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Long-range universality in cold two and three body collisions

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NIST and The University of Maryland

Thanks to many colleagues in theory and experiment
who have contributed to this work

<http://www.jqi.umd.edu/>



Joint
Quantum
Institute

Outline

Long-range universality of bound states and collisions:

- 1.Two-body (QDT picture)
2. Three-body
3. Molecules and reactive collisions
4. Species with complex or chaotic collisions

What do we mean by “universal”?

Answer: Independent of “short-range” details,
characterized by a few simple parameters.

Example: zero-range interaction proportional to
s-wave scattering length a

Only one parameter **a** depends on the “details”

s-wave scattering phase shift: $\tan \eta(k) \approx -ka$

Bound state energy: $E_b = -\frac{\hbar^2}{2\mu a^2}$

Threshold collision summary

s-wave

$$S(k) = e^{2i\eta(k)} \rightarrow e^{-2ik(a-ib)} \text{ as } k \rightarrow 0$$

Complex scattering length $a-ib$

Balakrishnan, et al., Chem. Phys. Lett. 280, 5 (1997)

Bohn and PSJ, Phys. Rev. A 56, 1486 (1997)

J. Hutson, New J. Phys. 9, 152 (2007)

E-dependent $\alpha(k) = a(k) - ib(k) = -\frac{\tan \eta(k)}{k} = \frac{1}{ik} \frac{1 - S(k)}{1 + S(k)}$

mean field	$\frac{4\pi\hbar}{\mu} a(0)n$
loss rate	$\frac{4\pi\hbar}{\mu} b(0)n$

Need variation with E away from E=0

- precision binding energy measurements
- lattice zero point energy $a(E_n)$
- few-body beyond “scattering length universality”
- finite temperature QDGs

Effective range expansion

$$k \cot \eta(E) = -\frac{1}{a} + \frac{1}{2} r_0 k^2$$

$$r_0 = 2.918 \bar{a} \frac{\bar{a}^2 + (a - \bar{a})^2}{a^2}$$

Flambaum, Gribakin, and Harabati,
Phys. Rev. A 59, 1998 (1999)
Gao, Phys. Rev. A 62, 050702 (2000)

$$\bar{a} = \frac{2\pi}{\Gamma(\frac{1}{4})^2} \left(\frac{2\mu C_6}{\hbar^2} \right)^{\frac{1}{4}}$$

Bound state corrections

$$E_b = -\frac{\hbar^2}{2\mu(a - \bar{a})^2}$$

Gribakin, Flambaum,
Phys. Rev. A 48, 446 (1993)
or improved formula of
Gao, J. Phys. B 37, 4273 (2004)

Still not good enough—we can do much better

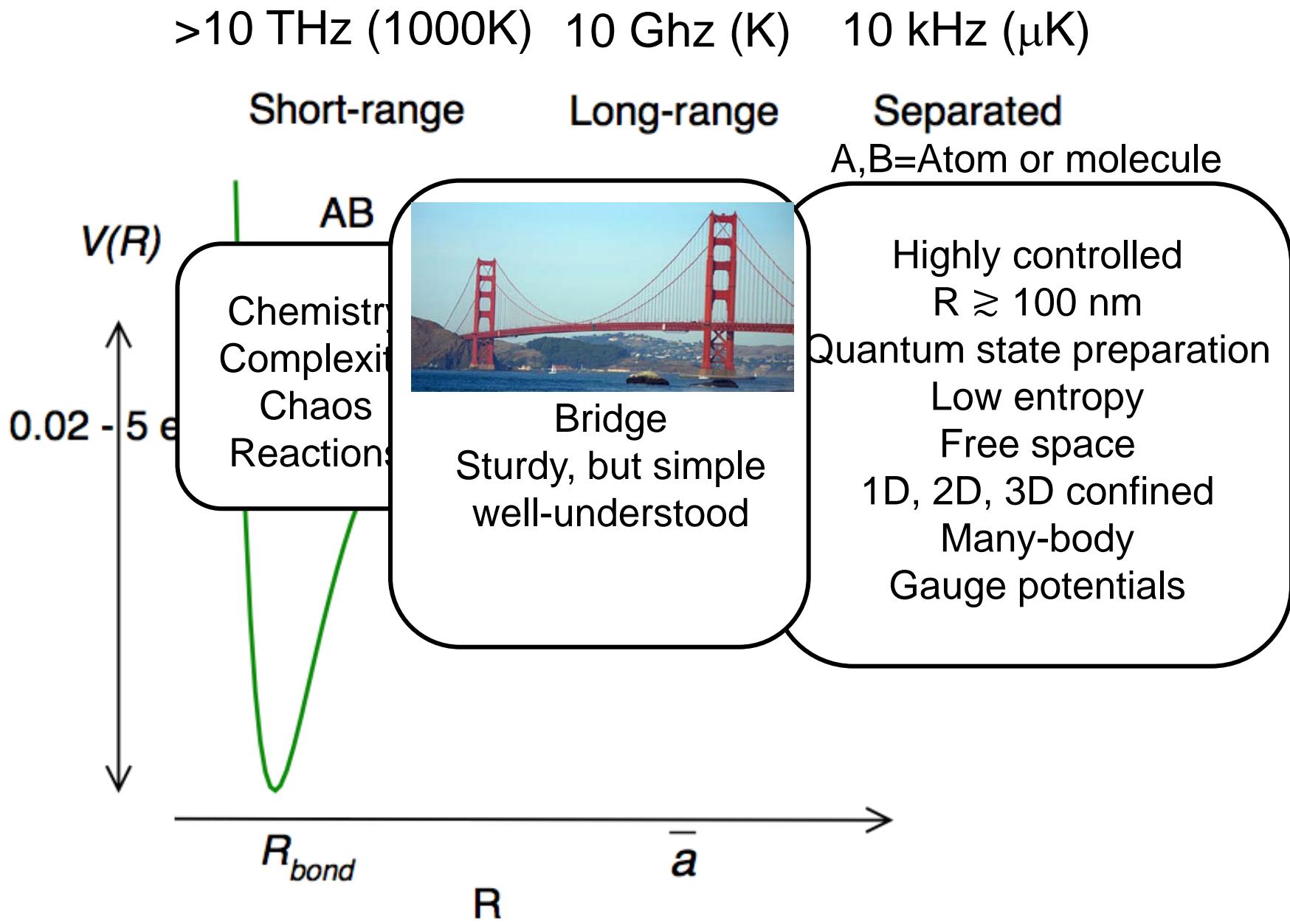


Fig. 2, PSJ Faraday Disc 142, 361 (2009)

>10 THz (1000K) 10 Ghz (K) 10 kHz (μ K)

Short-range

A few
Parameters (QDT)

Phase (a_{bg})
Feshbach strength (s_{res})
Reactivity (y)

Chemistry
Complexity
Chaos
Reactions

Long-range



Not unique
More than one way
To build a bridge
(Mies/PSJ/Hutson
Greene/Bohn, Gao)
Analytic
Numerical

Separated

A,B=Atom or molecule

Highly controlled
 $R \gtrsim 100$ nm
Quantum state preparation
Low entropy
Free space
1D, 2D, 3D confined
Many-body
Gauge potentials

R_{bond}

R

\bar{a}

Fig. 2, PSJ Faraday Disc 142, 361 (2009)

Quantum defect theory

1. Pick a **reference problem** we can solve

Classic example: Coulomb potential, H-like atom
or $p = 6$ or $p = 4$ potential

Independent solutions $f(R, E)$, $g(R, E)$

2. Parameterize dynamics by a **few “physical” QDT parameters**
subject to experimental fitting
and theoretical interpretation

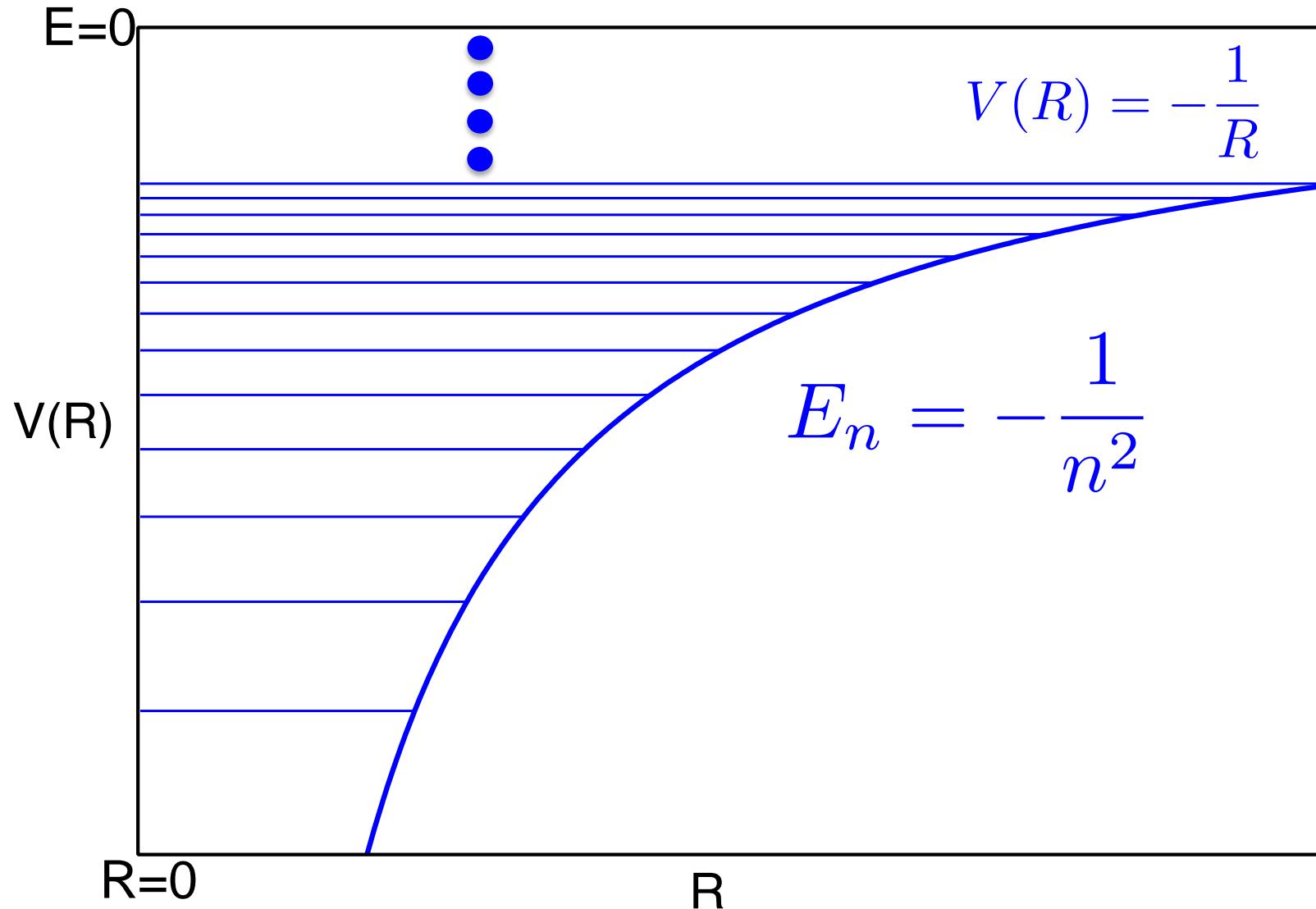
phase (diagonal, scattering length)

interactions (non-diagonal, multichannel)

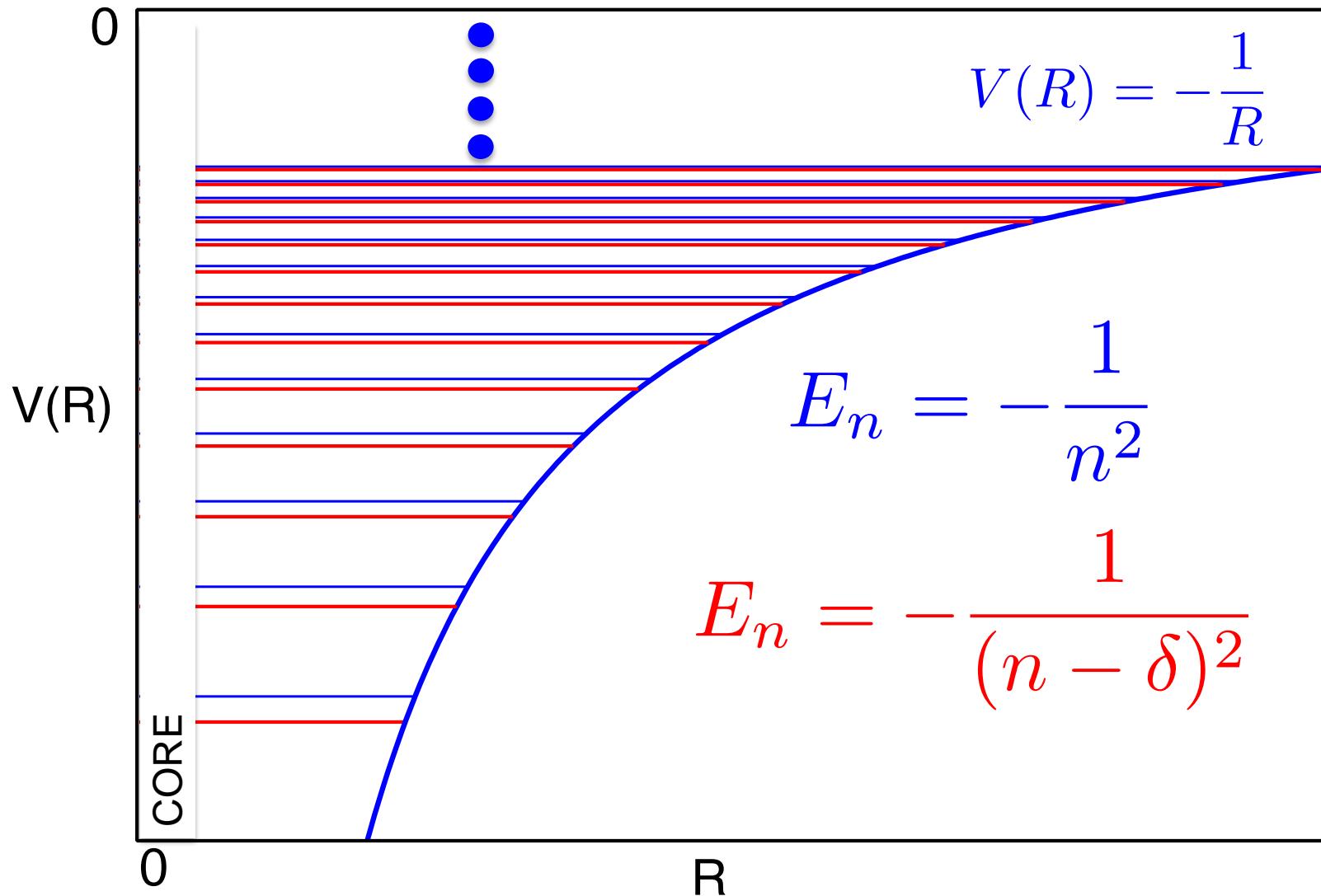
3. Use methods of QDT to calculate
bound and scattering states, resonances, cross sections, etc.

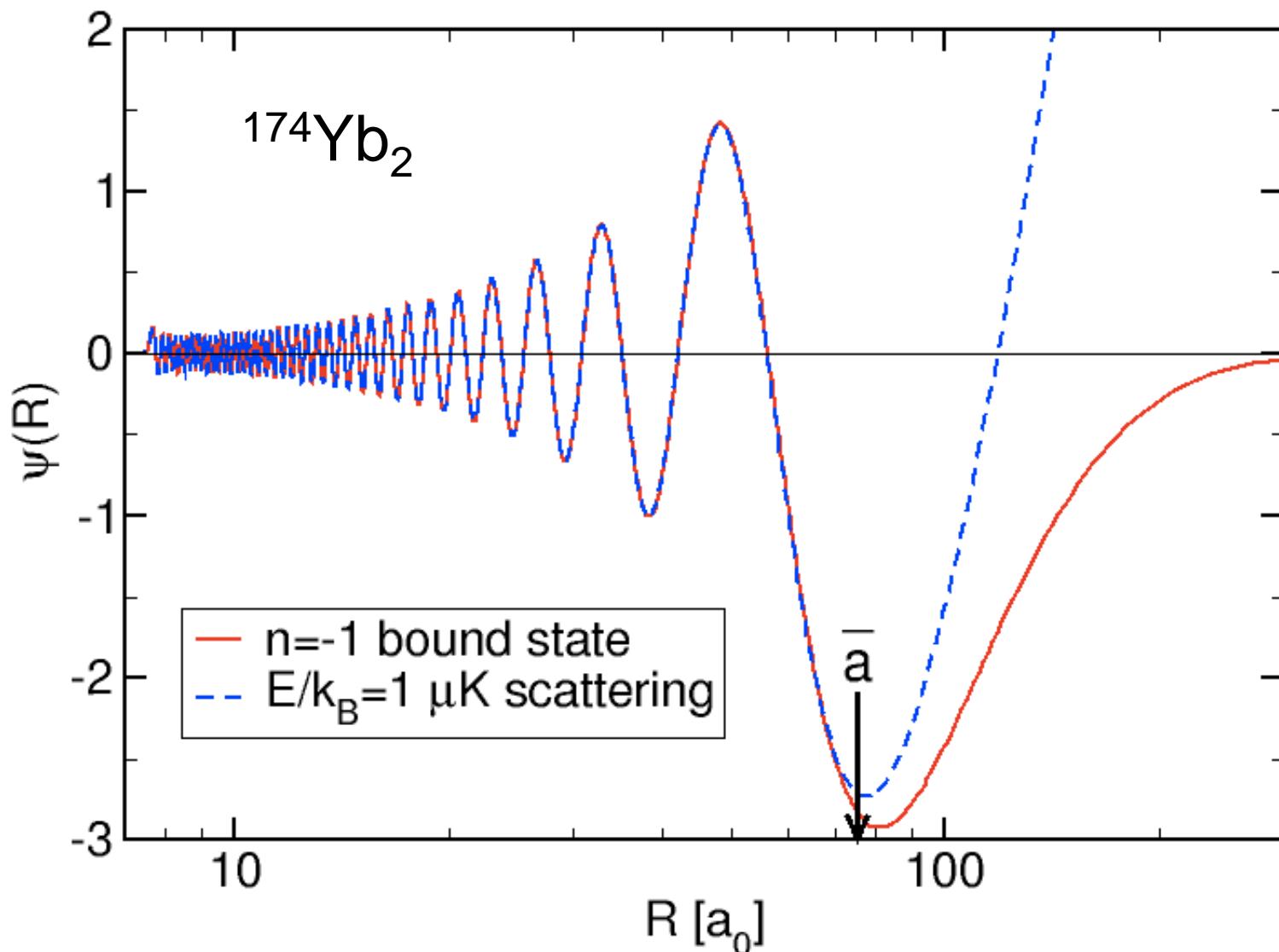
$$\Psi(R, E) = [f(R, E) + g(R, E) K] A$$

H atom



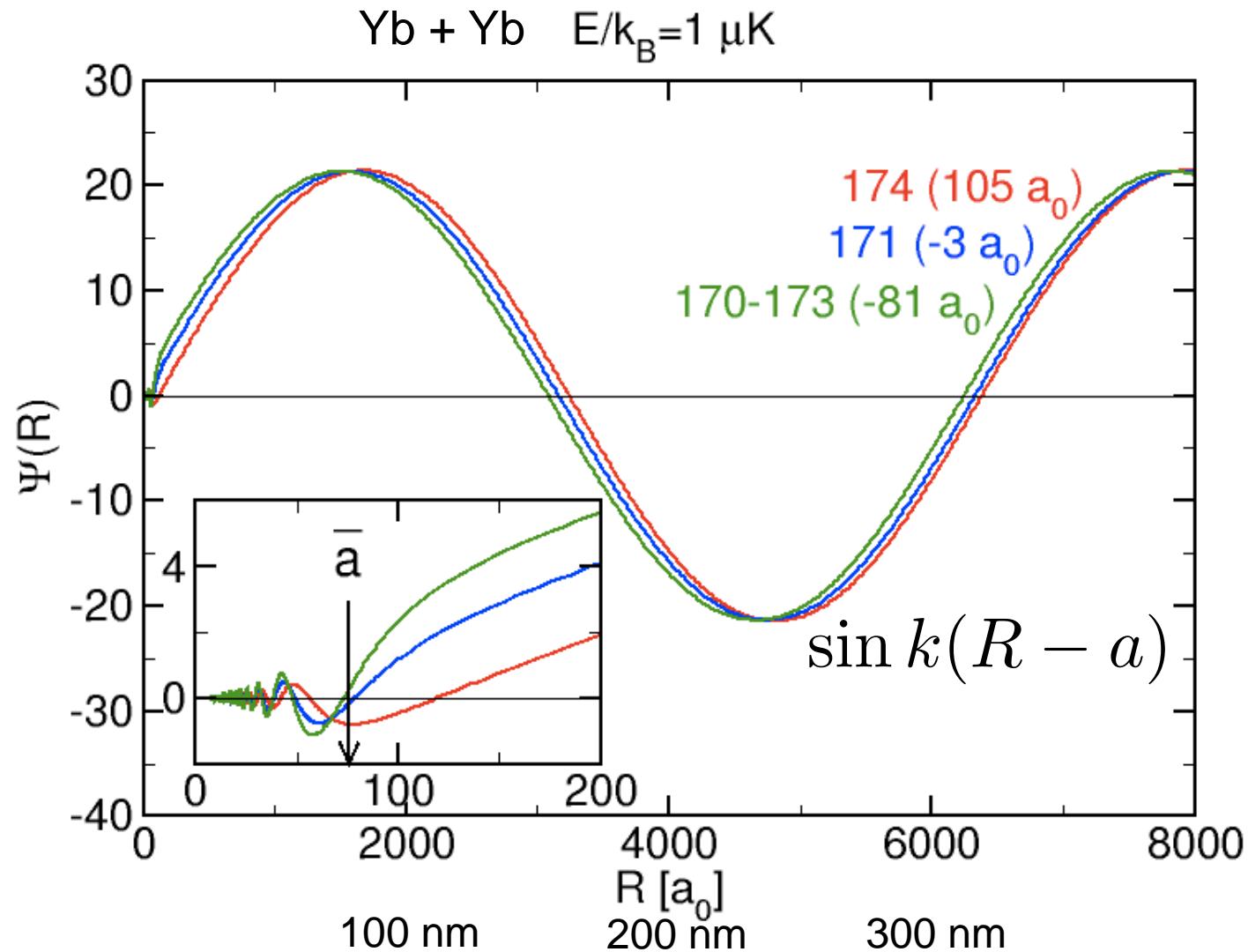
Multi-electron atom





From Krems et al, Cold Molecules, PSJ Chapter 6, arXiv:0902.1727

QDT: The s-wave scattering length + long range potential determine the scattering and bound state properties (of all partial waves) near threshold.



Van der Waals potential

Write the Schrödinger equation in length and energy units of

$$R_{\text{vdw}} = \frac{1}{2} \left(\frac{2\mu C_6}{\hbar^2} \right)^{\frac{1}{4}}$$
 or $\bar{a} = \frac{\Gamma(3/4)}{\Gamma(5/4)} R_{\text{vdW}} = 0.956 R_{\text{vdW}}$

Gribakin and Flambaum, Phys. Rev. A 48, 546 (1993)

$$E_{\text{vdw}} = \frac{\hbar^2}{2\mu R_{\text{vdw}}^2}$$

The potential becomes $-\frac{16}{r^6} + \frac{\ell(\ell+1)}{r^2}$ vdw units.

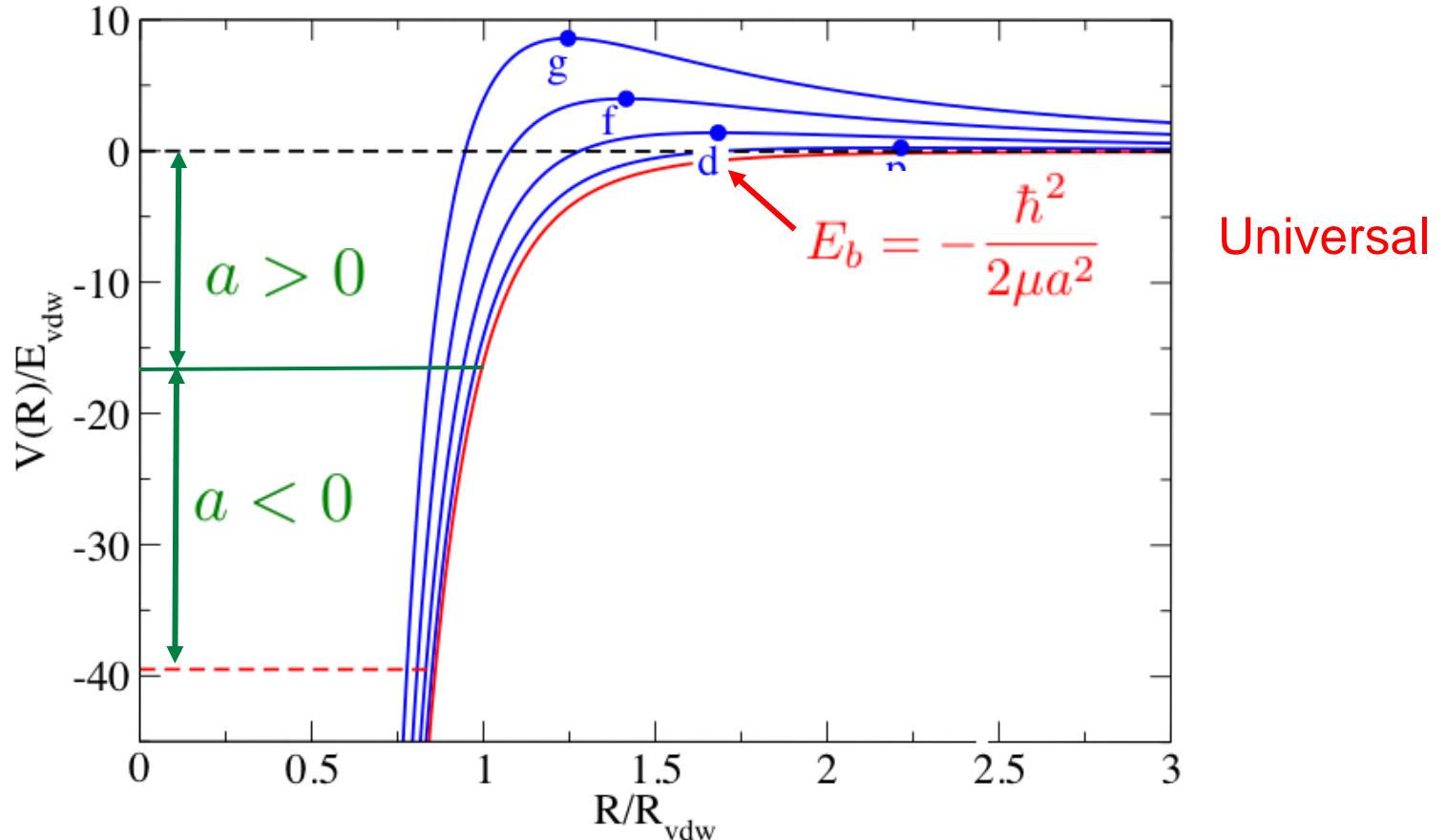
This potential has exact analytic solutions and many useful properties.

B. Gao, Phys. Rev. A 58, 1728, 4222 (1998) + series of papers.

See Jones, et al., Rev. Mod. Phys. 78, 483 (2006) Photoassociation

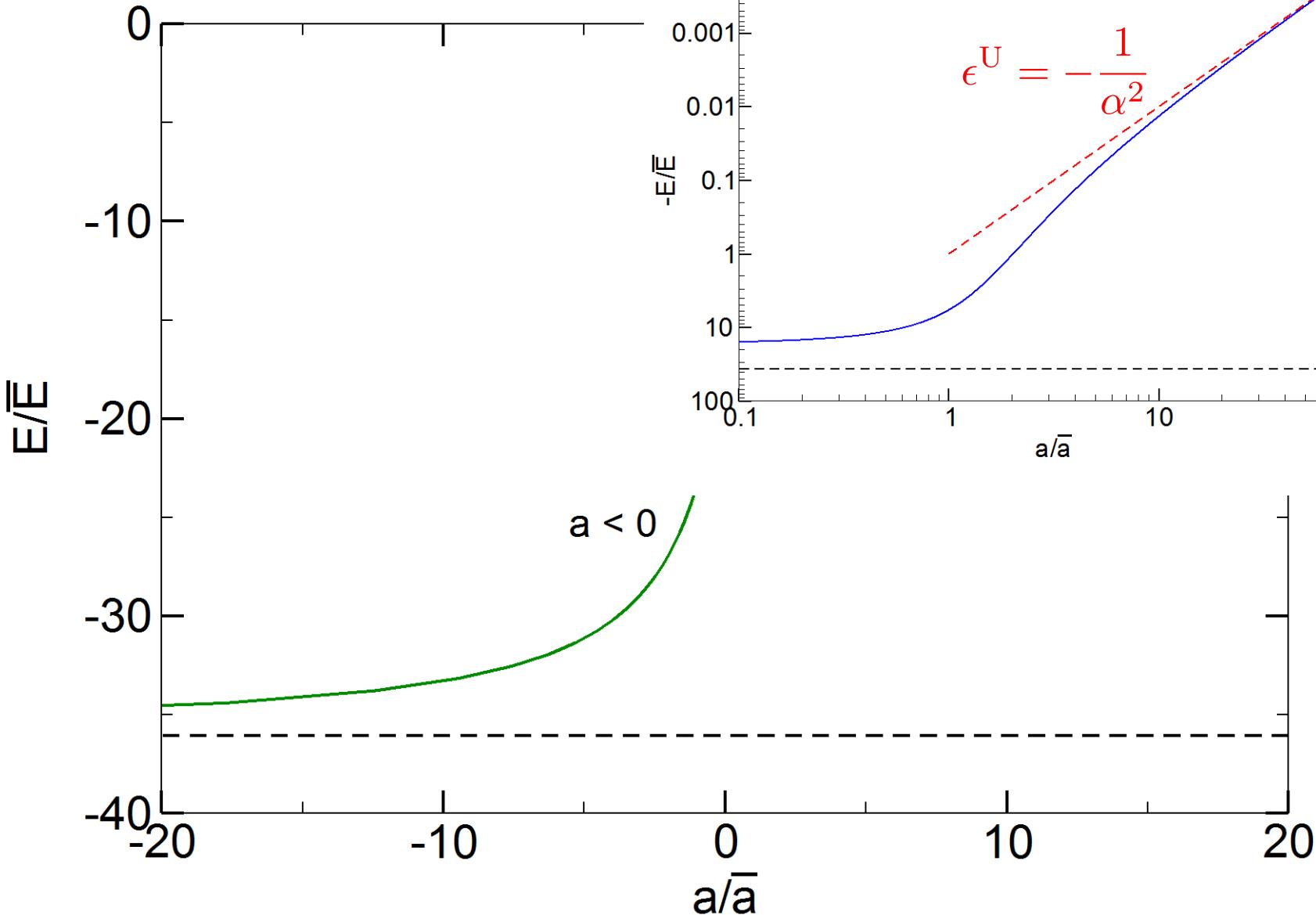
and Chin et al., Rev. Mod. Phys. 82, 1225 (2010) Feshbach resonances

“Size” of vdW potential

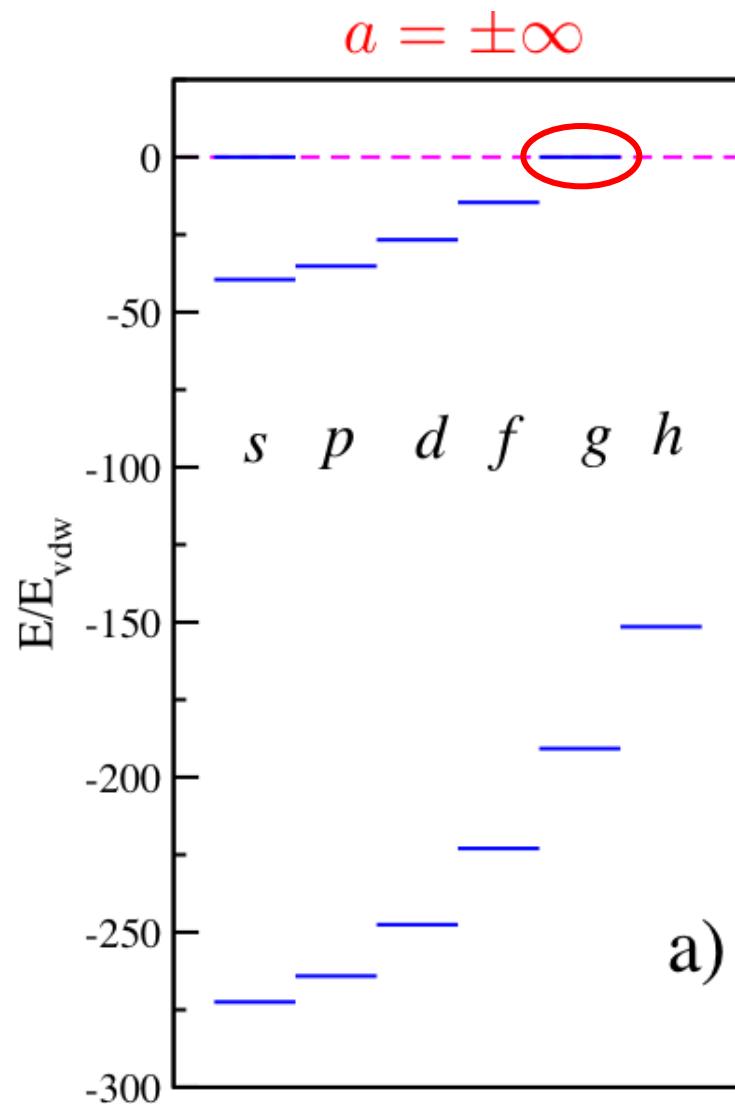


Jones, et al., Rev. Mod. Phys. 78, 483 (2006)

Universal energy of last (-1) bound state of vdW potential



Universal vdW bound state spectrum: depends on a



Gao, Phys. Rev. A 62, 050702 (2000); Chin et al, Rev. Mod. Phys. 82, 1225 (2010)

Resonant Scattering Picture

following U. Fano, Phys. Rev. 124, 1866 (1961); see Chin et al, RMP (2010)

Bound state

$$\frac{|n\rangle}{\text{V}}$$

Continuum



Closed channel
(Resonance)

Open channel
(Background)

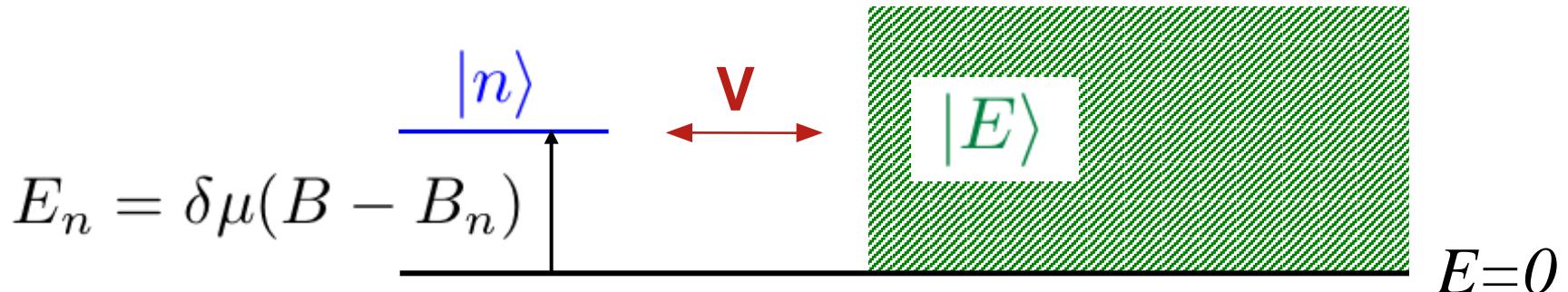
$$\eta(E) = \eta_{\text{bg}} + \eta_{\text{res}}(E)$$

$$\eta_{\text{res}} = -\tan^{-1} \frac{\frac{1}{2}\Gamma_n}{E - E_n - \delta E_n}$$

width $\Gamma_n = 2\pi|\langle n|V|E\rangle|^2$

shift $\delta E_n = \int \frac{|\langle n|V|E'\rangle|^2}{E_n - E'} dE'$

Threshold Resonant Scattering



$$\eta(E, B) = \eta_{\text{bg}}(E) - \tan^{-1} \frac{\frac{1}{2}\Gamma(E)}{E - E_n - \delta E_n(E)}$$

As $E \rightarrow 0$ $\eta_{\text{bg}} = -ka_{\text{bg}}$

$$\frac{1}{2}\Gamma_n(E) = (ka_{\text{bg}}) \delta\mu \Delta_n$$

$$a(B) = a_{\text{bg}} \left(1 - \frac{\Delta_n}{B - B_0} \right)$$

Shifted $B_0 = B_n + \delta B_n$

$$s_{\text{res}} = \frac{a_{\text{bg}}}{\bar{a}} \frac{\Delta\mu_{\text{diff}}}{\bar{E}}$$

Universal form of vdW threshold Feshbach scattering

$$\eta(\epsilon, B) = \eta_{\text{bg}}(\epsilon) - \tan^{-1} \left(\frac{\kappa s_{\text{res}}}{\epsilon - m_{\text{dif}}(B - B_0)} \right)$$

Strength



$$\epsilon = E/\bar{E} \quad \eta_{\text{bg}} = -\kappa \alpha_{\text{bg}} \quad \kappa = k \bar{a} \quad m_{\text{dif}} = \frac{\mu_{\text{dif}}}{\bar{E}}$$

“background”
scattering length

$\eta_{\text{bg}}(\epsilon)$, $s_{\text{res}}(\epsilon)$, $B_0(\epsilon)$ are UNIVERSAL functions of ϵ , α_{bg}

PSJ & Gao, Atomic Physics 20 (ICAP 2006), ed. by C. Roos, H. Haffner, and R. Blatt
(available at arxiv:0609013).

More accurate than the effective range expansion:
Blackley, Hutson, PSJ, PRA 89, 042701 (2014)

Extension to multiple overlapping resonances:
Jachymski & PSJ, PRA 88, 052701 (2013)

Universal three-body properties

Universal few-body physics

Universality of the Three-Body Parameter for Efimov States in Ultracold Cesium

M. Berninger,¹ A. Zenesini,¹ B. Huang,¹ W. Harm,¹ H.-C. Nägerl,¹ F. Ferlaino,¹ R. Grimm,^{1,2} P. S. Julienne,³ and J. M. Hutson⁴

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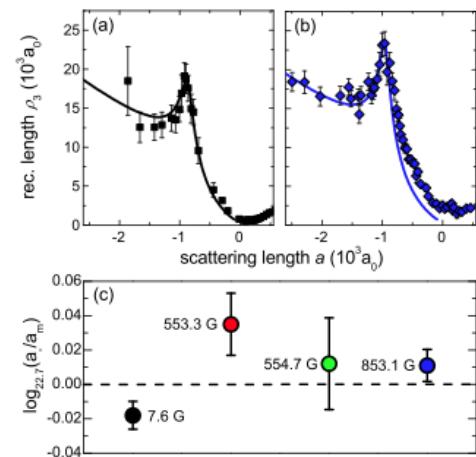
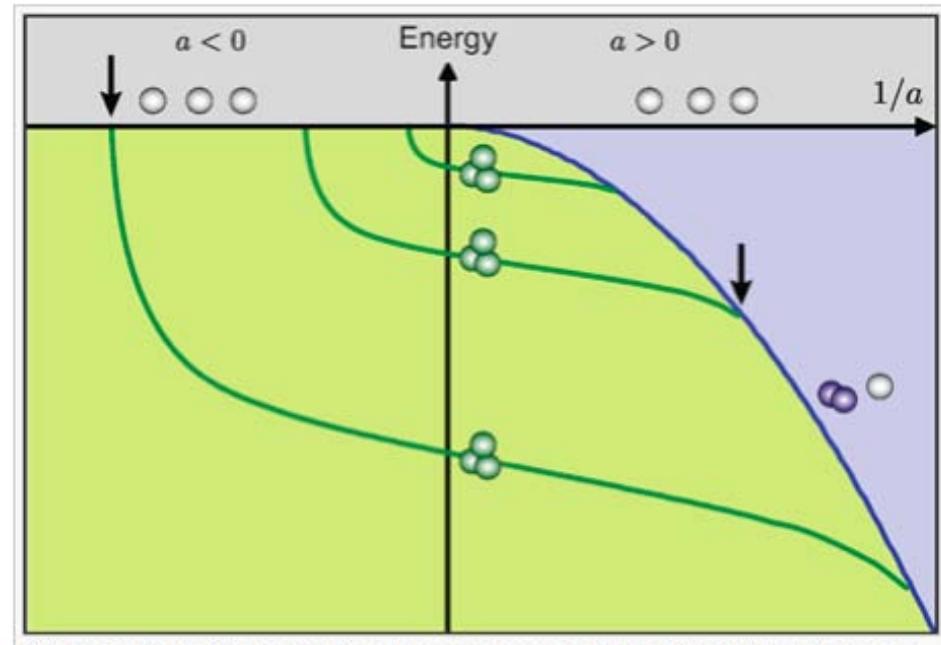
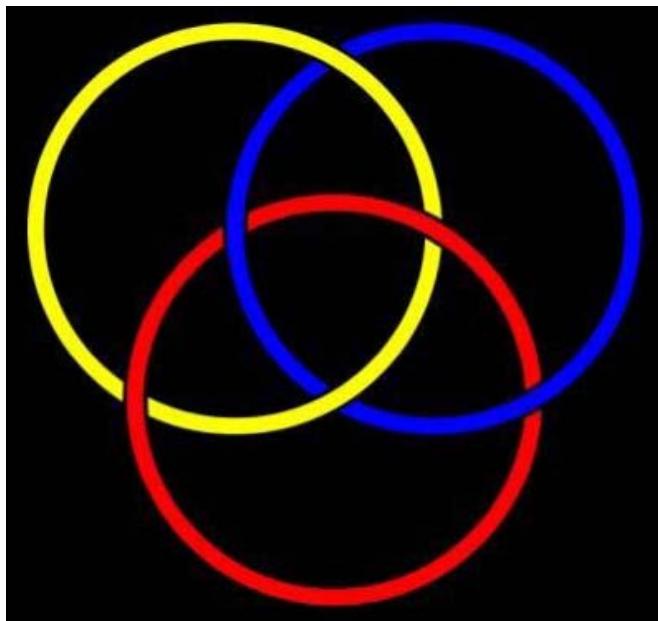
(Dated: July 27, 2011)

Phys. Rev. Lett. 107, 120401 (2011)

Original Cs experiment: T. Kraemer et al., Nature 440, 315 (2006)

Theory: J. Wang et al, Phys. Rev. Lett. 108, 263001 (2012)

Naidon et al, Phys. Rev. Lett. 112, 105301 (2014)

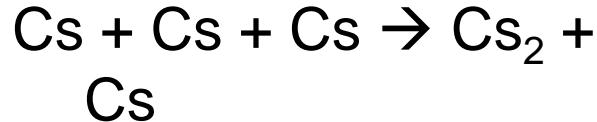


3-Body recombination of 3 alkali-metal atoms

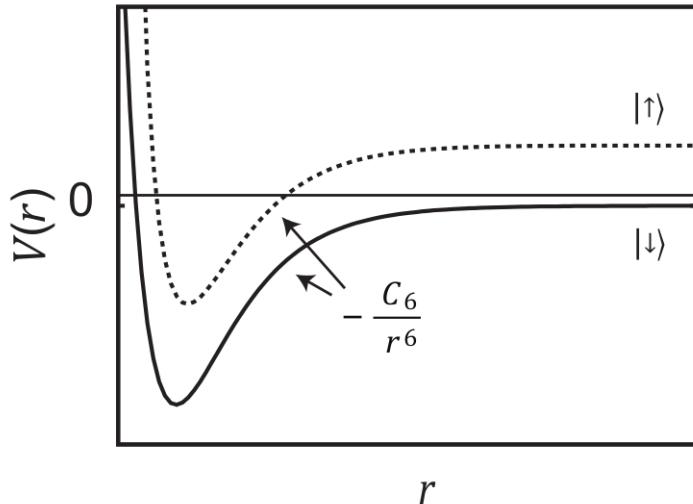
Universal van der Waals physics

Use only known 2-body physics + additive pairwise van der Waals potentials for 3 atoms

Y. Wang & PSJ, Nat. Phys. 10, 768 (2014)



Two-channel Cs + Cs interaction: “Exact” 2-body Feshbach model

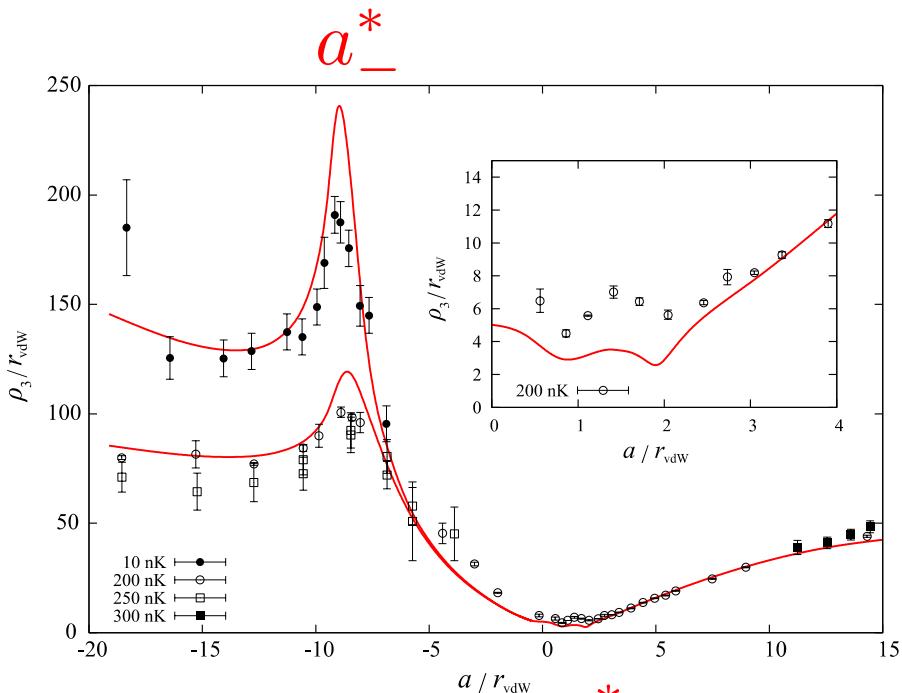


2-channel numerical model
Using s_{res} , α_{bg} , m_{dif} for Cs-Cs

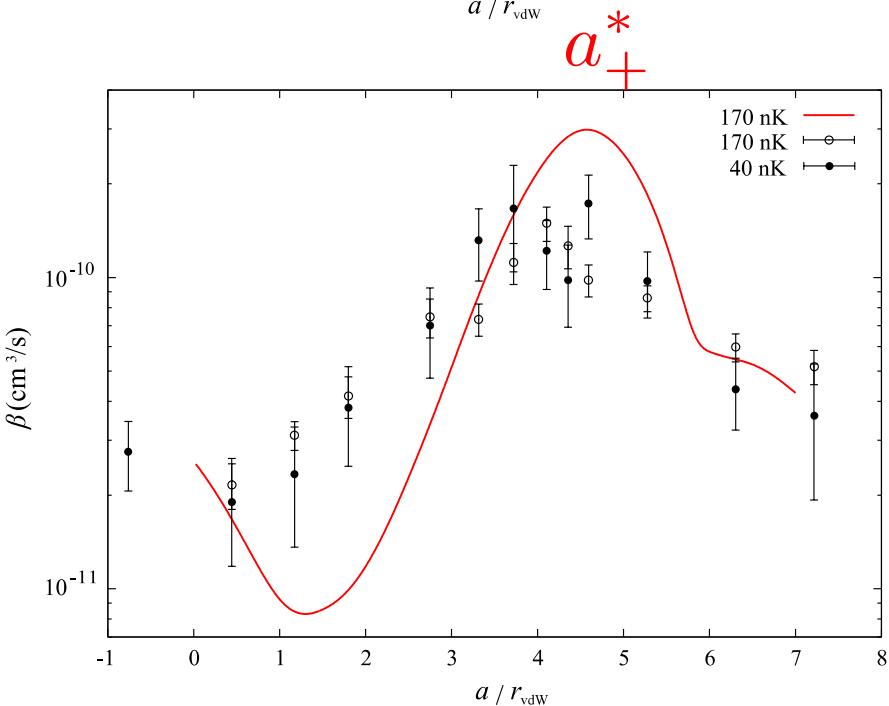
6-12 Lennard-Jones potentials
+ short-range coupling
Mies (2000), PSJ(2006)

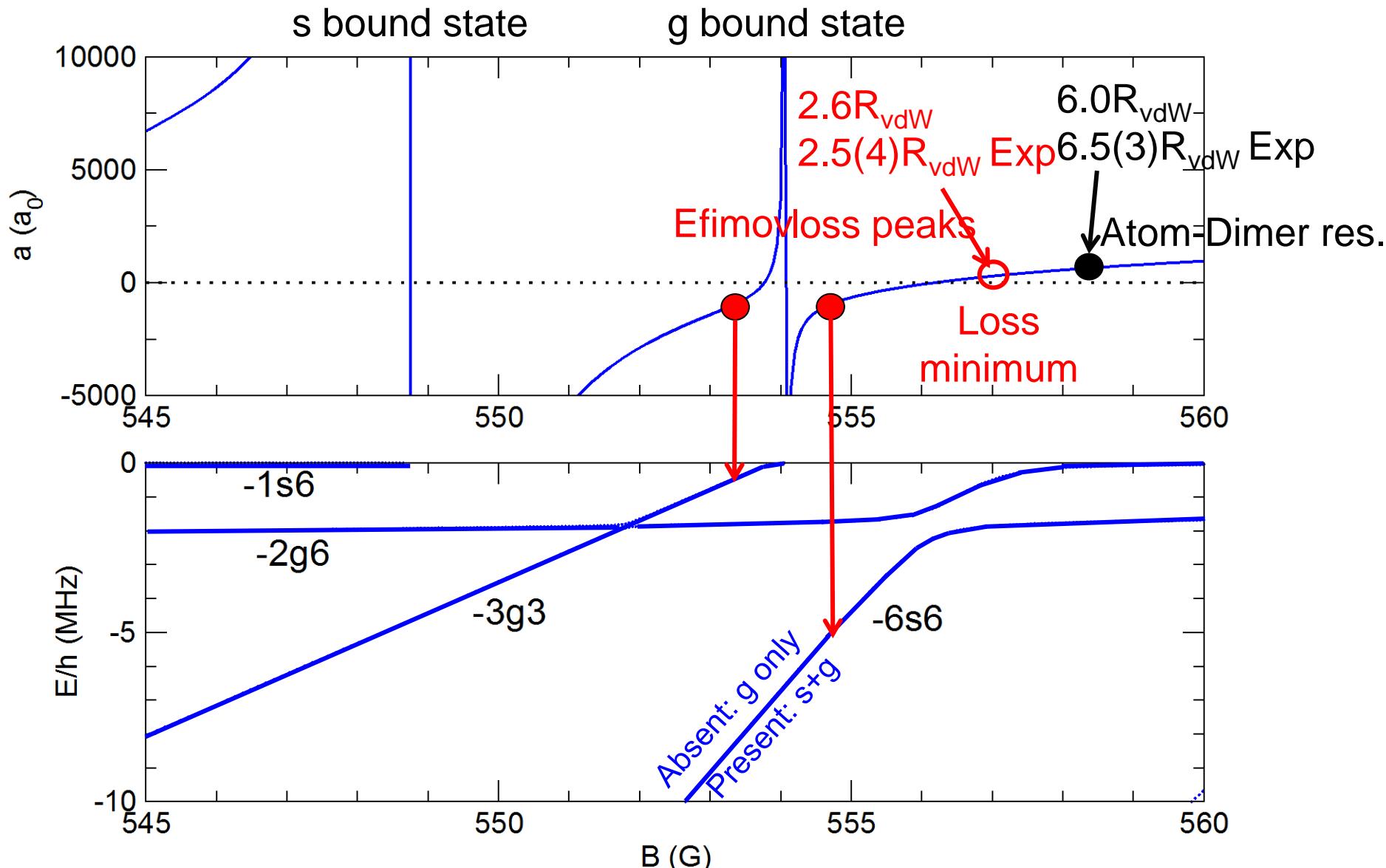
Number of bound states can be varied, $N = 2$ to 4.

Numerically solve 3B equations in hyperspherical basis



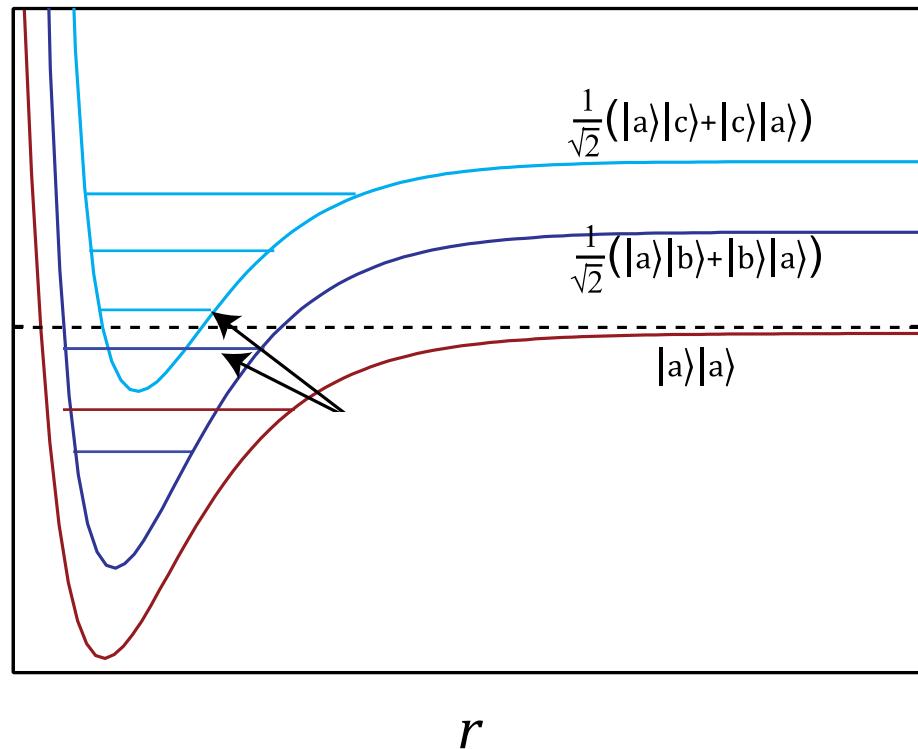
T. Kraemer et al., Nature 440, 315 (2006)





Cs coupled channels model from Berninger, et al, Phys. Rev. A 87, 032517 (2013)
 Cs overlapping resonances (multiple s_{res}): Jachymski, PSJ, PRA 88, 052701(2013)
 Efimov measurements: Innsbruck, Phys. Rev. Lett. 107, 120401 (2011);
 Few-Body Syst. 51, 113–133 (2011); Phys. Rev. A 90, 022704(2014)
 3-body theory: Wang and PSJ, Nat. Phys. 10, 768(2014)

3-channel 2-resonance model for universal van der Waals 3-body physics



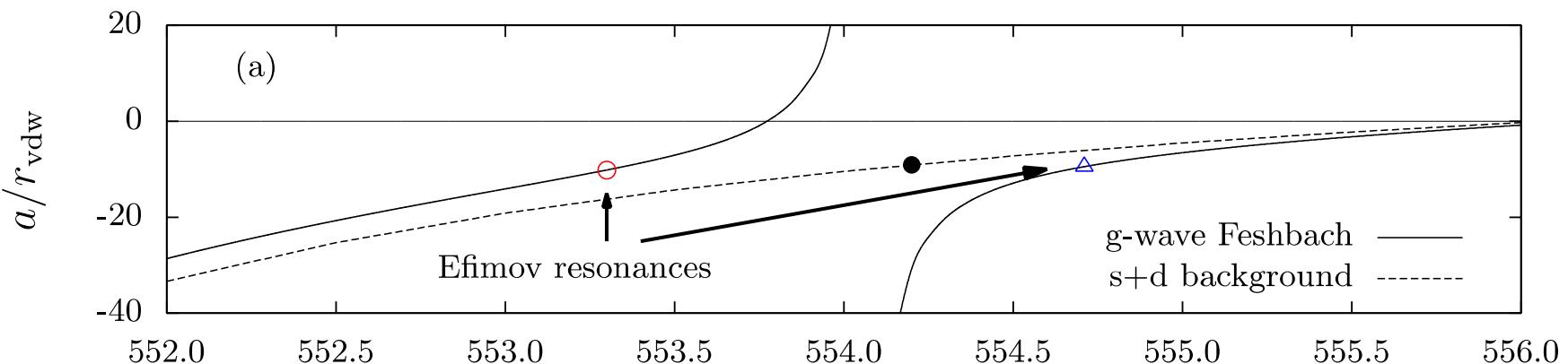
Simultaneously describes two closed channels with overlapping resonances

States of recombination come from BOTH channels.

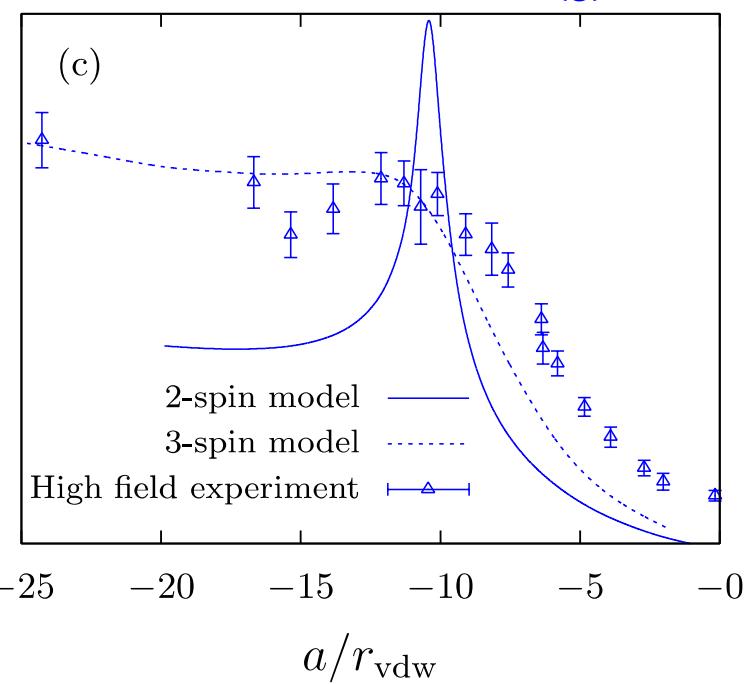
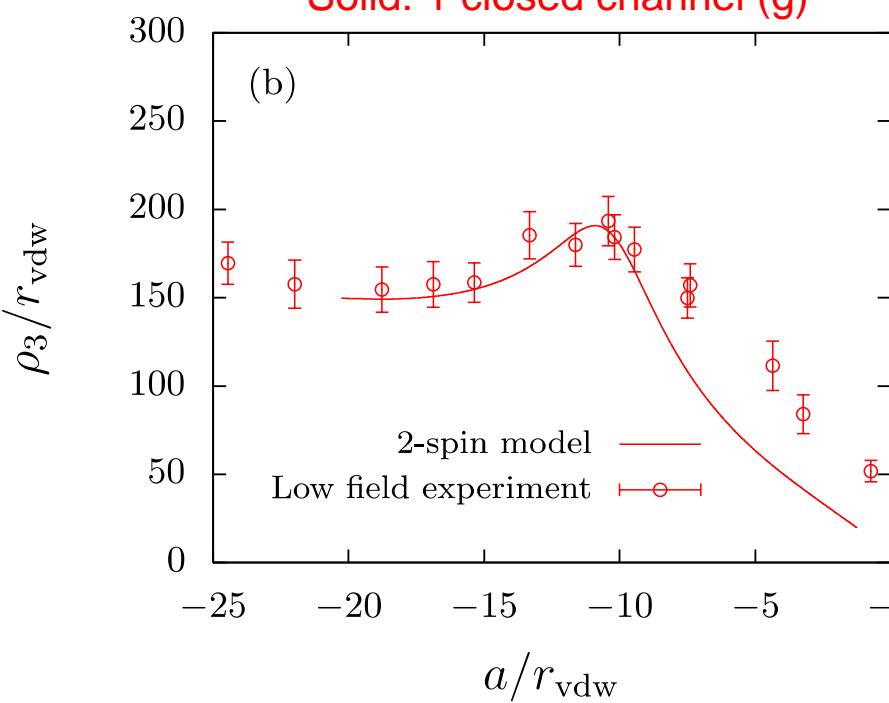
Necessary for universal van der Waals physics near 554G Cs resonance

$$a(B) = a_{\text{bg}} \left(1 - \frac{\Delta_1}{B - B_1}\right) \left(1 - \frac{\Delta_2}{B - B_2}\right)$$

Jachymski, PSJ, PRA 88, 052701(2013)
Y. Wang & PSJ, Nat. Phys. 10, 768 (2014)

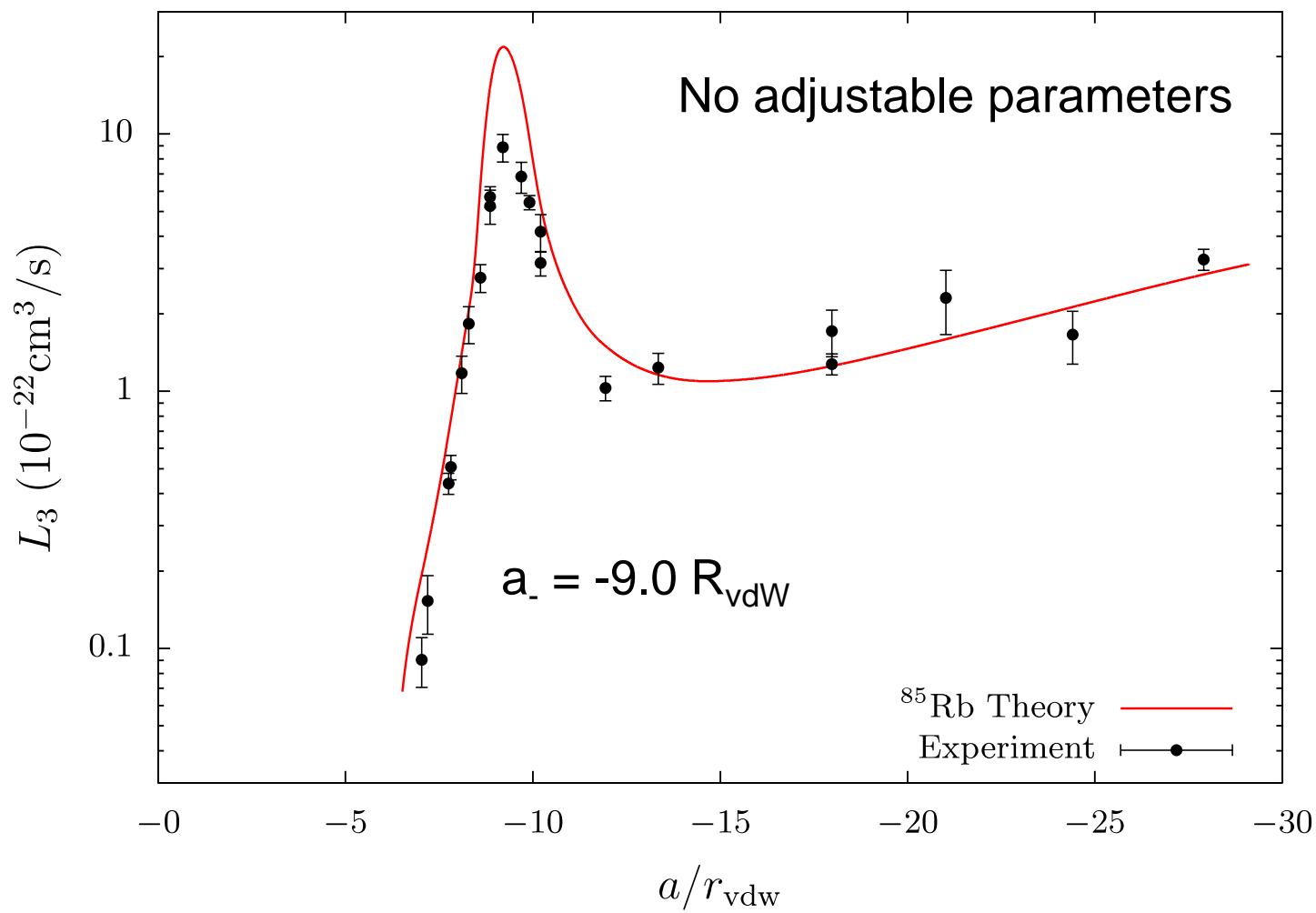


B (Gauss) Dashed: 2 closed channels ($s+g$)
Solid: 1 closed channel (g)



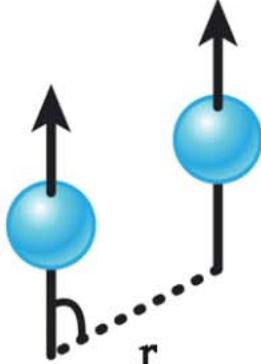
Calculations: no adjustable parameters

2-spin model, ^{85}Rb , $s_{\text{res}} = 28$, $\alpha_{\text{bg}} = -5.4$



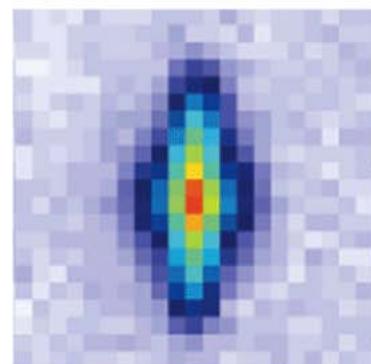
JILA data: Wild, et al., Phys. Rev. Lett. 108, 145305 (2012)

Universality in complex or chaotic collisions

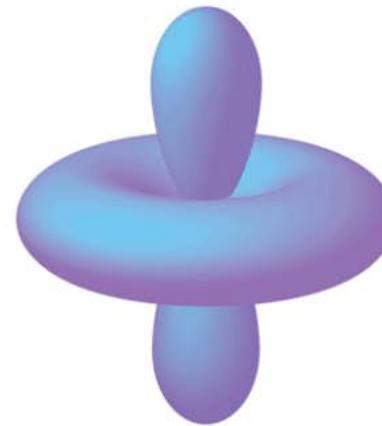


$4f^{12}6s^2\ ^3H_6$
 $7\mu_B$

Ferlaino group
 Innsbruck
 ^{168}Er
 30000 atoms
 order 100nK



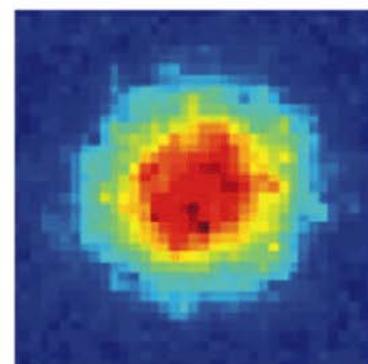
(b)



(a)

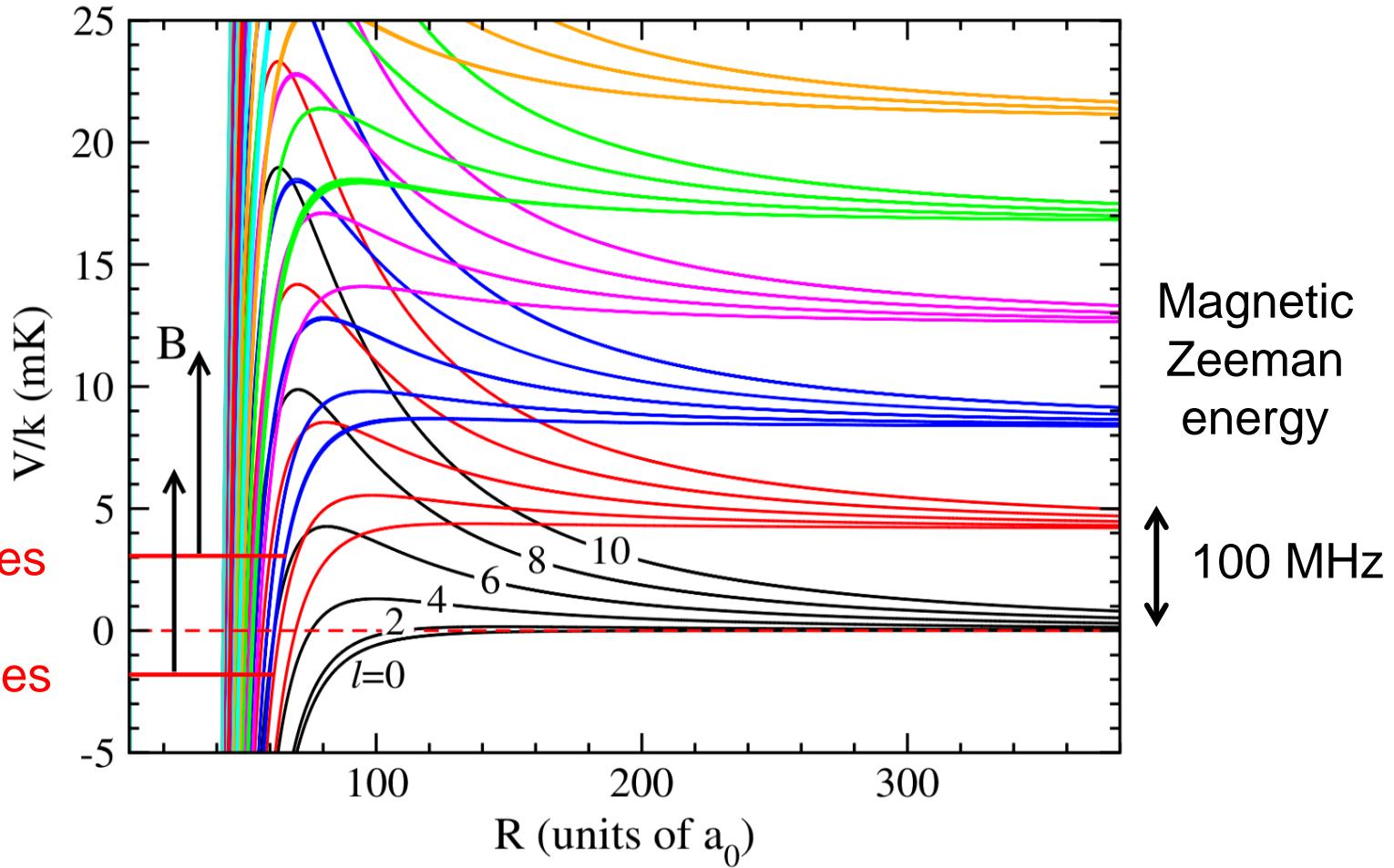
$4f^{10}6s^2\ ^5I_8$
 $10\mu_B$

Lev group
 Stanford
 ^{161}Dy
 6000 atoms
 $64\text{nK } T/T_F=0.2$



(c)

Mixed states
of many
partial waves



Diagonal potential energy curves for $^{164}\text{Dy} + ^{164}\text{Dy}$ at $B = 50\text{G}$

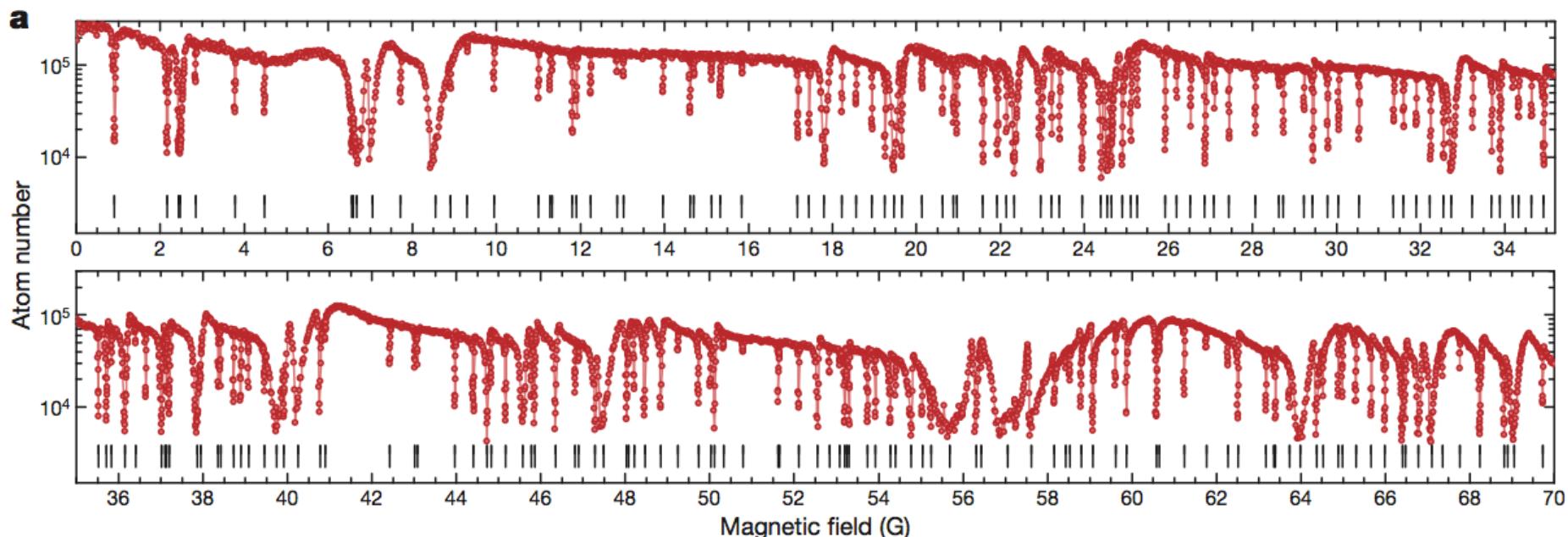
Asymptotic $|(j_1 j_2) j m_j, l m_l\rangle$ channels, $m_j + m_l = -16$, $0 \leq l \leq 10$

From Petrov, Tiesinga, Kotchigova, PRL 109, 103002 (2012)

Quantum chaos in ultracold collisions of gas-phase erbium atoms

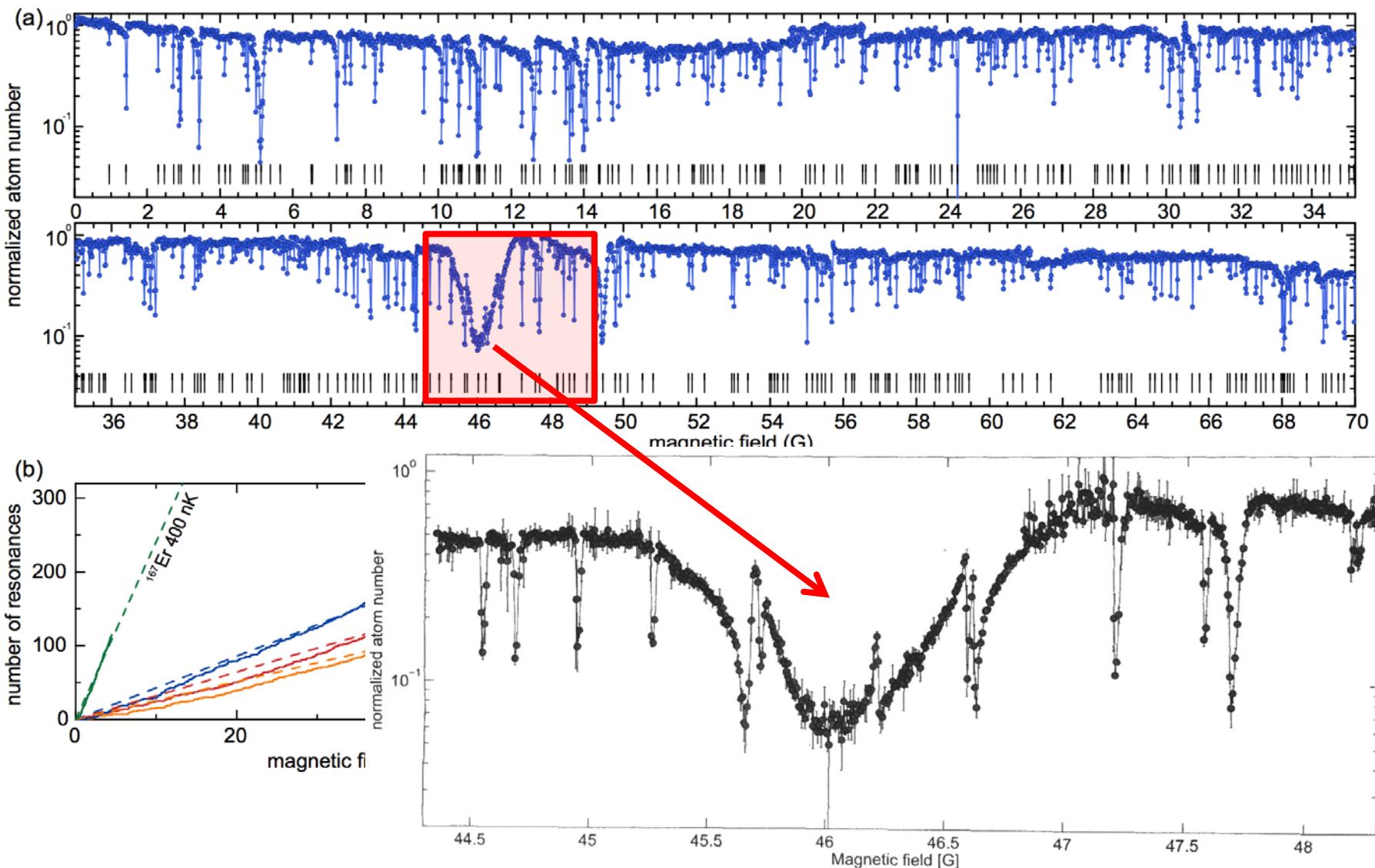
$$\dot{n} = -L_3 n^3$$

Albert Frisch¹, Michael Mark¹, Kiyotaka Aikawa¹, Francesca Ferlaino¹, John L. Bohn², Constantinos Makrides³, Alexander Petrov^{3,4,5} & Svetlana Kotochigova³



^{168}Er in ground state ${}^3\text{H}_6(m=-6)$

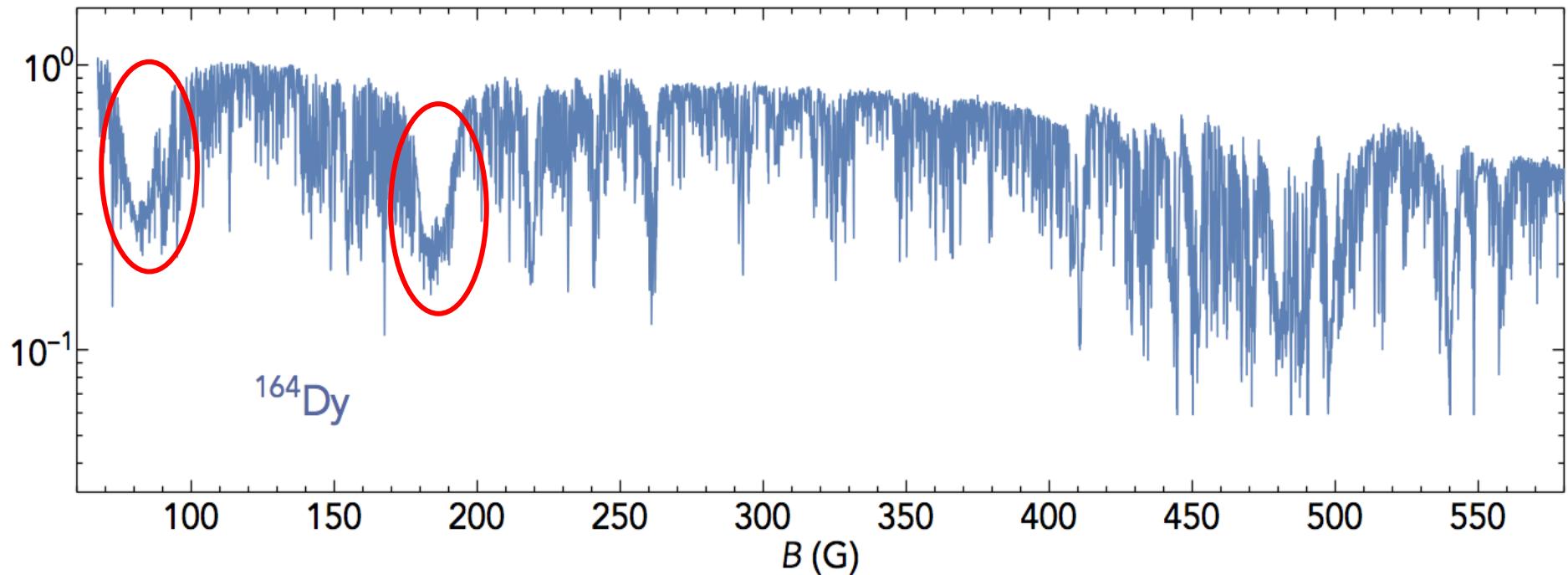
^{164}Dy M=-8



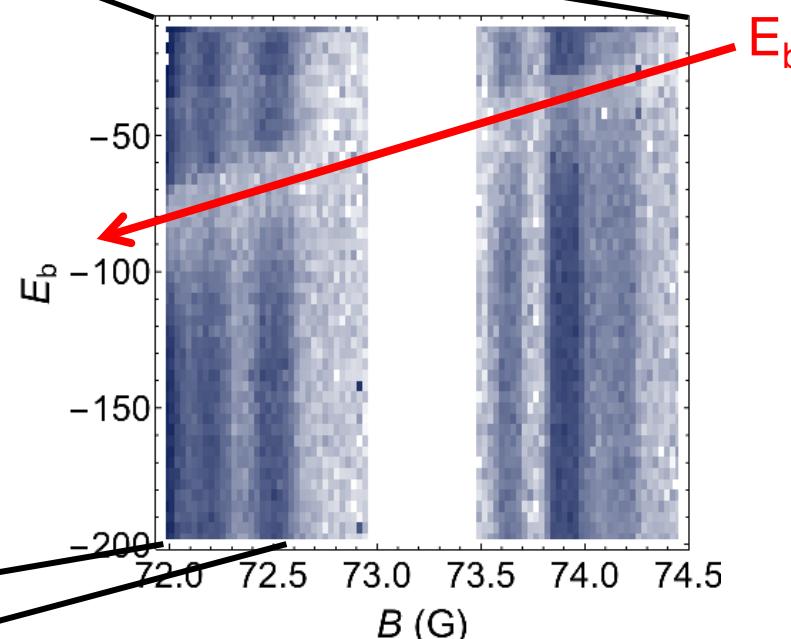
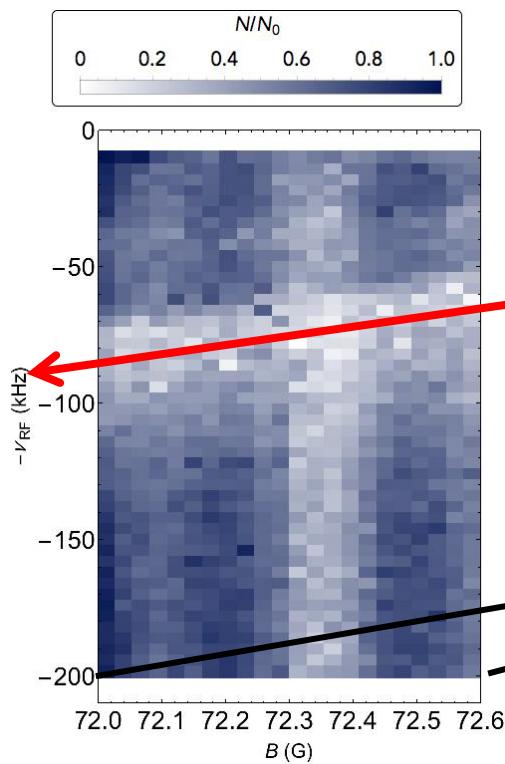
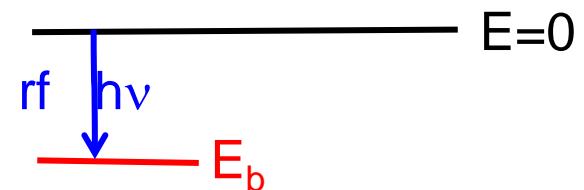
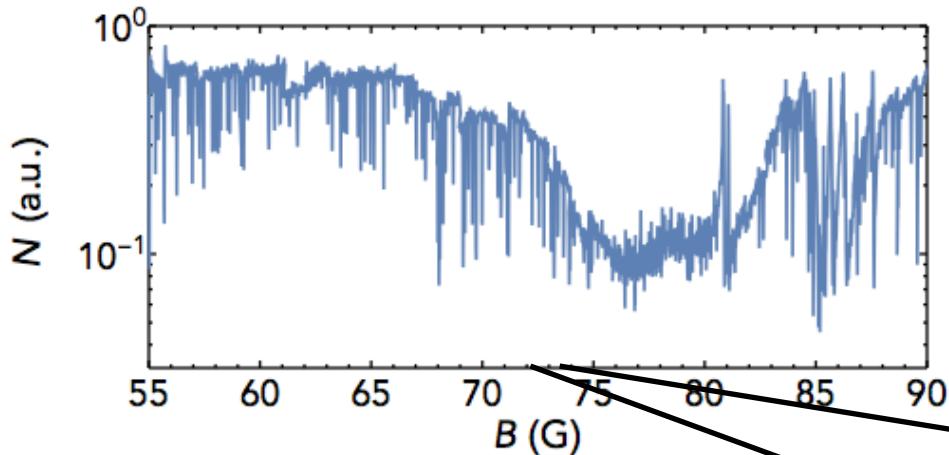
Pfau/Ferlaino/Kotochigova collaboration, arxiv:1506.05221

Feshbach resonances in Dysprosium

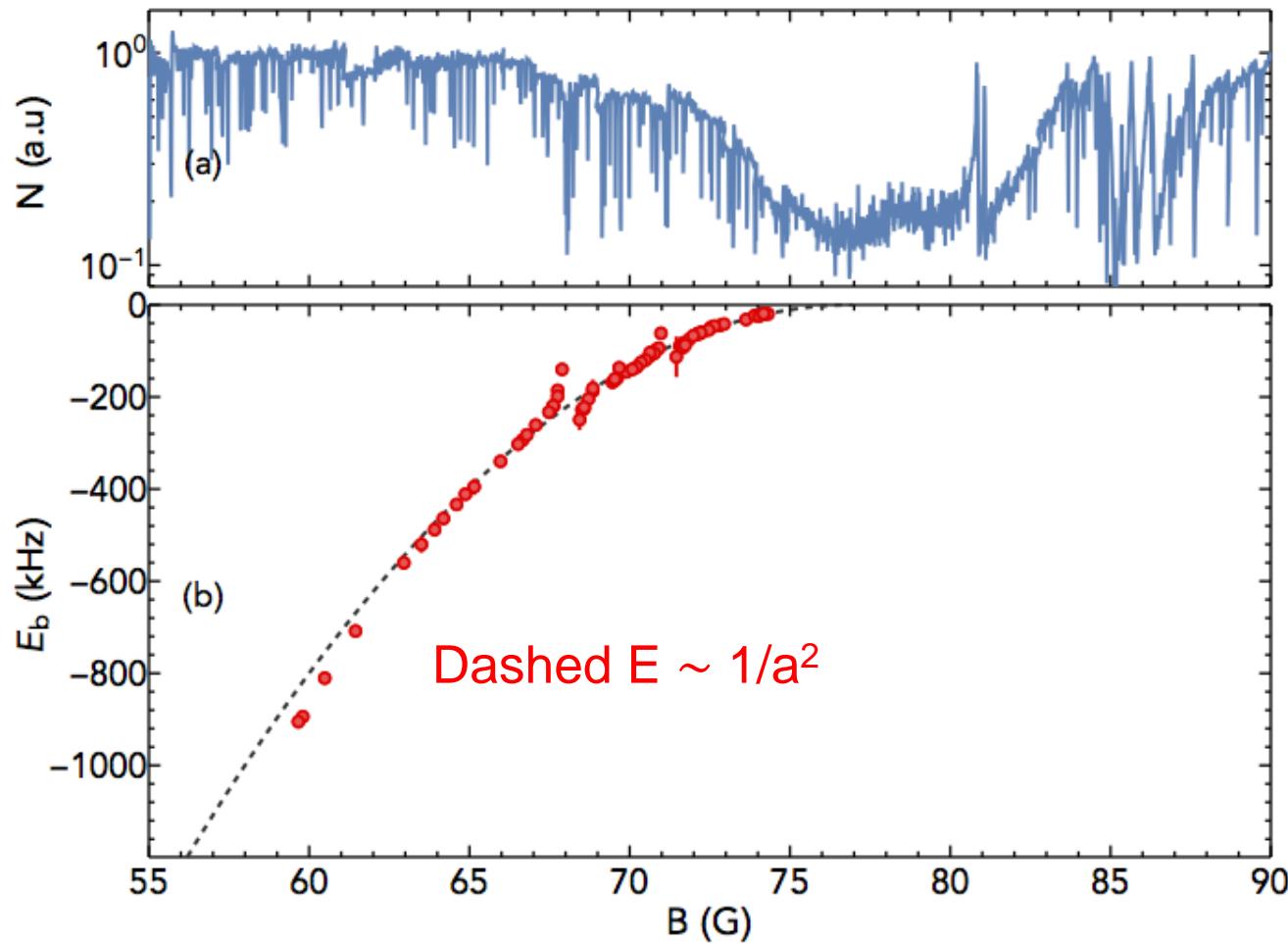
Atom-loss spectroscopy of ^{164}Dy at high field, observation of several broad features



T. Maier, I. Ferrier-Barbut, H. Kadau, M. Schmitt, M. Wenzel, C. Wink, T. Pfau,
K. Jachymski, and PSJ, Phys. Rev. A, 92, 060702 (2015).

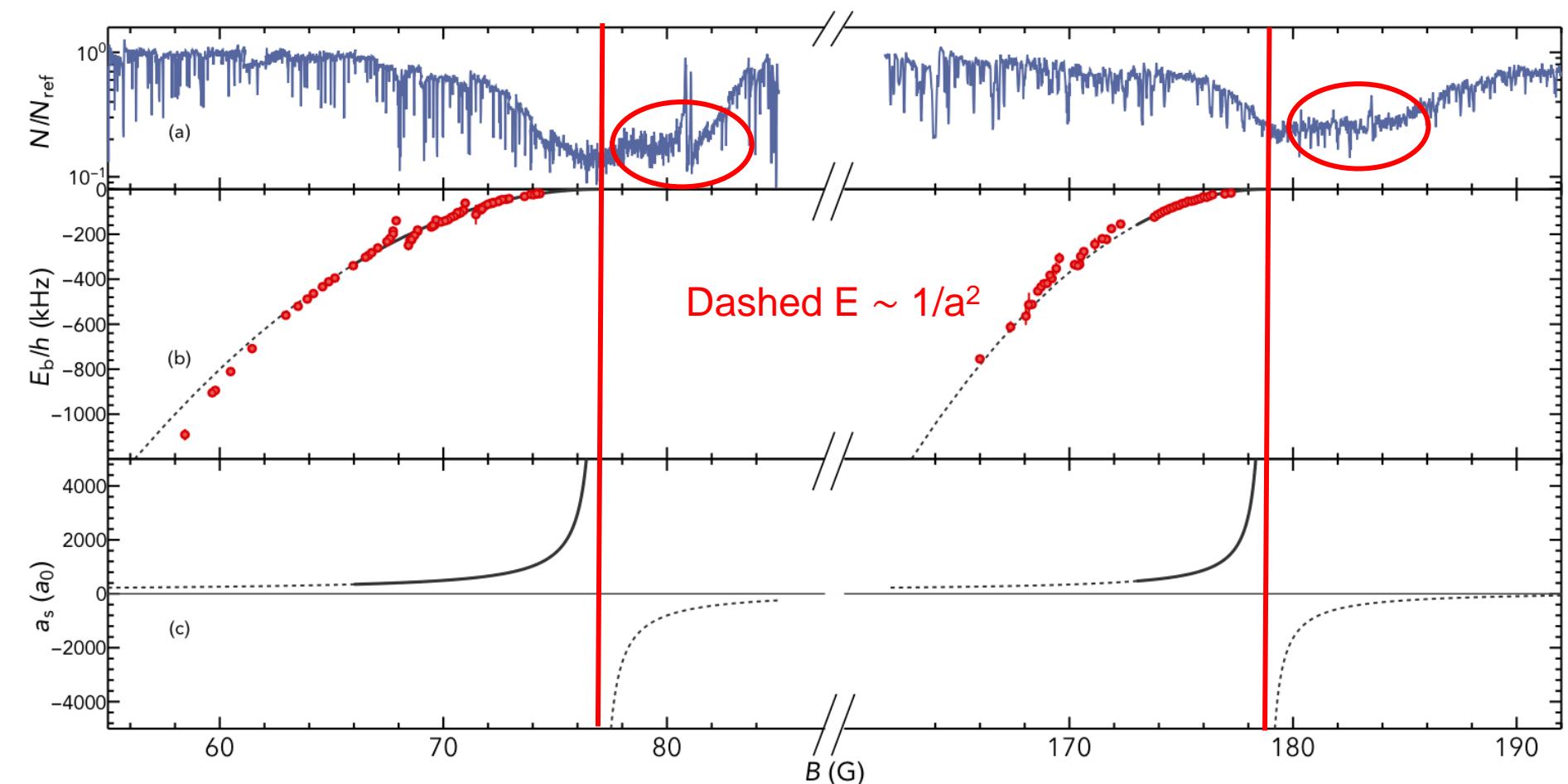


Slide thanks to Igor Ferrier and Tilman Pfau



Signature of “universal” s-wave halo state

Slide thanks to Igor Ferrier and Tilman Pfau



Signature of “universal” s-wave state

Scattering of two structureless dipoles

Bohn, Cavagnero, Ticknor, New. J. Phys. 11, 055039 (2009)

s-wave: $p=6$ van der Waals

d-wave: $p=3$ dipolar + $p=6$ van der Waals

s-d off-diagonal coupling → adiabatic s-wave $p=4$ term

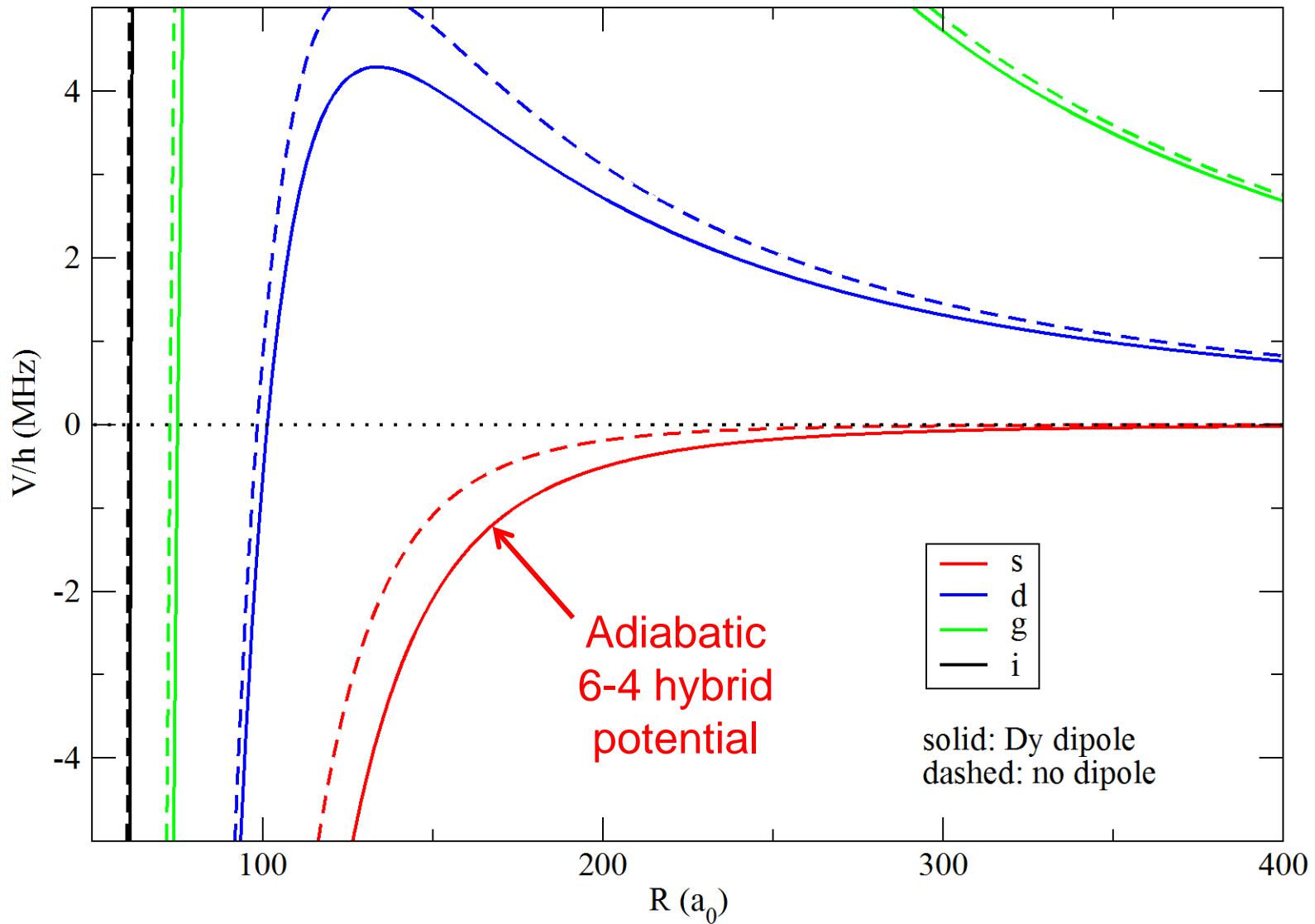
Dy+Dy: $R_6 = 154 a_0$ $E_6/h = 0.93 \text{ MHz}$ or $\approx 50 \mu\text{K}$

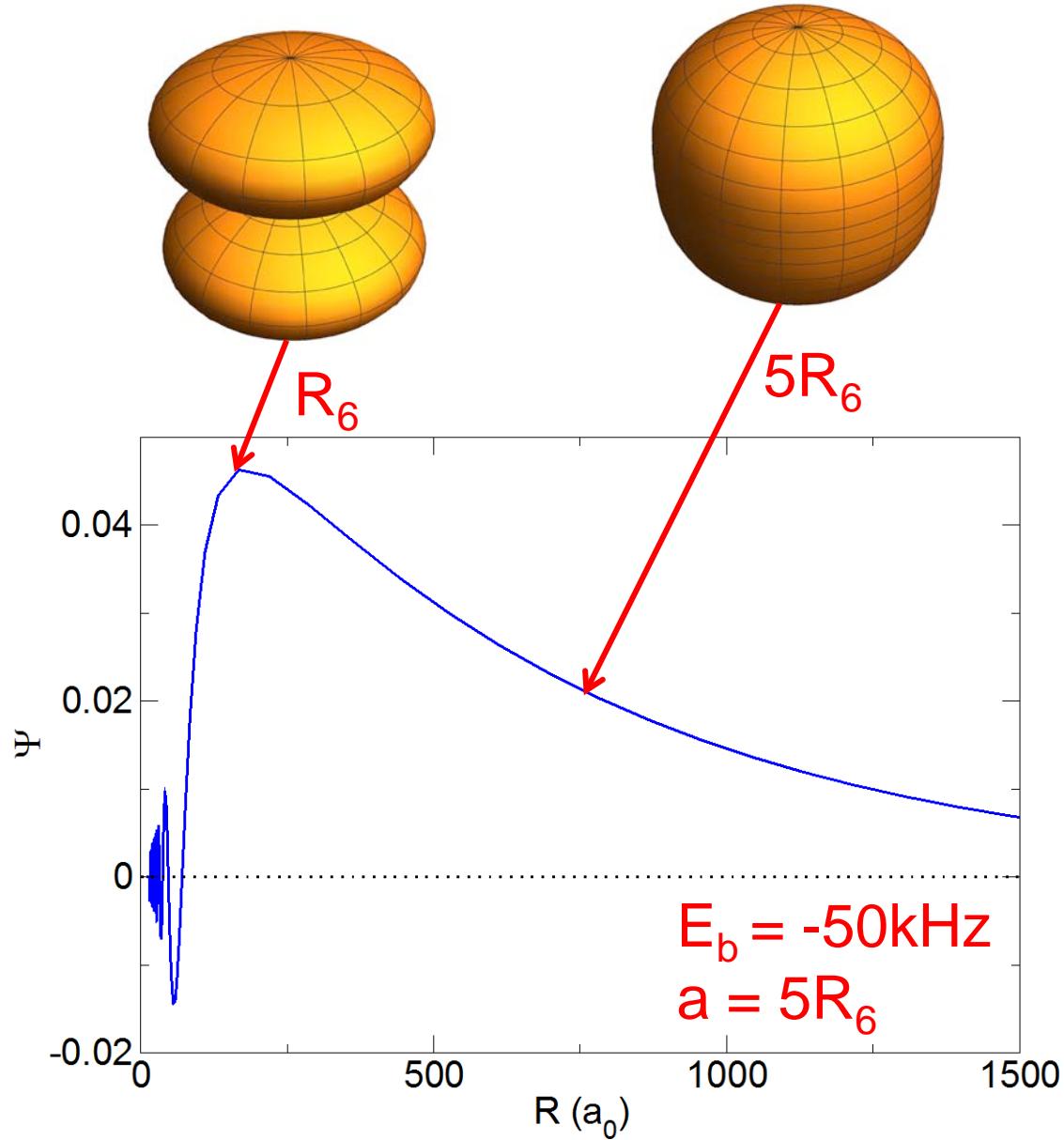
$$R_3 = 194 a_0 \text{ (Bohn definition)}$$
$$R_4 = 143 a_0$$

Solve s-wave (1-channel) scattering for a single potential to get “universal” $E(a)$ function in reduced units of E/E_p and a/R_p .

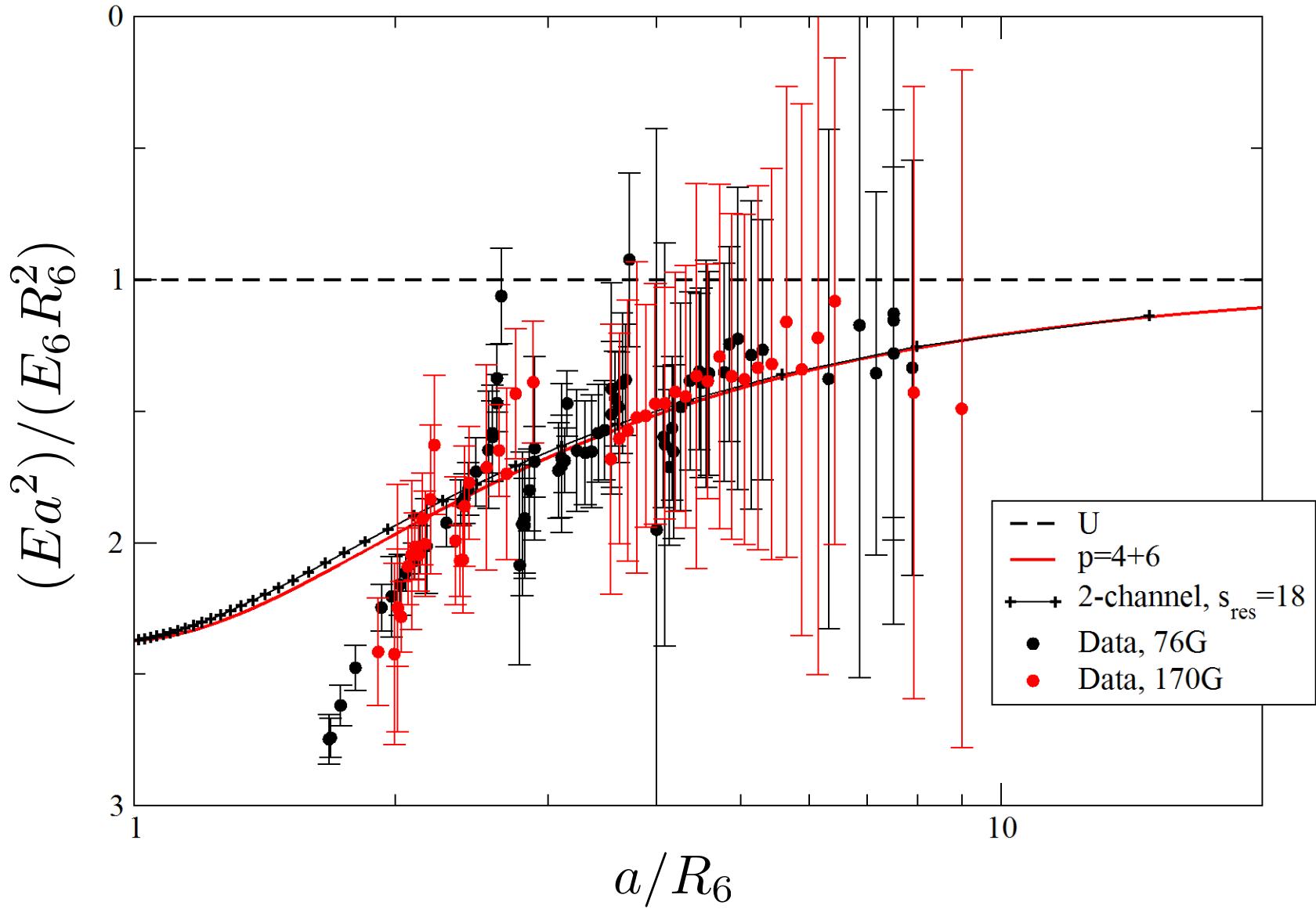
Do this for a $p=6$ potential, or a $p=3$ potential or
a p=6 + p=4 “hybrid” potential
(scale with E_6, a_6).

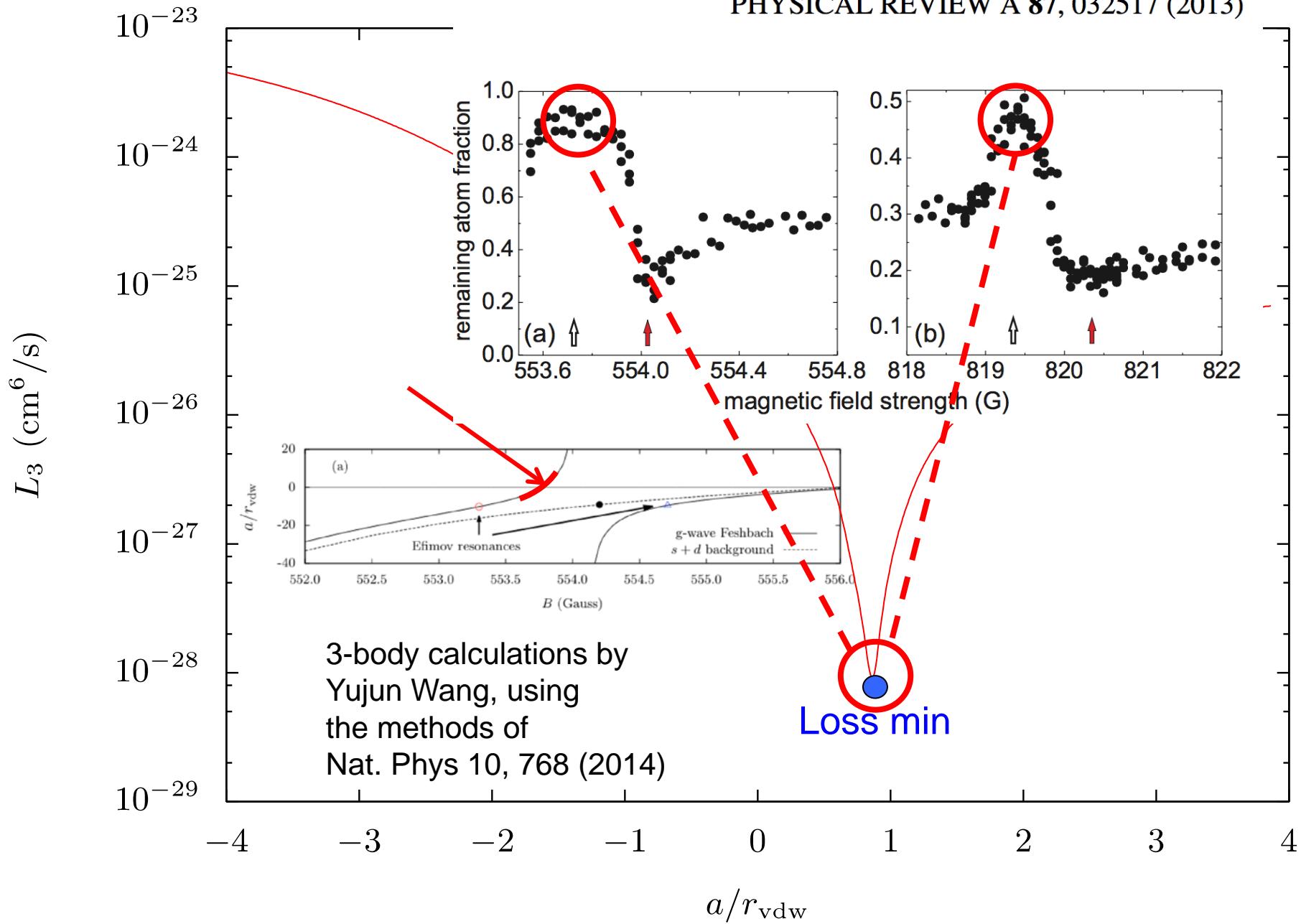
Dy_2 adiabatic potentials

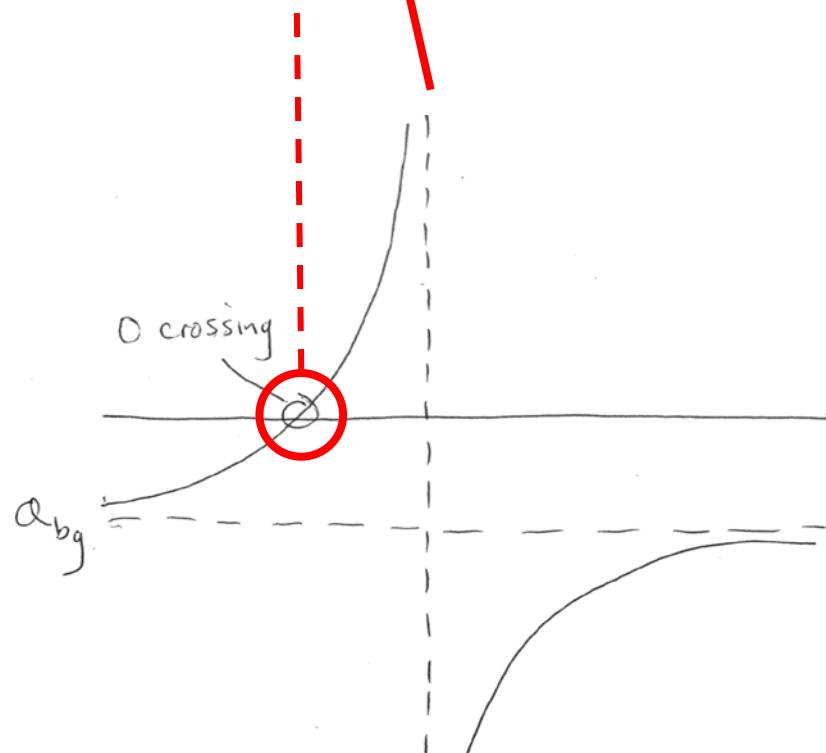
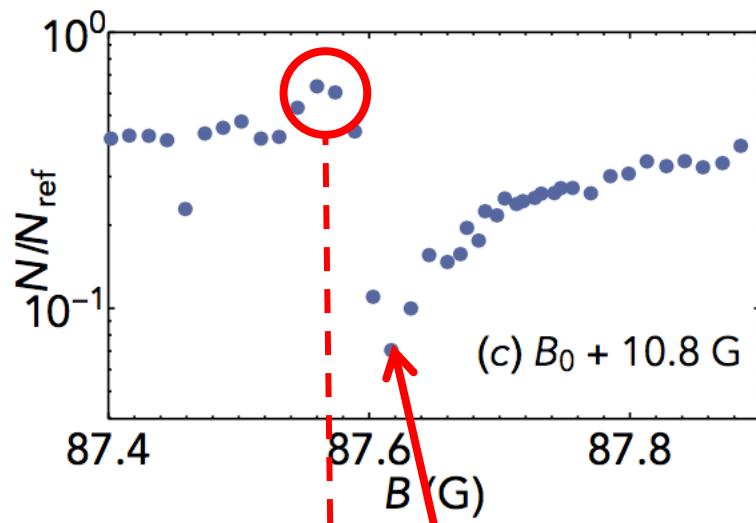
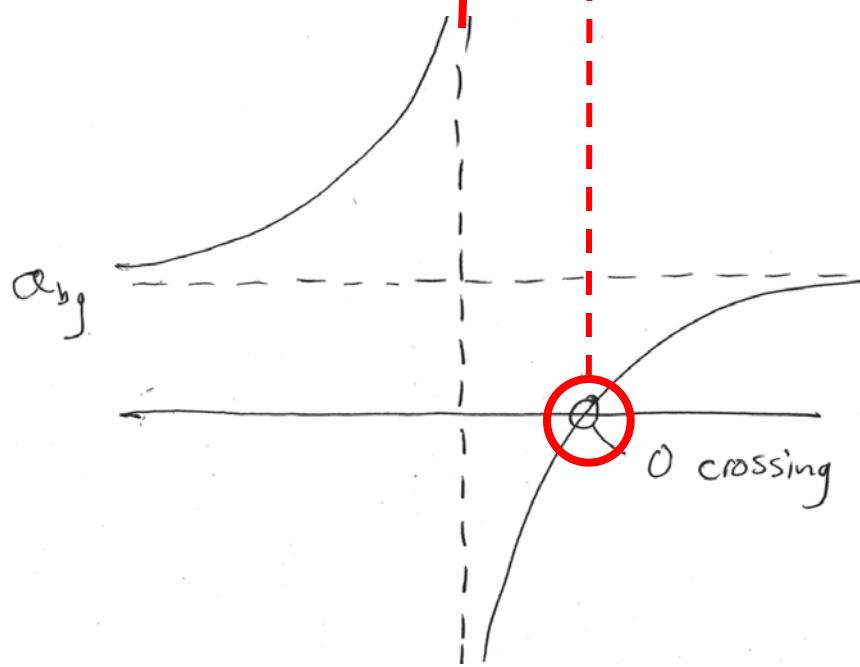
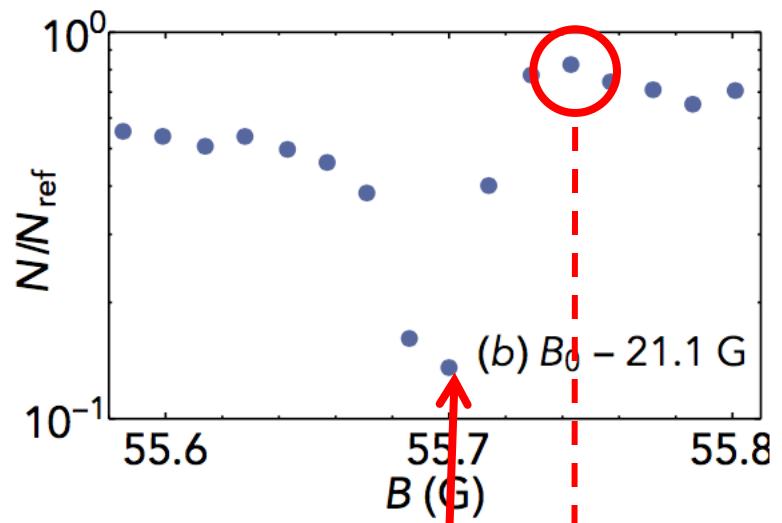




Ea^2 versus a in scaled units

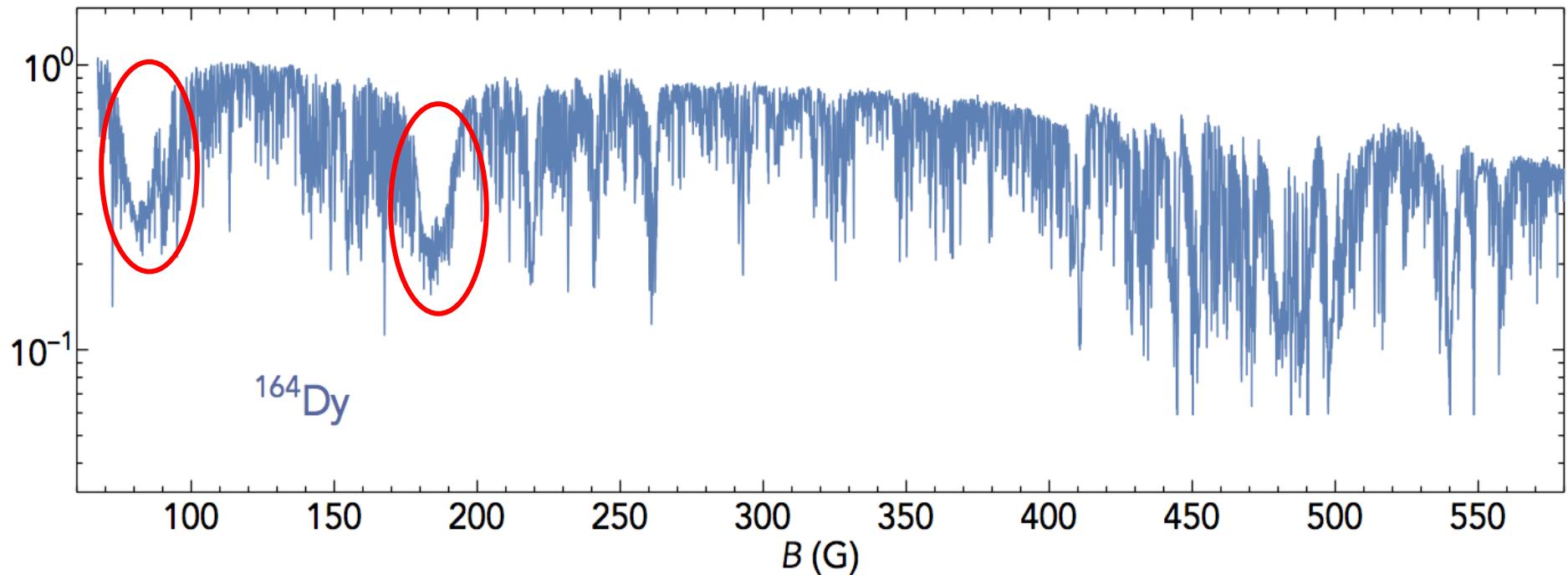






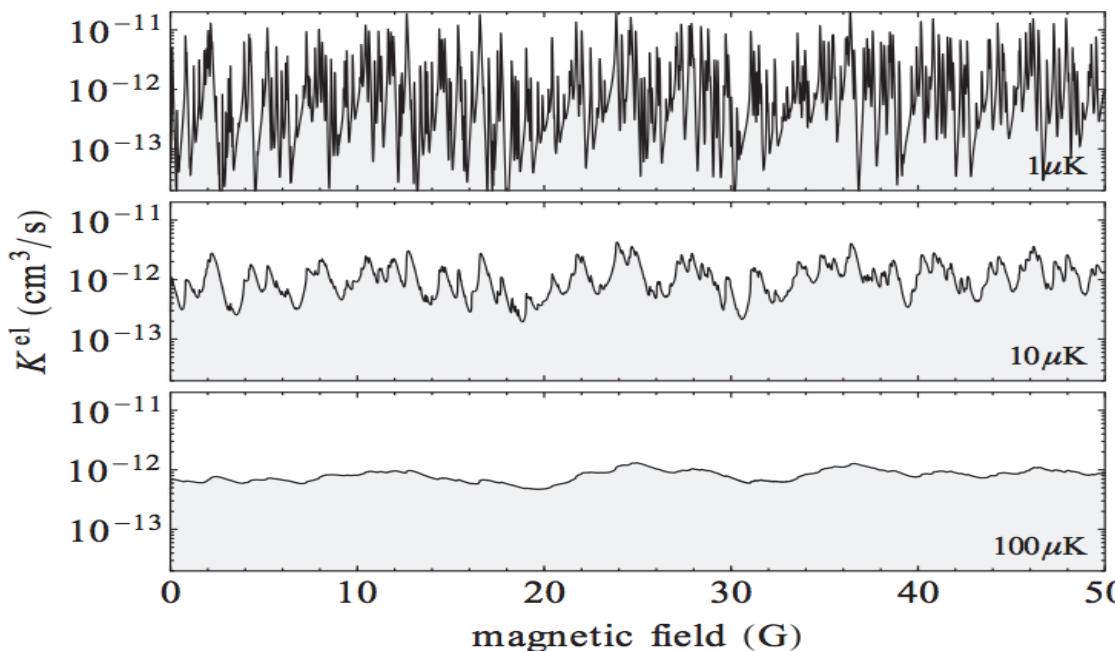
Feshbach resonances in Dysprosium

Atom-loss spectroscopy of ^{164}Dy at high field, observation of several broad features

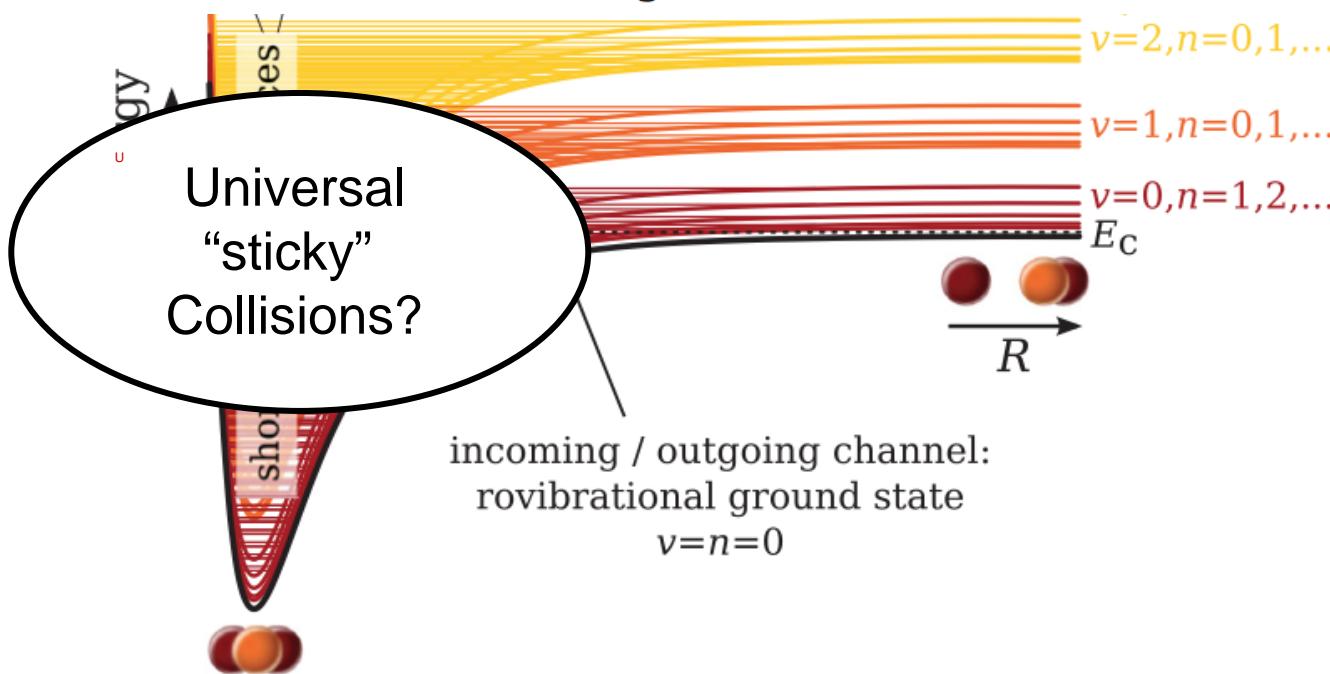


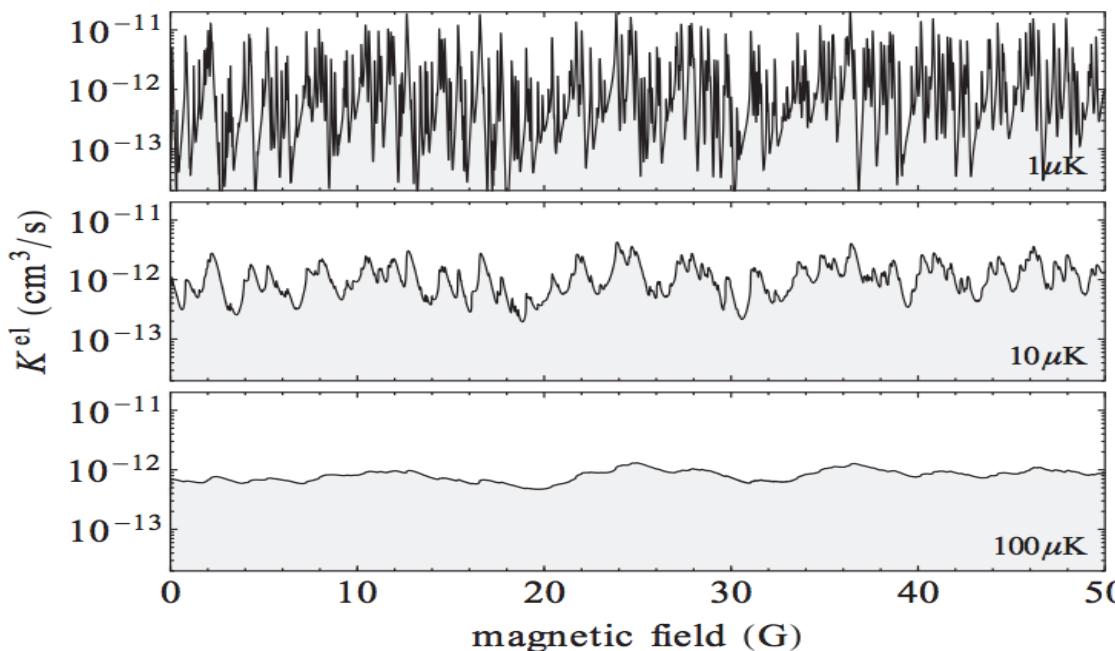
Patterned complexity

T. Maier, I. Ferrier-Barbut, H. Kadau, M. Schmitt, M. Wenzel, C. Wink, T. Pfau,
K. Jachymski, and PSJ, Phys. Rev. A, 92, 060702 (2015).

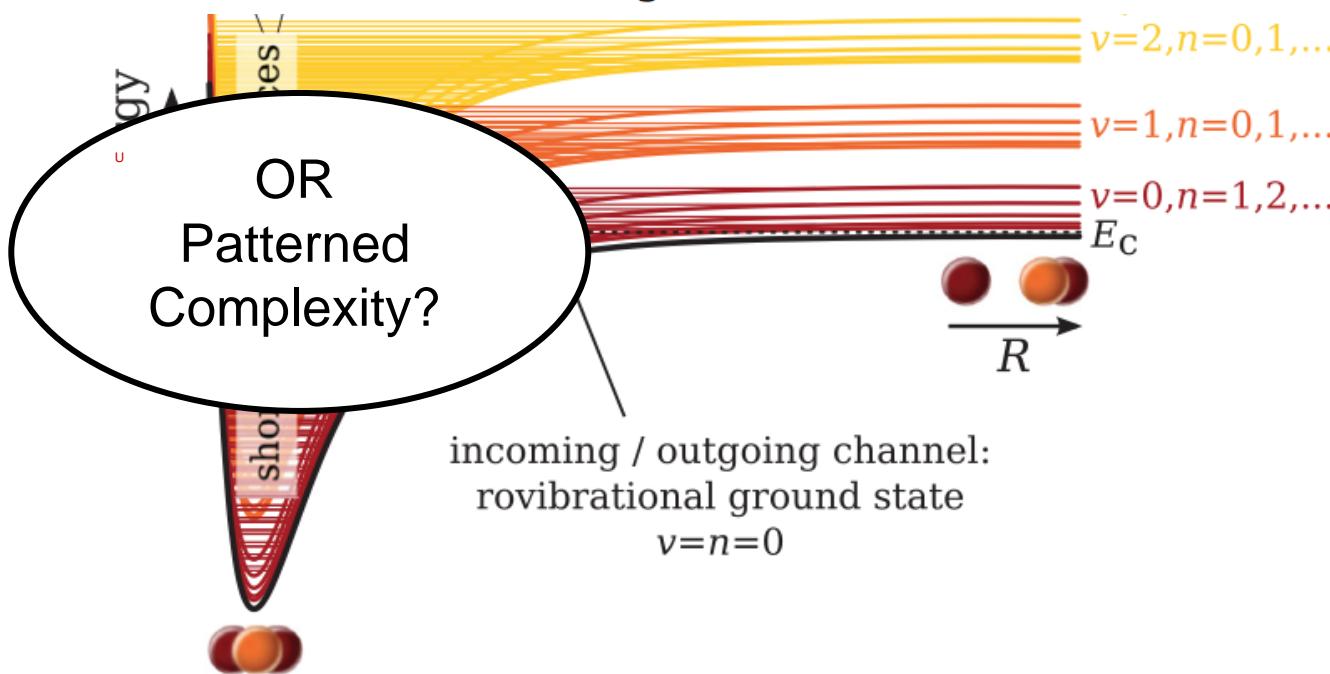


Toy
Statistical
model
 $\text{Rb} + \text{KRb}$





Toy
Statistical
model
 $\text{Rb} + \text{KRb}$



The End