



Alexander von Humboldt

Stiftung / Foundation



# Accurate van der Waals interactions from ground-state electron density

Alexandre Tkatchenko

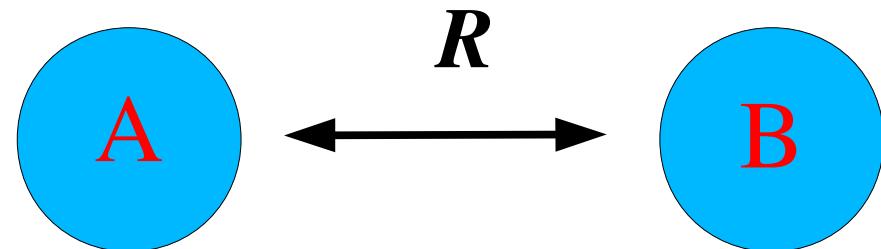
Theory Department, Fritz-Haber-Institut der MPG

Berlin, Germany

[tkatchen@fhi-berlin.mpg.de](mailto:tkatchen@fhi-berlin.mpg.de)

EXCITCM09, Nov 2, 2009

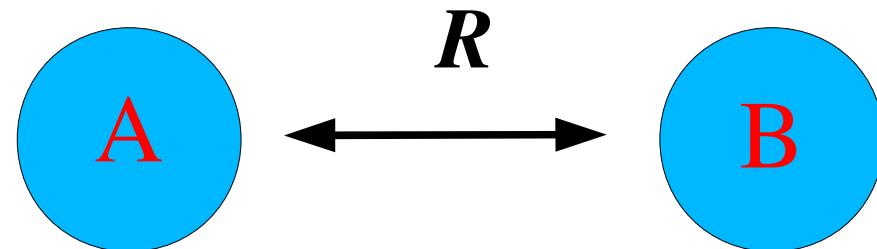
# Back to Basics



No orbital overlap (large enough ***R***):

$$E_{\text{A-B}}(R) = - \left( \frac{C_6}{R^6} + \frac{C_8}{R^8} + \frac{C_{10}}{R^{10}} + \dots \right)$$

# Back to Basics



No orbital overlap (large enough  $R$ ):

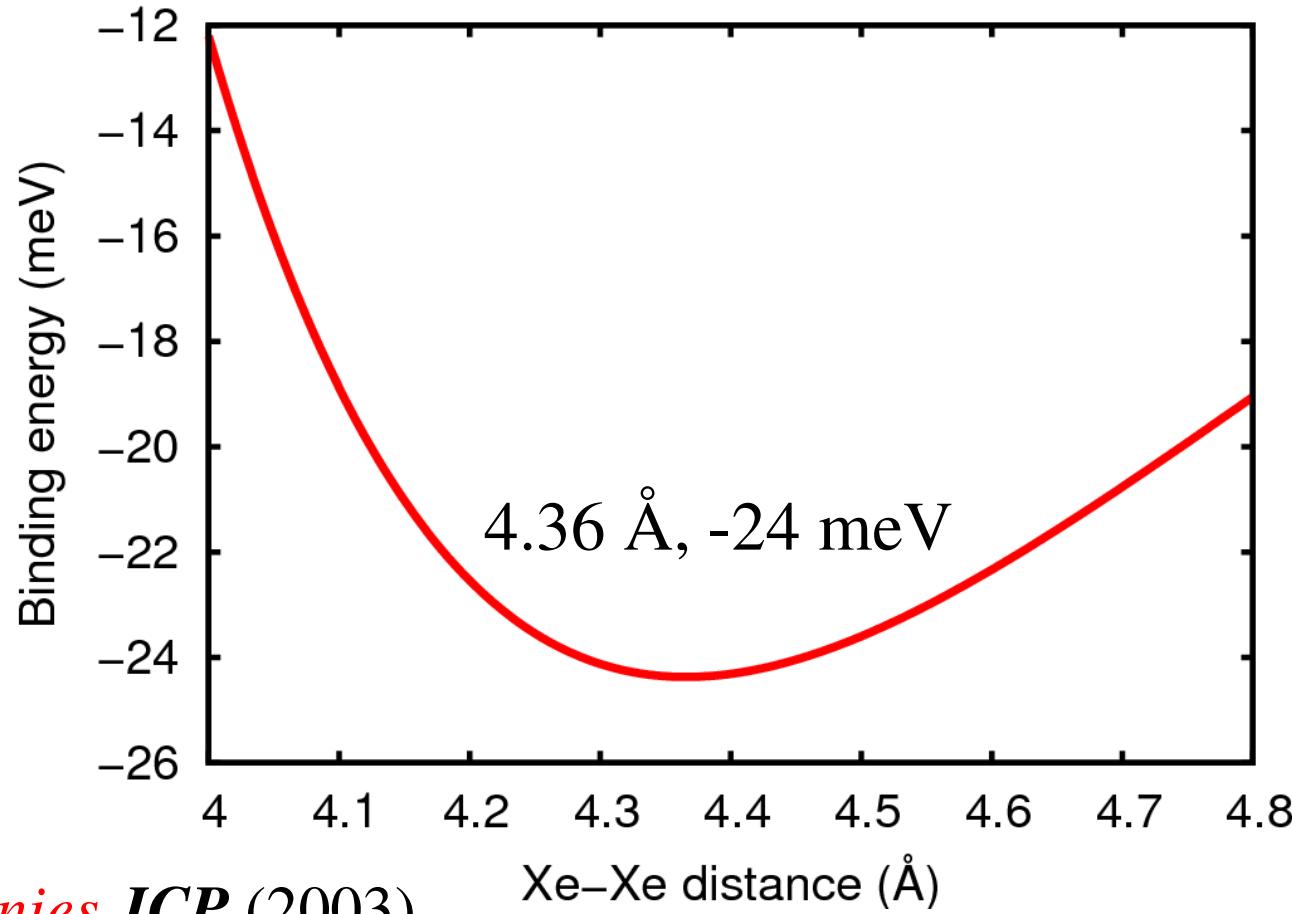
$$E_{A-B}(R) = - \left( \frac{C_6}{R^6} + \frac{C_8}{R^8} + \frac{C_{10}}{R^{10}} + \dots \right)$$

20% error in  $C_n$  coefficients leads to 20% error in binding energy

# Role of vdW asymptotics at equilibrium distance

$$E_{\text{Xe-Xe}}(R) = E_{\text{Hartree-Fock}}(R) - \left( f_6(R) \frac{C_6}{R^6} + f_8(R) \frac{C_8}{R^8} + f_{10}(R) \frac{C_{10}}{R^{10}} + \dots \right)$$

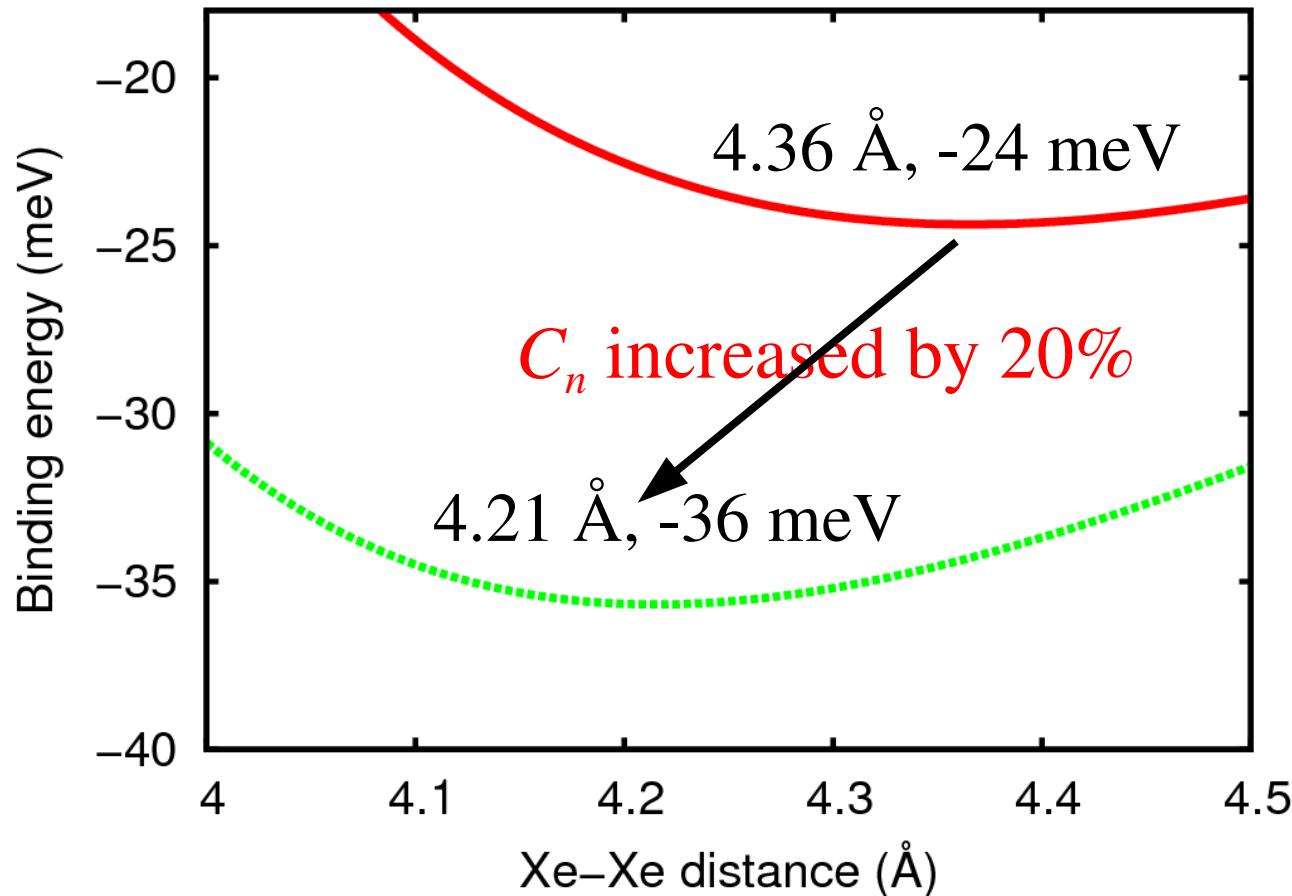
Xe-Xe



# Role of vdW asymptotics at equilibrium distance

$$E_{\text{Xe-Xe}}(R) = E_{\text{Hartree-Fock}}(R) - \left( f_6(R) \frac{C_6}{R^6} + f_8(R) \frac{C_8}{R^8} + f_{10}(R) \frac{C_{10}}{R^{10}} + \dots \right)$$

Xe-Xe

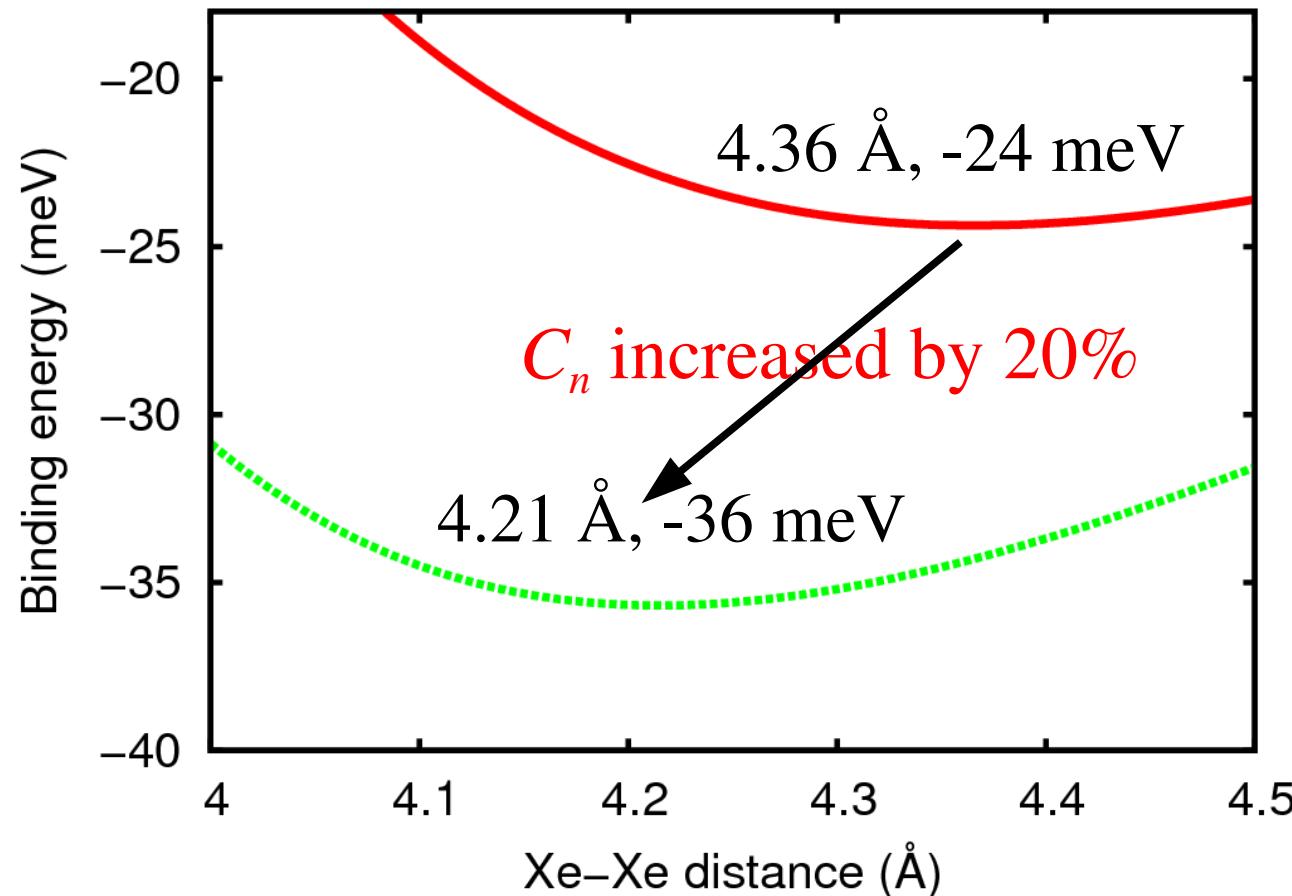


# Asymptotics do matter: RG dimers

$\sim 60\%$      $\sim 30\%$      $\sim 10\%$

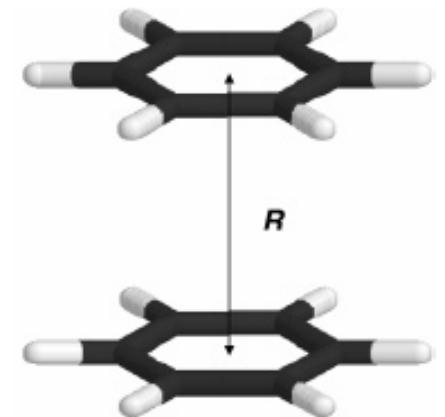
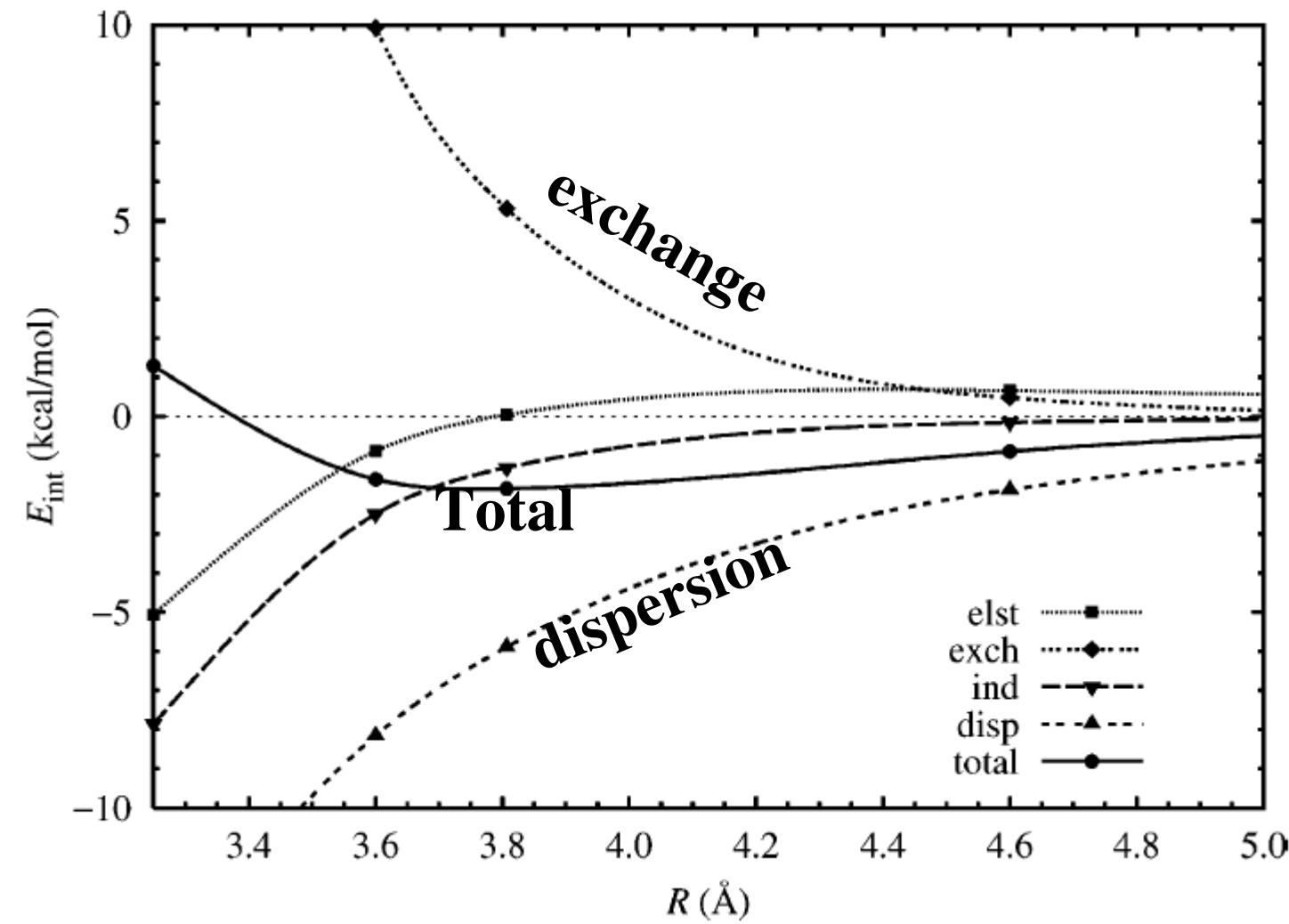
$$E_{\text{Xe-Xe}}(R) = E_{\text{Hartree-Fock}}(R) - \left( f_6(R) \frac{C_6}{R^6} + f_8(R) \frac{C_8}{R^8} + f_{10}(R) \frac{C_{10}}{R^{10}} + \dots \right)$$

Xe-Xe



Increasing  $C_n$  by **20%** increases the binding energy by **50%**

# Asymptotics do matter: Benzene dimer



SAPT:  
*Podeszwa, Bukowski,  
Szalewicz, JPCA (2006)*

Error of **20%** in the dispersion energy leads to **50%** error in binding

Can we get dispersion coefficients  
(from ground-state theory)  
to better than 5% ?

# Accurate reference $C_6$ database<sup>(1)</sup>

$$S(k) = \int_{E_0}^{\infty} E^k \left( \frac{df}{dE} \right) dE$$

Differential dipole oscillator strength  
(DOS)

Transition energy (ground state--excited state)

<sup>(1)</sup> W. J. Meath and co-workers (1977-present)

# Accurate reference $C_6$ database<sup>(1)</sup>

$$S(k) = \int_{E_0}^{\infty} E^k \left( \frac{df}{dE} \right) dE$$

Experimental data (DOSD)

$$S(k) = \sum_{i=1}^n \epsilon_i^k f_i, \quad k = 0, -1, \dots, -2n+1$$

Theory (pseudo-DOSD)

$$C_{6AB} = \frac{3}{2} \sum_{i,j} \frac{f_i^A f_j^B}{\epsilon_i^A \epsilon_j^B (\epsilon_i^A + \epsilon_j^B)}$$

Using different sets of exp.  
 $S(k)$  data,  $C_{6AB}$  is typically  
accurate to **1-2%**

<sup>(1)</sup> W. J. Meath and co-workers (1977-present)

# Accurate reference $C_6$ database<sup>(1)</sup>

$$S(k) = \int_{E_0}^{\infty} E^k \left( \frac{df}{dE} \right) dE$$

Experimental data (DOSD)

$$S(k) = \sum_{i=1}^n \epsilon_i^k f_i, \quad k = 0, -1, \dots, -2n+1$$

Theory (pseudo-DOSD)

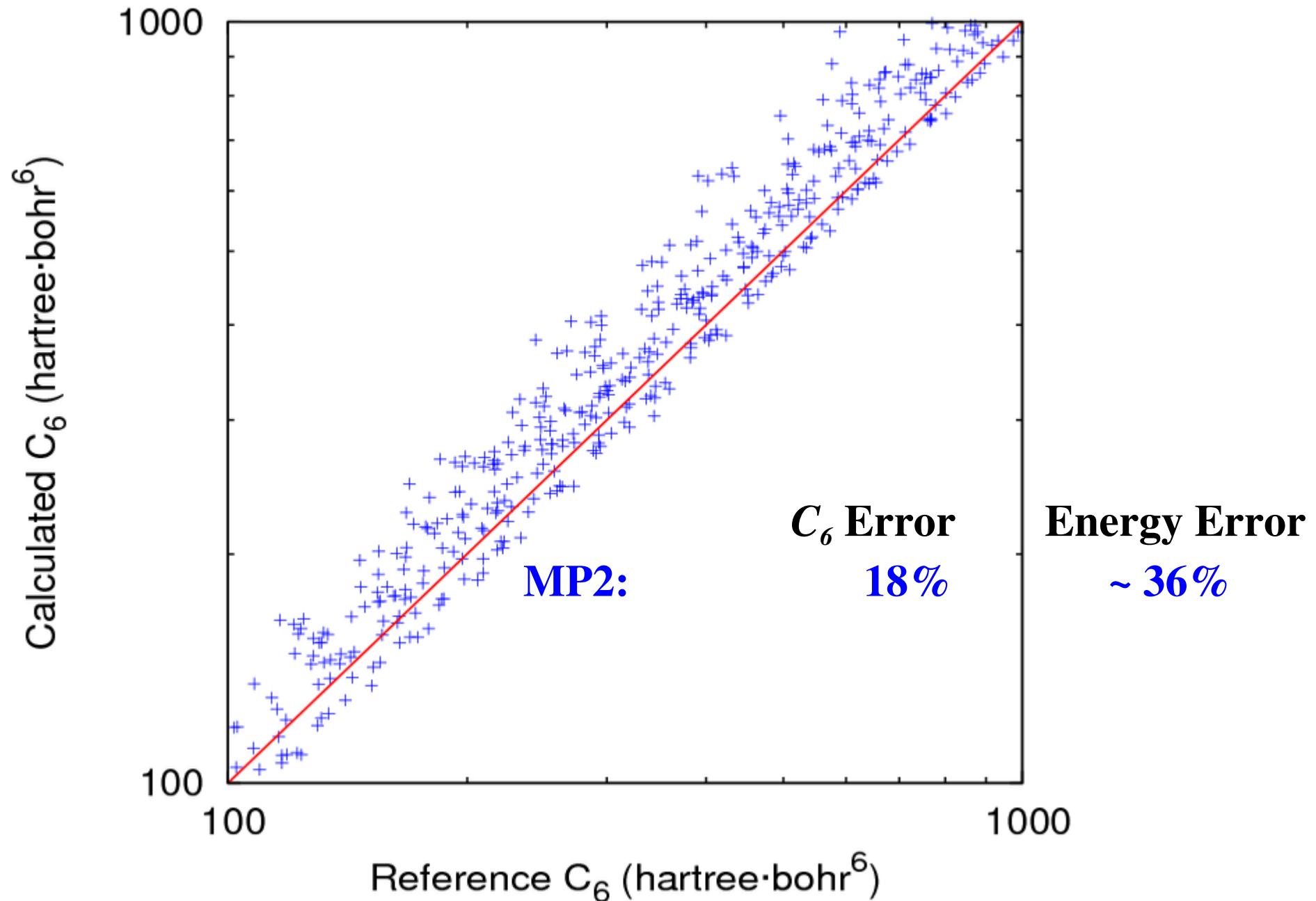
$$C_{6AB} = \frac{3}{2} \sum_{i,j} \frac{f_i^A f_j^B}{\epsilon_i^A \epsilon_j^B (\epsilon_i^A + \epsilon_j^B)}$$

Using different sets of exp.  
 $S(k)$  data,  $C_{6AB}$  is typically  
accurate to **1-2%**

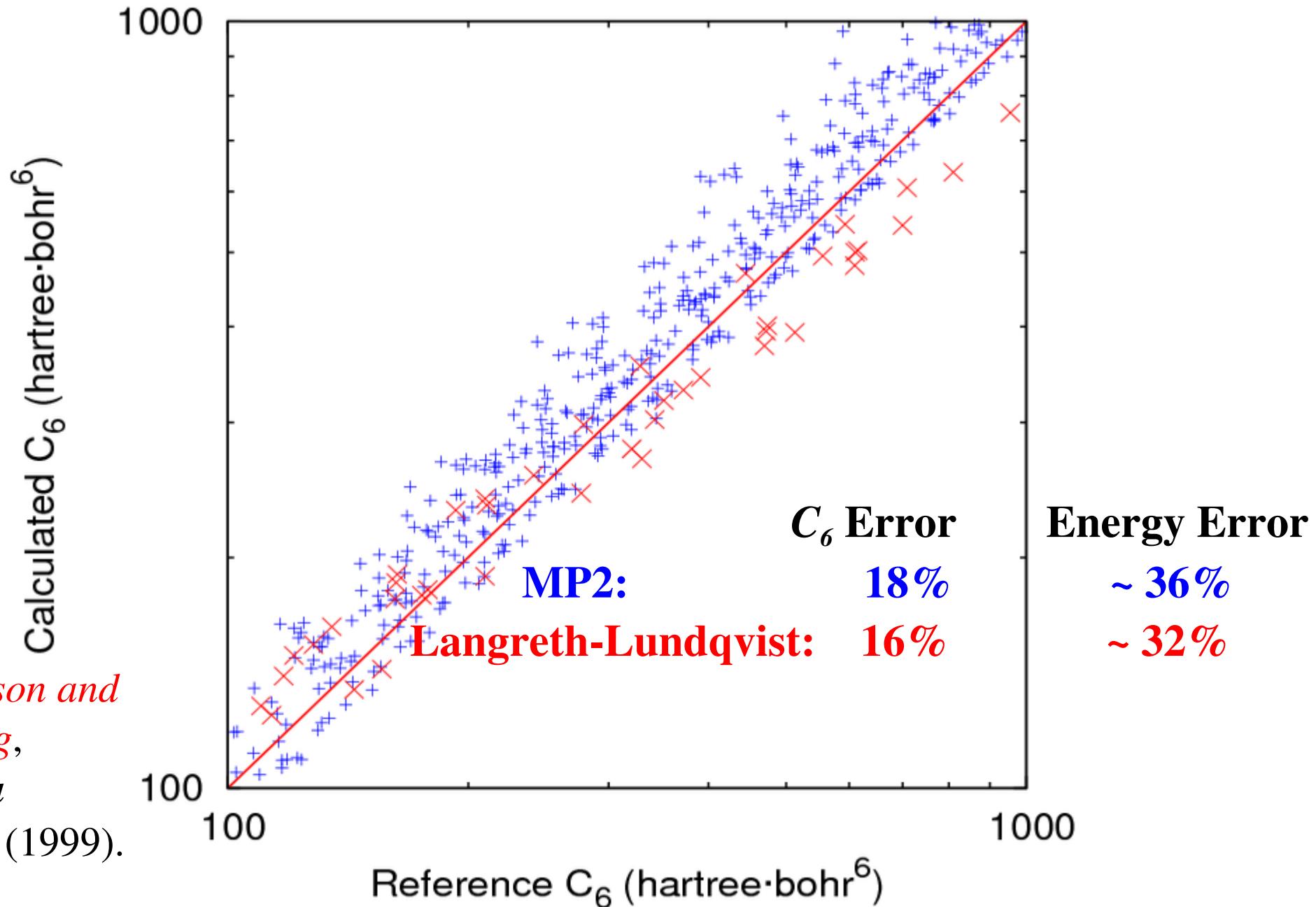
*Data available for atoms, alkanes, alkenes, alkynes, alcohols,  
 $H_2$ ,  $N_2$ ,  $H_2S$ ,  $NH_3$ ,  $SO_2$ ,  $COS$ ,  $CO_2$ ,  $CS_2$ ,  $SiH_4$ ,  $CCl_4$ , etc.  
(50 atoms and molecules – 1225 interaction pairs)*

<sup>(1)</sup> W. J. Meath and co-workers (1977-present)

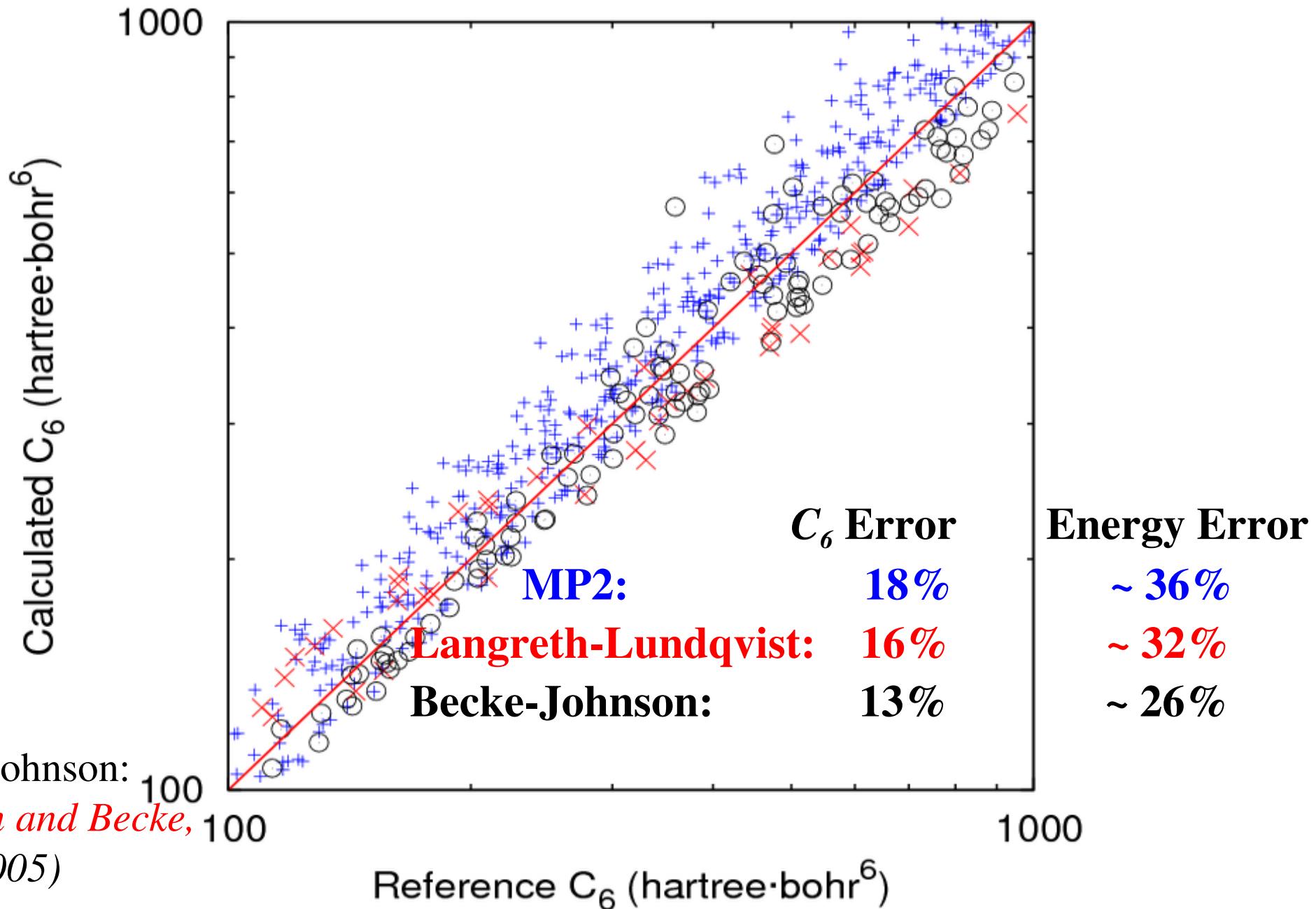
# Can we get accurate dispersion coefficients ?



# Can we get accurate dispersion coefficients ?

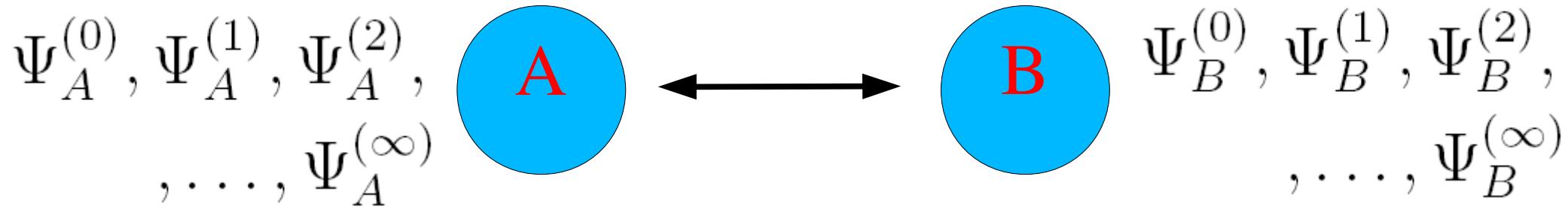


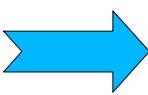
# Can we get accurate dispersion coefficients ?



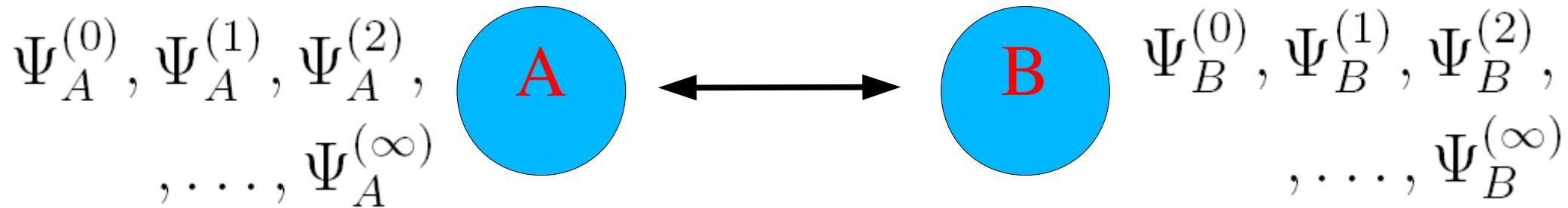
Let us start from  
*first principles* ...

# Atom-atom dispersion energy



No multipole moments on  $A$  and  $B$    $E_{int} = E_{disp}$

# Atom-atom dispersion energy



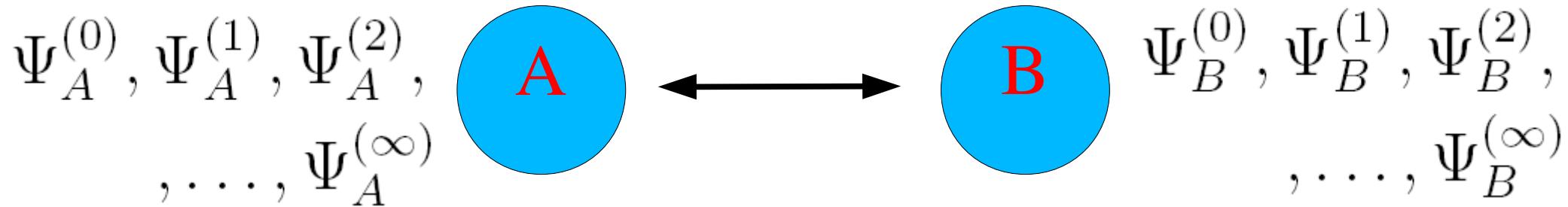
No multipole moments on  $A$  and  $B$   $E_{int} = E_{disp}$

$$E_{disp} = - \sum_{i \neq 0, j \neq 0} \frac{\langle 00 | H' | ij \rangle \langle ij | H' | 00 \rangle}{\epsilon_i^A + \epsilon_j^B - \epsilon_0^A - \epsilon_0^B} \quad H' = \sum_{a \in A} \sum_{b \in B} \frac{Z_a Z_b}{r_{ab}}$$

To leading order (long-range approximation):

$$E_{disp}^\infty = - \frac{6}{R_{AB}^6} \sum_{i \neq 0, j \neq 0} \frac{\langle 0 | \mu^A | i \rangle^2 \langle 0 | \mu^B | j \rangle^2}{\epsilon_i^A + \epsilon_j^B - \epsilon_0^A - \epsilon_0^B}$$

# Atom-atom dispersion energy



No multipole moments on  $A$  and  $B$   $\rightarrow E_{int} = E_{disp}$

$$E_{disp} = - \sum_{i \neq 0, j \neq 0} \frac{\langle 00 | H' | ij \rangle \langle ij | H' | 00 \rangle}{\epsilon_i^A + \epsilon_j^B - \epsilon_0^A - \epsilon_0^B} \quad C_6 \text{ coefficient}$$

To leading order (long-range approximation):

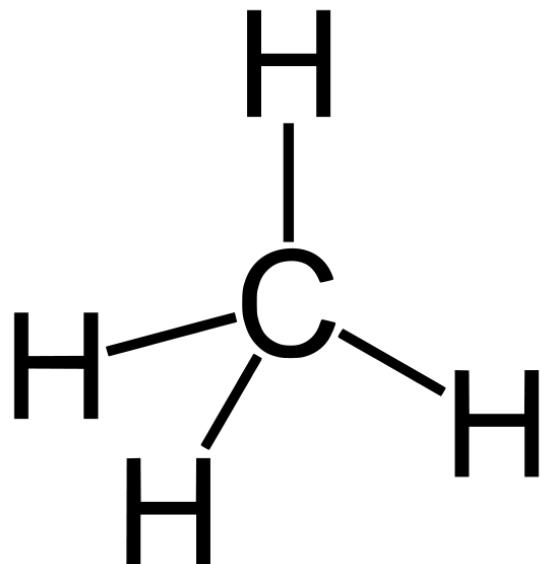
$$E_{disp}^\infty = - \frac{6}{R_{AB}^6} \sum_{i \neq 0, j \neq 0} \frac{\langle 0 | \mu^A | i \rangle^2 \langle 0 | \mu^B | j \rangle^2}{\epsilon_i^A + \epsilon_j^B - \epsilon_0^A - \epsilon_0^B}$$

# TS-vdW method: Basic ingredients

- 1) Only *relative* polarizability is used – avoiding inaccurate absolute polarizability of (semi-) local DFT or Hartree-Fock.
- 2) *Ab initio* reference *free-atom*  $C_6$  coefficients (*Chu and Dalgarno*).
- 3) Accurate combination rule for heteronuclear coefficients.
- 4) *Atom-in-a-molecule/solid* polarizability from Hirshfeld partitioning of the electron density (*Becke and Johnson*).
- 5)  $C_6 = C_6[n(r)]$ :  $C_6$  becomes a functional of the electron density.

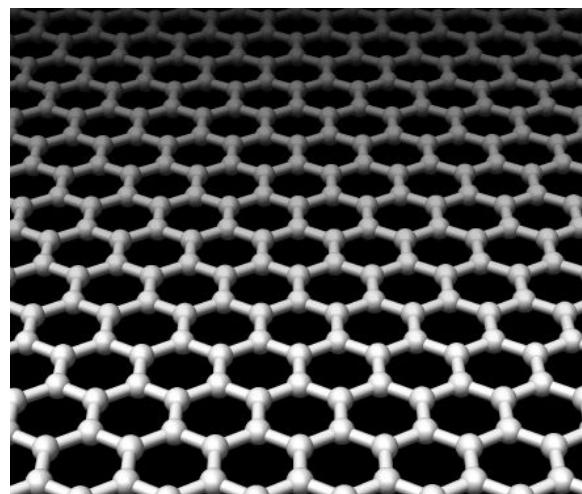
*A. Tkatchenko and M. Scheffler PRL (2009).*

# $C_6$ is a functional of the density (Carbon-Carbon $C_6$ coefficient)



**24.1** HartreeBohr<sup>6</sup>

$sp^3$



**33.0** HartreeBohr<sup>6</sup>

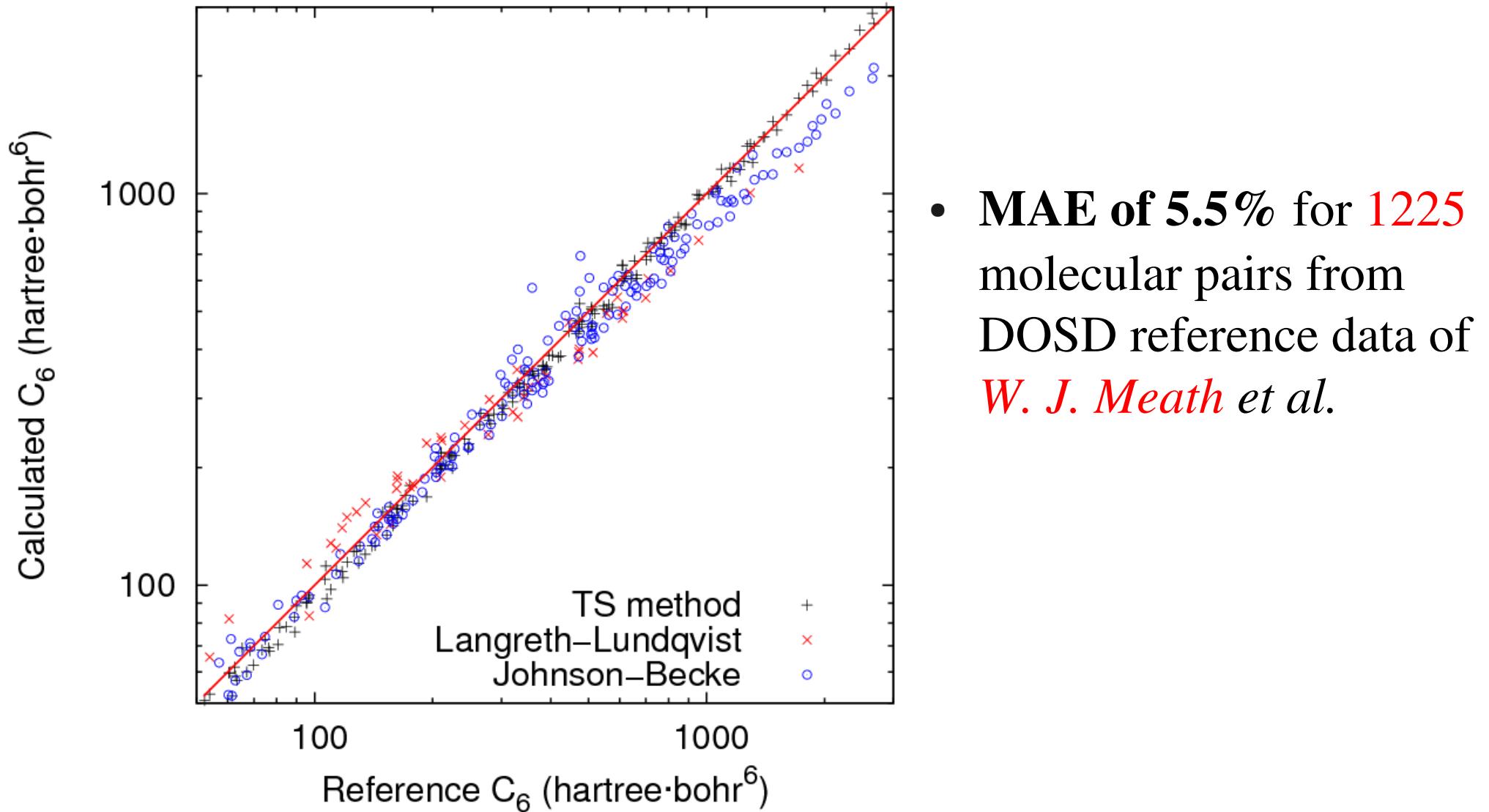
$sp^2$



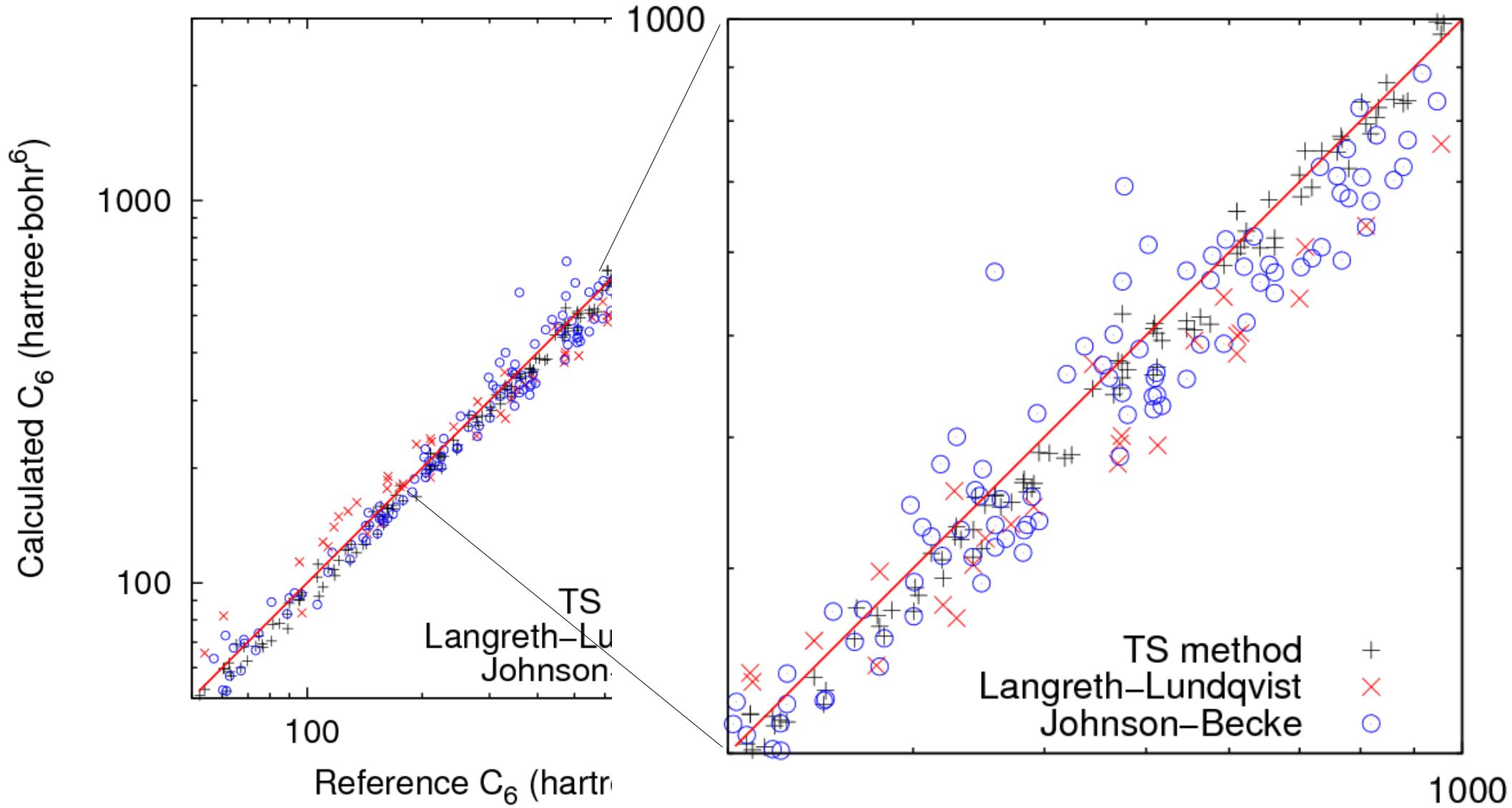
**38.6** HartreeBohr<sup>6</sup>

$sp^3$

# Performance of TS-vdW for molecules

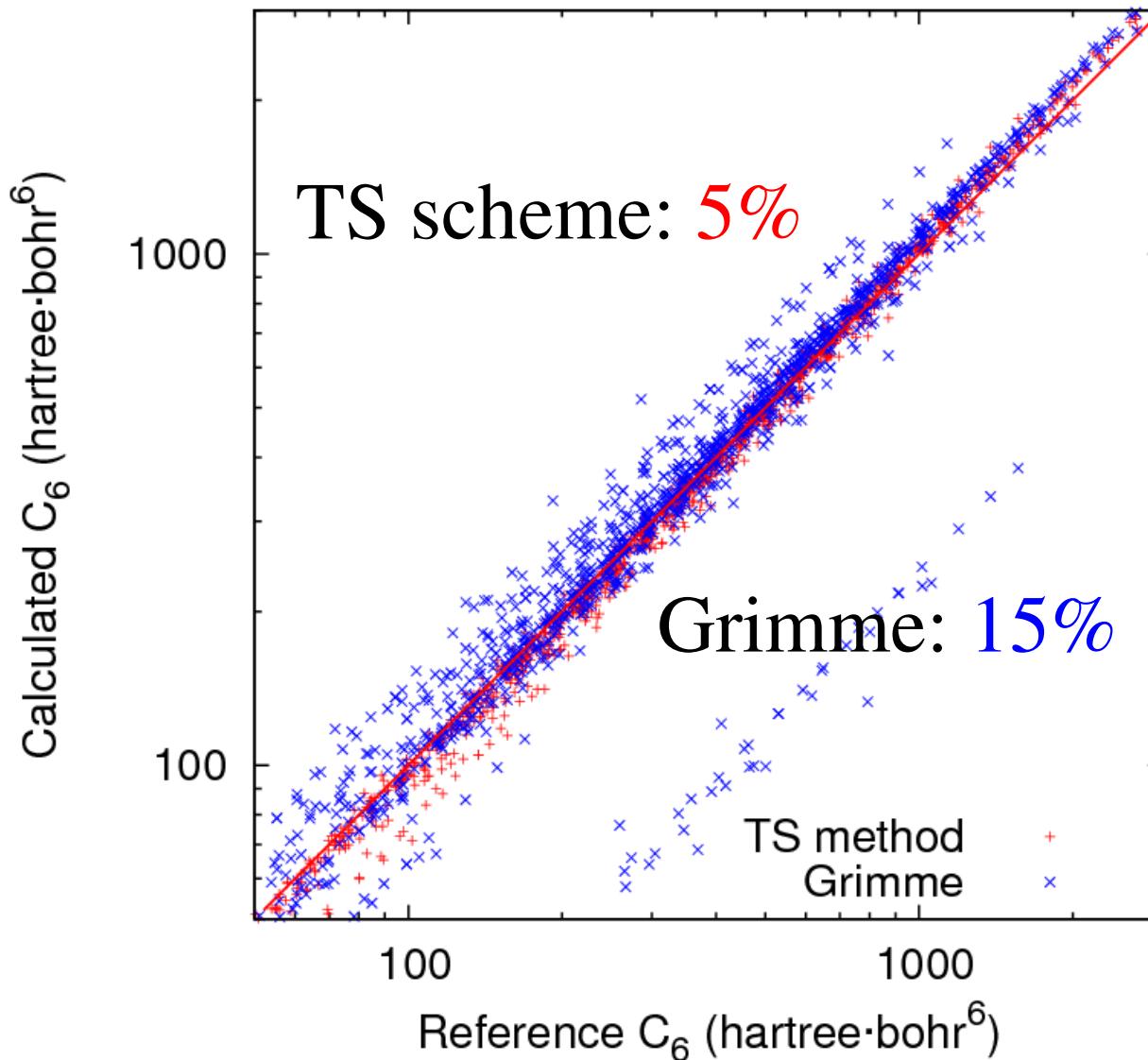


# Performance of TS-vdW for molecules



Results depend negligibly (1% deviation) on the employed  $xc$  functional

# Comparison to fixed $C_6$ schemes



... and we can treat all elements in the periodic table

# **DFT+vdW<sup>(1)</sup> & MP2+ΔvdW<sup>(2)</sup>**

Leading dispersion term (or  $\Delta C_6$  term) is added to DFT or MP2 total energy, damped at short distances,

$$E_{vdW} = - \sum_A \sum_{B>A} f_{damp}(R_{AB}, R_A^0, R_B^0) C_{6AB} R_{AB}^{-6}$$

$$R_{eff}^0 = \left( \frac{V^{eff}}{V^{free}} \right)^{1/3} R_{free}^0$$

- Effective vdW parameters are functionals of the electron density:

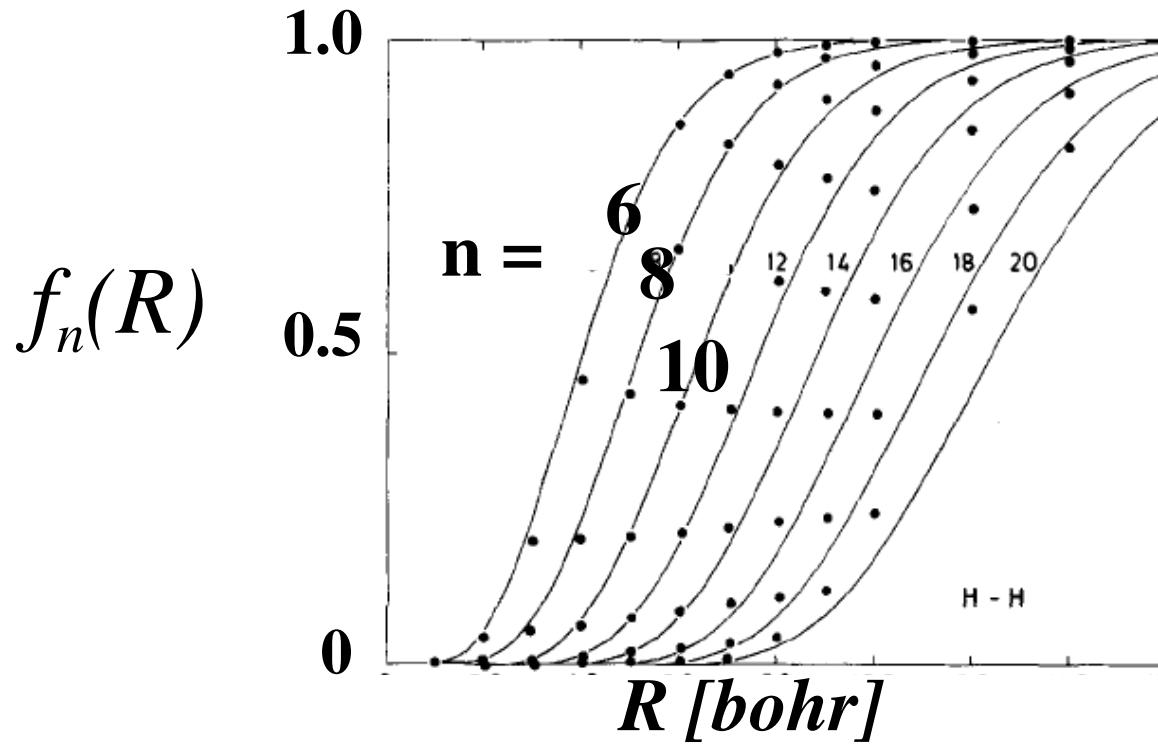
$$C_6 = C_6[n(\mathbf{r})], \quad R_{vdW} = R_{vdW}[n(\mathbf{r})]$$

- One damping parameter for DFT (for each functional), no parameters for MP2

<sup>(1)</sup>**DFT+vdW:** Tkatchenko and Scheffler PRL (2009)

<sup>(2)</sup>**MP2+ΔvdW:** Tkatchenko, DiStasio Jr., Head-Gordon, Scheffler JCP (2009)

# Short-range damping function



**Damping function:**

- Analytical form known *exactly* for dispersion energy in H-H
- Law of corresponding states for other spherical atoms (rare-gas, alkaline earth, mercury)
- Parameterized when used with DFT or quantum chemical methods

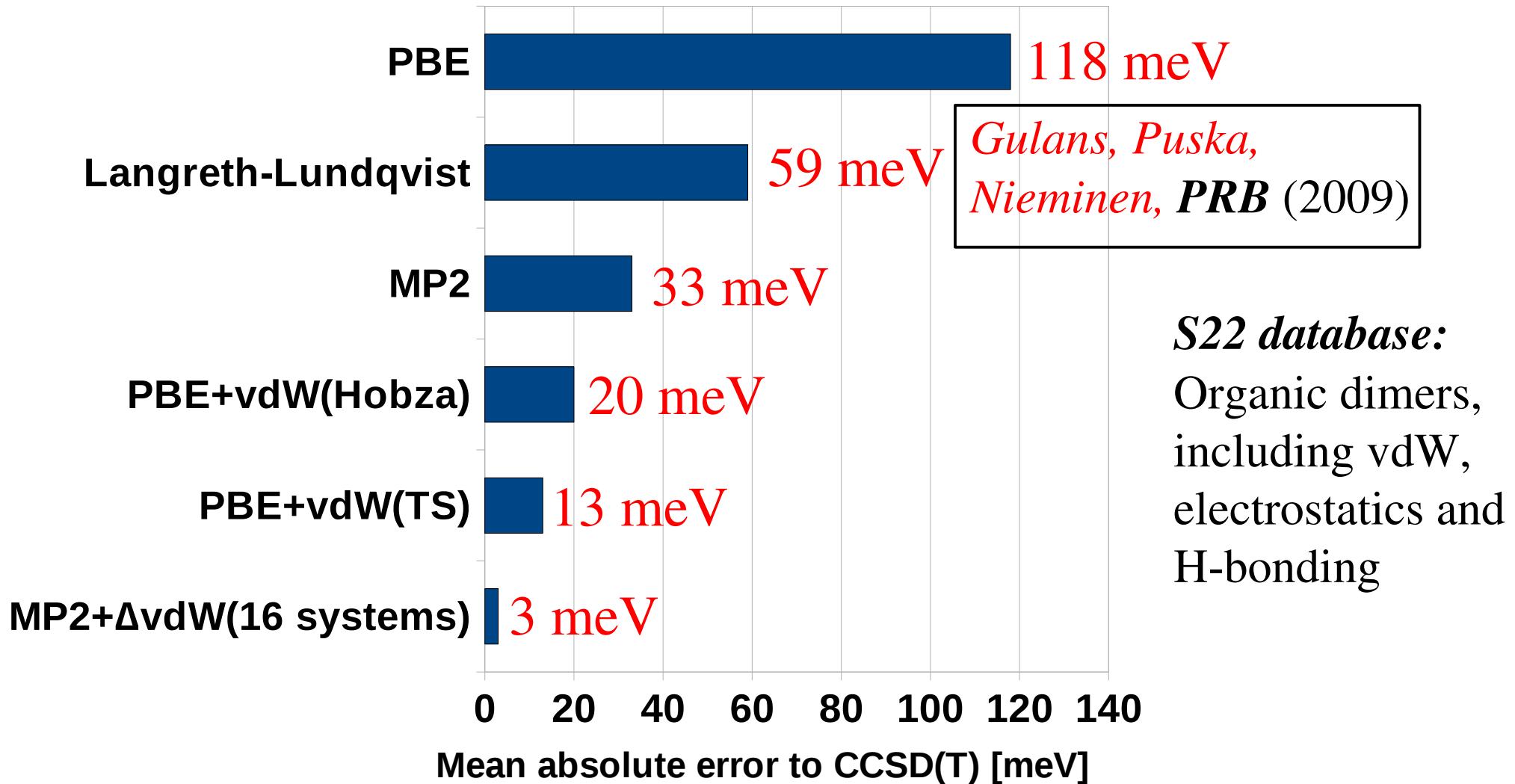
Analytic:

*Aziz, Koide, Meath, Tang, Toennies, et al.*

HF&DFT:

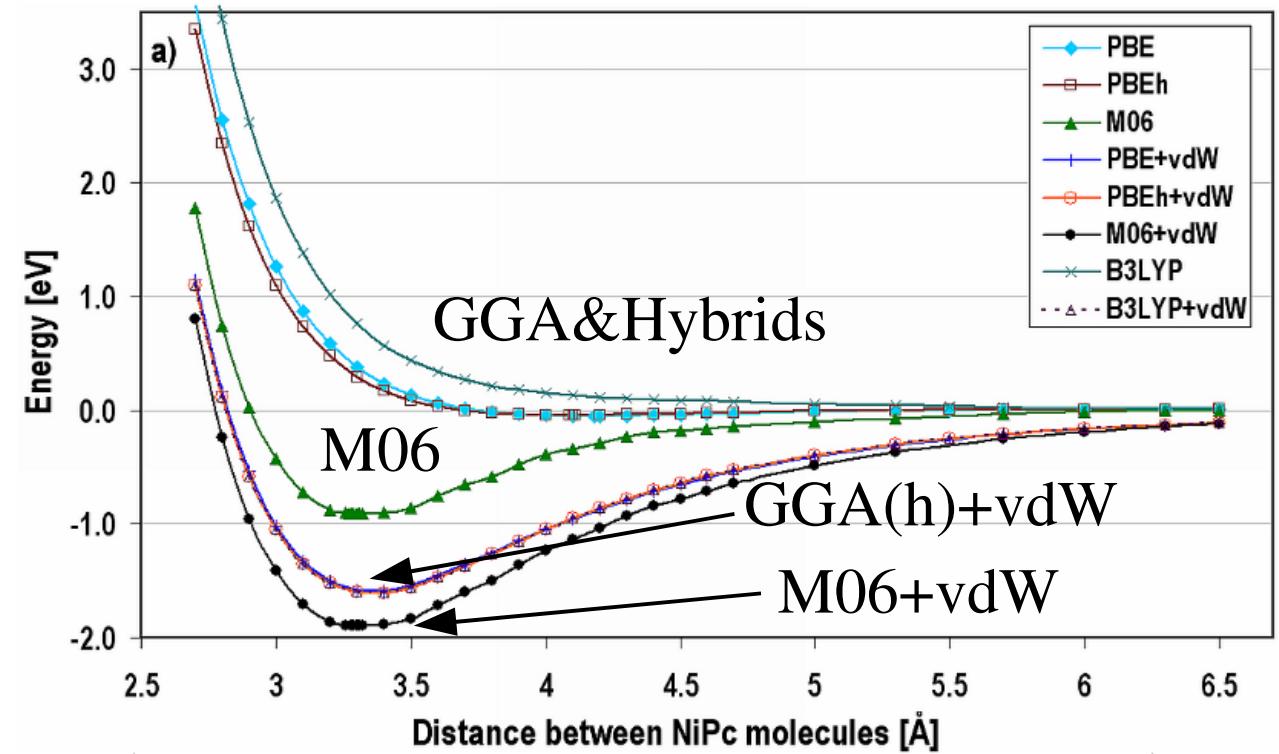
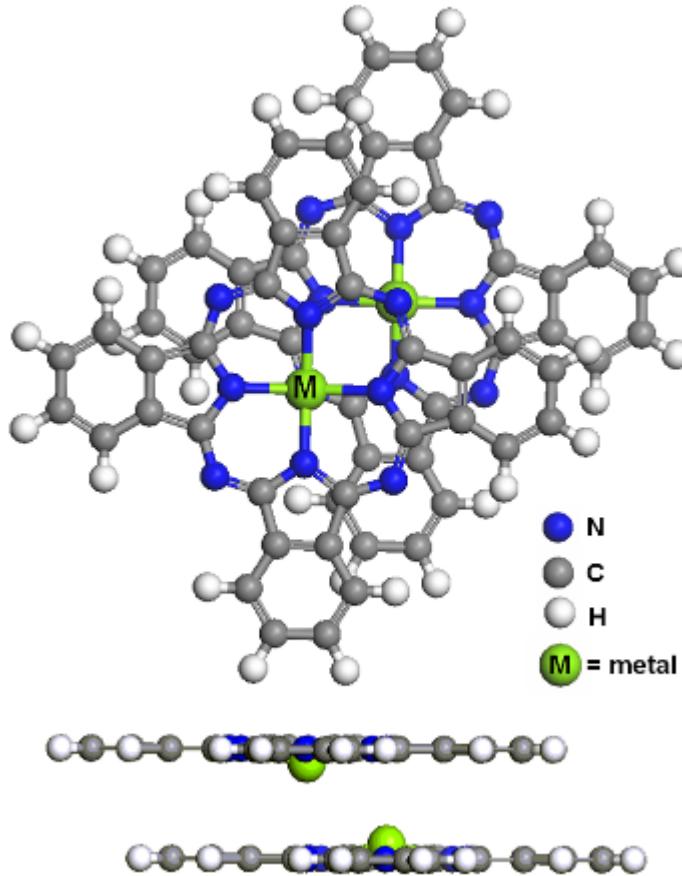
*Scoles, Yang, Grimme, Bechstedt, Hobza, et al.*

# Performance for intermolecular interactions (S22\*)



\**Jurecka, Sponer, Cerny, Hobza PCCP 8, 1985 (2006)*

# Importance of vdW asymptotics (large molecules)



# Outlook

- Extending the method to solids and adsorption  
(dielectric screening in **metals**, etc.)
- Short-range damping function
- Anisotropic and non-additive many-body vdW interactions



# Summary

Alexander von Humboldt  
Stiftung / Foundation

- Developed accurate *first-principles* method for the long-range van der Waals interactions
- Can be coupled to DFT ( $DFT+vdW$ ) or MP2 ( $MP2+\Delta vdW$ )
- Consistently accurate results for a variety of applications
- Can be coupled to hybrid functionals for improved electronic structure (*Marom, Tkatchenko, Scheffler, Kronik, submitted*) or followed by  $GW$  calculations