

Dynamics of the current eruption of Mount St. Helens: volcanic extrusion regulated by granular friction

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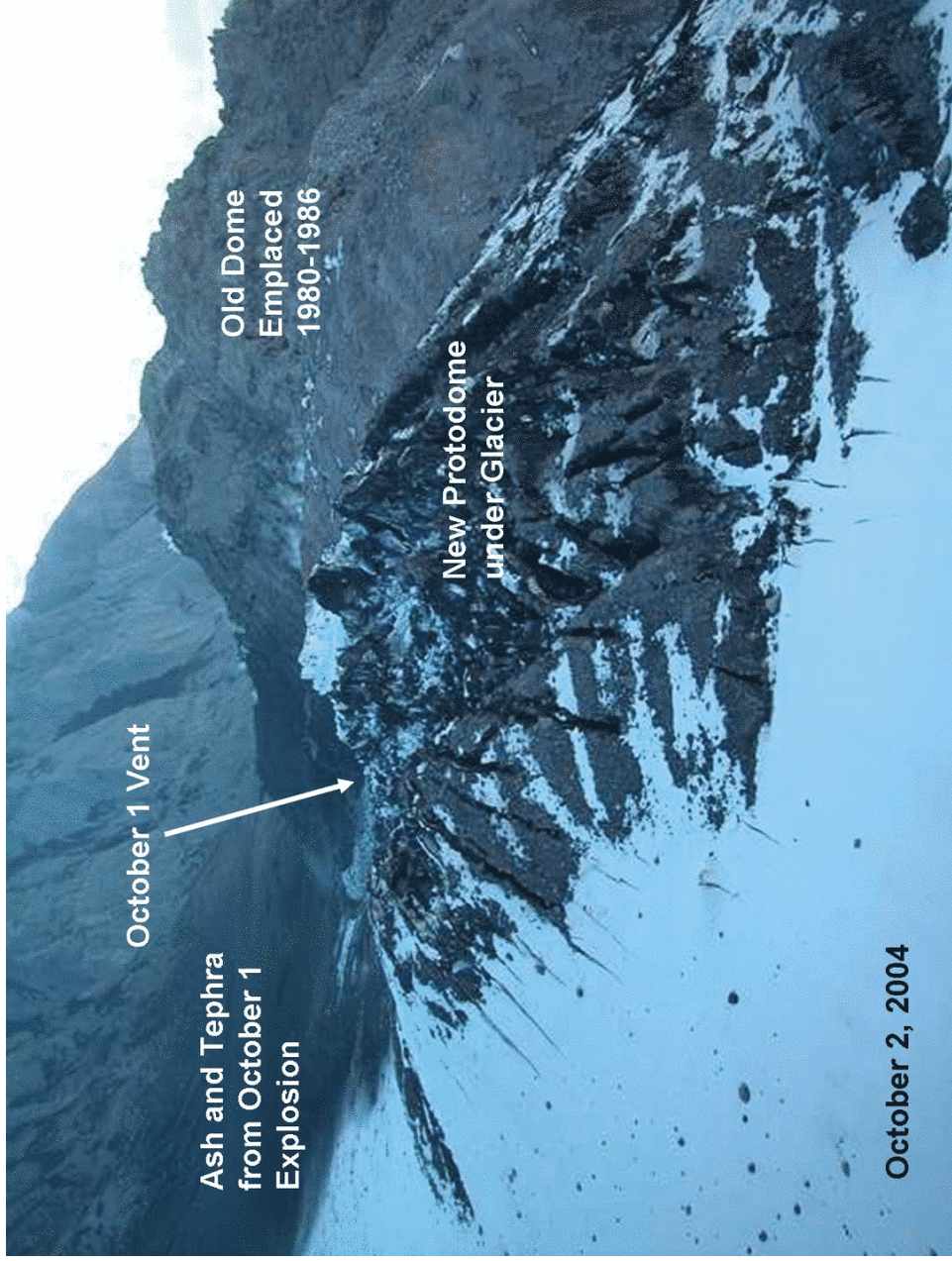
April 26, 2005

Onset of the Eruption

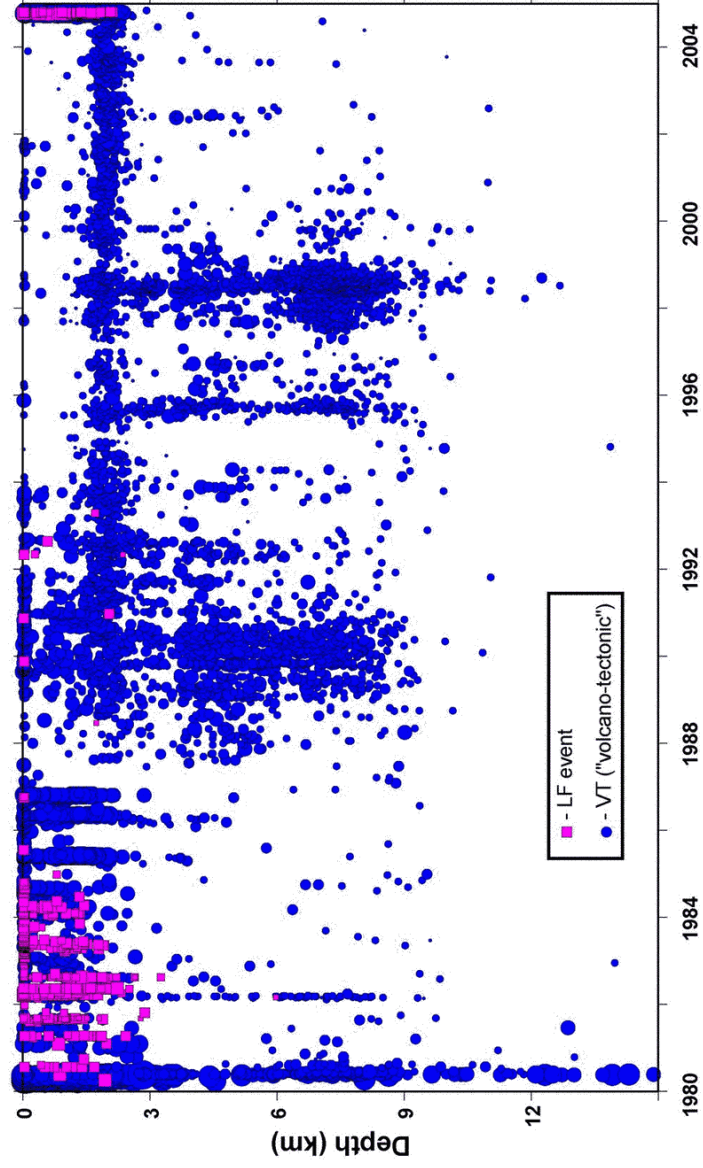
October 1, 2004



Old Dome
Emplaced
1980-1986

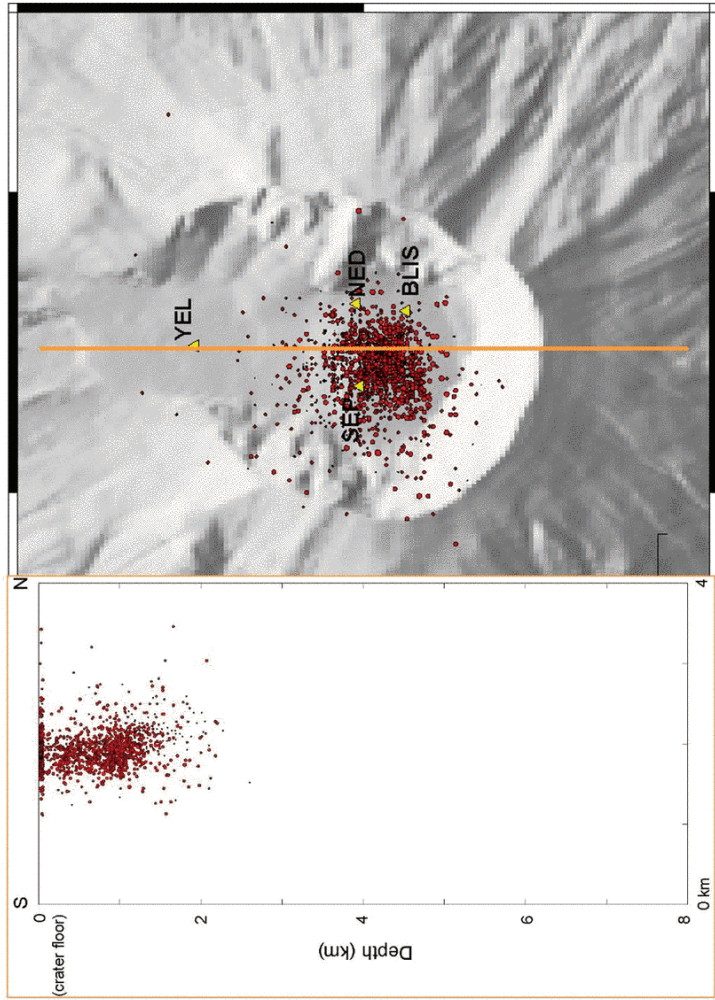


Mount St. Helens seismicity, 1980-2004

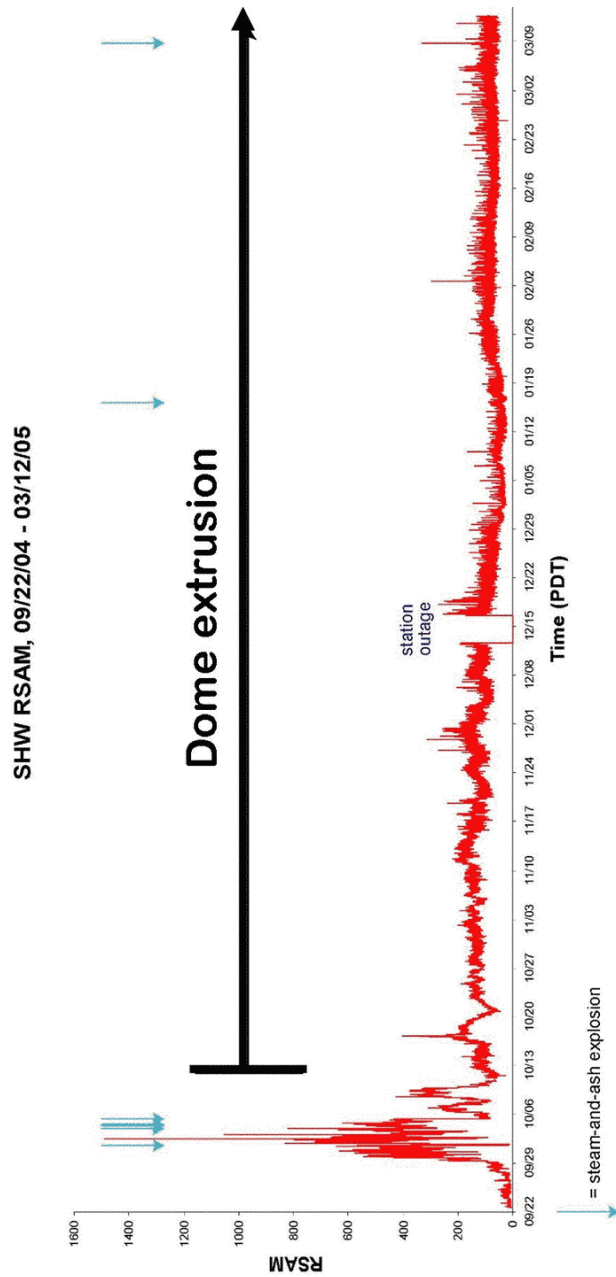


➤ 15,303 located eqs through 09/2004

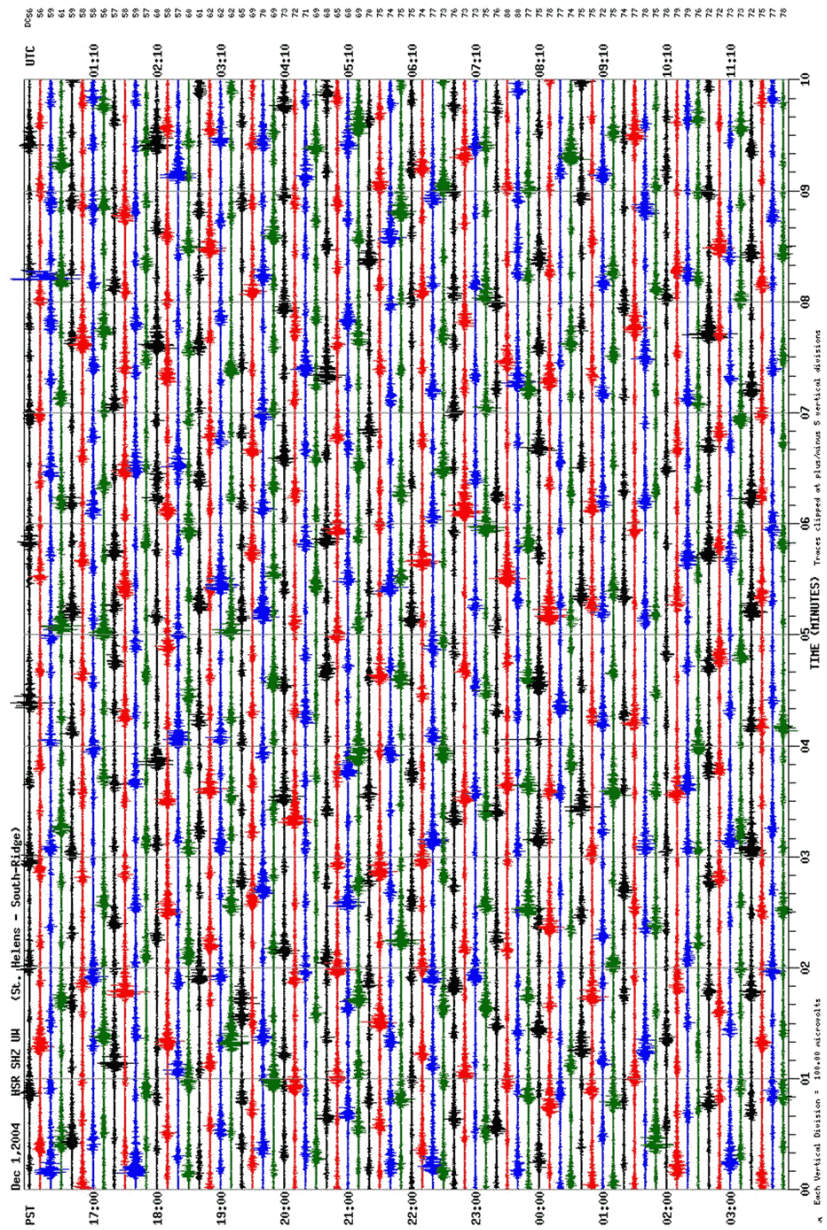
Sept.– Dec. 2004 Mount St. Helens Seismicity distribution



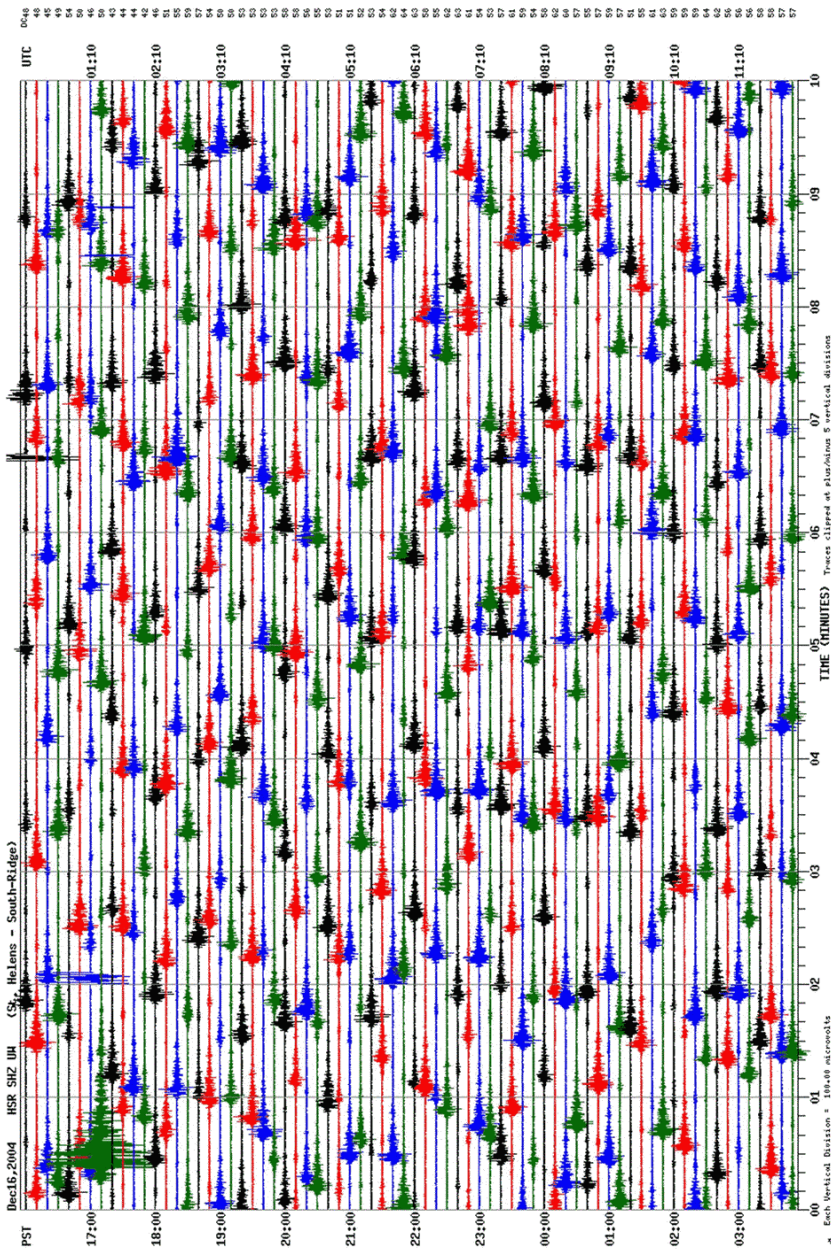
Mount St. Helens seismic intensity, 9/2004 to 3/2005

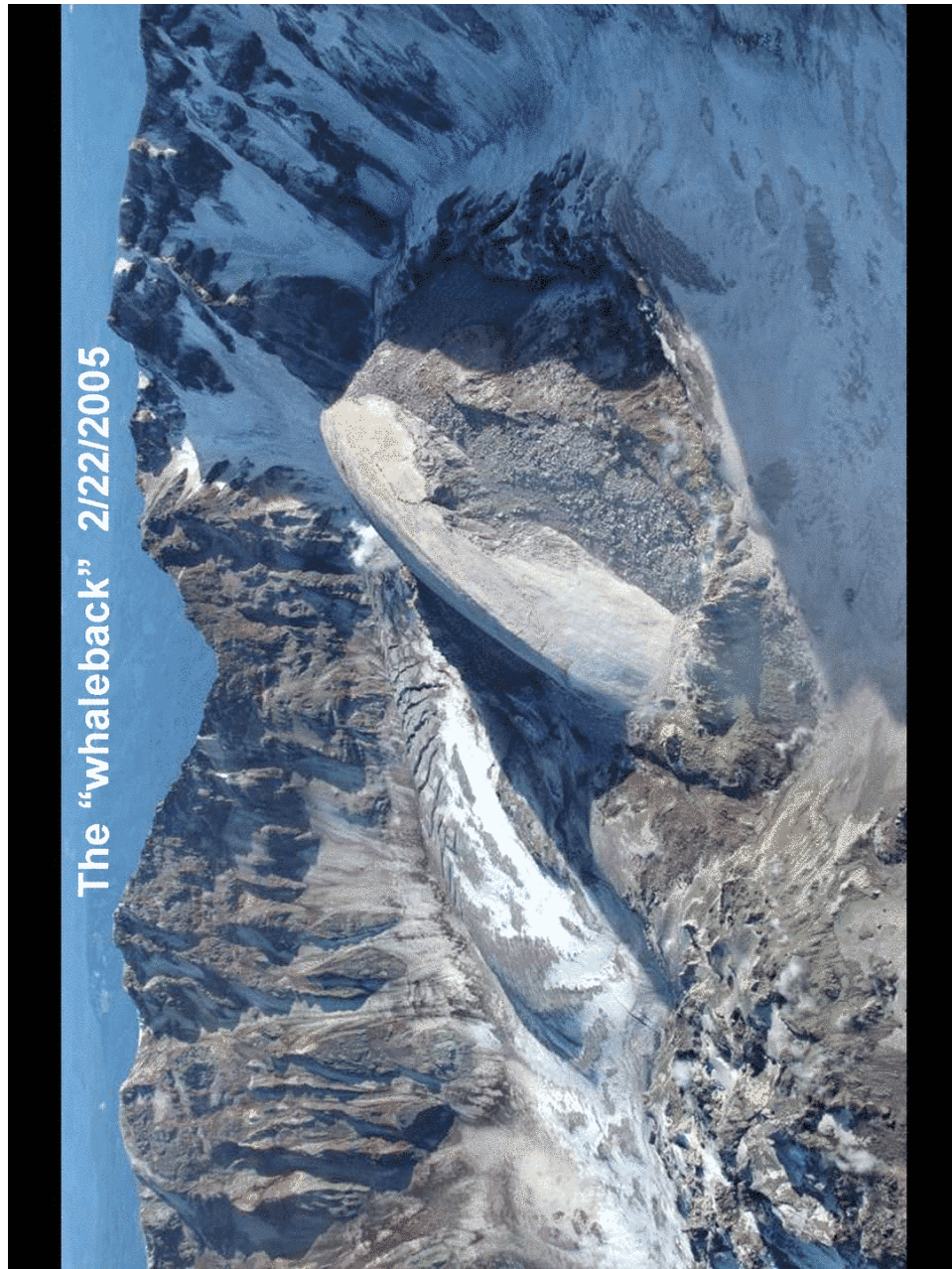


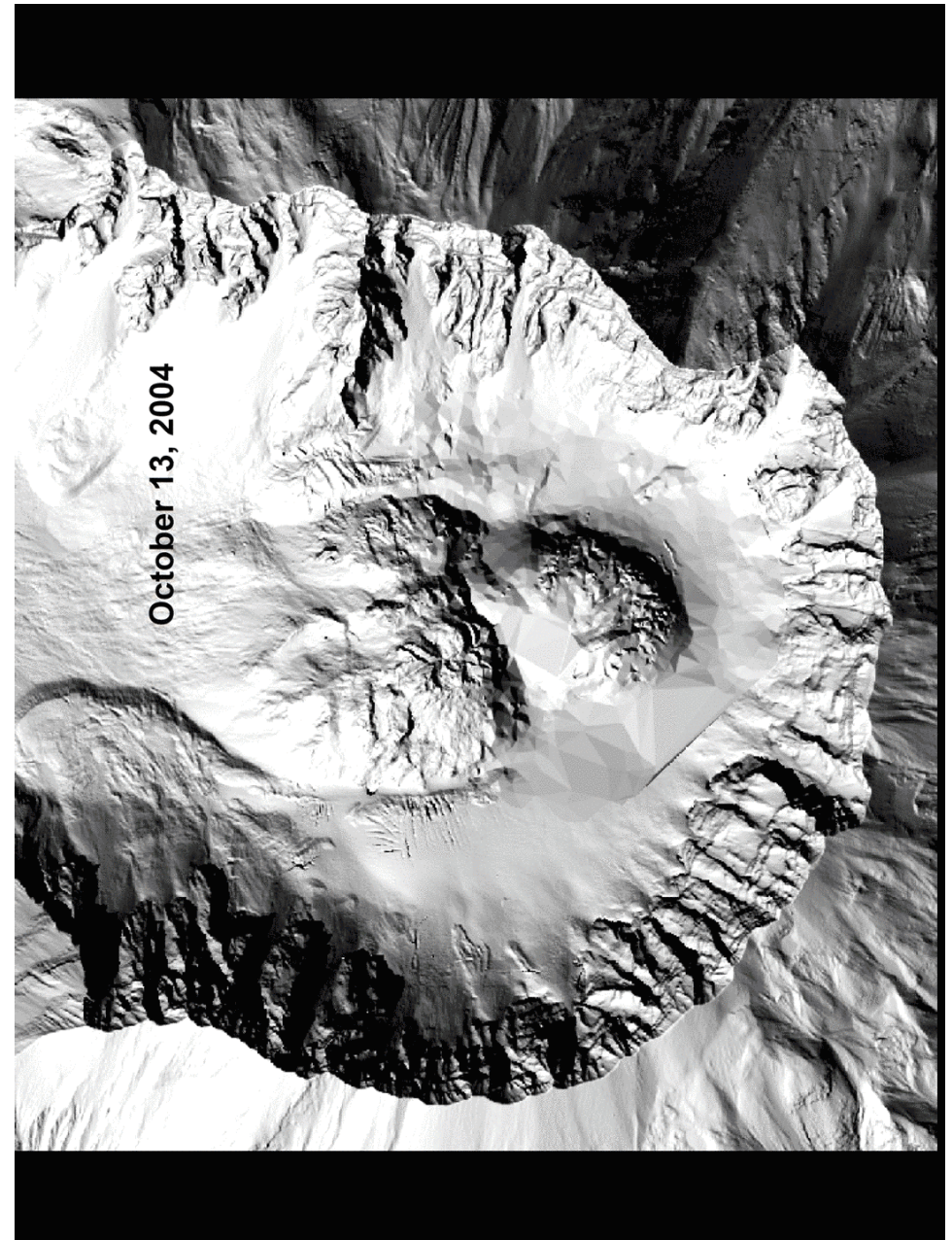
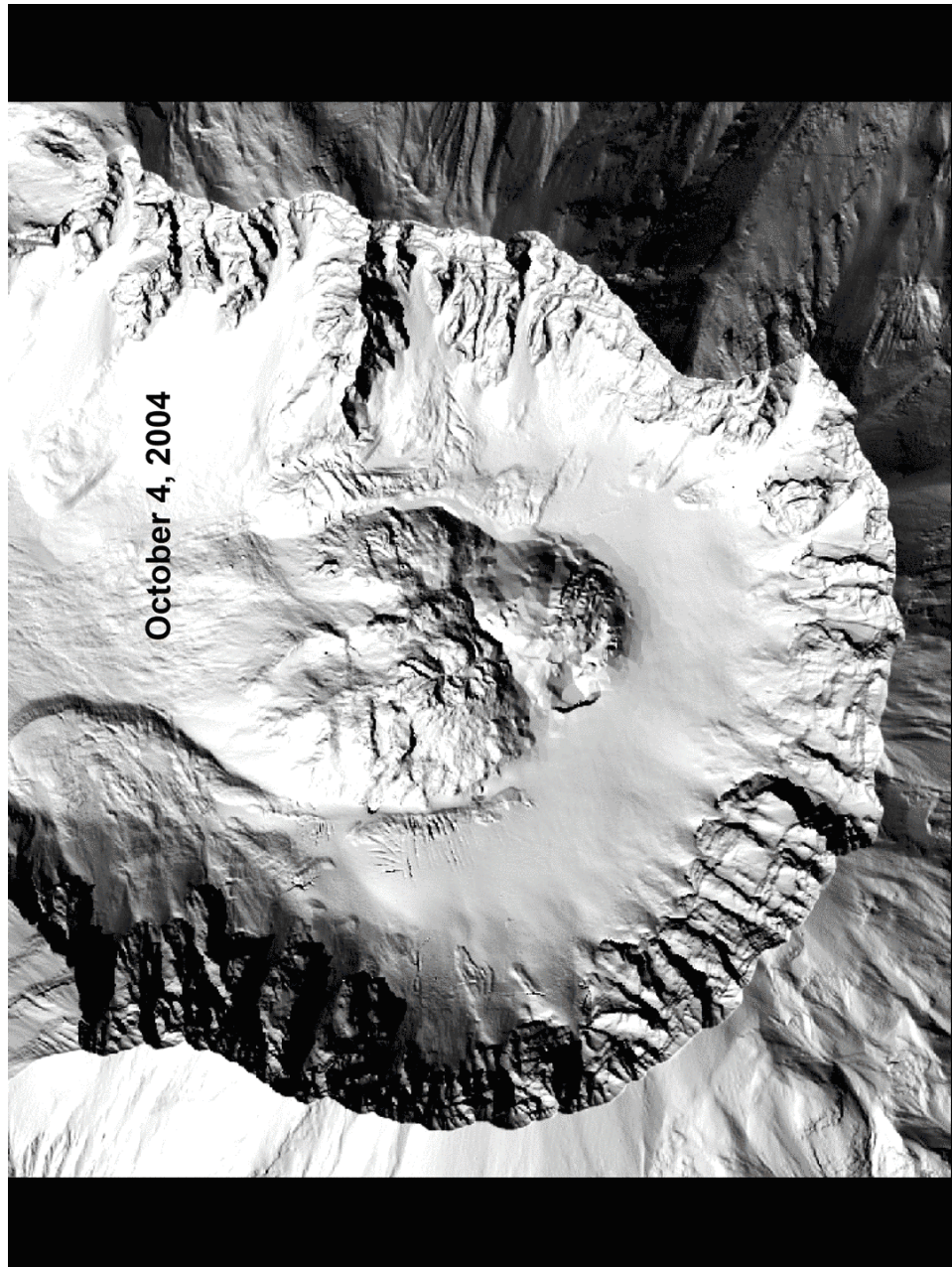
Repetitive “drumbeat” earthquakes, > 1 million to date

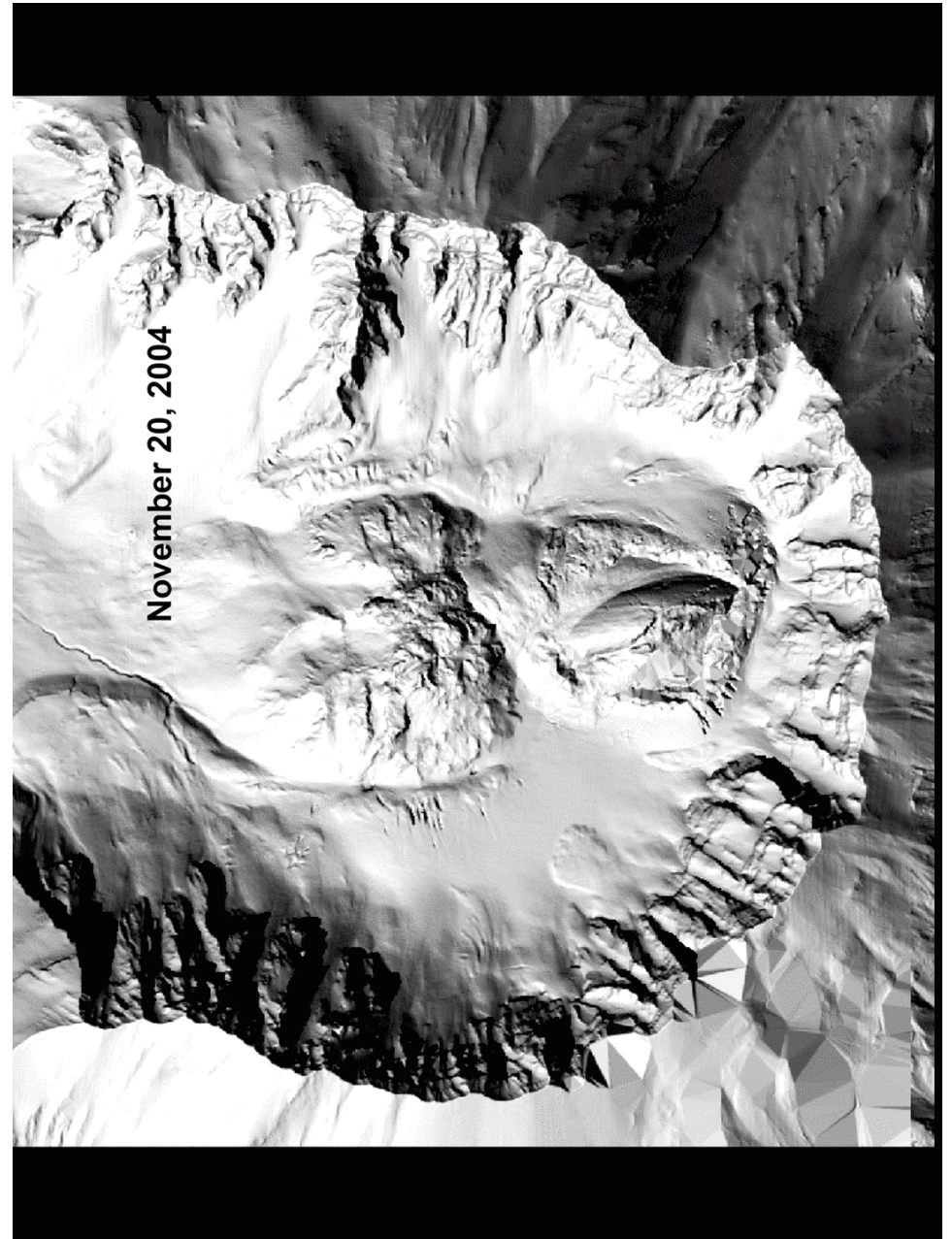


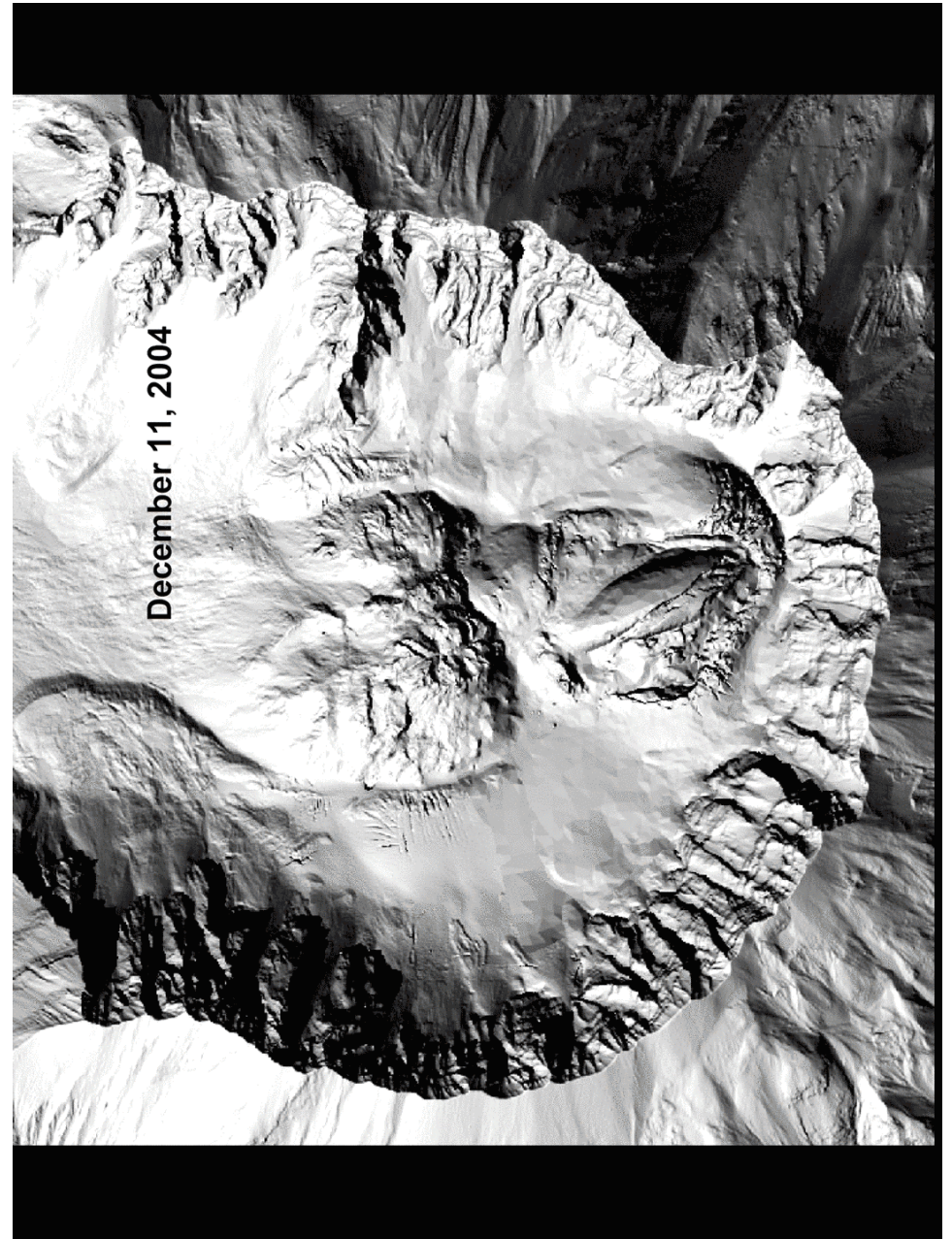
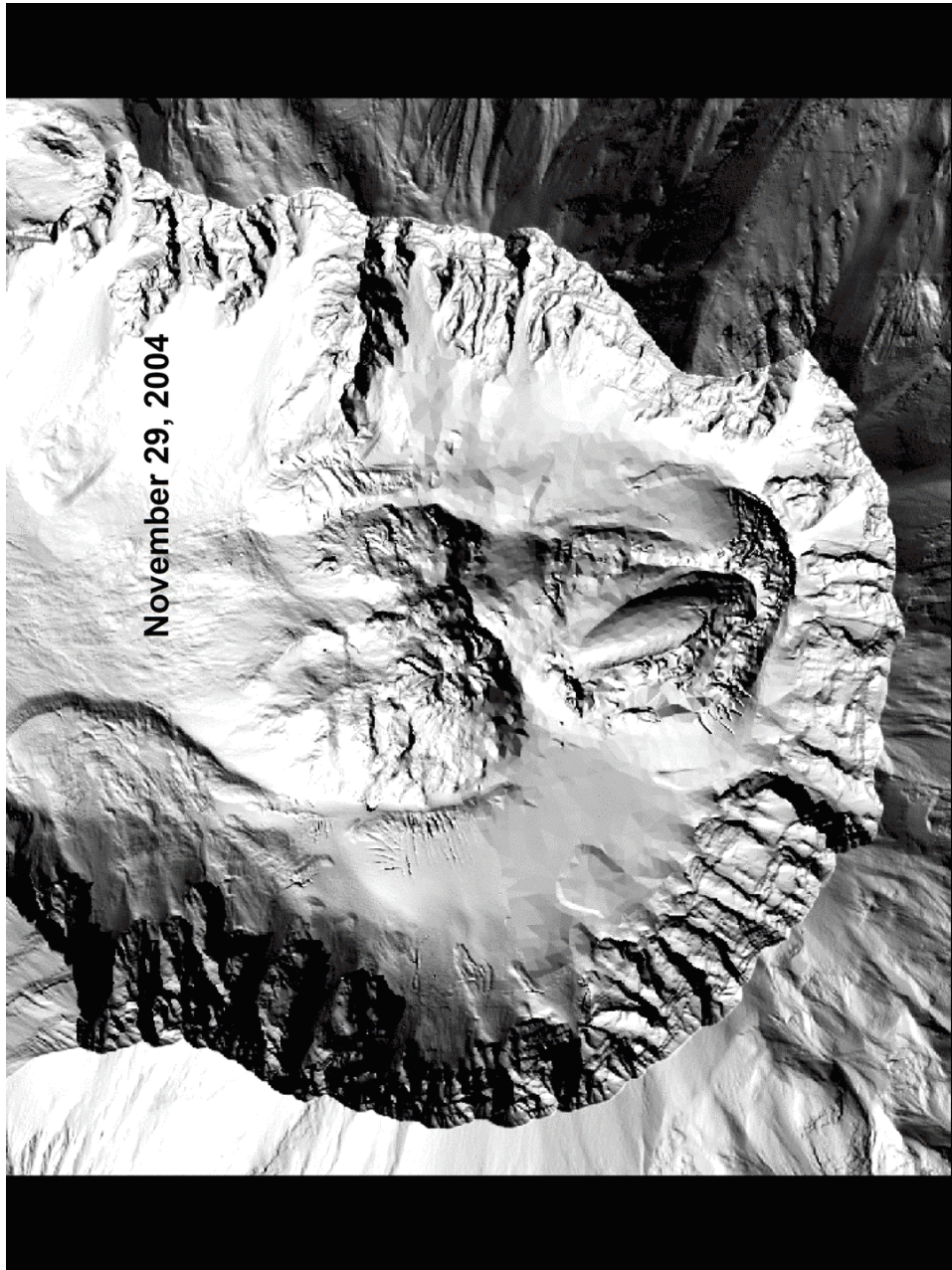
Drumbeats change only subtly with time

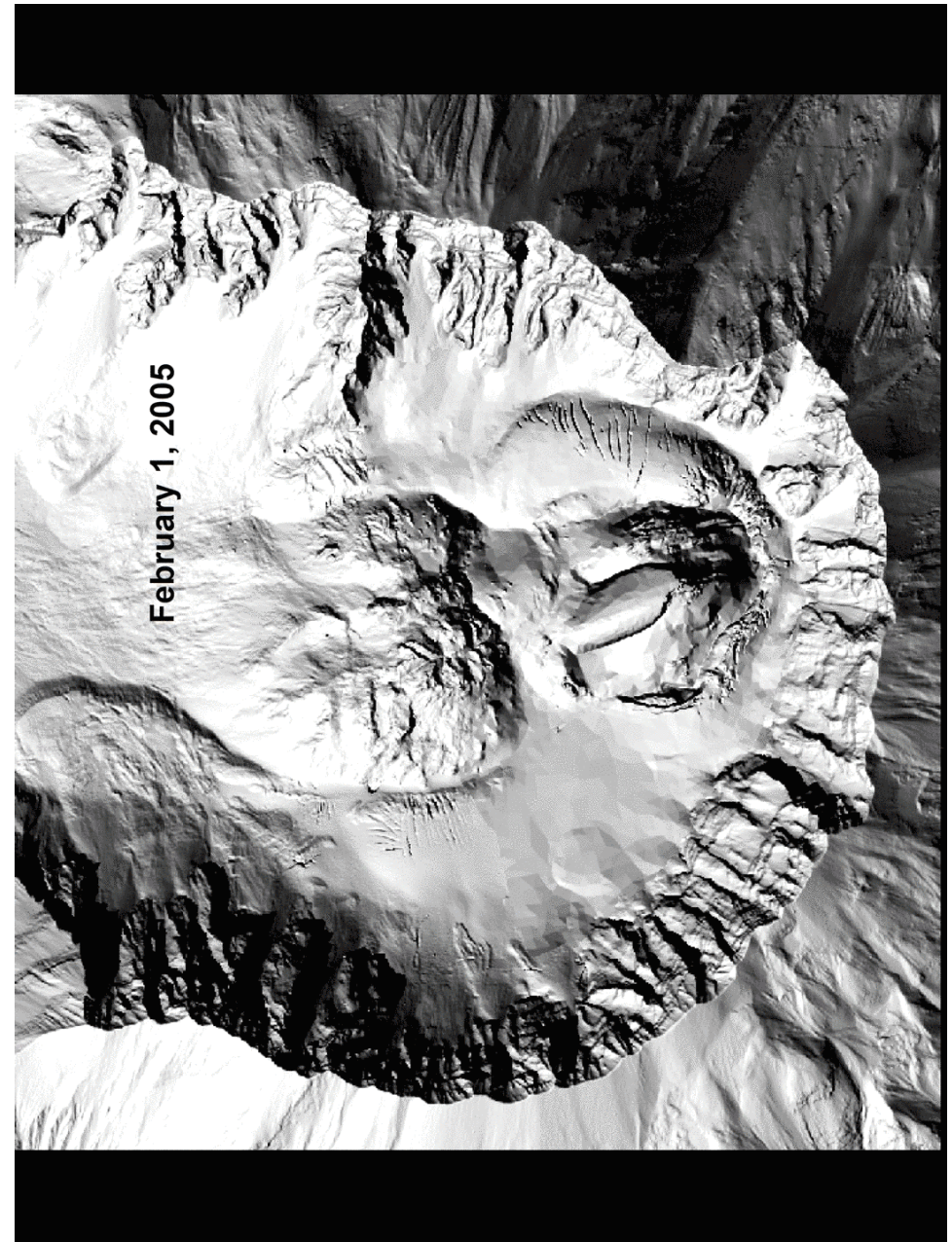
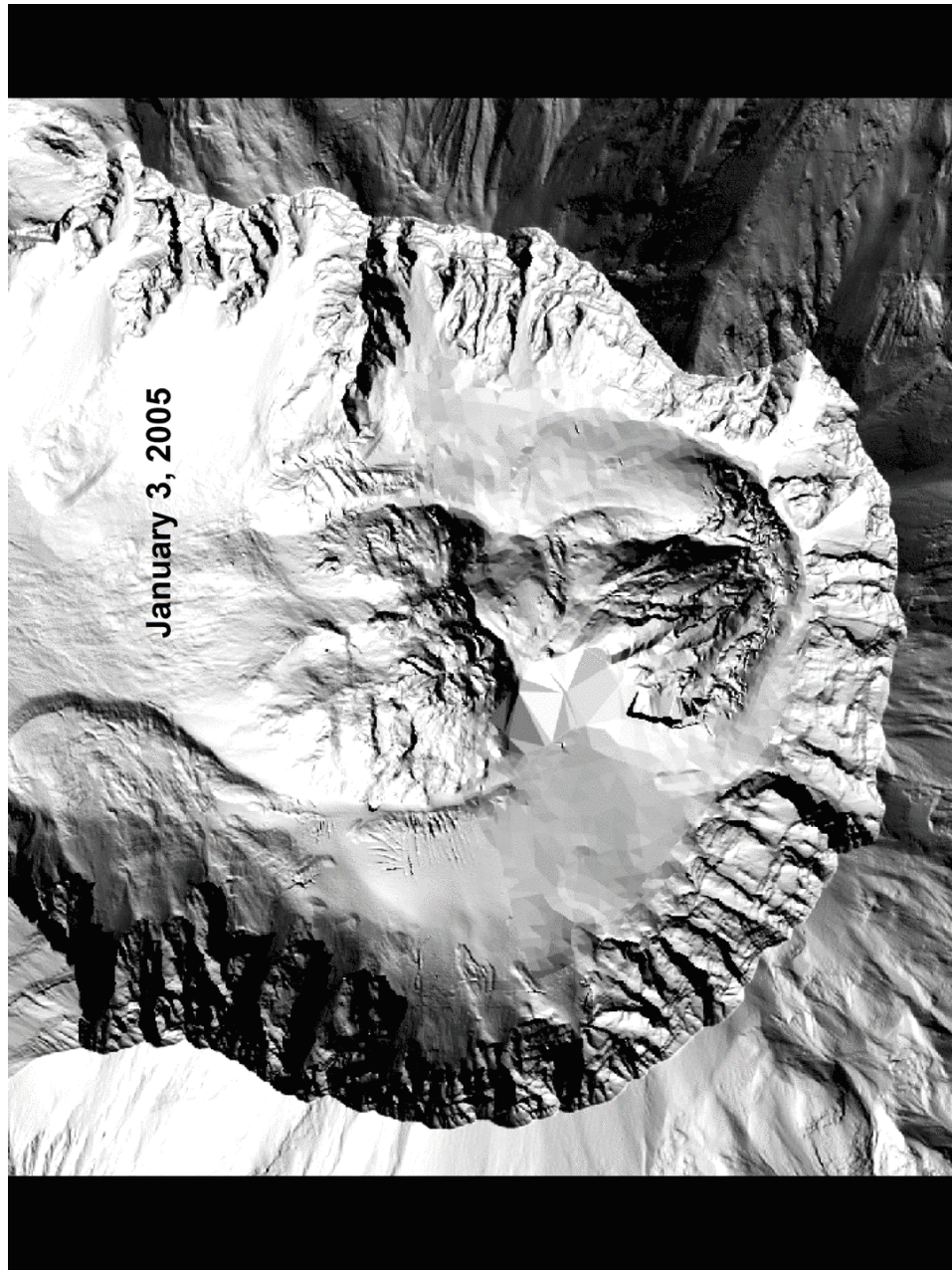


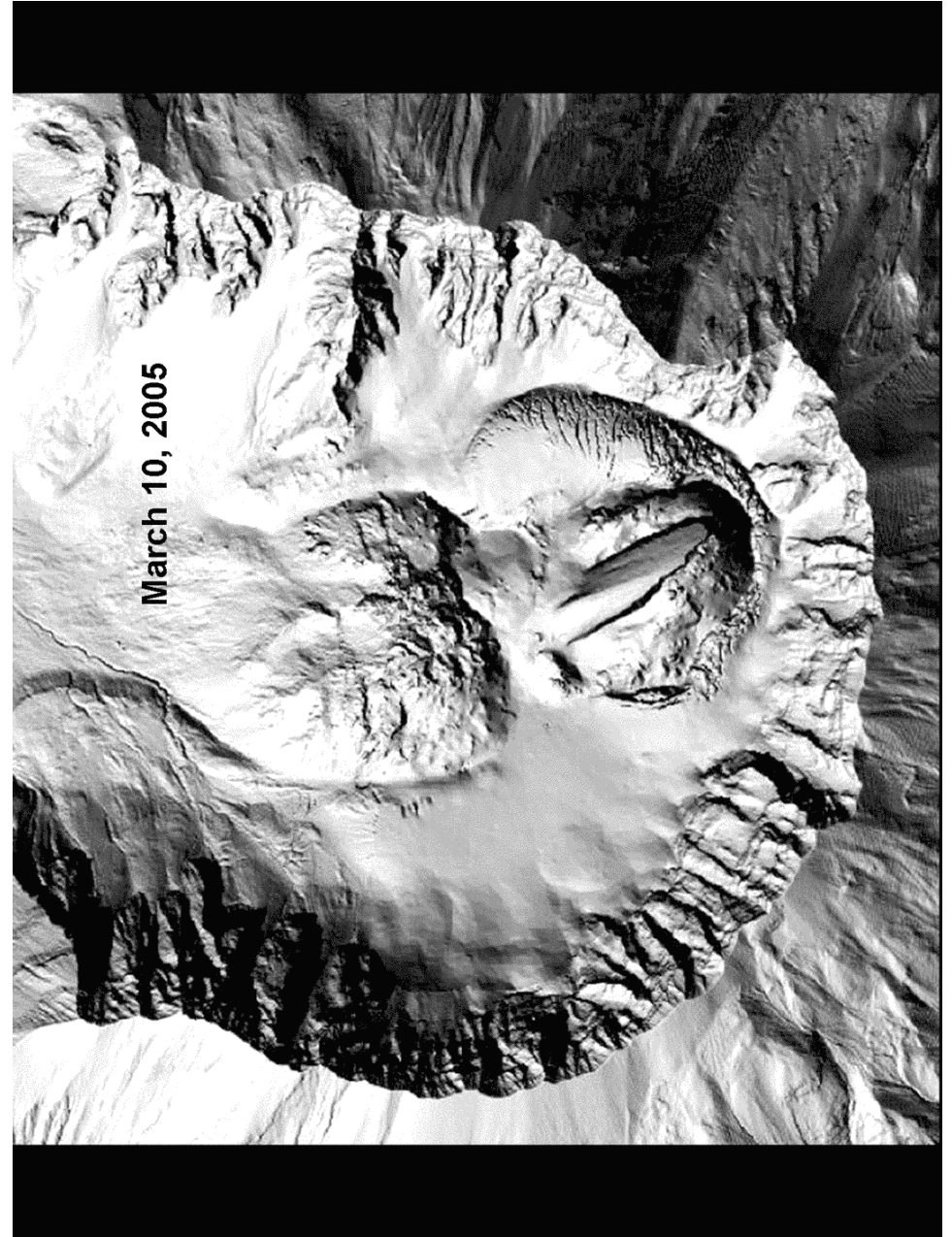
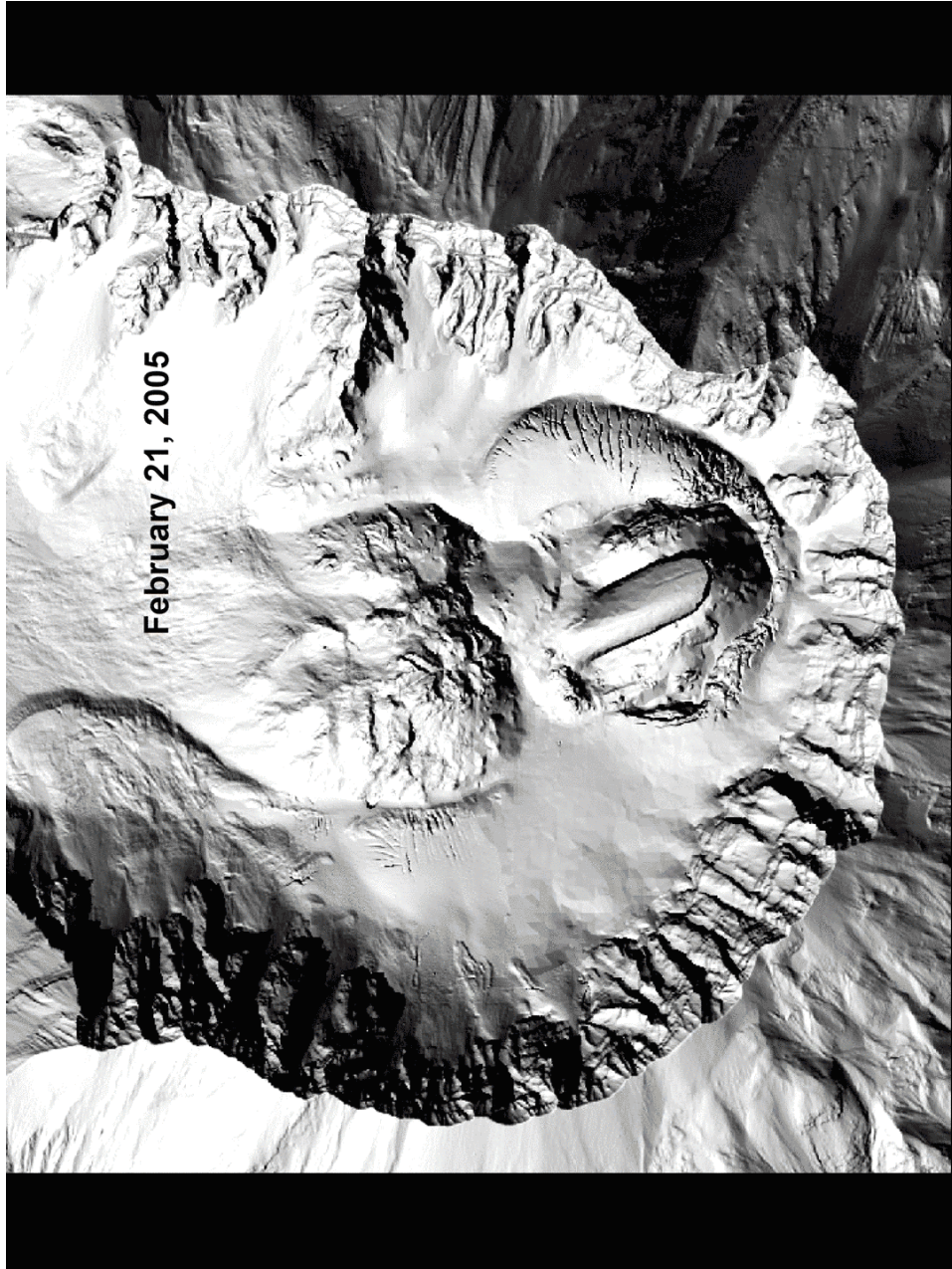


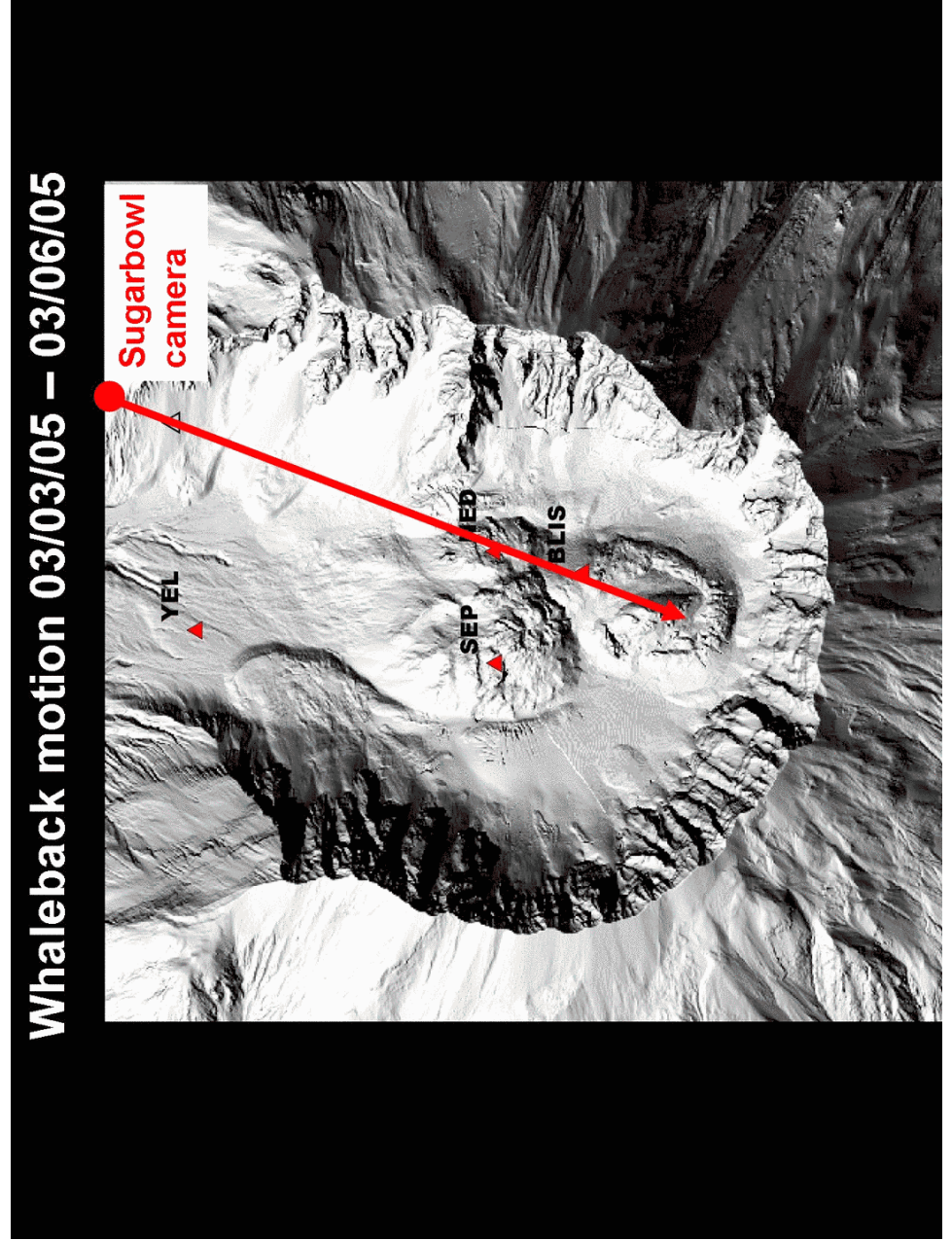
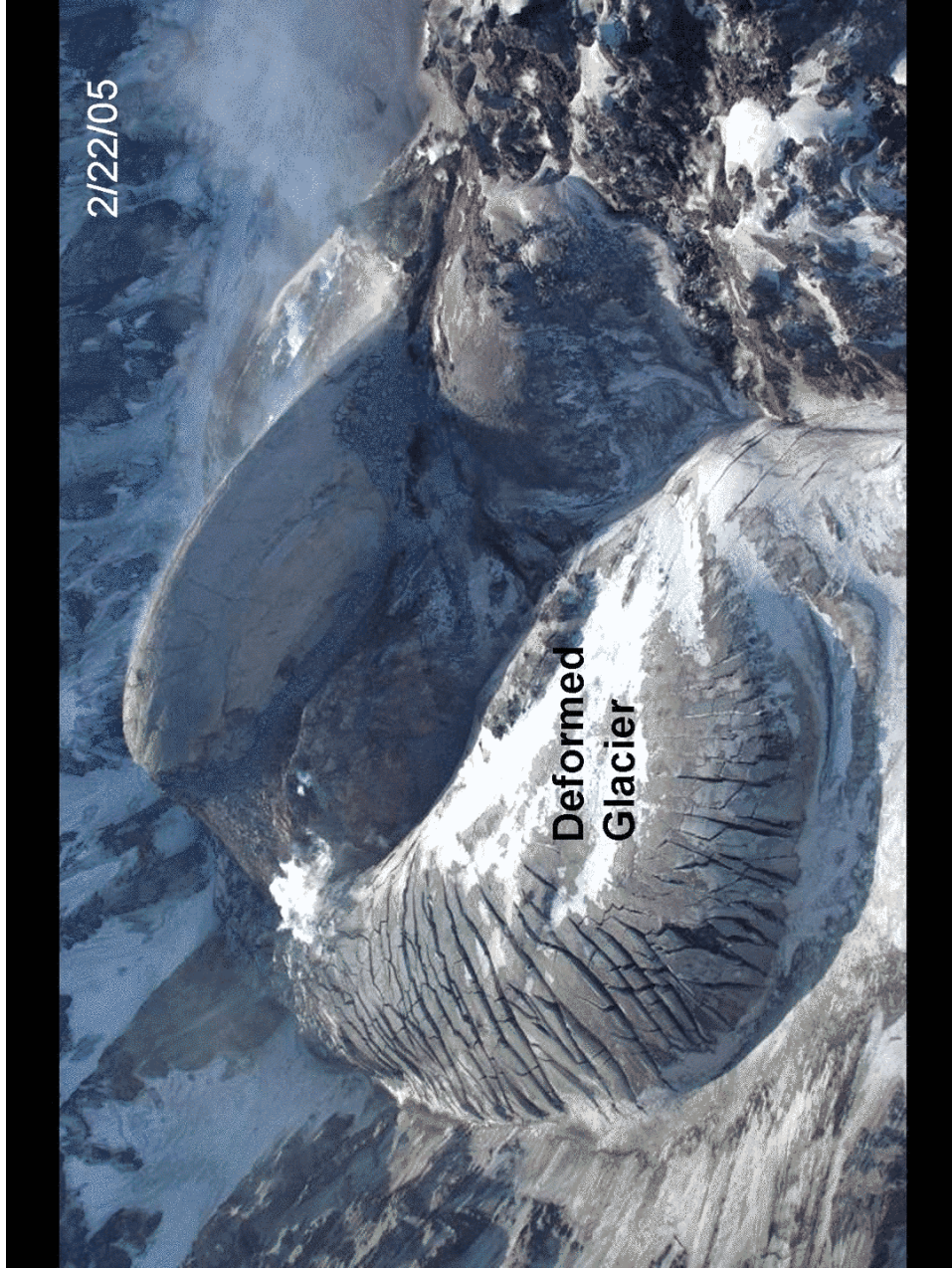




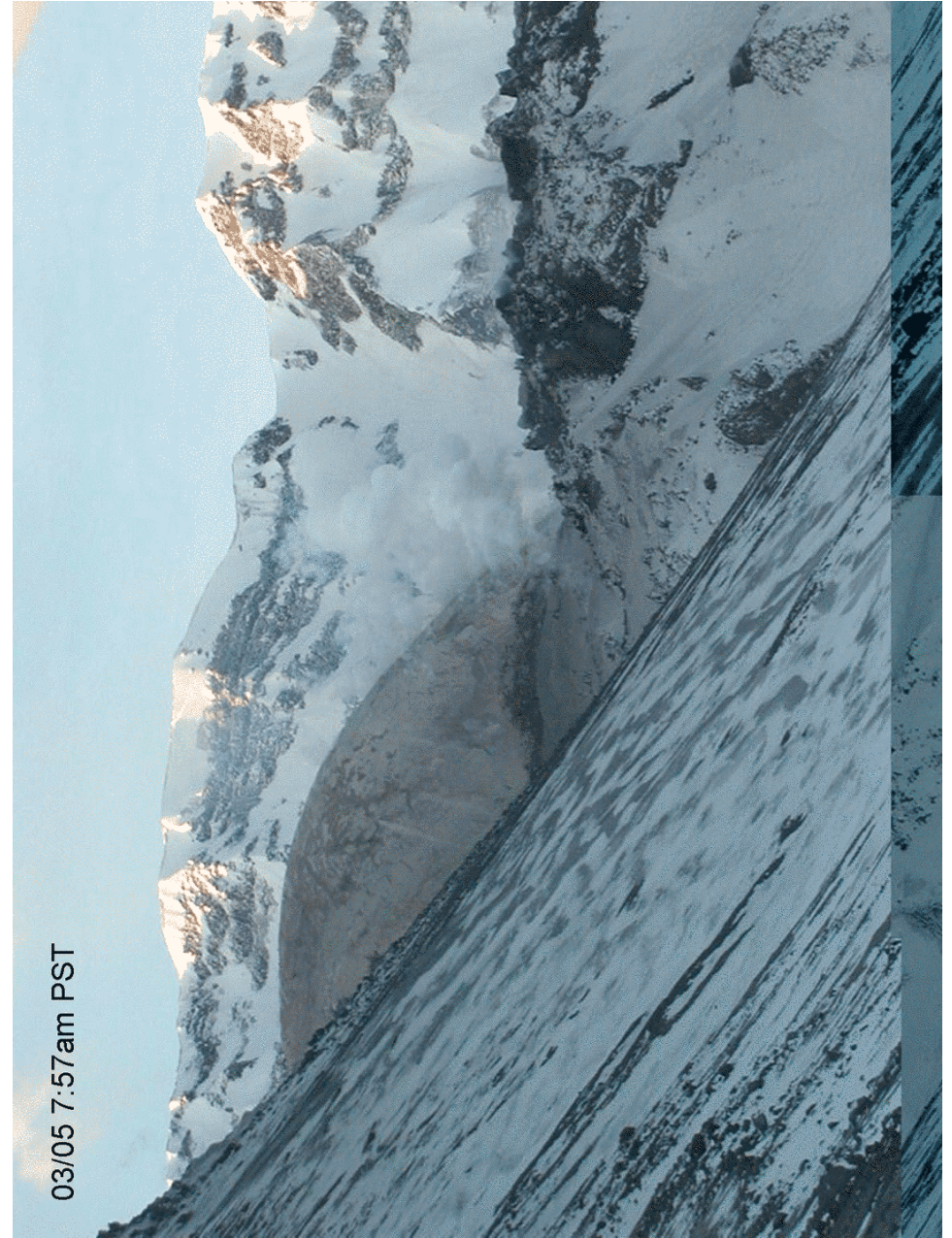




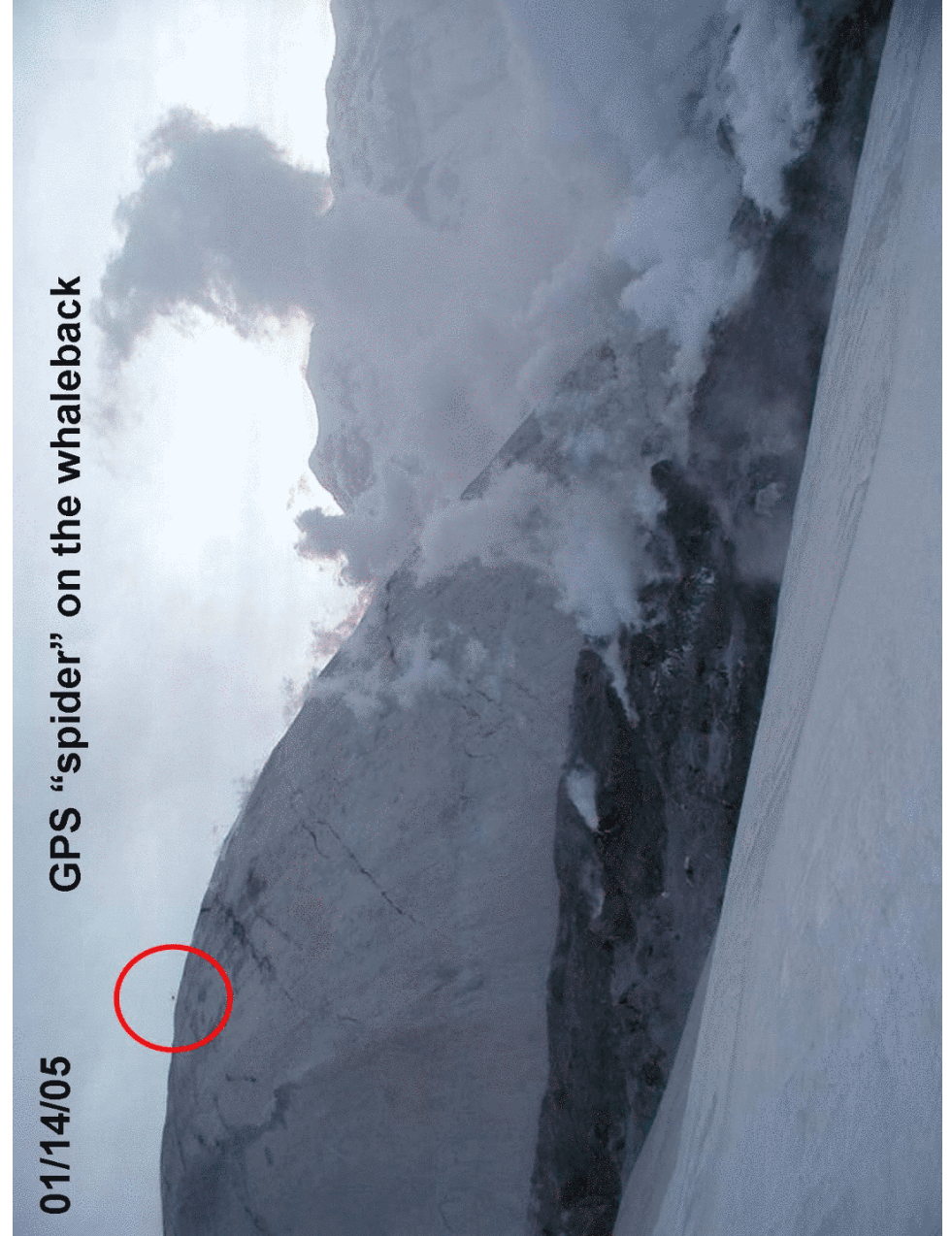






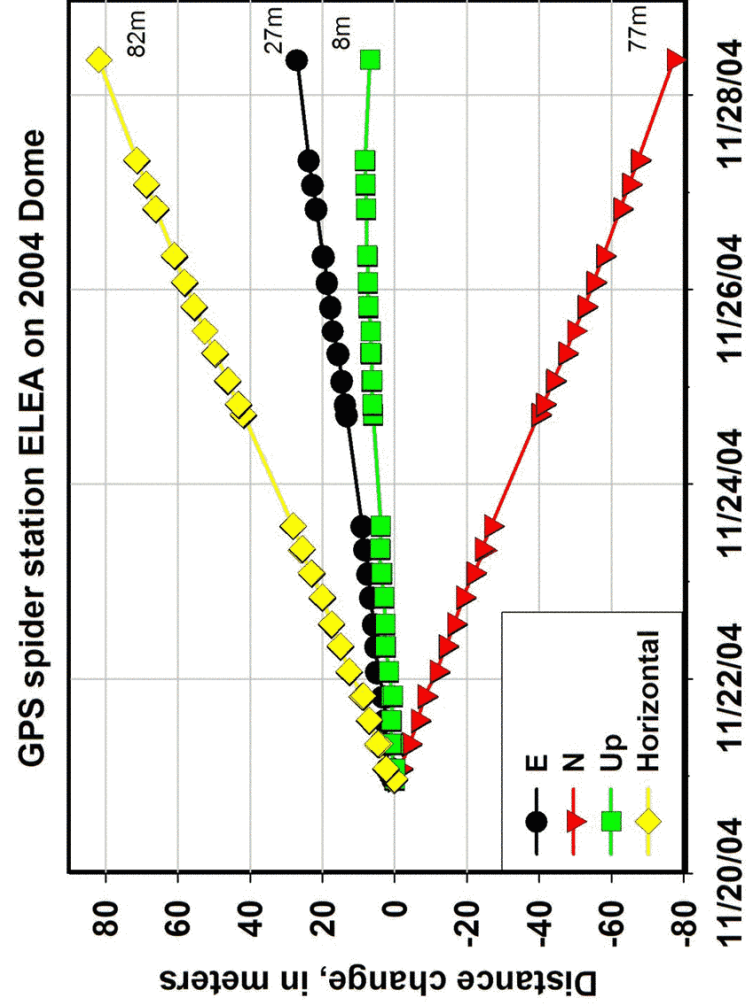




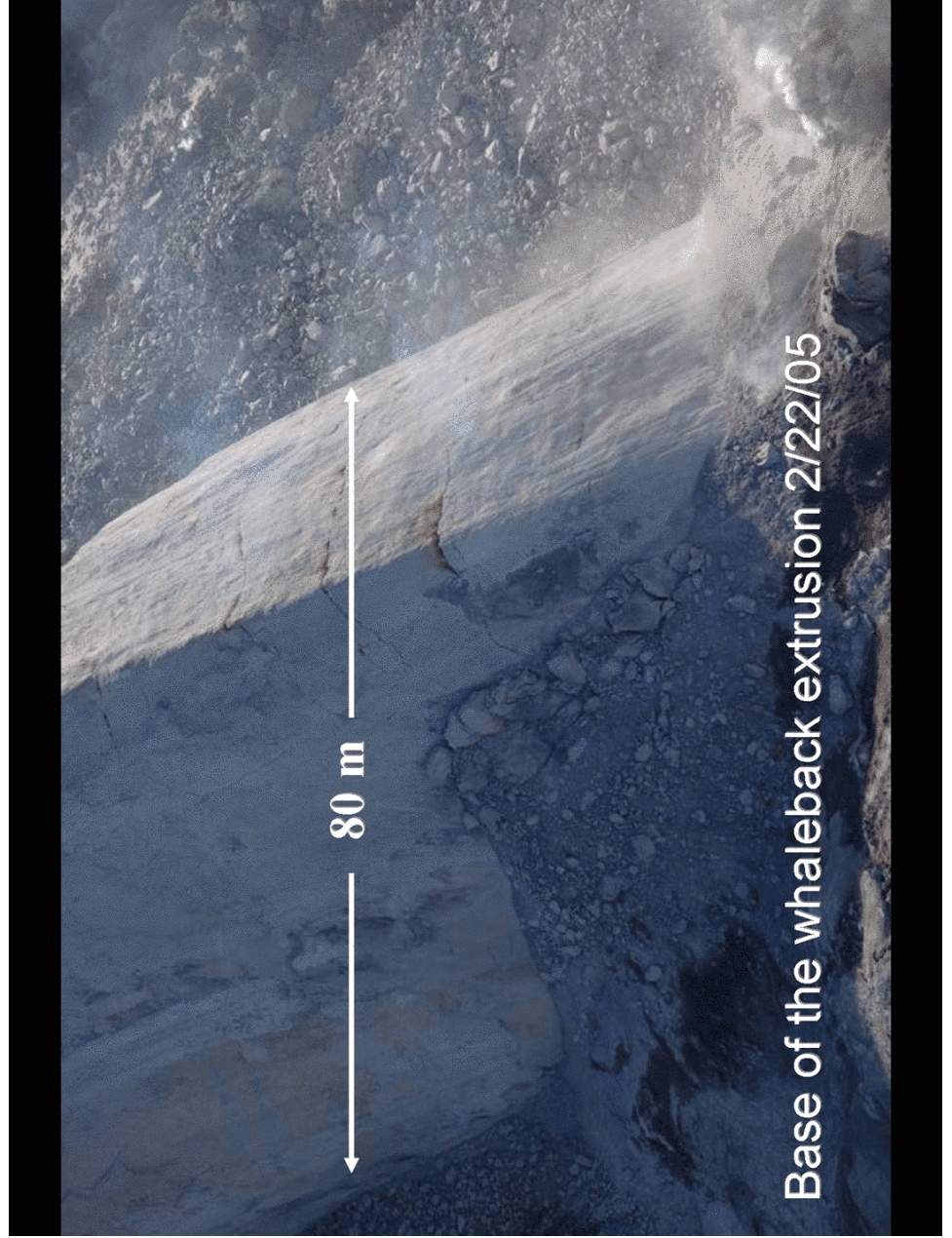
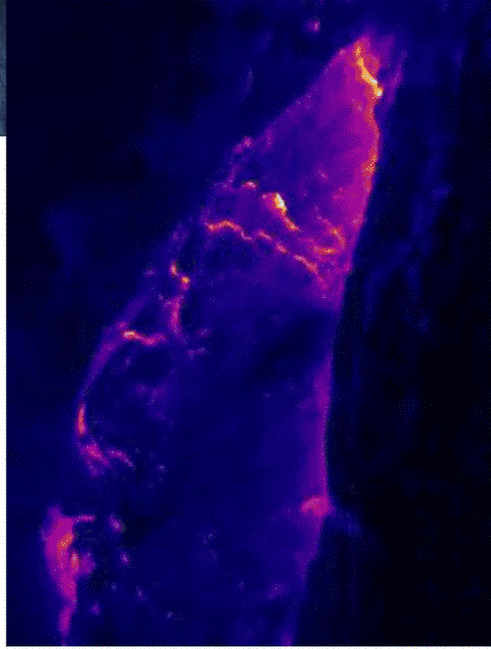
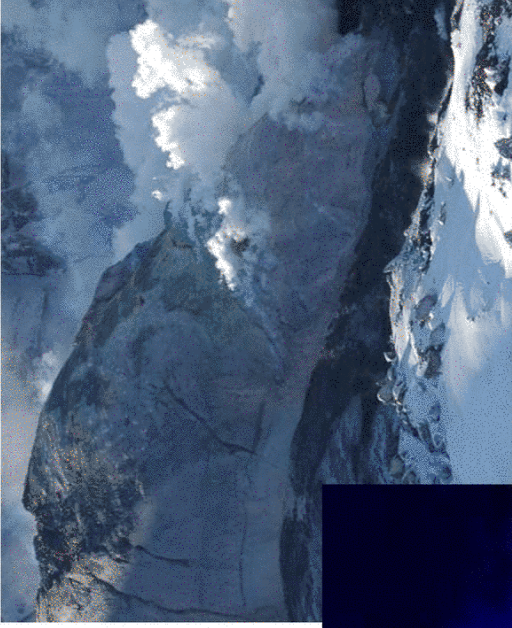


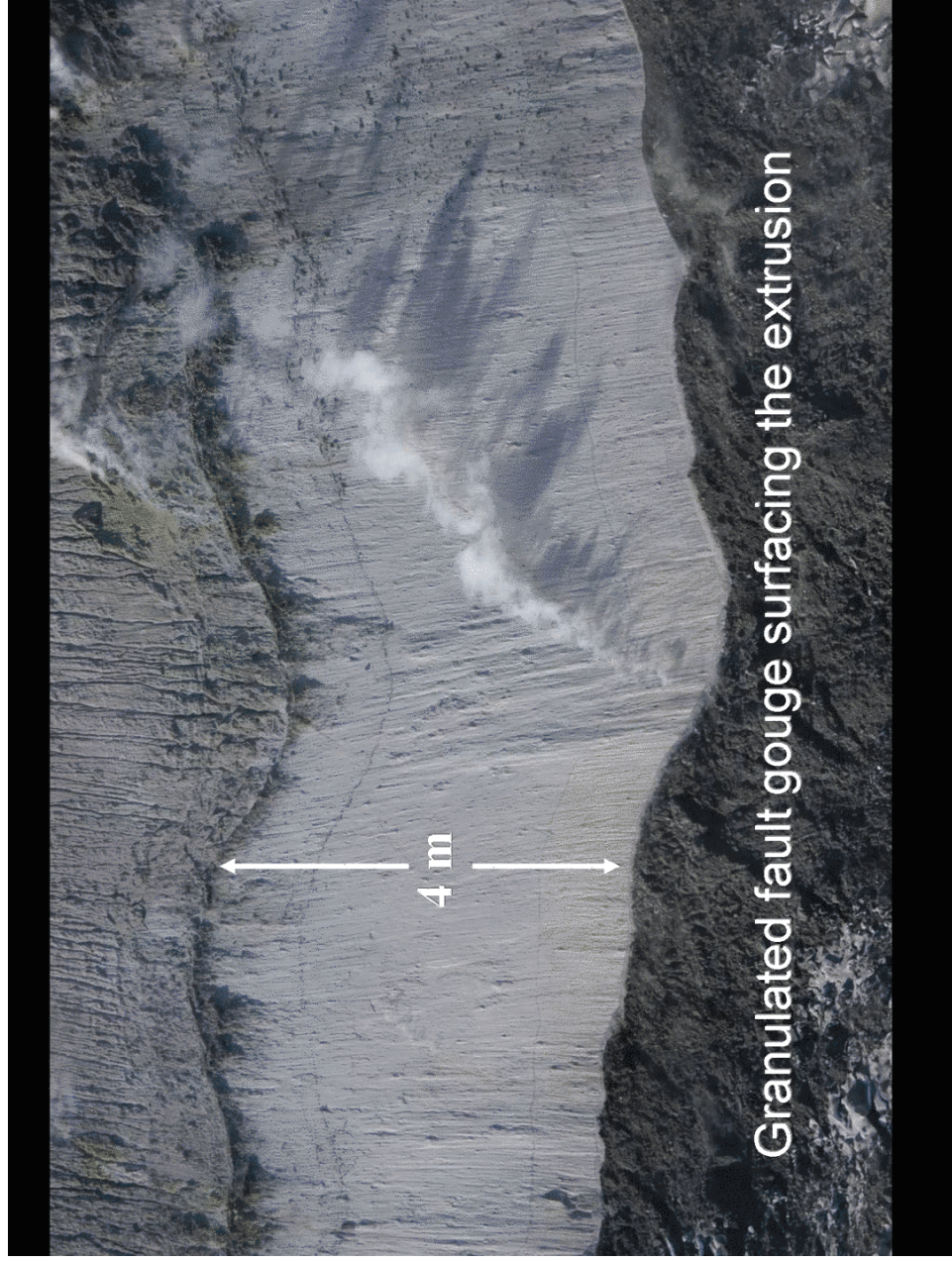


Whaleback motion 11/21/04 – 11/29/04



FLIR (Forward Looking Infrared) imagery

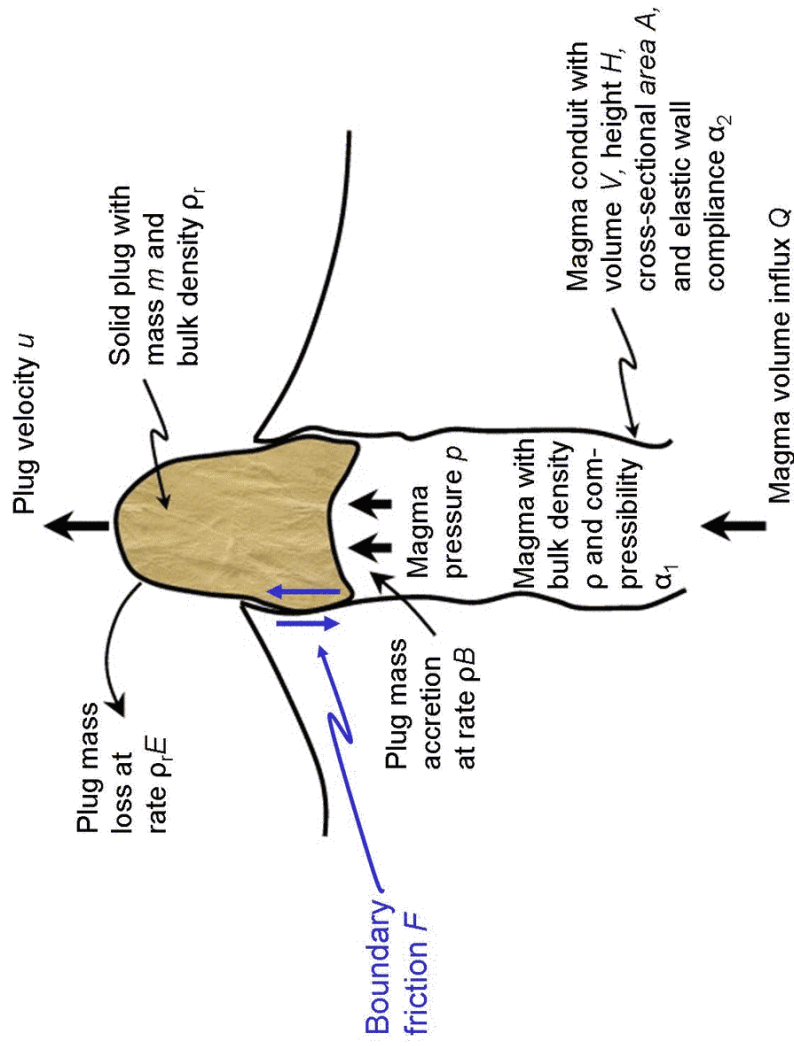




FACTS: Quasi-steady, solid-state extrusion at long-term rates ranging from 1 - 3 m³/s is accompanied by repetitive, quasi-periodic small earthquakes. The extruded plug is surfaced with striated, granulated fault gouge.

HYPOTHESIS: Extrusion is driven by magma pressure pushing against the base of an accreting solid plug that slips frictionally and incrementally against the conduit wall rock. Repetitive stick-slip cycles generate quasi-periodic small earthquakes.

SPASM Model (Seismogenic Plug of Ascending, Solidifying Magma)



Model Components

Physical Conservation Laws

- Linear momentum of solid plug with evolving mass
- Mass of solid plug (evolves due to basal accretion and surface erosion)
- Linear momentum of compressible conduit fluid (laminar flow in tube)
- Mass of conduit fluid (evolves due to change in density and conduit volume)

Constitutive "Laws"

- Elastic compliance of conduit walls
- Compressibility of conduit fluid
- Friction rule for solid plug interacting with conduit walls

Simplifying Assumptions

- One-dimensional dynamics
- System is driven by constant influx of magma at depth

Reduced system of governing equations

Evolution of Plug velocity

$$\frac{du}{dt} = -g + \frac{1}{m_0 + \kappa t} \left[pA - \kappa u - \frac{F(u/u_{ref})}{\kappa} \right]$$

possible strong nonlinearity in friction rule

Evolution of Magma pressure on plug base

$$\frac{dp}{dt} = \frac{1}{\alpha_1 + \alpha_2} \frac{Q - Au}{V}$$

weak nonlinearity due to quotient u/V

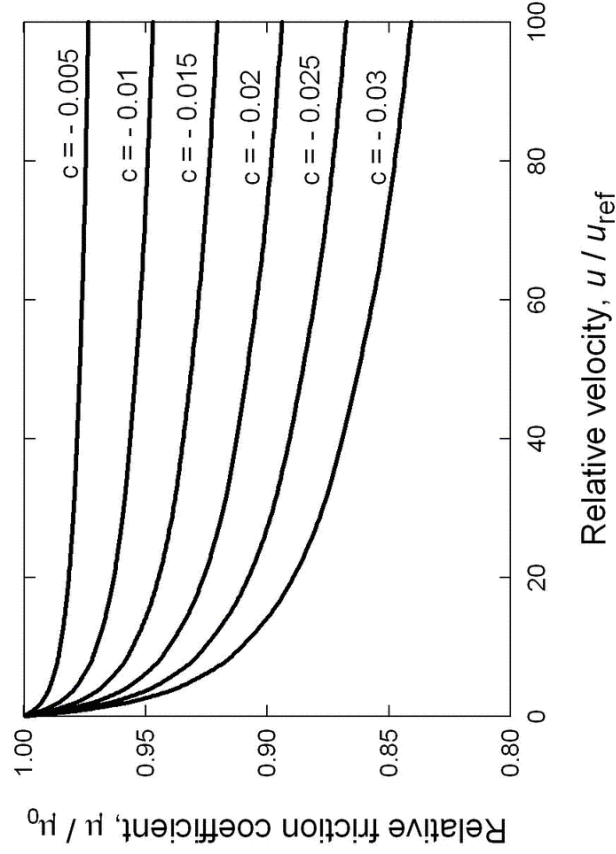
Evolution of Volume of magma-filled conduit

$$\frac{dV}{dt} = \frac{\alpha_1}{\alpha_1 + \alpha_2} Au + \frac{\alpha_2}{\alpha_1 + \alpha_2} Q - B$$

Nonlinear friction rule for sliding of solid plug

$$F(u/u_{ref}) = \lambda m_0 g \mu = \lambda m_0 g \mu_0 [1 + c \sinh^{-1}(u/u_{ref})]$$

For $u/u_{ref} < 1$, behavior approximates linear; for $u/u_{ref} > 1$, behavior approximates logarithmic



Part 1. Analytical Results

Exact steady-state solution
(for case with $Q = Au$, $\kappa = 0$)

$$u_0 = \frac{Q}{A} - p_0 = \frac{m_0 g C F_0}{A} - V_0$$

In this case plug velocity, plug mass, plug friction, conduit volume, and conduit pressure remain constant, and forces remain balanced.

Linearize about this steady state to examine behavior of small perturbations in plug velocity $u'(t)$. Assume rate-dependence of friction is linear, i.e. $dF/du = \lambda m_0 g \mu_0 C / u_{ref} = C$

Find velocity perturbations obey damped, forced oscillator equation

$$\left(\frac{1}{m_0} \frac{d^2 u}{dt^2} + \frac{1}{m_0} \left(\frac{2\kappa C F_0}{u_0} \right) \frac{du}{dt} + \frac{1}{m_0 (\alpha_1 - \alpha_2)} \frac{A^2}{V_0} u \right) = \frac{g \kappa}{m_0}$$

Acceleration rate
Damping rate
Restoring rate
Forcing rate

For simplest case with zero κ (no plug mass change) and zero damping ($C = 0$, no rate-dept. friction) this equation has a simple periodic solution

$$u = u_0 \cos(t/t_0) \quad \text{where} \quad t_0 = \frac{[m_0 V_0 (\alpha_1 - \alpha_2)]^{1/2}}{A}$$

Period of free oscillations is $T_0 = \frac{2\pi [m_0 V_0 (\alpha_1 \ominus \alpha_2)]^{1/2}}{A}$

For a cylindrical plug and conduit this period can be expressed as

$$T_0 = 2\pi [(\alpha_1 \ominus \alpha_2) \rho_r H_{plug} H_{conduit}]^{1/2}$$

Insert some appropriate numbers for Mount St. Helens:

bulk compressibility $\alpha_1 + \alpha_2 = 10^{-8} \text{ Pa}^{-1}$

plug-rock density $\rho_r = 2400 \text{ kg/m}^3$

plug height $H_{plug} = 1 \text{ km}$

conduit height $H_{conduit} = 8 \text{ km}$

These numbers yield free oscillation period $T = 87 \text{ s}$, similar to common earthquake recurrence interval during 2004-2005 eruption of Mount St. Helens

Stability of quasi-steady extrusion
subject to small oscillations

Case I : plug mass is constant ($\kappa = 0$)

Stability parameter

$$D = (T/4\pi) g \lambda \mu_0 \frac{c}{u_{ref}}$$

Oscillations decay if $D > 0$; grow if $D < 0$

Instability can be manifested as either growing oscillations or runaway extrusion, depending on the value of D .

Stability of quasi-steady extrusion
subject to small oscillations

Case II : plug mass evolves ($\kappa \neq 0$)

Stability parameter

$$C/\kappa = \frac{1}{\kappa} m_0 g \lambda \mu_0 \frac{c}{u_{ref}}$$

- $\kappa > 0, C > 0$ unconditionally stable
- $\kappa < 0, C < 0$ unconditionally unstable
- ★ $\kappa > 0, C < 0$ stable oscillations if $-C/\kappa < 3/2$
- $\kappa < 0, C > 0$ stable oscillations if $-C/\kappa > 3/2$

This is the only case in which persistent stick-slip cycles occur !

Part 2. Computational Results for Nonlinear System

$$\frac{du}{dt} = -g + \frac{1}{m_0 + \kappa t} [pA - \kappa u - F(u/u_{ref})]$$

$$\frac{dp}{dt} = \frac{1}{\alpha_1 + \alpha_2} \frac{Q - Au}{V}$$

$$\frac{dV}{dt} = \frac{\alpha_1}{\alpha_1 + \alpha_2} Au + \frac{\alpha_2}{\alpha_1 + \alpha_2} Q - B$$

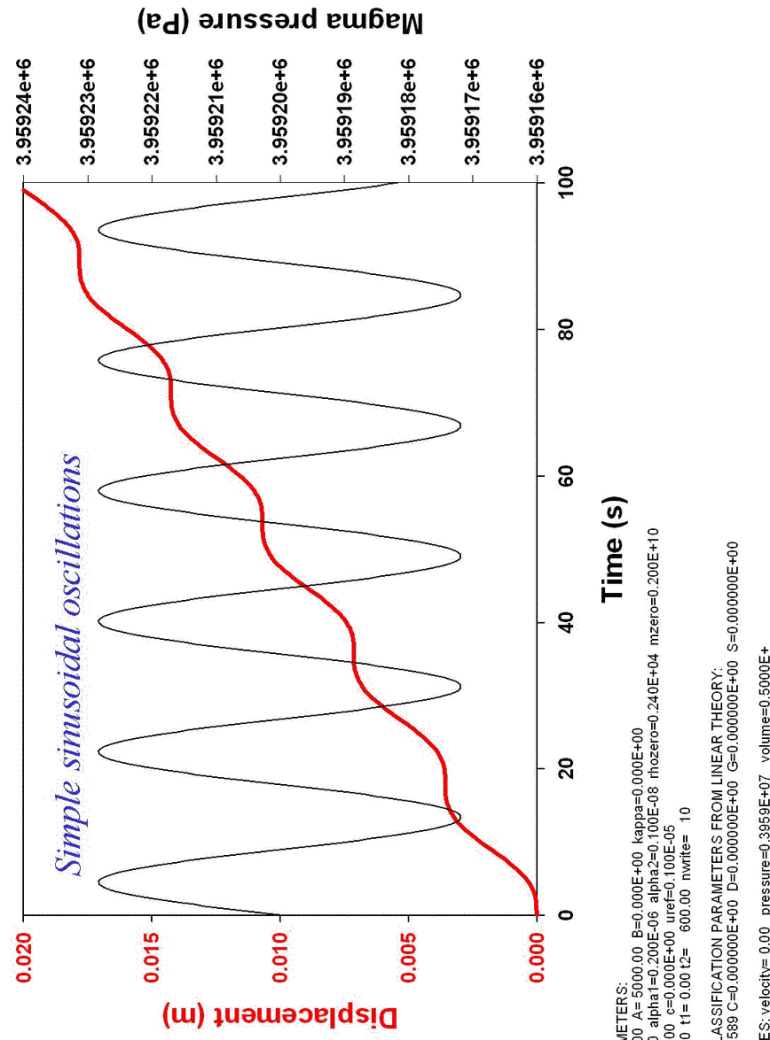
$$F(u/u_{ref}) = \lambda m_0 g \mu = \lambda m_0 g \mu_0 [1 + c \sinh^{-1}(u/u_{ref})]$$

First, fix values of all parameters except c at reasonable values:

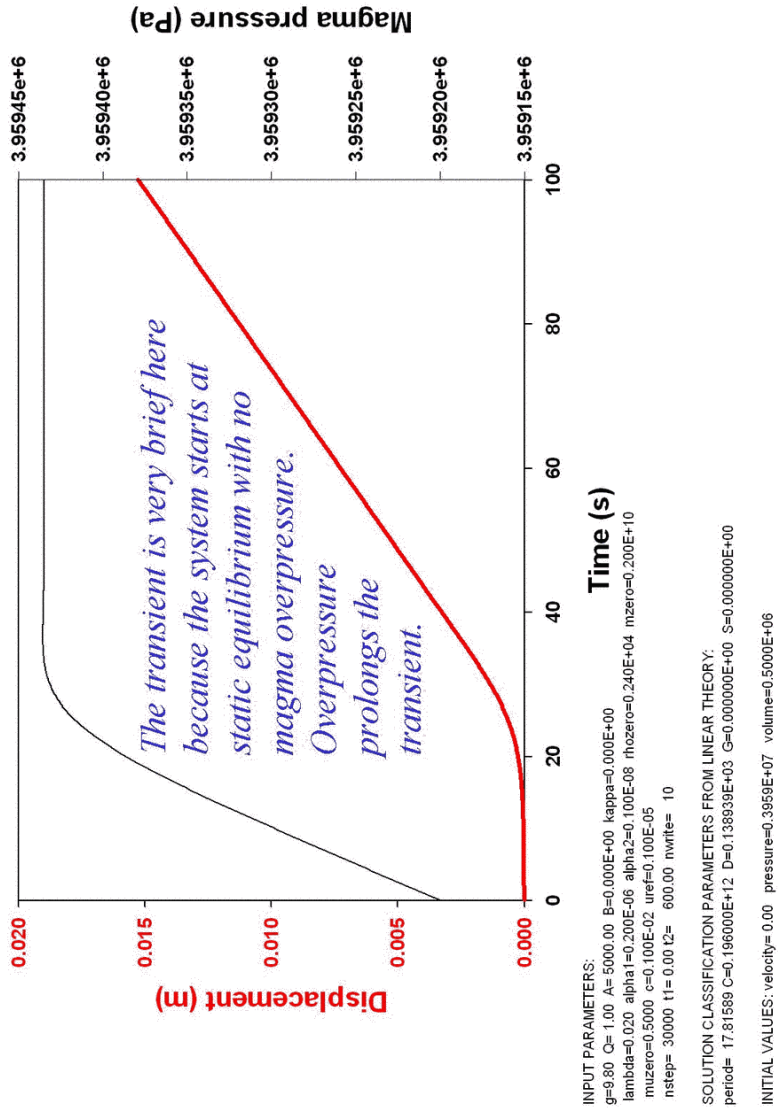
- Magma influx rate $Q = 1 \text{ m}^3/\text{s}$
- Plug base area $A = 5000 \text{ m}^2$ (80 m diameter circle)
- Plug mass $m = 2 \times 10^9 \text{ kg}$
- Plug mass growth rate $\kappa = 0$
- Magma compressibility $\alpha_1 = 2 \times 10^{-7} \text{ Pa}$
- Conduit compliance $\alpha_2 = 1 \times 10^{-9} \text{ Pa}$
- Static friction coefficient $\mu_0 = 0.5$
- Effective stress coefficient $\lambda = 0.02$ (friction can support 2% of plug weight)
- Reference sliding rate $u_{ref} = 1 \times 10^{-6} \text{ m/s}$ (*implies strong nonlinearity*)

Then, use Runge-Kutta method to generate numerical results for different values of c

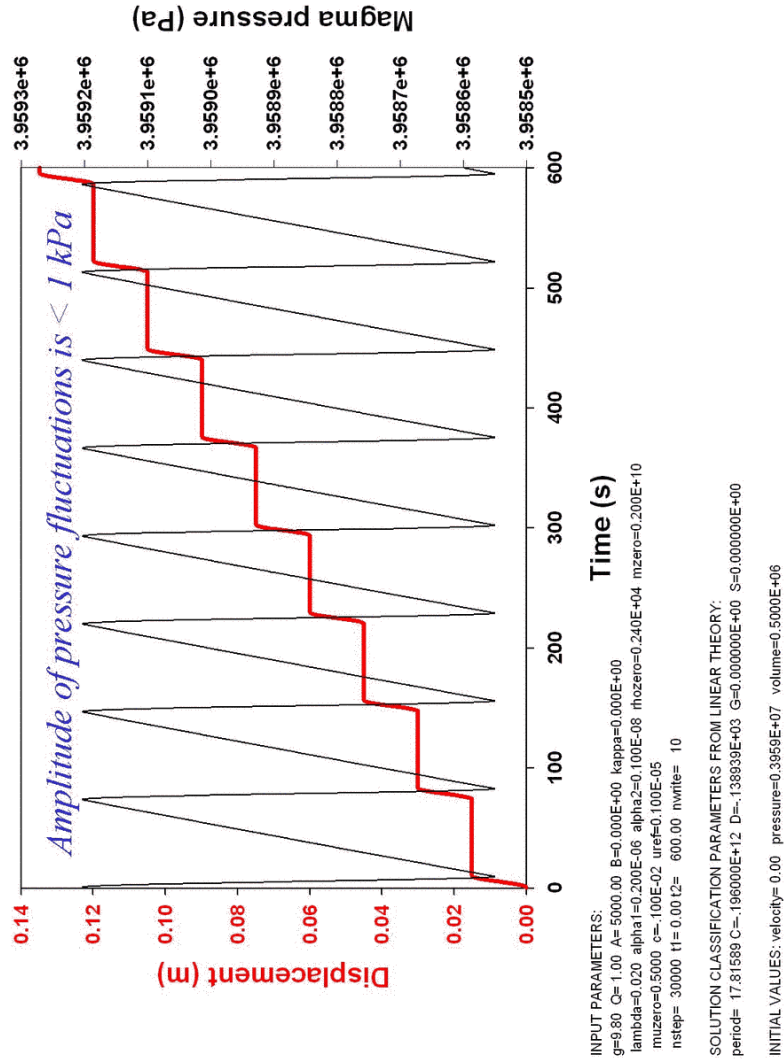
Baseline solution for $c = 0$ reproduces analytical results for linearized system, substantiating the inference that the nonlinearity involving uV is weak.



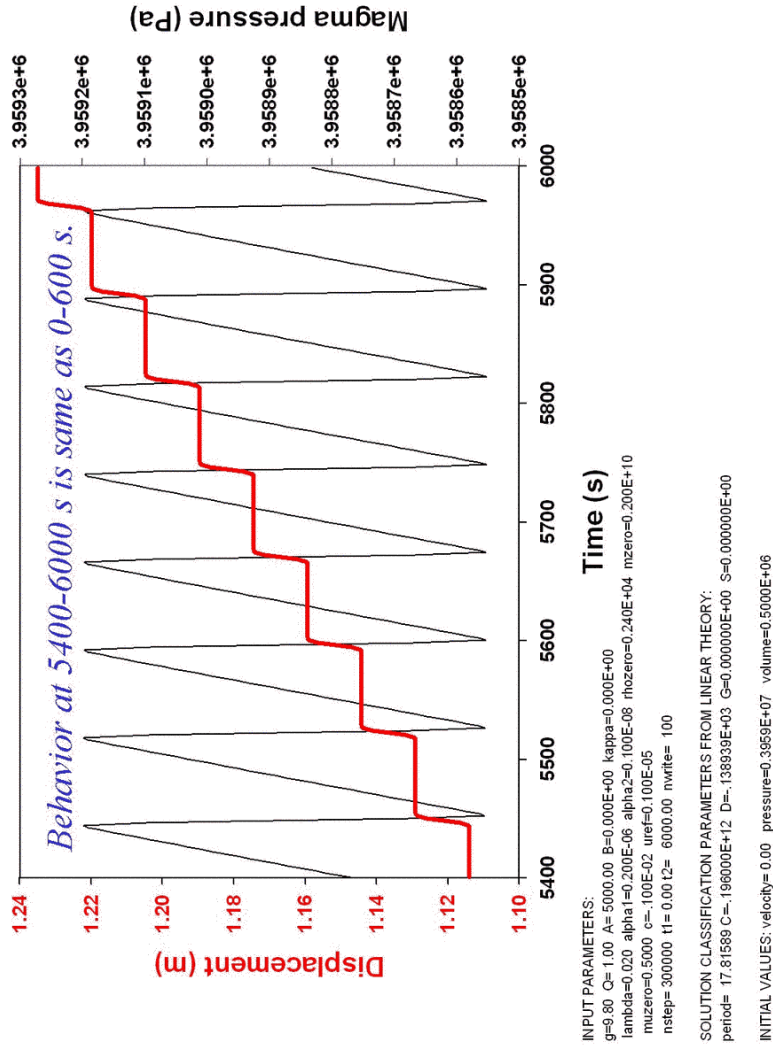
In rate-strengthening case with $c = 0.001$, first a transient occurs in which forces equilibrate. Then the system reaches stable, dynamic equilibrium. This occurs for both linear (i.e., “viscous”) and nonlinear rate strengthening.



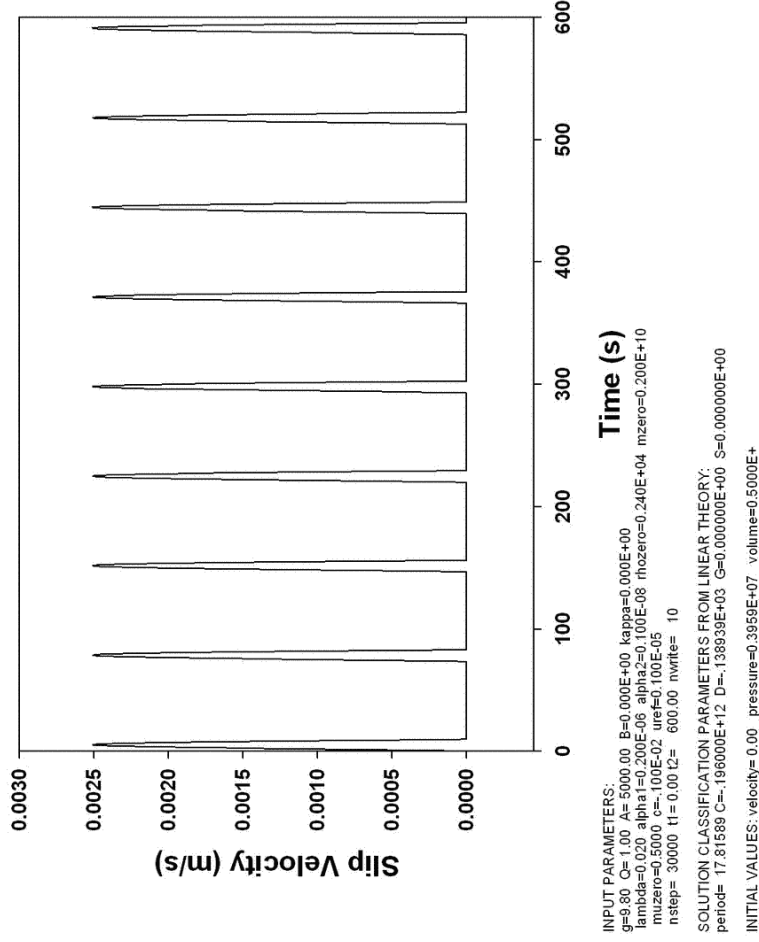
In rate-weakening case with $c = -0.001$, the system exhibits stick-slip oscillations with periods longer than those predicted by the linear theory



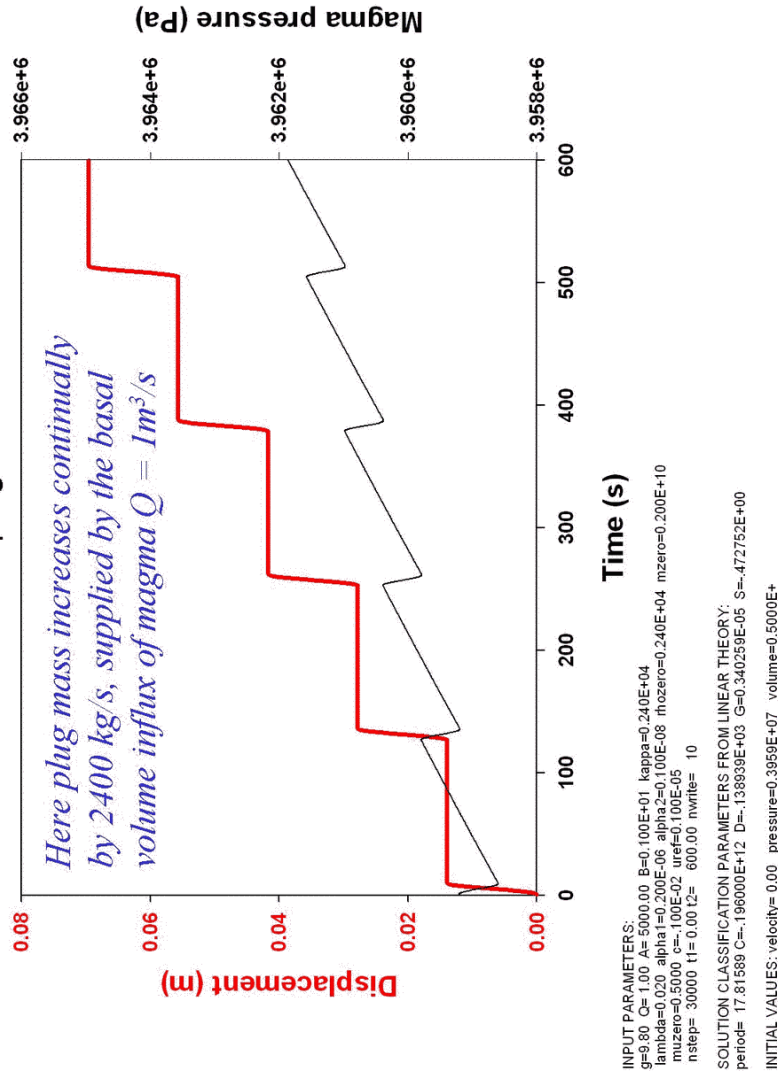
The amplitude of the stick-slip oscillations does not grow with time, contrary to predictions of the linear theory. The nonlinearity stabilizes the oscillations by limiting positive feedback between slip velocity and frictional resistance. This stabilization may help explain the repetitive earthquake pattern at MSH



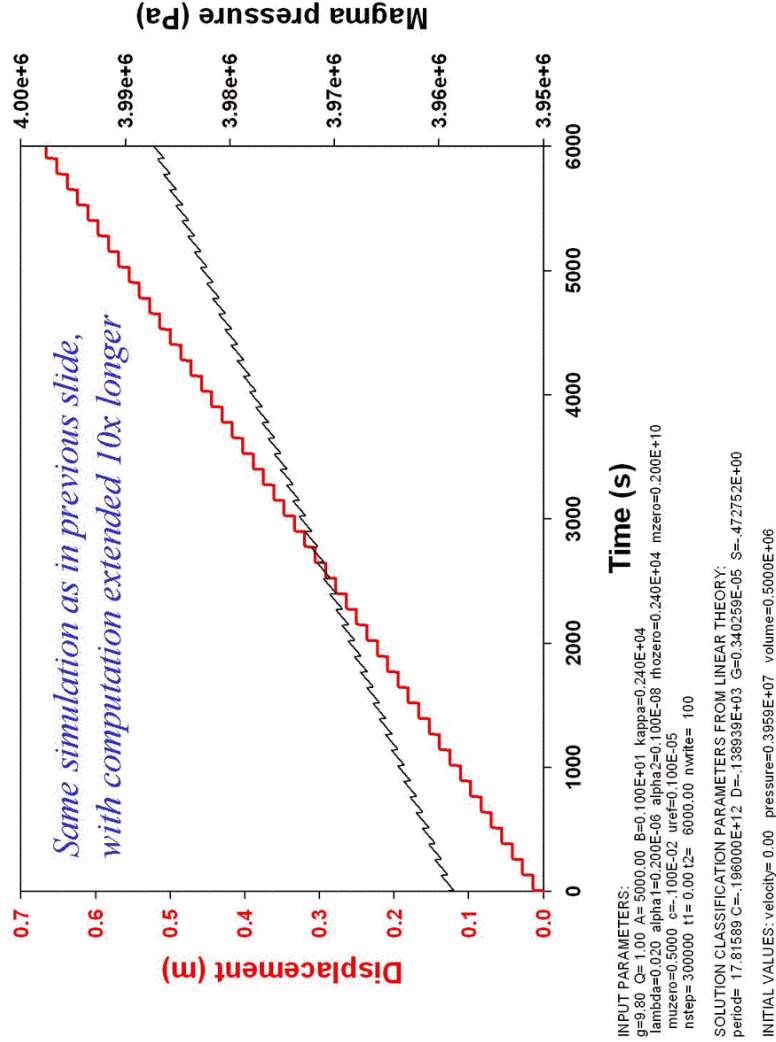
Velocity “spikes” accompany the slip events depicted in previous slides. Work done during each of these spikes is about 3×10^8 J, similar to the energy release in a M 2.2 earthquake. Of course, some of the work may be done aseismically, resulting in smaller earthquakes.



Now, include mass accretion at base of plug, with no compensating erosion of surface. The added mass lengthens the oscillation period, reduces the extrusion rate, and causes the magma pressure to gradually increase. The converse occurs if there is net loss of plug mass.



... and the beat goes on forever unless magma influx changes or one or more parameters evolve to change the system response. Note, however, that the magma pressure is gradually increasing...



Normalized damped, forced oscillator equation

$$(1 + Gt) \frac{d^2u}{dt^2} + 2GIR \frac{du}{dt} + u = \frac{GI}{A}$$

Dimensionless parameters also control numerical solutions

$$I \approx \frac{t_0 g}{Q/A} \quad O(1000) \text{ or larger applies to MSH}$$

$$G \approx \frac{t_0}{m_0/\kappa} \quad O(1) \text{ or smaller applies to MSH}$$

$$R \approx \lambda \mu_0 c \frac{Q/A}{u_{ref}} \quad O(1) \text{ or smaller probably applies to MSH}$$

Conclusions

1. Nonlinear rate-weakening boundary friction can explain stick-slip oscillations during extrusion of magma plug.
2. Use of realistic (i.e. best-guess) parameter values produces stick-slip oscillations with roughly the correct period and amplitude to correspond to repetitive “drumbeat” earthquakes at MSH.
3. Nonlinearity of the friction rule stabilizes stick-slip oscillations and allows them to occur “forever” if plug mass remains constant or if it evolves.
4. If plug mass changes due to net erosion or accretion, the oscillation period and amplitude and the surface extrusion rate change in response.
5. Long-term, oscillatory behavior of the system is remarkably stable unless magma influx changes, friction evolves, or magma pressure increases sufficiently to cause a hydrofracturing “blowout.”