#### Localization and Percolation The Relationship Between Simulated Systems

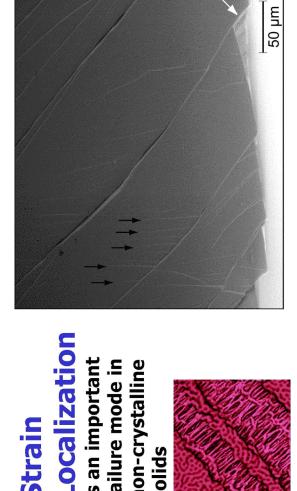
#### Yunfeng Shi Michael L. Falk

University of Michigan Materials Science and Engineering









s an important ailure mode in non-crystalline

Strain

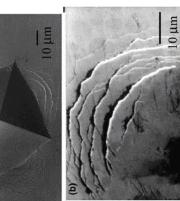
SEM Image of Shear Bands Formed in Bending Metallic Glass, Hufnagel, El-Deiry, Vinci (2000)

Craze in a block copolymer (PCHE/PE) film, Veeco NanoTheatre(2004)



#### Response Indentation for Characterizing Mechanical Metallic Glass





Mat. Sci. Eng. A (2005) A.L. Greer., A. Castellero, S.V. Madge, I Walker, J.R. Wilde

metallic alloys"

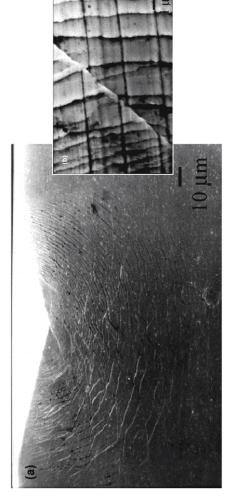
"Hardness and plastic deformation in a bulk metallic glass" Acta Materialia (2005)

U. Ramamurty, S. Jana, Y. Kawamura, K. Chattopadhyay



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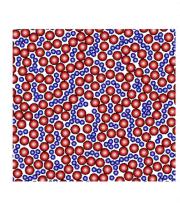
### Metallic Glass Mechanical Response Indentation for Characterizing

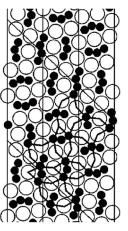


"Hardness and plastic deformation in a bulk metallic glass" Acta Materialia (2005)
U. Ramamurty, S. Jana, Y. Kawamura, K. Chattopadhyay



## **D Simulation System**







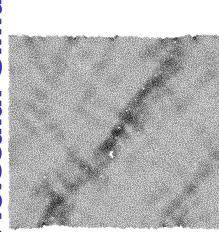


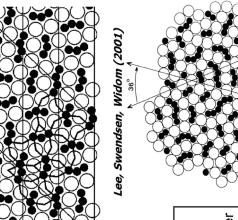
Binary system with quasi-crystalline packing (Lancon et al, Europhys. Lett, 198)

45:55 composition, 20° 000 atoms  $^{\bullet}$  T<sub>MCT</sub>  $\approx$  0.325

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Widom, Strandburg, Swendsen (19





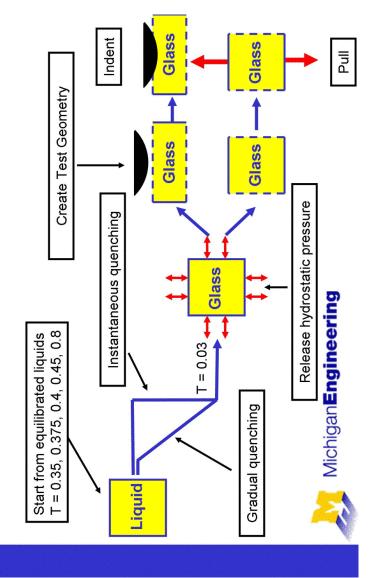
systems can be studied at a larger

Simple system that exhibits shear localization

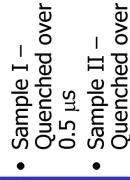
Michigan**Engineering** и

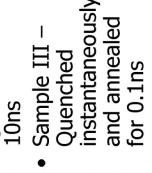
Widom, Strandburg, Swendsen (1987)

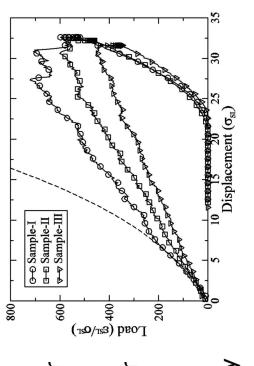
# Preparation of Glasses



## Glasses Nanoindentation of Three









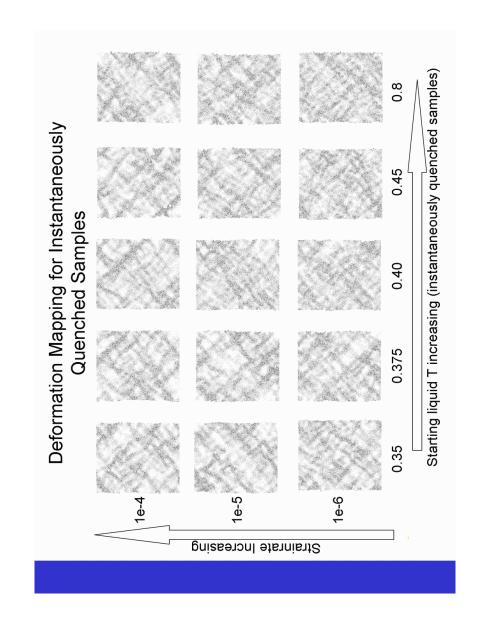
#### 25 nm Glasses 0.1ns Anneal Sample III **Shear Strain** -2.5 Nanoindentation of Three **Color Denotes Deviatoric** 10ns Quench Sample II Michigan **Engineering** 0.5µs Quench Sample I

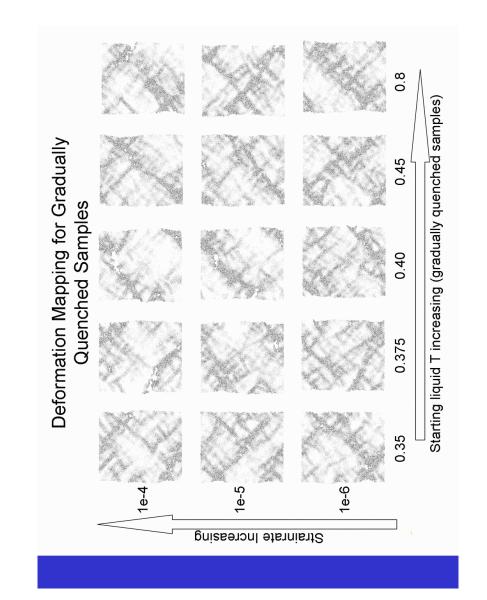
#### Glasses 11.5 Nanoindentation of Three 11.5



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**–** 25 nm





Potential Energy per Atom ( $\varepsilon_{SL}$ )

# Quantification of Shear Localization

**Participation Ratio** 

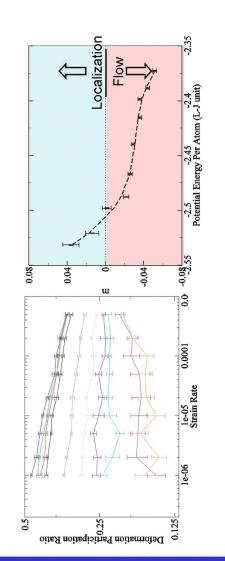
than the nominal strain Participation Ratio: Percentage of material with a local shear strain larger

favors homogenous quenched samples Low strain rate instantaneously deformation in

 $2 \times 10^{-3}$  $= 2 \times 10^{-}$ 

> favors inhomogeneous gradually quenched Low strain rate deformation in samples.

0.3 Deformation Participation Ratio Sensitivity of DPR Strain-rate





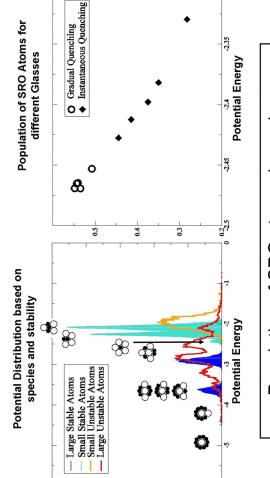
Е MichiganEngineering

m < 0: homogenous deformation → 0 and system size → For &

0: localized deformation

# MichiganEngineering

# Local Structural Analysis

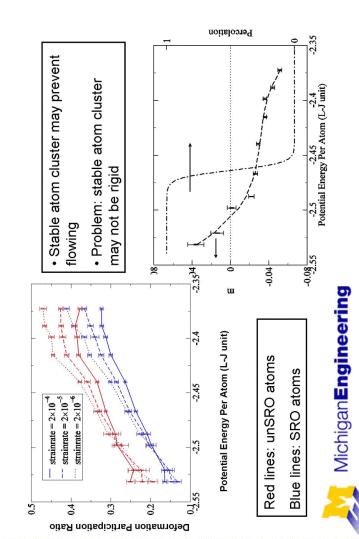


108

Population of SRO atoms depends on the thermal history of the glassy system.

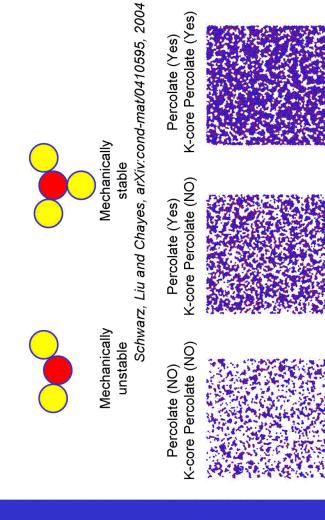
#### **Nanoindentation** Sample-II Sample-] Structural Signature of Michigan Engineering Sample-II Sample-I



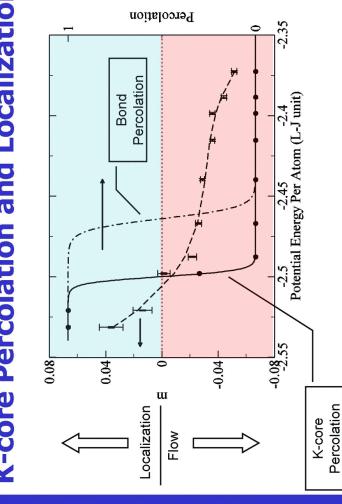


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# K-core Percolation of SRO



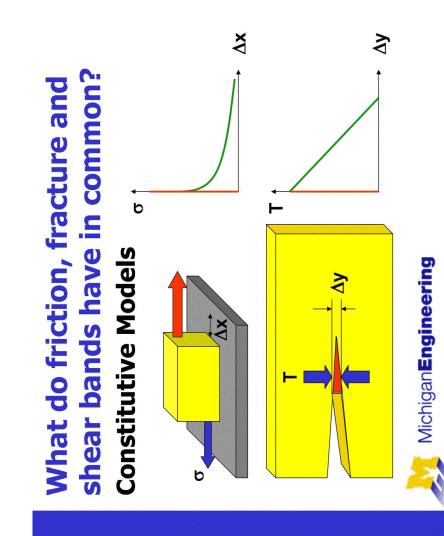
# K-core Percolation and Localization



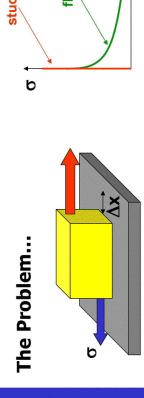
#### Summary

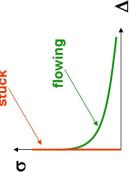
- Simulated glasses with higher degrees of SRO demonstrated a stronger tendency toward localization.
- In more rapidly quenched samples localization appears to decrease at lower strain rates
  - slowly quenched samples localization appears to increase at lower strain rates. In more
- stable backbone of material with quasi-crystal-The transition from homogeneous to localized K-core Percolation of a deformation in the quasi-static limit ike short range order. corresponds to the





## **Constitutive Models**



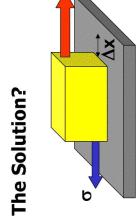


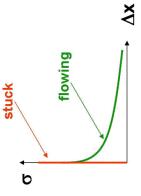
- Most of these systems involve a "stuck" or "jammed" state and a "flowing" or "slipping" state.
- Simple constitutive laws often include discontinuities.
- properties that can sensitively depend on the dynamic This complicates the analysis of instabilities and other transition between jammed and flowing states.



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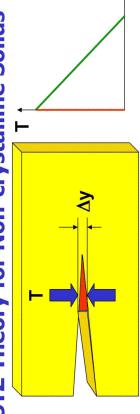
### **Formulations** Rate and State

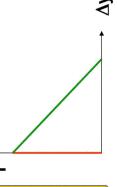




- freedom not represented in, e.g. Coulomb Friction Law These transitions arise due to internal degrees of
- essential physics of these hidden degrees of freedom. Rate and State formulations attempt to capture the
  - Examples: Dieterich (1978), Rice and Ruina (1983).
- What do these new degrees of (e.g. contact area in Dieterich) Michigan Engineering But the issue arises: freedom represent?

#### Theory for Non-crystalline Solids Rate and State Plastici

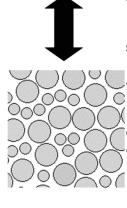




- A number of such models exist in the context of plastic deformation primarily aimed at modeling crystals.
- Crystal deformation is difficult to relate to microscopics because of the nature of the defect (D-2 object)
- some simplifying assumptions, e.g. that the defect is STZ theory attempts this for non-crystals by making D=0 and has zero mobility.

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#### Spaepen) Zones **Shear Transformation**







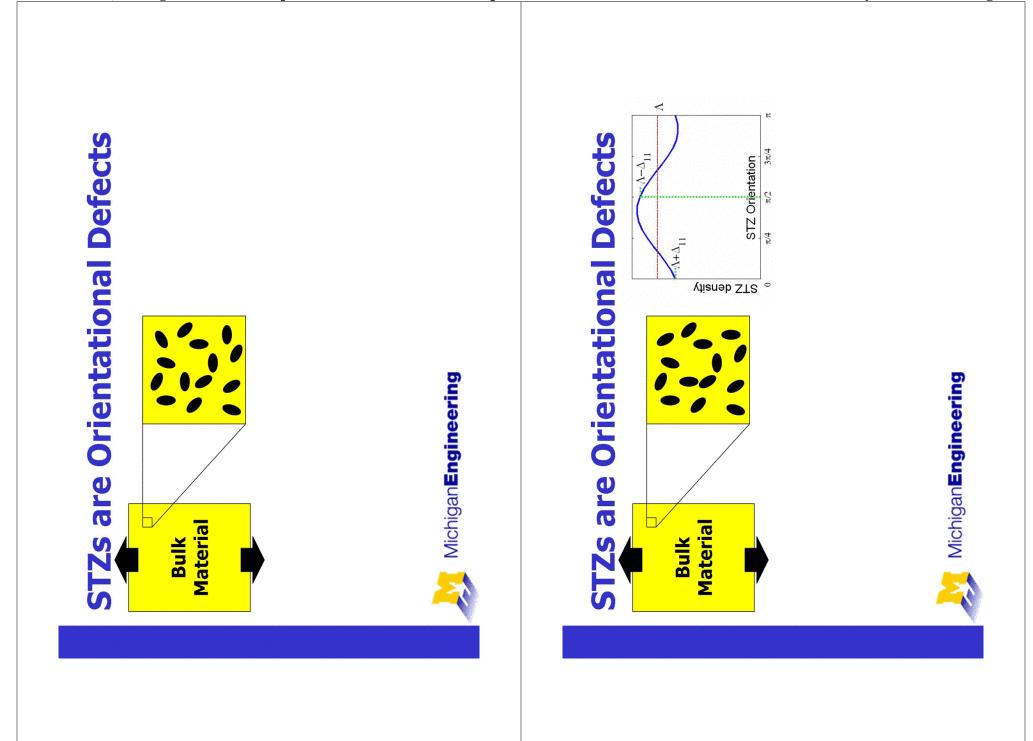


- Region may reverse its rearrangement if stress is reversed shortly thereafter
- Stress in opposite direction produces deformation amount of additional rearrangement

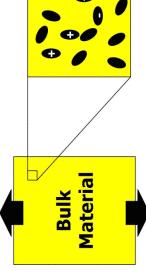
Deformation becomes permanent after some

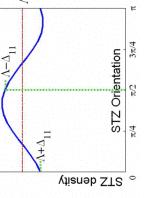
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in different regions



## **Orientational Defects** are





Simplifying to shear only along x and y principal axes:

$$n_{\scriptscriptstyle +} = \frac{n_{\scriptscriptstyle \infty}}{2} \left( \Lambda + \Delta_{11} \right)$$

$$oldsymbol{n}_{+}=rac{n}{2}\left(\Lambda+\Delta_{11}
ight)$$
 $oldsymbol{n}_{-}=rac{n_{\infty}}{2}\Big(\Lambda+\Delta_{22}\Big)=rac{n_{\infty}}{2}\Big(\Lambda-\Delta_{11}\Big)$ 

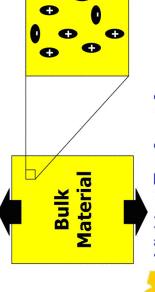


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### The STZ Model

 $\epsilon_0/n_\infty$ = volume per STZ Plastic Strain Rate Proportional to Flips  $\left( egin{aligned} \left( egin{aligned} s \end{array} 
ight) n_{_{-}} - R_{_{+}} \left( s 
ight) n_{_{+}} \end{aligned} \end{aligned} 
ight]$ R ි ද П





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Plastic Strain Rate Proportional to Flips

Master Equation for Densities 
$$\dot{n}_{\pm}=R_{\mp}n_{\mp}-R_{\pm}n_{\pm}+\left[\Gamma\left(s,n_{\pm}\right)+
ho(T)\right]\left[\frac{1}{2}n_{\infty}-n_{\pm}n_{\pm}\right]$$



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### The STZ Model

Plastic Strain Rate Proportional to Flips

$$\dot{\mathcal{E}}^{pl} = \frac{\mathcal{E}_0}{n_{\infty}} \left[ R_{-}(s) n_{-} - R_{+}(s) n_{+} \right]$$

 $\varepsilon_0/n_\infty$ = volume per STZ

Master Equation for Densities

$$\dot{n}_{\scriptscriptstyle \pm} = R_{\scriptscriptstyle \mp} n_{\scriptscriptstyle \mp} - R_{\scriptscriptstyle \pm} n_{\scriptscriptstyle \pm} + \left[ \overline{\Gammaig(s,n_{\scriptscriptstyle \pm}ig)} + oldsymbol{
ho}(T) 
ight] \left[ rac{1}{2} \, n_{\scriptscriptstyle \infty} - n_{\scriptscriptstyle \pm} 
ight]$$

For the state of the system to depend only upon the strain history in the low T, low rate limit, the same rate function must control both annihilation and creation



### The STZ Model

Plastic Strain Rate Proportional to Flips

$$\dot{\mathcal{E}}^{pl} = rac{\mathcal{E}_0}{n_{_{\! \infty}}} igg[ R_{_{\! -}}ig(sig) n_{_{\! -}} - R_{_{\! +}}ig(sig) n_{_{\! +}} igg] \quad rac{arepsilon_o n_{_{\! \infty}}}{n_{_{\! \infty}}} = ext{volume per STZ}$$

Master Equation for Densities

$$\dot{n}_{_{\pm}} = R_{_{\mp}}n_{_{\mp}} - R_{_{\pm}}n_{_{\pm}} + \left[ \left[ \Gamma\left(s,n_{_{\pm}}
ight) + 
ho(T) 
ight] \left[ rac{1}{2}n_{_{\infty}} - n_{_{\pm}} 
ight]$$

This provides a closure for  $\Gamma$  assuming that the First Law of Thermodynamics is obeyed and that  $\Gamma$  is proportional to the rate of dissipation.

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L. Pechenik (2003)

The STZ Model

Plastic Strain Rate Proportional to Flips

$$\dot{\mathcal{E}}^{pl} = rac{\mathcal{E}_0}{n_{_{\! \infty}}} igg[ R_{_{\! -}}ig(sig) n_{_{\! -}} - R_{_{\! +}}ig(sig) n_{_{\! +}} igg] \quad _{\! \epsilon_0 n_{_{\! \infty}}} = ext{volume per STZ}$$

Master Equation for Densities

$$\dot{n}_{\scriptscriptstyle \pm} = R_{\scriptscriptstyle \mp} n_{\scriptscriptstyle \mp} - R_{\scriptscriptstyle \pm} n_{\scriptscriptstyle \pm} + \left[\Gamma\left(s, n_{\scriptscriptstyle \pm}
ight) + 
ho\left(T
ight)
ight]\left[rac{1}{2}n_{\scriptscriptstyle \infty}
ight] - n_{\scriptscriptstyle \pm}$$

This is identified with the "Granular Temperature"  $\chi$ The  $n_{\infty}$  parameter is the ratio of annihilation to creation and sets the equilibrium defect density.



Langer (2004)

### The STZ Model

Plastic Strain Rate Proportional to Flips

Master Equation for Densities

$$\dot{n}_{\scriptscriptstyle \pm} = R_{\scriptscriptstyle \mp} n_{\scriptscriptstyle \mp} - R_{\scriptscriptstyle \pm} n_{\scriptscriptstyle \pm} + \left[ \Gamma \left( s, n_{\scriptscriptstyle \pm} 
ight) + 
ho \left( T 
ight) 
ight] \left[ rac{1}{2} e^{-1/art arkappa} - n_{\scriptscriptstyle \pm} 
ight] .$$

The  $n_{\infty}$  parameter is the ratio of annihilation to creation and sets the equilibrium defect density.

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Langer (2004)



TWICH INSULE INSTITUTE OF THE

### The STZ Model

Plastic Strain Rate Proportional to Flips

$$\dot{\mathcal{E}}^{pl} = rac{\mathcal{E}_0}{n_{_{\! \infty}}} \Big[ R_{_{\! -}}ig(sig) n_{_{\! -}} - R_{_{\! +}}ig(sig) n_{_{\! +}} \Big] \quad _{\! \epsilon_0 n_{_{\! \infty}}} = ext{volume per STZ}$$

Master Equation for Densities

$$\dot{n}_{\scriptscriptstyle \pm} = R_{\scriptscriptstyle \mp} n_{\scriptscriptstyle \mp} - R_{\scriptscriptstyle \pm} n_{\scriptscriptstyle \pm} + \left[\Gamma\left(s, n_{\scriptscriptstyle \pm}\right) + \rho\left(T\right)\right]\left[rac{1}{2}e^{-1/\chi} - n_{\scriptscriptstyle \pm}\right]$$

 $\chi$  has its own dynamics. Shear drives it to a high value  $(\chi_{\infty})$  while thermal fluctuations drive it to a low value  $(\chi_T = kT/E_{STZ})$ 



Langer (2004)

### The STZ Model

Plastic Strain Rate Proportional to Flips

$$\dot{\mathcal{E}}^{pl} = rac{\mathcal{E}_0}{n_\omega} igg[ R_{_-}ig( s ig) n_{_-} - R_{_+}ig( s ig) n_{_+} igg] \quad rac{\epsilon_0 n_\omega}{\epsilon_0 n_\omega} = ext{volume per STZ}$$

Master Equation for Densities

Master Equation for Densities 
$$\dot{n}_{\pm}=R_{\mp}n_{\mp}-R_{\pm}n_{\pm}+\Big[\Gammaig(s,n_{\pm}ig)+ig
hoig(Tig)\Big]\Big[rac{1}{2}e^{-1/\chi}-n_{\pm}\Big]$$

$$rac{ au_0 c_0}{arepsilon_0} \, \dot{\chi} = e^{-1/arkpi} \Gamma ig( \chi_\infty - \chi ig) + \kappa oldsymbol{
ho}ig( T ig) e^{-eta/arkpi} ig( \chi_T - \chi ig)$$

 $+\ell^2\nabla^2\chi$ 

by the diffusion of the granular temperature.

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Langer (2004)

# **Dynamic Transition from**

(Quasilinear Formulation) Hardening to Flow



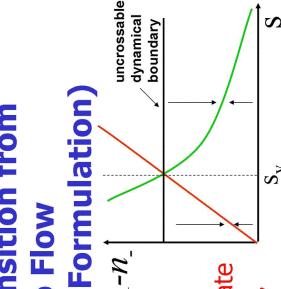
Jammed steady state  $n_{+}-n_{-}=(n_{+}+n_{-}) s/s_{y}$ 



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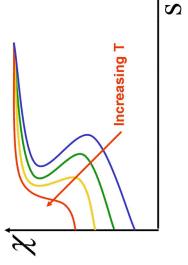
 $n_{+}-n_{-}=S_{V}/S, n_{+}+n_{-}=n_{\infty}$ 

Flowing steady state



#### (Quasilinear Formulation) **Dynamic Transition from** FIOW Hardening to

At low T, granular temperature,  $\chi$ , is double valued as a function of stress implying the possibility of 2 phase coexistence between jamming (creep) and flow.





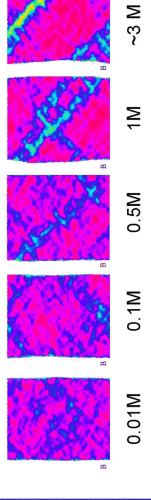
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# **Current Investigations**

- Is this a good model for the observed localization?
- Is this a unique or optimal model?
- Does \(\chi\) map directly onto %SRO?
- Can this model help quantify the role of internal degrees of freedom in the softening process?
- Does this have any bearing on localization in other (e.g., granular) materials?



#### How does this vary with dimensionality?



Repeated in 3D binary LJ glass using 50:50

Wahnstrom potential

Strain rate  $1 \times 10^{-5} t_0^{-1}$ 

Shown, Quench time: 3.33  $T_{MCT}$  to 0.1  $T_{MCT}$ .





