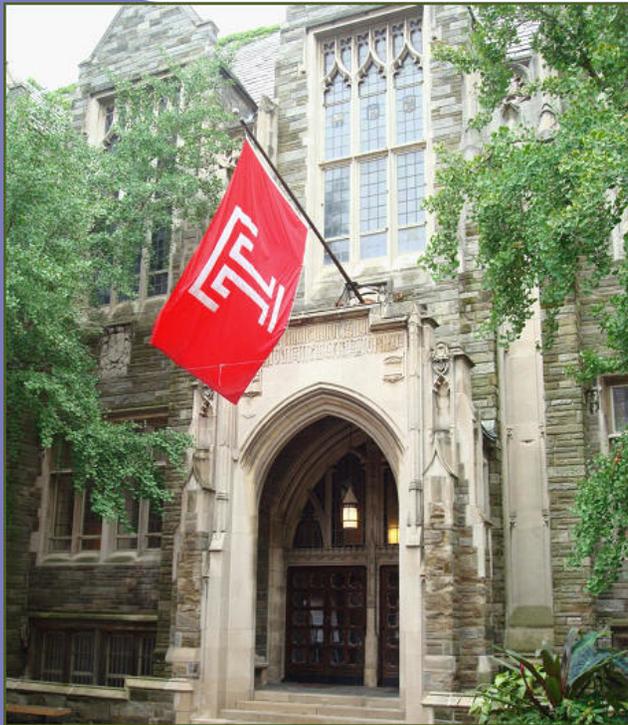


Three-Body Interactions with Magnetic Lanthanides



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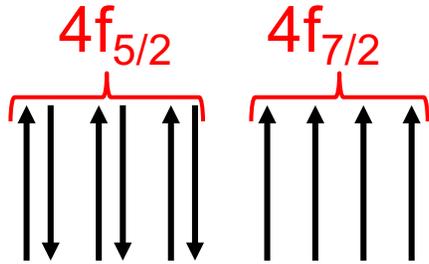
Introduction

- There has been a lot of excitement in using confined ultra-cold atoms and molecules to simulate many-body physics.
- Until recently, this type of study has *mainly* focused on alkali-metal atoms and their di-atomic molecules.
- Ultracold atomic physics is now poised to enter a new regime, where far-more complex atomic species can be cooled and studied.
- Magnetic lanthanide atoms with their large magnetic moment and large orbital angular momentum are extreme examples of such species.

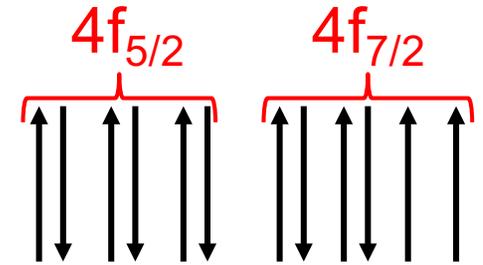
Dy and Er: submerged 4f-shell atoms

Dy ($4f^{10} 6s^2 ({}^5I_8)$)

Er ($4f^{12} 6s^2 ({}^3H_6)$)



$$\mu = 10\mu_B$$



$$\mu = 7\mu_B$$

PERIODIC TABLE
Atomic Properties of the Elements

NIST
National Institute of Standards and Technology
Technology Administration, U.S. Department of Commerce

Physics Laboratory
Standard Reference Data Group
physics.nist.gov
www.nist.gov/srd

Frequently used fundamental physical constants
For the most accurate values of these and other constants, visit physics.nist.gov/constants
1 second = 9192 631 770 periods of radiation corresponding to the transition between the two hyperfine levels of the ground state of ${}^{133}\text{Cs}$

speed of light in vacuum c 299 792 458 m s⁻¹ (exact)
Planck constant h 6.626 070 15 × 10⁻³⁴ J s (exact) ($h = h/2\pi$)
Boltzmann constant k 1.380 658 × 10⁻²³ J K⁻¹

Solids
 Liquids
 Gases
 Artificially Prepared

1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	
1	H	He															He	
2	Li	Be															Ne	
3	Na	Mg															Ar	
4	K	Ca	Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ga	Ge	As	Se	Br	Kr
5	Rb	Sr	Y	Zr	Nb	Mo	Tc	Ru	Rh	Pd	Ag	Cd	In	Sn	Sb	Te	I	Xe
6	Cs	Ba															Rn	
7	Fr	Ra																
			Ac	Th	Pa	U	Np	Pu	Am	Cm	Bk	Cf	Es	Fm	Md	No	Lr	

Based upon ${}^{12}\text{C}$. () indicates the mass number of the most stable isotope. For a description of the data, visit physics.nist.gov/data

NIST SP 966 (September 2003)

Some of the electrons in the 4f shell are aligned in the same direction creating a large magnetic moment and a large spin.

First experiments with ultracold Dy and Er

The first realization of a quantum-degenerate gas of **Er** was achieved by F. Ferlaino's group.

*Aikawa et al. (PRL **108**, 210401 (2012))*

The Stanford group of B. Lev cooled **Dy** atoms to ultracold temperatures and also created BECs.

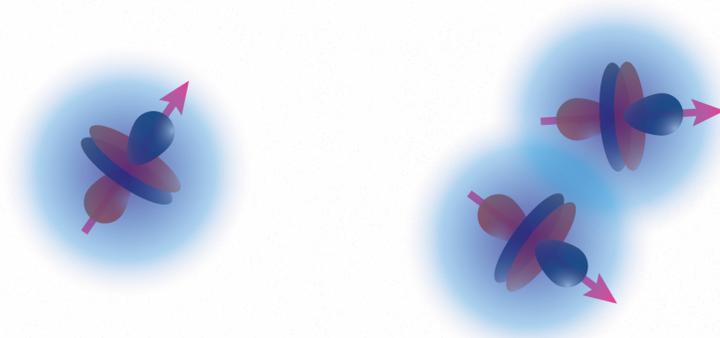
*Lu et al. (PRL **107**, 190401 (2011))*

A growing number of other experimental groups started to work with these atoms.

Detection of resonances by three-body recombination

- Ultracold **bosonic Er** and **Dy** atoms are confined in an optical dipole trap and are in their energetically-lowest Zeeman sublevel.
- Inelastic **three-body recombination** causes atom loss from the trap.

Three colliding atoms form a weakly-bound dimer and atom



- At resonance, **the recombination process is enhanced** due to the coupling between the open and closed channel leading to a resonant increase of the atomic loss.
- We identify the magnetic field values with the maximum loss as the positions of Fano-Feshbach resonances.

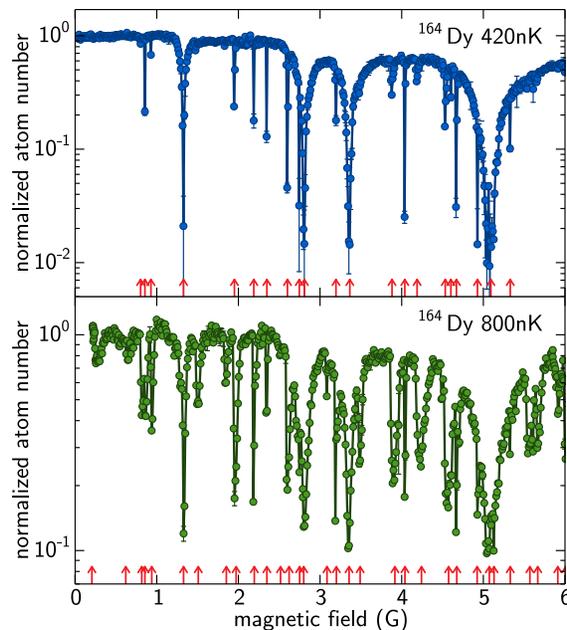
Extreme temperature sensitivity

We studied the extreme sensitivity of three-body recombination rate to the temperature of the atoms.

In magnetic lanthanides this phenomenon was first observed by Lev's group in Dy.

Phys. Rev. A 89, 020701 (2014)

Atom-loss spectrum for bosonic ^{164}Dy



← 3 res. per G at 420 nK

← 5 res. per G at 800 nK

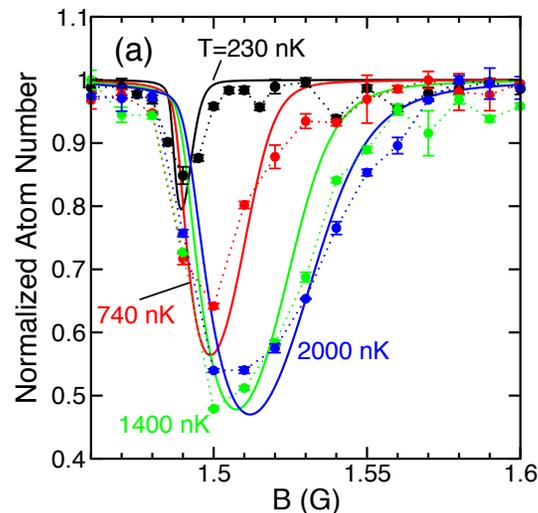
Number of resonances is increased by ~50 %

Trimer model

In atom-loss spectrum of ^{168}Er the increase of resonance number with temperature was only 25%.

In addition, the new atom-loss features show a dramatic broadening and a shift of maximum loss to larger B fields with increasing T.

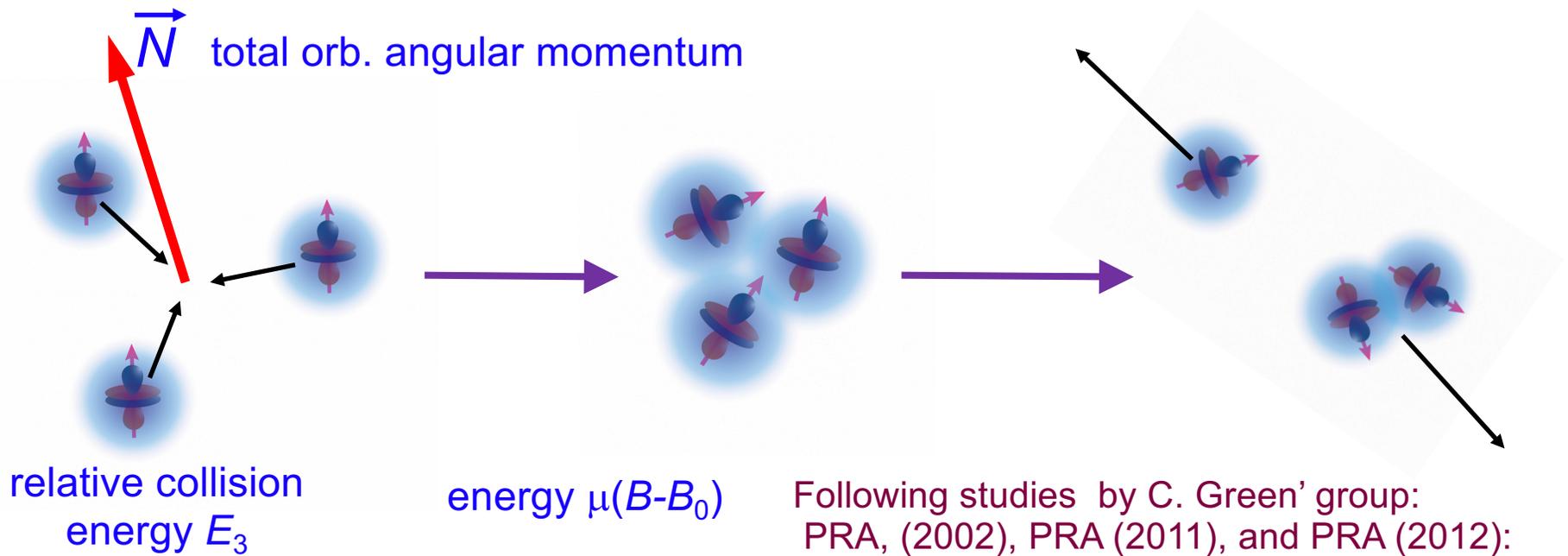
Resonances that already exist at the lowest T do not show these behaviors.



- We describe “trimer” model that suggests that this T dependence is due to scattering processes in the **d-wave entrance channel of the trimer**. (total orbital angular momentum $N=2$ of the trimer)
- This is despite of the fact that the **two-body d-wave centrifugal barrier** ($250 \mu\text{K}$) is a 100x larger than our highest T.

d-wave Feshbach resonances were detected in loss spectrum of **Cr**.
Laburthe-Tolra group (PRA 79, 032706 (2009))

Trimer model



$$\Gamma(E_3) \propto E_3^{N+2}$$

$$\Gamma_{\text{br}} \text{ constant for small } E_3$$

- Lineshape of the resonance

$$|S(E_3, B)|^2 = \frac{\Gamma(E_3)\Gamma_{\text{br}}}{(E_3 - \mu(B - B_0))^2 + (\Gamma_{\text{tot}}(E_3)/2)^2}$$

Trimer model

- Three-body recombination rate coefficient is

$$L_3(E_3, B) \propto v_3 \frac{\pi^2}{k_3^5} |S(E_3, B)|^2 \quad (\text{cm}^6/\text{s})$$

relative velocity

relative wavenumber

- Thermal averaging of rate coefficients leads to additional broadening

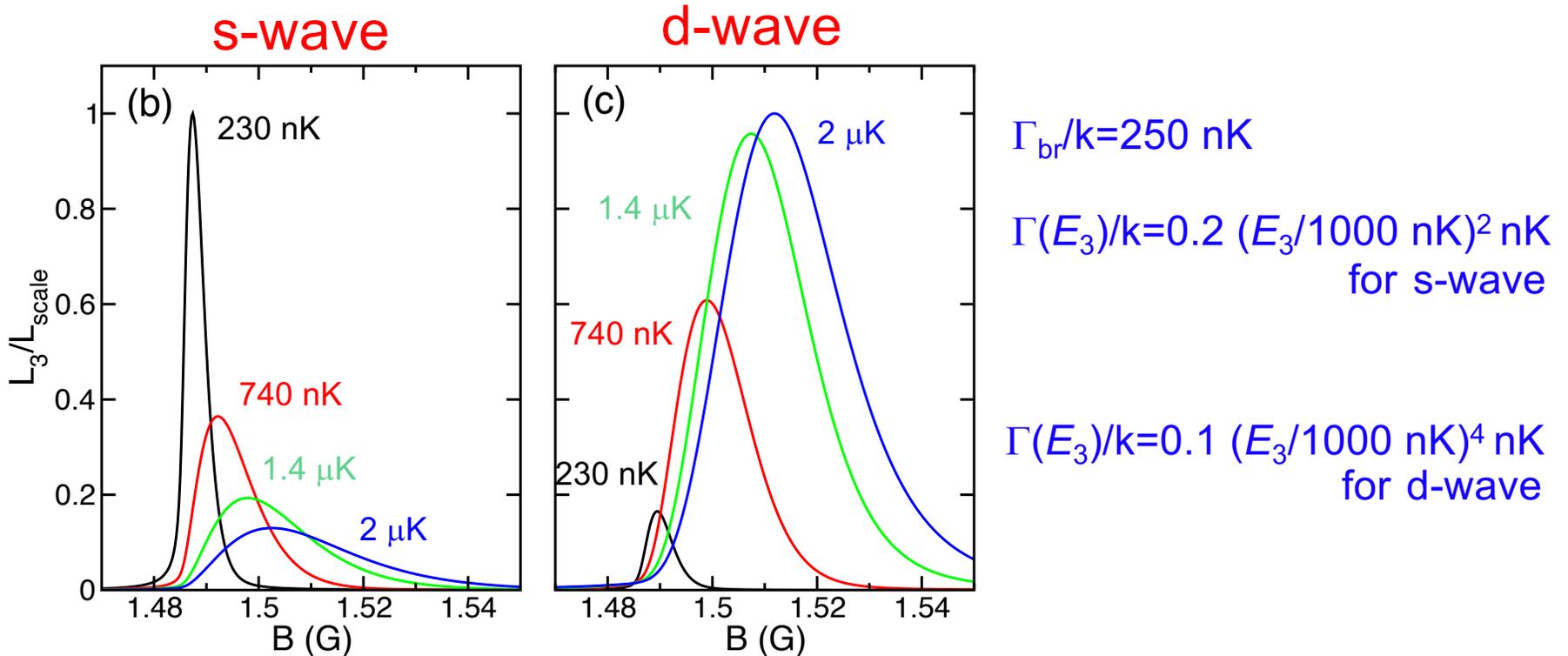
$$L_3(T, B) = \frac{1}{Z} \int_0^\infty E^2 dE L_3(E, B) e^{-E/kT}$$

- For $kT \gg \Gamma(E_3), \Gamma_{\text{br}}$ we derive $L_3(T, B) \propto (kT)^{N-1}$

Trimer simulation

- Loss rate coeff. L_3 for the **N=0 (s-) and N=2 (d-) wave** entrance channel for several temperatures.

scaled in relative units



- A striking difference is the temperature-dependence of the s- and d-wave entrance-channel resonances. The d-wave case agrees with experiment.

Conclusion

- I described the extreme sensitivity to the temperature of the atom-loss spectra and three-body recombination in scattering of magnetic lanthanides.
- Here, I shown that entrance channels with **zero** and **non-zero** relative orbital angular momentum N lead to line shapes with different temperature behavior.
- Those with $N=0$ (s-wave) entrance channels have sharply decreasing recombination rates with T , whereas those with $N=2$ (d-wave) entrance channels have an increasing recombination rate.

Our group members



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Eite Tiesinga
NIST, JQI, and QuICS