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Nature of Overstabilities in Dilute Plasmas



Tamara Bogdanović University of Maryland

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and collaborators Chris Reynolds and Steven Balbus

Role of thermal conduction in dilute plasmas

 ICM plasma is dilute and weakly magnetized-- charged particles are nearly freely streaming along the lines of magnetic field.

$$\begin{split} r_g &= \frac{3.1 \times 10^8 \,\mathrm{cm}}{Z} \left(\frac{T_g}{10^8 \,\mathrm{K}} \right)^{1/2} \left(\frac{m}{m_e} \right)^{1/2} \left(\frac{B}{1 \,\mu\mathrm{G}} \right)^{-1} \\ \lambda_e &= \lambda_i \approx 23 \,\mathrm{kpc} \left(\frac{T_g}{10^8 \,\mathrm{K}} \right)^2 \left(\frac{\mathrm{n_e}}{10^{-3} \,\mathrm{cm}^{-3}} \right)^{-1} \end{split}$$

• Anisotropic conduction alters classic condition for convection.

(Balbus 00)

Schwarzschild criterion $(\partial S/\partial z > 0)$

• What are the implications for the ICM?



MHD instabilities and overstabilities

(Balbus & Reynolds 10)



MHD equations of magnetized plasma

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho v) = 0, \qquad \text{mass}$$

$$\rho \frac{D v}{D t} = \frac{(\nabla \times B) \times B}{4\pi} - \nabla P + \rho g, \qquad \text{momentum}$$

$$\frac{\partial B}{\partial t} = \nabla \times (v \times B), \qquad \text{induction}$$

$$\frac{D \ln P \rho^{-\gamma}}{D t} = -\frac{\gamma - 1}{P} [\nabla \cdot Q + \rho \mathcal{L}], \qquad \text{entropy}$$

$$\rho \mathcal{L} \simeq 2 \times 10^{-27} n_e^2 T^{1/2} \operatorname{erg cm}^{-3} \operatorname{s}^{-1},$$

 $\chi \simeq 6 \times 10^{-7} T^{5/2} \,\mathrm{erg} \,\mathrm{cm}^{-1} \,\mathrm{s}^{-1} \,\mathrm{K}^{-1}.$ Spitzer conduction

Local WKB perturbations (Wentzel-Kramers-Brillouin)



"Instability" vs. "overstability"



$\kappa \equiv \chi T/P$ diffusivity (cm²/s) in terms of Spitzer conductivity

Stability criteria

thermal stability (Field criterion):

$$T\Theta_{T|P} + (\boldsymbol{k} \cdot \boldsymbol{b})^2 \kappa > 0,$$

thermal/heat flux driven instabilities:

$$\left(\frac{\gamma-1}{\gamma}\right)\kappa(\boldsymbol{k}\cdot\boldsymbol{b})^2\left(\left(1-2b_z^2\right)k_{\perp}^2+2b_xb_zk_xk_z\right)\cdot\frac{g}{k^2}\frac{d\ln T}{dz}+(\boldsymbol{k}\cdot\boldsymbol{v}_A)^2a_1>0.$$

radiative cooling/heat flux driven overstabilities:

$$T\Theta_{T|P} + \kappa (\mathbf{k} \cdot \mathbf{b})^2 \mathcal{R} > 0,$$

$$\mathcal{R} = f (dInT, dInp) < 1$$
 "reduction factor"

Interplay of thermal stability and overstabilities

$T\Theta_{T P} + (\boldsymbol{k} \cdot \boldsymbol{b})^2 \kappa > 0,$	$T\Theta_{T P} + (\boldsymbol{k} \cdot \boldsymbol{b})^2 \kappa < 0,$
$T\Theta_{T P} + \kappa (\boldsymbol{k} \cdot \boldsymbol{b})^2 \mathcal{R} > 0,$	$T\Theta_{T P} + \kappa (\boldsymbol{k} \cdot \boldsymbol{b})^2 \mathcal{R} > 0.$
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Interplay of thermal stability and overstabilities

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Interplay of thermal stability and overstabilities





 $T\Theta_{T|P} + (\boldsymbol{k} \cdot \boldsymbol{b})^2 \kappa > 0,$

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 $T\Theta_{T|P} + (\boldsymbol{k} \cdot \boldsymbol{b})^2 \kappa < 0,$

 $T\Theta_{T|P} + \kappa (\boldsymbol{k} \cdot \boldsymbol{b})^2 \mathcal{R} < 0$

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Heat-flux Driven Overstability (HFO)



Thermally stable and overstable to HFO

Thermally unstable and overstable to HFO



~2500 t_{dyn}

~200 t_{dyn}

HFO growth and dispersion relation



Overstability driven by radiative cooling (RCO)





Thermally stable and overstable to RCO

 $\omega_{cool} > \omega_{dyn} \approx \omega_{cond}$

Thermally unstable and overstable to RCO

 $\omega_{\text{cool}} > \omega_{\text{dyn}} > \omega_{\text{cond}}$



~400 t_{dyn}

~100 t_{dyn}

RCO growth and dispersion relation



Conclusions & prospects

- MHD overstabilities related to the well known MHD instabilities.
- Analytic theory: Need for thorough understanding of plasma processes (e.g., anisotropic viscosity).
- Simulations: Understanding the relative importance of individual plasma instabilities/overstabilities and their connection.
- Observations: Spectro-polarimetric measurements, measurements of temperature profiles and metallicity in cluster outskirts, measurements of turbulence, other yet to be realized methods...