



Numerical simulations of sediment transport and aeolian ripples

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Goal & outline

- Microscopic mechanisms of transport
- Empirical laws and modeling assumptions

- **Model**

O. Durán, B. Andreotti and P. Claudin, Phys. Fluids **24**, 103306 (2012).

- Transport layer and sediment flux

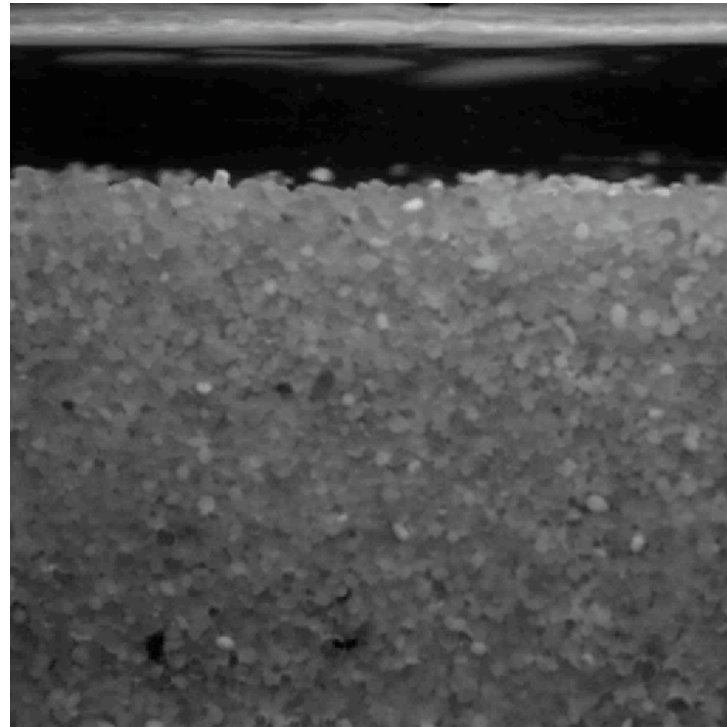
- Scaling laws

- **Aeolian ripples**

O. Durán, P. Claudin and B. Andreotti , submitted.

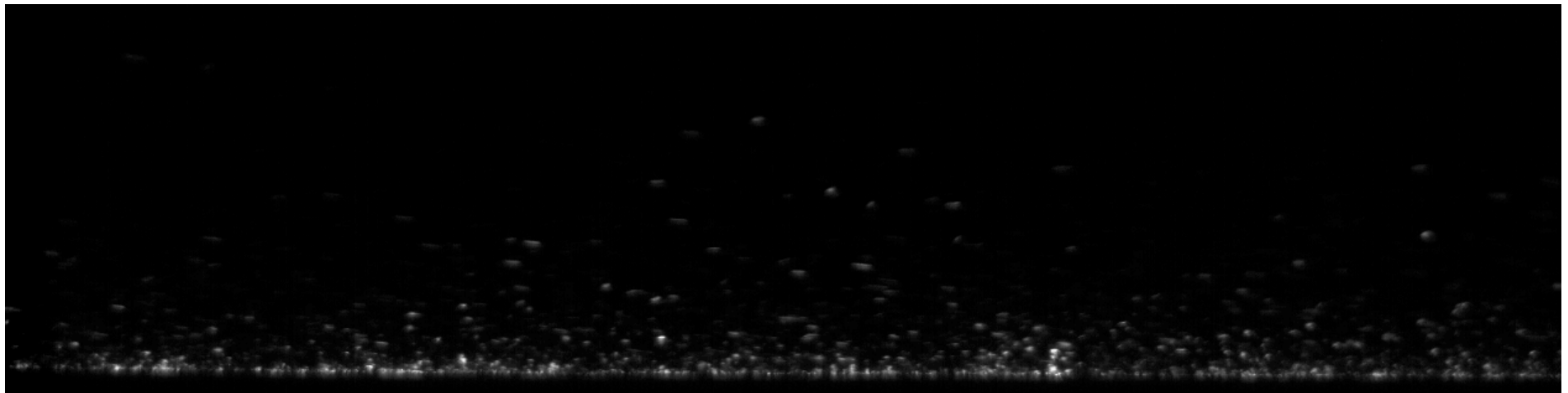
Sediment transport

Sub-aqueous bed load



Aeolian saltation

From E. Lajeunesse (IPGP)



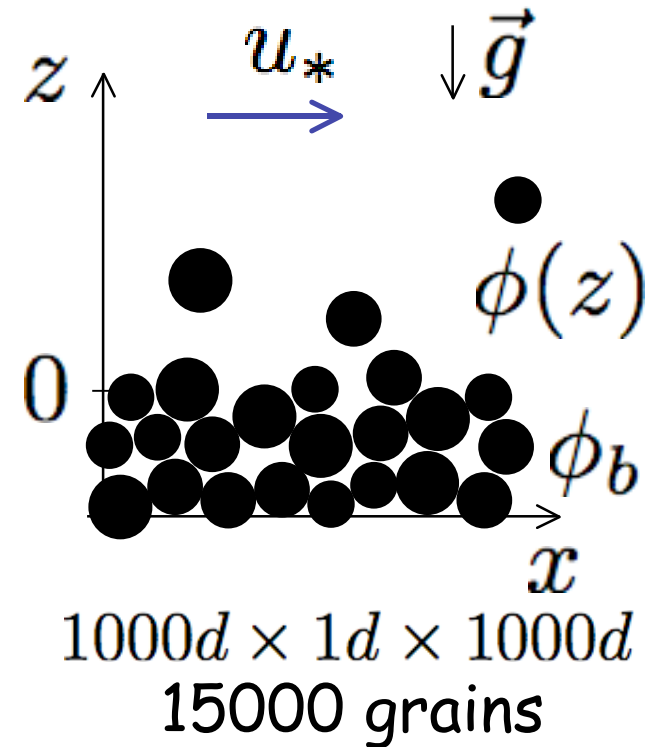
Champ: 12 cm x 3 cm

Acquisition 1kHz

The numerical model

- Principle:

GRAINS (DEM)	$m \frac{d\vec{u}^p}{dt} = \dots$
FLUID (RANS)	$\rho_f (1 - \phi) \frac{D\vec{u}^f}{Dt} = \dots$
COUPLING	$\vec{f}_{\text{fluid}}^p = \vec{f}_{\text{drag}}^p + \vec{f}_{\text{Arch}}^p$



- Control parameters:

$$\Theta = \frac{\rho_f u_*^2}{(\rho_p - \rho_f)gd}$$

Sields

$$\rho_p / \rho_f$$

Density ratio

$$\mathcal{R} = \frac{d}{\nu} \sqrt{\left(\frac{\rho_p}{\rho_f} - 1 \right) gd}$$

Grain Reynolds

The grains

$$m \frac{d\vec{u}^p}{dt} = m \vec{g} + \sum_q \vec{f}^{p,q} + \vec{f}_{\text{fluid}}^p$$

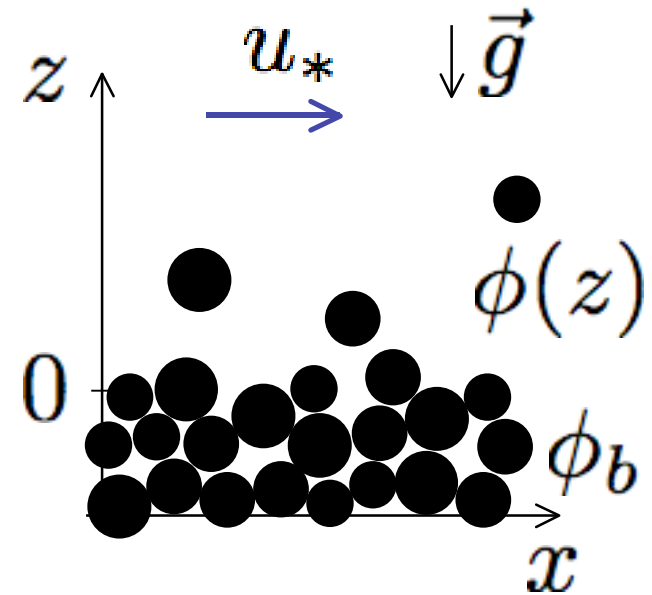
$$I \frac{d\vec{\omega}^p}{dt} = \frac{d}{2} \sum_q \vec{n}^{p,q} \times \vec{f}^{p,q}$$

$$\vec{f}_{\text{fluid}}^p = \vec{f}_{\text{drag}}^p + \vec{f}_{\text{Arch}}^p$$

$$\vec{f}_{\text{drag}}^p = \frac{\pi}{8} \rho_f d^2 C_d(R_u) |\vec{u} - \vec{u}^p| (\vec{u} - \vec{u}^p)$$

$$C_d(R_u) = \left(\sqrt{C_d^\infty} + \sqrt{R_u^c / R_u} \right)^2$$

$$\vec{f}_{\text{Arch}}^p = \frac{\pi}{6} d^3 \text{div} \sigma^f$$



$$R_u = |\vec{u} - \vec{u}^p| d / \nu$$

Hydrodynamics

RANS

$$\rho_f(1 - \phi)D_t u_i = -\partial_i p^f + \rho_f(1 - \phi)g_i + \partial_j \tau_{ij}^f - F_i$$

$$\partial_z p^f = -\rho_f g$$

$$\partial_z \tau^f = F_x$$

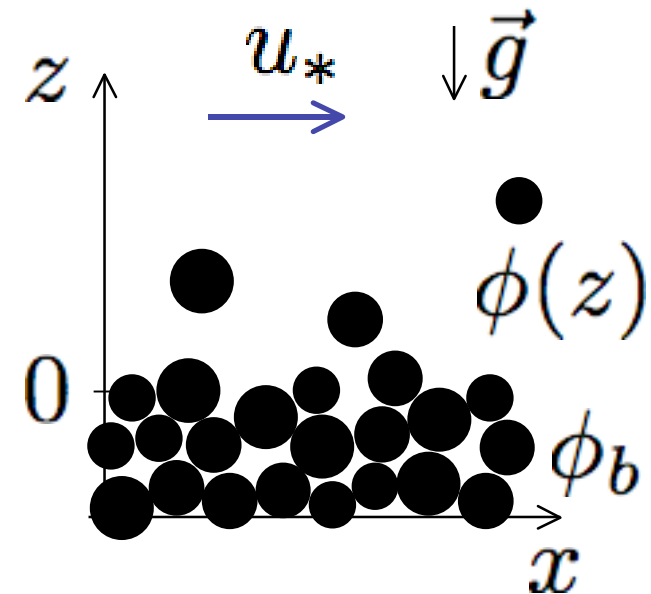
CLOSURE

$$\tau^f = \rho_f(\nu + \ell^2 |\partial_z u|) \partial_z u$$

mixing length

COUPLING

$$\vec{F}(z) = \frac{1}{A dz} \left\langle \sum_{p \in \{z; z+dz\}} \vec{f}_{\text{fluid}}^p \right\rangle$$



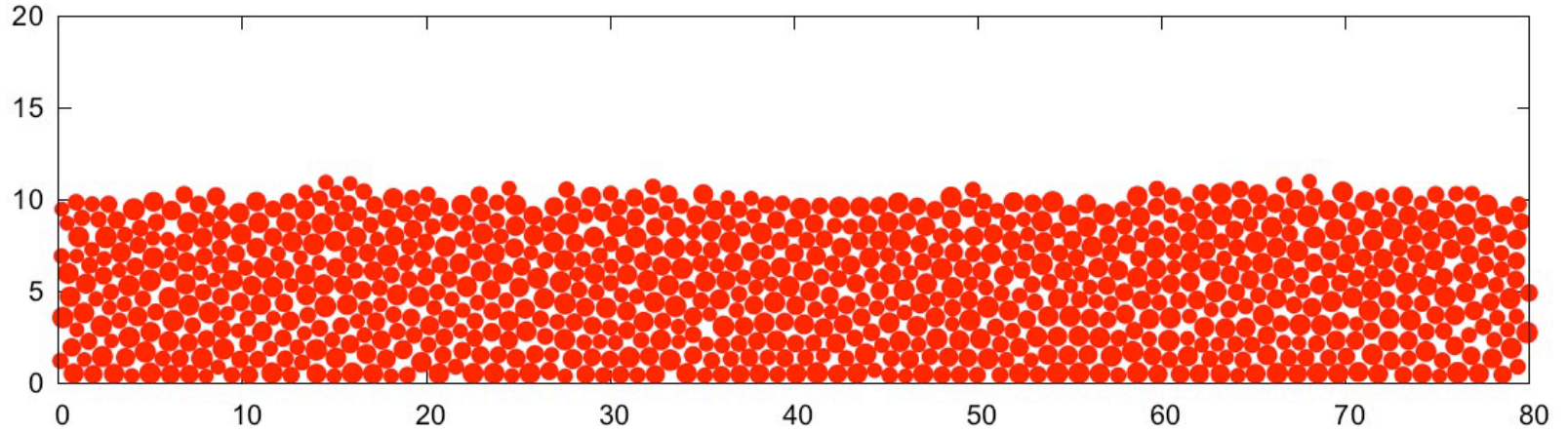
Movies

flow →

$$\sqrt{\Theta/\Theta_d} = 2$$

Water

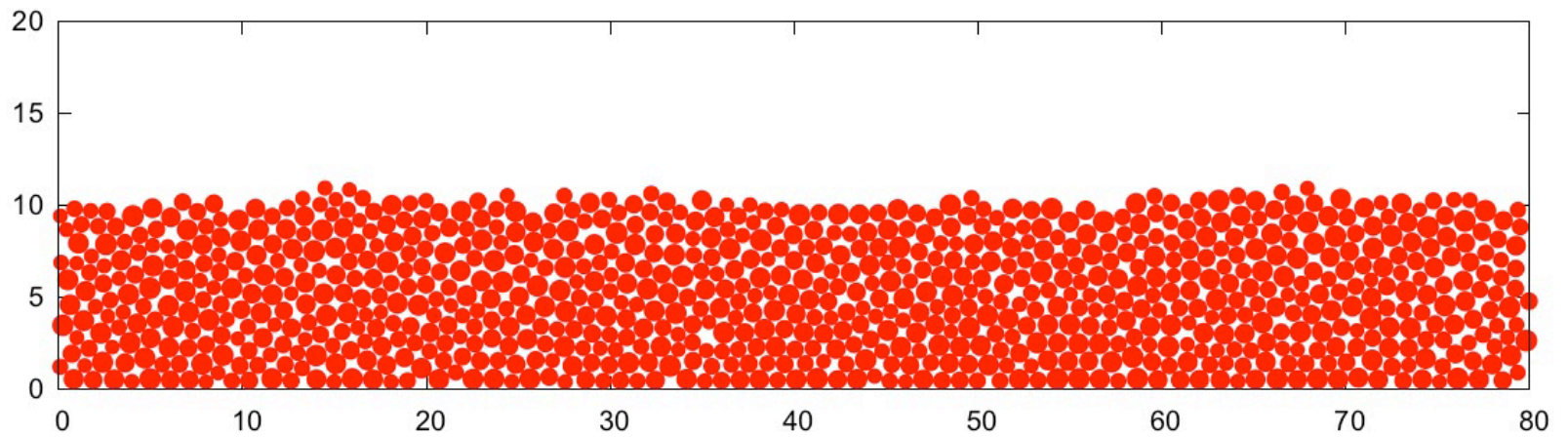
$$\frac{\rho_p}{\rho_f} = 2$$



ϕ_b

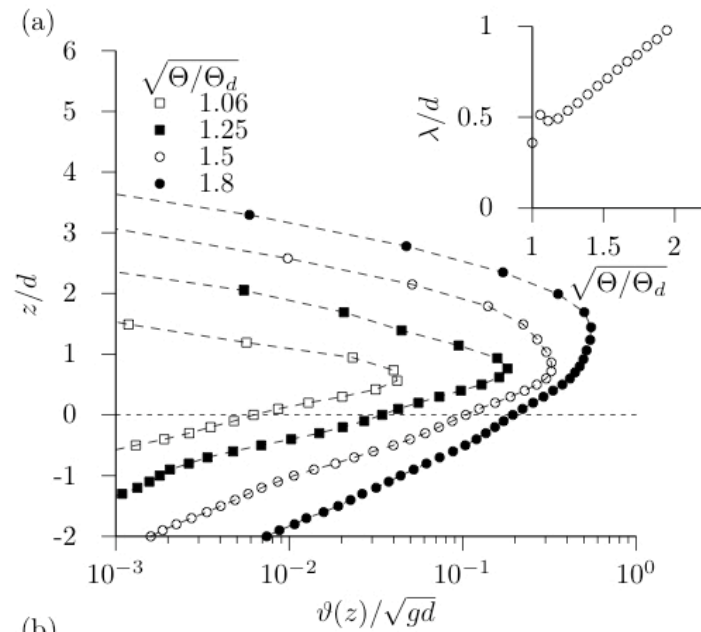
Air

$$\frac{\rho_p}{\rho_f} = 2000$$

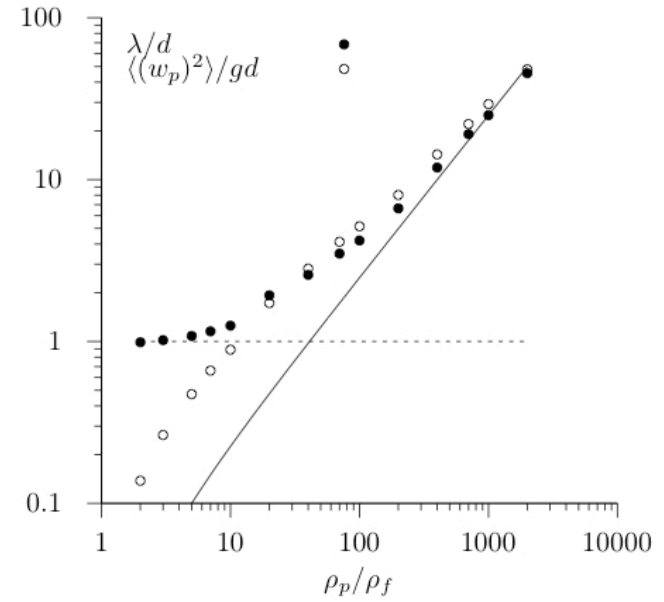
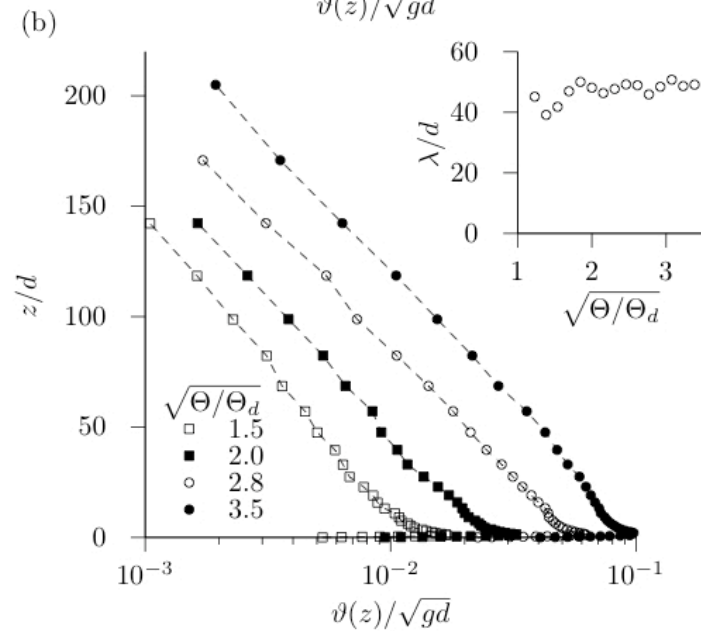


Transport layer

Water



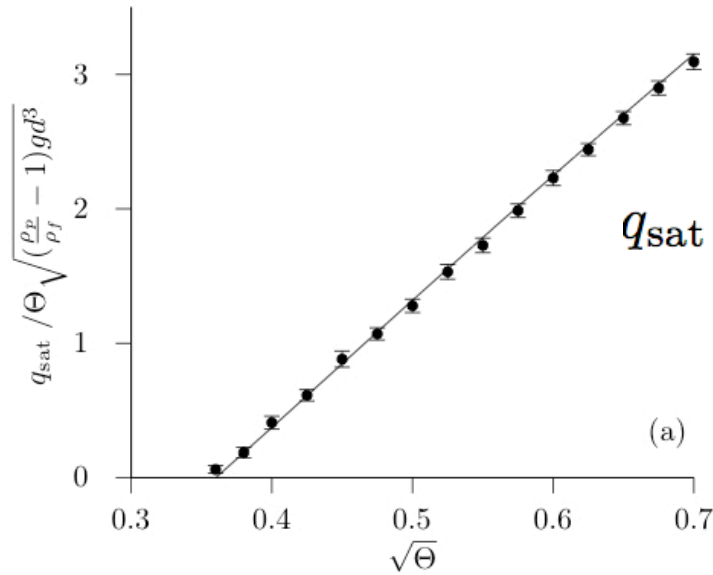
Air



Saturated fluxes

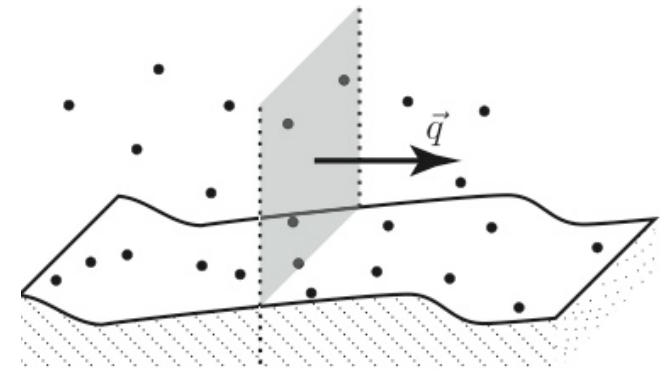
$$q_{\text{sat}} = \frac{1}{A\phi_b} \frac{\pi}{6} d^3 \sum_p u^p$$

Water

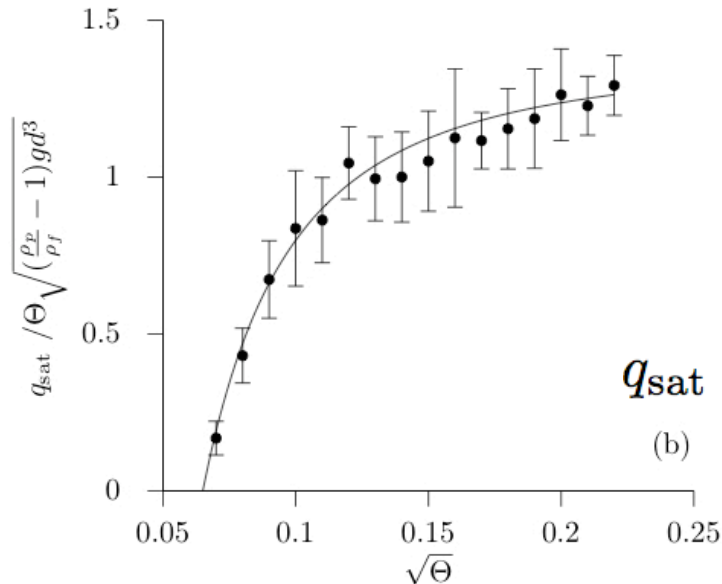


$$q_{\text{sat}} = \frac{u_d d}{\phi_b \mu_d} (\Theta - \Theta_d) \left(\sqrt{\frac{\Theta}{\Theta_d}} - \sqrt{\frac{\mu_d}{\mu_s}} \right)$$

Meyer-Peter & Müller (1948)



Air

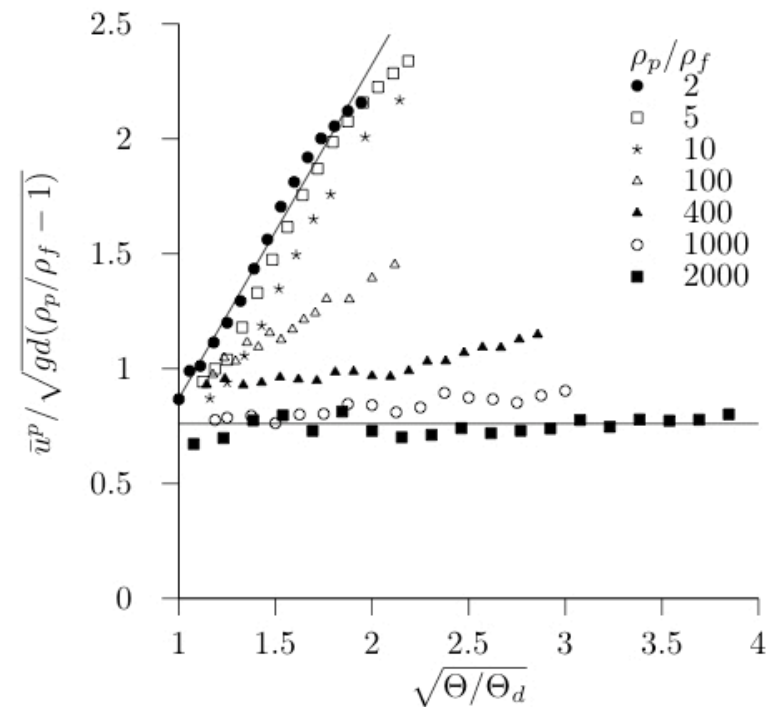
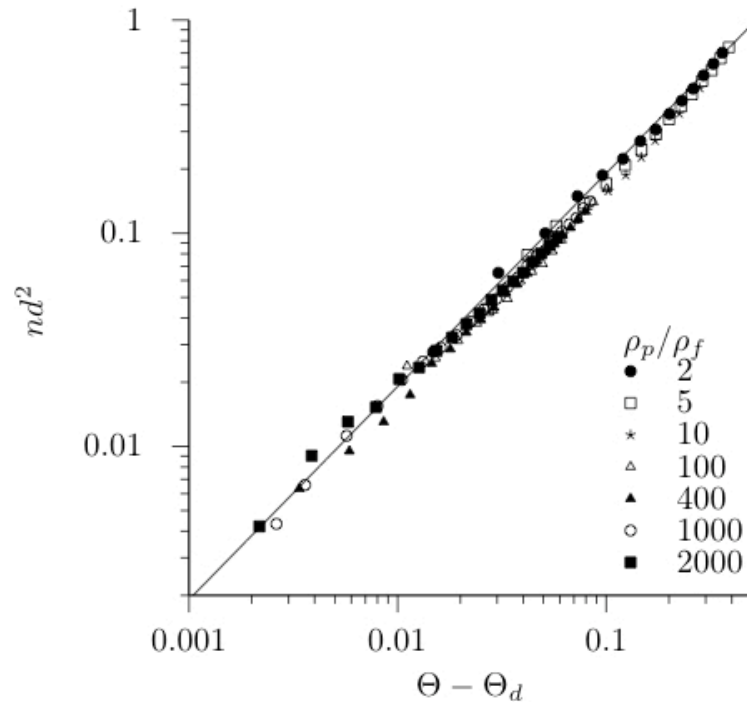


$$q_{\text{sat}} \propto (1 - \rho_f/\rho_p) u_d d (\Theta - \Theta_d)$$

Rasmussen et al. (1991, 1996, 1999)

Flux decomposition

$$q_{\text{sat}} = \frac{1}{\phi_b} \frac{\pi}{6} d^3 n \bar{u}^p$$



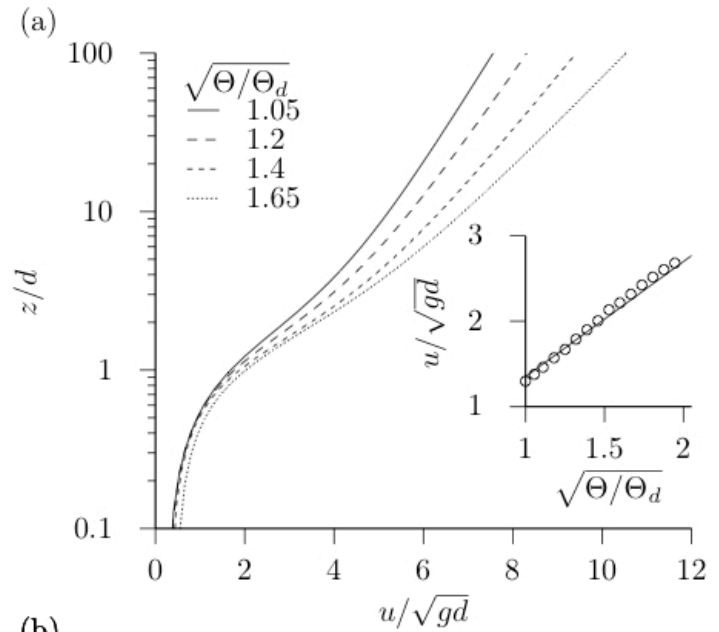
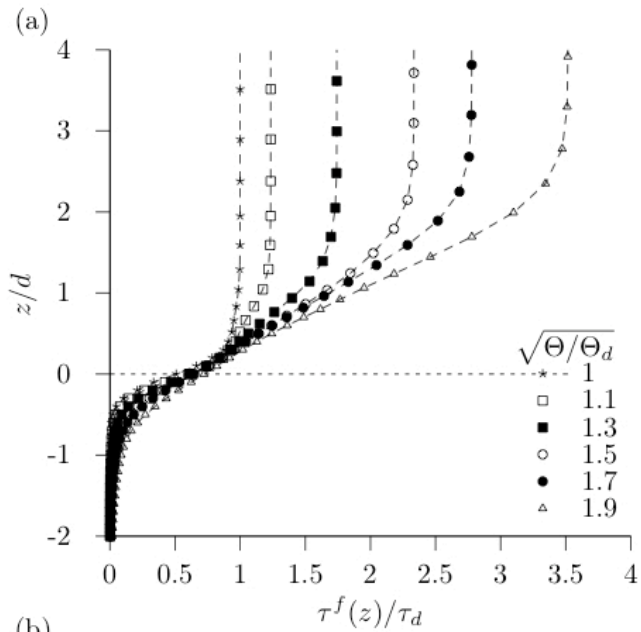
$$n = \frac{\left(\sum_p u_p\right)^2}{A \sum_p u_p^2}$$

$$\bar{u}^p = \frac{\sum_p u_p^2}{\sum_p u_p}$$

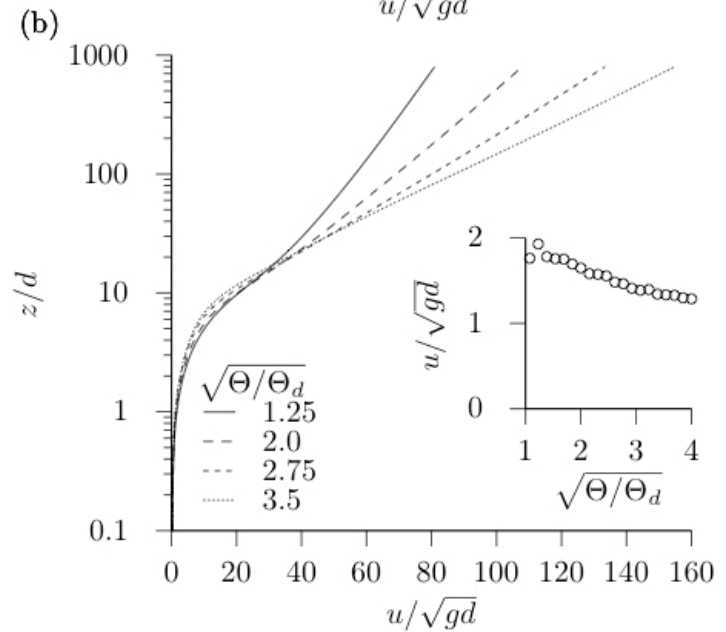
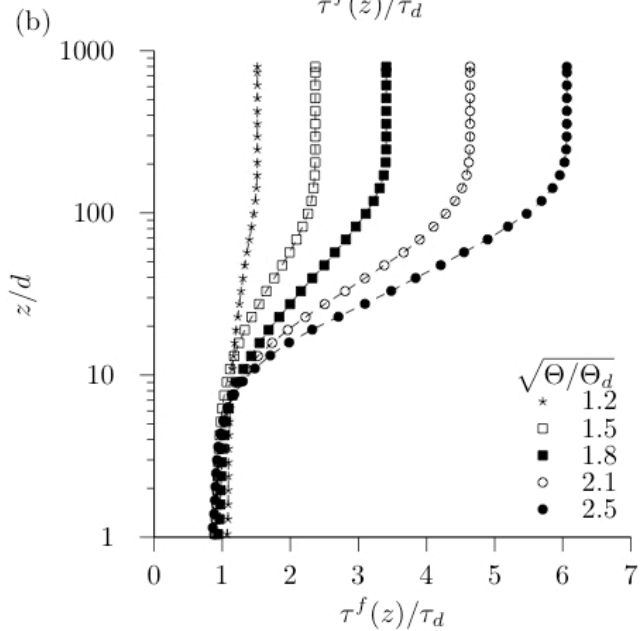
Lajeunesse et al. (2010)

Grains' feedback on the flow

Water



Air



Interpretation

$$q_{\text{sat}} = \frac{1}{\phi_b} \frac{\pi}{6} d^3 n \bar{u}^p$$

Water

- Bagnold's approach (1956)

$$\tau^p = n f_d \quad \tau^p = \rho_f u_*^2 - \tau_d$$
$$f_d \simeq C_d \rho_f (u - \bar{u}^p) d^2 \simeq \mu (\rho_p - \rho_f) g d^3$$

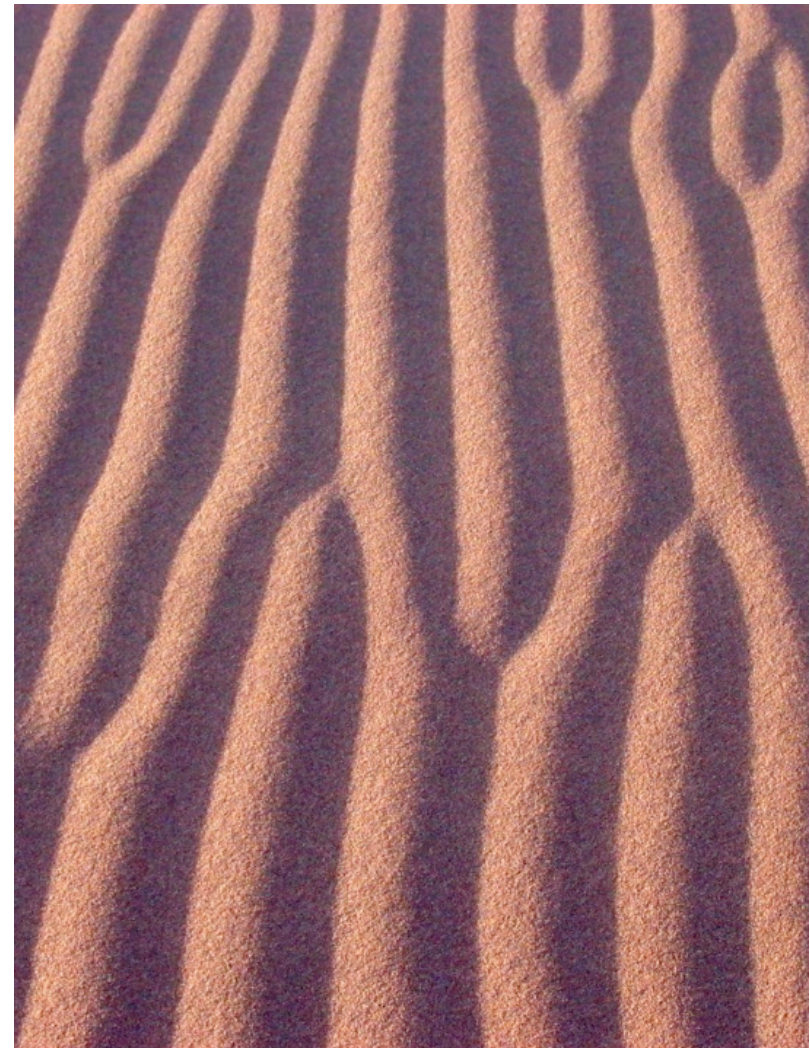
$$n \simeq (\Theta - \Theta_d) / d^2$$
$$\bar{u}^p \simeq u_d \left(\sqrt{\Theta / \Theta_d} - 1 \right)$$

Air

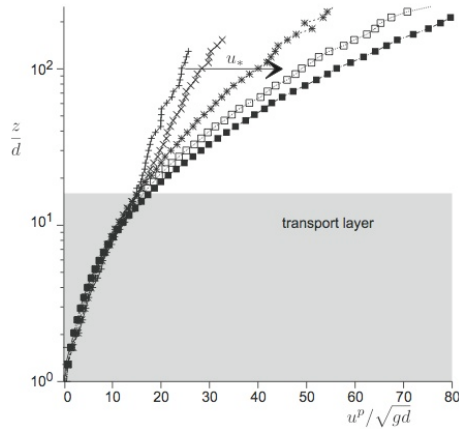
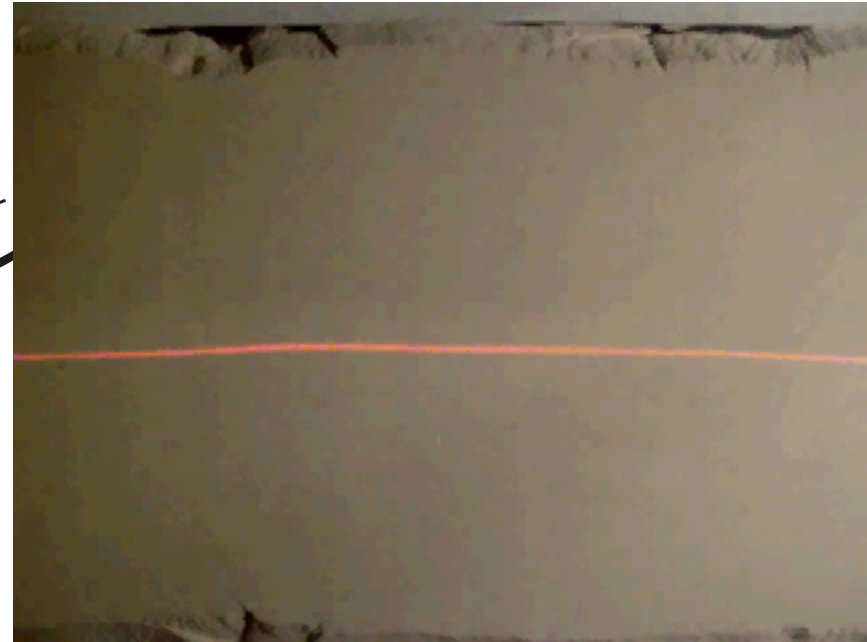
- Approach of Owen (1964) and Ungar & Haff (1987)

$$\bar{u}^p \simeq u_d \quad \text{replacement capacity} = 1$$

Aeolian ripples

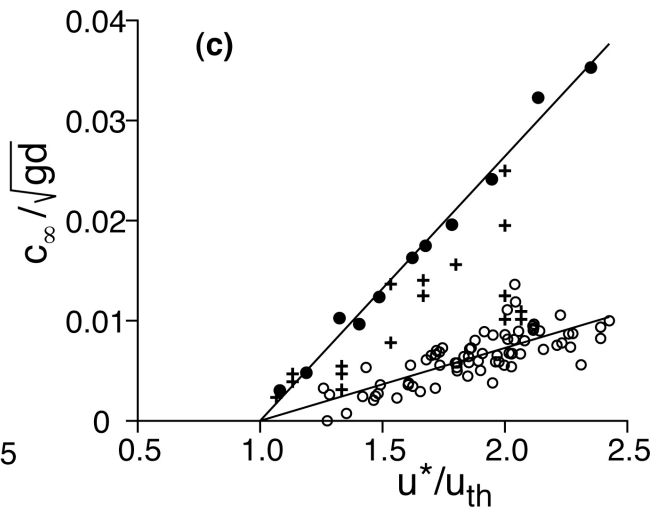
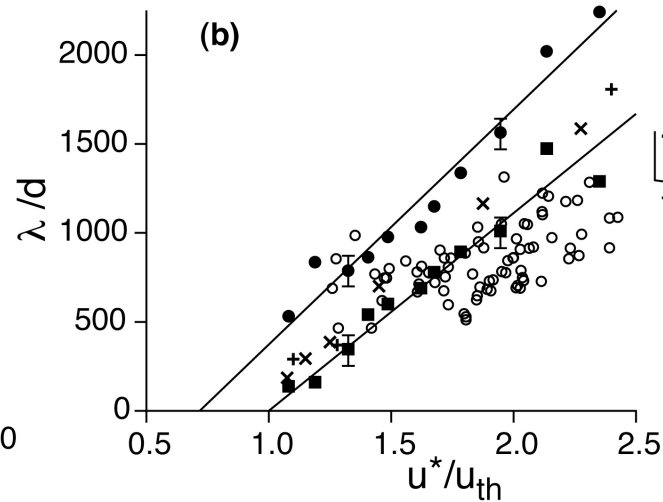
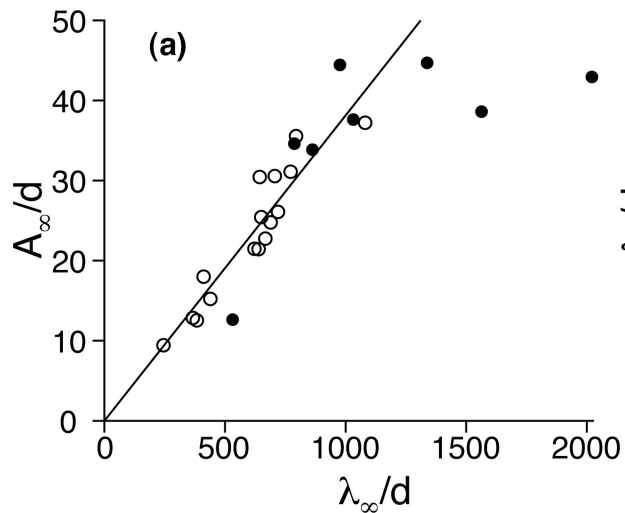


Ripple size and speed



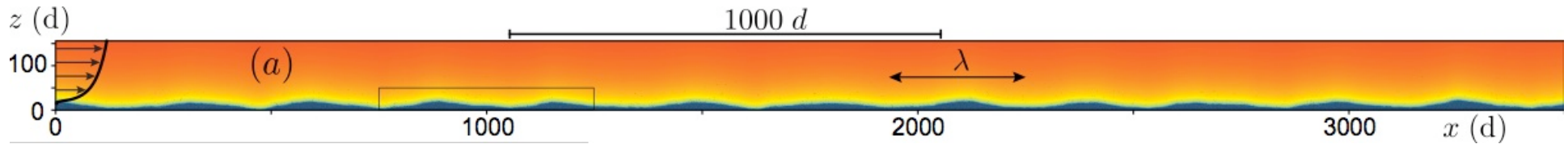
Anderson 1987
Csahok et al. 2000
Yizhaq 2004

Screening mechanism



Andreotti et al. 2006

Numerical ripples

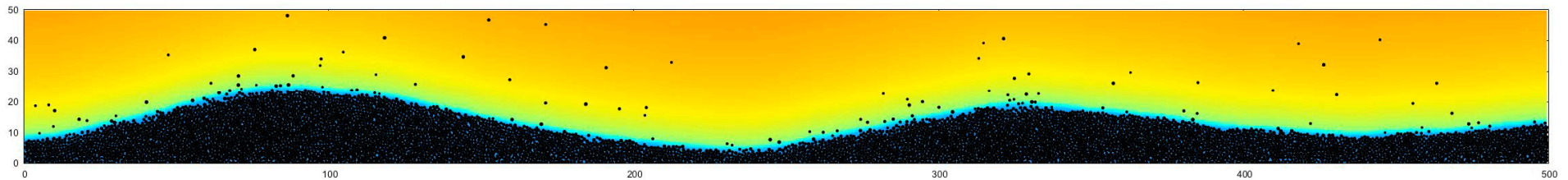


$$\rho_p / \rho_f = 500$$

$$\mathcal{R} = 22$$

$$\Theta / \Theta_{\text{th}} = 1-25$$

$3400d \times 1d \times 1000d$
45000 grains



Bed elevation profile

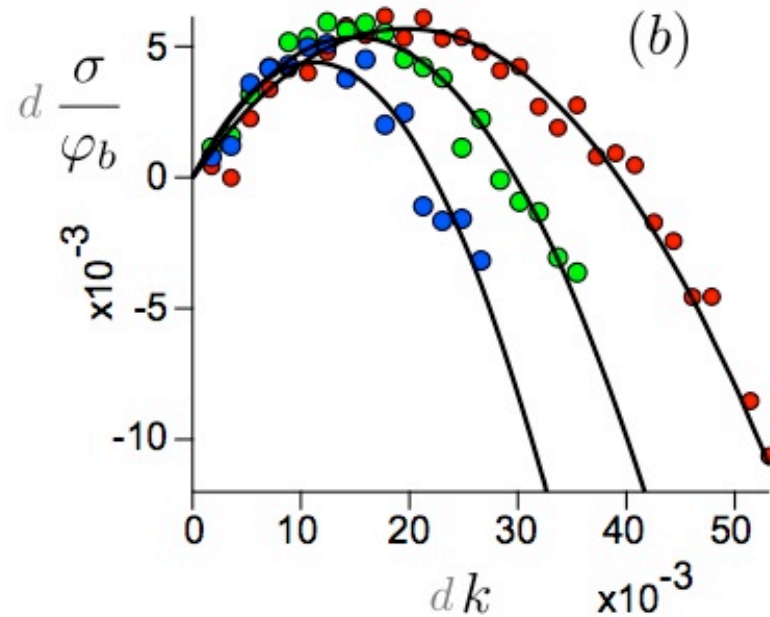
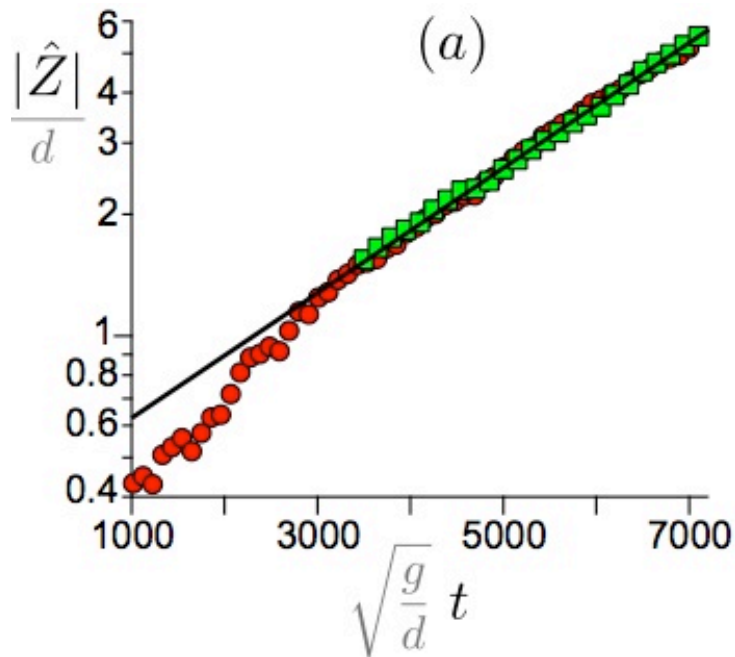
$$Z(x) = \frac{1}{\phi_b} \int_0^\infty \left[\phi(x, z) - \frac{1}{A} \int_0^A \phi(x, z) dx \right] dz$$

Dispersion relation

From a flat bed

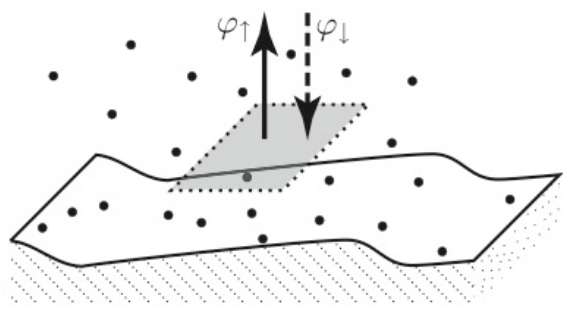
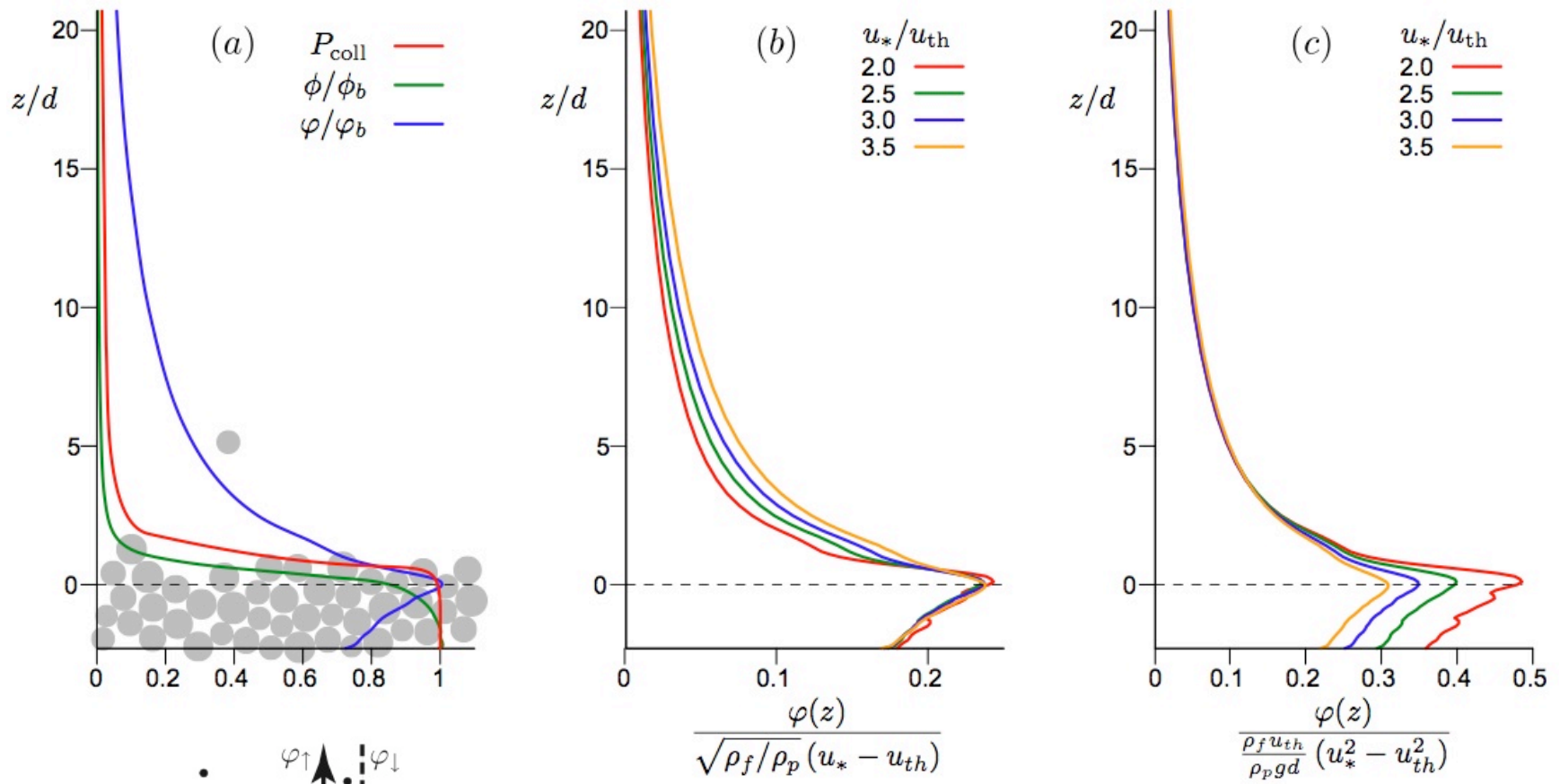
From a modulated bed e^{ikx}

$$\rightarrow |\hat{Z}|(t) = |\hat{Z}|(0) e^{\sigma t}$$



- $u_*/u_{th} = 3$
- $u_*/u_{th} = 4$
- $u_*/u_{th} = 5$

Flat bed: upper and basal layers



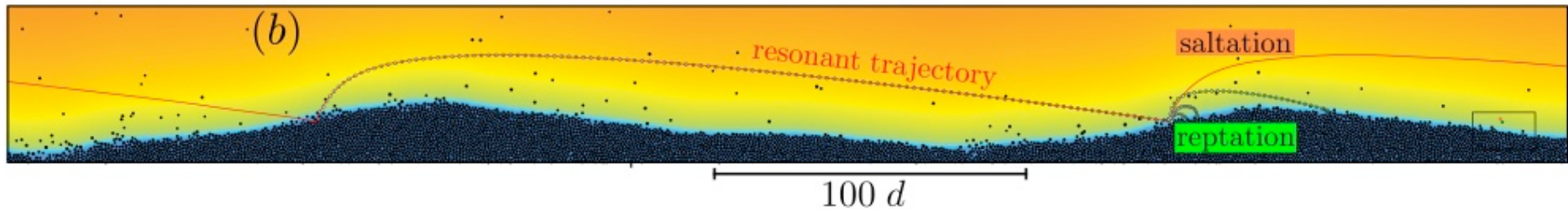
collisional (gaseous) layer

$$\tau^p \sim \rho_p u_b^2$$

$$\phi_b \sim u_b$$

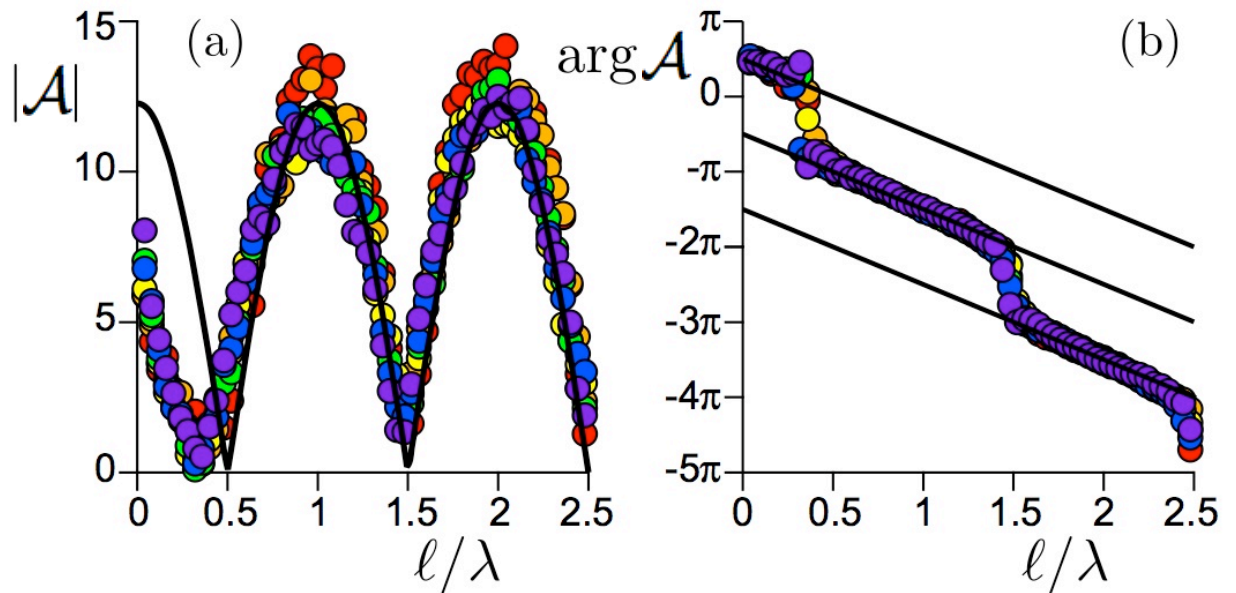
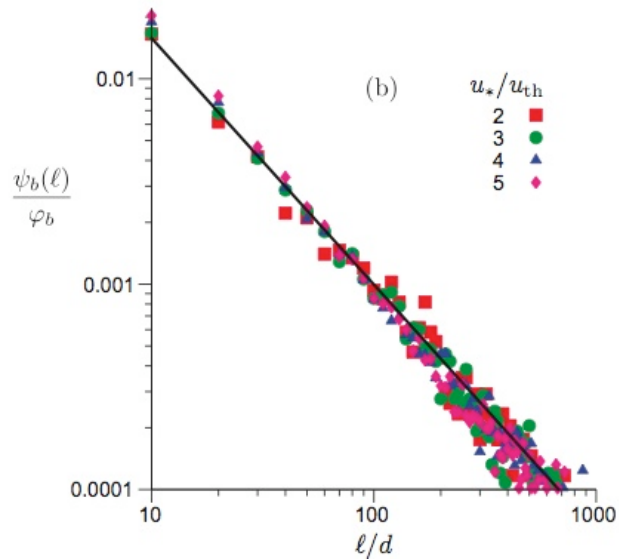
$$\varphi = \frac{1}{\phi_b A dz} \frac{\pi}{6} d^3 \sum_{p \in \{z, z+dz\}} u_{\uparrow}^p$$

Resonant trajectories



Hop length distribution

$$\psi(\ell)$$



$$\hat{\psi}(\ell) = \mathcal{A} \psi_b(\ell) i k \hat{Z}$$

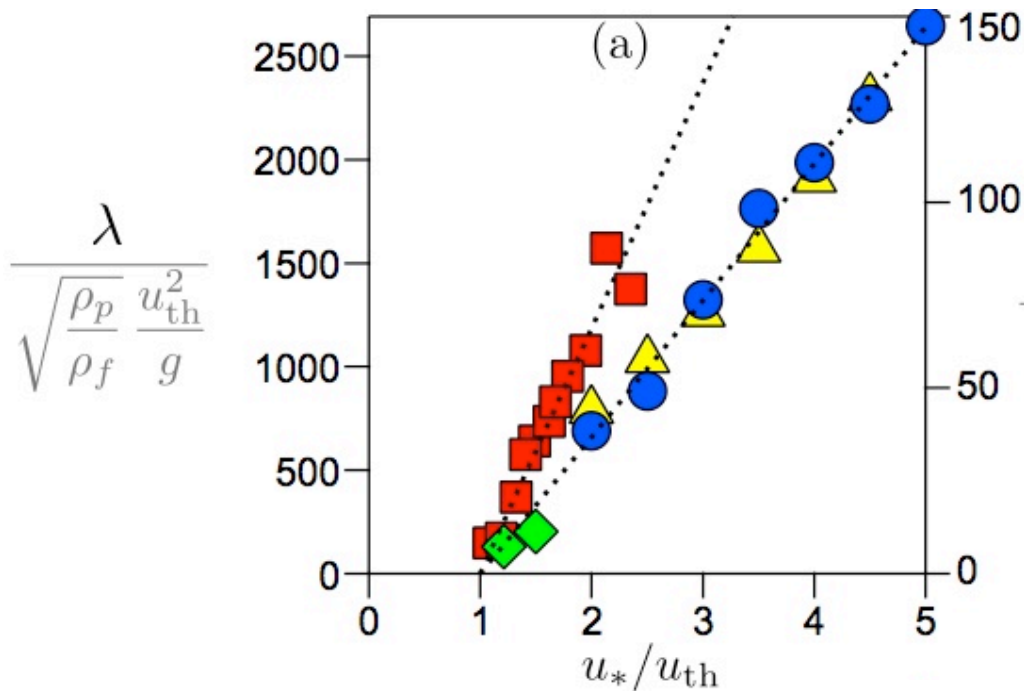
Ripple instability (1)

Gravity-induced diffusion of transport (saltons)

Growth rate

$$\sigma = a\varphi_b k - bq_s k^2$$

Small-trajectory grains (reptons) ejected by resonant grains



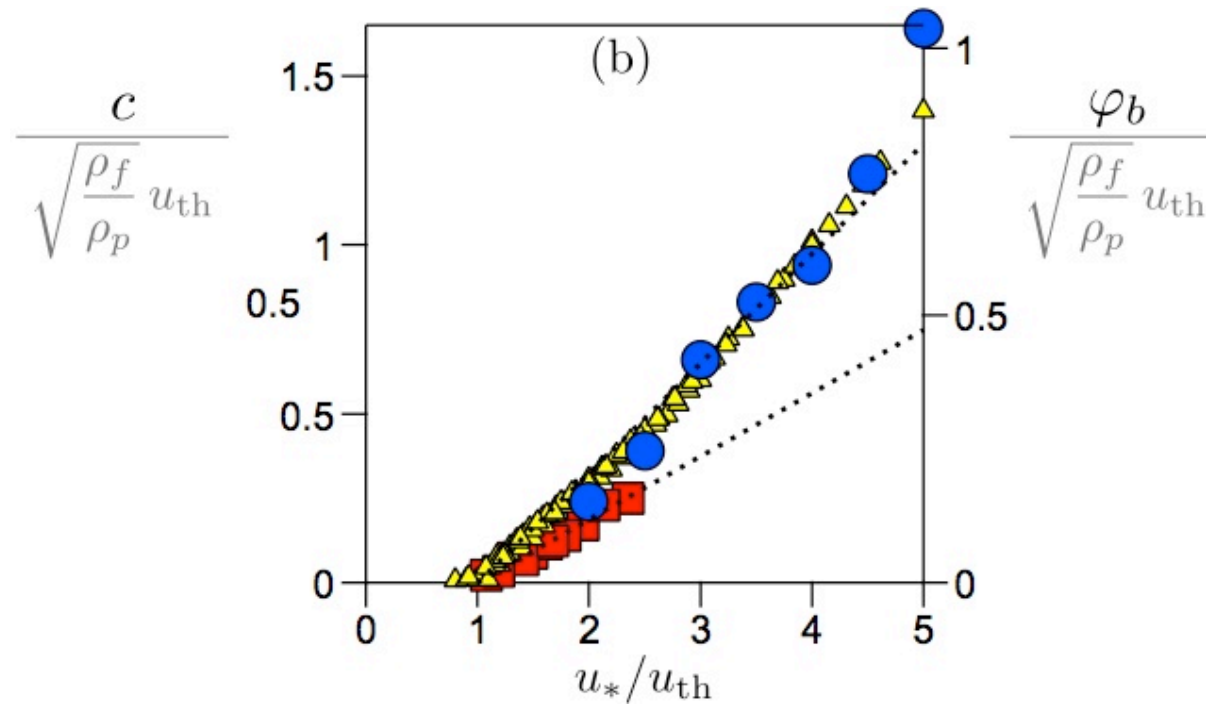
$$\lambda \sim q_s / \varphi_b$$

■ Experimental data from Andreotti et al. (2006)

Ripple instability (2)

Propagation speed

$$c \sim \varphi b$$



Summary & perspectives

- Mechanisms and scaling laws for sediment transport
- Mechanisms and scaling laws for aeolian ripples

- Non-homogeneous situations
transients → saturation length L_{sat}
- Transition to from bed load to suspended load

Supplementary slides

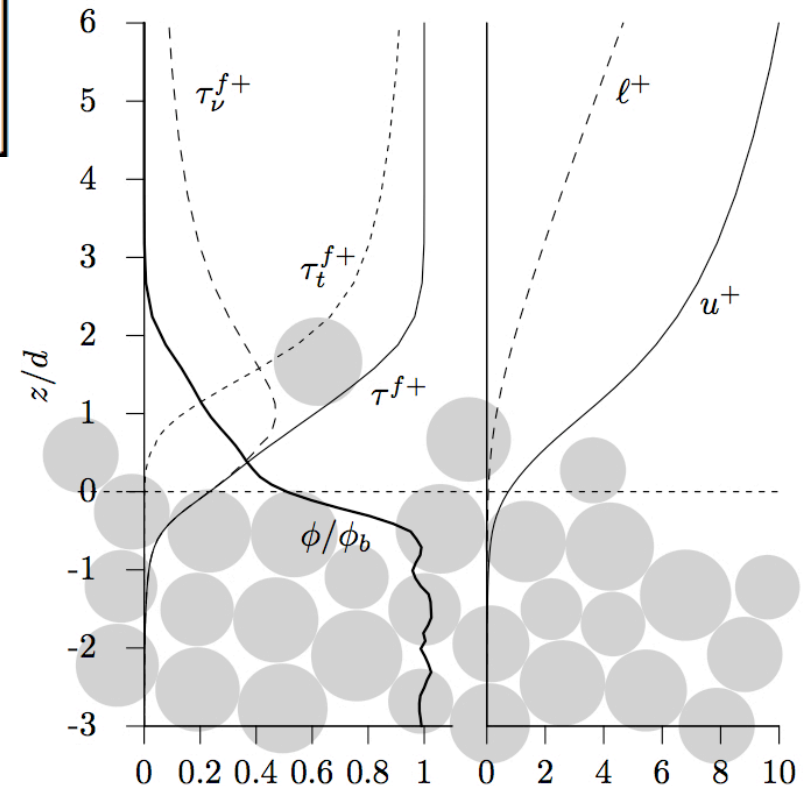
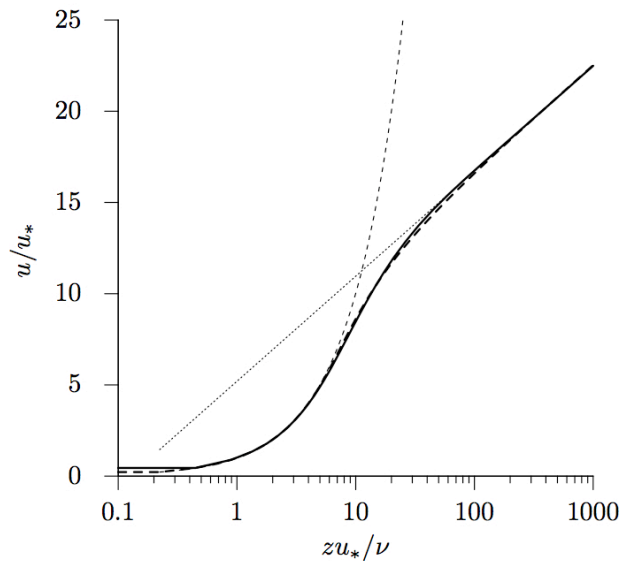
Mixing length

$$\tau^f = \rho_f (\nu + \ell^2 |\partial_z u|) \partial_z u$$

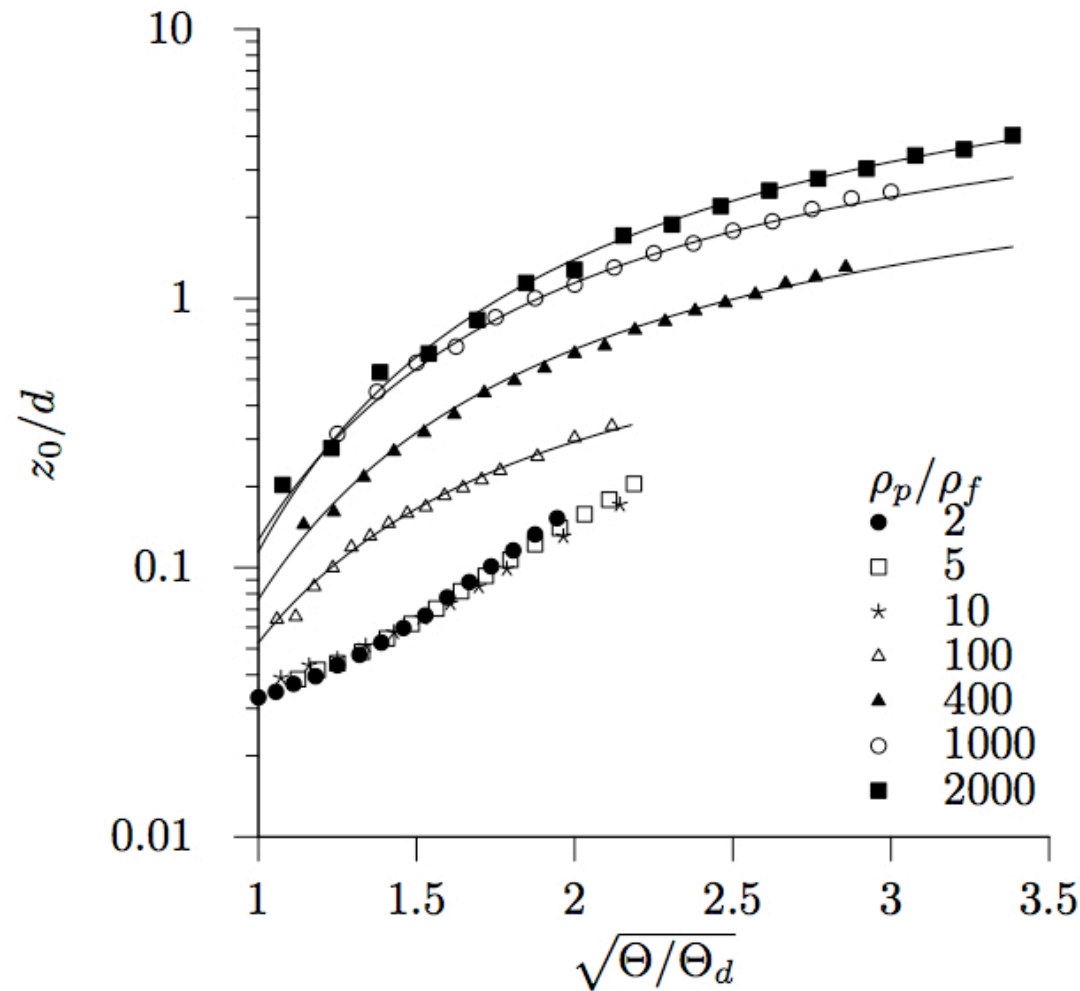
$$\ell = \kappa z \left[1 - \exp \left(-\frac{1}{R_{vD}} \frac{zu_*}{\nu} \right) \right]$$

van Driest (1956)

$$\partial_z \ell = \kappa \left[1 - \exp \left(-\sqrt{\frac{1}{R_c} \left(\frac{u\ell}{\nu} \right)} \right) \right]$$



Hydrodynamical roughness

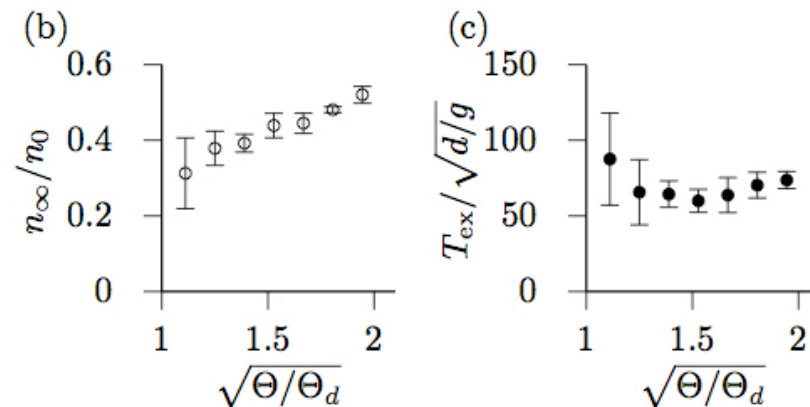
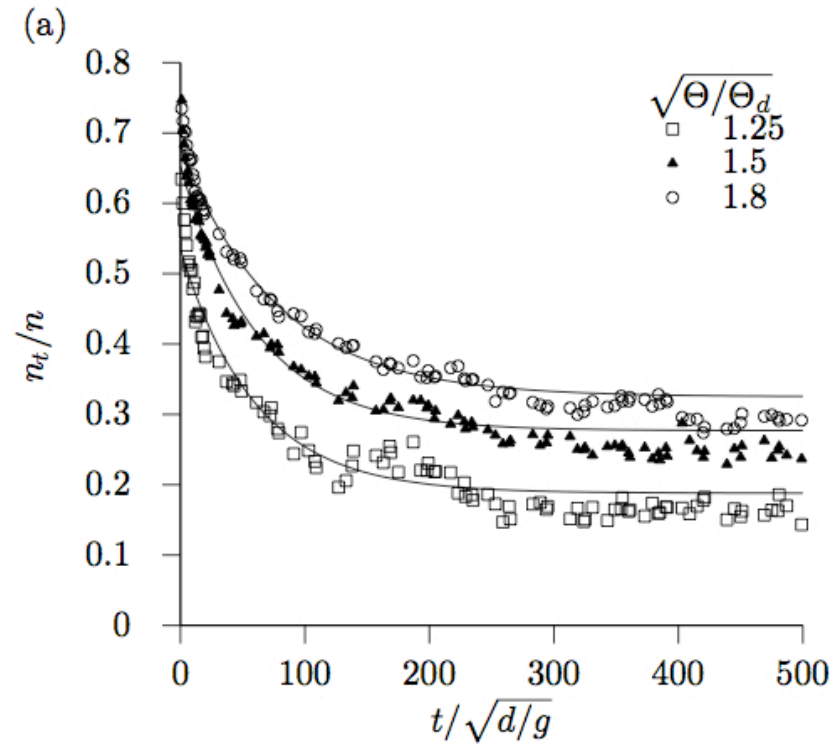


$$u = \frac{u_*}{\kappa} \ln \frac{z}{z_0}$$

Bagnold's focal point $z_0 = H_f \exp(-\kappa U_f / u_*)$

Exchange time

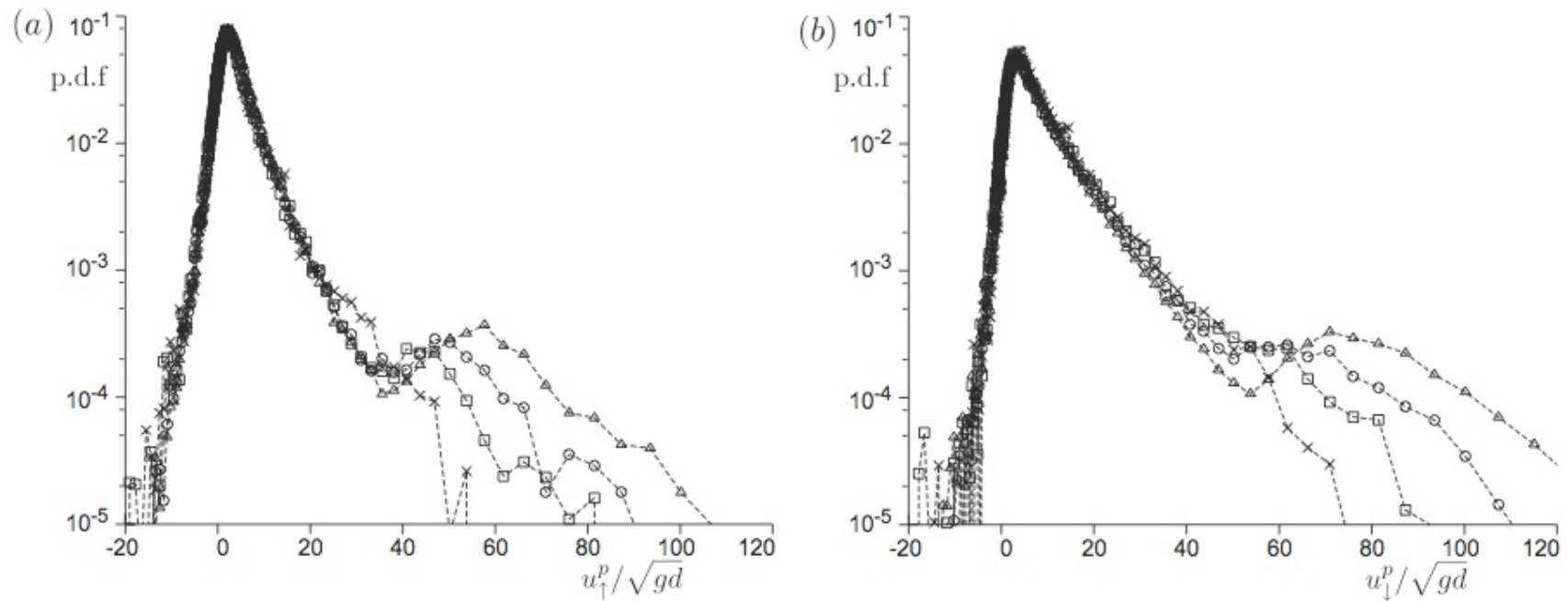
$$n_t = \frac{\left(\sum_{p \in \varepsilon} u_p \right)^2}{A \sum_{p \in \varepsilon} u_p^2}$$



$$n_t(t) = (n_0 - n_\infty) \exp(-t/T_{\text{ex}}) + n_\infty$$

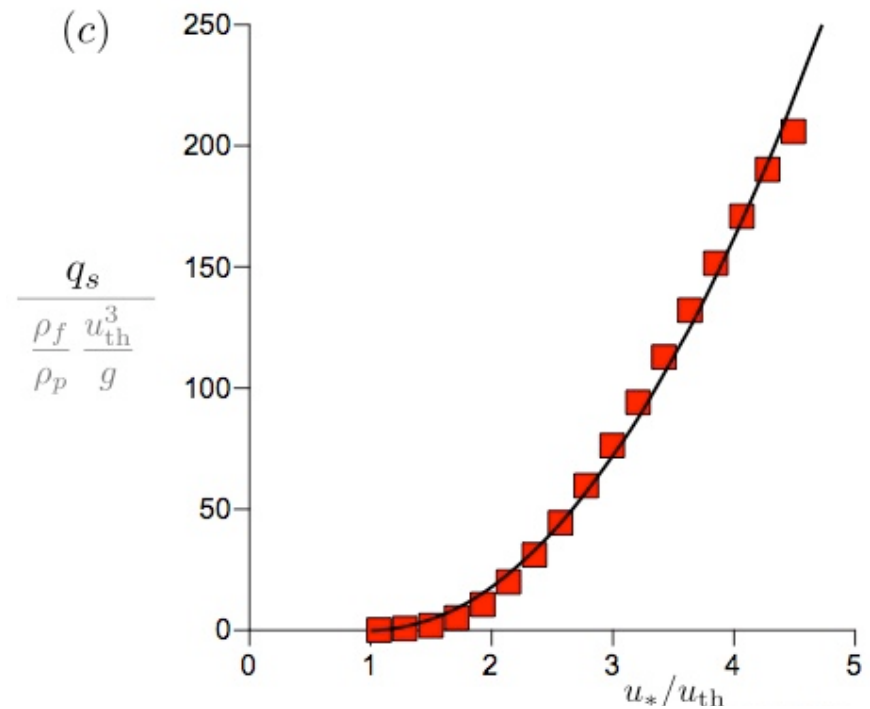
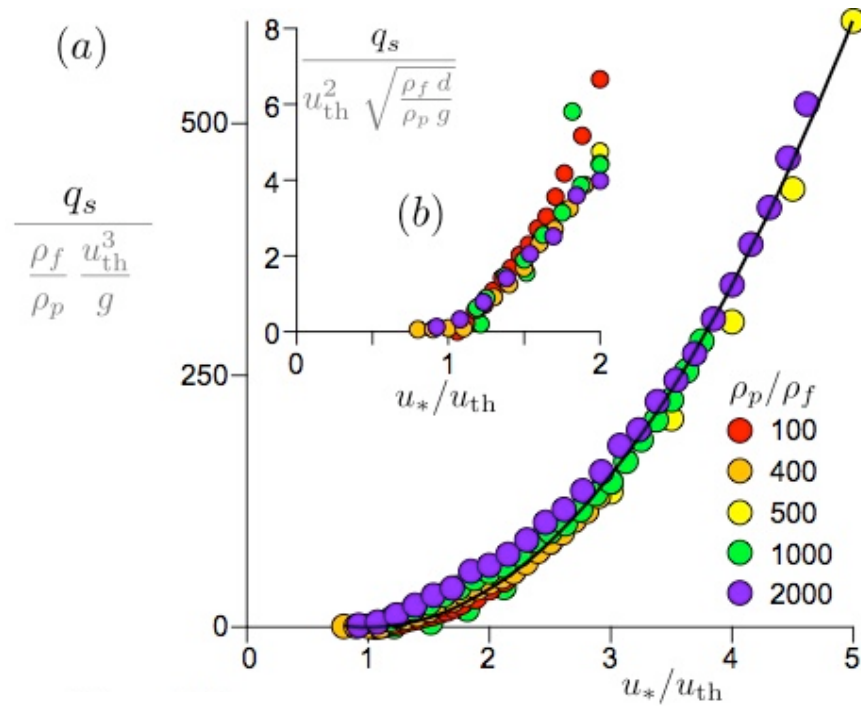
Particle velocity distribution

Aeolian transport



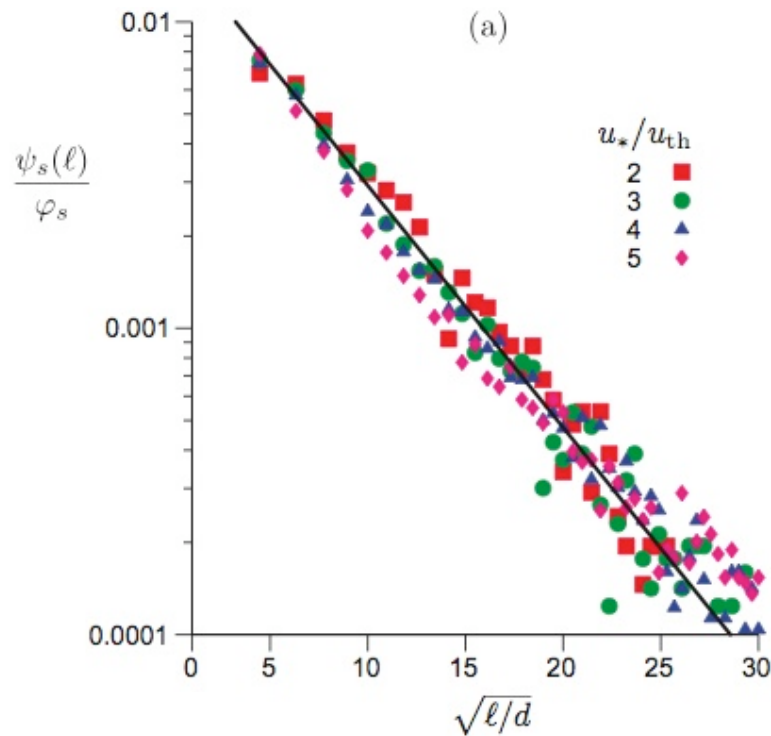
$$u/u_{th} = 2.94 (\times), 4.41 (\square), 5.88 (\circ), 7.35 (\triangle)$$

Sediment flux scaling



- Experimental data from Iversen & Rasmussen (1999)

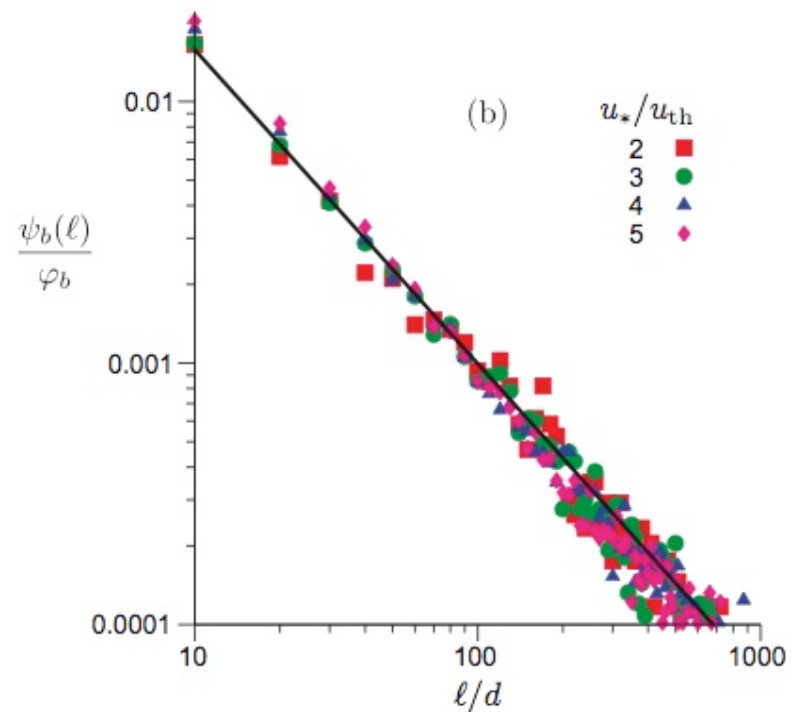
Hop length distribution



Upper transport layer

$$\psi_s(\ell) = \varphi_s \ell_s^{-1} \exp(-\sqrt{\ell/\ell_s})$$

$$\ell_s \simeq 30d$$

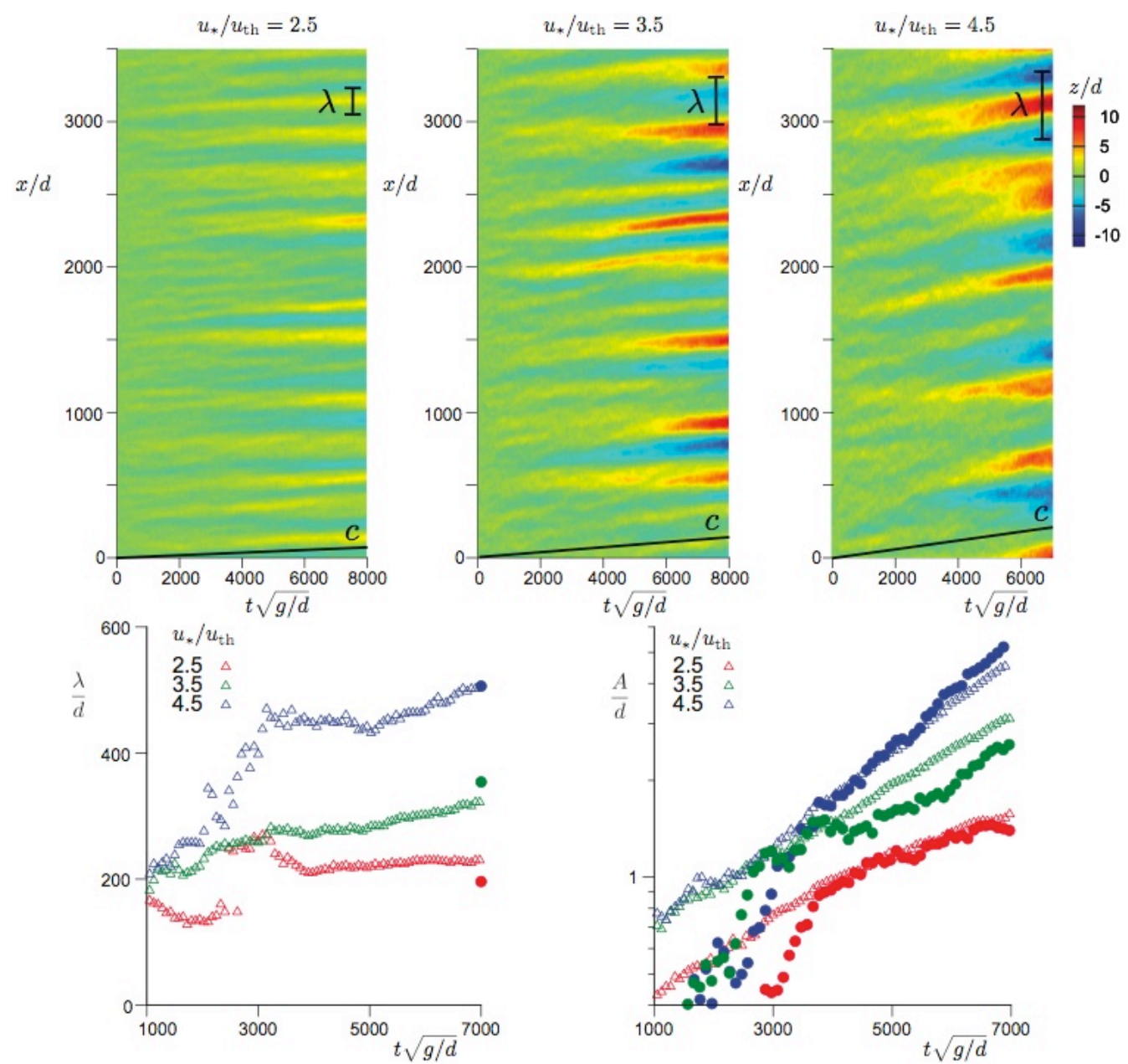


Basal collisional layer

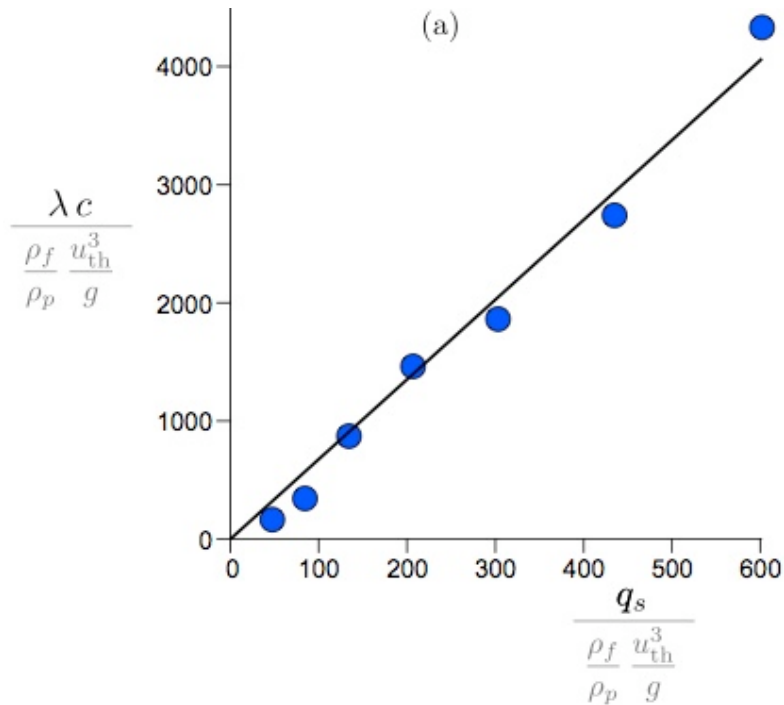
$$\psi_b(\ell) = \frac{\varphi_b}{d} \left(\frac{d}{\ell} \right)^\alpha$$

$$\alpha \simeq 1.2$$

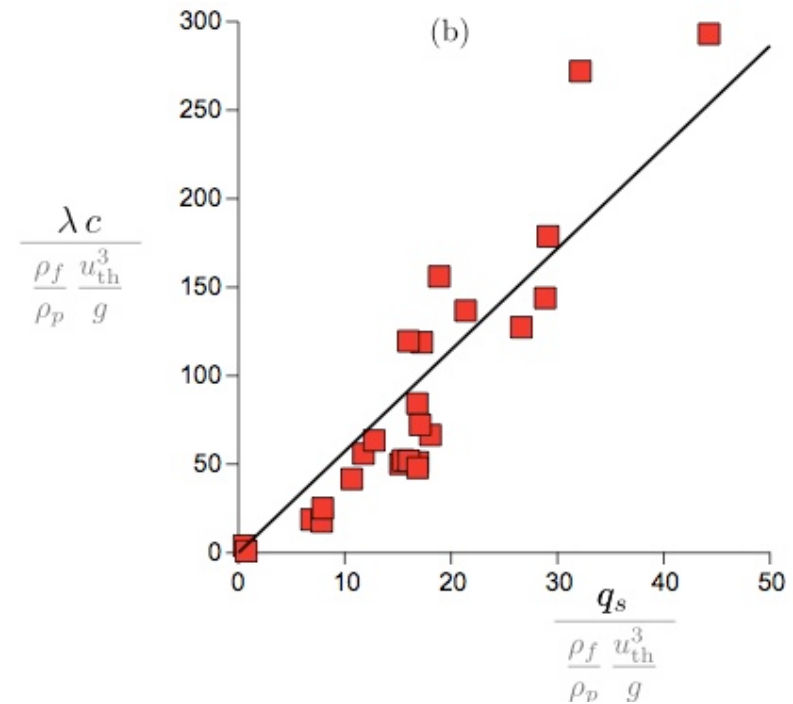
Ripple time evolution



A proxy for the sediment flux



Wind tunnel measurements



Field measurements

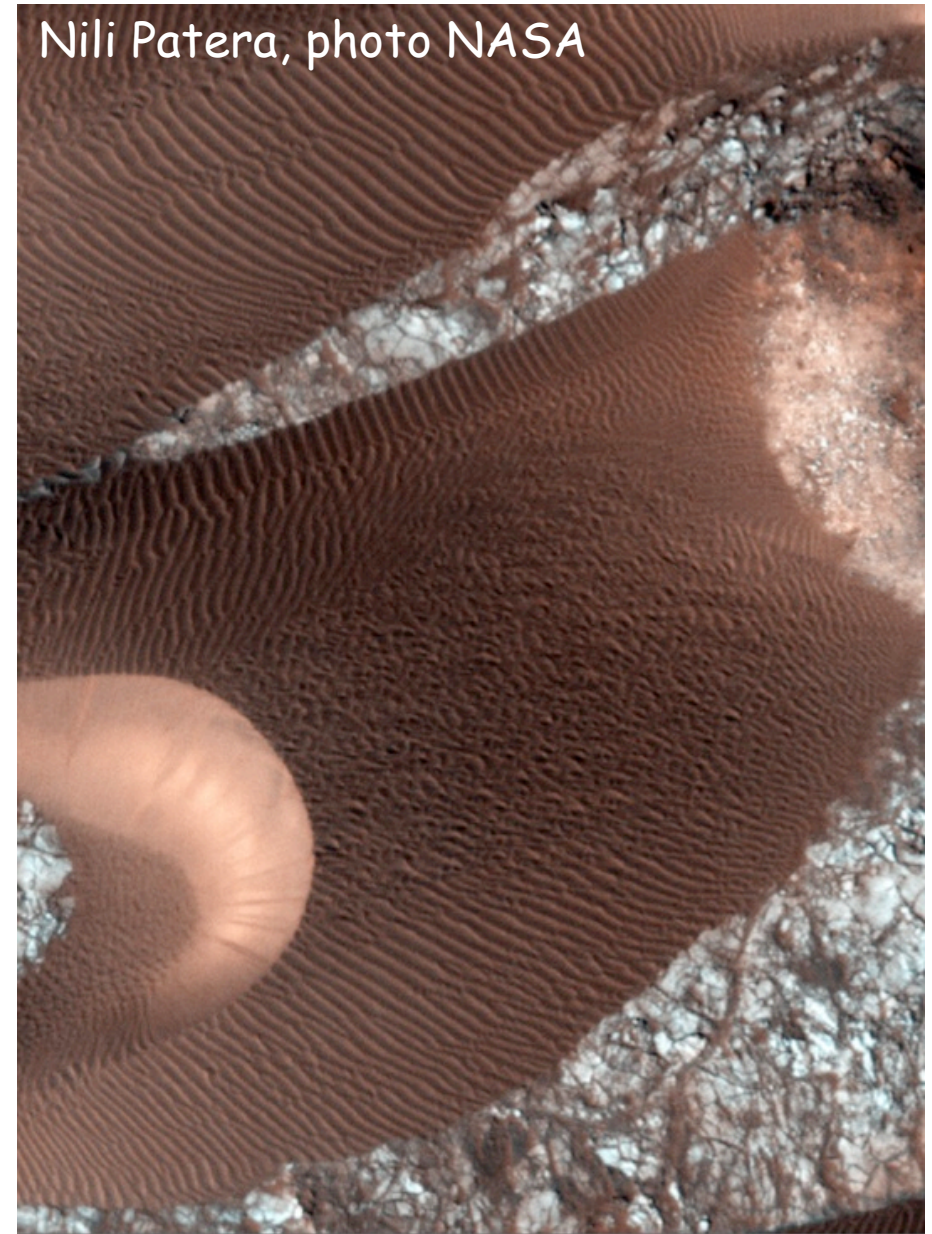
data from Andreotti et al. 2006

Ripples on Mars

Nili Patera, photo NASA

Parameter	Earth	Mars
Gravity acceleration g (m ² /s)	9.8	3.7
Atmosphere density ρ_f (kg/m ³)	1.2	$1.5\text{--}2.2 \cdot 10^{-2}$
Atmosphere viscosity ν (m/s ²)	$1.5 \cdot 10^{-5}$	$6.3 \cdot 10^{-4}$
Grain diameter d (μm)	165–185 (a)	100–200
Grain density ρ_p (kg/m ³)	2650	3000
Static to dynamic threshold ratio u_{st}/u_{th}	1.8 (b)	2.9–3.8 (b)
Grain Reynolds number \mathcal{R}	22	1–4
Mature ripple wavelength λ (m)	0.14 (c)	4–5 (d) 4.5 (e) 2–4 (f)

Numerical values of various parameters. Notes: (a) Grain size on dunes of Atlantic Sahara (Elbelrhiti et al. 2008); (b) Values from Claudin & Andreotti (2006); (c) Value for $u_*/u_{th} \simeq 1.4$ and $d = 180 \mu\text{m}$ (Andreotti et al., 2006); (d) Data from Silvestro et al. (2010); (e) Data from Bridges et al. (2012a); (f) Data from Bridges et al. (2012b).



$$\lambda \sim q_s / \varphi b$$