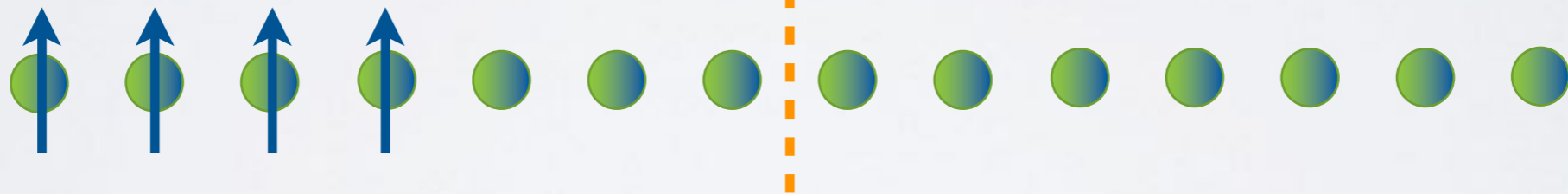
$$|\Phi\rangle = |Z_{\pm}\rangle, |X_{\pm}\rangle$$

propagation of correlations

upper bounded

$$I(\ell : N - \ell) = S(\ell) + S(N - \ell) - S(N)$$

$$\leq \ell \qquad \leq N - \ell \qquad = N - L_0$$



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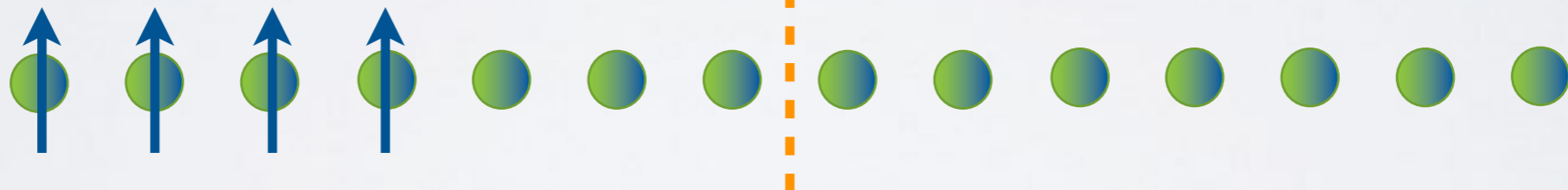
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$$\begin{array}{ccc} \leq \ell & \leq N - \ell & = N - L_0 \end{array}$$

yet shows difference between many-body and single particle localized

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exact calculation possible in the non-interacting and 1-bit models

propagation of correlations

non-interacting case: quadratic fermionic model

$$\rho_{\Phi} \propto |\Phi\rangle\langle\Phi| \otimes Id^{\otimes \frac{L-L_0}{2}} \quad |\Phi\rangle = |Z_{\pm}\rangle, |X_{\pm}\rangle$$

propagation of correlations

non-interacting case: quadratic fermionic model

exact evolution: single parameter $\mathcal{V}_\ell = \sum_{r=0}^{\ell-1} |\langle r|U(t)|0\rangle|^2$
probability to the left

$$\rho_\Phi \propto |\Phi\rangle\langle\Phi| \otimes Id^{\otimes \frac{L-L_0}{2}} \quad |\Phi\rangle = |Z\pm\rangle, |X\pm\rangle$$

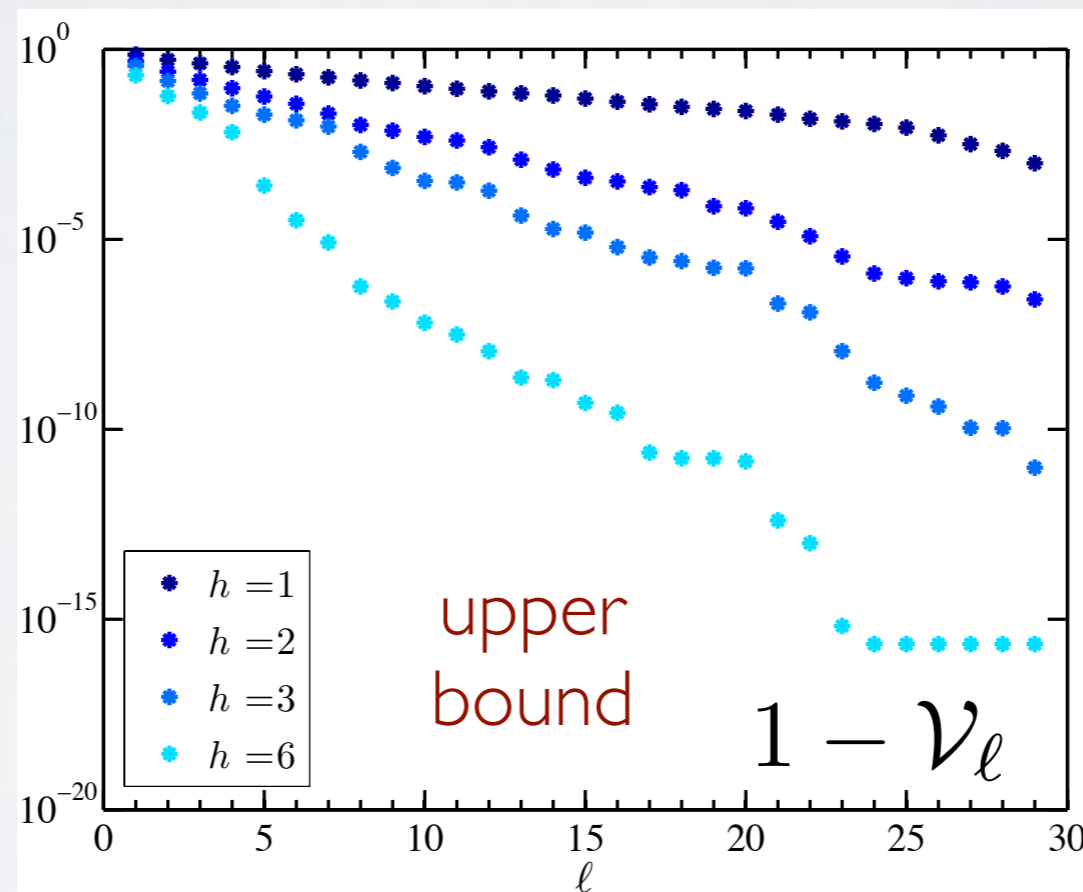
propagation of correlations

non-interacting case: quadratic fermionic model

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probability to
the left

probability to the
right



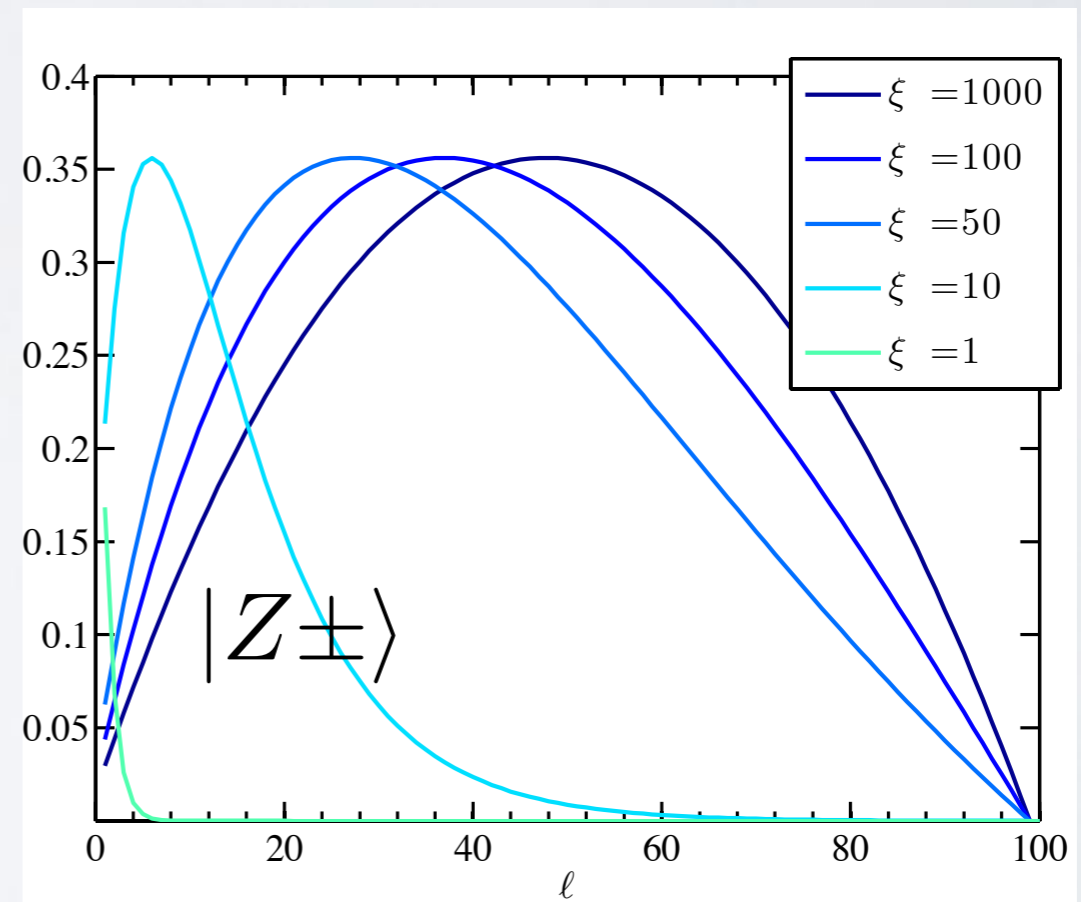
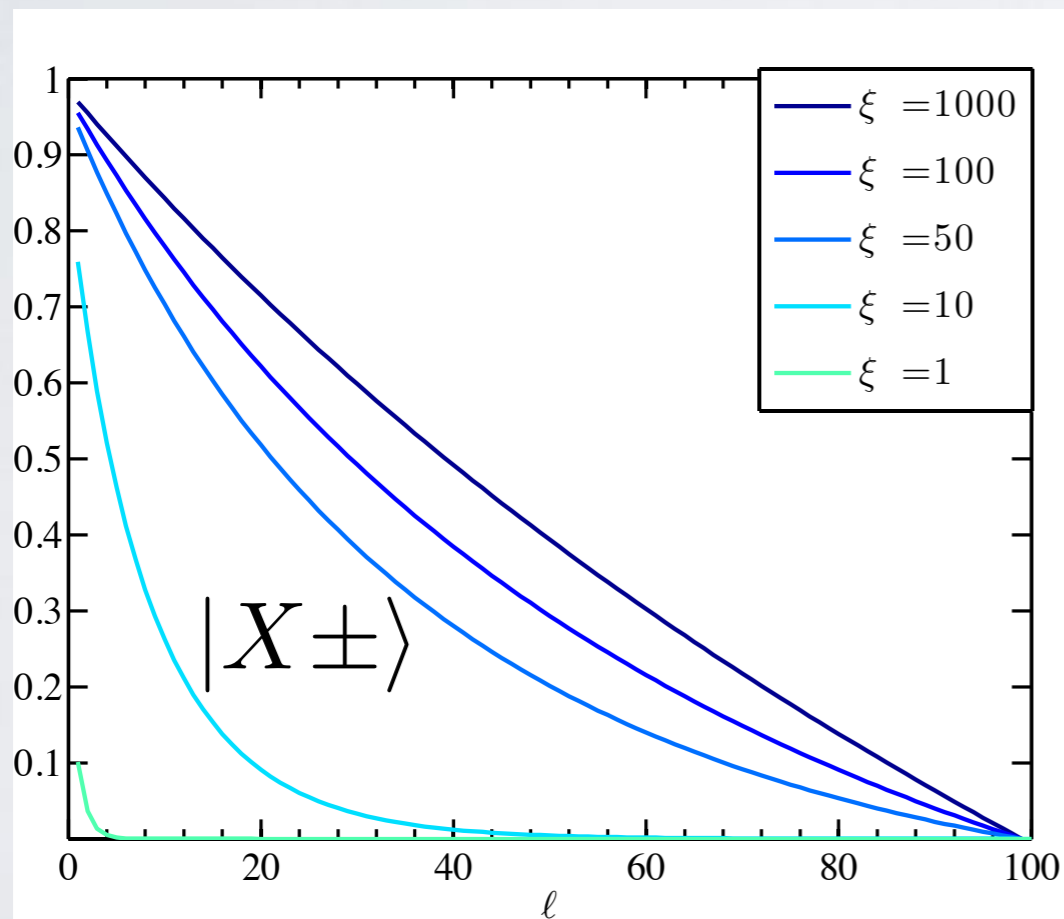
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propagation of correlations

non-interacting case: localization $h > 0$

asymptotic value of mutual information

assuming exponential decay with ξ



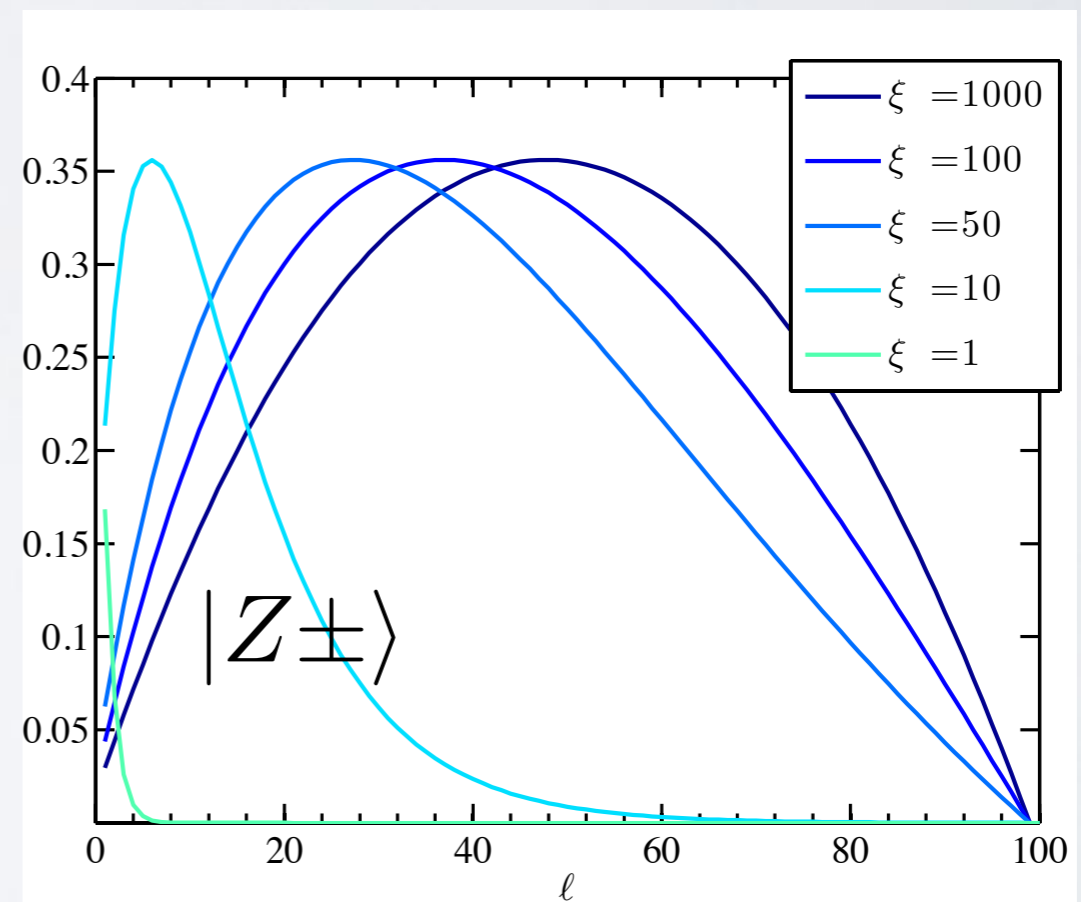
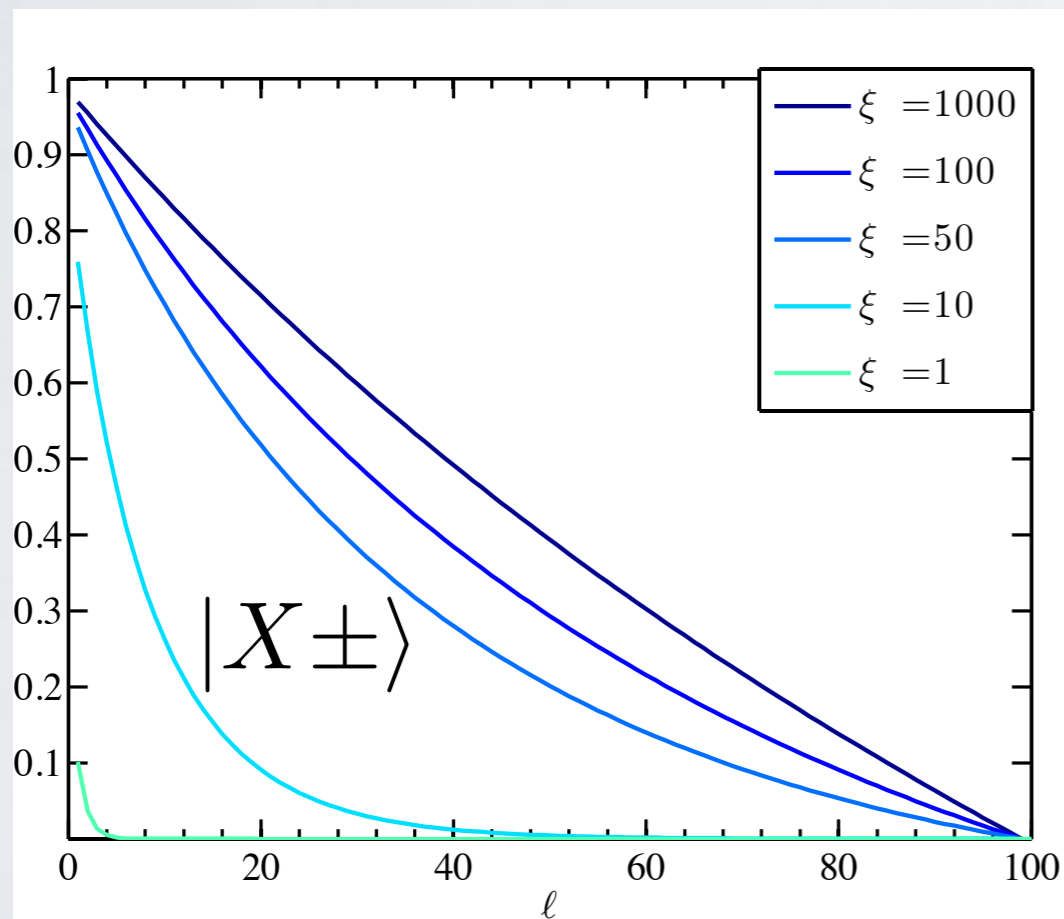
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propagation of correlations

non-interacting case: localization $h > 0$

$L = 50, h = 1$



$$\rho_{\Phi} \propto |\Phi\rangle\langle\Phi| \otimes Id^{\otimes \frac{L-L_0}{2}}$$

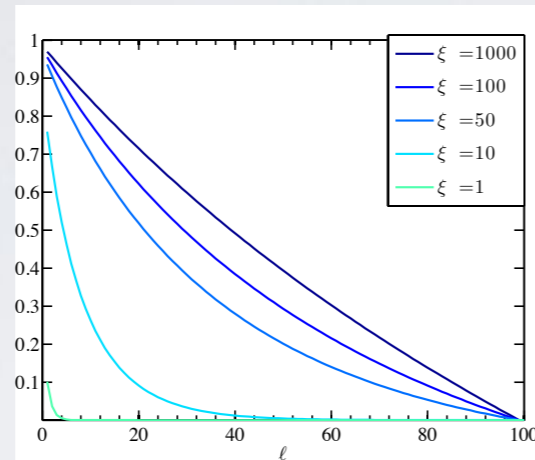
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propagation of correlations

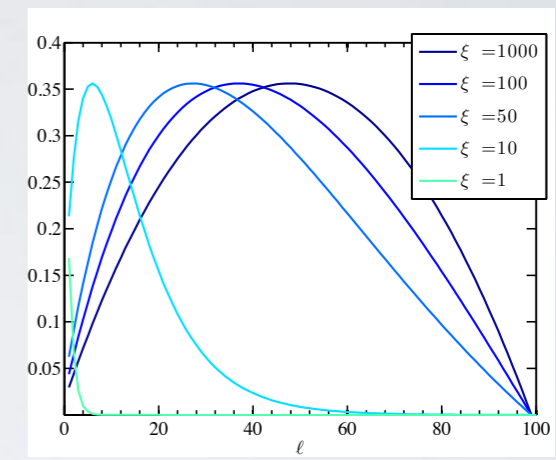
non-interacting case: localization $h > 0$

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$|X_{\pm}\rangle$



$|Z_{\pm}\rangle$



$$\rho_{\Phi} \propto |\Phi\rangle\langle\Phi| \otimes Id^{\otimes \frac{L-L_0}{2}}$$

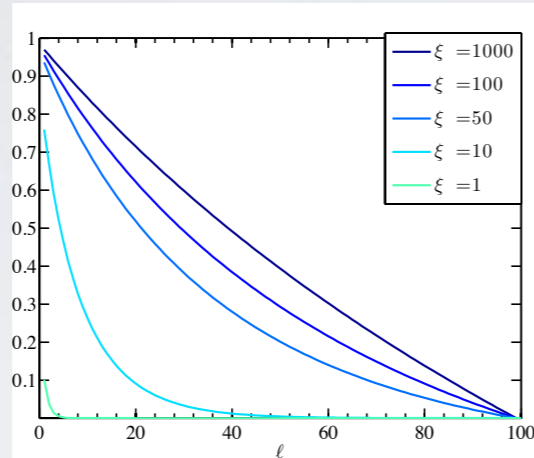
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propagation of correlations

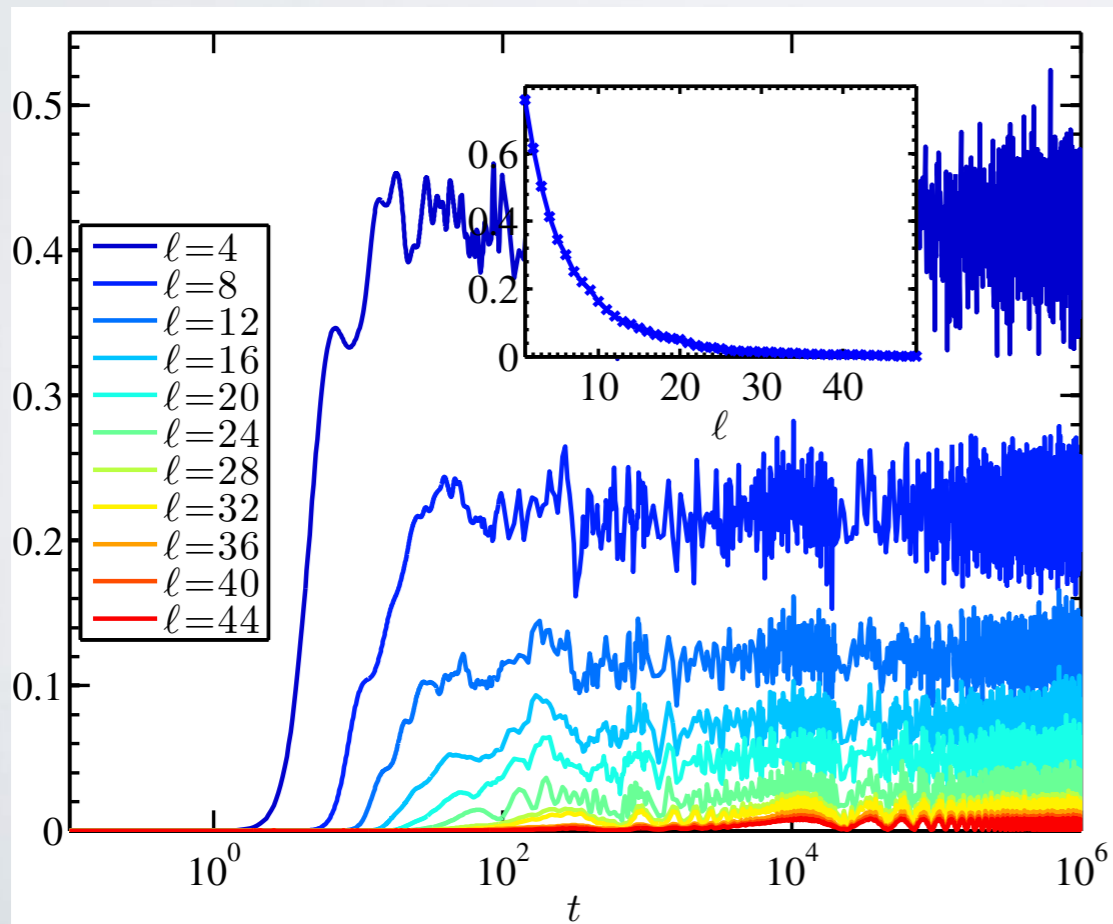
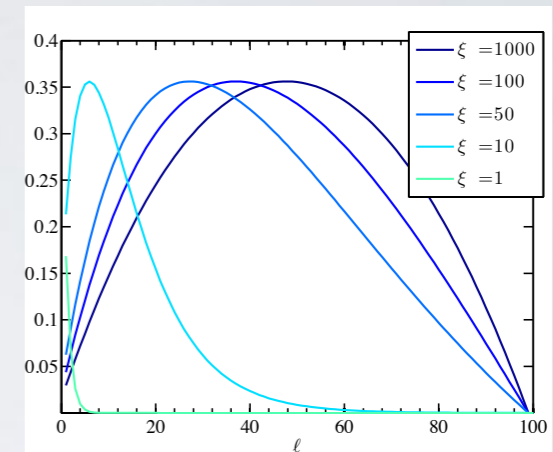
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$|X_{\pm}\rangle$



$|Z_{\pm}\rangle$



$$\frac{L - L_0}{2}$$

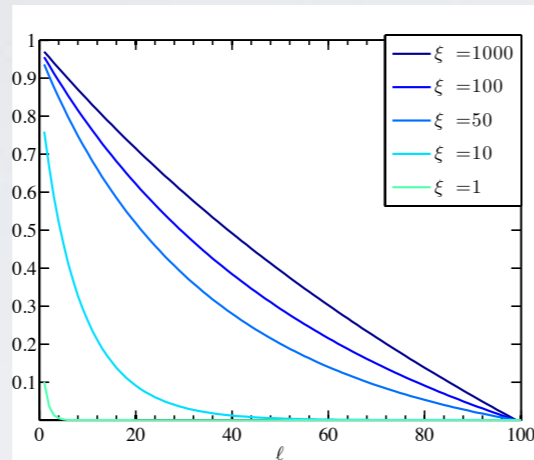
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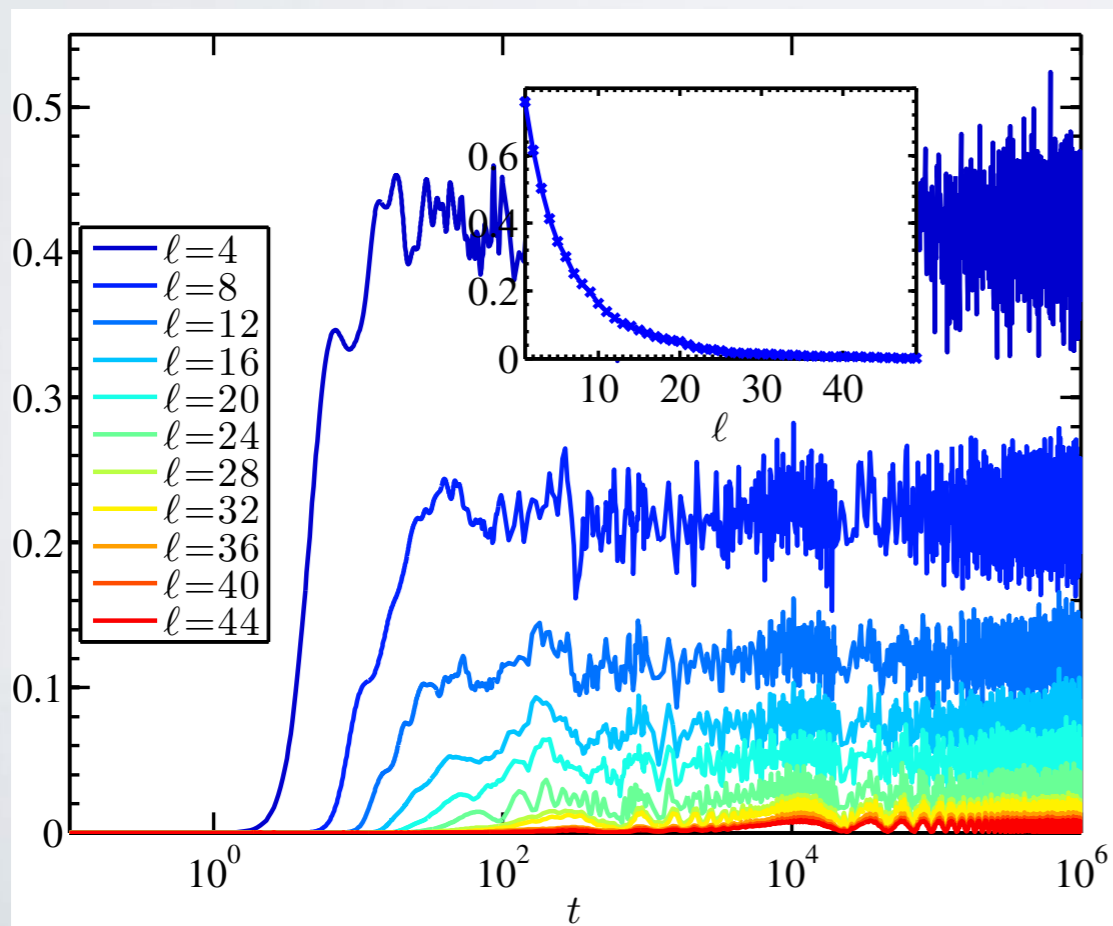
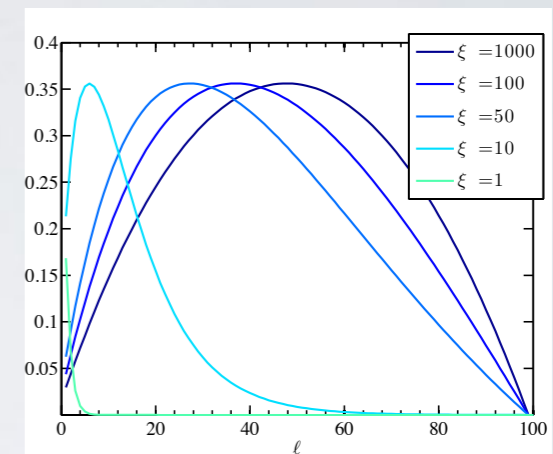
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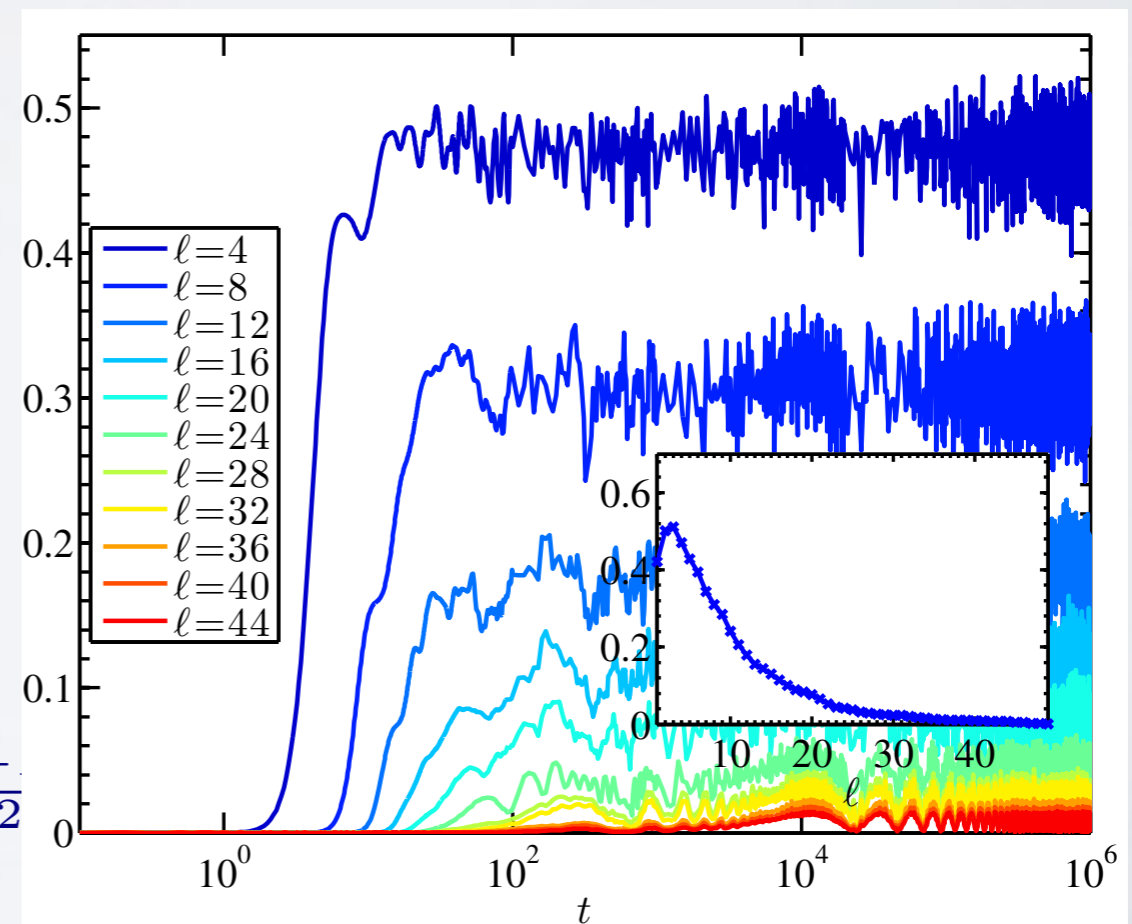
$|X_{\pm}\rangle$



$|Z_{\pm}\rangle$



$\frac{L-\ell}{2}$

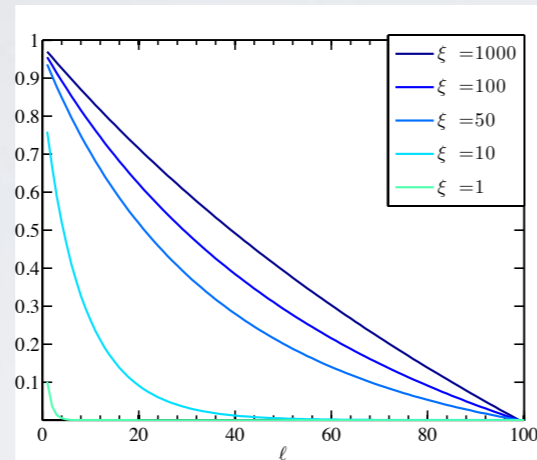


propagation of correlations

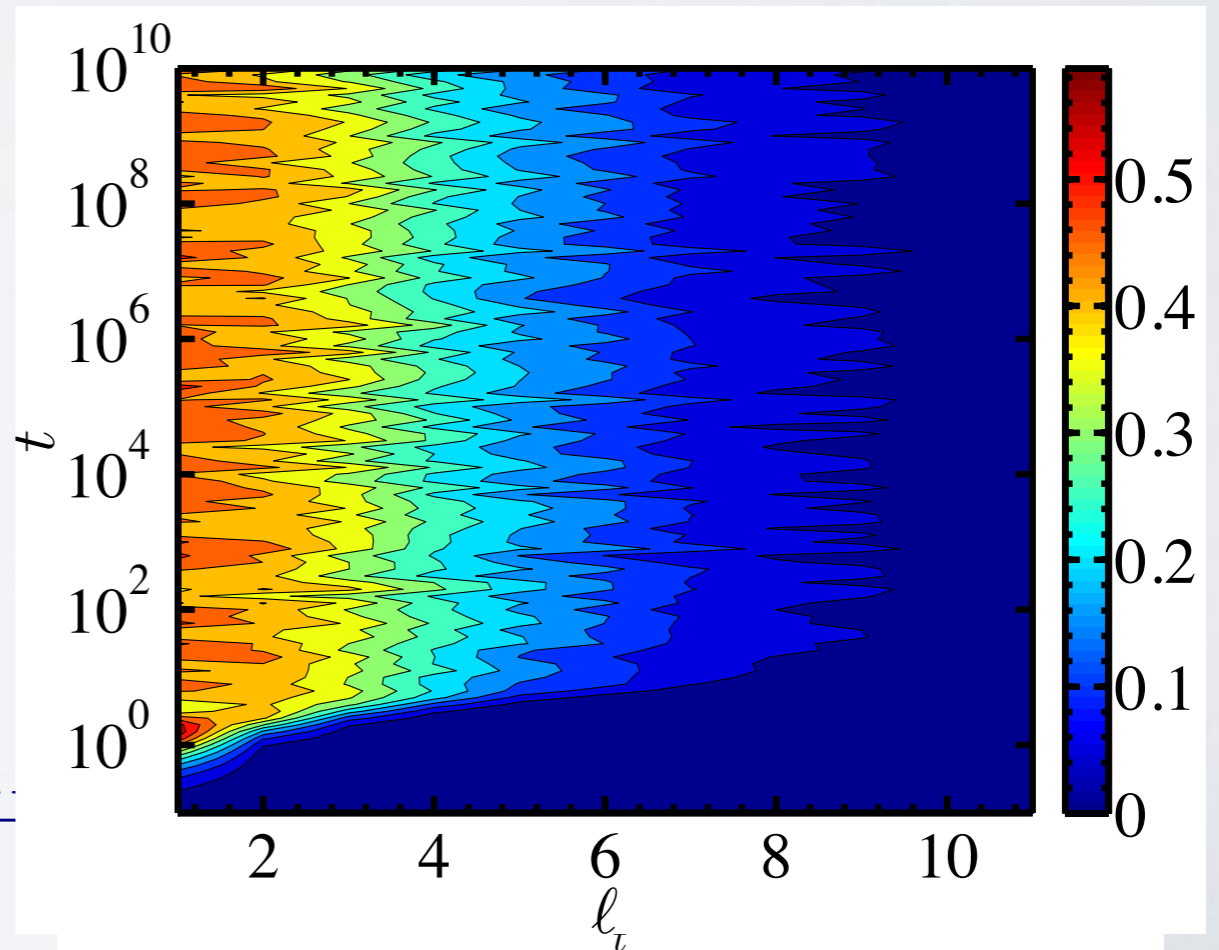
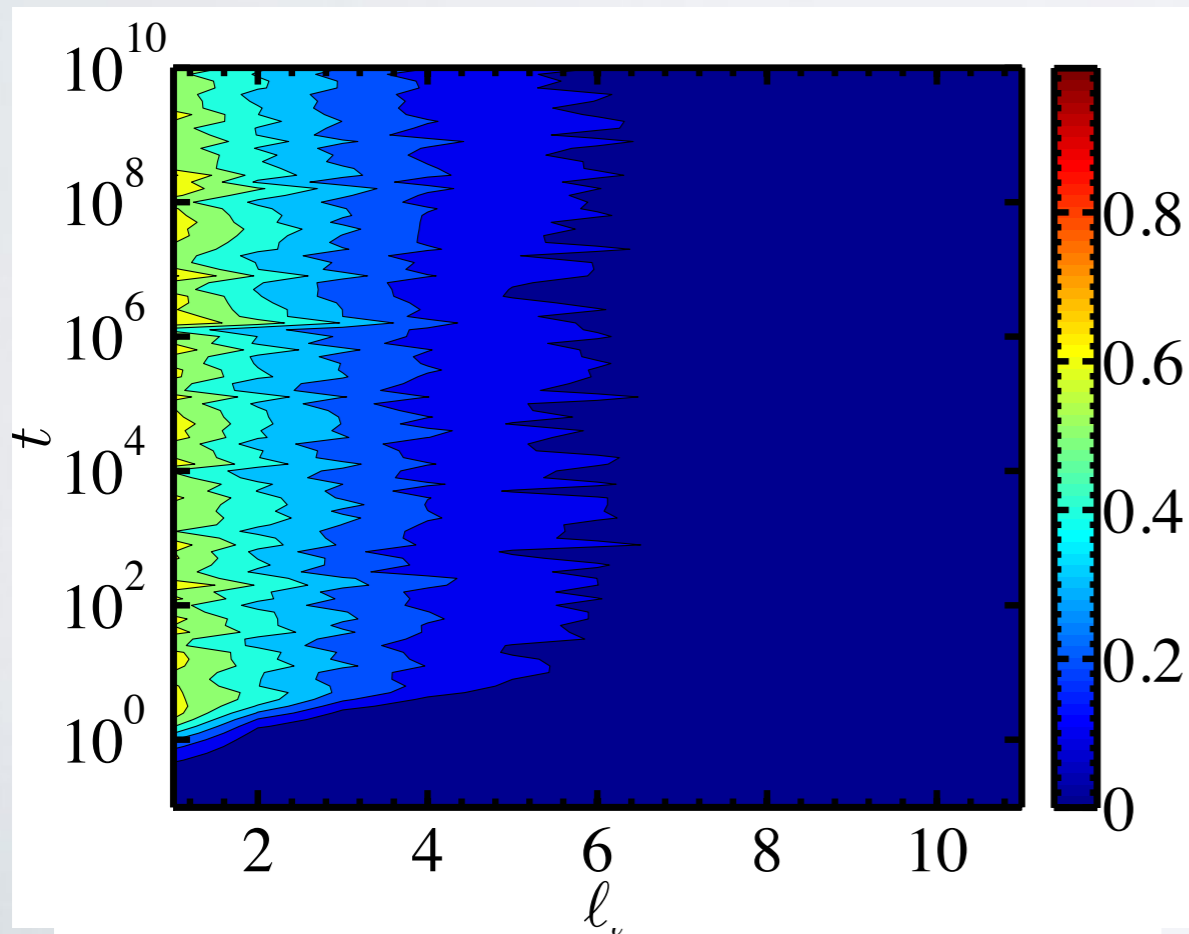
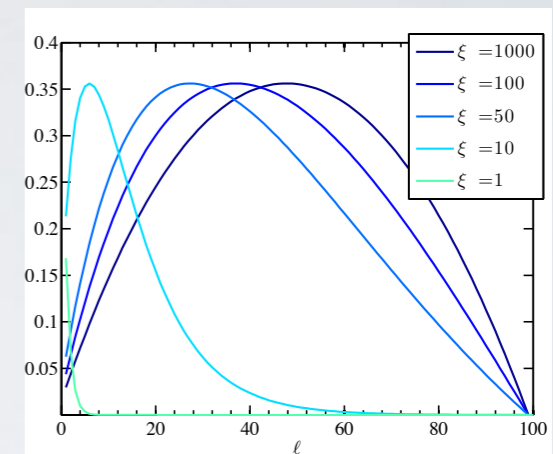
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propagation of correlations

interacting case: l-bit model

propagation of correlations

interacting case: l-bit model

$$H_{\text{eff}} = \sum_{i=0}^{N-1} \epsilon_i \tau_z^{[i]} + \sum_{i,j=0}^{N-1} K_{ij}^{(2)} \tau_z^{[i]} \tau_z^{[j]} + \sum_{i,j,k=0}^{N-1} K_{ijk}^{(3)} \tau_z^{[i]} \tau_z^{[j]} \tau_z^{[k]} + \dots,$$

propagation of correlations

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initial states

$$\rho_{\Phi} \propto |\Phi\rangle\langle\Phi| \otimes Id^{\otimes \frac{L-L_0}{2}} \quad |\Phi\rangle = |Z\pm\rangle, |X\pm\rangle$$

propagation of correlations

interacting case: l-bit model

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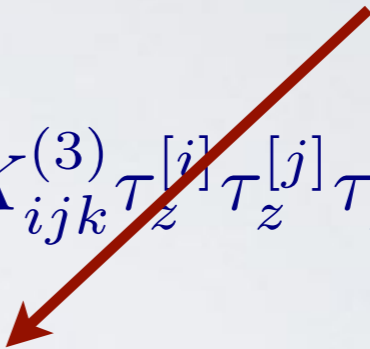
initial states: simple model

$$\rho_{X_+} \approx \frac{1 + \tau_x^{[0]}}{2} \otimes Id^{\otimes \frac{L-L_0}{2}}$$

propagation of correlations

interacting case: l-bit model

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exponentially decreasing

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propagation of correlations

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only terms
involving $\tau_z^{[0]}$

exponentially
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propagation of correlations

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single parameter

$$x(\ell, t) = \prod_{k=\ell}^{N-1} \cos(2tK_{0k}^{(2)})$$

propagation of correlations

interacting case: l-bit model

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initial states: simple model

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single parameter

$$x(\ell, t) = \prod_{k=\ell}^{N-1} \cos(2tK_{0k}^{(2)})$$

$$K_{0k}^{(2)} \approx e^{-k/\xi} \Rightarrow x(\ell, t) \approx 1 - 2t^2(N - \ell)e^{-2\ell/\xi}$$

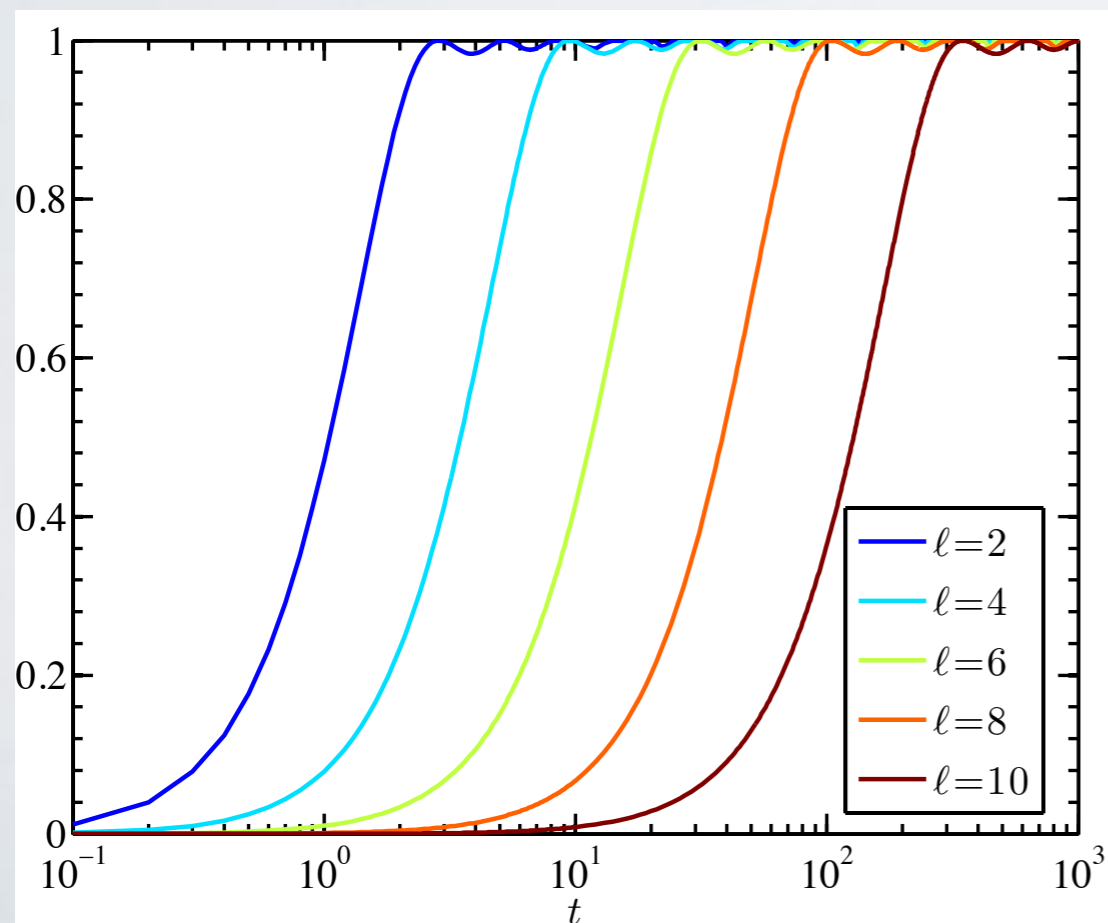
can be close to 0
takes exponential time

propagation of correlations

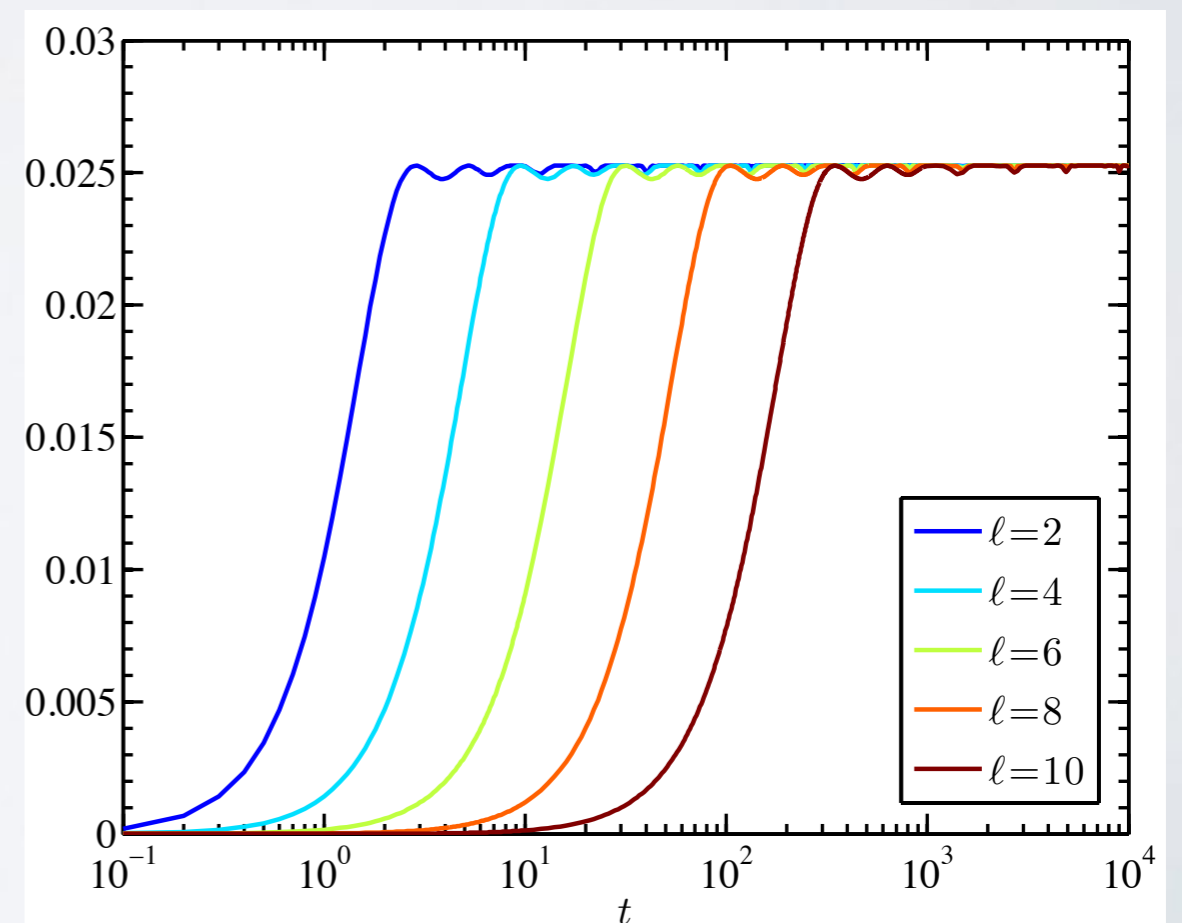
interacting case: l-bit model

$$\xi = 10$$

$|X_{\pm}\rangle$



$|Z_{\pm}\rangle$



propagation of correlations

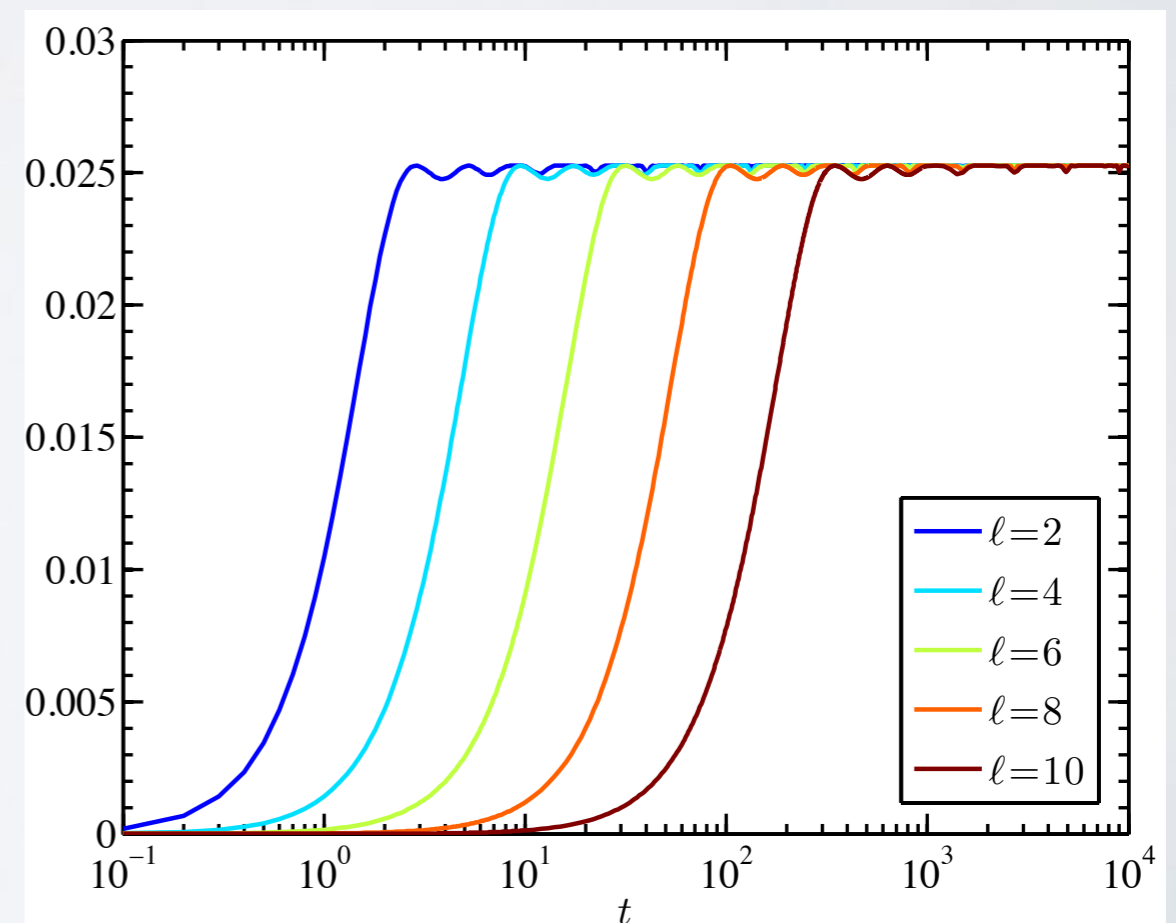
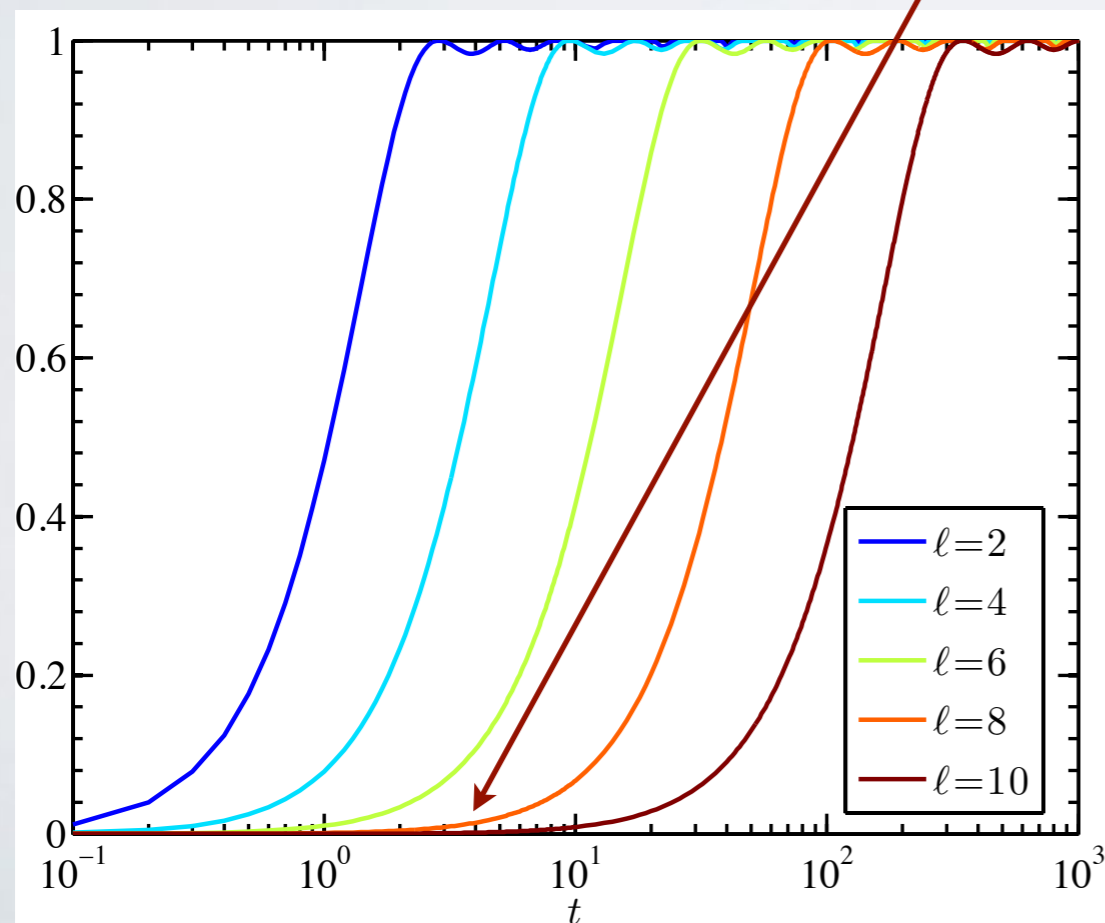
interacting case: l-bit model

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$|X_{\pm}\rangle$

exponential time
to reach a cut

$|Z_{\pm}\rangle$



propagation of correlations

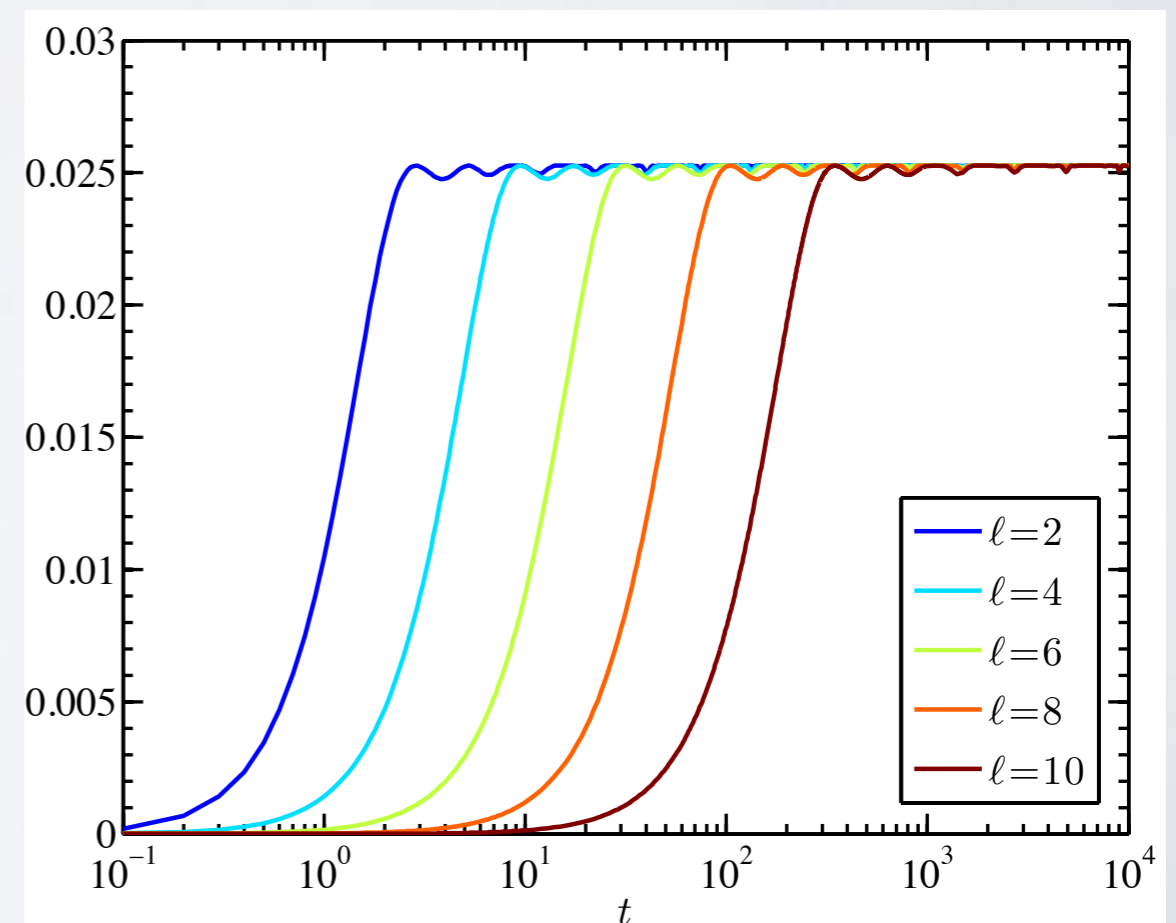
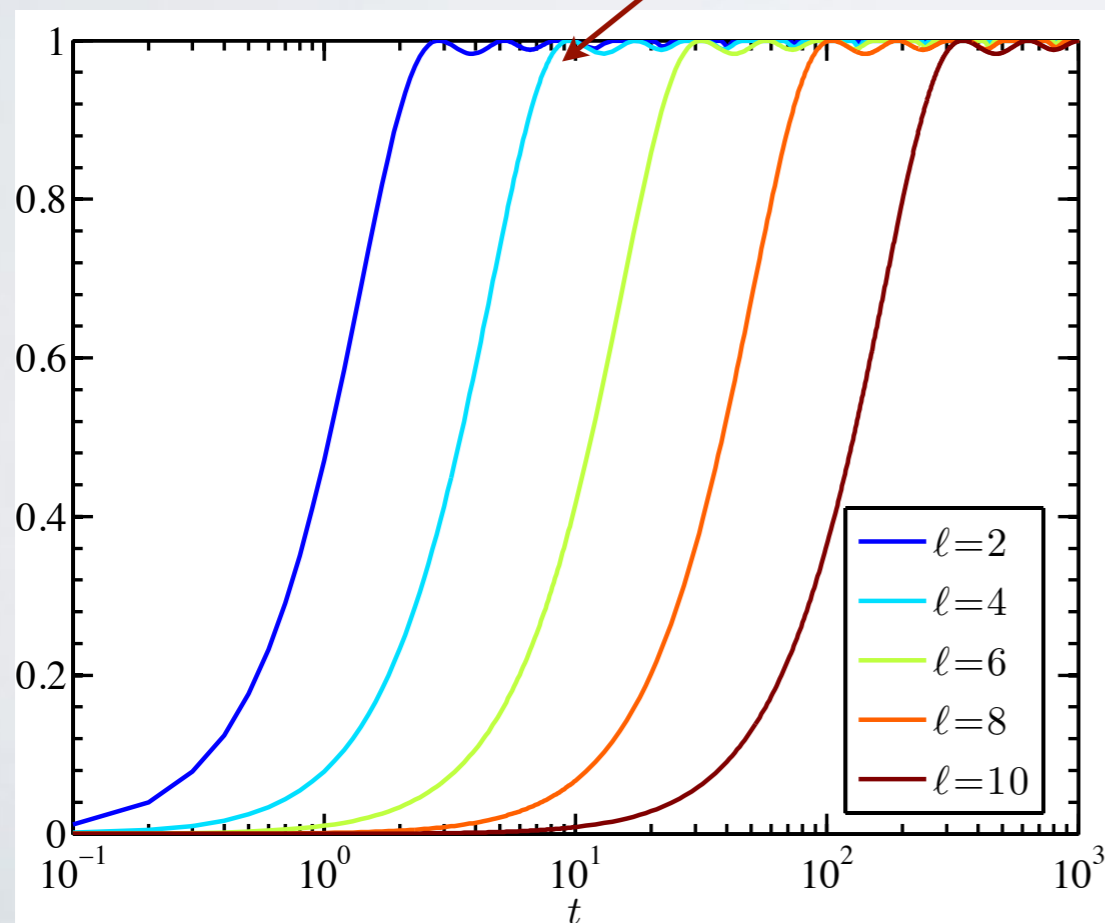
interacting case: l-bit model

$$\xi = 10$$

can reach the largest value

$|X_{\pm}\rangle$

$|Z_{\pm}\rangle$

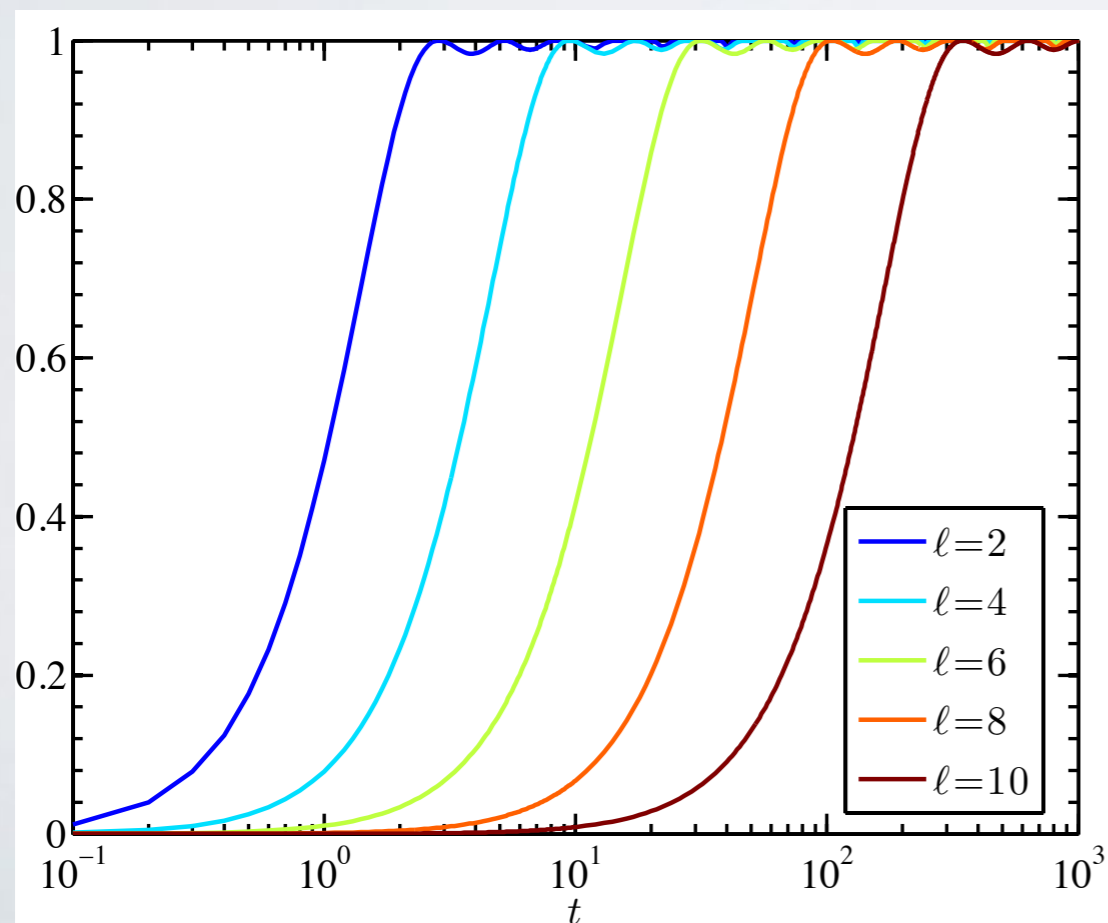


propagation of correlations

interacting case: l-bit model

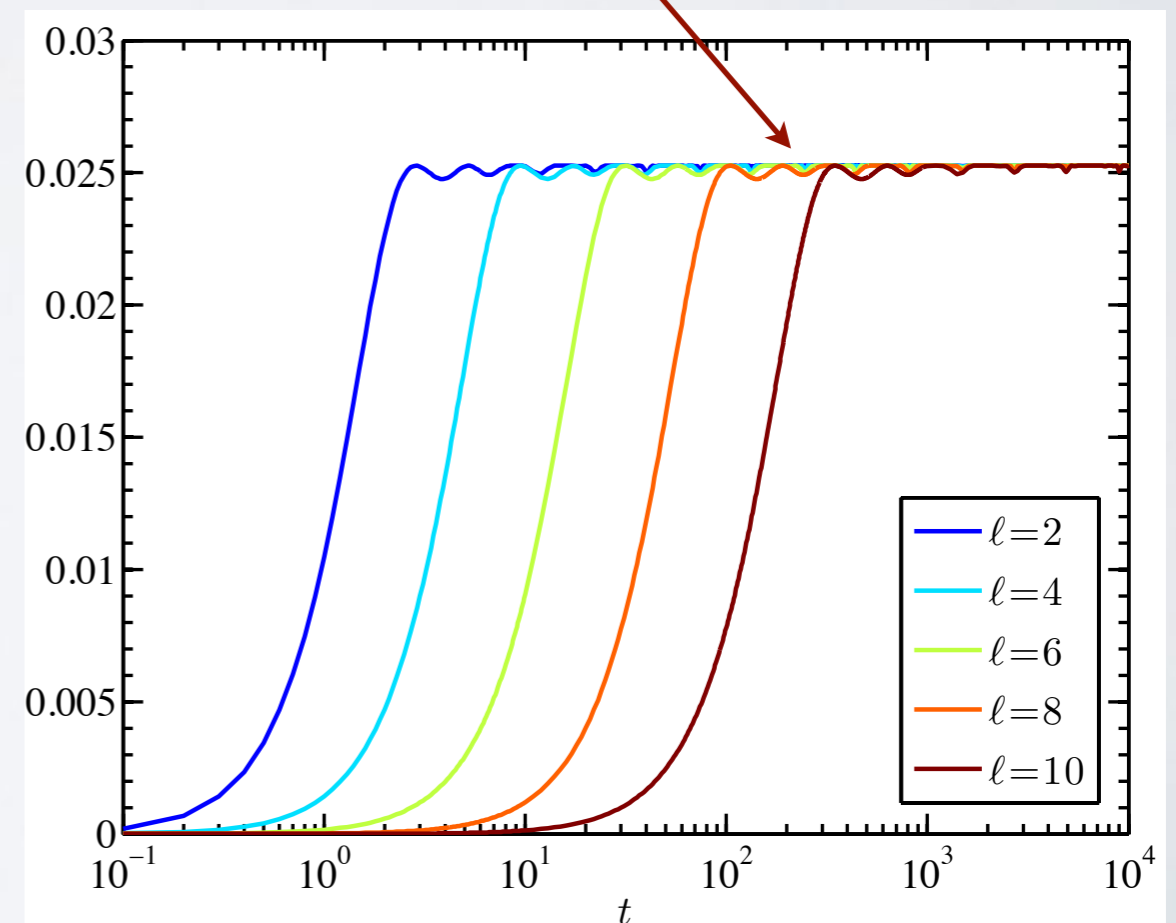
$\xi = 10$

$|X_{\pm}\rangle$



Z case from $\tau_x^{[0]}$
upper bounded

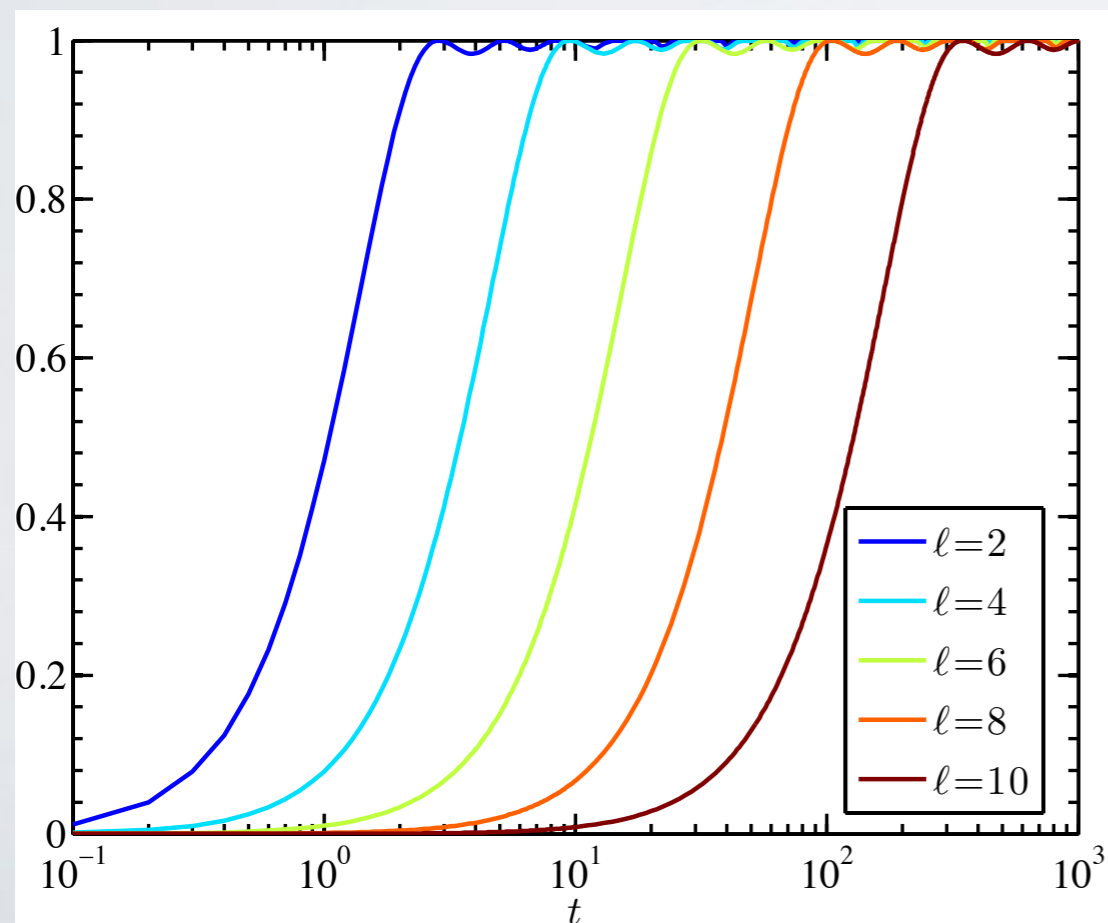
$|Z_{\pm}\rangle$



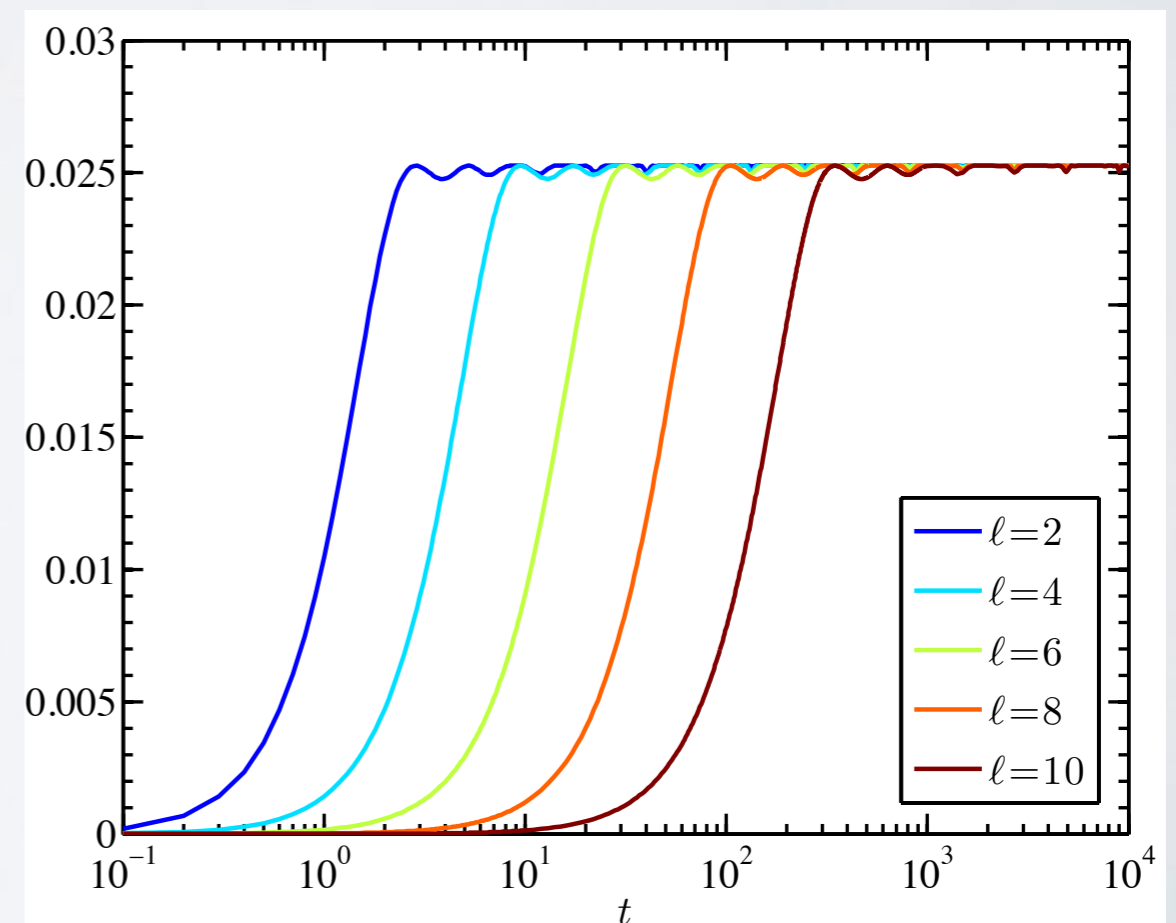
propagation of correlations

interacting case: l-bit model

$|X_{\pm}\rangle$



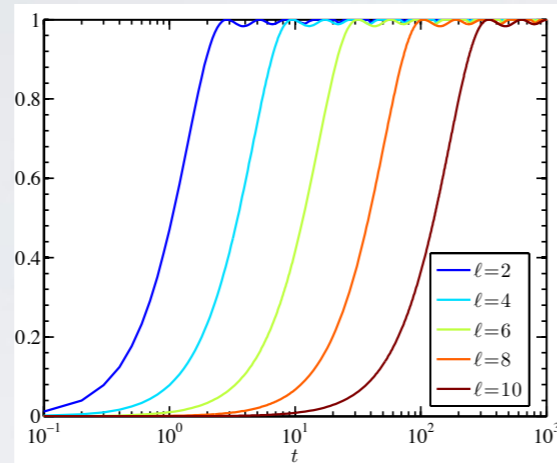
$|Z_{\pm}\rangle$



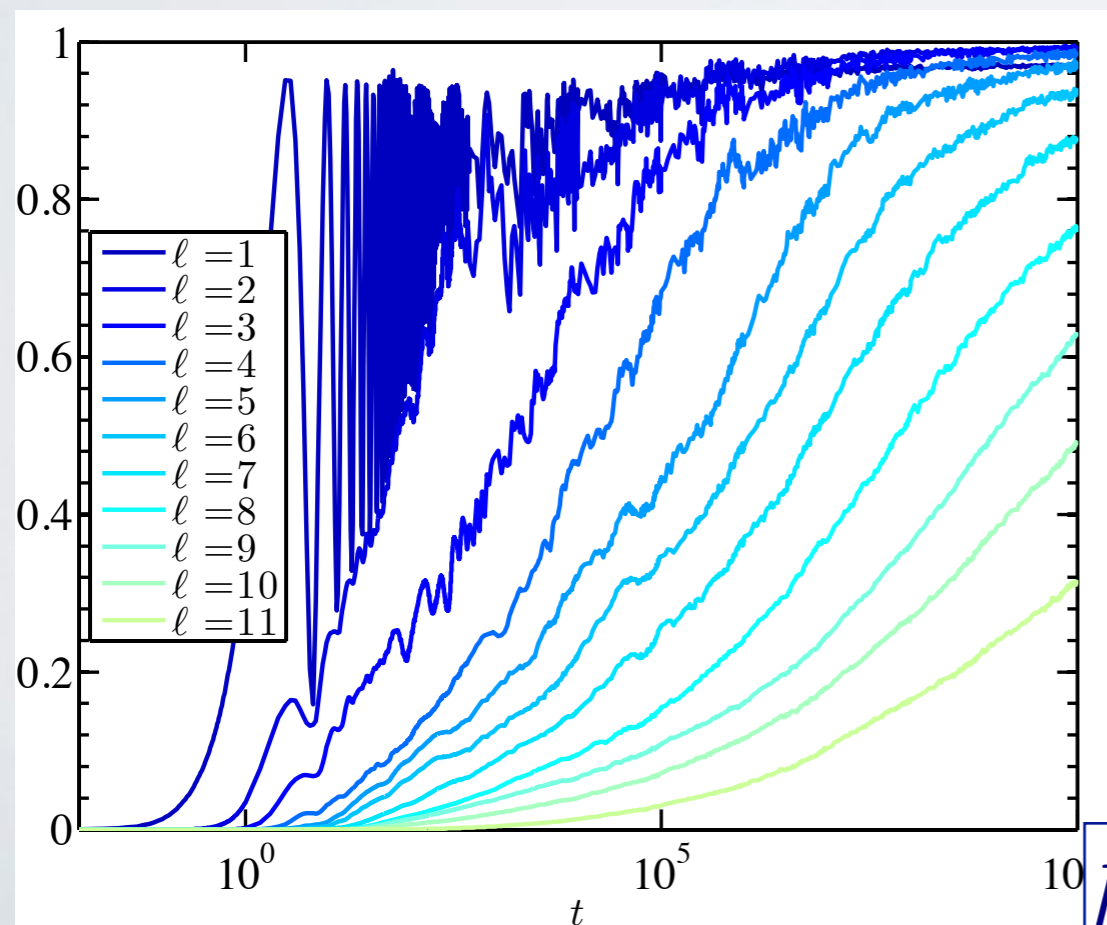
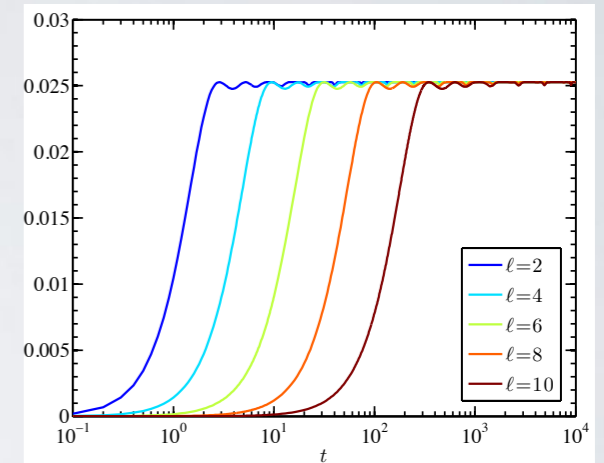
propagation of correlations

interacting case: l-bit model

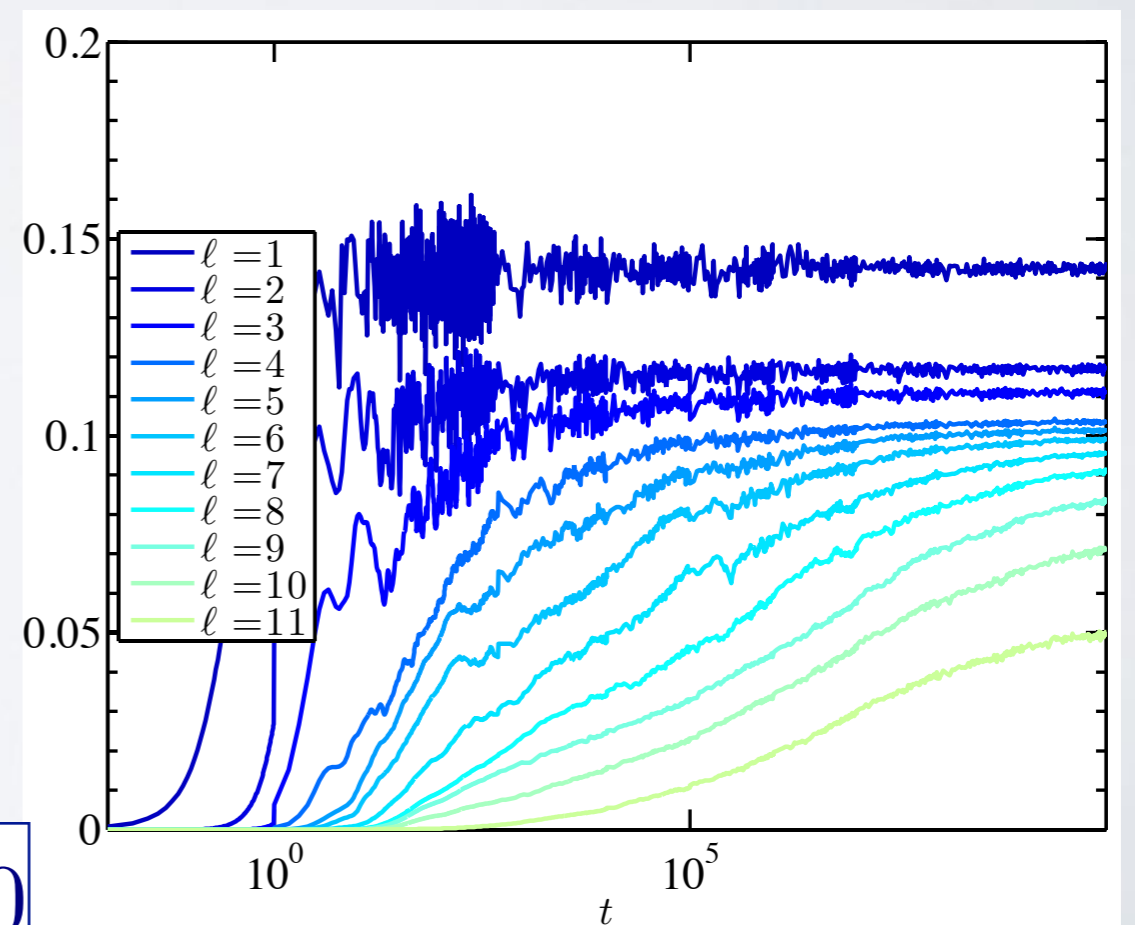
$|X_{\pm}\rangle$



$|Z_{\pm}\rangle$



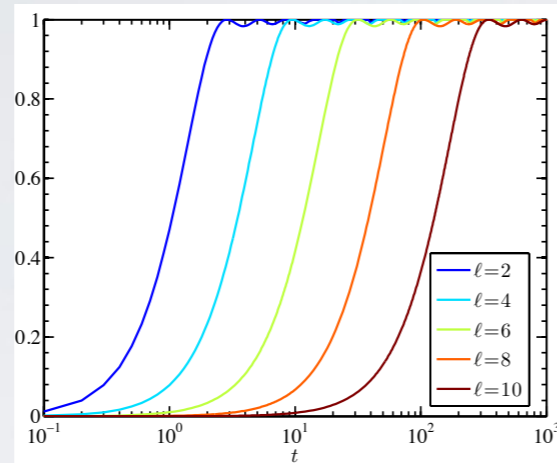
$h = 10$



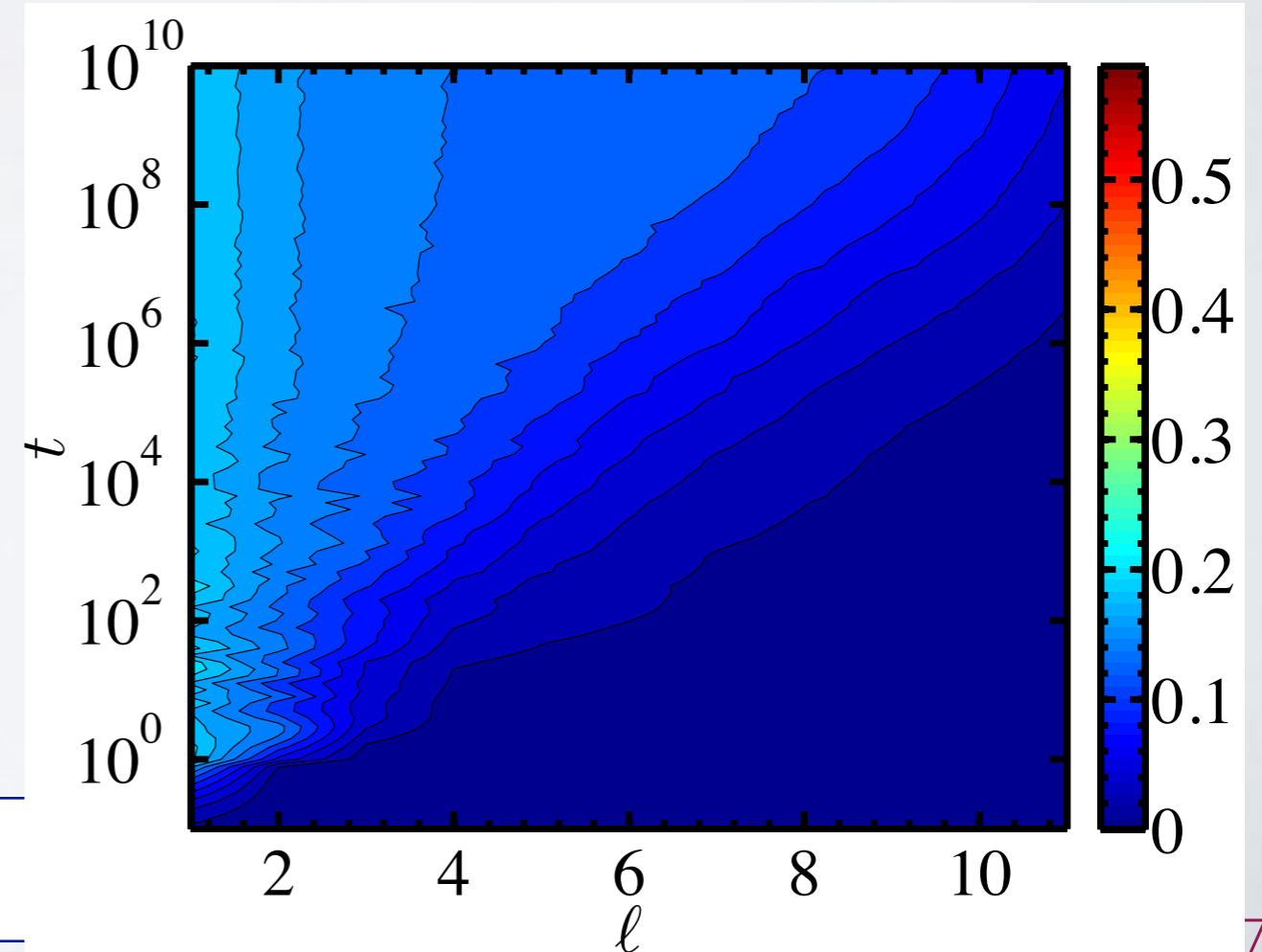
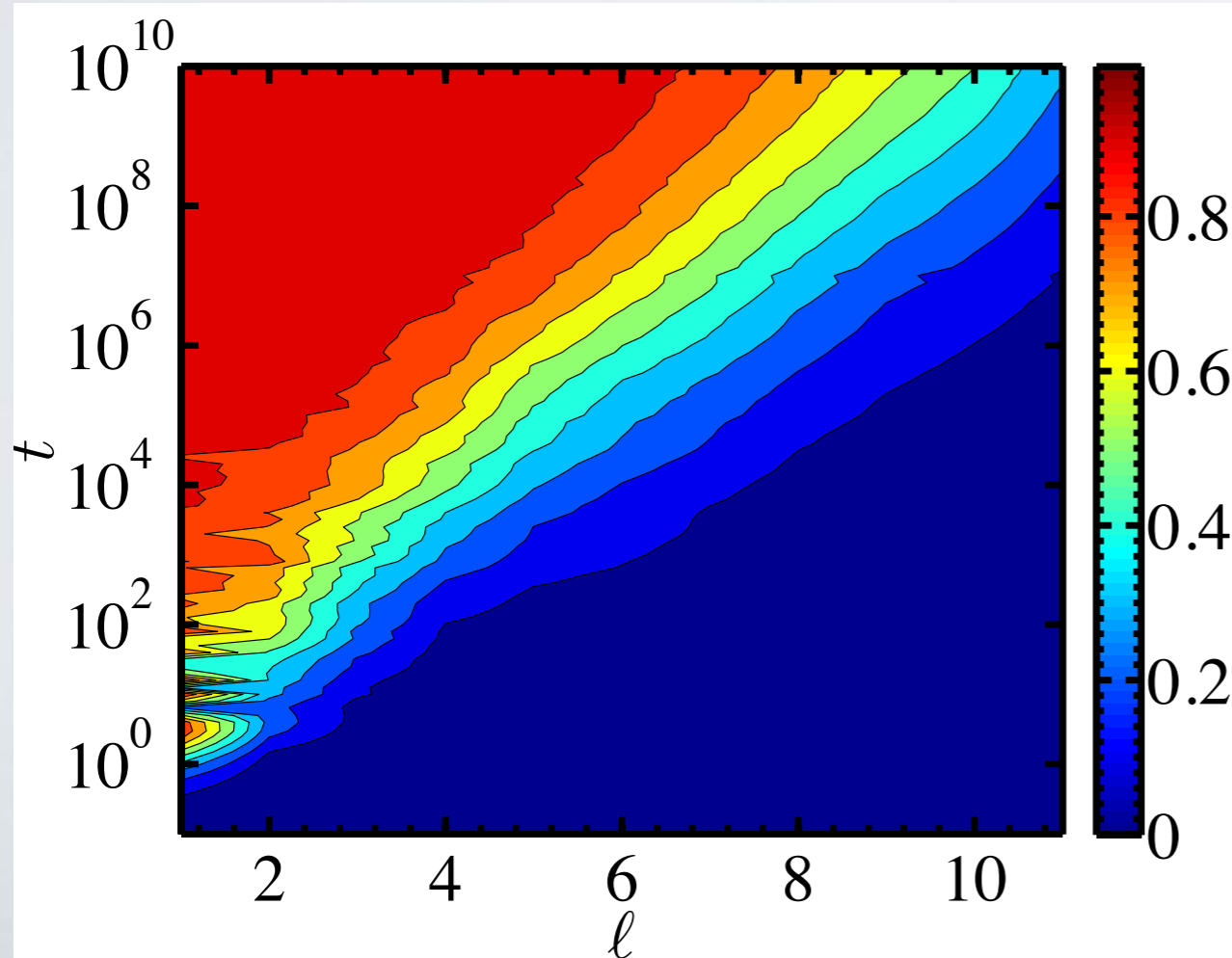
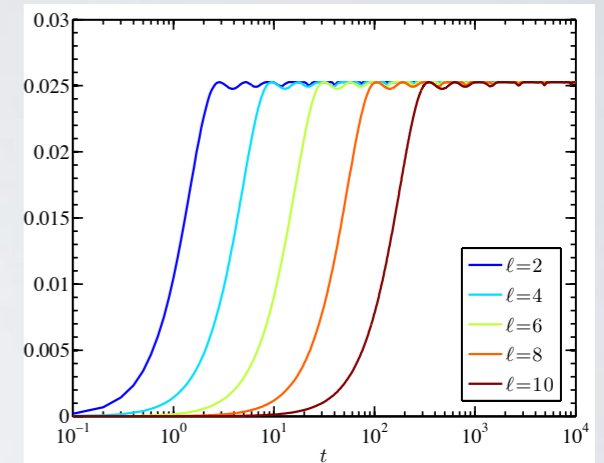
propagation of correlations

interacting case: l-bit model

$|X_{\pm}\rangle$



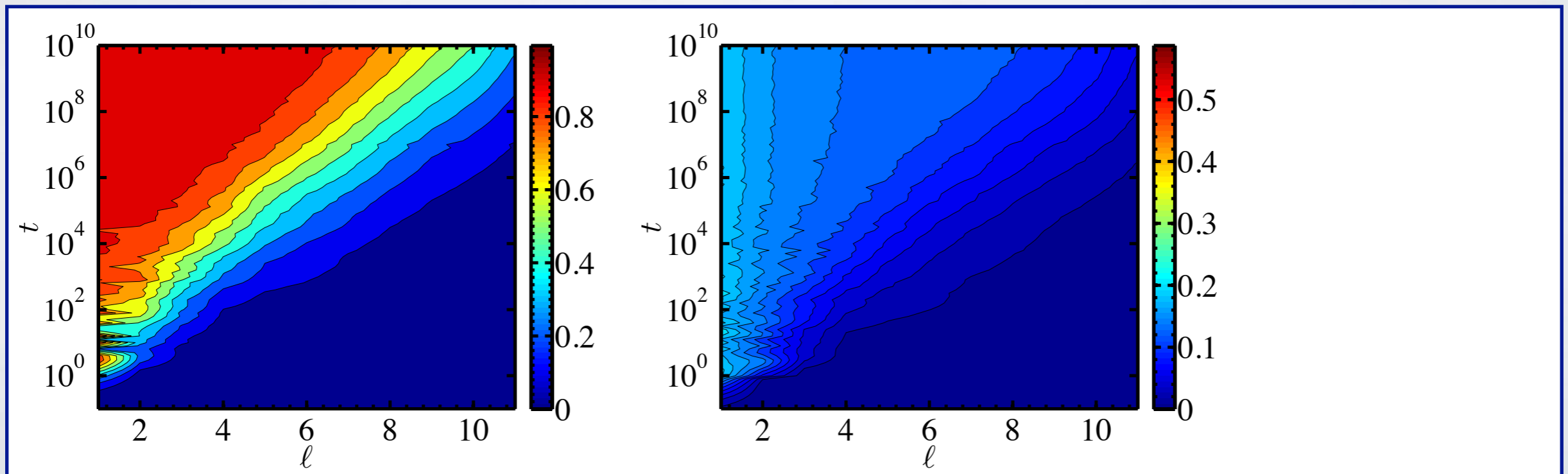
$|Z_{\pm}\rangle$



propagation of correlations

$|X_{\pm}\rangle$

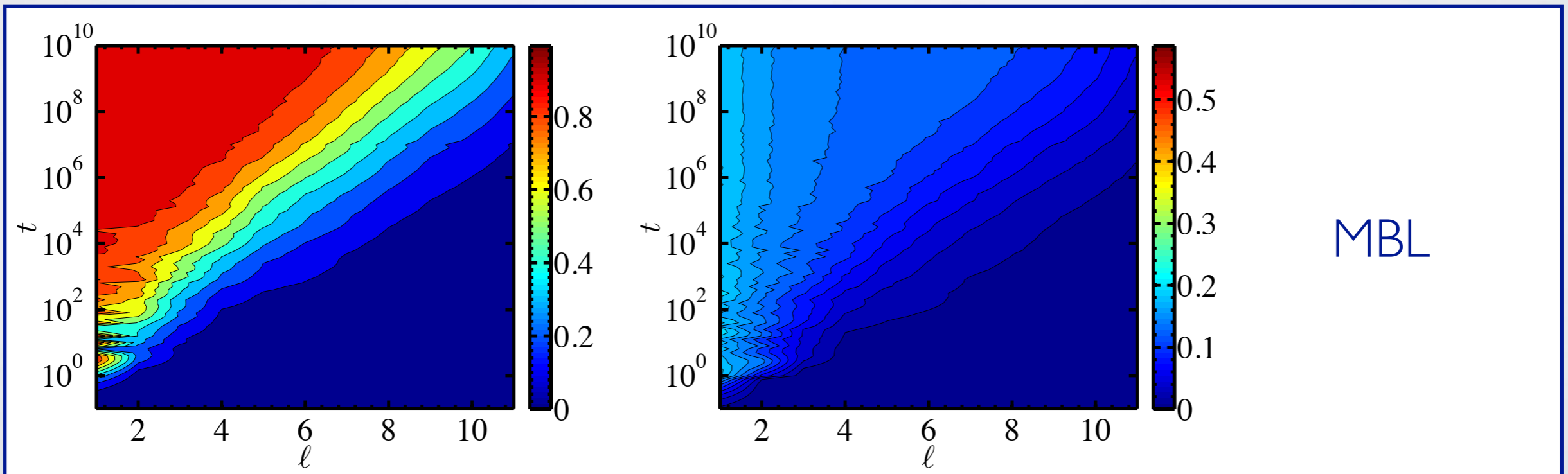
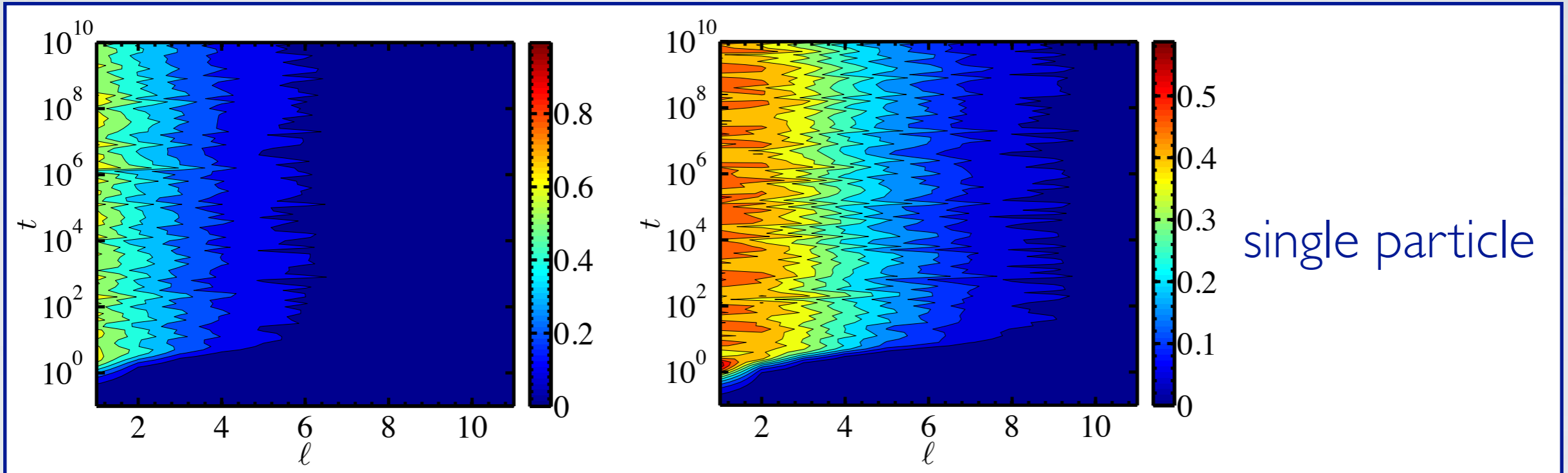
$|Z_{\pm}\rangle$



propagation of correlations

$|X_{\pm}\rangle$

$|Z_{\pm}\rangle$



Some questions we are asking

dynamics of
mixed states

propagation of correlations



quantum memory features



simulability with MPO

Hamiltonian
properties

local conserved quantities

quantum memory

$$\rho_{\Phi} \propto |\Phi\rangle\langle\Phi| \otimes Id^{\otimes \frac{L-L_0}{2}}$$



quantum memory

$$\rho_{\Phi} \propto |\Phi\rangle\langle\Phi| \otimes Id^{\otimes \frac{L-L_0}{2}}$$



could be used to encode a qubit

quantum memory

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quantify potential as quantum memory

recovery fidelity

quantum memory

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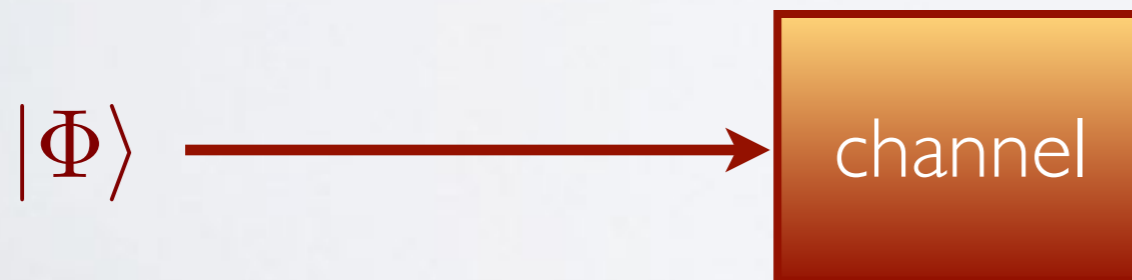
quantum memory

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quantify potential as quantum memory

recovery fidelity



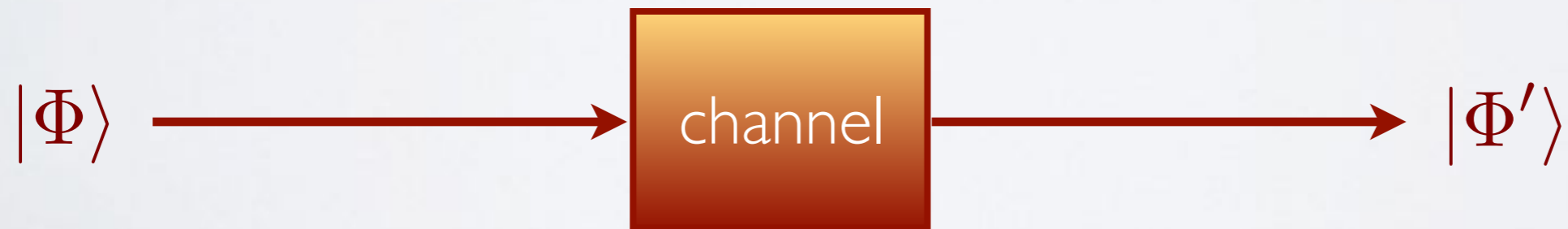
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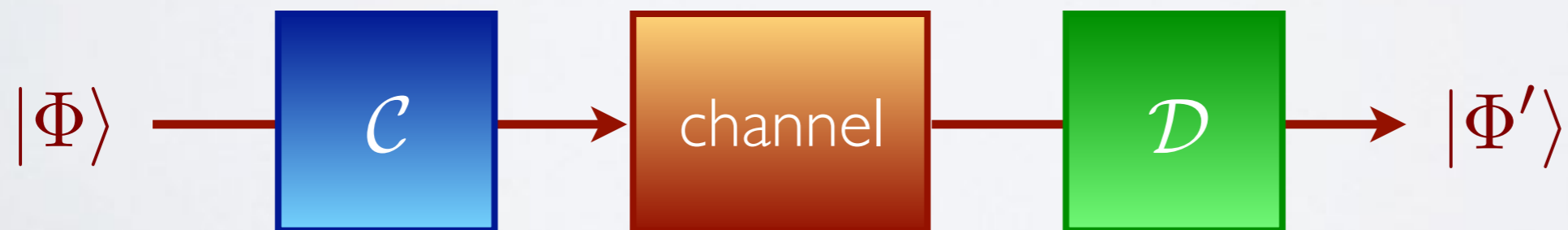
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recovery fidelity



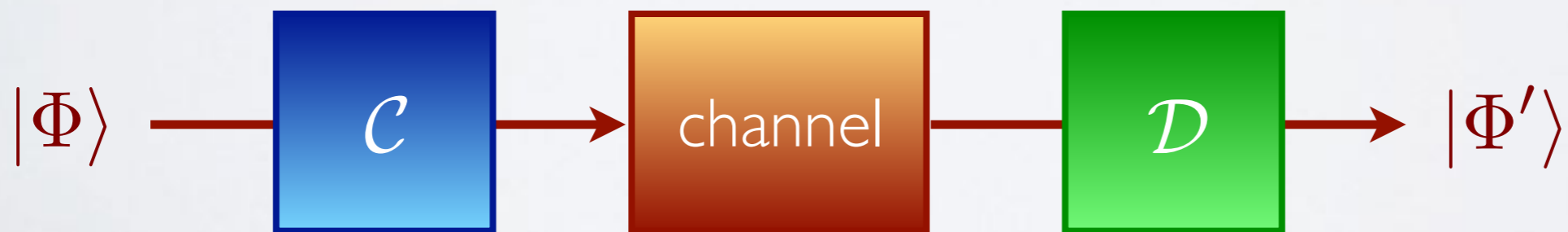
quantum memory

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quantify potential as quantum memory

recovery fidelity



preparing
the state

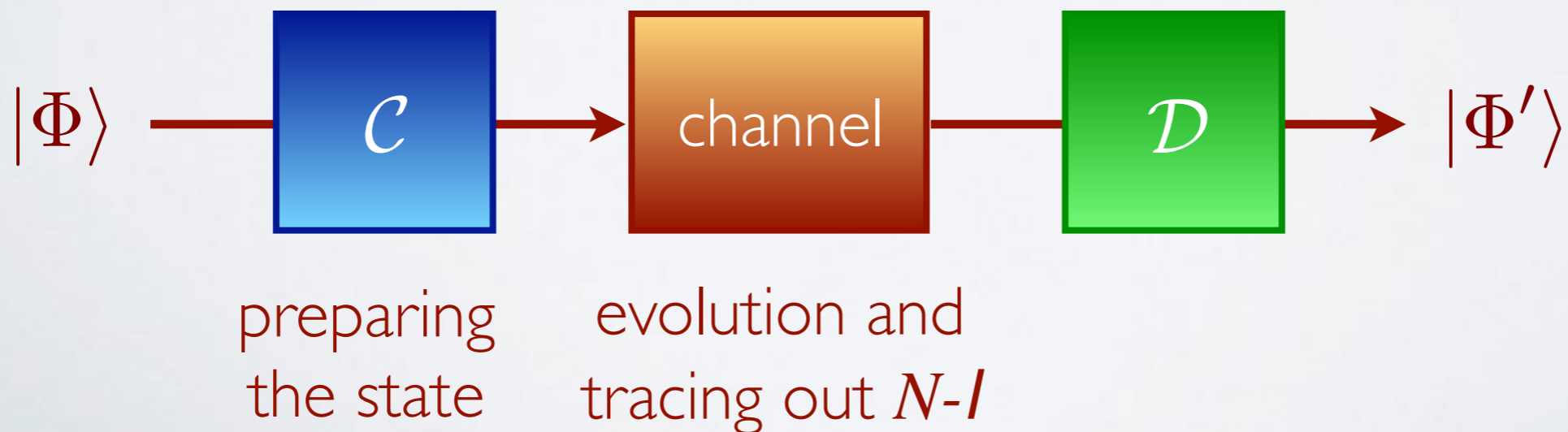
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quantify potential as quantum memory

recovery fidelity



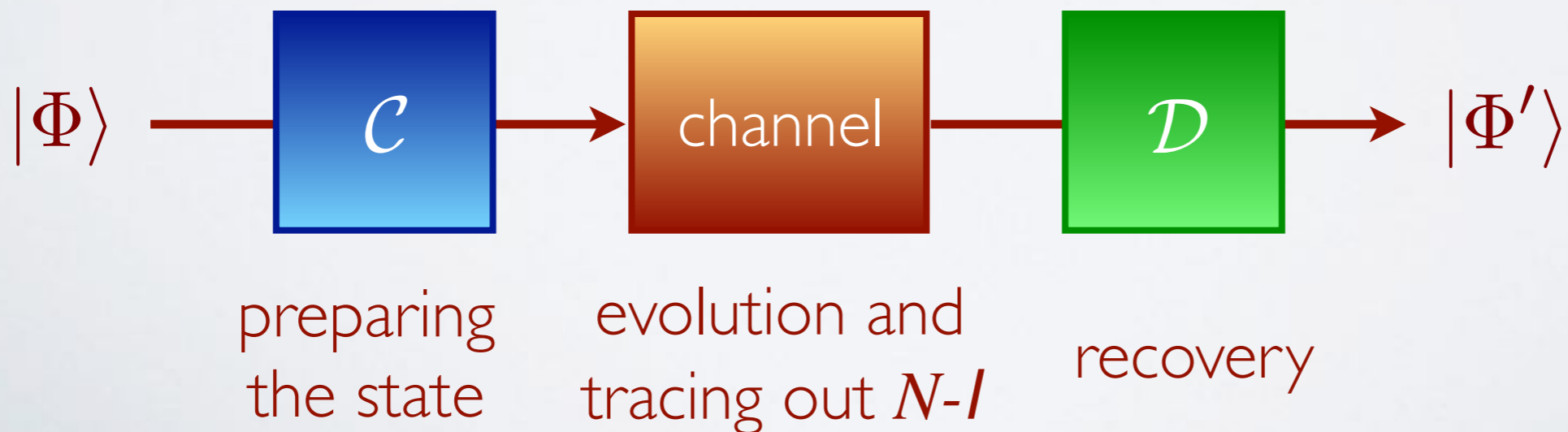
quantum memory

$$\rho_{\Phi} \propto |\Phi\rangle\langle\Phi| \otimes Id^{\otimes \frac{L-L_0}{2}}$$



quantify potential as quantum memory

recovery fidelity



quantum memory

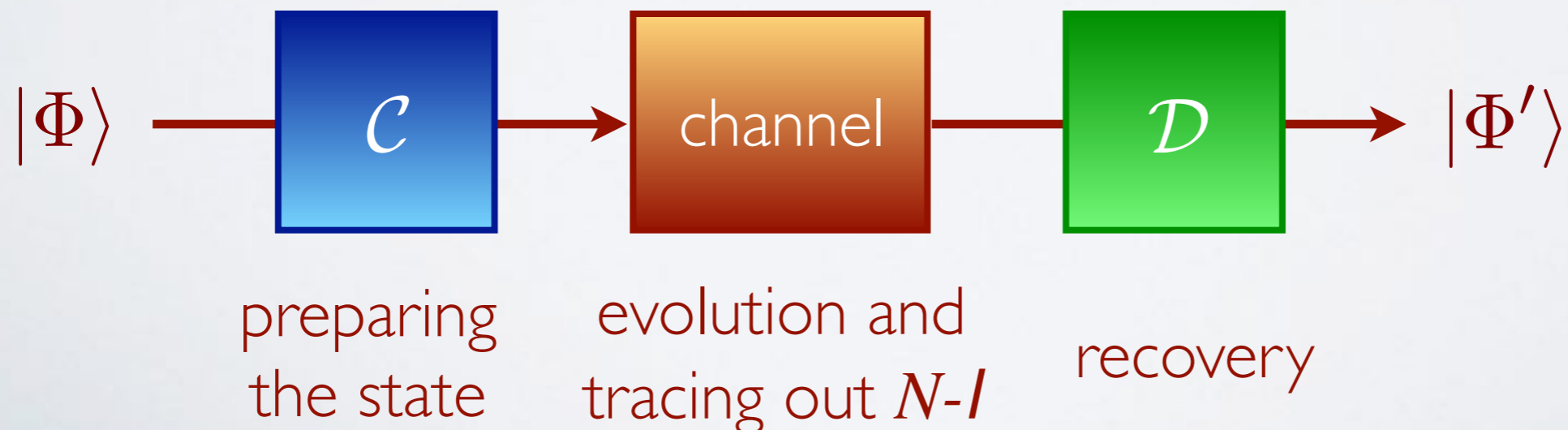
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quantify potential as quantum memory

recovery fidelity

how well we can recover a qubit state



quantum memory

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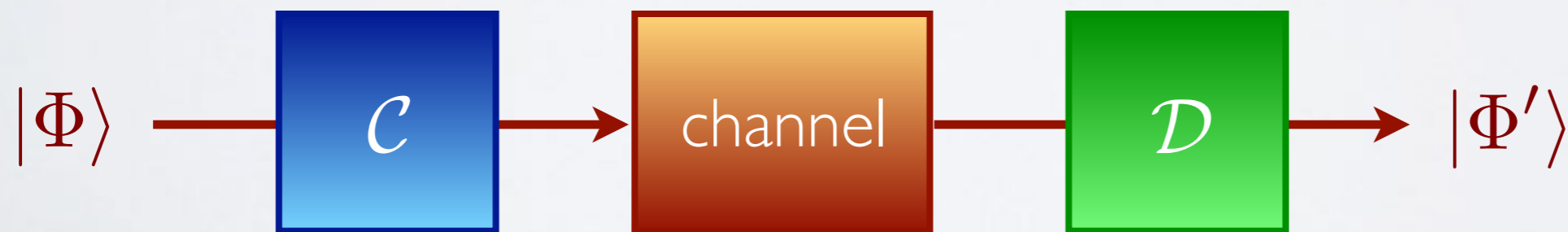
quantify potential as quantum memory

recovery fidelity

how well we can recover a qubit state

determined by distinguishability of orthogonal pairs $|X_{\pm}\rangle$

$|Z_{\pm}\rangle$



preparing
the state

evolution and
tracing out $N-1$

recovery

non-interacting case

$$\text{disting}_X = 2\sqrt{\mathcal{V}_l}$$

$$\text{disting}_Z = 2\mathcal{V}_l$$

$$\mathcal{V}_l = \sum_{r=0}^{\ell-1} |\langle r|U(t)|0\rangle|^2$$

non-interacting case $h > 0 \Rightarrow \xi$

$$\text{disting}_X = 2\sqrt{\mathcal{V}_l} \geq 2\sqrt{1 - 2N(N - l)e^{-l/\xi}}$$

$$\text{disting}_Z = 2\mathcal{V}_l \geq 2 - 2CN(N - l)e^{-l/\xi}$$

quantum memory

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can store quantum information

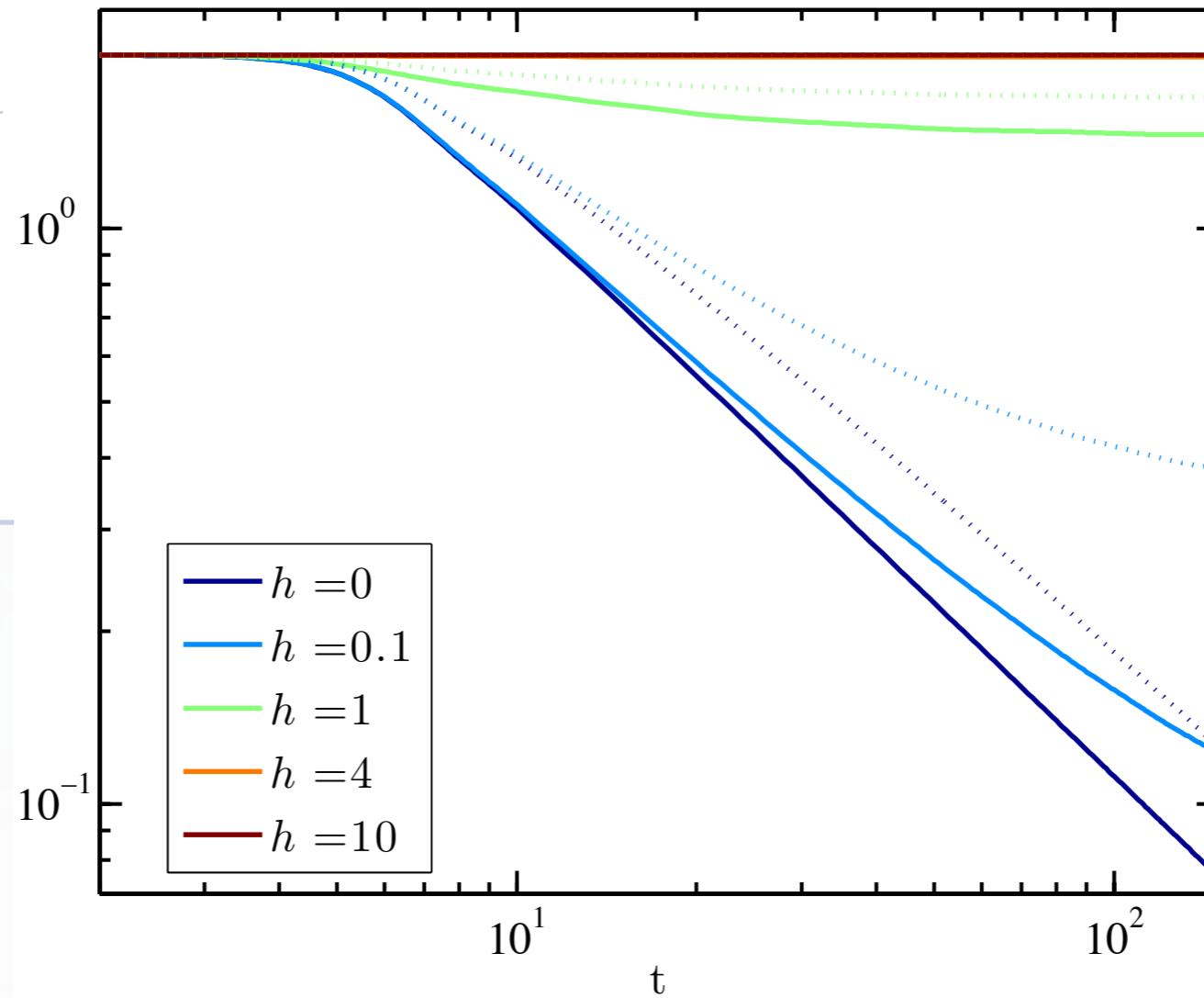
quantum memory

non-interacting case

$$h \rightarrow 0 \rightarrow \xi$$

disting_X

disting_Z



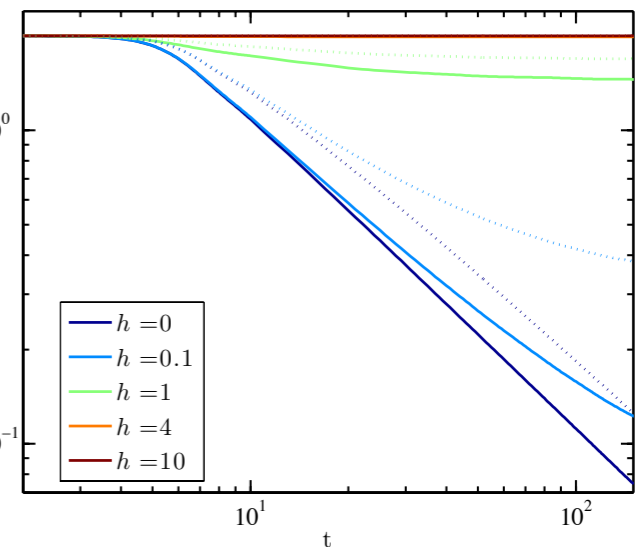
information

quantum memory

non-interacting case $h > 0 \Rightarrow \xi$

$$\text{disting}_X = 2\sqrt{\mathcal{V}_l} \geq 2\sqrt{1 - 2N(N - l)e^{-\xi}}$$

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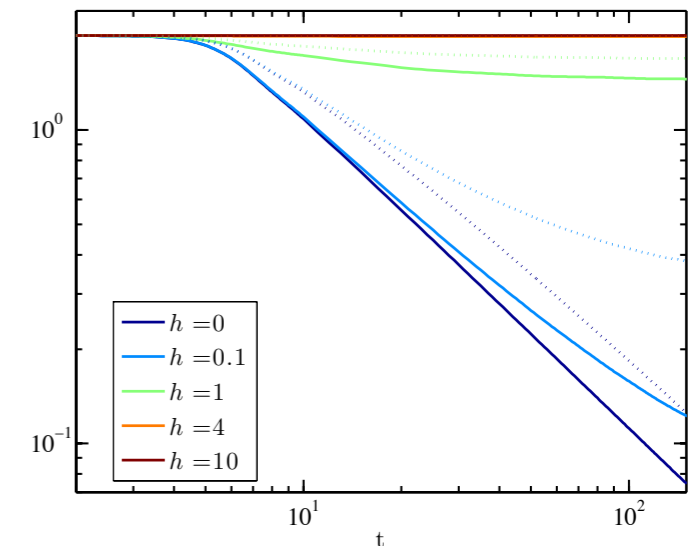
can store quantum information

quantum memory

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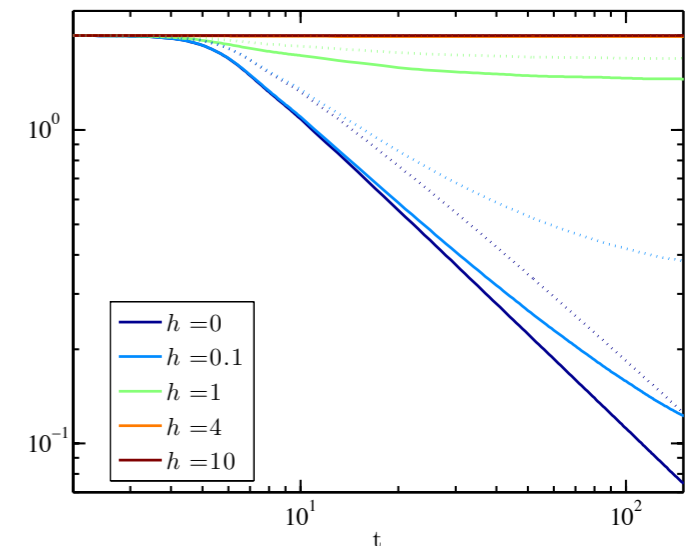
can store quantum information

quantum memory

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can store quantum information

interacting case: l-bits

$$\text{disting}_X \approx 2|x(\ell, t)|$$

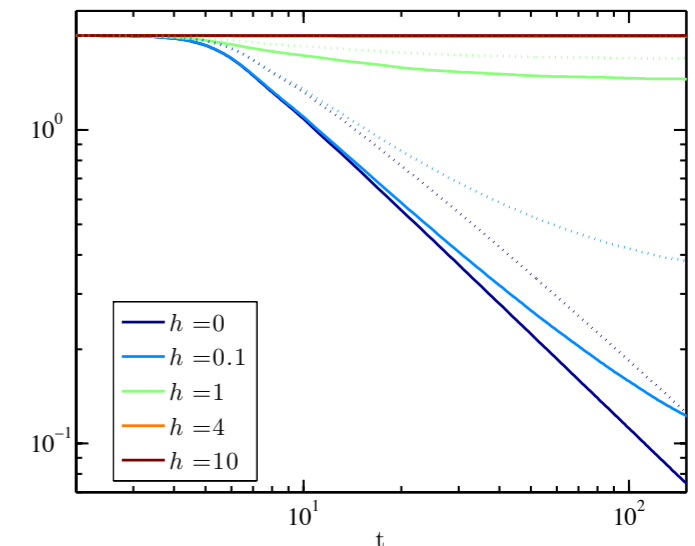
$$\text{disting}_Z \approx 2\sqrt{1 - \alpha^2(1 - x(\ell, t))^2}$$

$$x(\ell, t) = \prod_{k=\ell}^{N-1} \cos(2tK_{0k}^{(2)})$$

non-interacting case $h > 0 \Rightarrow \xi$

$$\text{disting}_X \geq 2\sqrt{1 - 2N(N - \ell)e^{-\ell/\xi}}$$

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can store quantum information

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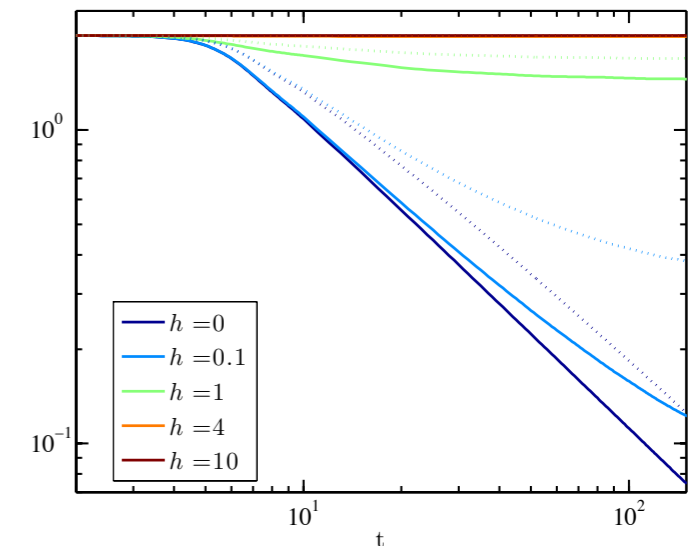
$$\text{disting}_Z \approx 2\sqrt{1 - \alpha^2(1 - x(\ell, t)^2)} \geq 2\sqrt{1 - \alpha^2}$$

quantum memory

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can store quantum information

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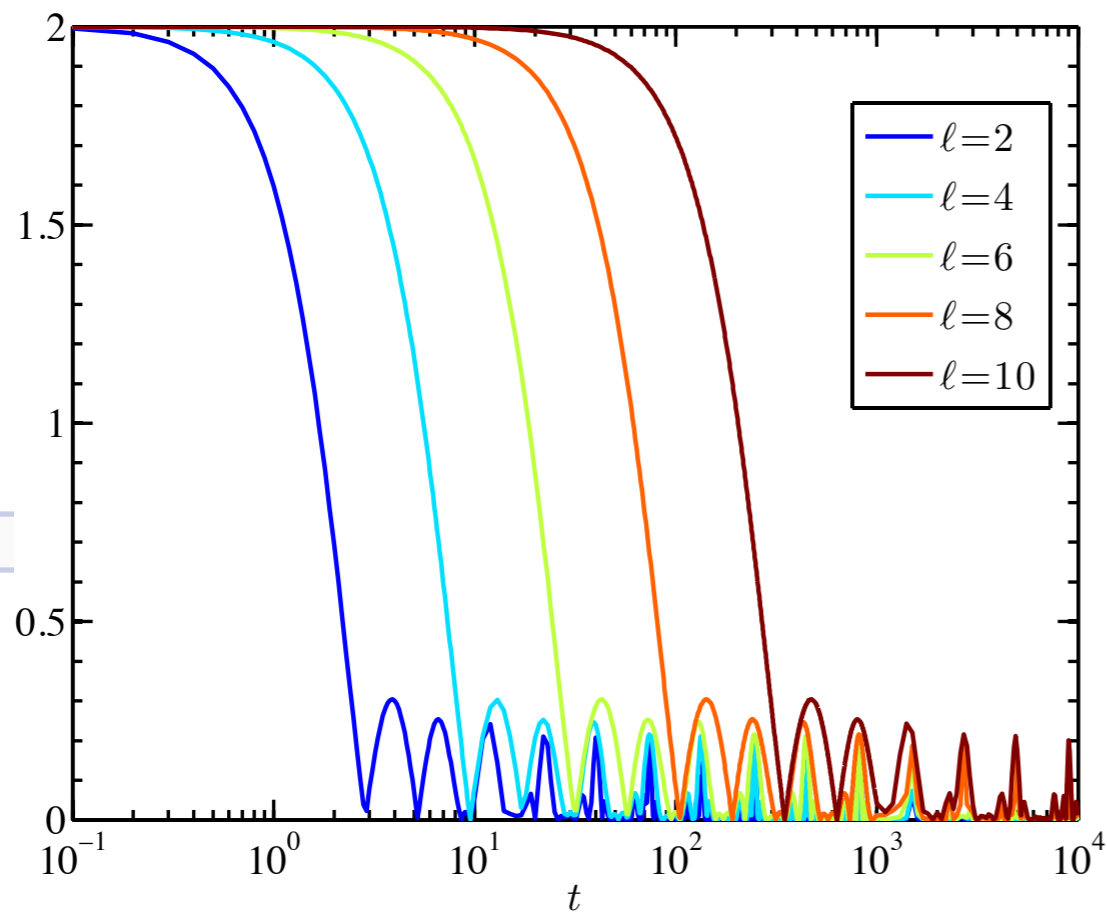
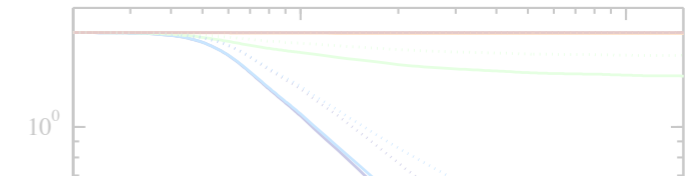
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can store only classical information

quantum memory

non-interacting case $h > 0 \Rightarrow \xi$

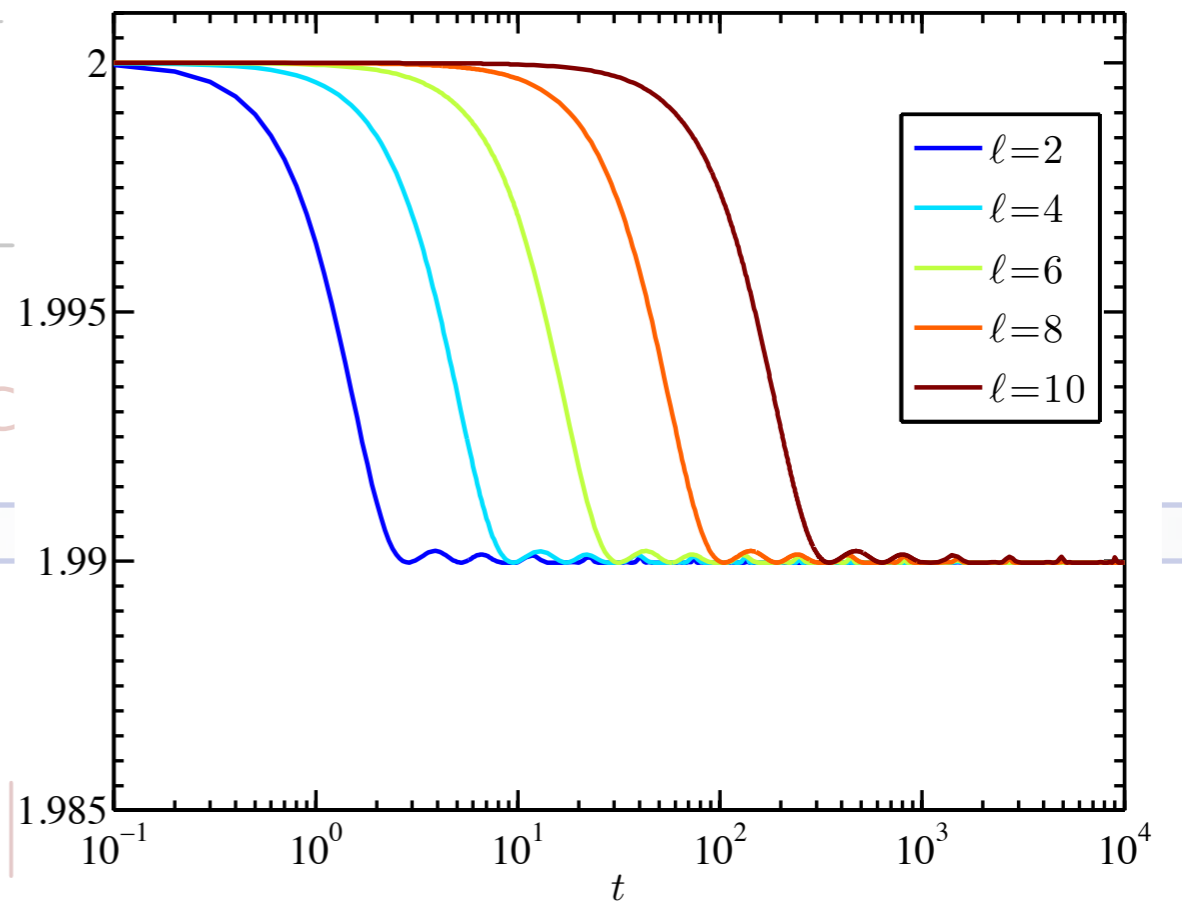


N

—

C

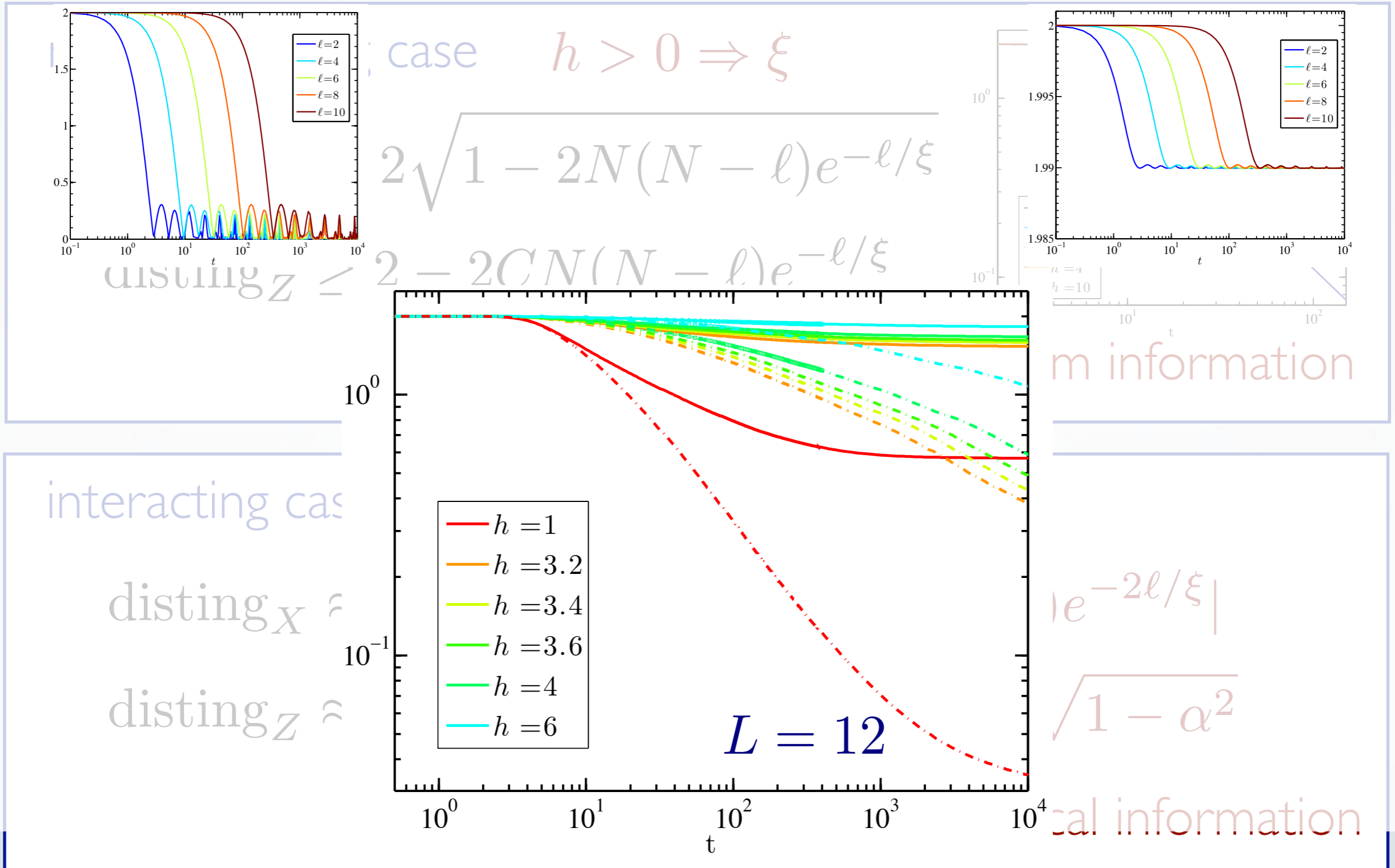
$2|$



$$\text{disting}_Z \approx 2\sqrt{1 - \alpha^2(1 - x(\ell, t)^2)} \geq 2\sqrt{1 - \alpha^2}$$

can store only classical information

quantum memory

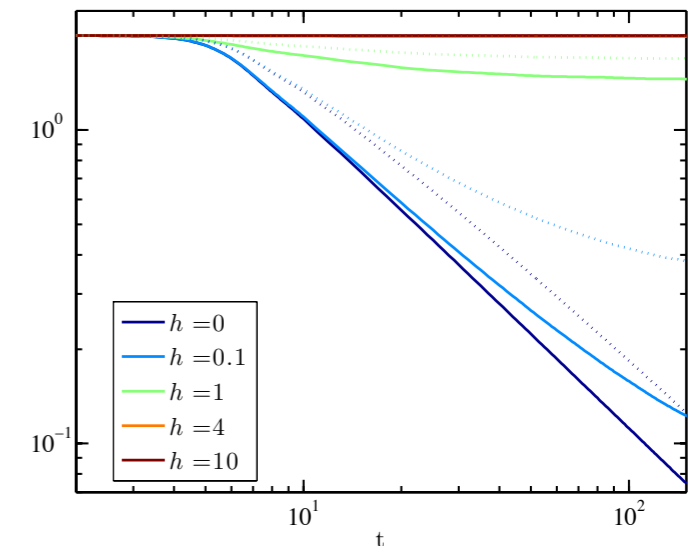


quantum memory

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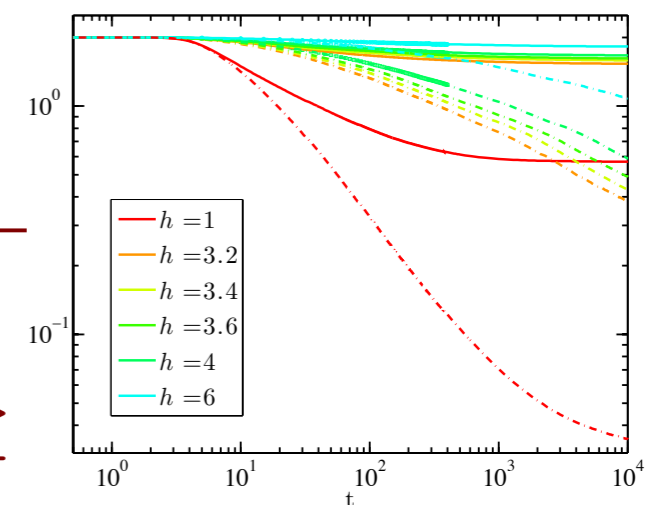


can store quantum information

interacting case: l-bits

$$\text{disting}_X \approx 2|x(\ell, t)| \approx 2|1 - 2t^2(N -$$

$$\text{disting}_Z \approx 2\sqrt{1 - \alpha^2(1 - x(\ell, t)^2)} \geq$$



can store only classical information

Some questions we are asking

dynamics of mixed states

propagation of correlations



quantum memory features



simulability with MPO



Hamiltonian properties

local conserved quantities

simulability with MPO

simulability with MPO

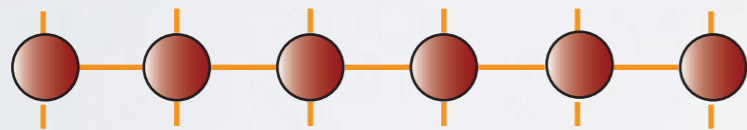
Q: how do errors behave?

simulability with MPO

Q: how do errors behave?

estimating error

best approximation

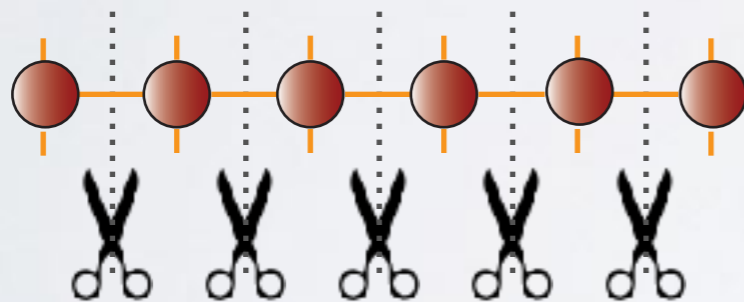


simulability with MPO

Q: how do errors behave?

estimating error

best approximation



simulability with MPO

Q: how do errors behave?

estimating error

best approximation



MPO
approximation
with smaller D
for long chains
only Frobenius
norm



simulability with MPO

localization \Leftrightarrow errors in TN simulation

$$L = 20$$

$$D_{\max} = 120$$



simulability with MPO

localization \Leftrightarrow errors in TN simulation

D for constant error

$$L = 20$$

$$D_{\max} = 120$$



simulability with MPO

localization \Leftrightarrow errors in TN simulation

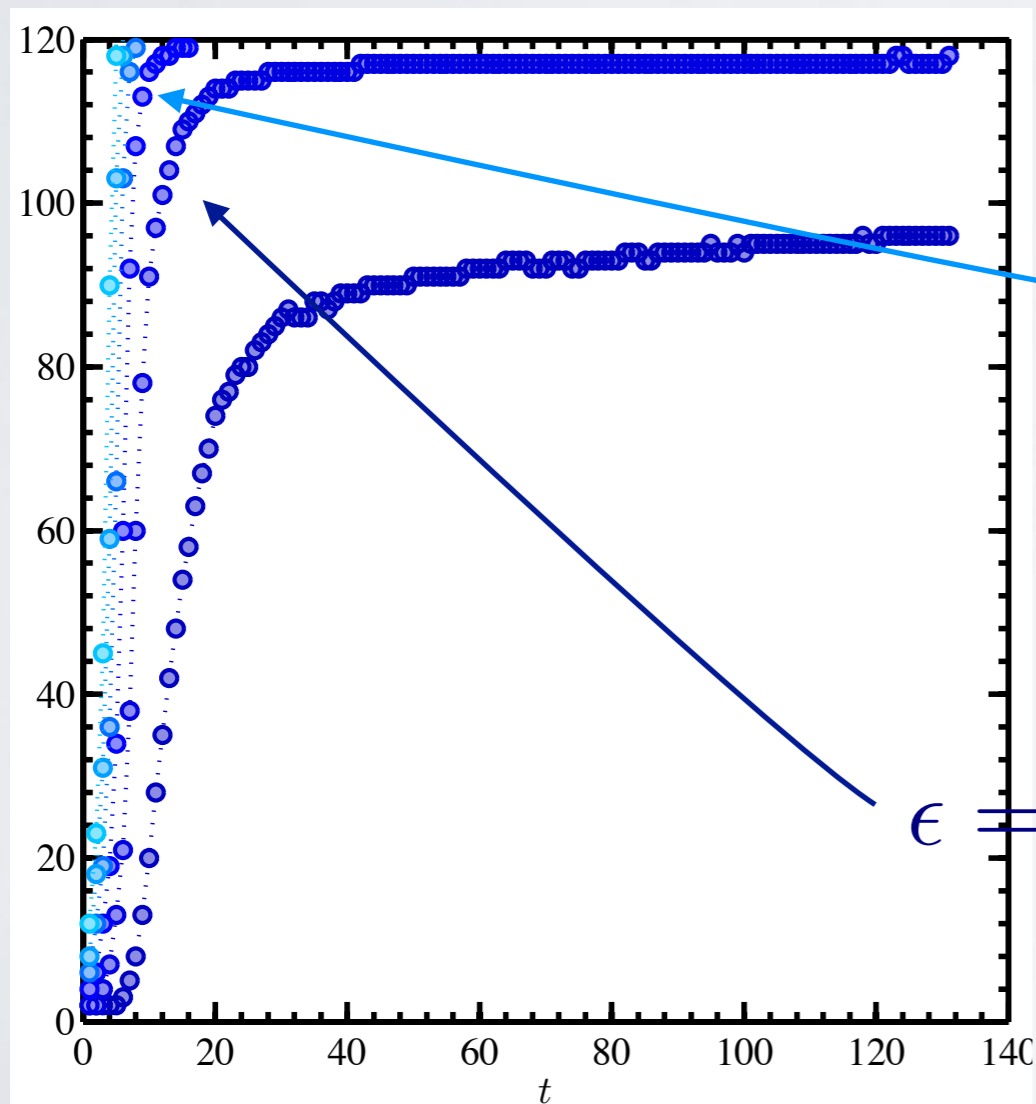
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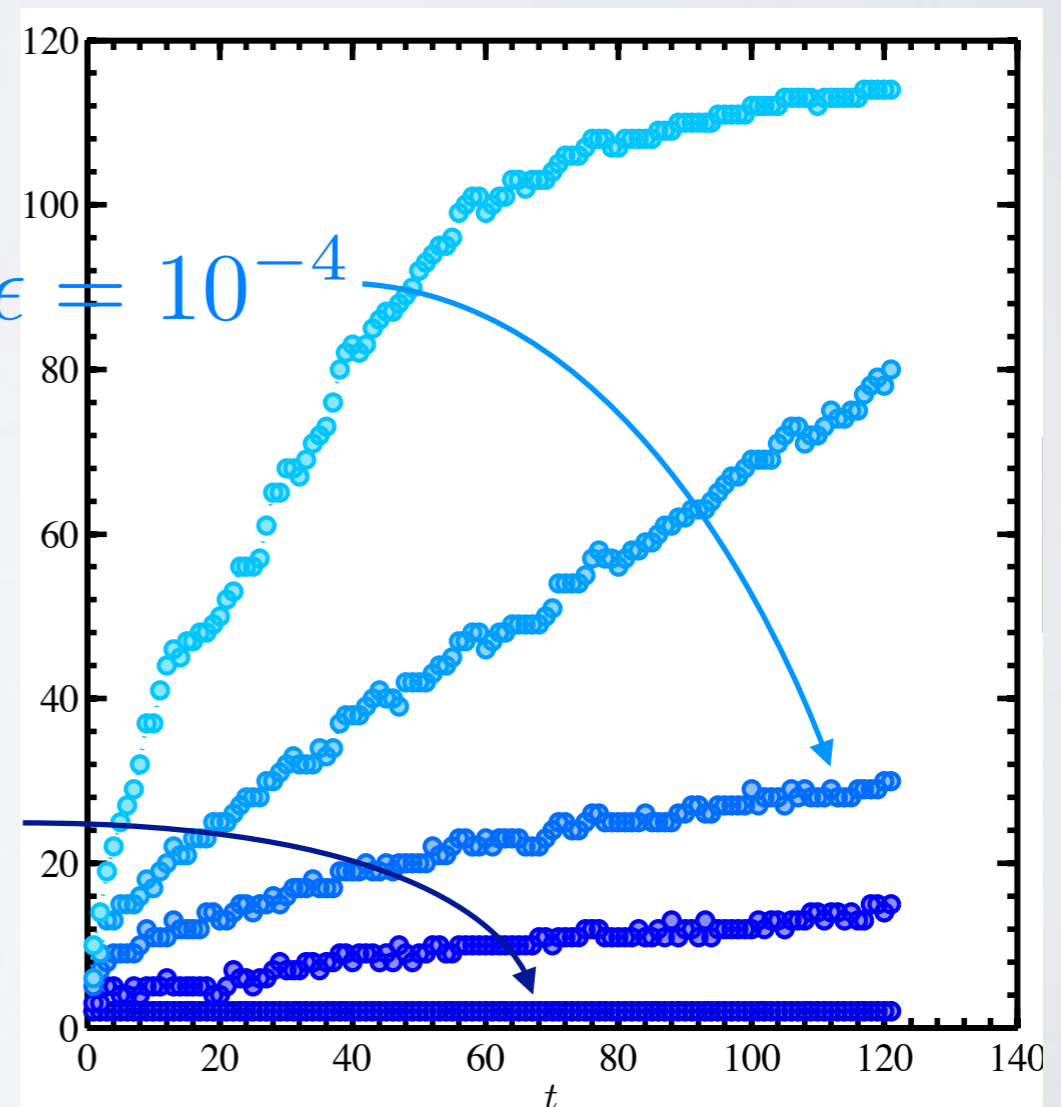
D for constant error

$$h = 0.5$$

$$h = 7.0$$



time



time



simulability with MPO

localization \Leftrightarrow errors in TN simulation

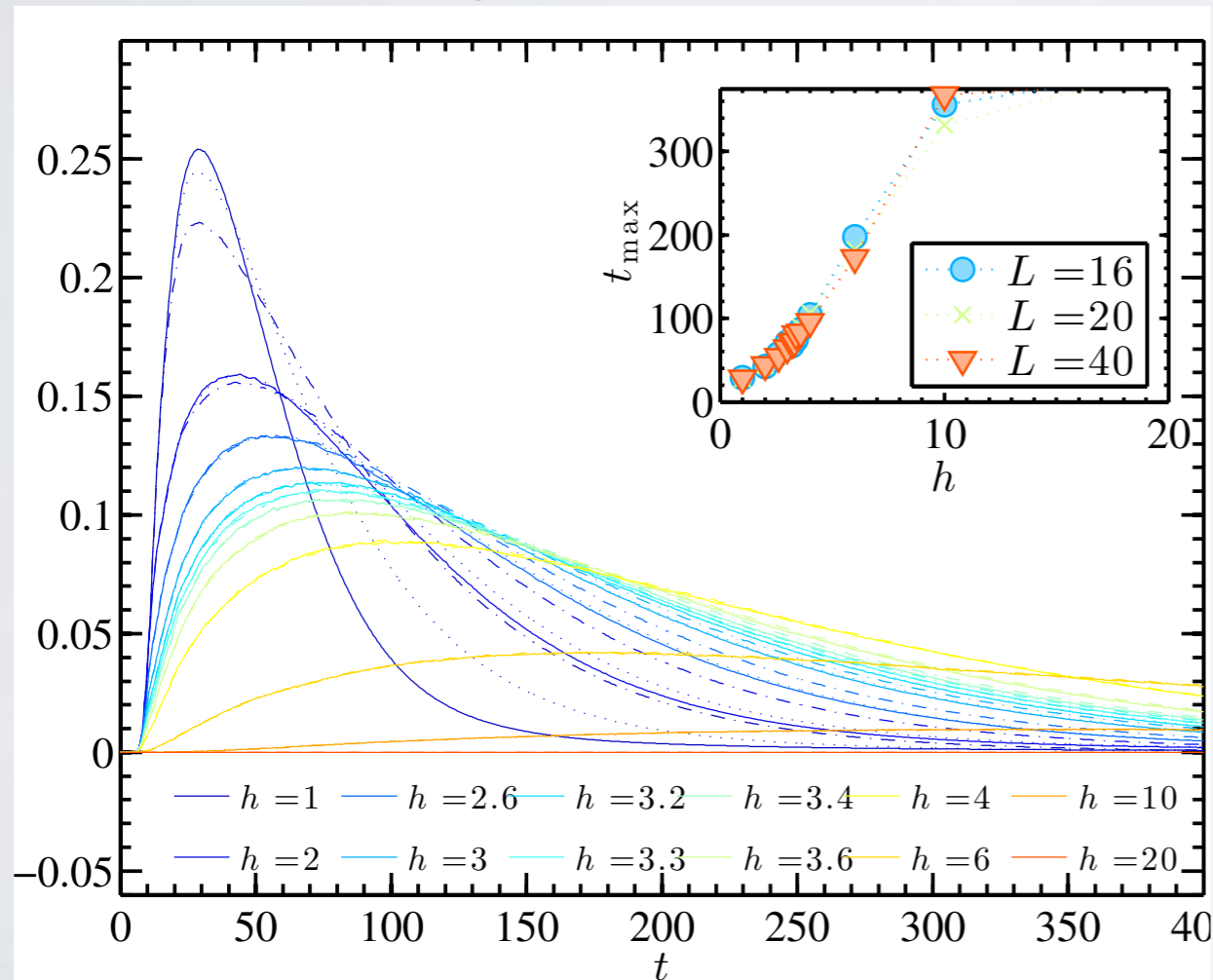


simulability with MPO

localization \Leftrightarrow errors in TN simulation

larger systems (Hilbert-Schmidt distance)

single instance



time

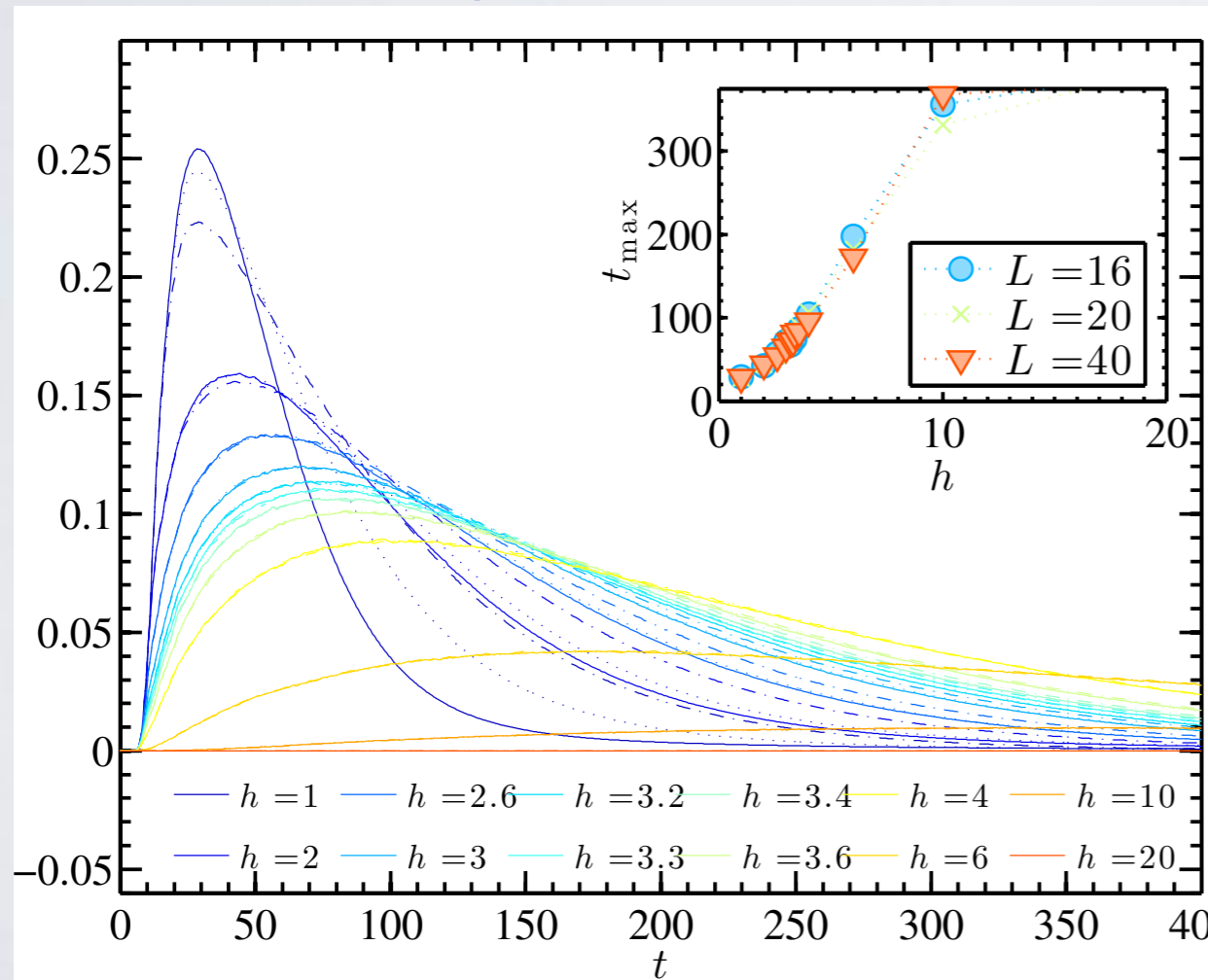


simulability with MPO

localization \Leftrightarrow errors in TN simulation

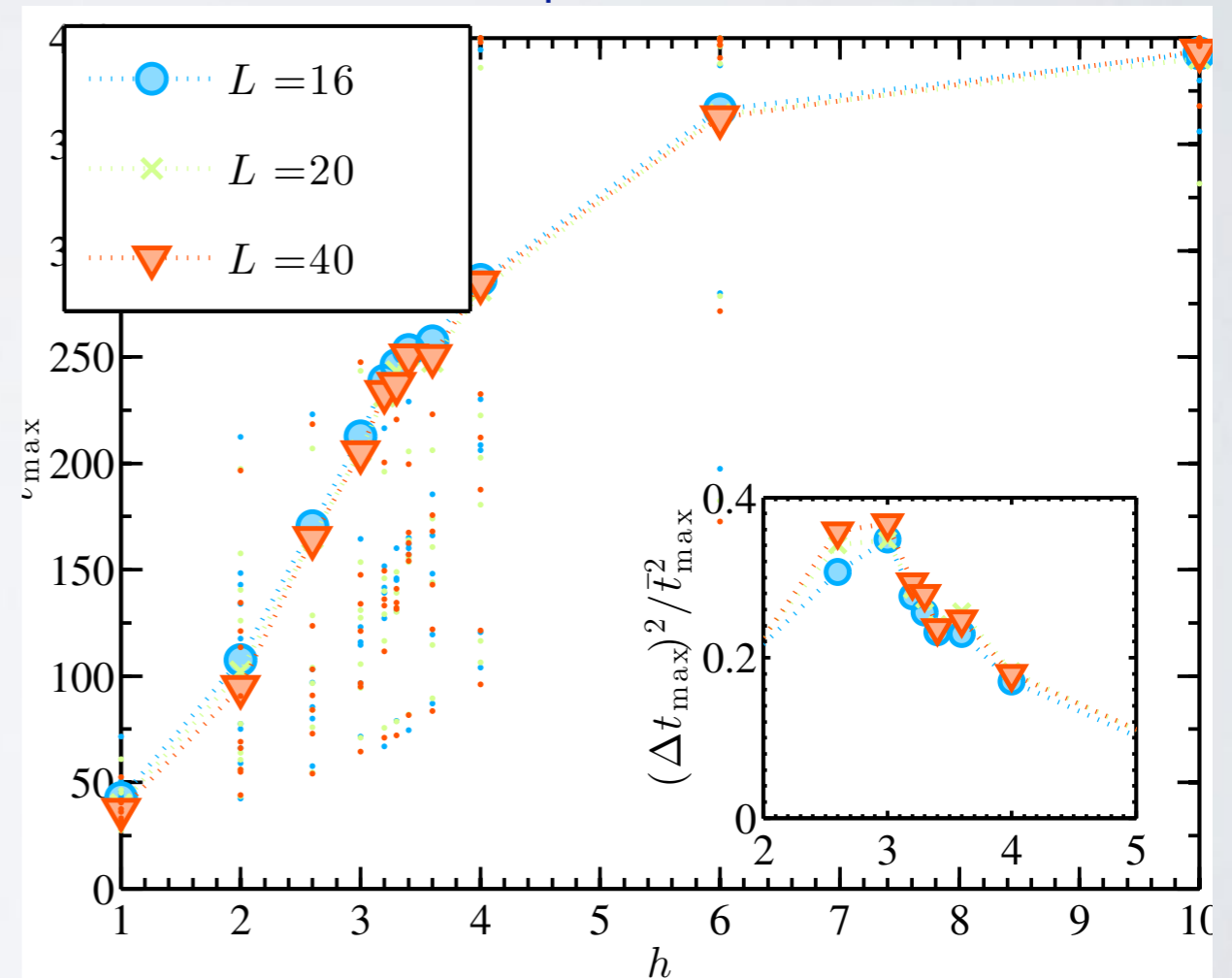
larger systems (Hilbert-Schmidt distance)

single instance



time

multiple realizations



disorder

Some questions we are asking

dynamics of mixed states

propagation of correlations ✓

quantum memory features ✓

simulability with MPO ✓

Hamiltonian properties

local conserved quantities ←

(ALMOST) LOCAL CONSERVED OPERATORS

What are the slowest evolving (local) operators?

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$$\frac{dA(t)}{dt} = i[H, A(t)]$$

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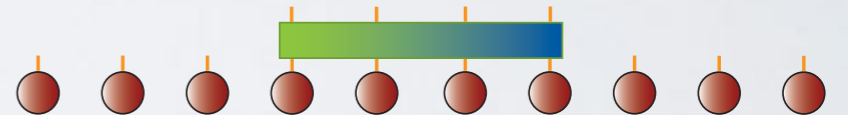
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(ALMOST) LOCAL CONSERVED OPERATORS

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operator acting on M sites

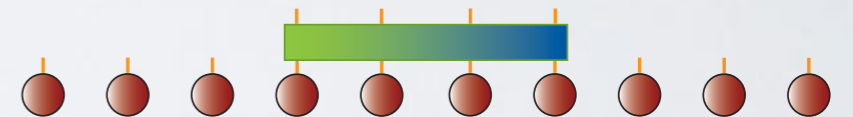


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Goal: minimizing $\|[H, A_M]\|$

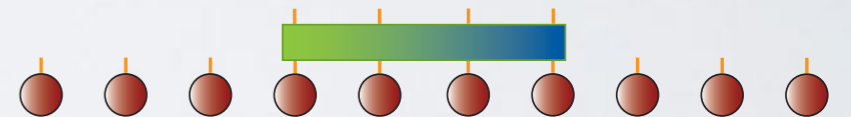
$$\lambda_M = \min_{A_M} \frac{\|[A_M, H]\|_2^2}{\|A_M\|_2^2}$$

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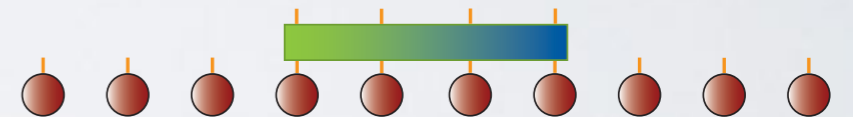


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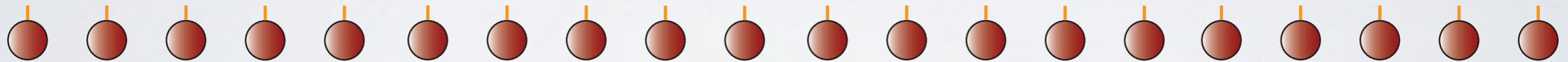
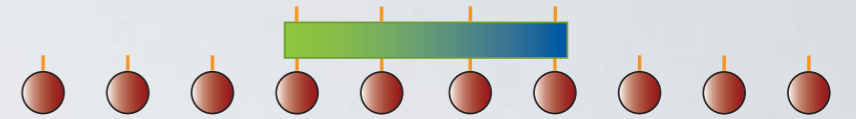
numerically with ED
and TNS

operator acting on M sites



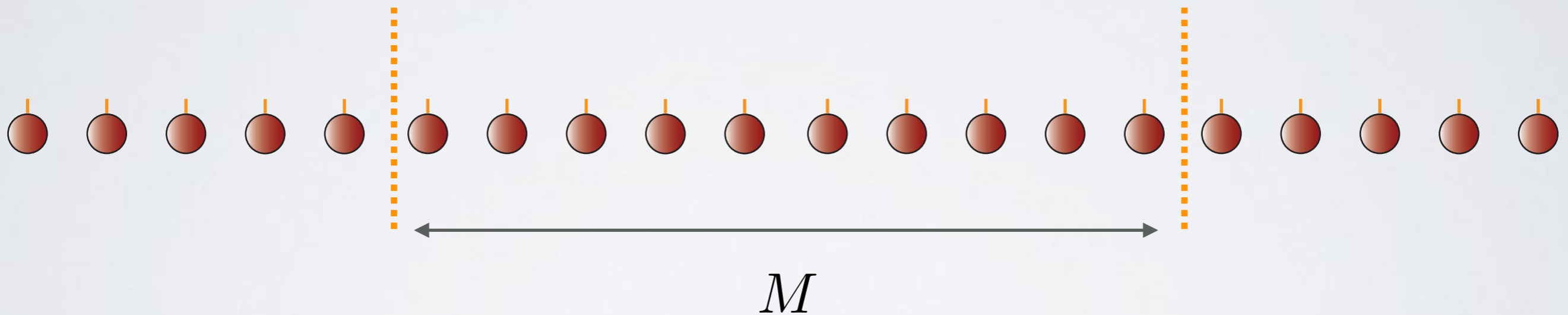
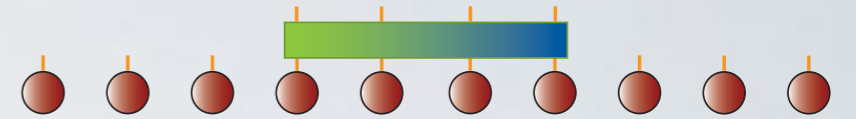
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operator acting on M sites



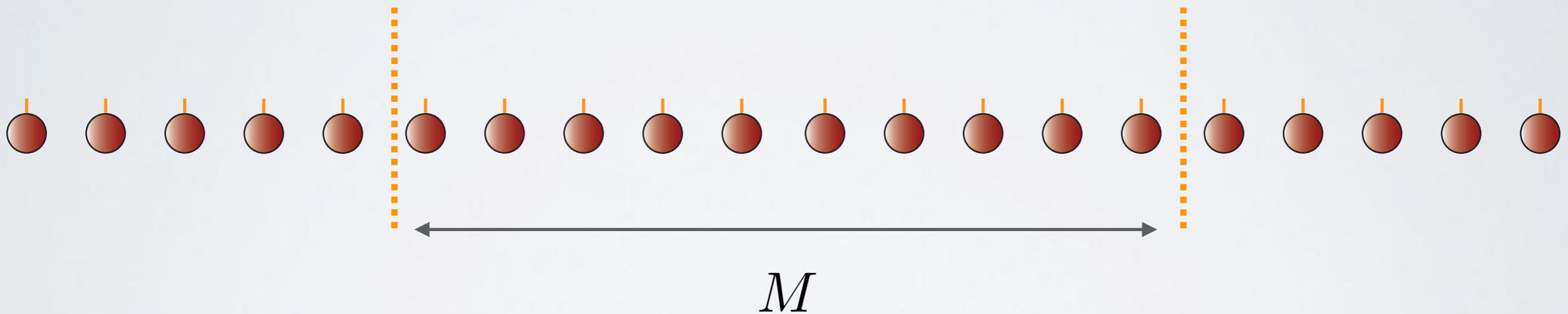
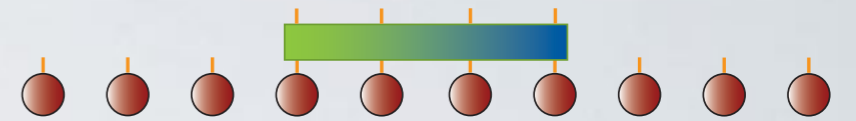
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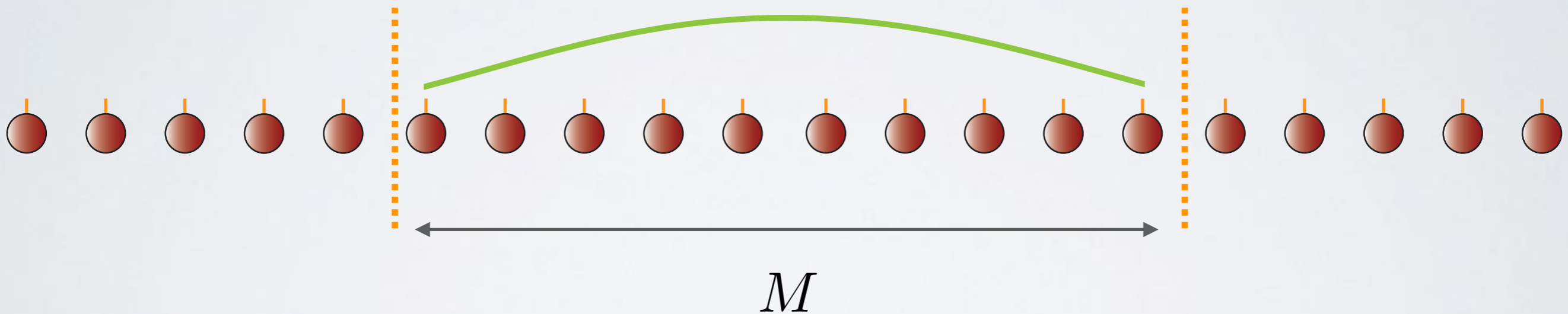
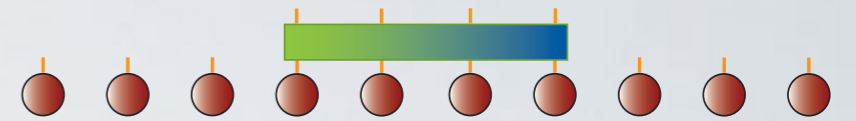
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operator acting on M sites



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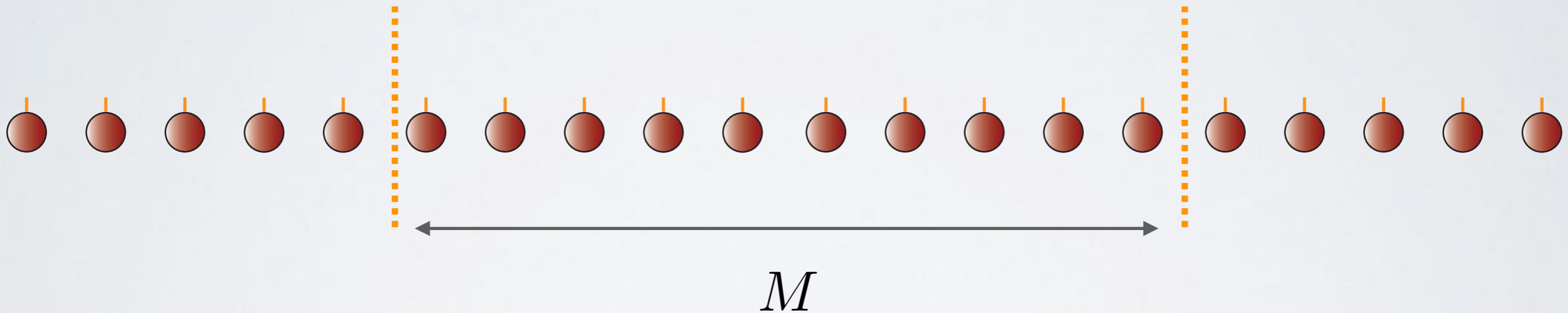
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$$\lambda_M = \min_{A_M} \frac{\|[A_M, H]\|_2^2}{\|A_M\|_2^2}$$

in the localized regime: I-bit model

$$H_{\text{eff}} = \sum_{i=0}^{N-1} \epsilon_i \tau_z^{[i]} + \sum_{i,j=0}^{N-1} K_{ij}^{(2)} \tau_z^{[i]} \tau_z^{[j]} + \sum_{i,j,k=0}^{N-1} K_{ijk}^{(3)} \tau_z^{[i]} \tau_z^{[j]} \tau_z^{[k]} + \dots,$$



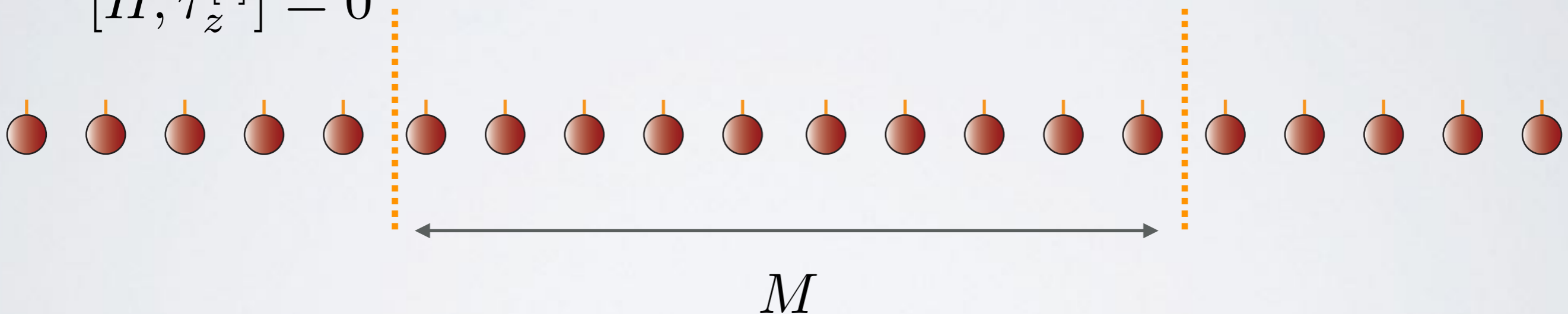
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in the localized regime: I-bit model

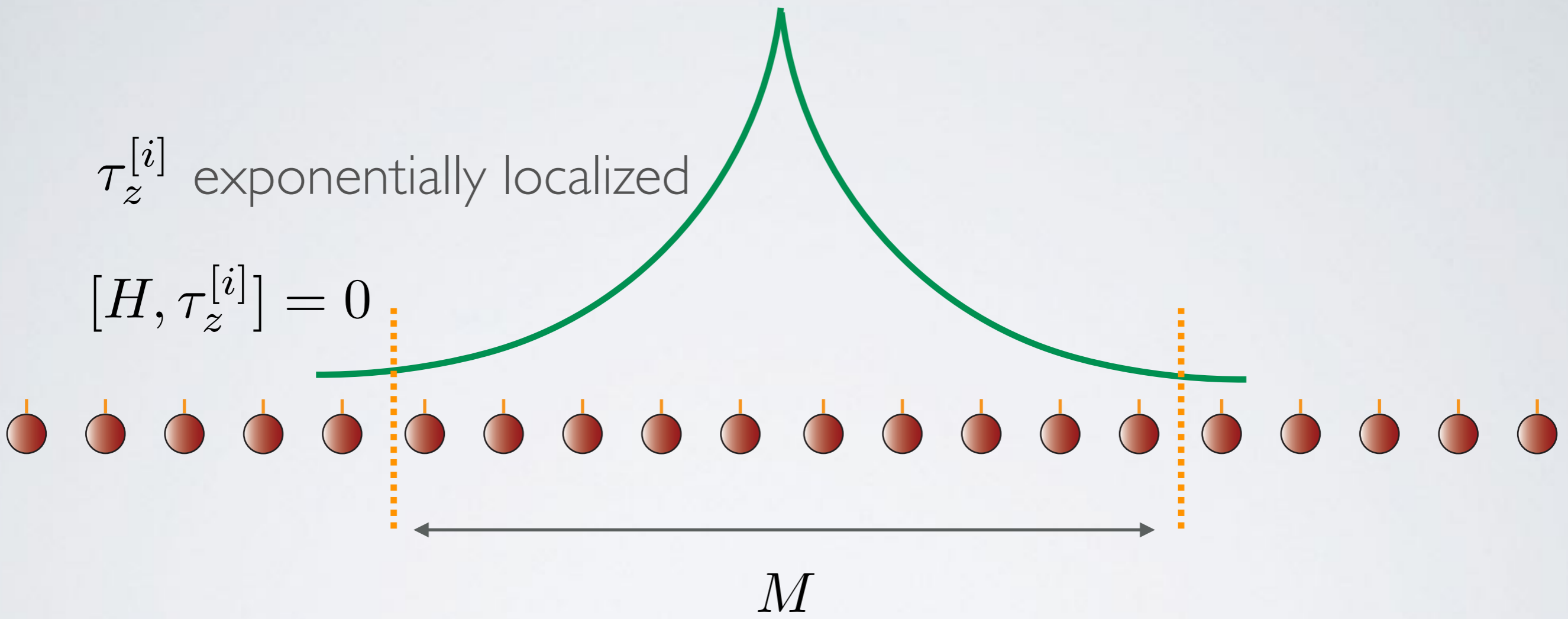
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$\tau_z^{[i]}$ exponentially localized

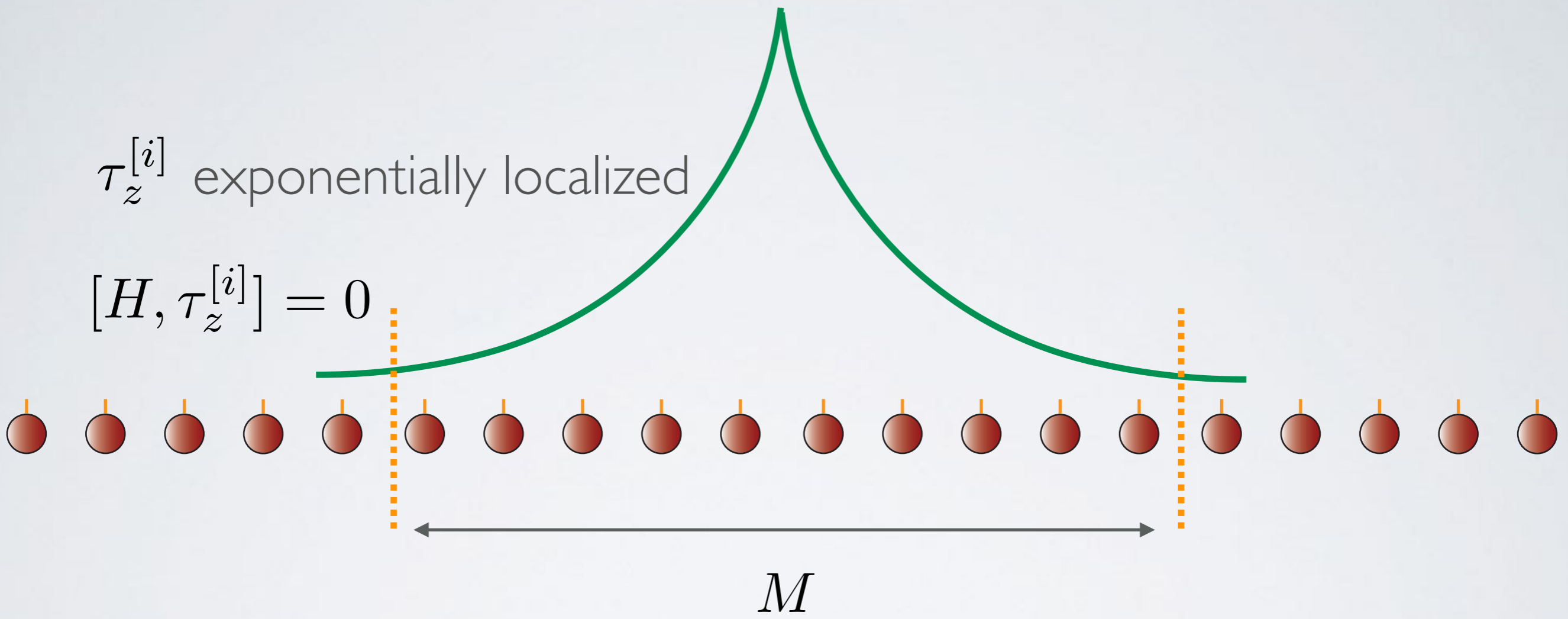
$$[H, \tau_z^{[i]}] = 0$$



$$\lambda_M = \min_{A_M} \frac{\|[A_M, H]\|_2^2}{\|A_M\|_2^2}$$



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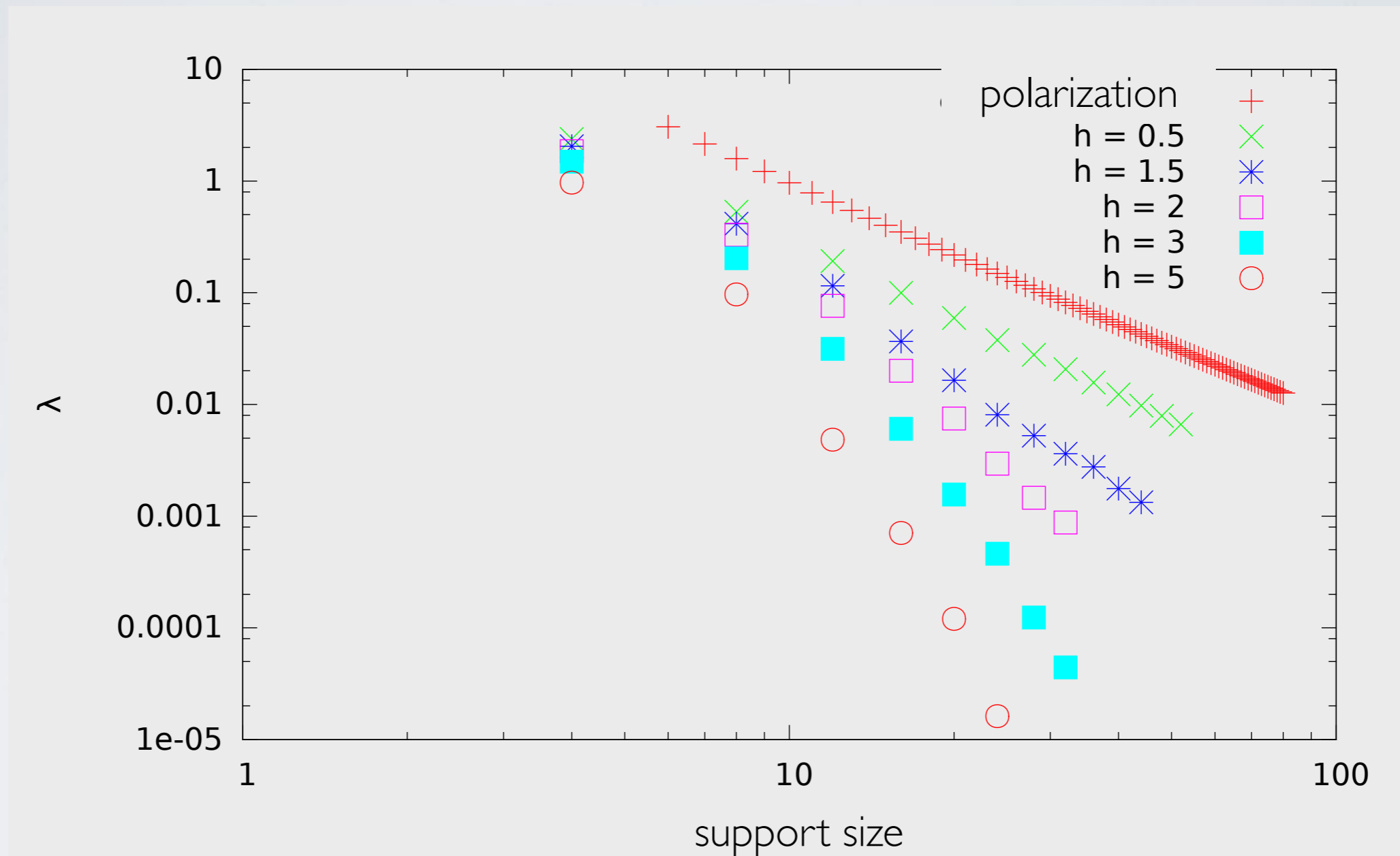
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truncated support \rightarrow expect exponentially small

see also Chandran et al. PRB 2015

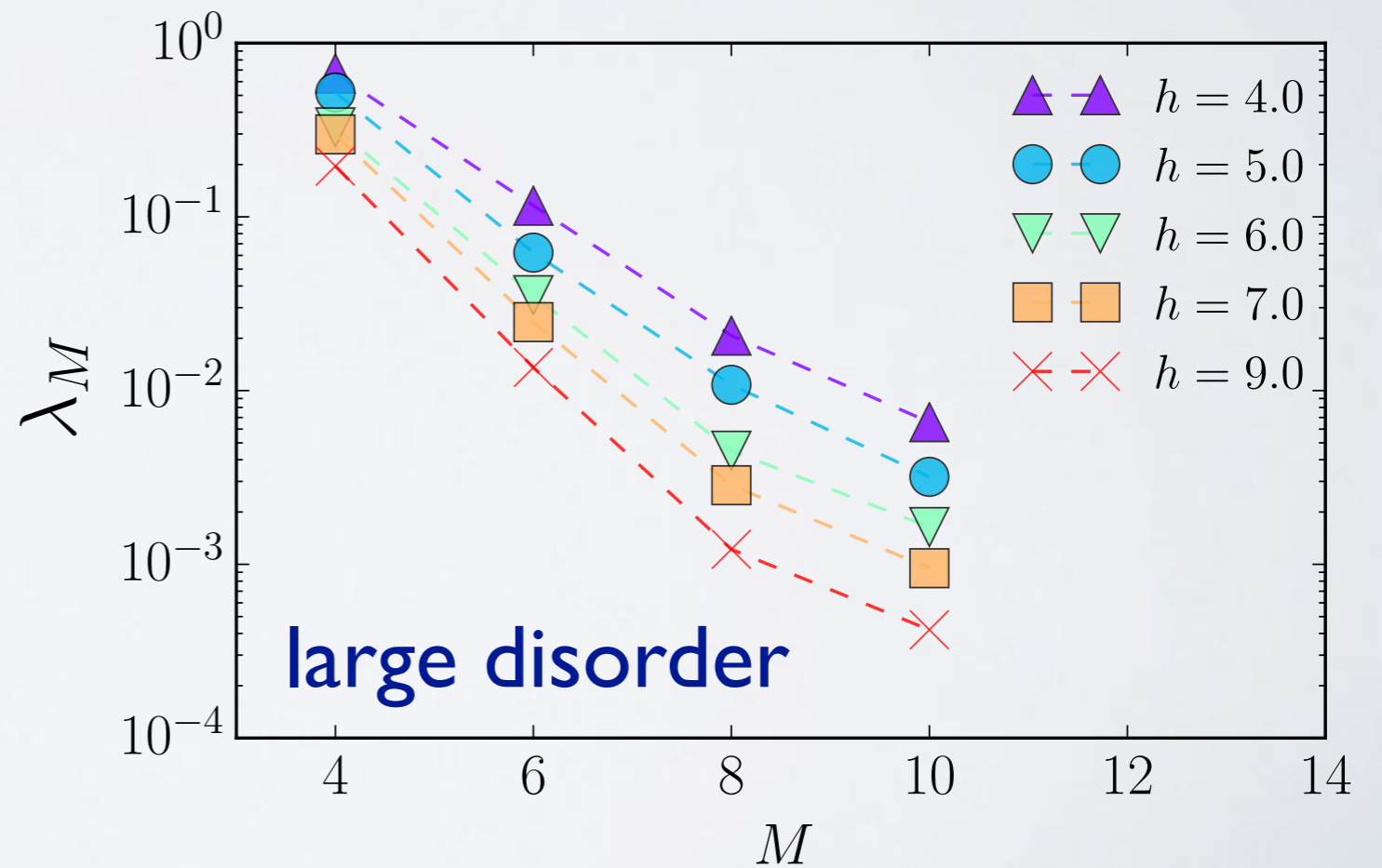
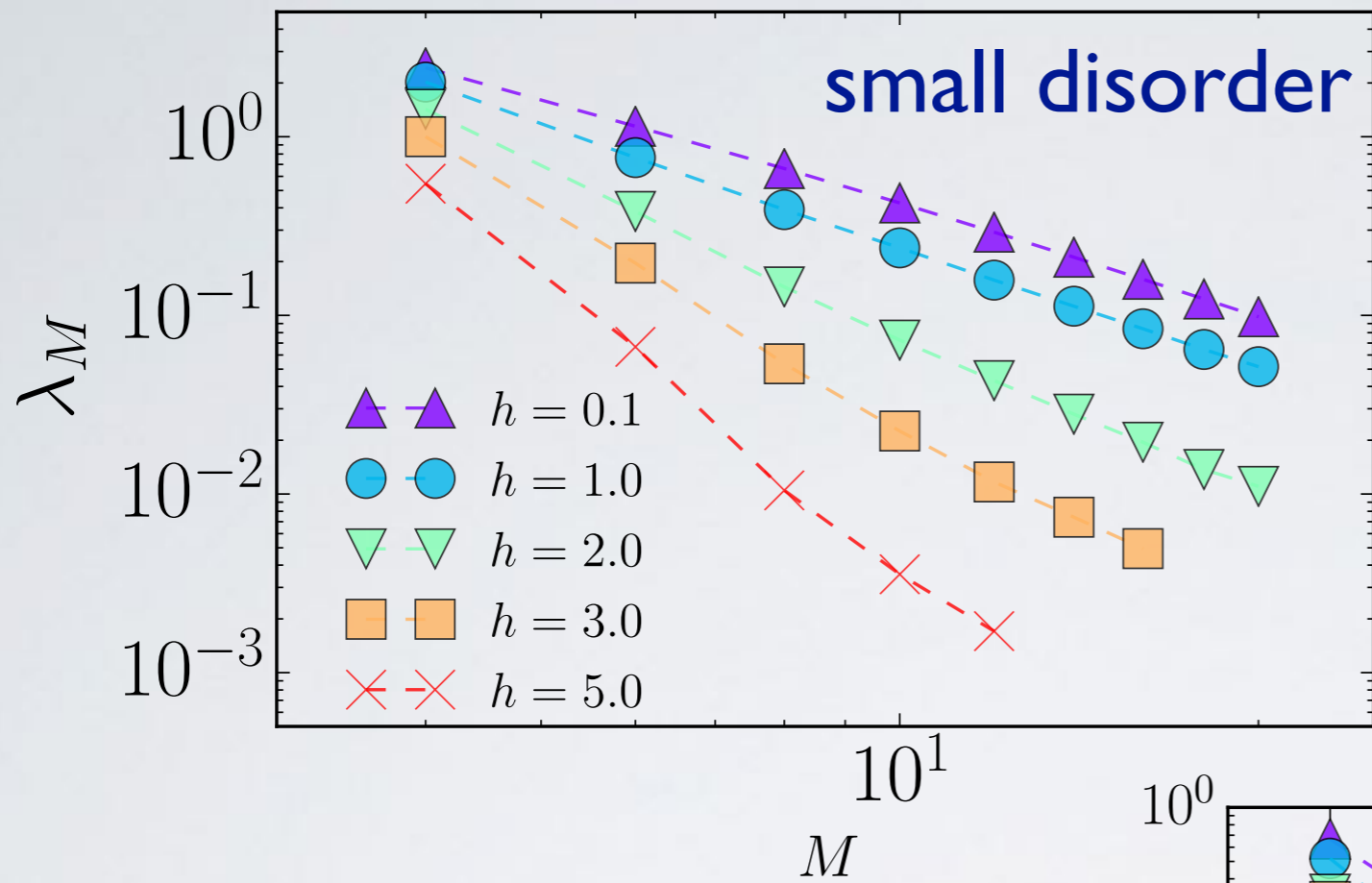
N. Pancotti et al PRB 97, 094206 (2018)

ALMOST CONSERVED QUANTITIES & MBL

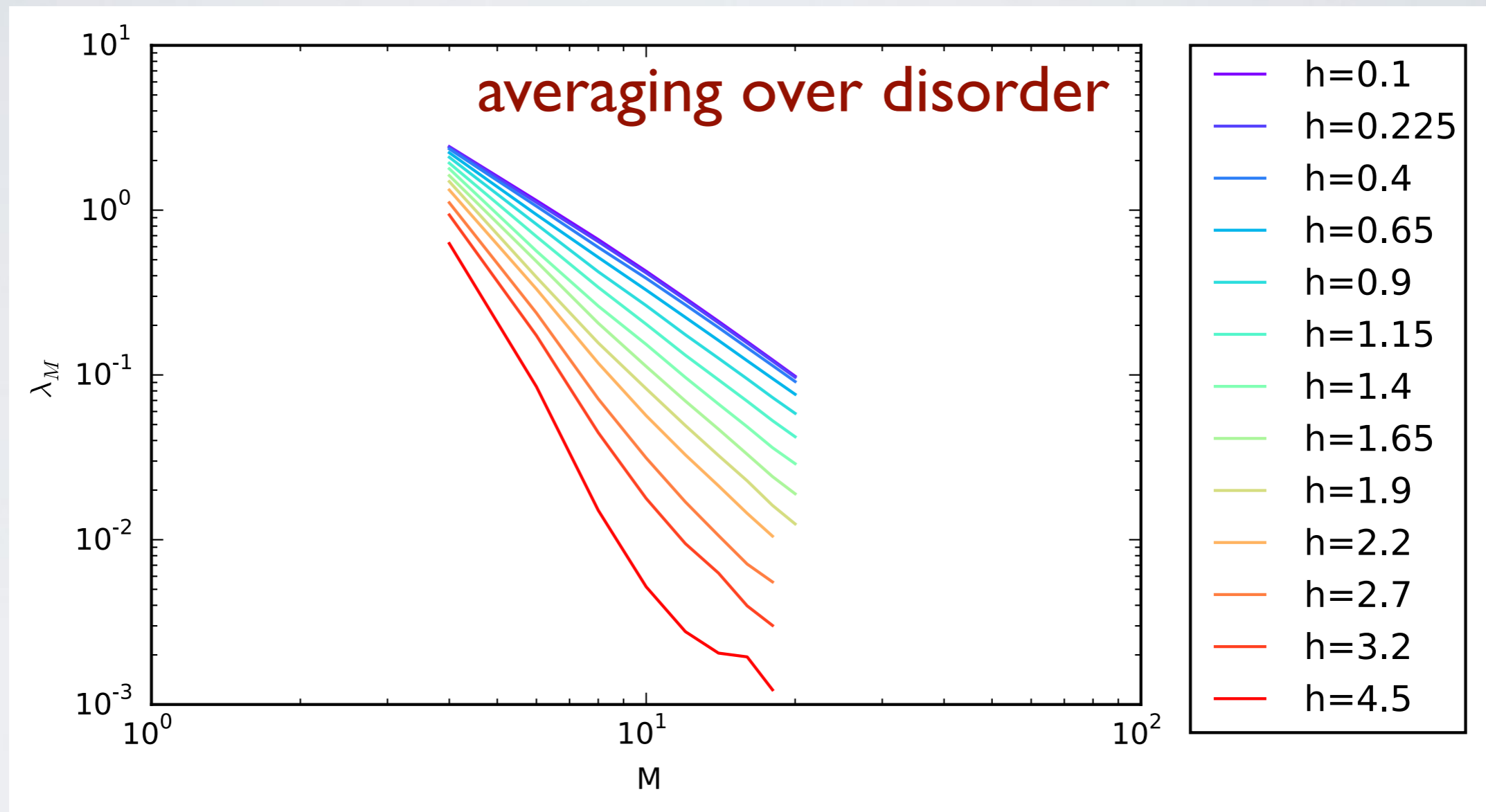


example: single disorder realization

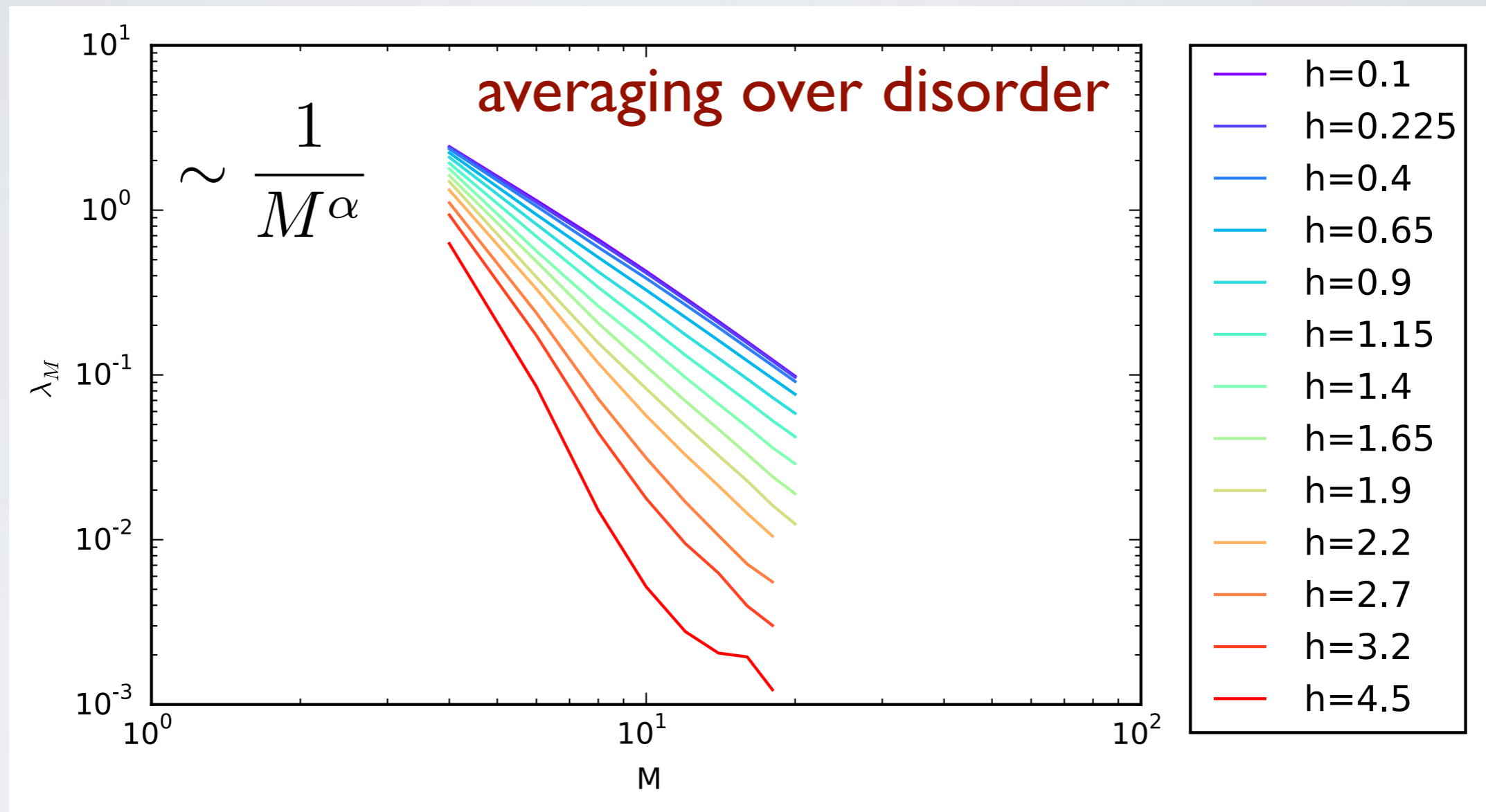
averaging over disorder



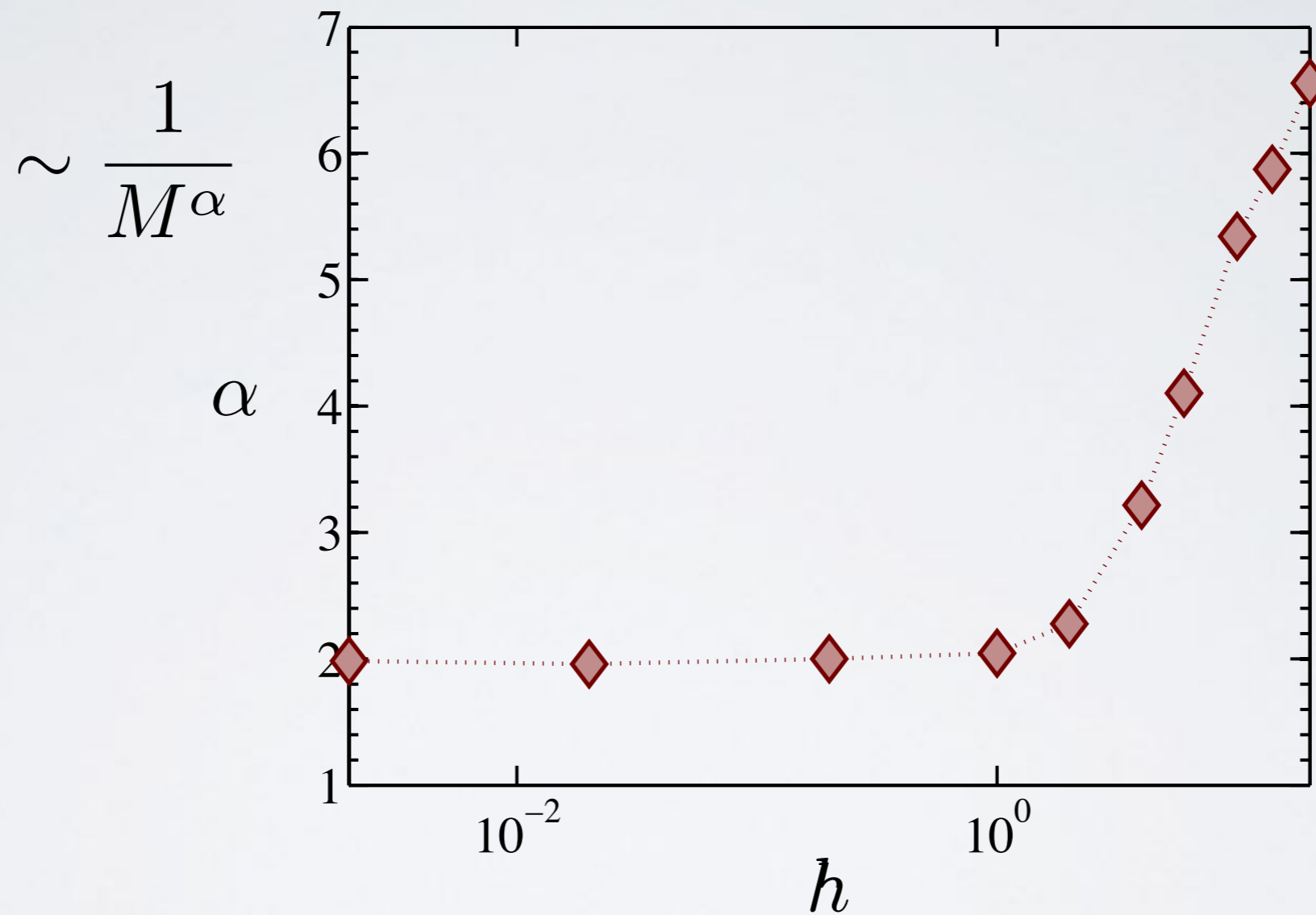
ALMOST CONSERVED QUANTITIES & MBL



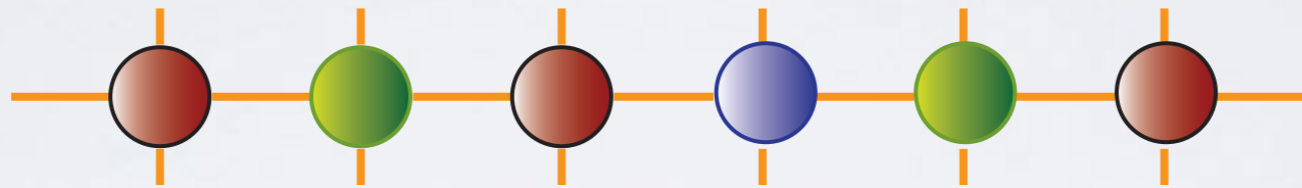
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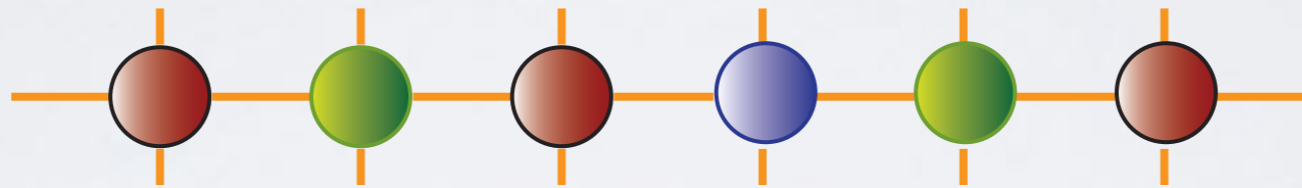
ALMOST CONSERVED QUANTITIES & MBL



constructive method



constructive method

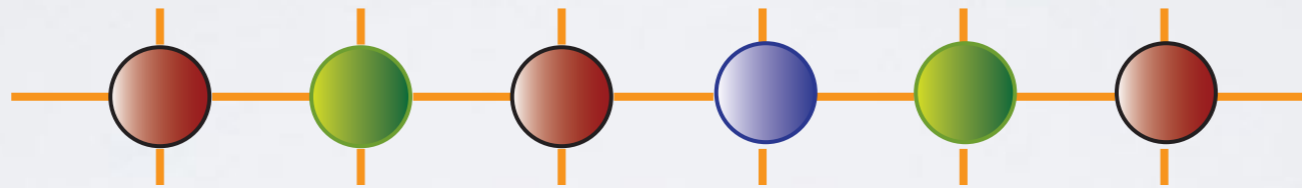


analyze weight of components with different support

$$\sigma_i^{[m]} \otimes \dots \otimes \sigma_j^{[m+d]}$$

constructive method

efficient! → simple
projector on MPO

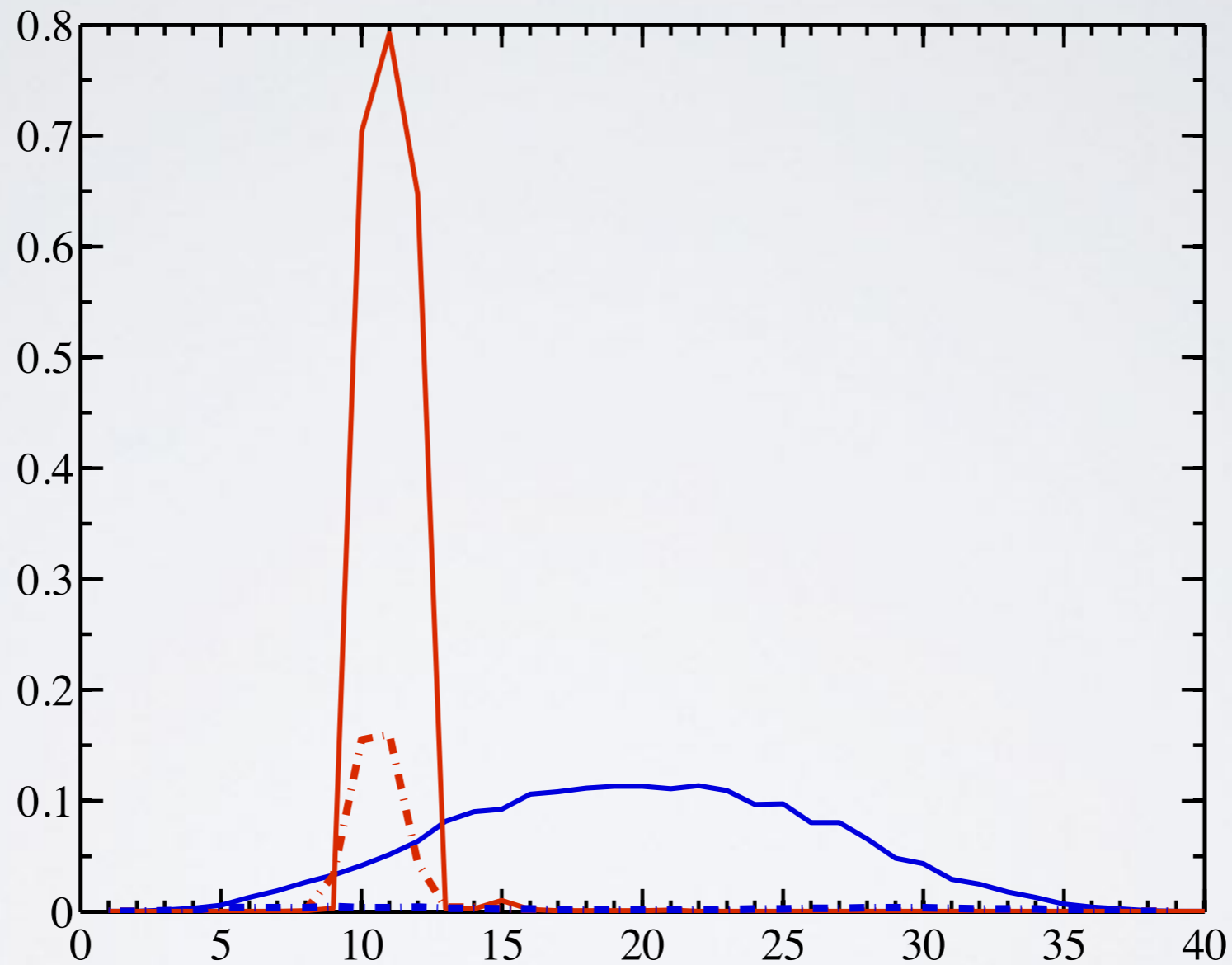


analyze weight of components with different support

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composition of slow operators: how local?

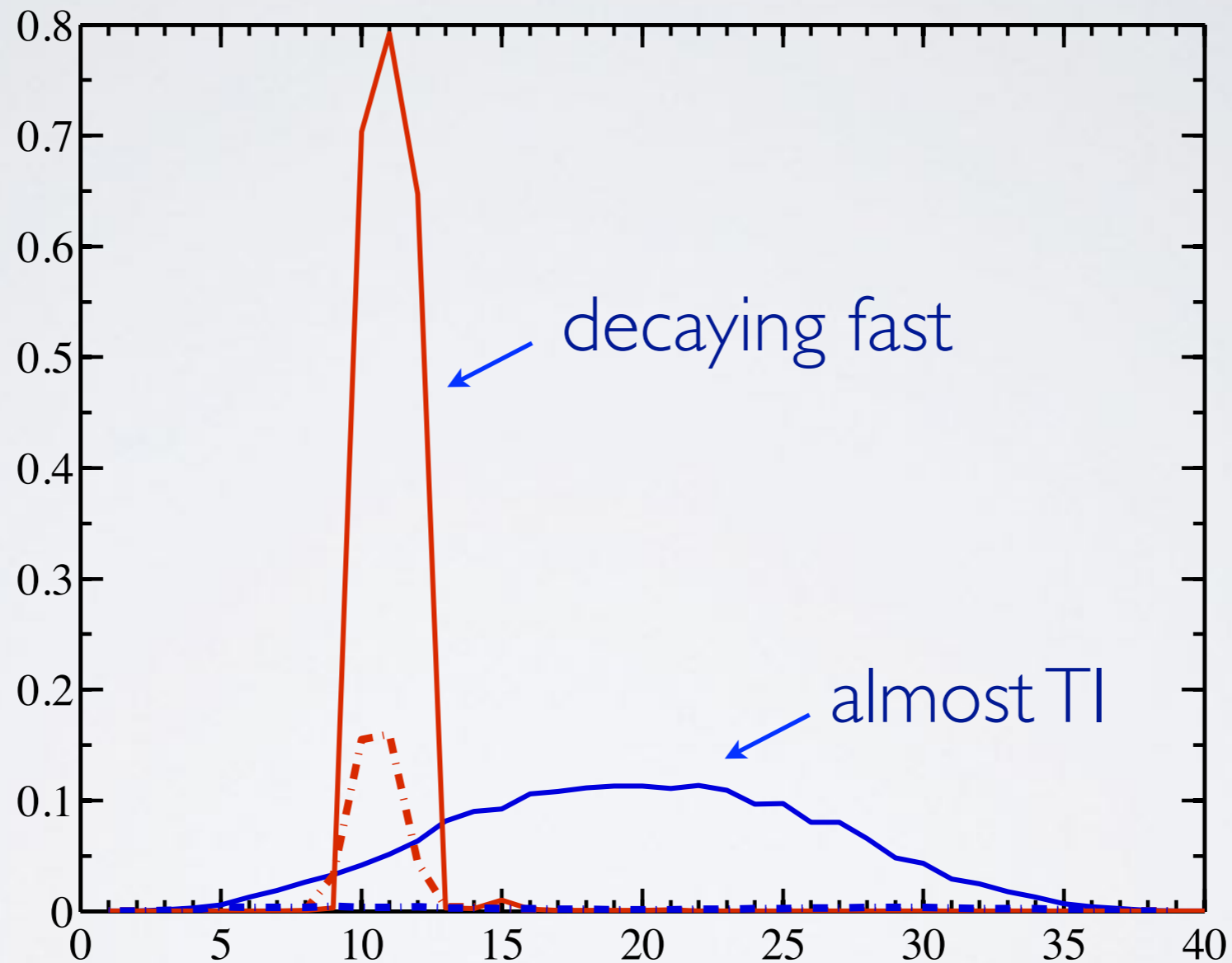
landscape of terms with fixed range



single realization $M = 40$

composition of slow operators: how local?

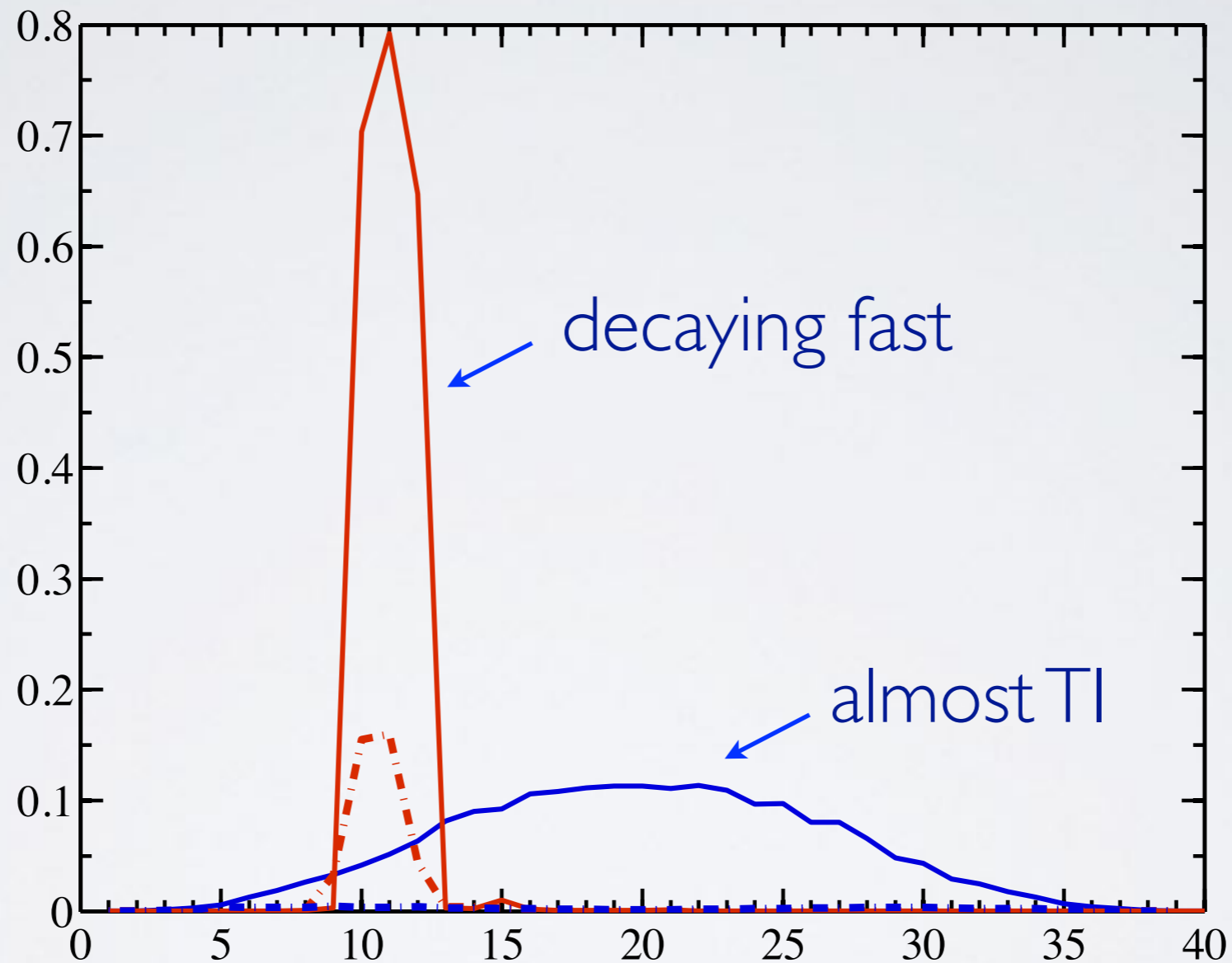
landscape of terms with fixed range



single realization $M = 40$

composition of slow operators: how local?

landscape of terms with fixed range



and much more information

in the statistics!

single realization $M = 40$

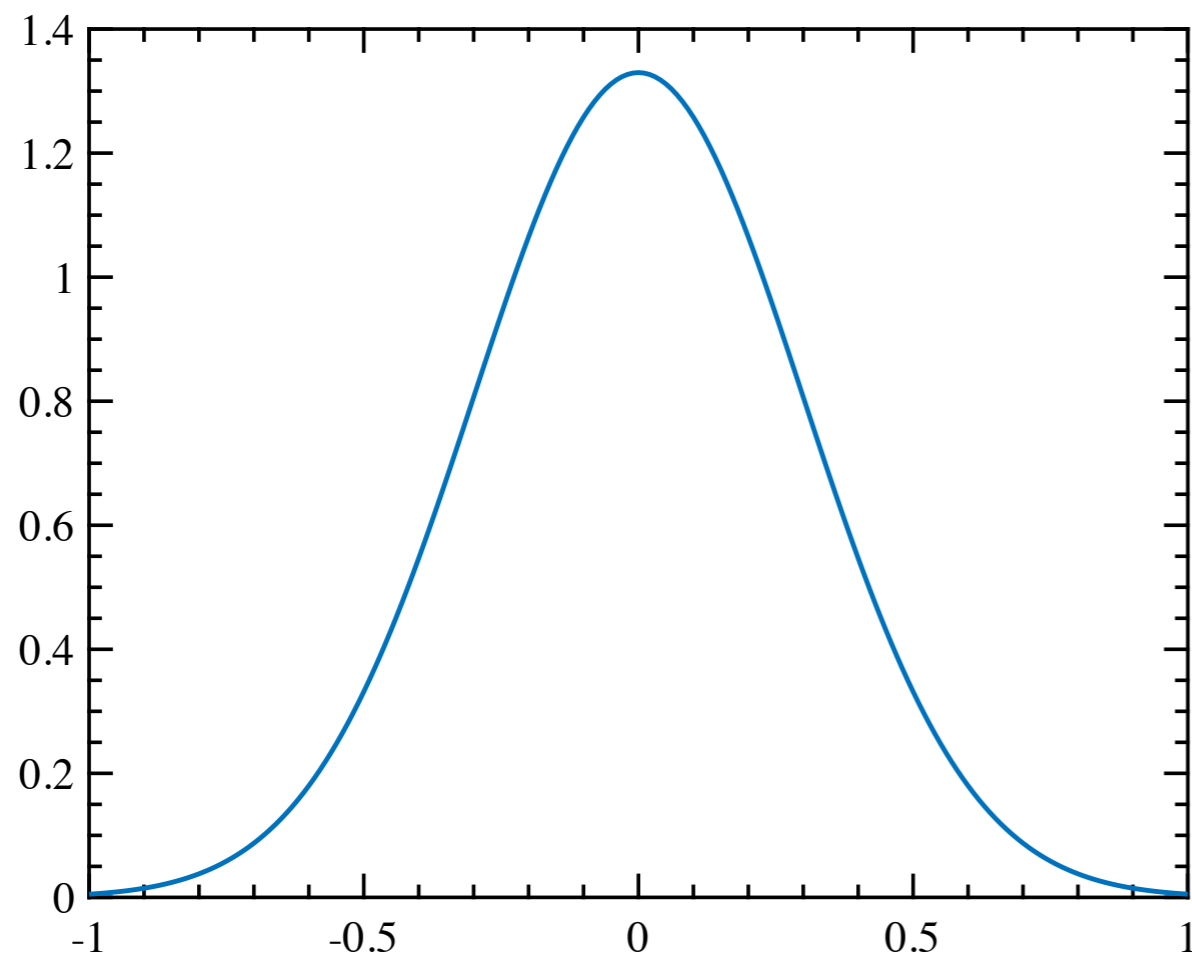
Statistics of small commutators

Well described by Extreme Value Theory

Statistics of small commutators

Well described by Extreme Value Theory

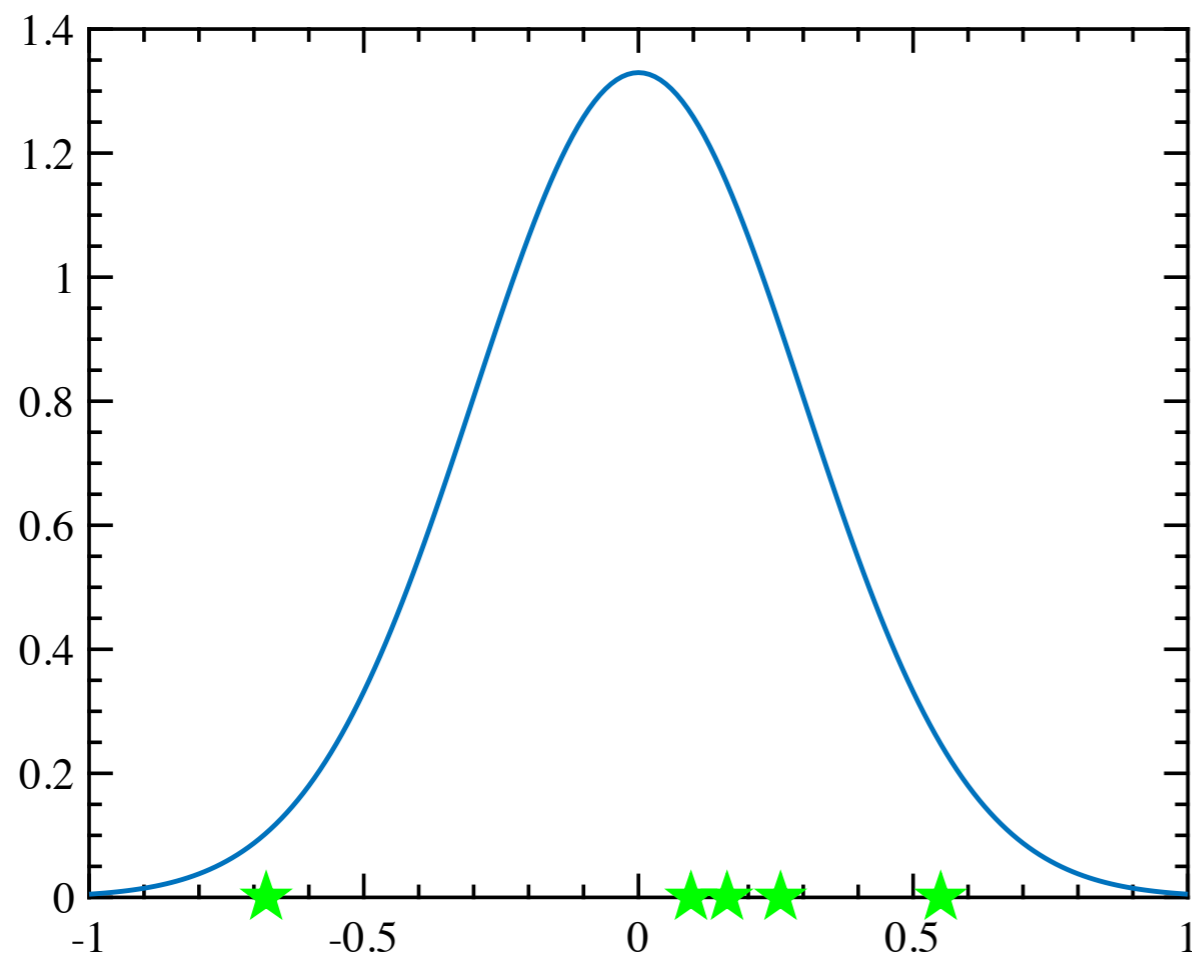
Q: extreme values when sampling from a pdf



Statistics of small commutators

Well described by Extreme Value Theory

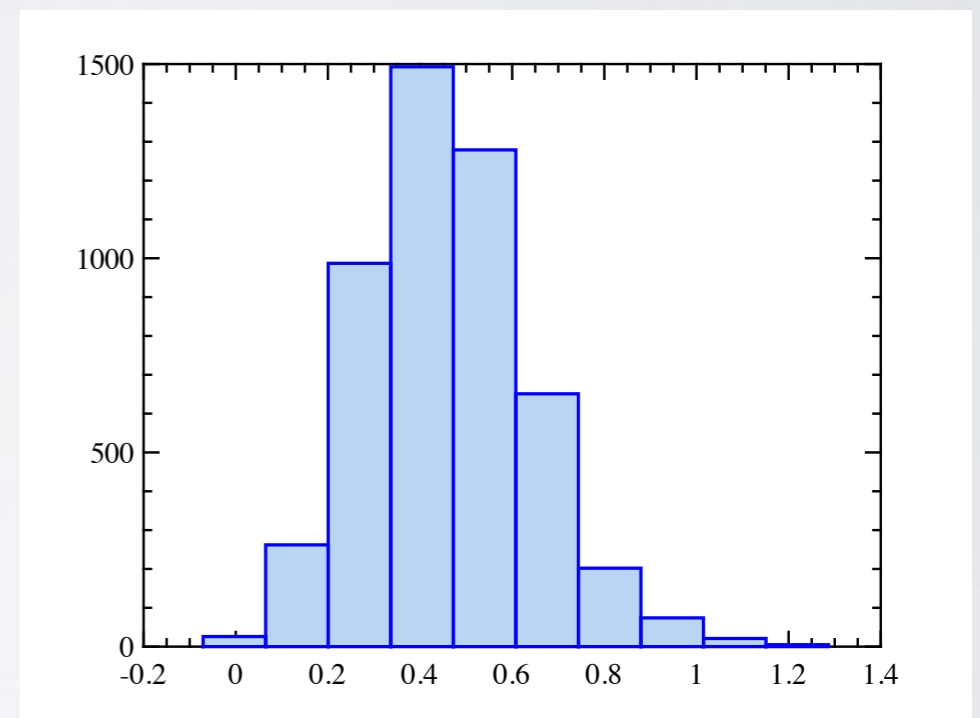
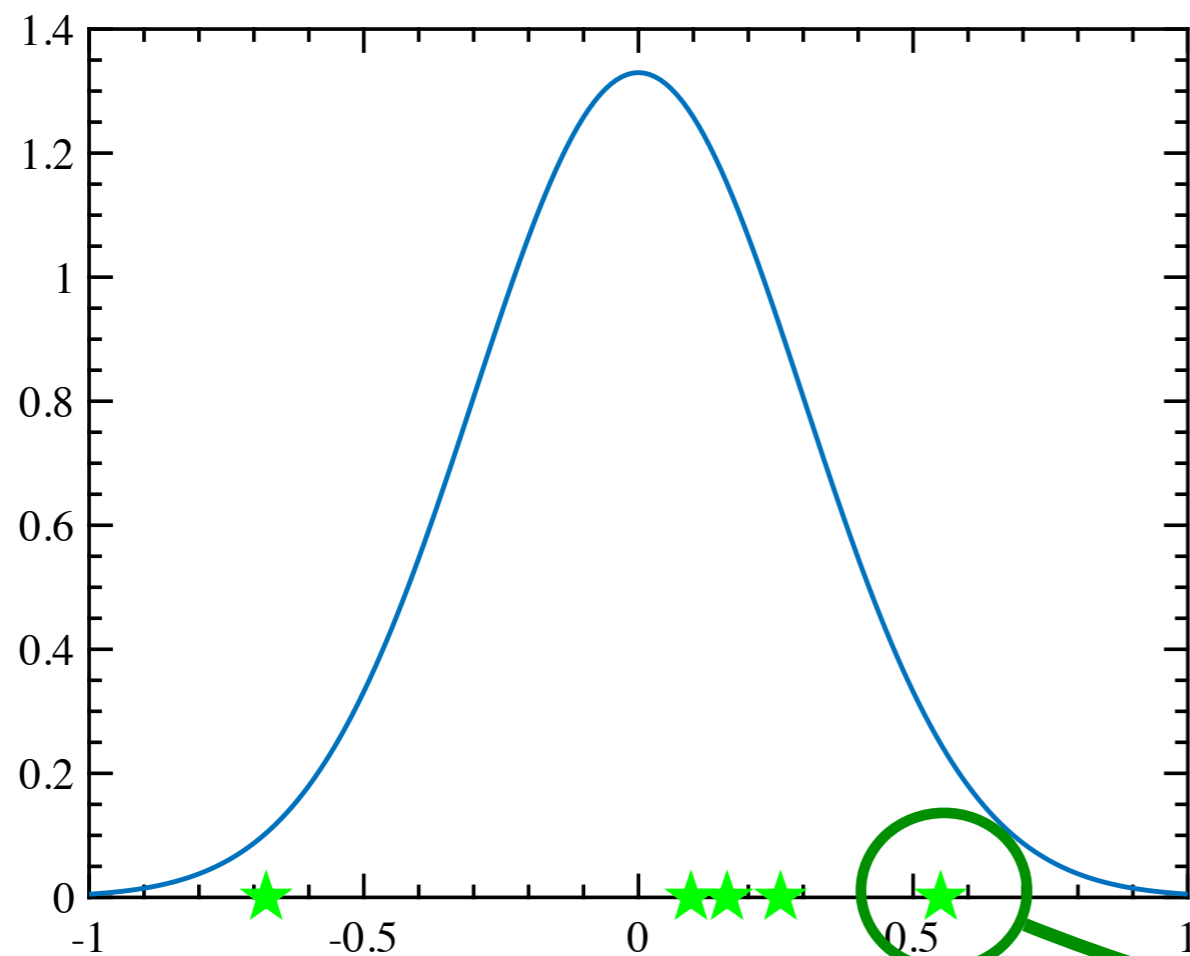
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Well described by Extreme Value Theory

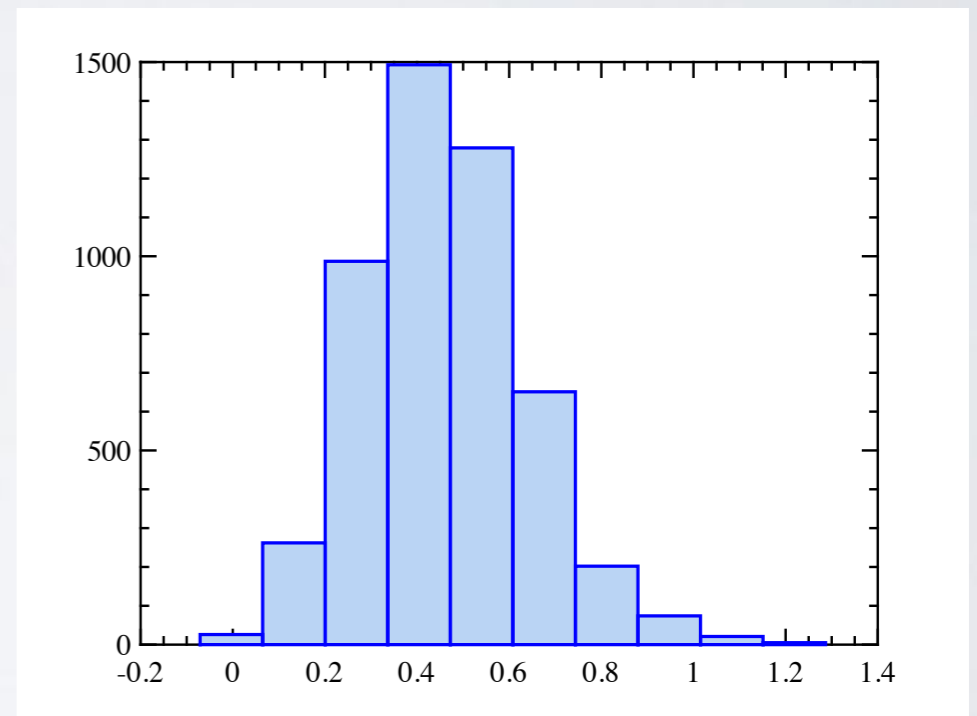
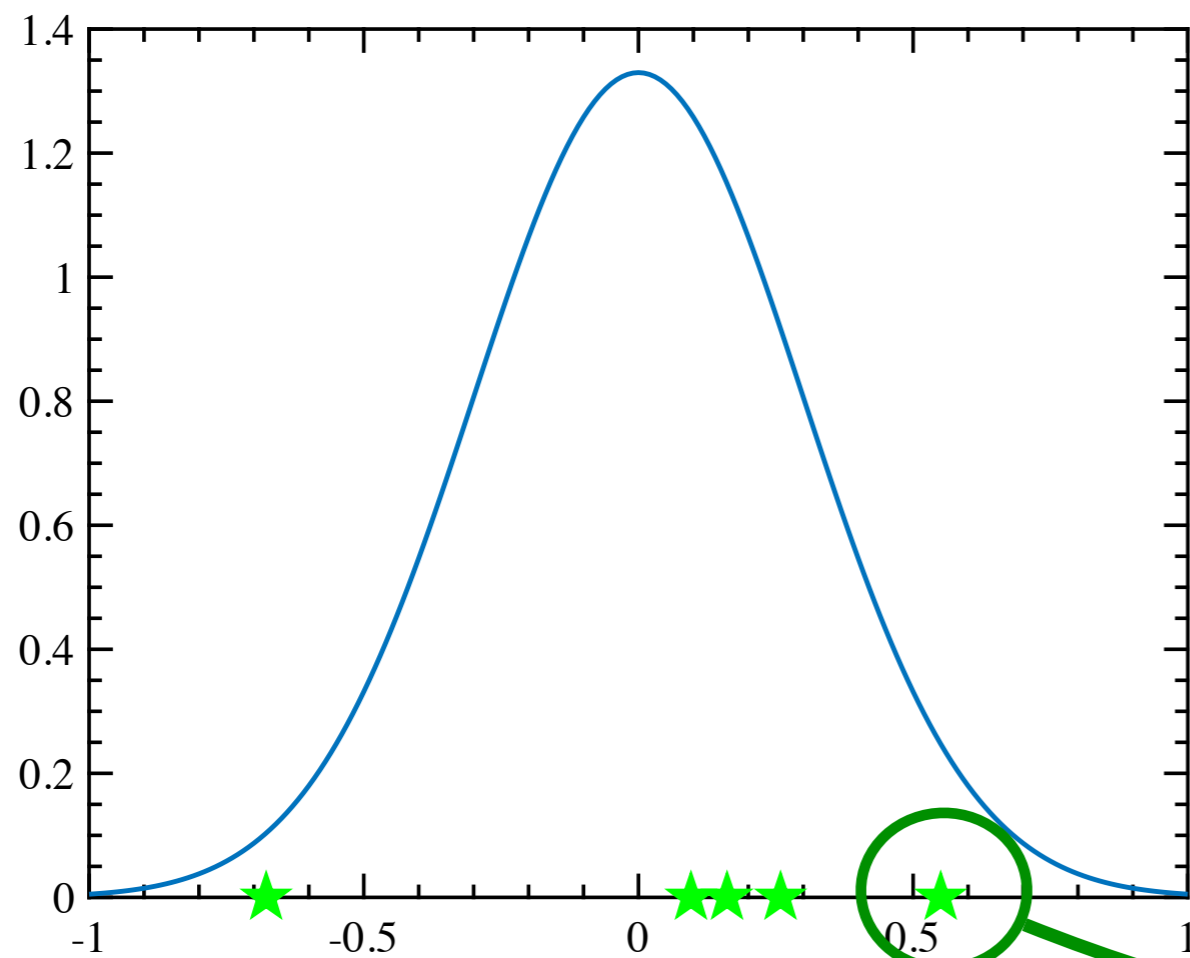
Q: extreme values when sampling from a pdf



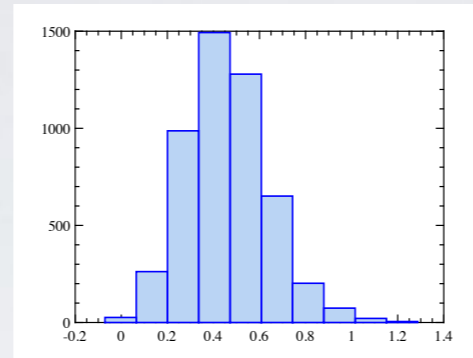
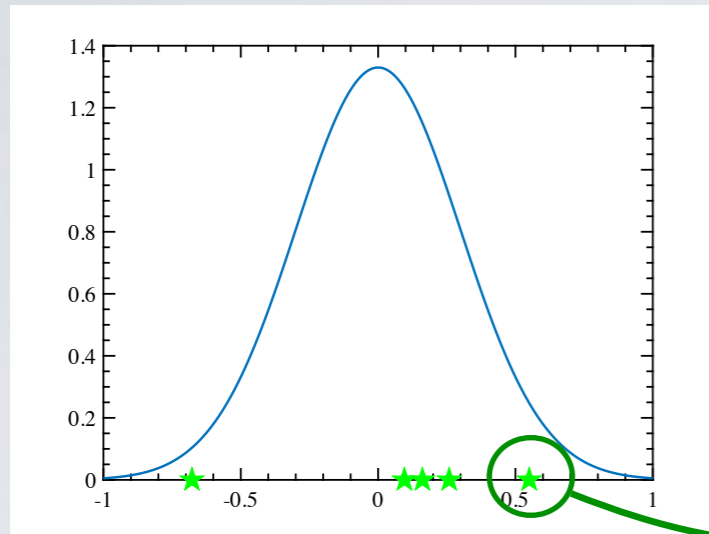
Statistics of small commutators

EVT

when is there a limit,
it is of the form



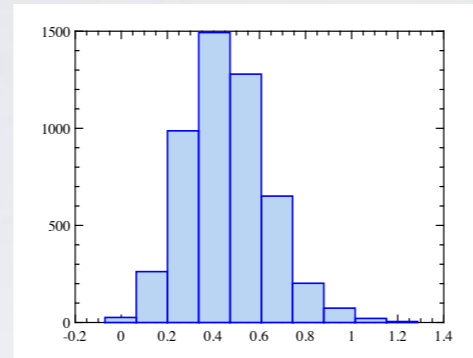
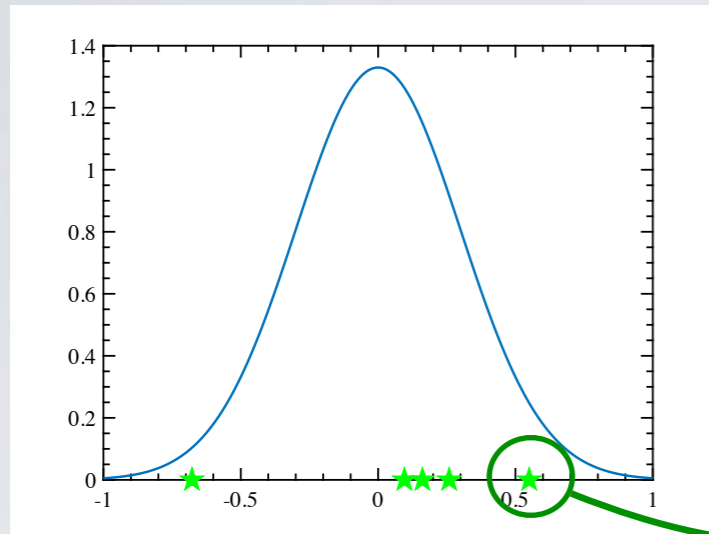
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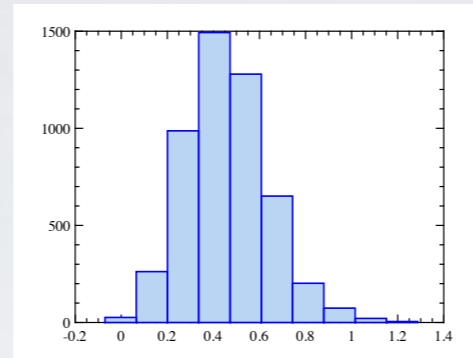
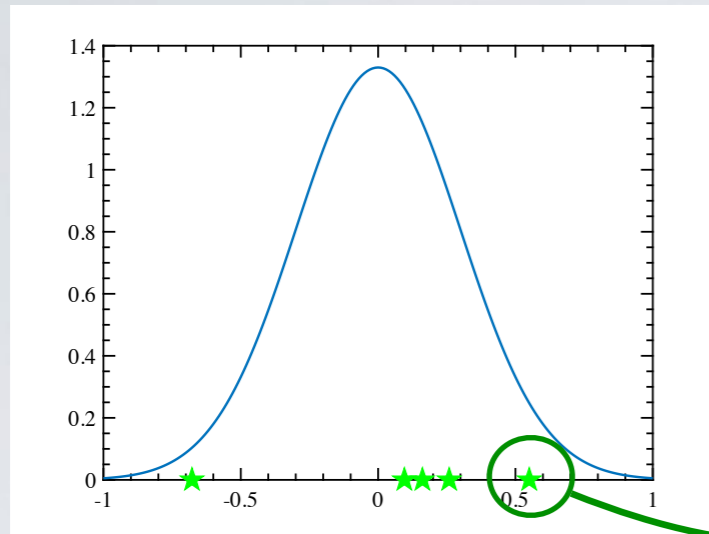
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GEV

$$G_{\zeta}(y) = \exp \left[- (1 + \zeta y)^{-\frac{1}{\zeta}} \right] \quad \text{rescaled and centered}$$

CDF for extrema

Statistics of small commutators



EVT

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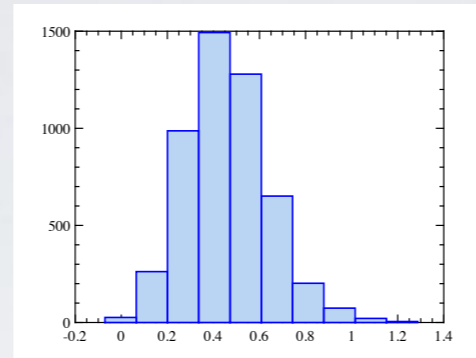
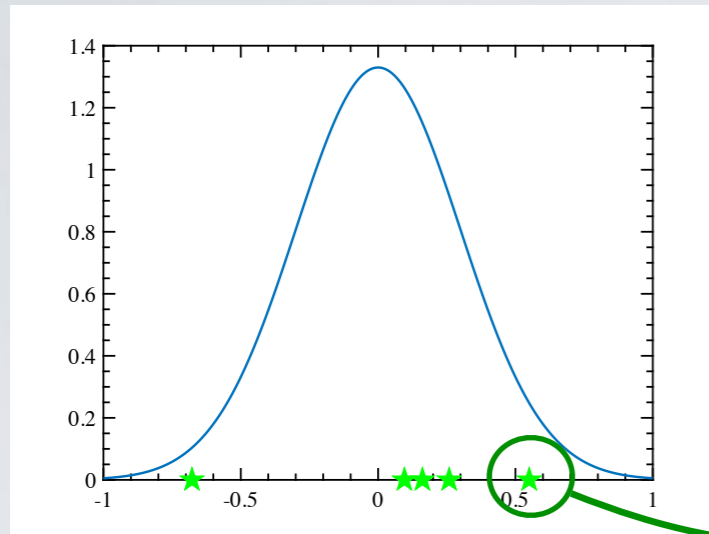
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CDF for extrema

three subfamilies

Statistics of small commutators



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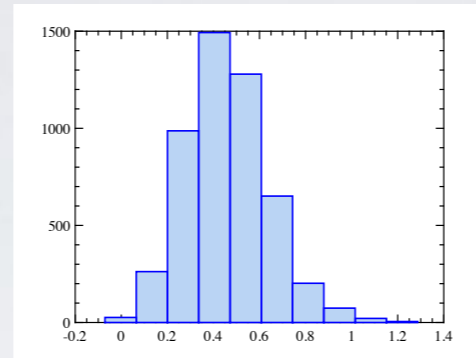
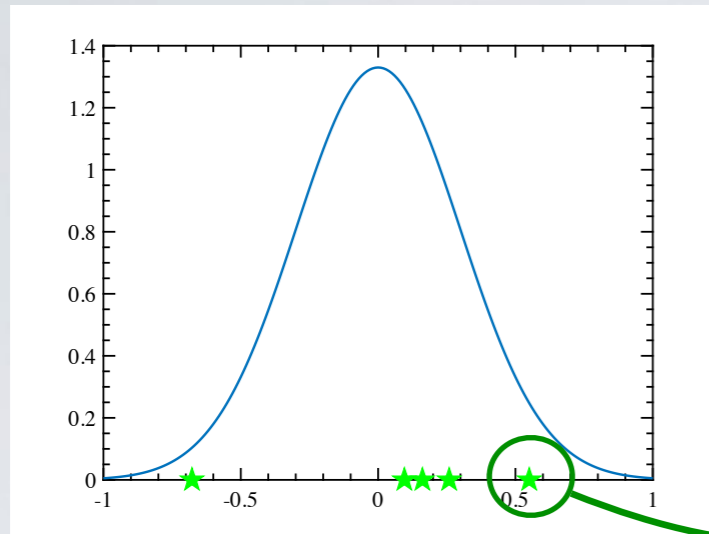
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$\zeta > 0$ Fréchet: polynomial tails

Statistics of small commutators



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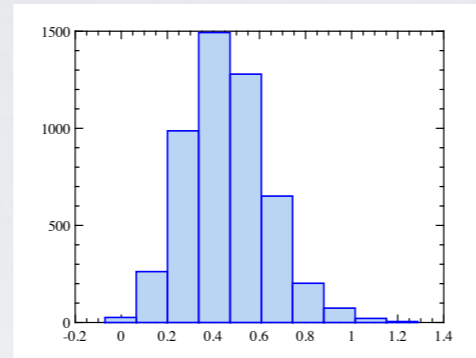
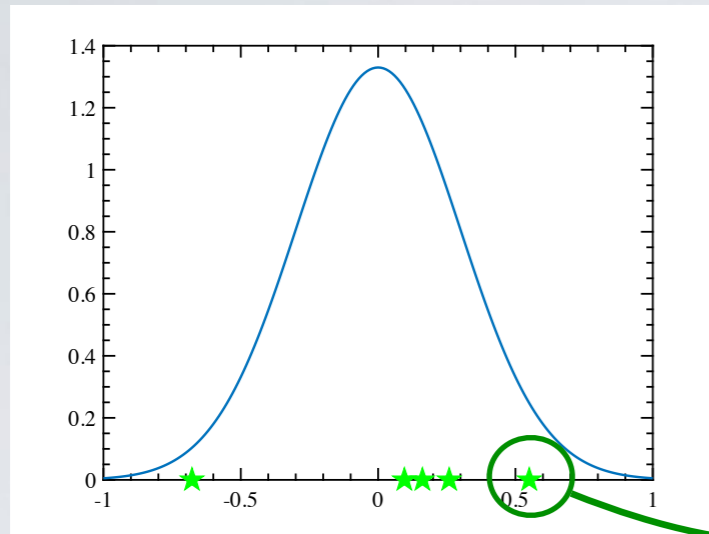
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$\zeta > 0$ Fréchet: polynomial tails

$\zeta = 0$ Gumbel: exponential tails $G_0(y) \rightarrow \exp[-\exp(-y)]$

Statistics of small commutators



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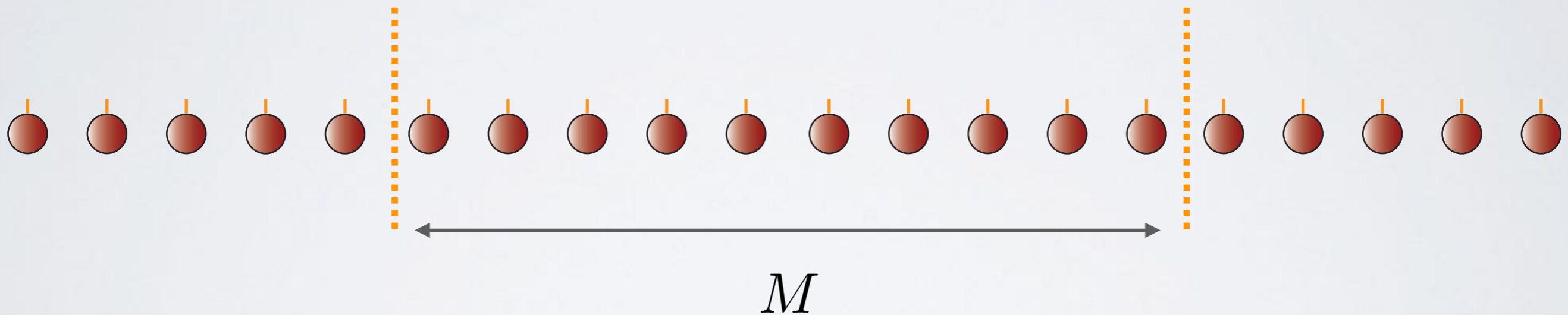
$\zeta = 0$ Gumbel: exponential tails $G_0(y) \rightarrow \exp[-\exp(-y)]$

$\zeta < 0$ Weibull: bounded light tails

Statistics of small commutators

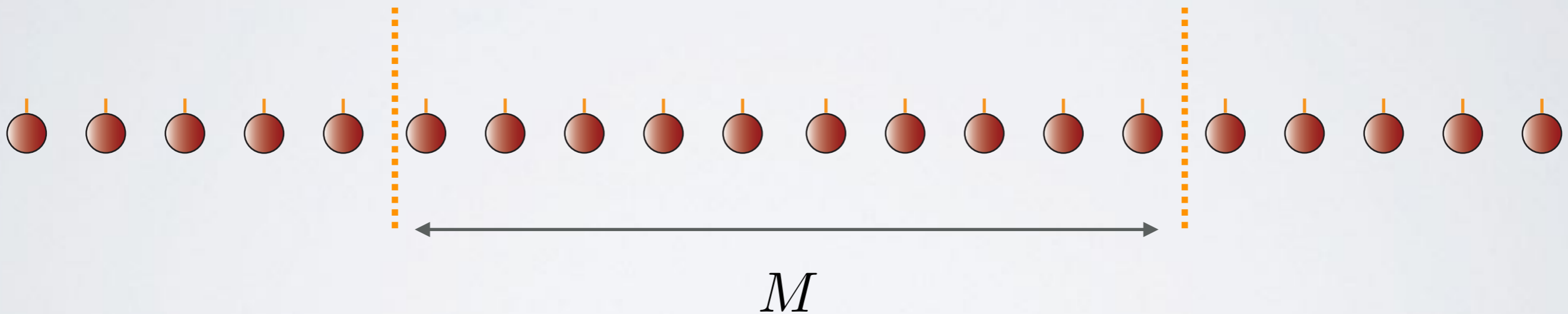


Statistics of small commutators



$$\lambda_M = \min_{A_M} \frac{\|[A_M, H]\|_2^2}{\|A_M\|_2^2}$$

Statistics of small commutators

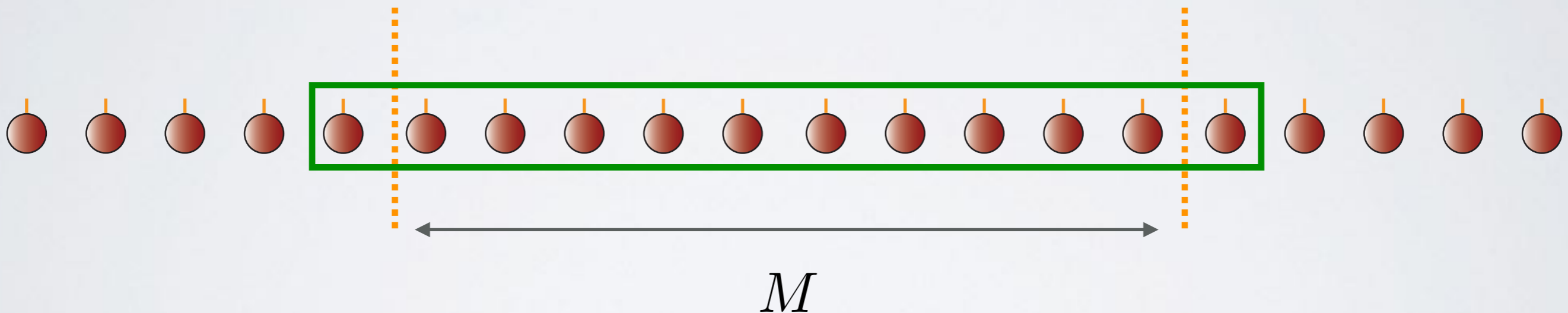


$$\lambda_M = \min_{A_M} \frac{\|[A_M, H]\|_2^2}{\|A_M\|_2^2}$$

Statistics of small commutators

minimum eigenvalue of an effective
Hamiltonian on vectorized operators

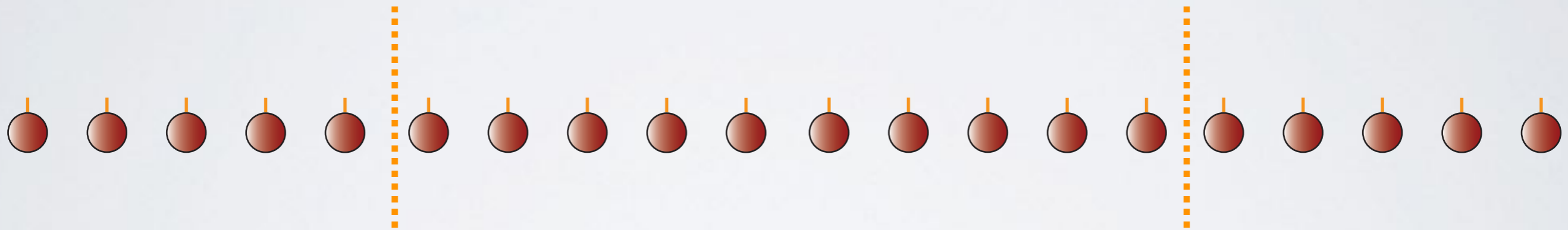
$$H_{\text{eff}} \approx \left(H \otimes \mathbb{I} - \mathbb{I} \otimes H^T \right)_{M+2}^2$$



$$\lambda_M = \min_{A_M} \frac{\| [A_M, H] \|_2^2}{\| A_M \|_2^2}$$

Statistics of small commutators

EVT rare regions affect the distribution of commutators

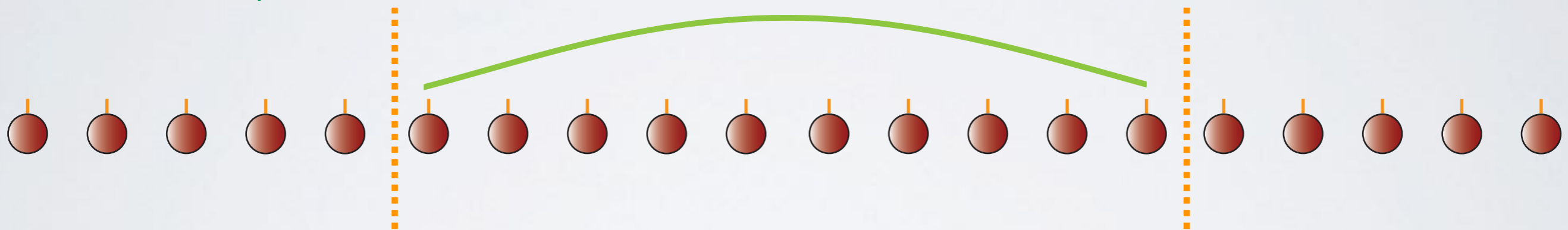


Statistics of small commutators

EVT rare regions affect the distribution of commutators

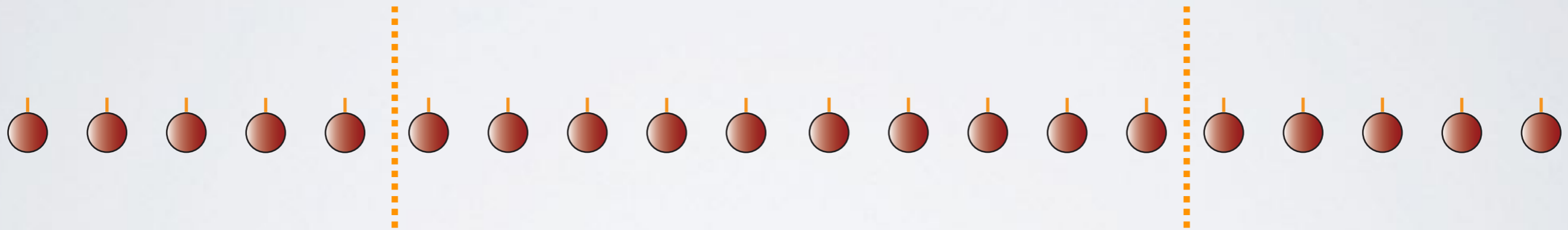
typical in thermal
phase

power law



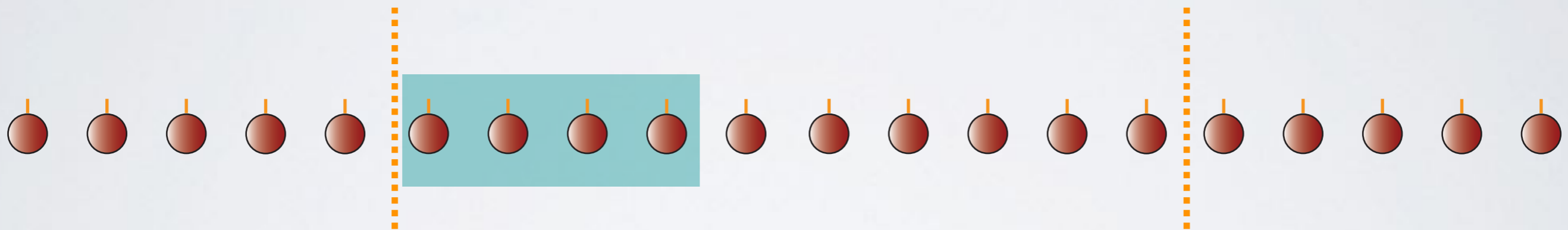
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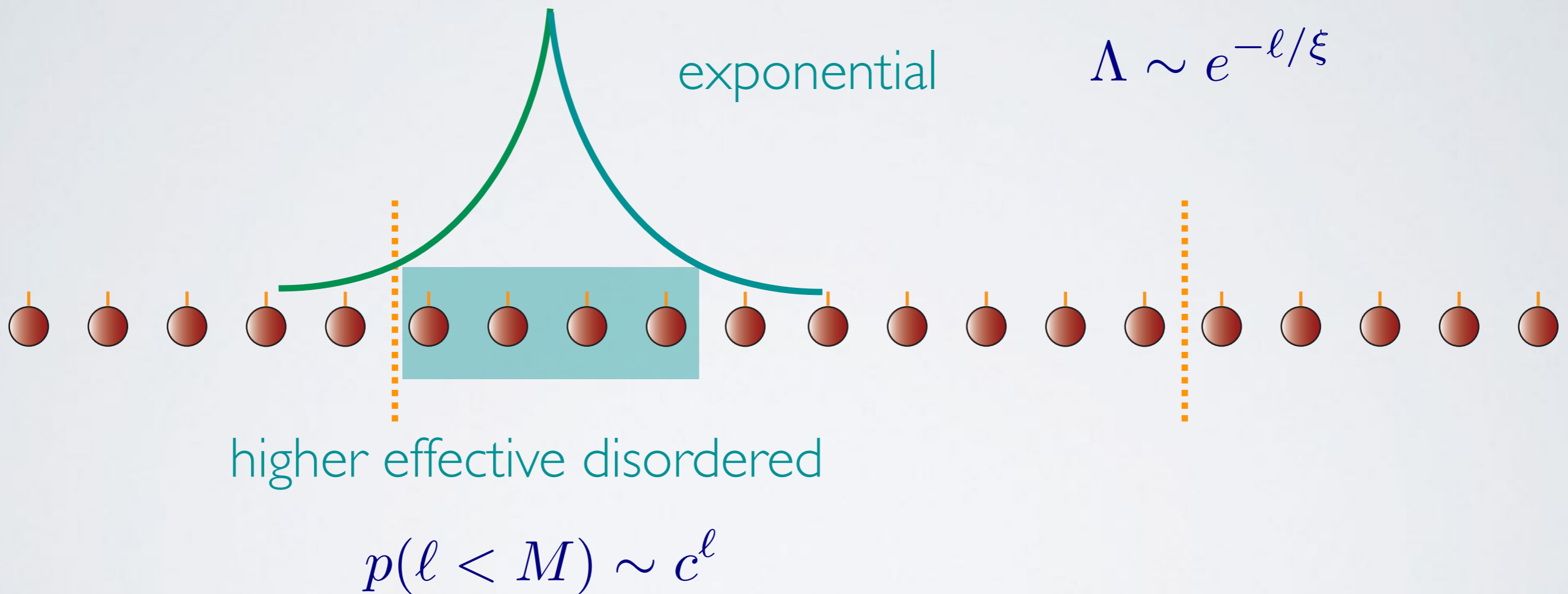


higher effective disordered

$$p(\ell < M) \sim c^\ell$$

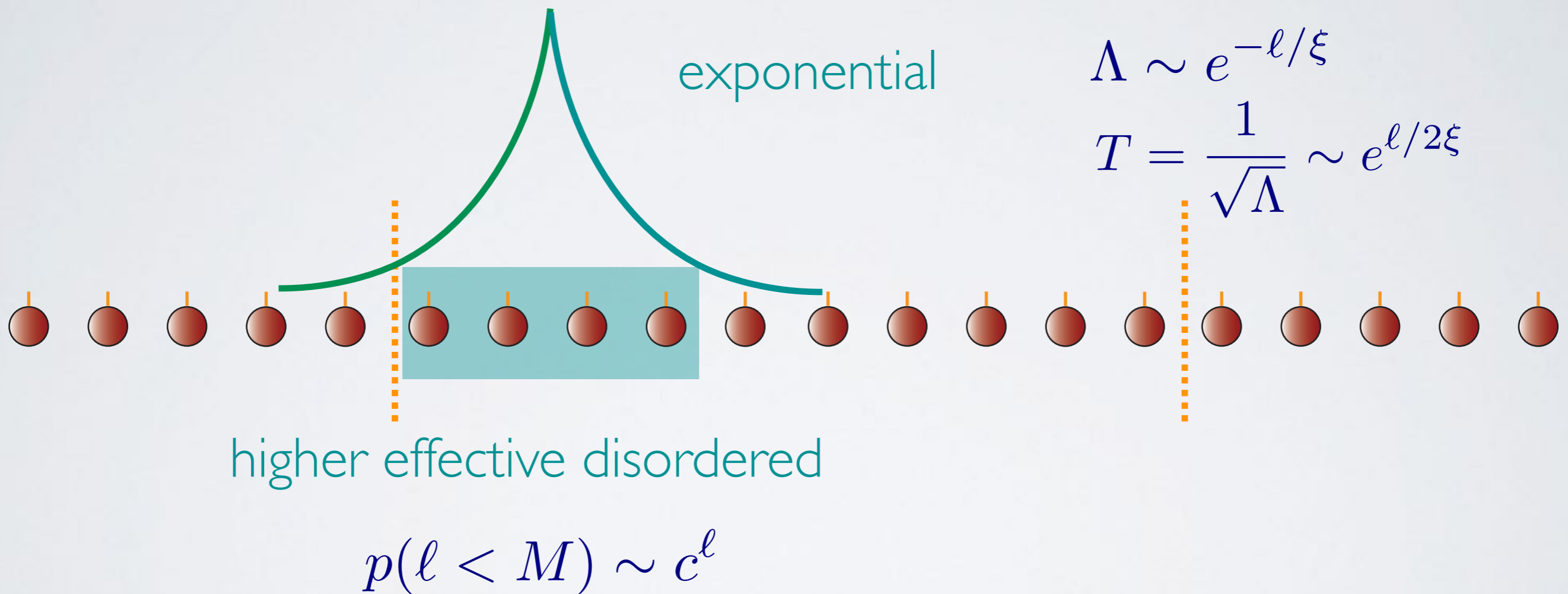
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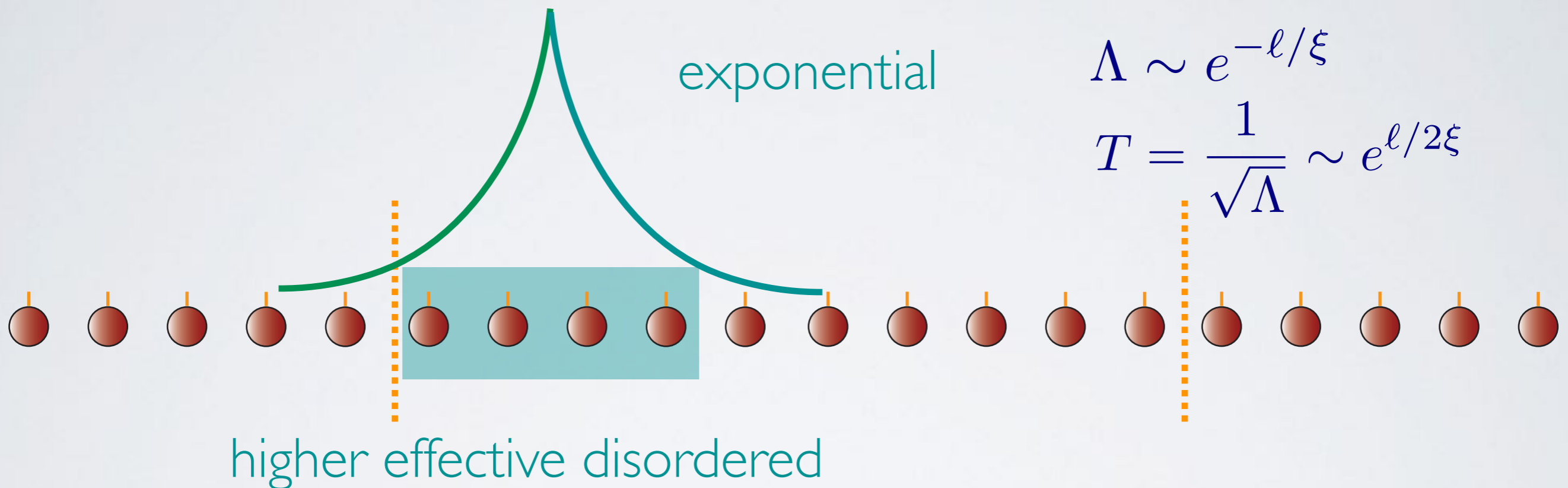
Statistics of small commutators

EVT rare regions affect the distribution of commutators



Statistics of small commutators

EVT rare regions affect the distribution of commutators



$$\Lambda \sim e^{-\ell/\xi}$$

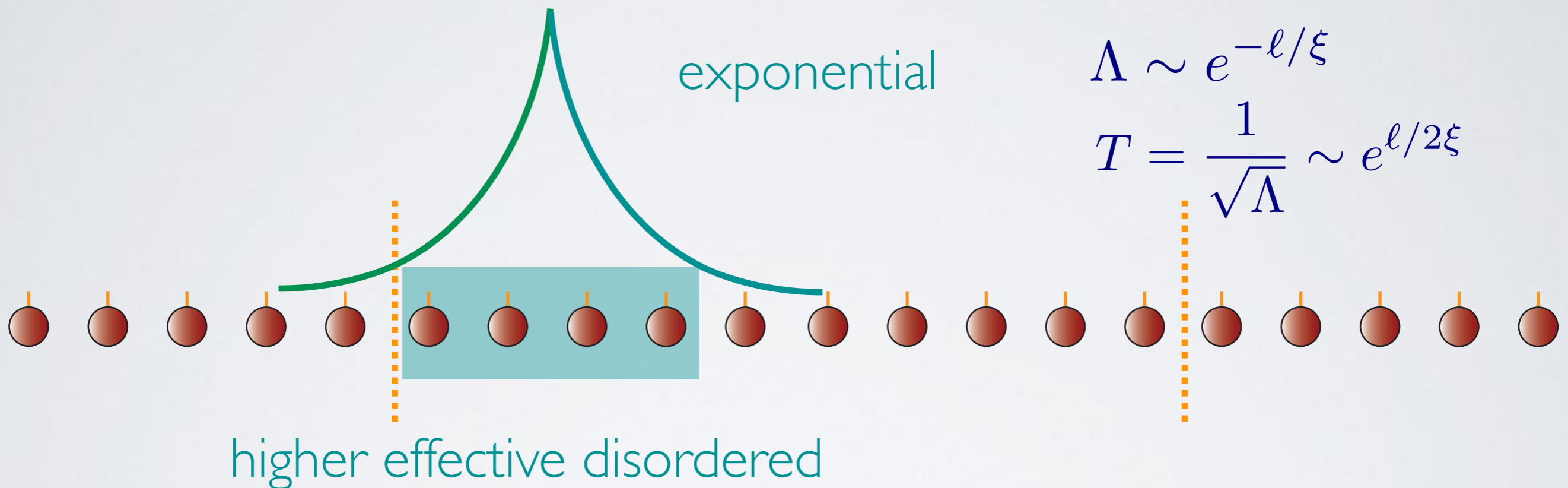
$$T = \frac{1}{\sqrt{\Lambda}} \sim e^{\ell/2\xi}$$

$$p(\ell < M) \sim c^\ell$$

$$p(T) \propto T^{-2\xi|\ln c|-1}$$

Statistics of small commutators

EVT rare regions affect the distribution of commutators



$$\Lambda \sim e^{-\ell/\xi}$$

$$T = \frac{1}{\sqrt{\Lambda}} \sim e^{\ell/2\xi}$$

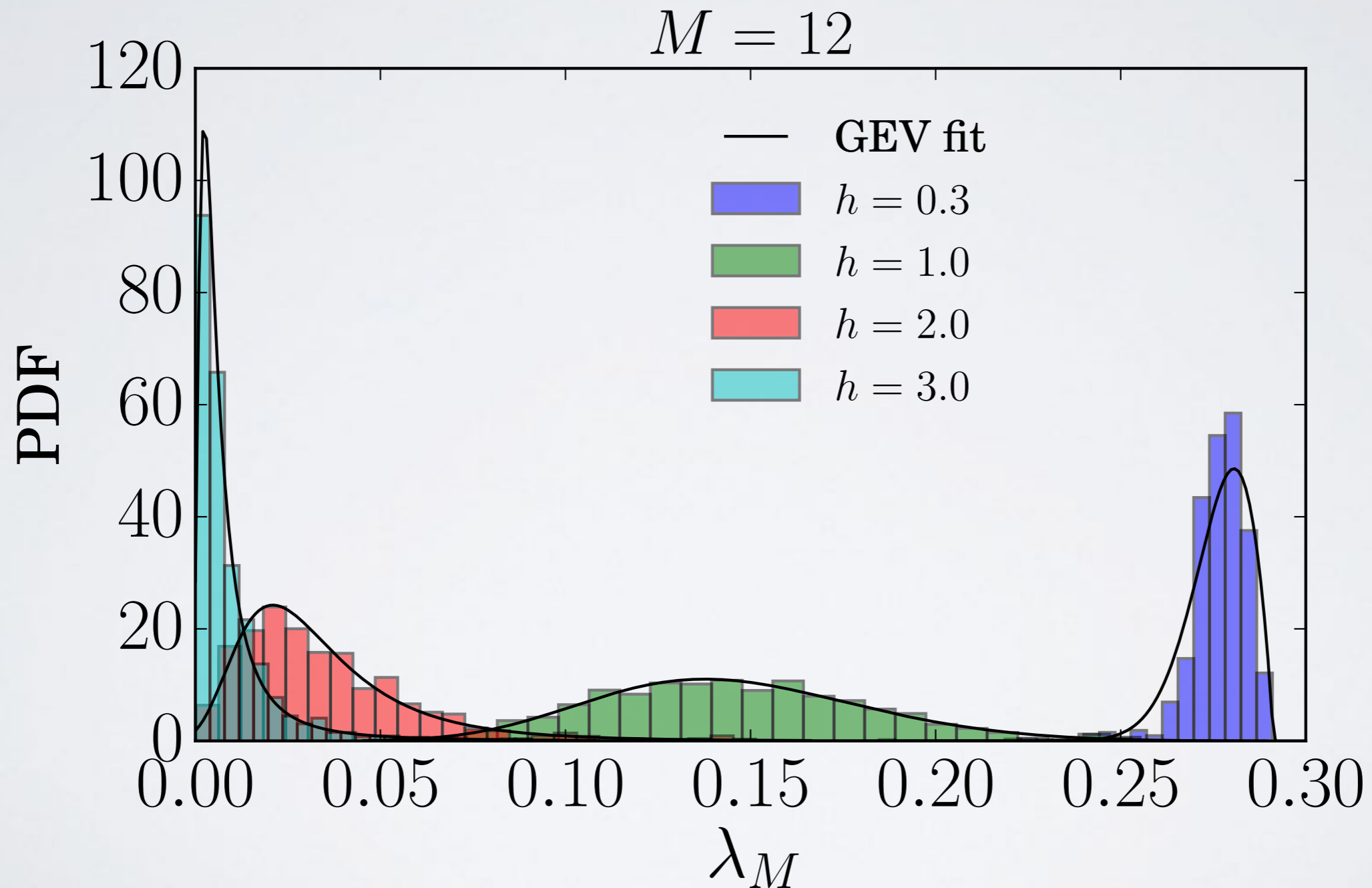
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Fréchet distribution

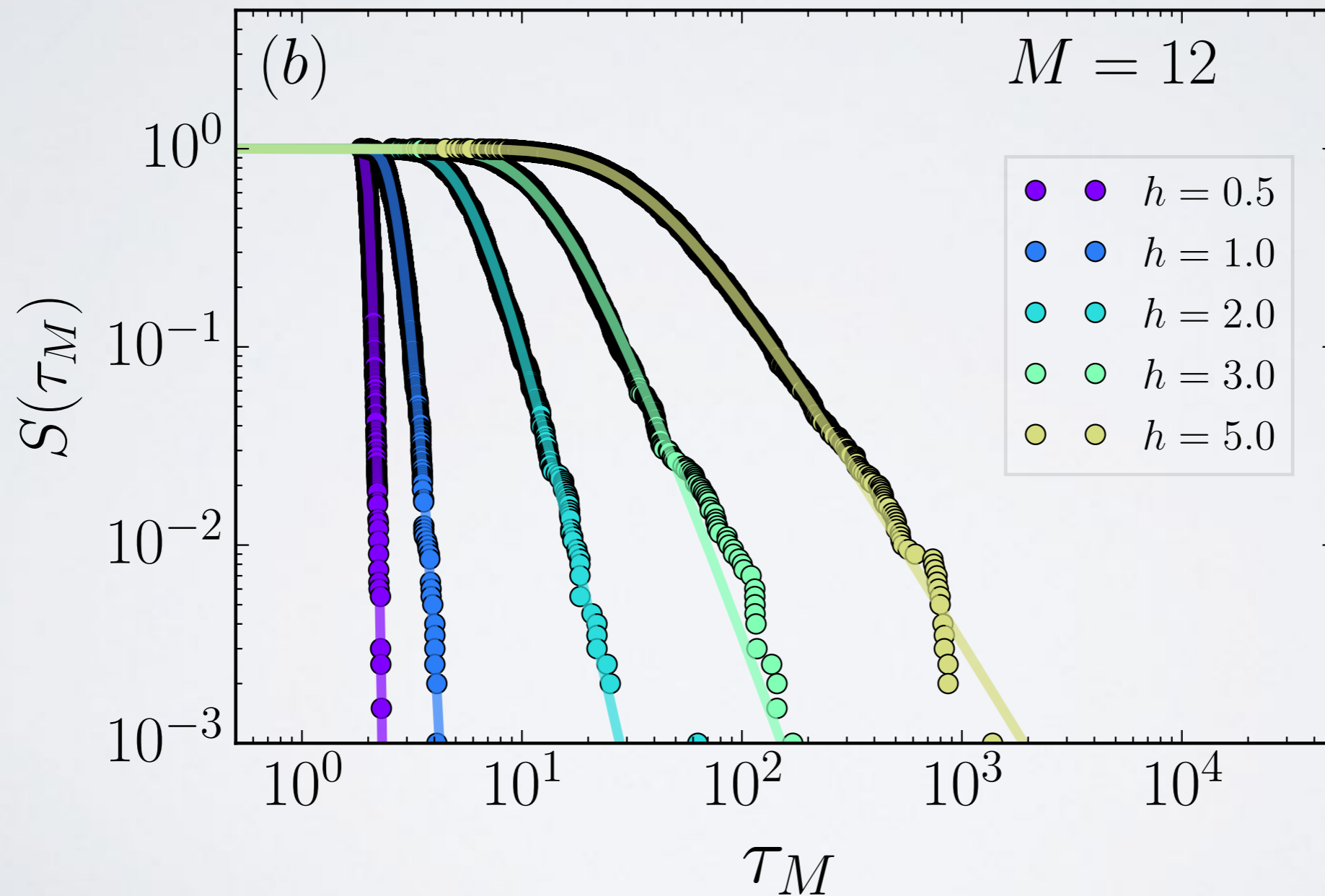
Statistics of small commutators

good fit to generalised extreme value distribution



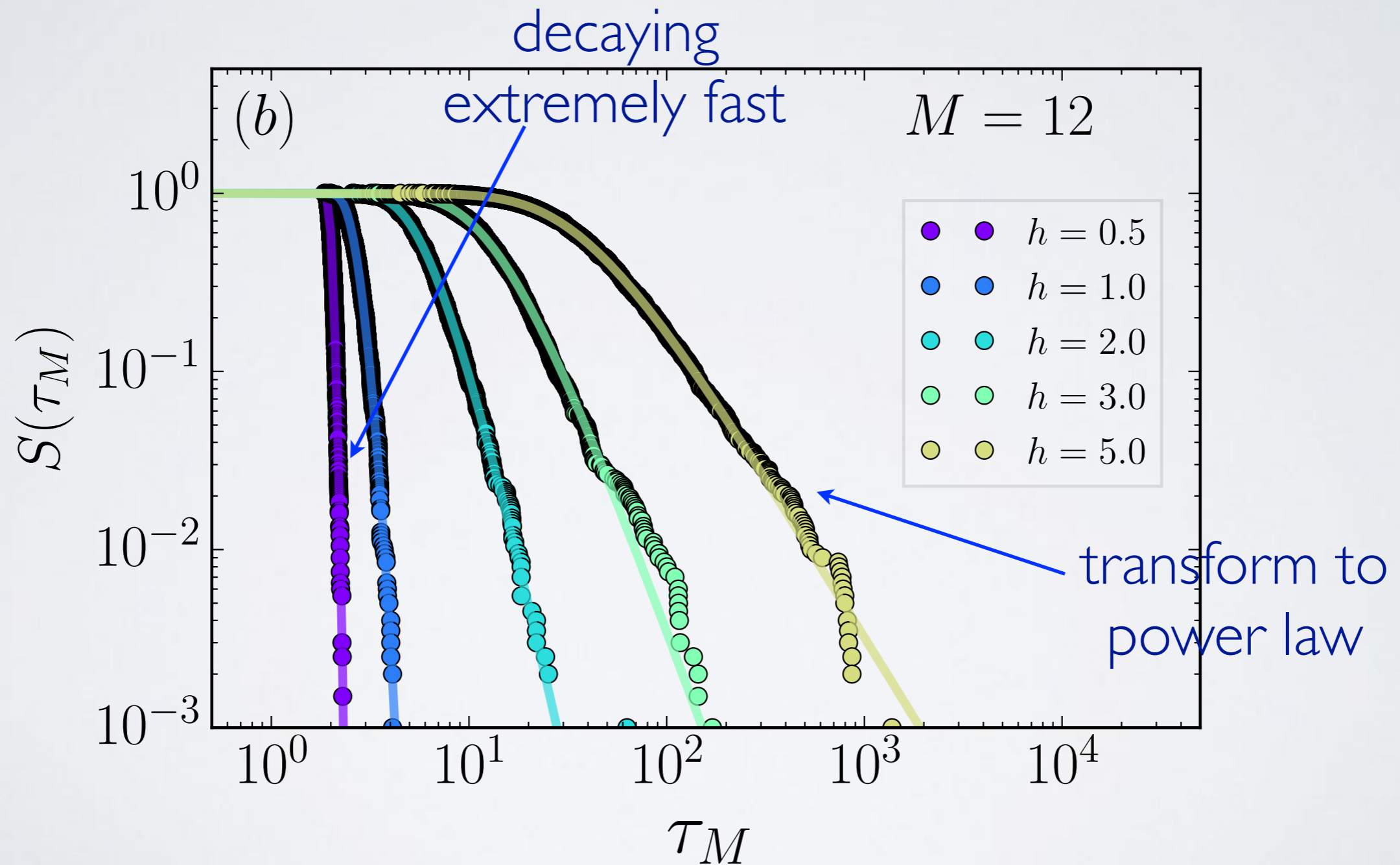
Statistics of small commutators

survival rate (probability of extremely large values)



Statistics of small commutators

survival rate (probability of extremely large values)



TO CONCLUDE

MBL effects on the dynamics of very mixed states:
quantum information perspective

MCB, N. Yao, S. Choi, M. Lukin, J.I. Cirac PRB 96, 174201 (2017)

N. Pancotti et al PRB 97, 094206 (2018)

TO CONCLUDE

MBL effects on the dynamics of very mixed states:
quantum information perspective

simulability with MPO \longrightarrow truncation errors
sensitive to localization

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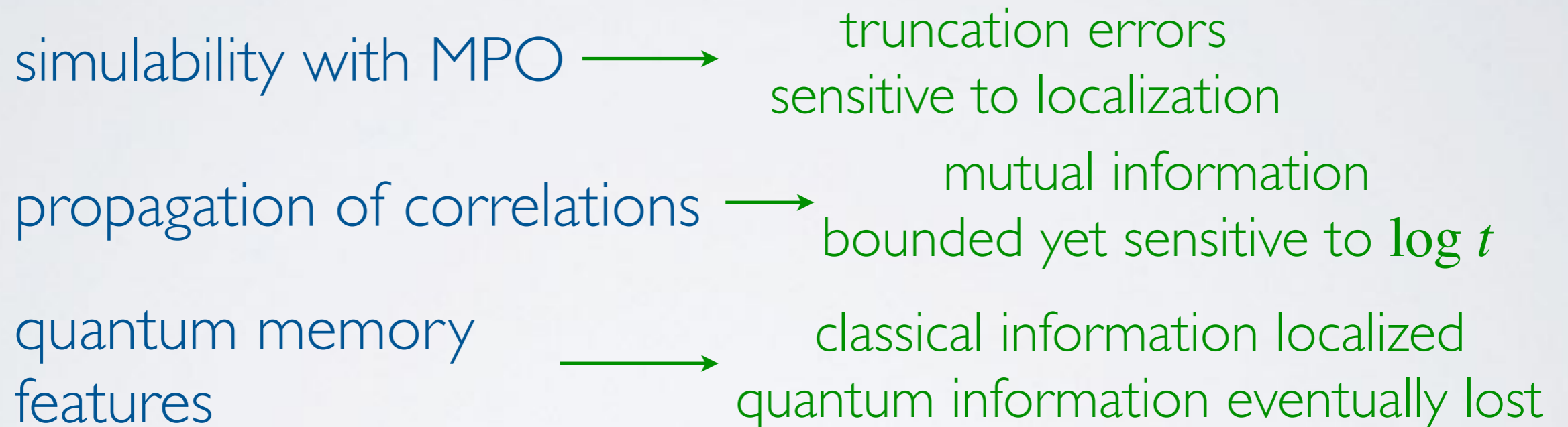
propagation of correlations \longrightarrow mutual information
bounded yet sensitive to $\log t$

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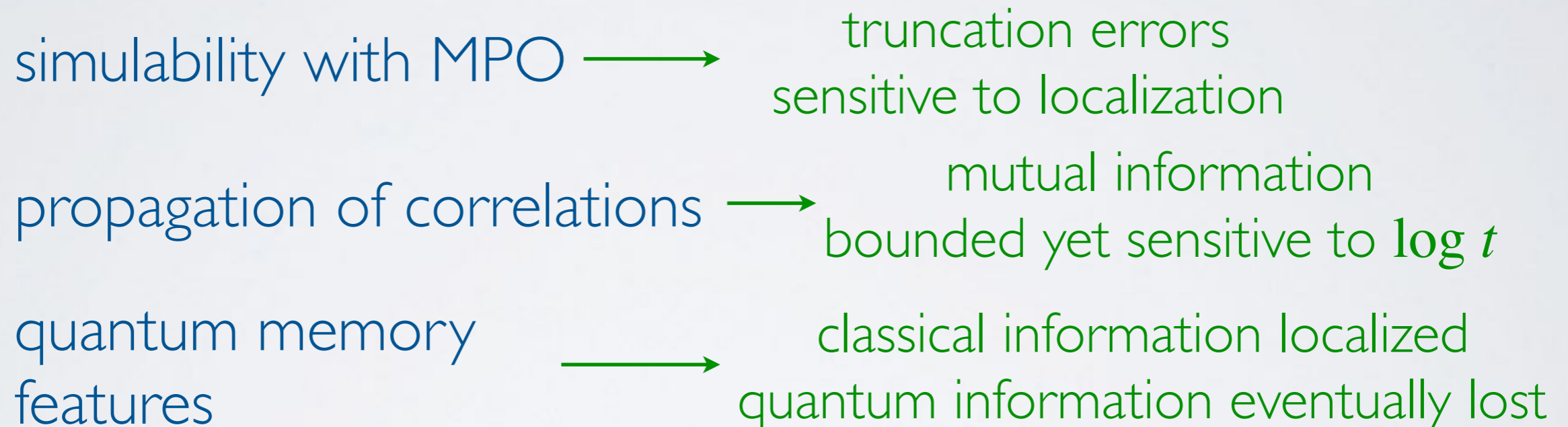


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THANKS

MBL effects on the dynamics of very mixed states:
quantum information perspective

simulability with MPO \longrightarrow truncation errors
sensitive to localization

propagation of correlations \longrightarrow mutual information
bounded yet sensitive to $\log t$

quantum memory
features \longrightarrow classical information localized
quantum information eventually lost

MCB, N. Yao, S. Choi, M. Lukin, J.I. Cirac PRB 96, 174201 (2017)

almost conserved
operators \longrightarrow signatures of localization, and
rare regions in the statistics

N. Pancotti et al PRB 97, 094206 (2018)