

Quench physics with few-body atomic systems in an optical lattice

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Background

- Prepare the superfluid ground-state of ultra-cold atoms in a 2D or 3D optical lattice.
- On average a few-atoms per unit cell ($\langle n \rangle = 2.5$.)
- “Suddenly”, turn off tunneling between unit cells. Then let atoms evolve and interact.

- In each site we end up with a few-body state of the kind

$$|0\rangle + e^{-iE_1 t} |1\rangle + e^{-iE_2 t} |2\rangle + e^{-iE_3 t} |3\rangle + \dots$$

think coherent state

missing amplitudes

- Interferometer can measure differences of (multiple) energies E_n . (Momentum distribution.)



Bosonic examples

- Paul, Johnson and Tiesinga,
Phys. Rev. A **93**, 043616 (2016).
- Paul and Tiesinga,
Phys. Rev. A **92**, 023602 (2015).

Fermionic example

- Nuske, Mathey and Tiesinga,
Phys. Rev. A **94**, 023607 (2016).

Bosonic example

- Atoms in a cubic 3D lattice with weak interactions
- Quench prepares in each site the state

$$|0\rangle + e^{-iE_1 t}|1\rangle + e^{-iE_2 t}|2\rangle + e^{-iE_3 t}|3\rangle + \dots$$

- In 2009 we studied the elastic three-body contribution!

See NJP **11**, 093022 (2009)

E_1 ← single-particle energy

$$E_2 = 2E_1 + U_2$$

$$E_3 = 3E_1 + 3U_2 + U_3$$

Three-body contribution!

Observed by S. Will, ... I. Bloch, Nature **465**, 197 (2010)

Three-body strength

- Imagine “weak” delta-function

$$g \delta(\mathbf{r}_1 - \mathbf{r}_2) (\partial/\partial \mathbf{r}) \mathbf{r}$$

don't mind the units

interaction between atoms. Then

$$U_2 \propto g \ll \Delta E_1$$

energy scale of trap

$$U_3 = c_0 \frac{(U_2)^2}{\Delta E_1}$$

a constant (known for a harmonic trap)

- Can we create situation where $U_2 \ll U_3$?
- As $U_3 \sim (U_2)^2$, the answer seems to be no

Zero-point energy

- We realize that collisions do not occur at zero collision energy and

$$(g + V \nabla_{\text{rel.}}^2 + \dots) \delta(\mathbf{r}_1 - \mathbf{r}_2)$$

collision-energy dependent
effective-range correction

- Now we find

$$U_2 \propto (g + V \Delta E_1)(1 + \dots) \rightarrow 0$$

when $V \cdot \Delta E_1 = -g$

and still
$$U_3 = d_0 \frac{g^2}{\Delta E_1}$$

Based on Phys. Rev. A **93**, 043616 (2016)

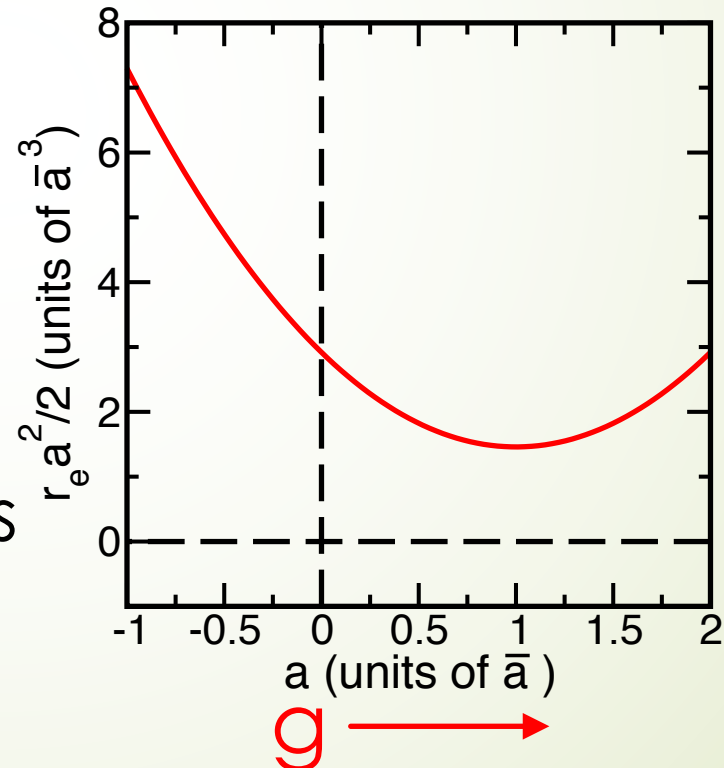
How realistic is this?

- Spin-less atoms like those in second column, ytterbium, or even alkali-metal atoms well away from resonances.
- van-der-Waals physics gives relationship between g and V .

- parabola !
- V is positive!

- The only valid case is ^{88}Sr with “small” $g < 0$.

$$a = -2a_0$$



Feshbach resonances

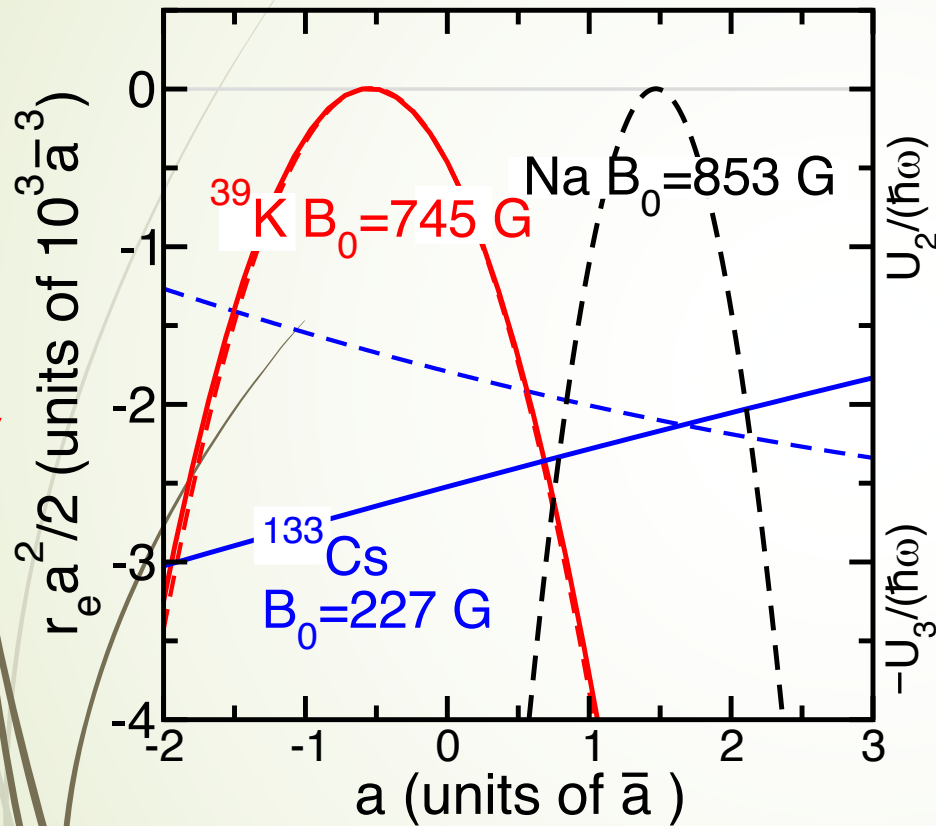
- near resonances the relationship between V and g is more favorable.
- Still a (inverted) parabola, but also $V < 0$ and “large”.
- We can now satisfy the relationship for “reasonable” $g > 0$.

Scattering lengths on order of van der Waals length

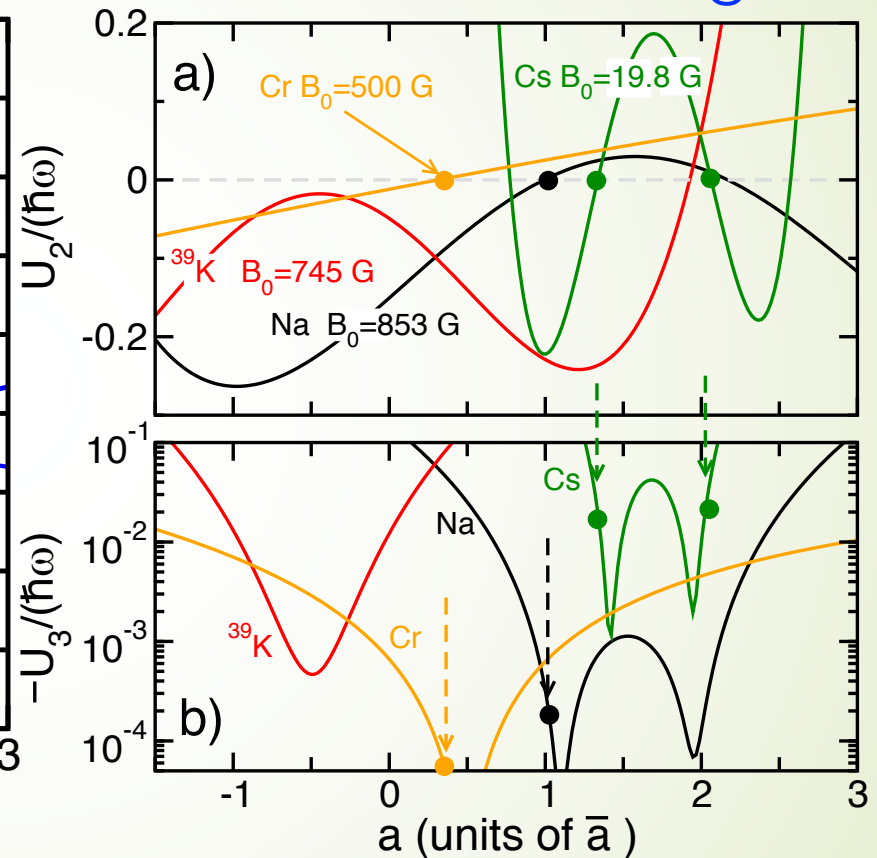
- Within the theory we have restrictions on how large V can be.
- Nevertheless, we suggest to try a few actual resonances. Both broad and narrow.

Feshbach resonances

Effective range



Interaction strength



Caveat

- Three-body recombination for reasonable g

$$K_3 = 25 \frac{\hbar^2 \bar{a}^4}{m} \quad \text{when} \quad |a| \sim \bar{a}$$

- For ^{88}Sr the experiment will be hard.
- For resonances loss puts upper limits on the depth of the lattice

$$\frac{|a|}{\bar{a}} \gg \frac{\bar{a}}{\ell}$$

harmonic oscillator length

Fermionic atoms in a lattice

- Spin up/down fermions hop around between wells (with energy J).
- Occupy only lowest vibr. level in each well.
- Only up and down atoms interact (with strength U_2 .)

See Phys. Rev. A **94**, 023607 (2016).

Superfluid are Cooper pairs

- For attractive $U_2 \sim J$, a BCS ground state of paired up/down atoms with momentum k and $-k$.
- Formally,

$$|BCS\rangle = e^{\sum_k u_k a_k^\dagger b_{-k}^\dagger} |0\rangle$$

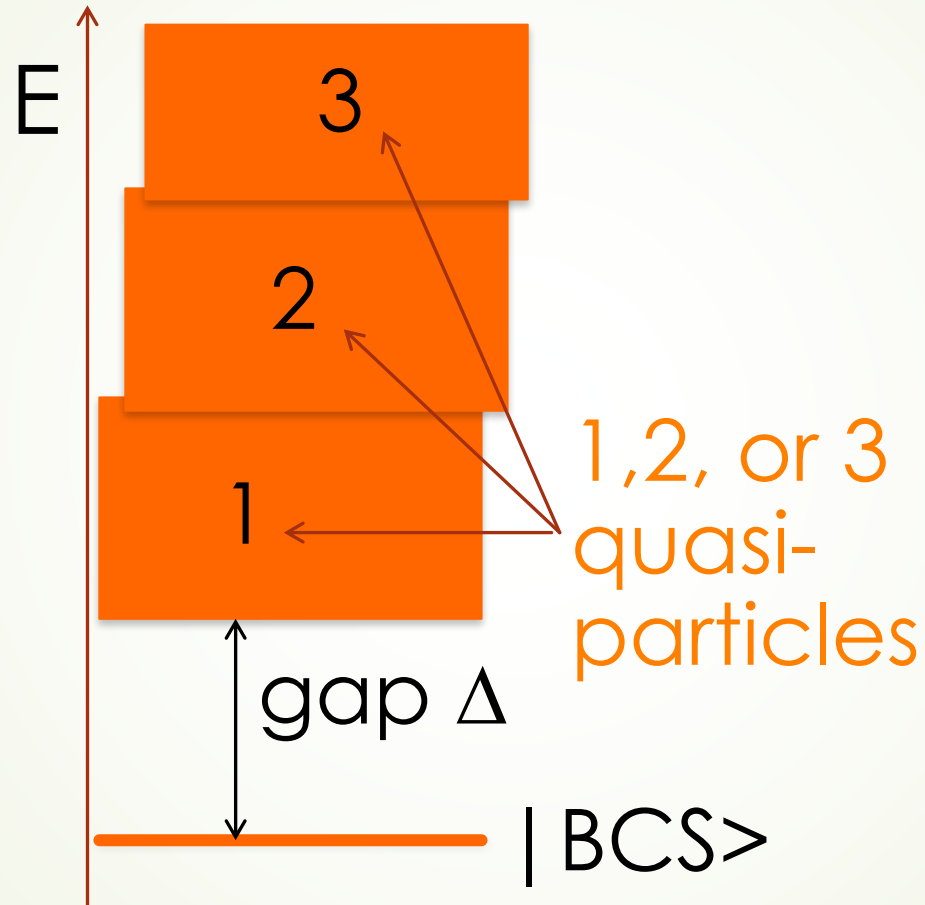
(missing band index)

Diagram illustrating the BCS ground state equation. The equation is $|BCS\rangle = e^{\sum_k u_k a_k^\dagger b_{-k}^\dagger} |0\rangle$. Arrows point from labels to the operators: "up" points to a_k^\dagger , "down" points to b_{-k}^\dagger , and "vacuum" points to $|0\rangle$.

- More importantly in each well/site

$$|0\rangle + |\uparrow\rangle + |\downarrow\rangle + |\uparrow\downarrow\rangle$$

Spectrum near half filling



From Zhao & Paramakanti, PRL **97**, 230404 (2006)

Now Quench to

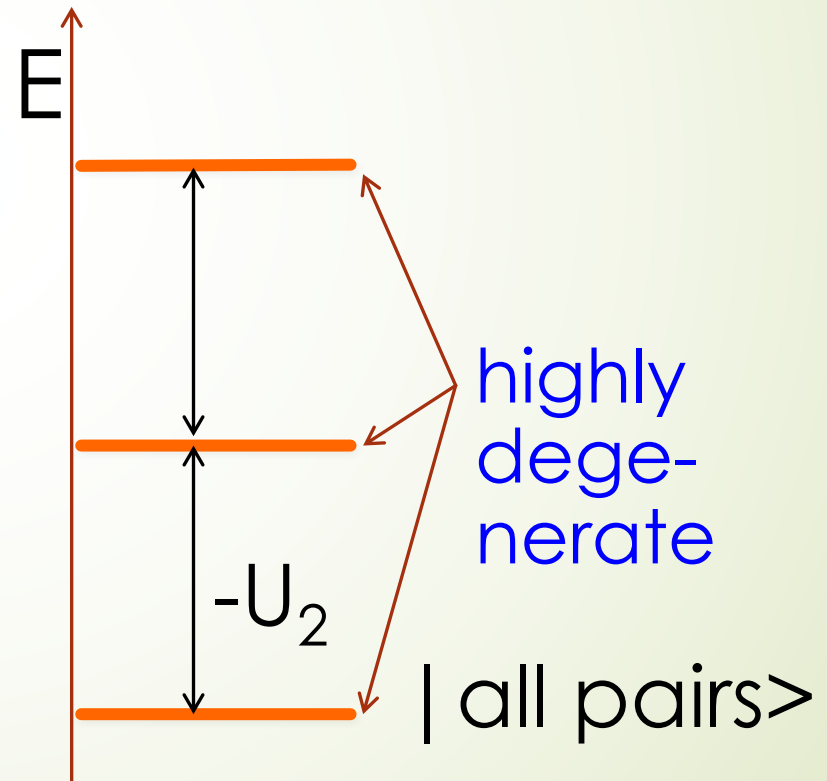
$$\frac{1}{2} U_2 b_i^\dagger a_i^\dagger b_i a_i$$

a, b -- spin up/down
 i -- unit cell

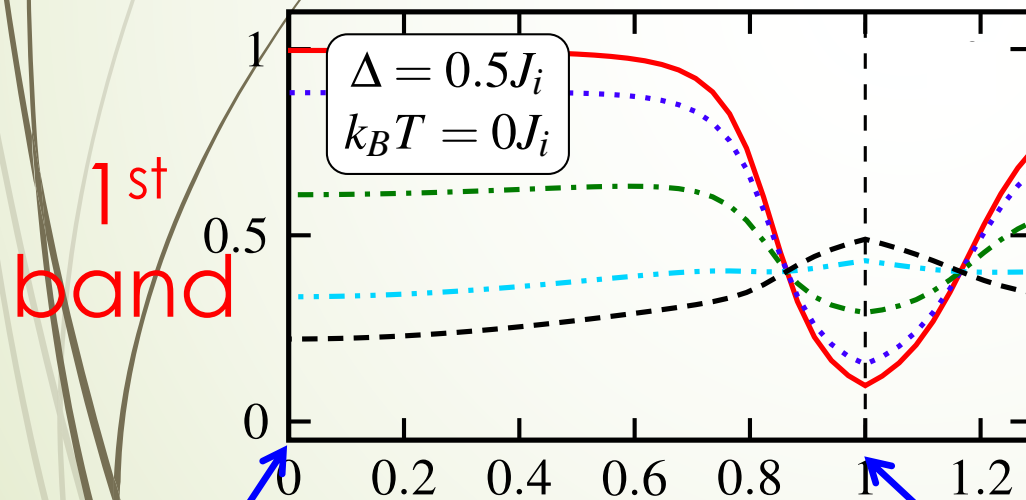
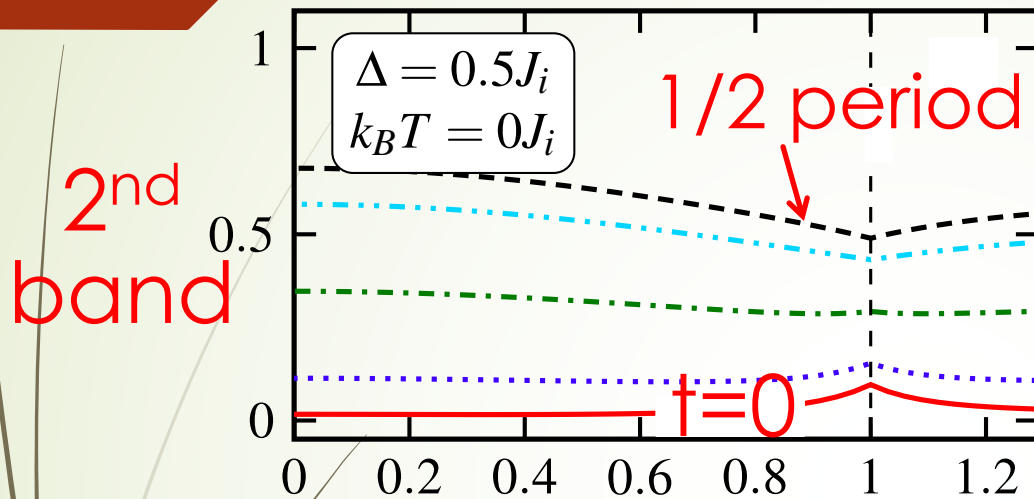
- Spectrum is exactly solvable

- Near half filling, ground state has *every other site* occupied with 2 atoms, 1 spin up and 1 spin down.

- Excited states have broken pairs.



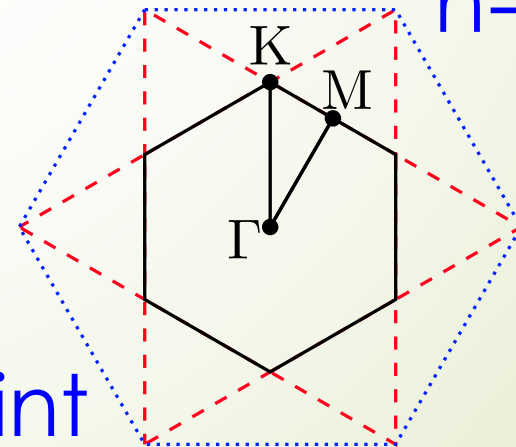
Momentum distribution in graphene



Γ point k (units of $\frac{4\pi}{3a}$)

K point

- U_2 periodic increase of population in 2nd band.
 - but also true for gap $\Delta=0$.
- $n=0.45$



Conclusions

- Two examples of few-body physics with optical lattices using atom-number superposition states.
- References
 - Phys. Rev. A **94**, 023607 (2016)
 - Phys. Rev. A **93**, 043616 (2016)
 - Phys. Rev. A **92**, 023602 (2015)