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The Simons Collaboration on the Many Electron Problem

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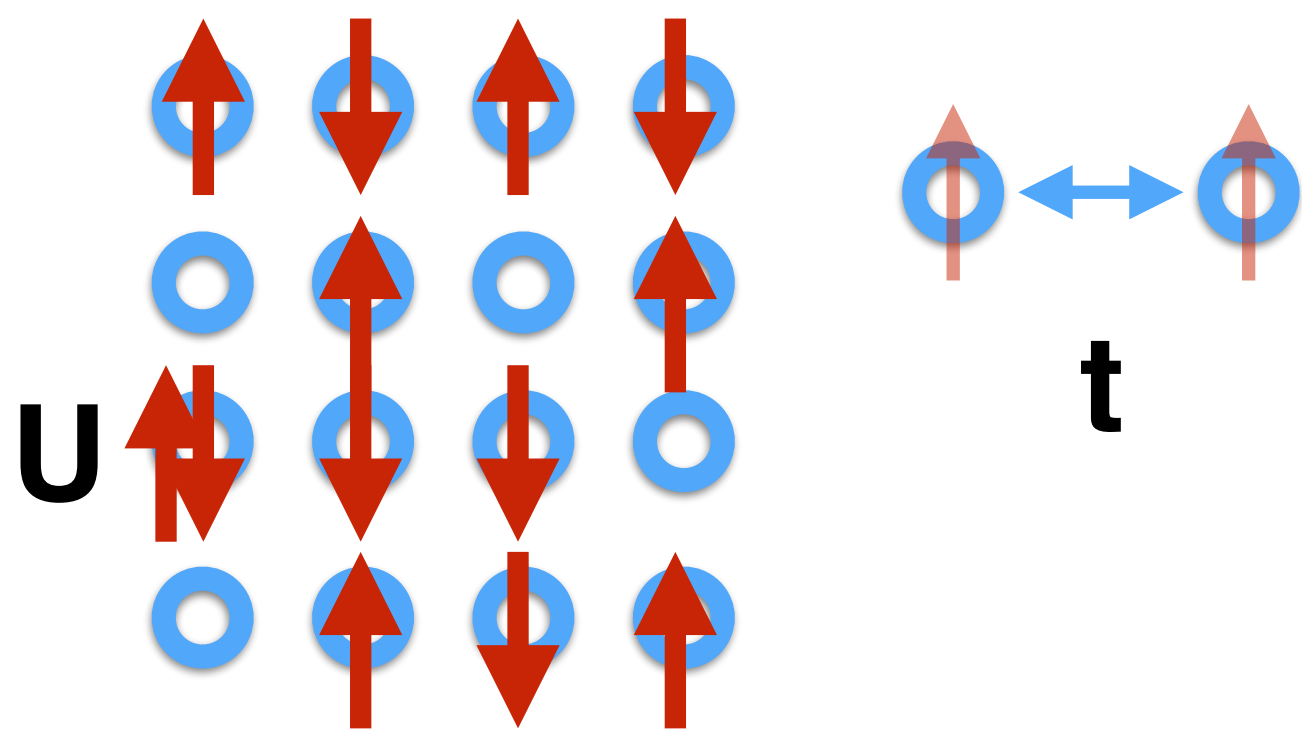
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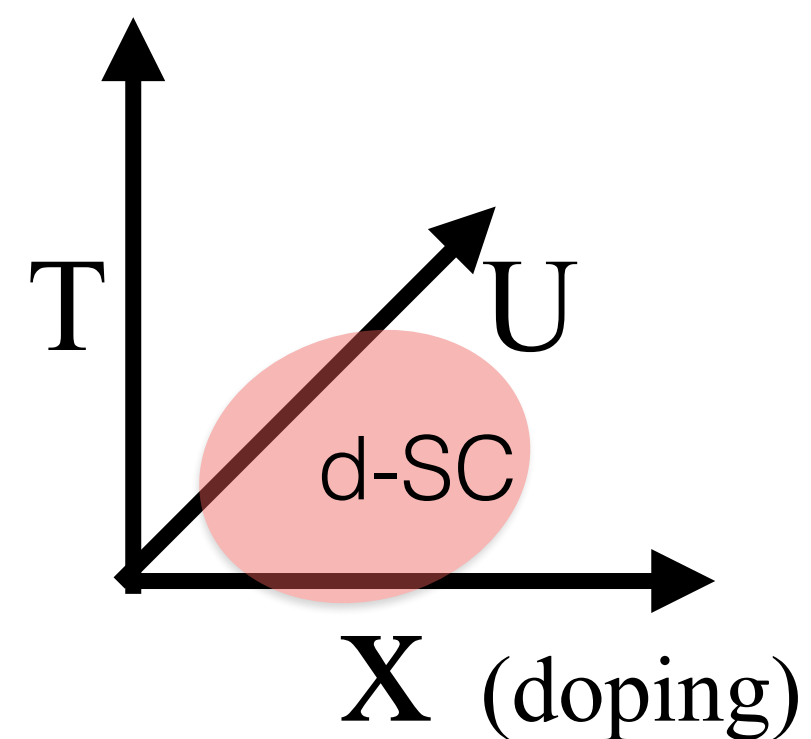
- Brief Intro to the Hubbard model
- Four powerful simulation methods
- Results for $U=8$, $1/8$ doping:
 - Consensus on the phase of the ground state
 - Stripes with nearly degenerate wavelengths

The 2D Hubbard model



$$H = - \sum_{\langle ij \rangle, \sigma} t_{ij} (c_{i\sigma}^\dagger c_{j\sigma} + c_{j\sigma}^\dagger c_{i\sigma}) + U \sum_i n_{i\uparrow} n_{i\downarrow}.$$

Perhaps the most important model in condensed matter physics—widely regarded as the starting point for understanding the high- T_c superconductors



t sets the energy scale, so only one parameter in H
 You can also vary the temperature and doping.
 The phase diagram specifies: $\text{phase}(T, x, U)$

For the high- T_c cuprates:
 $U/t \sim 8, \quad 0 < x < 0.3,$
 $T_{sc} < \sim J/12 \sim t/40$
 $T_{ps} < \sim J \sim t/3$

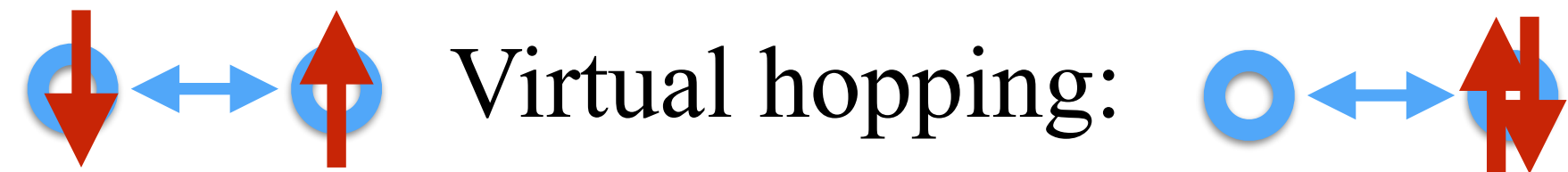
The t and U terms compete.
 Bandwidth $W = 8t$, maximal competition $W \sim U$, or $U=8t$ (cuprates!)

The model is more easily understood for small U or large U .
 Small U : quasiparticles, Fermi surfaces, diagrams, self energies....
 Large U : exchange, mapping to Heisenberg and t - J models (which are still hard to solve!)

The 2D Hubbard model—large U/t



U



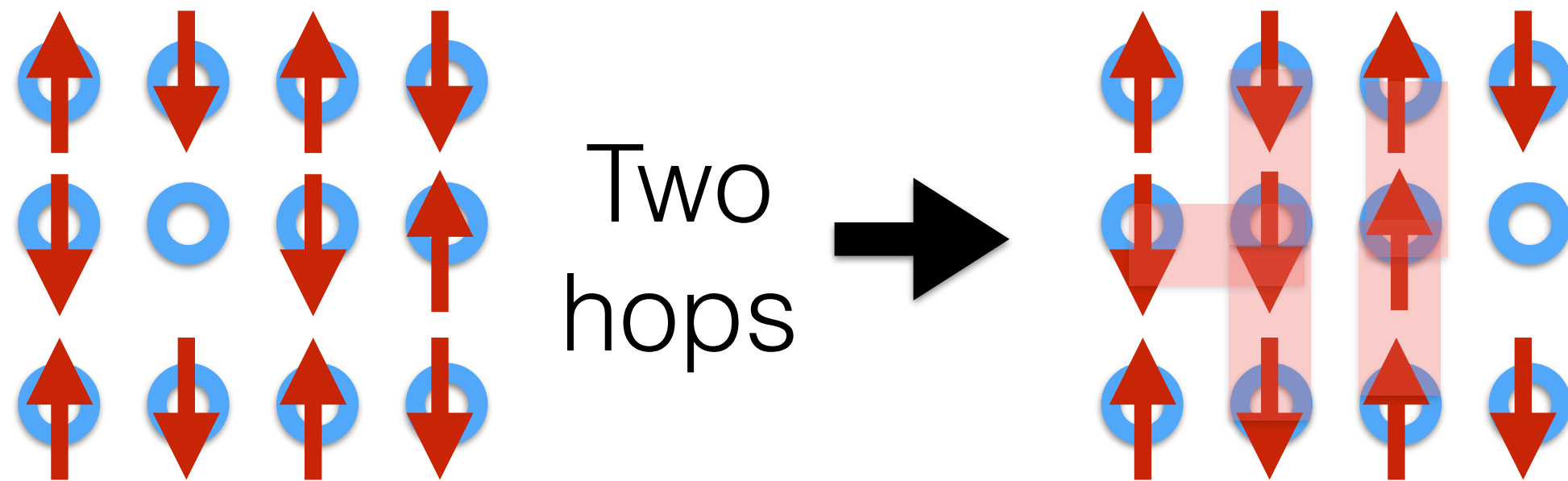
Perturbative effect: new exchange interaction

$$J \vec{S}_i \cdot \vec{S}_j \quad J \approx 4t^2 / U$$

favoring antiferromagnetism

t-J model: keep original hopping for holes, replace U term by $J \vec{S}_i \cdot \vec{S}_j$

Frustrated hole hopping in an antiferromagnet:

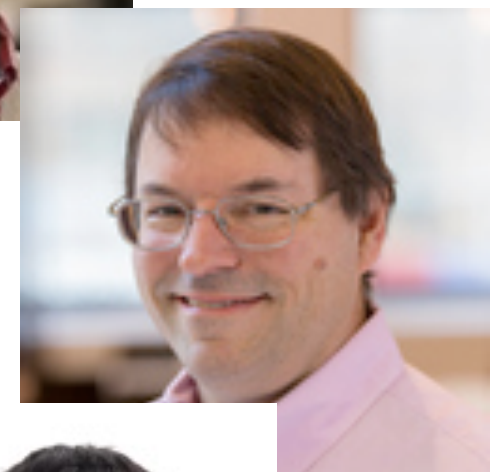
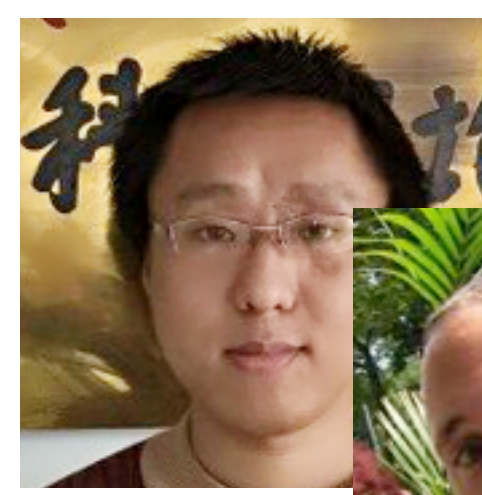


Single hole hopping is highly “frustrated”

Pairing: one hole can follow the other, erasing ferromagnetic bonds

Stripes: a vertical line of holes can hop together without frustration

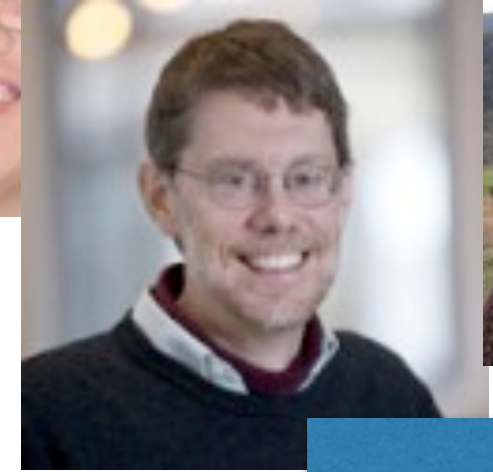
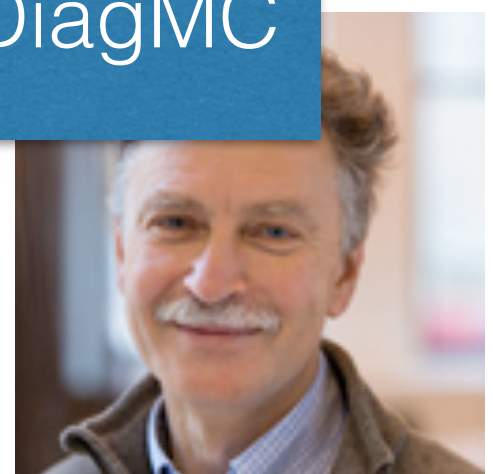
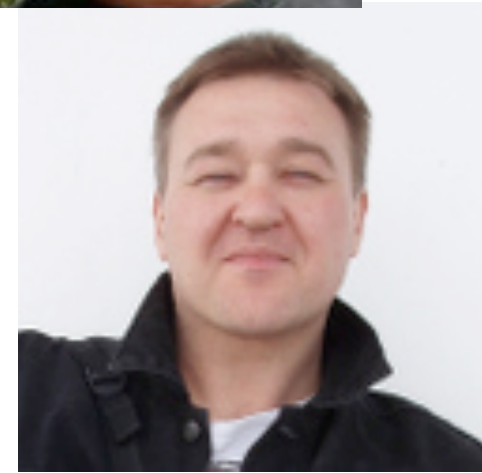
Generic feature of frustrated systems: new states can appear that are not directly favored by H , if they relieve frustration



DMRG

DiagMC

DCA



DF

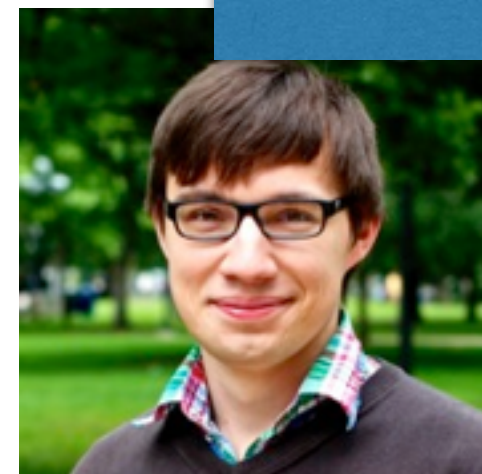


First Hubbard Benchmark

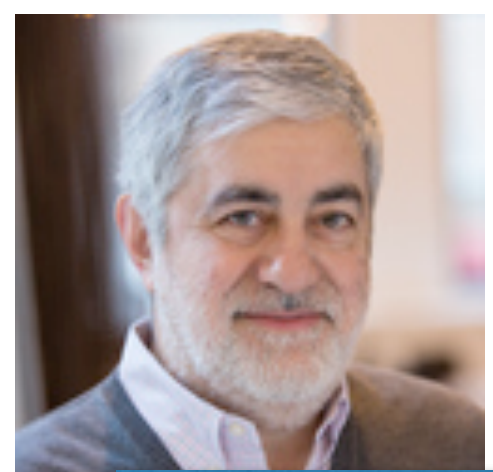
paper

$$H = - \sum_{\langle ij \rangle, \sigma} t_{ij} (c_{i\sigma}^\dagger c_{j\sigma} + c_{j\sigma}^\dagger c_{i\sigma}) + U \sum_i n_{i\uparrow} n_{i\downarrow}.$$

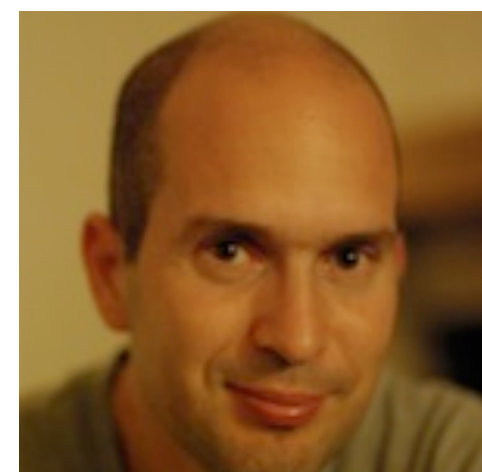
MRPHF



VMC & Fixed Node DMC



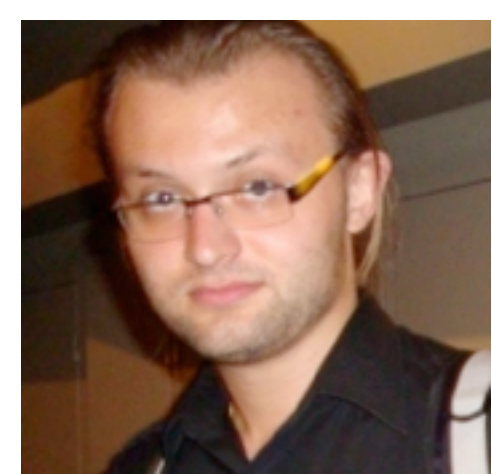
UCCSD
UCCSDT(Q)



DMET



AFQMC



Solutions of the Two Dimensional Hubbard Model: Benchmarks and Results from a Wide Range of Numerical Algorithms

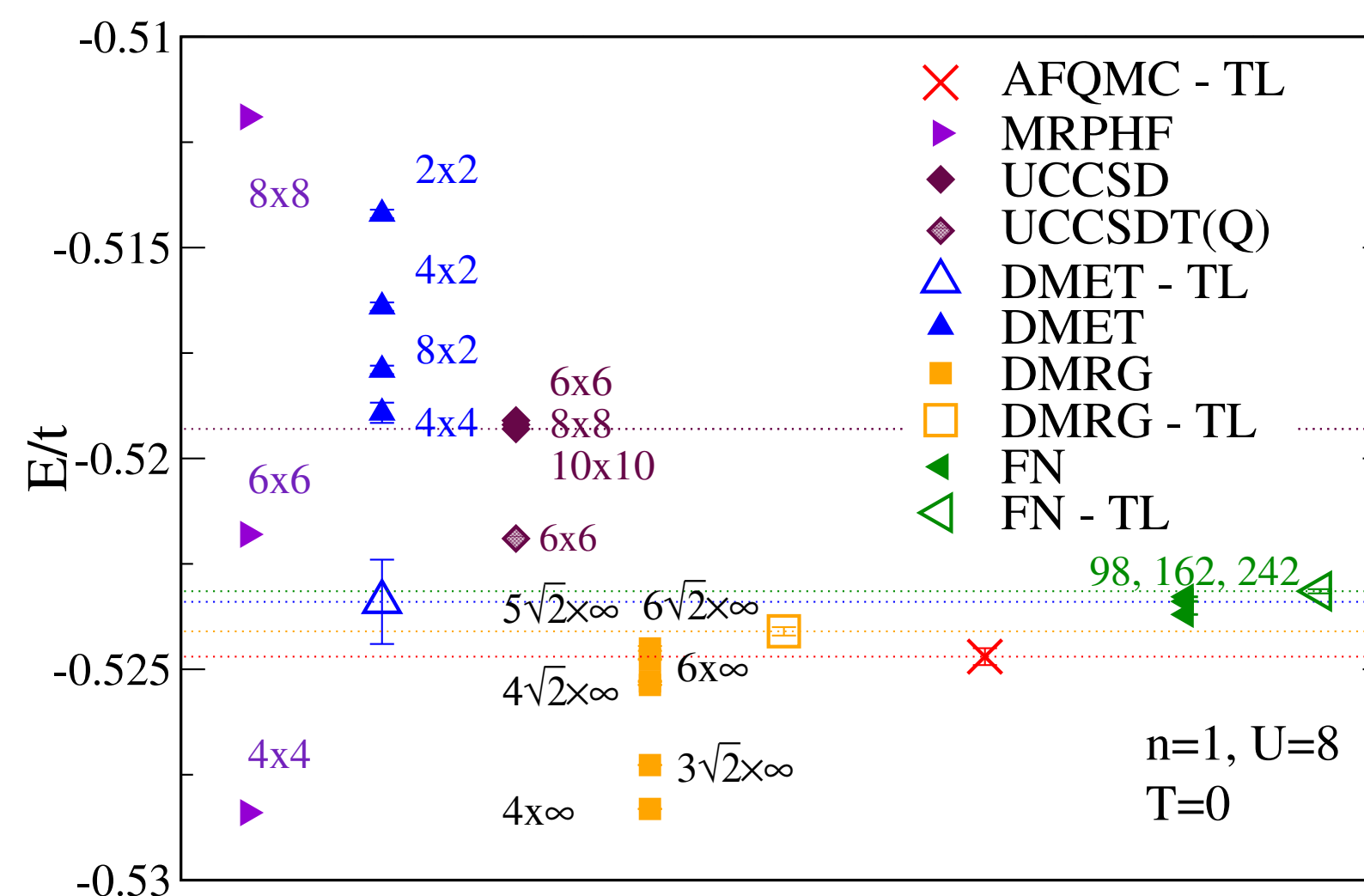
J. P. F. LeBlanc,¹ Andrey E. Antipov,¹ Federico Becca,² Ireneusz W. Bulik,³ Garnet Kin-Lic Chan,⁴ Chia-Min Chung,⁵ Youjin Deng,⁶ Michel Ferrero,⁷ Thomas M. Henderson,^{3,8} Carlos A.

We studied a wide range of doping, U/t , temperature—in thermodynamic limit. For most of parameter space, there was good agreement in energies and other properties.

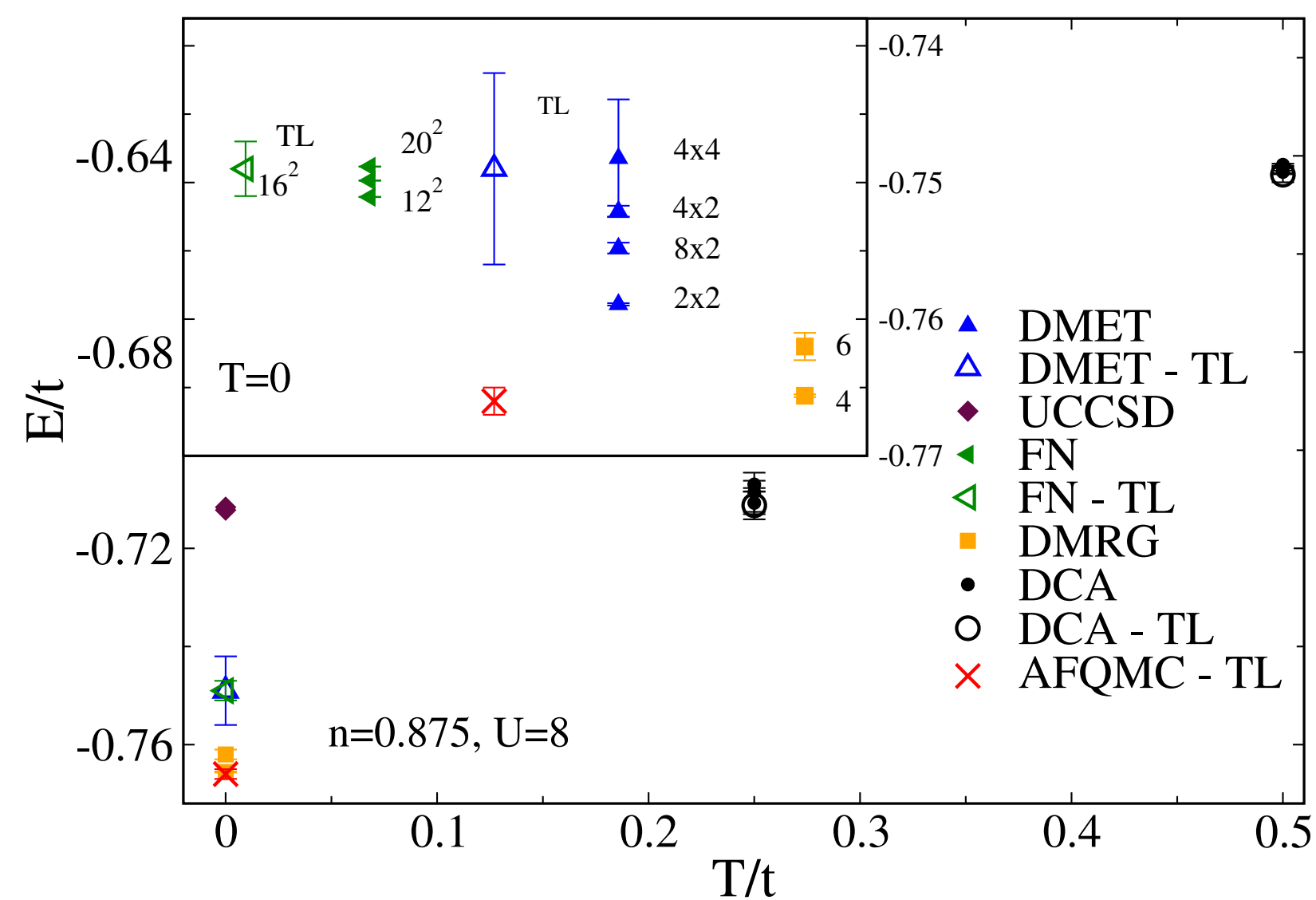
Energy turned out to be an good gauge of the quality of a simulation, correlating with other measures. Errors in E were usually positive—(semi)variational(ish).

For parameters relevant to the cuprates ($U/t \sim 8$), the finite T methods could not go to low enough T to compete in determining the ground state.

At $T=0$, the best different methods gave similar energies, but different states.



Ground state energies, **half-filling**, $U=8$



Energies vs T , **1/8 doping**, $U=8$

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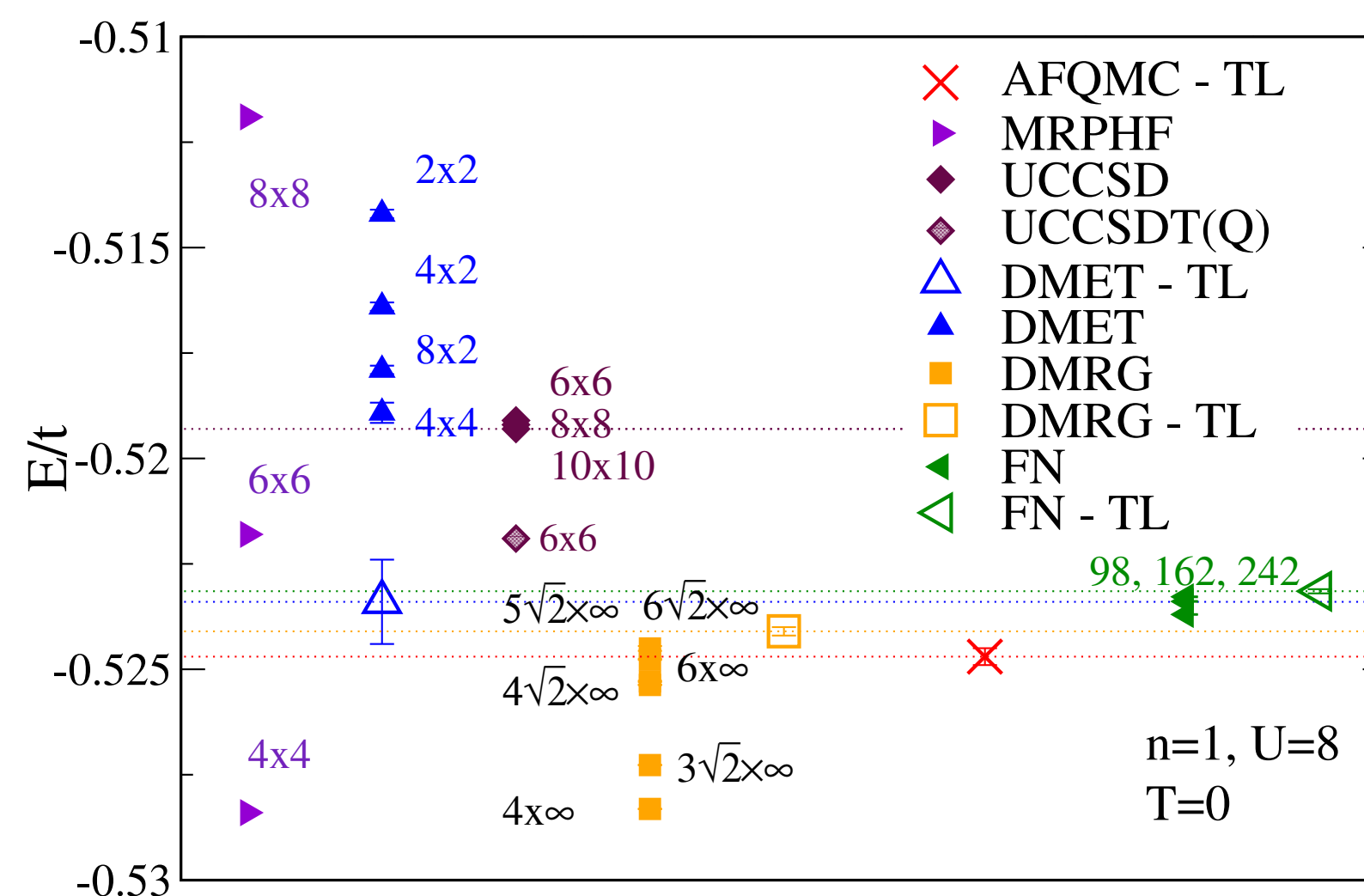
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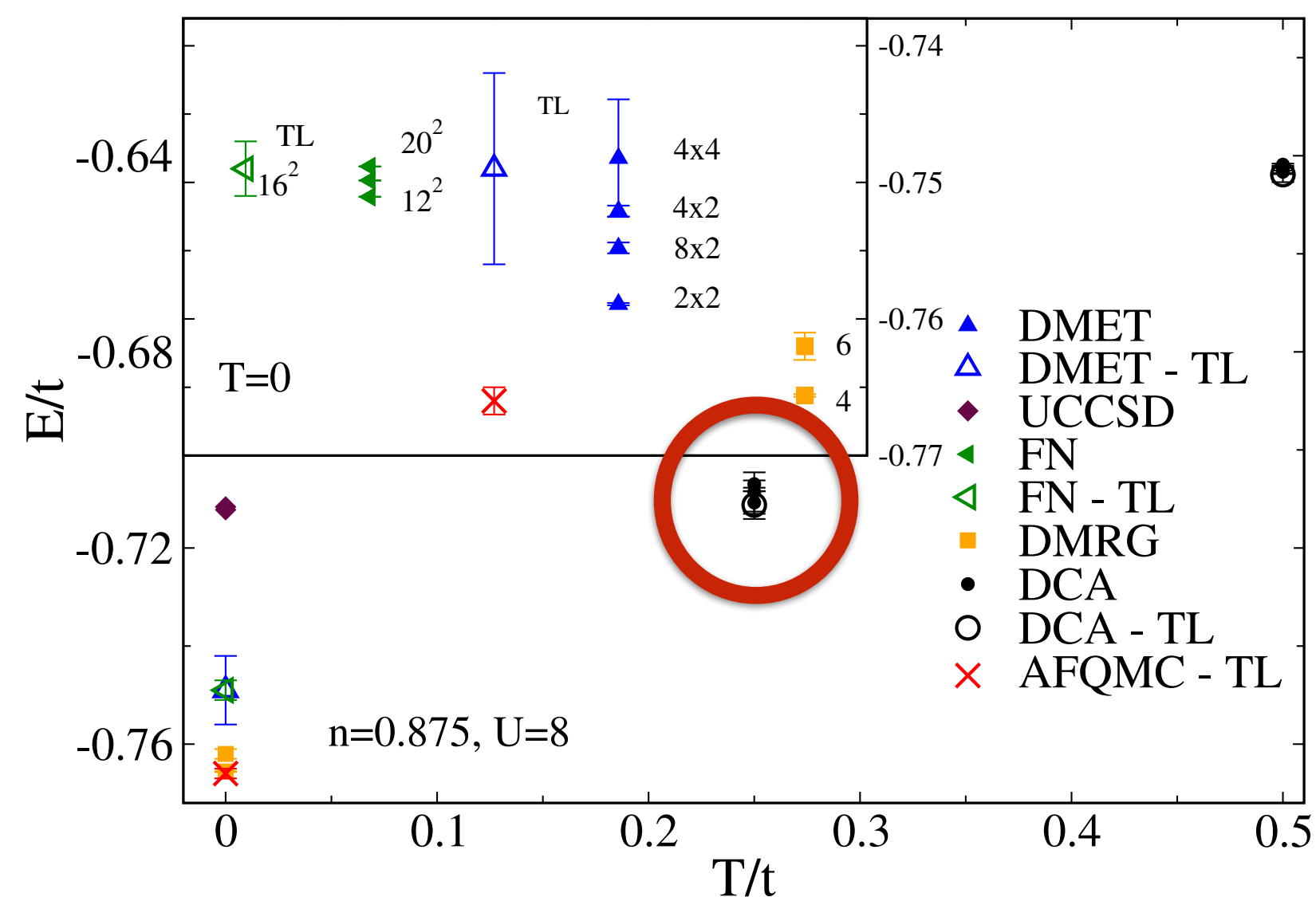
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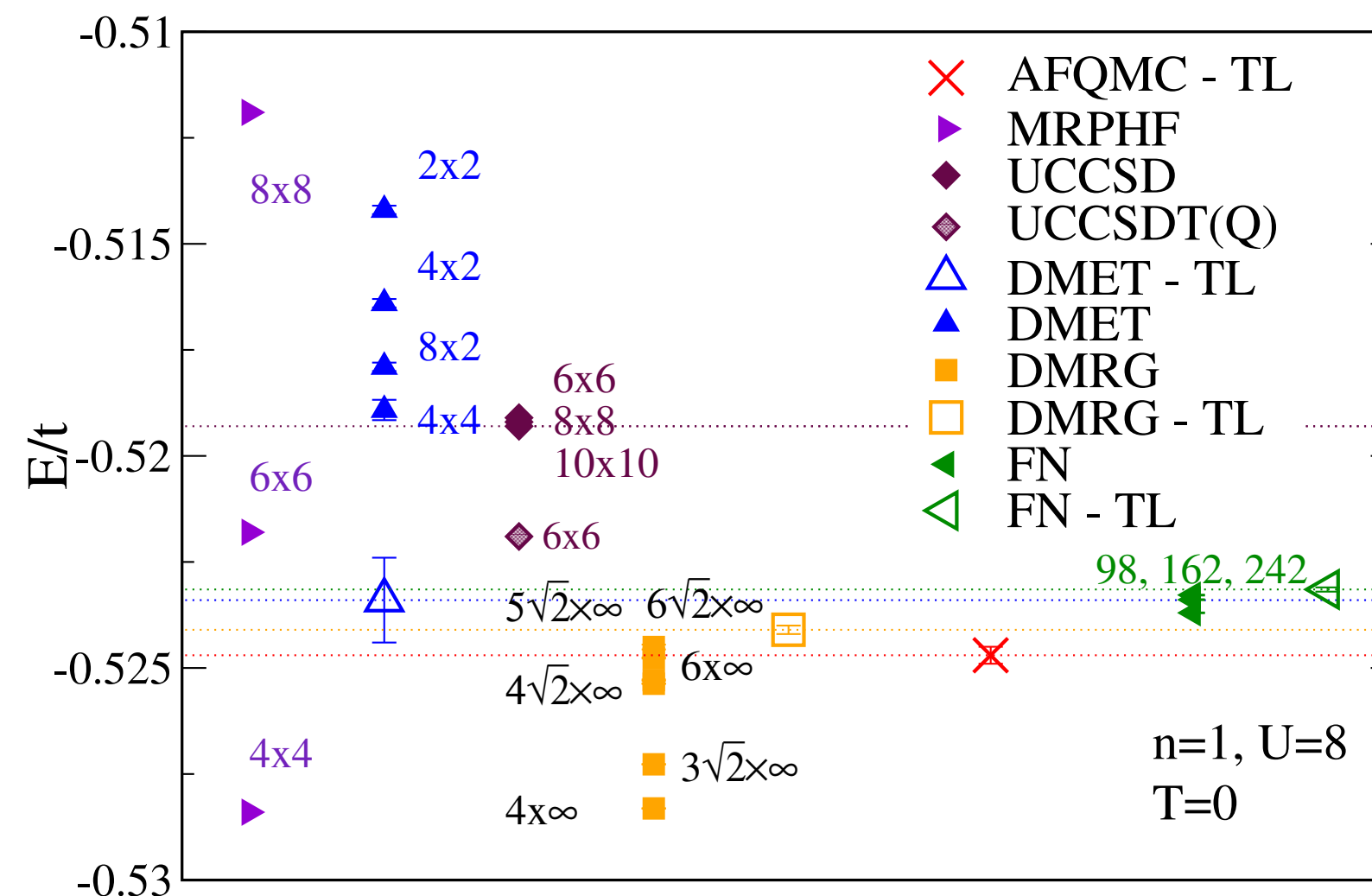
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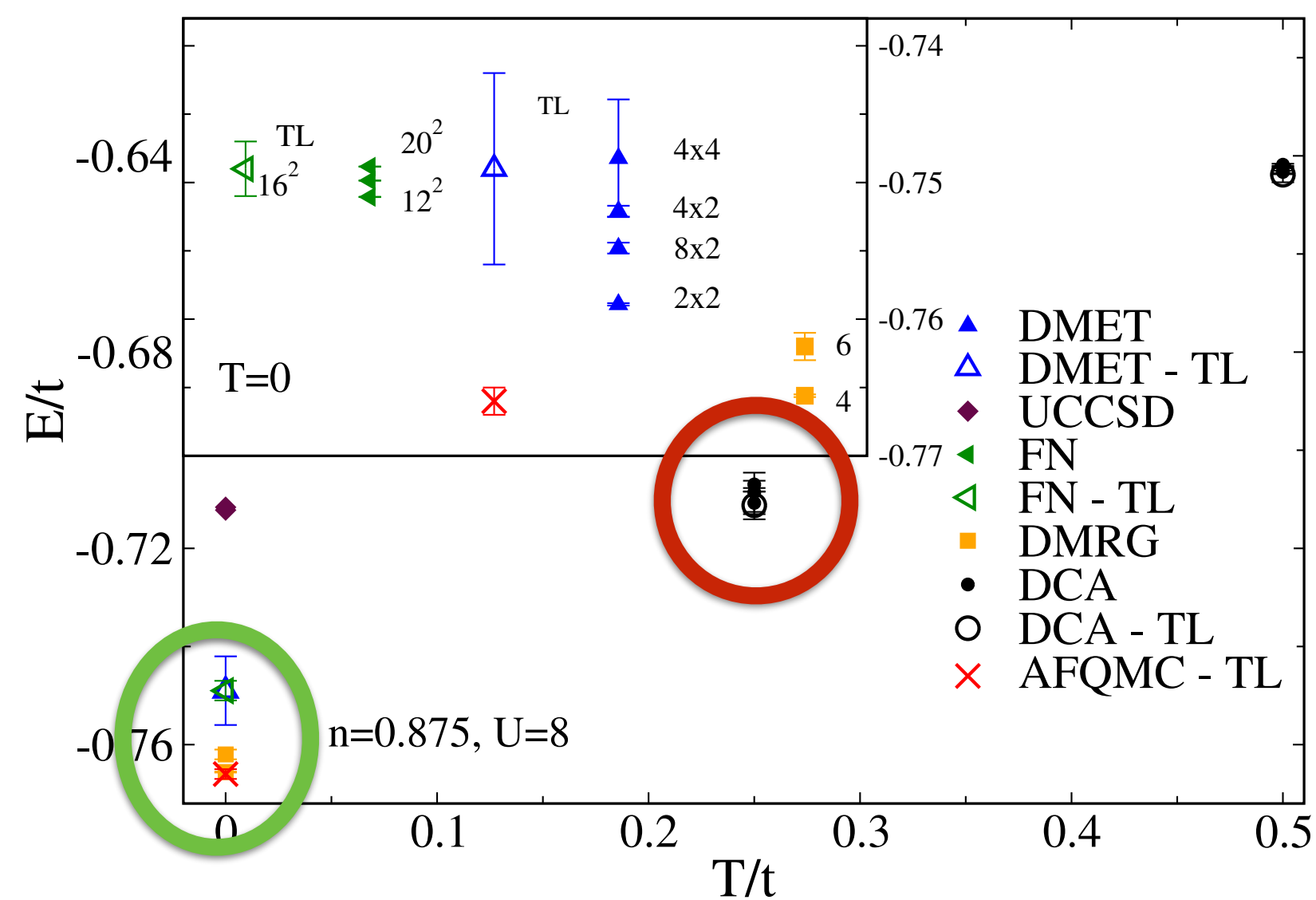
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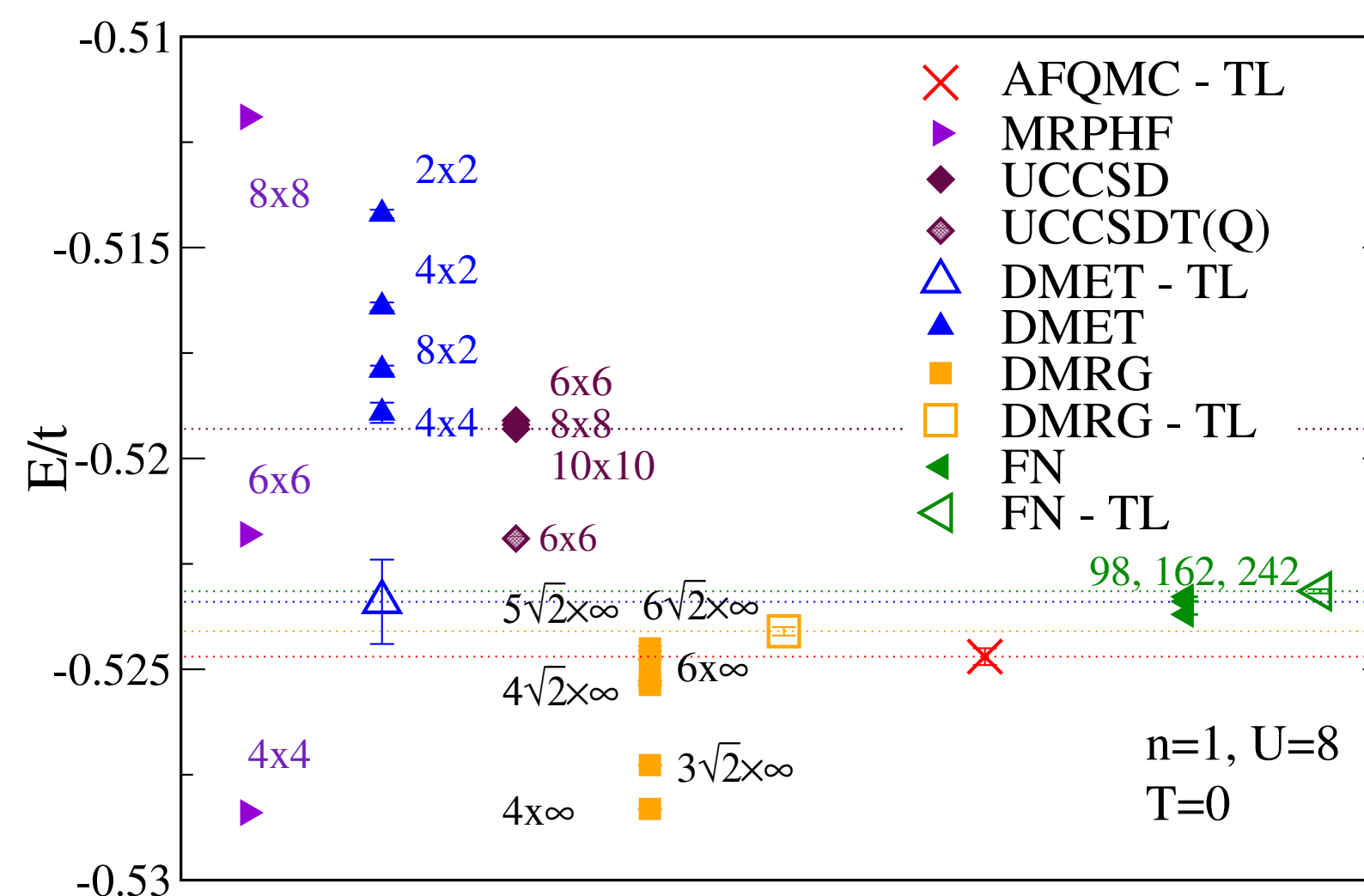
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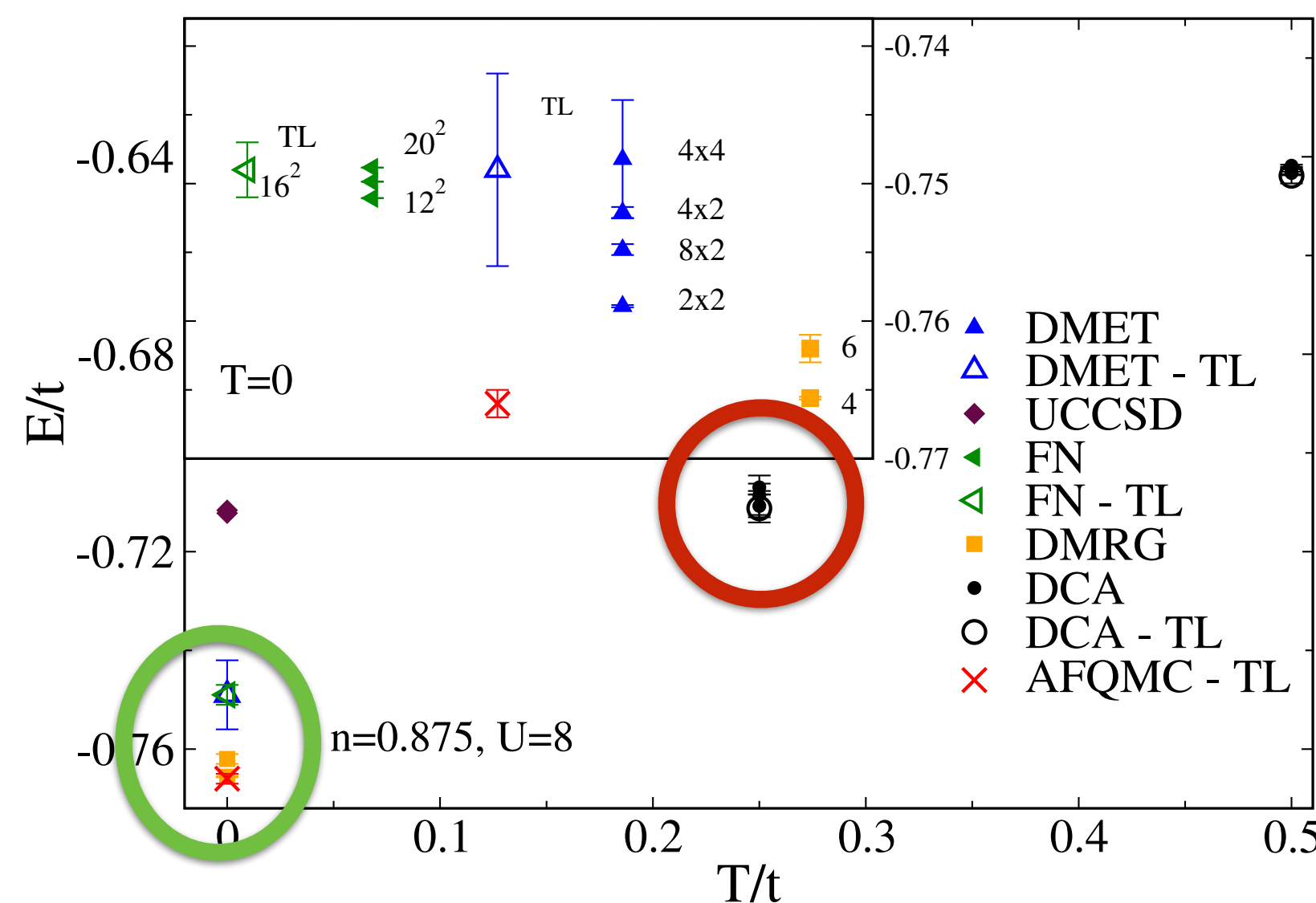
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Ground state energies, **half-filling**, $U=8$



Energies vs T , **1/8 doping**, $U=8$

A smaller group of us then decided to focus on this toughest region to see if we could resolve differences

Stripe order in the underdoped region of the two-dimensional Hubbard model

T=0, U=8, 1/8 doping

maximum uncertainty in the phase, maximum inhomogeneity

To appear in *Science*

Bo-Xiao Zheng^{1,2}, Chia-Min Chung³, Philippe Corboz⁴, Georg Ehlers⁵,
Ming-Pu Qin⁶, Reinhard M. Noack⁵, Hao Shi⁶, Steven R. White³,
Shiwei Zhang⁶, Garnet Kin-Lic Chan^{1*}

What is the ground state phase?

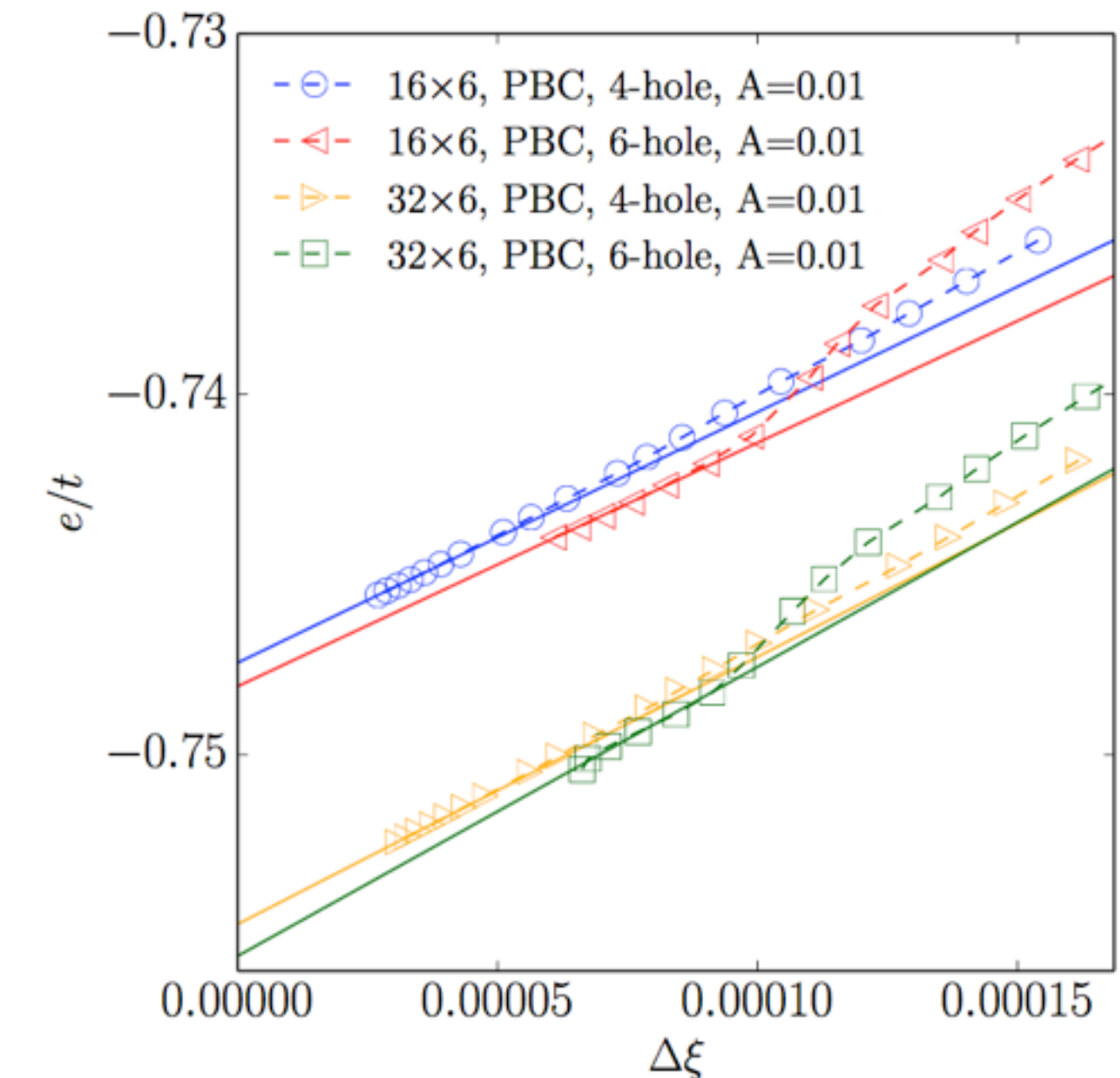
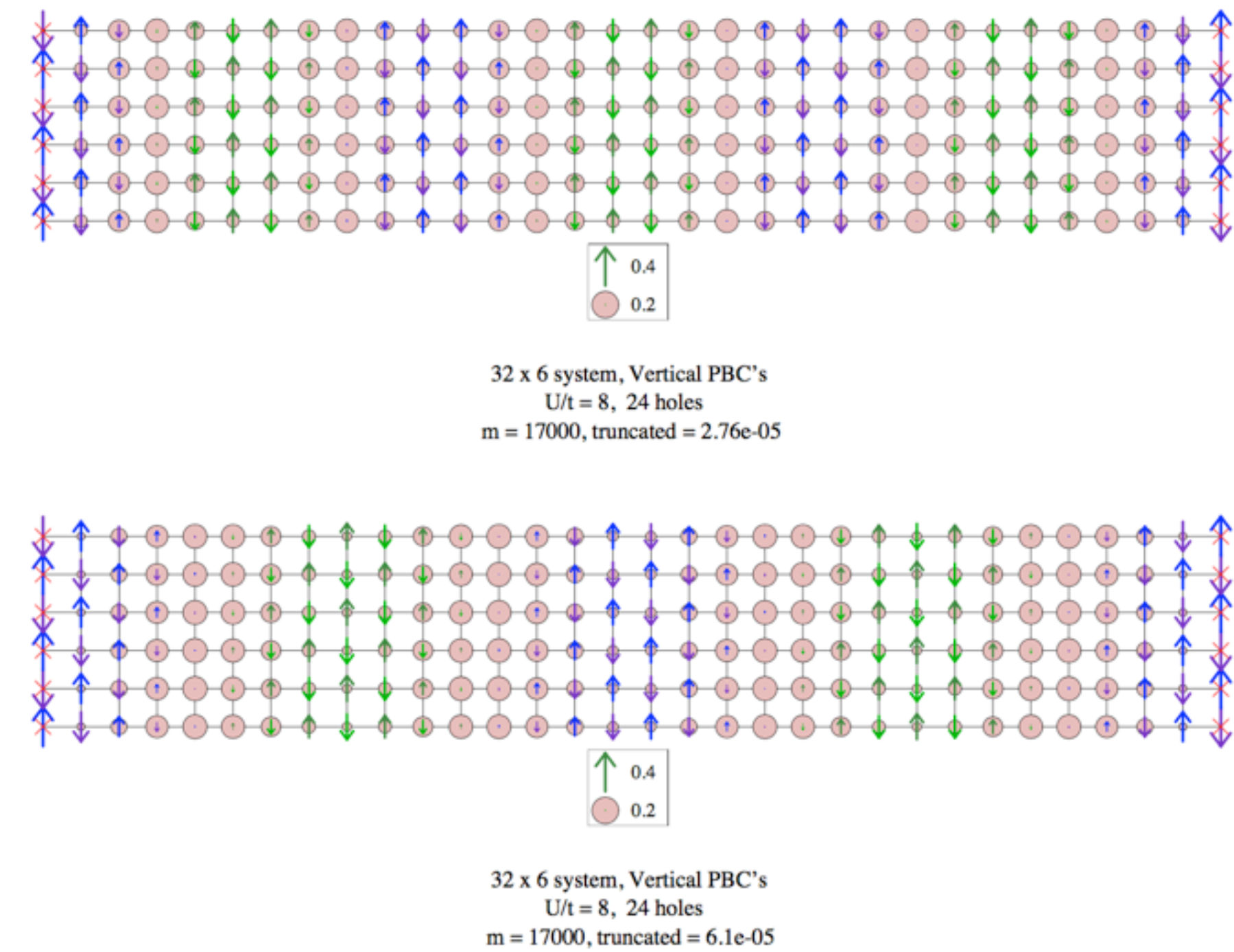
- Density matrix renormalization group (DMRG) (real-space and hybrid real/momentum space)
 - Uncontrolled errors: finite cylinder size
- Density Matrix Embedding theory (DMET) Like cluster DMFT but entanglement-based, not frequency-based. Cluster solved with DMRG (Garnet Chan)
 - Uncontrolled errors: finite cluster size
- Infinite projected entangled pair states (IPEPS)
 - Uncontrolled errors: finite bond dimension and extrapolation
- Constrained Path Monte Carlo (CPMC, AFQMC)
 - Uncontrolled errors: Constraint based on trial wavefunction

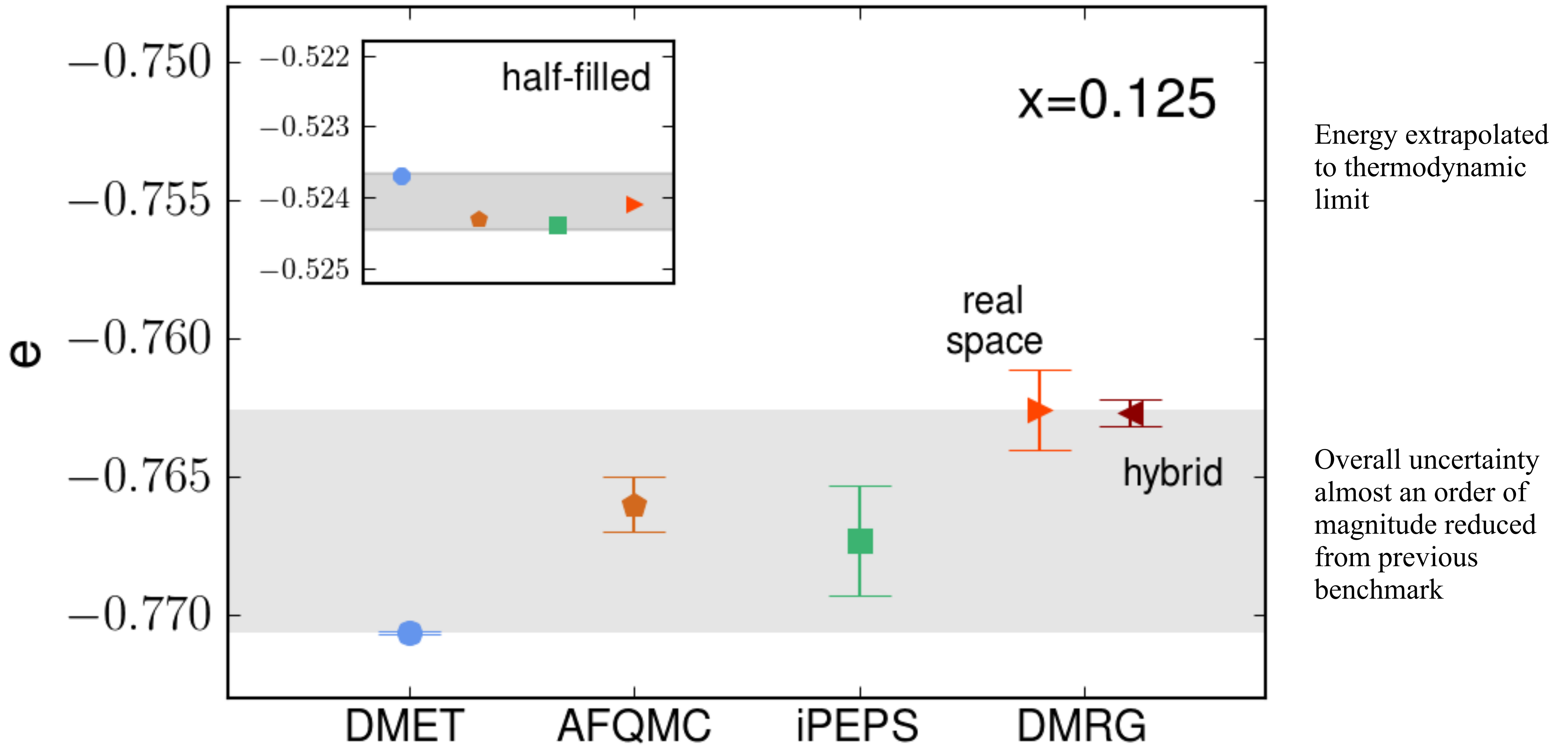
By focusing one one point in the phase diagram, we were all able to improve our results substantially—longer runs, better techniques, new techniques.

A key aspect of the work is that the uncontrolled errors are very different. Thus, if multiple methods agree, we can have high confidence we have the right answer.

Cross Validation: example

- DMRG: performed on cylinders; difficulty increases exponentially with cylinder width w . $w=4$: nearly exact, $w=6$, highly precise, $w=8$: out of reach
- CPMC: The constraint error is surprisingly small even with a noninteracting trial state. But no internal determination of size of error; hard to improve systematically.
- Cross validation: check CPMC error with DMRG on $w=4, 6$ (very small). Then we can trust CPMC on $w=8-12$. We can also estimate an energy correction from DMRG
- Both methods were initially giving stripes, but with different wavelengths. We were able to trace this to a problem in the DMRG analysis: filled stripes were metastable but higher energy at small bond dimension, but dropped below for large m ! After fixing the DMRG, excellent agreement (all four methods).

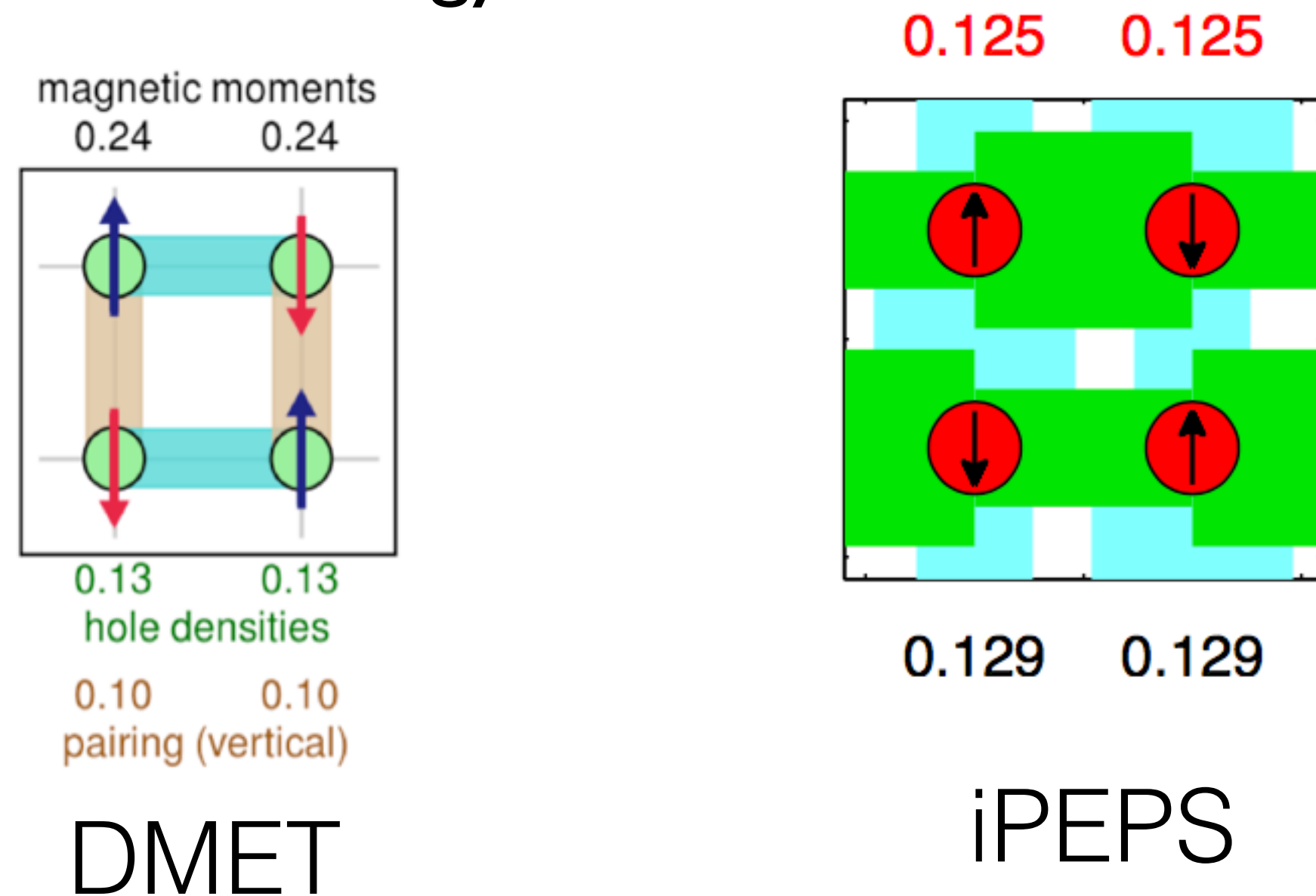
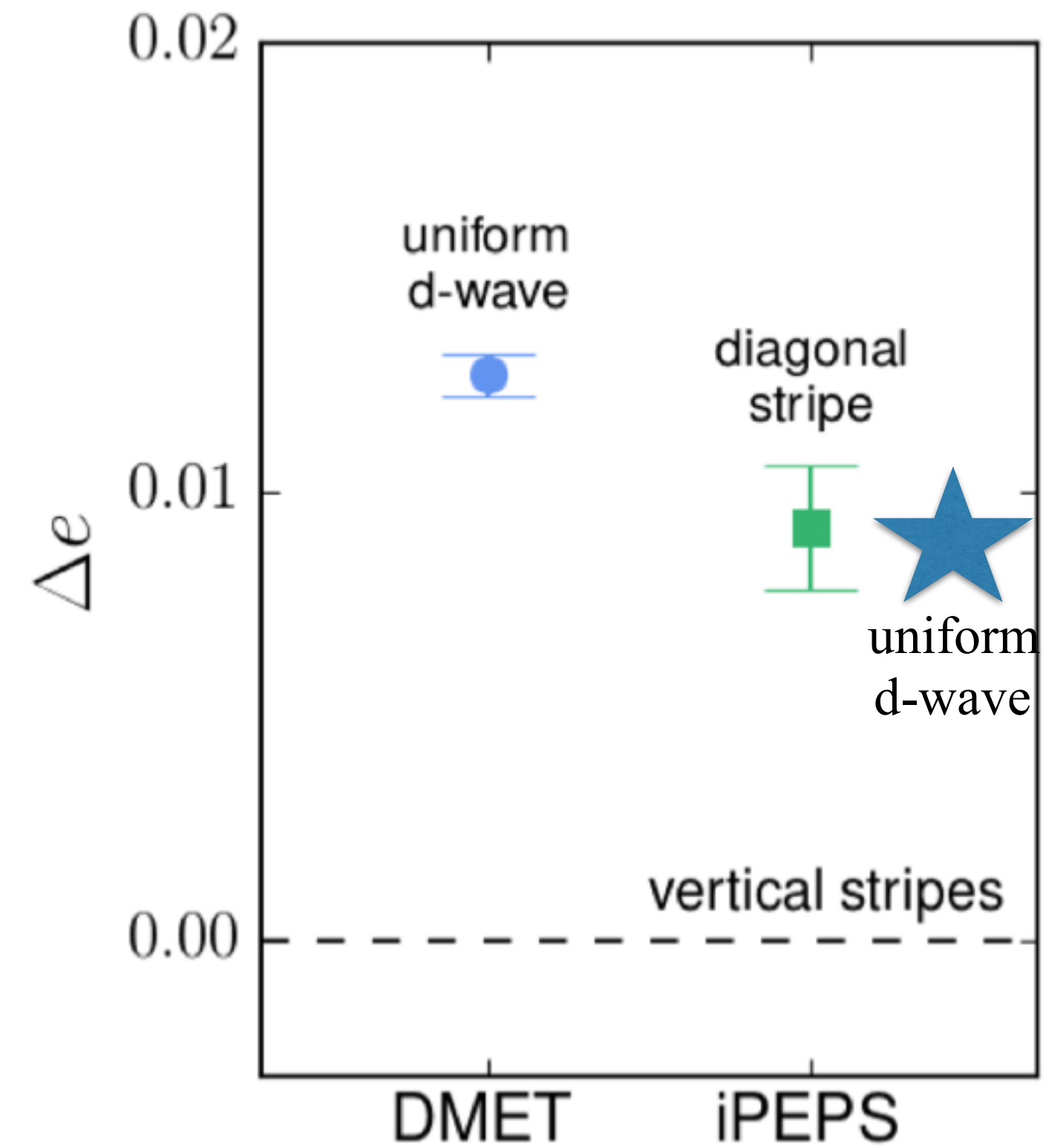




Error bars neglect systematic errors—that is what we need the comparison for.

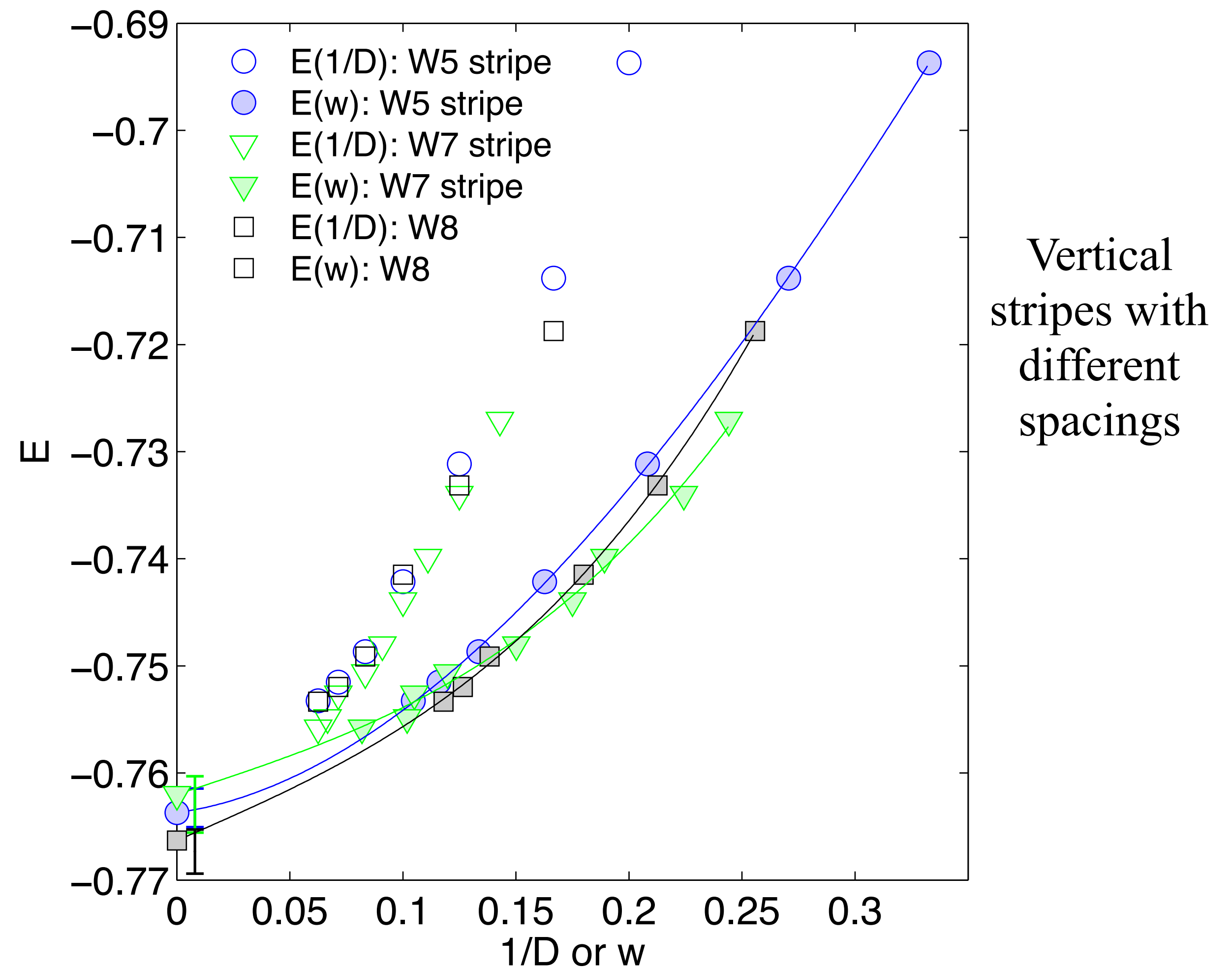
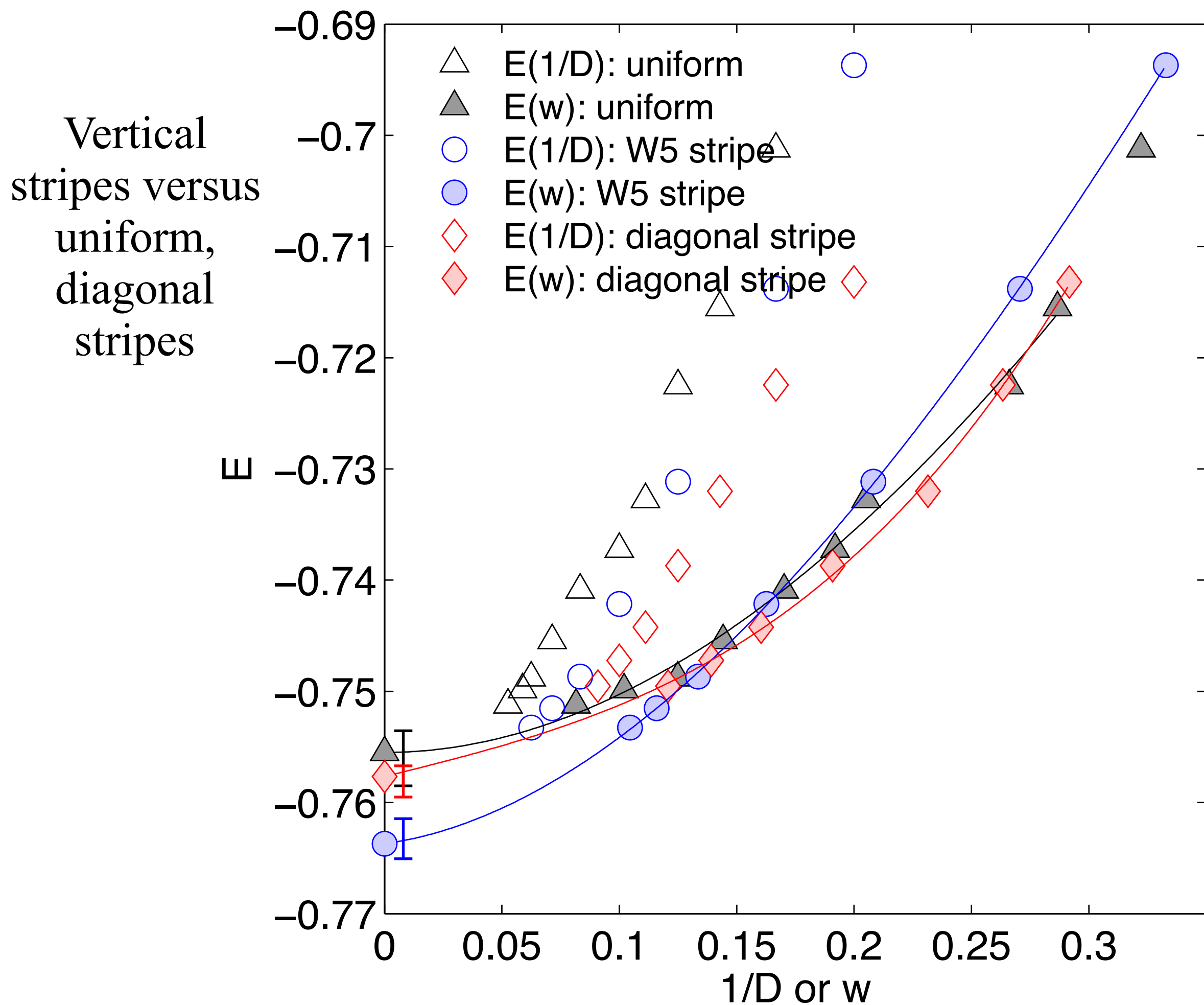
Uniform state versus stripes

- DMET and iPEPS both can be forced to give uniform states:
 - DMET has a cluster size. For a 2x2 cluster, no stripe patterns can form
 - iPEPS similarly has a cluster that is repeated to infinity. A 2x2 cluster cannot have stripes
- DMRG always gives stripes. Currently no way to force a uniform state. CPMC also gives stripes as lowest energy state.



Both uniform states from DMET and iPEPS show d-wave pairing

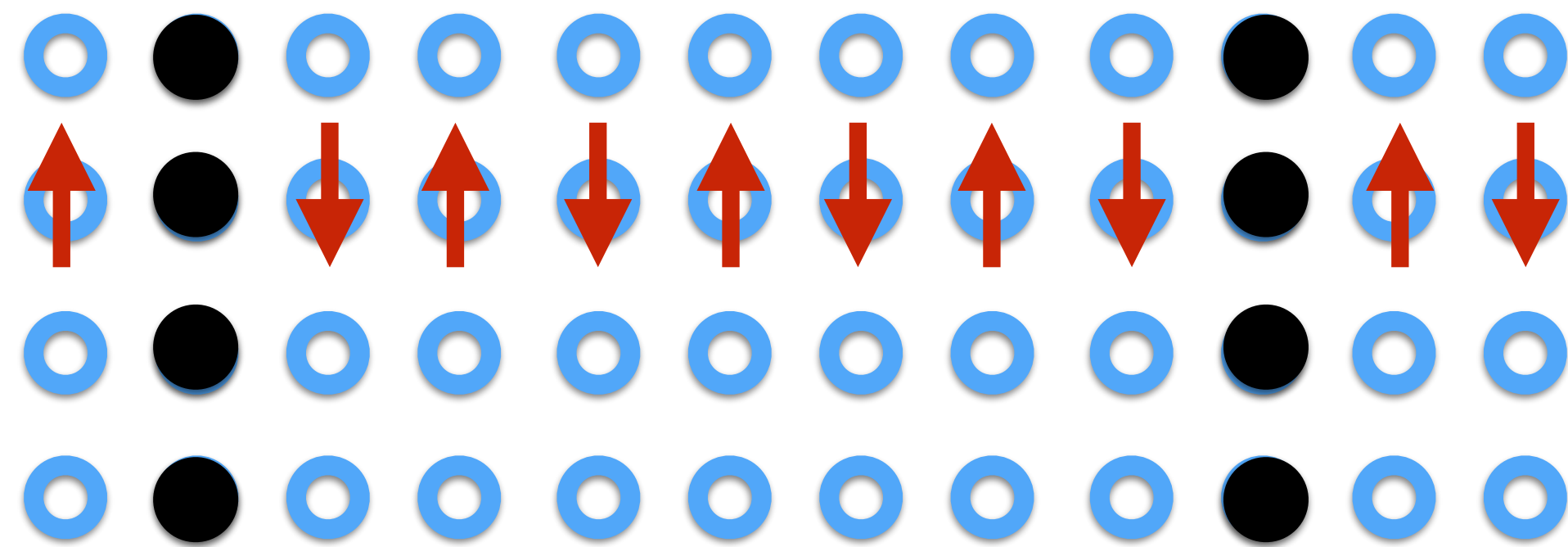
New iPEPS Energy Extrapolation method (Corboz)



Uniform, diagonal stripes higher in energy

Near degeneracy for vertical stripes with different spacings

Vertical stripes: filling, wavelength

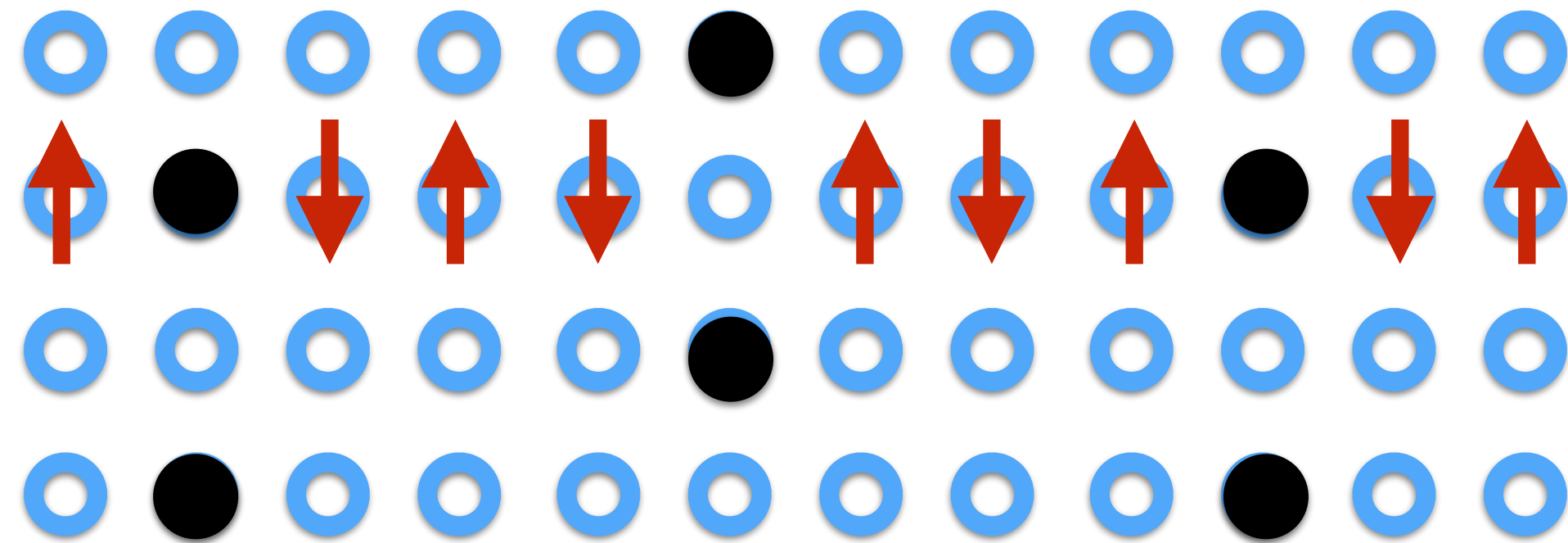


Filled Stripe

$$f = 1$$

$$\lambda = 8$$

The magnetic wavelength is 2λ



Half filled Stripe

$$f = 1/2$$

$$\lambda = 4$$



$$E = -6.1321$$

$$m = 40$$

Filled stripes were found with Hartree Fock in late 80's—but HF may not be not accurate

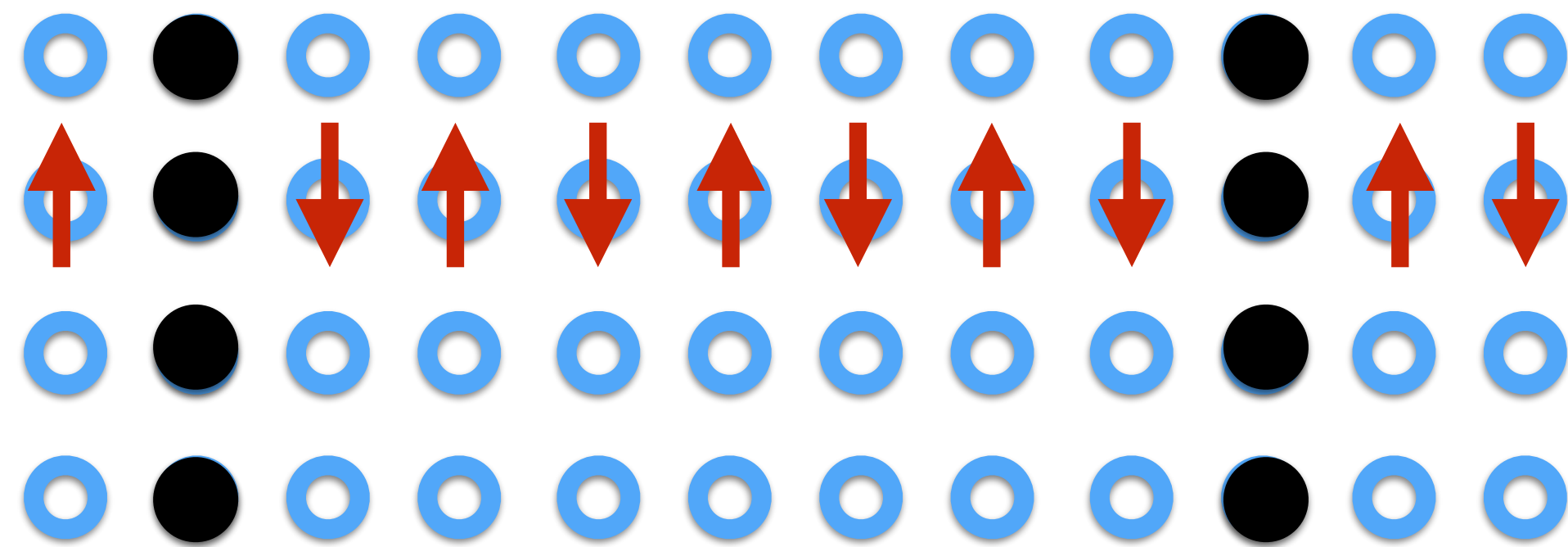
Zaanen

Half-filled stripes were found in some cuprates in the mid 90's. A few years later, DMRG on the t-J model showed half-filled stripes (White & Scalapino)

Tranquada

DMRG on the t-J model—
formation of two half-filled stripes

Vertical stripes: filling, wavelength

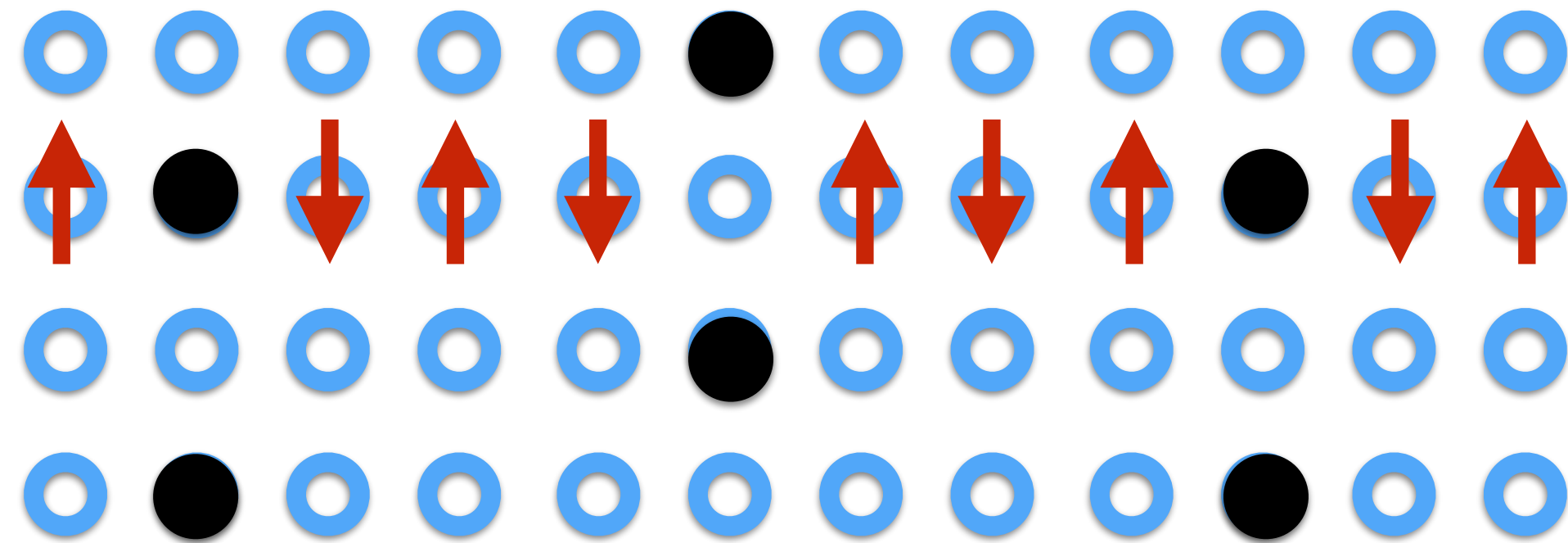


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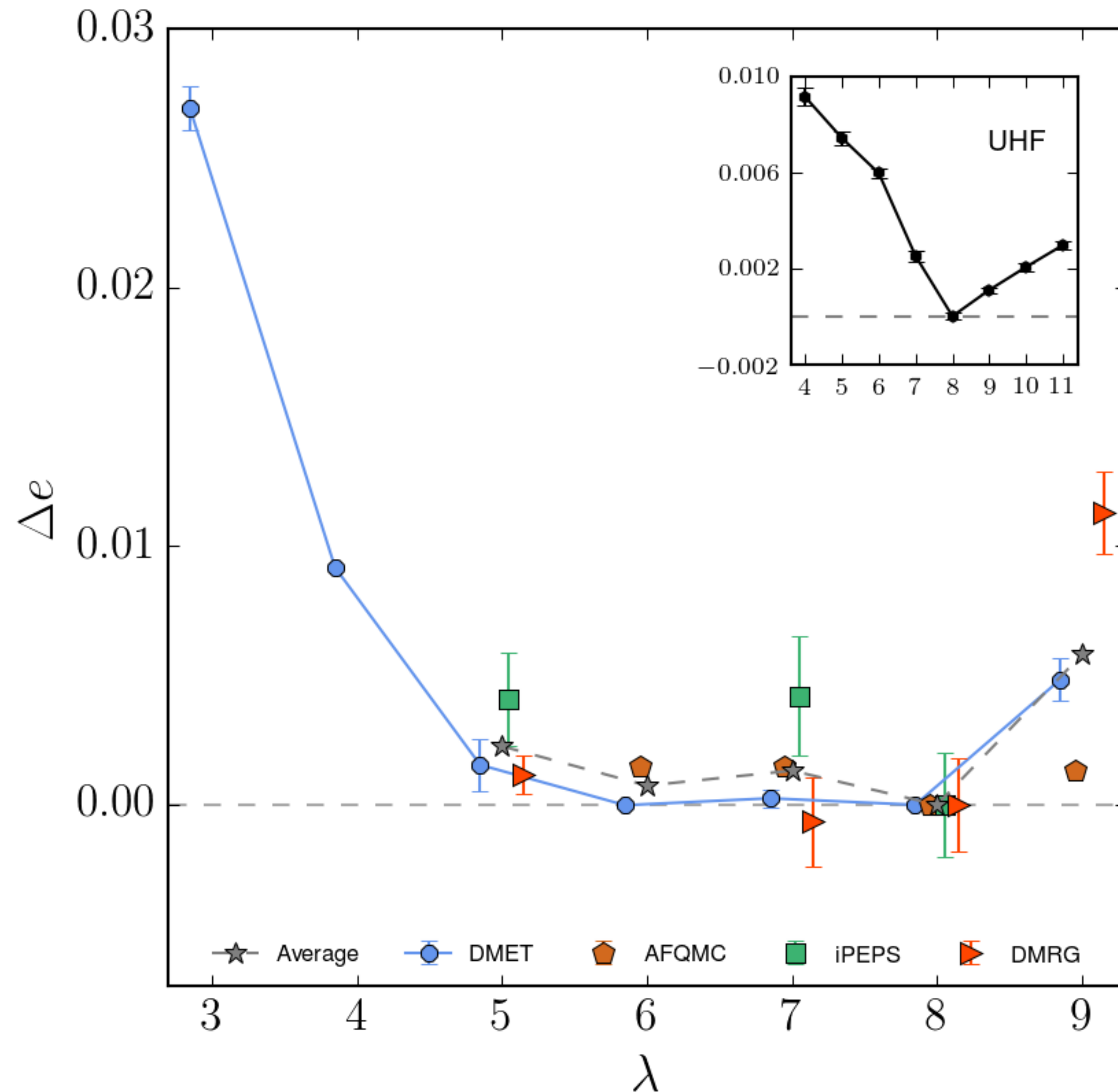
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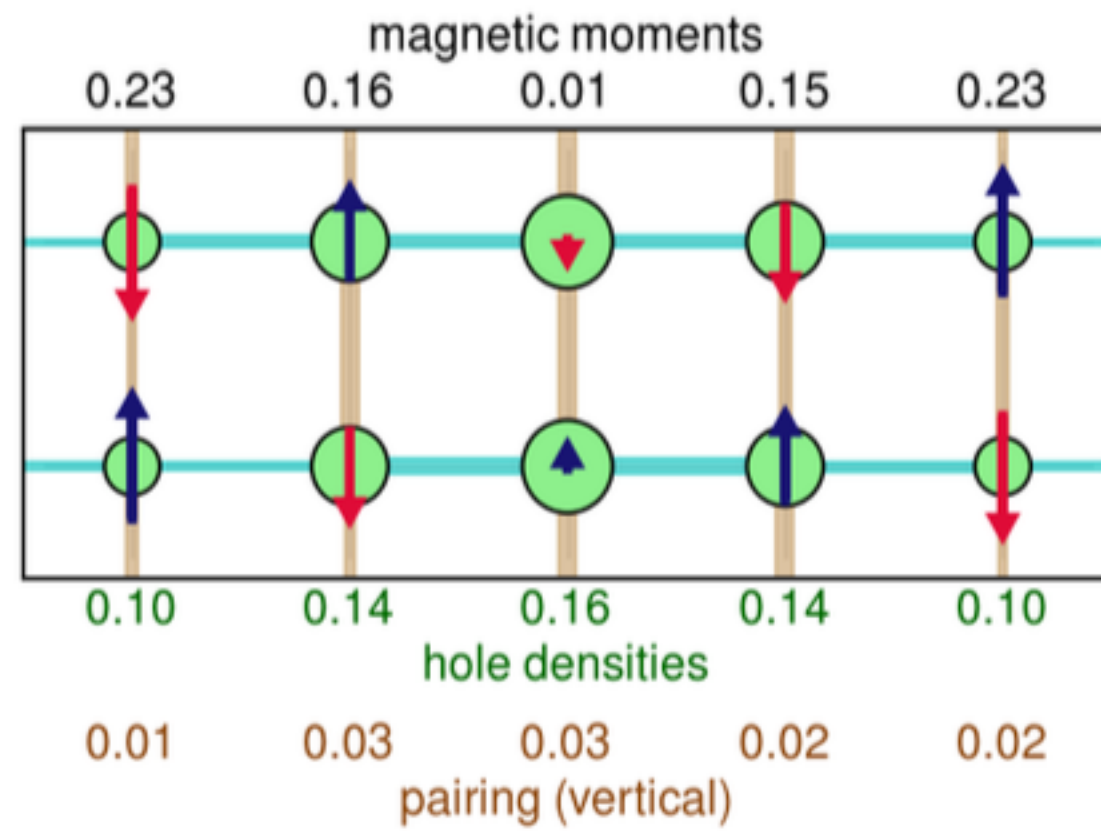
Vertical stripes: Energy versus wavelength



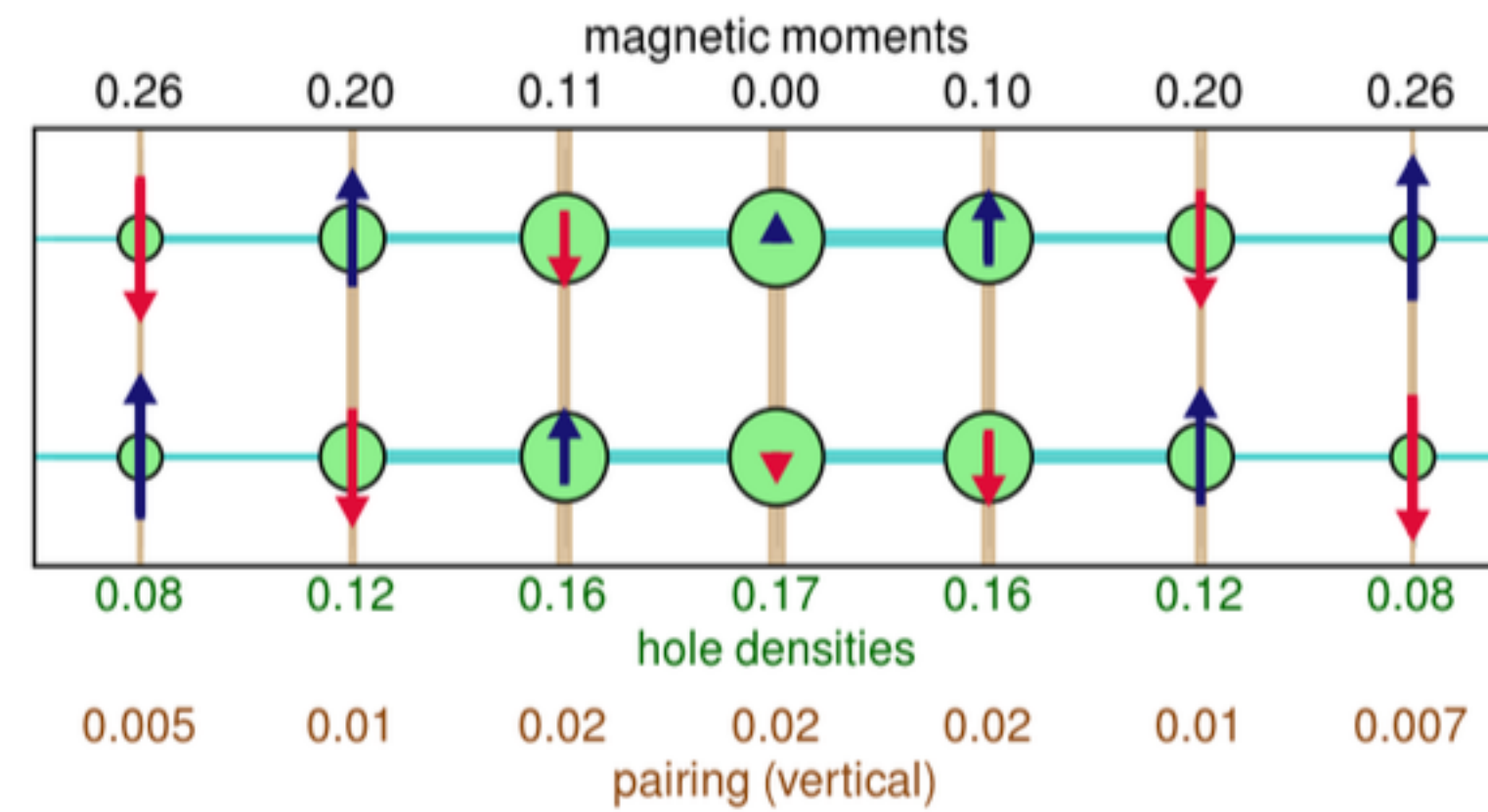
We find a remarkable near-degeneracy for states with different stripe wavelengths, with $\lambda=8$ very slightly lower in energy, and $\lambda=4$ significantly higher.

The near degeneracy likely points towards disordered stripes and/or fluctuating stripes.

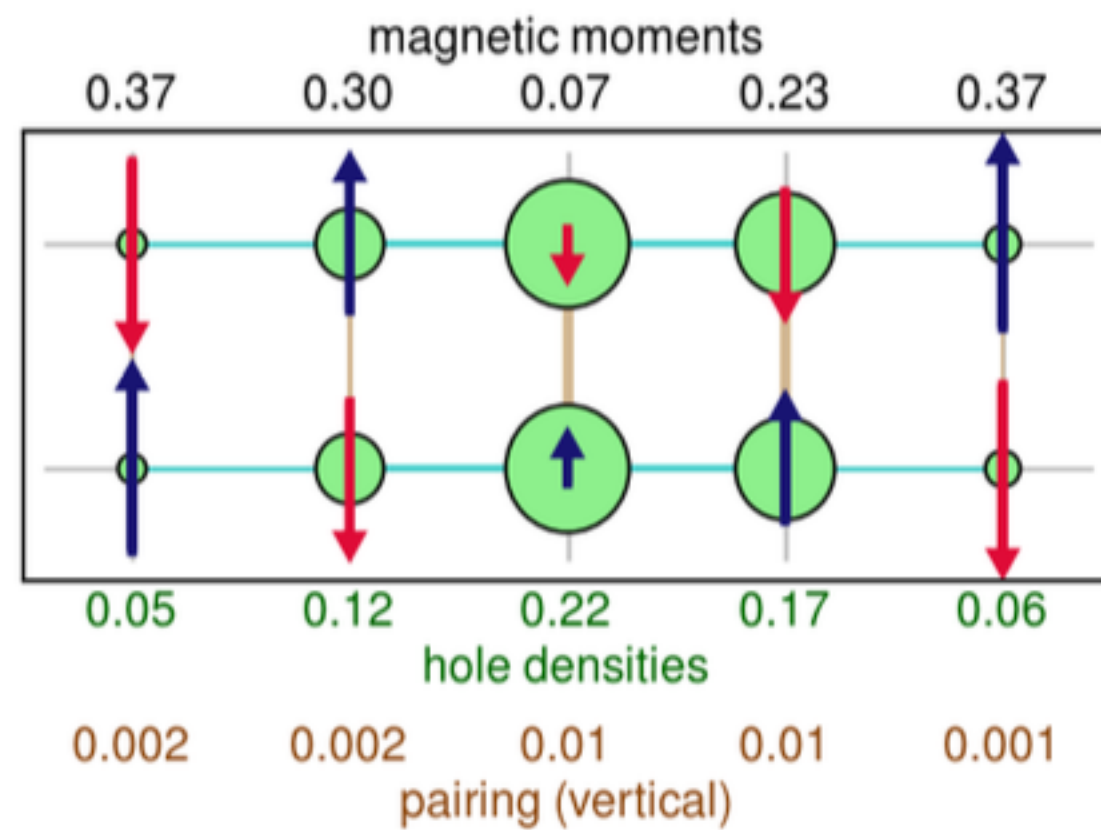
Pairing with partially filled stripes



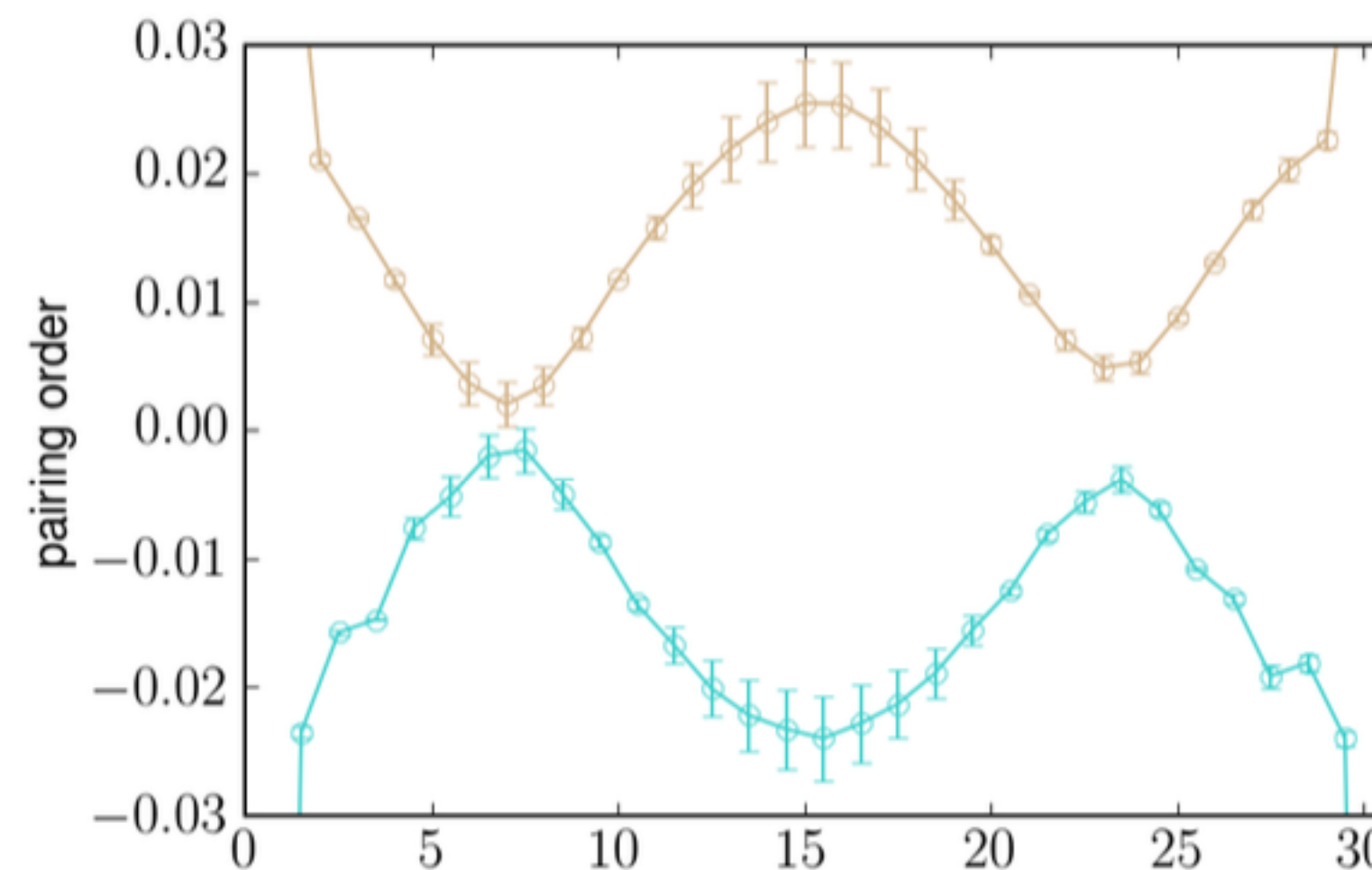
(a) iPEPS $\lambda = 5$



(b) iPEPS $\lambda = 7$



(c) DMET metastable $\lambda = 5$



(d) DMRG pairing order parameters.

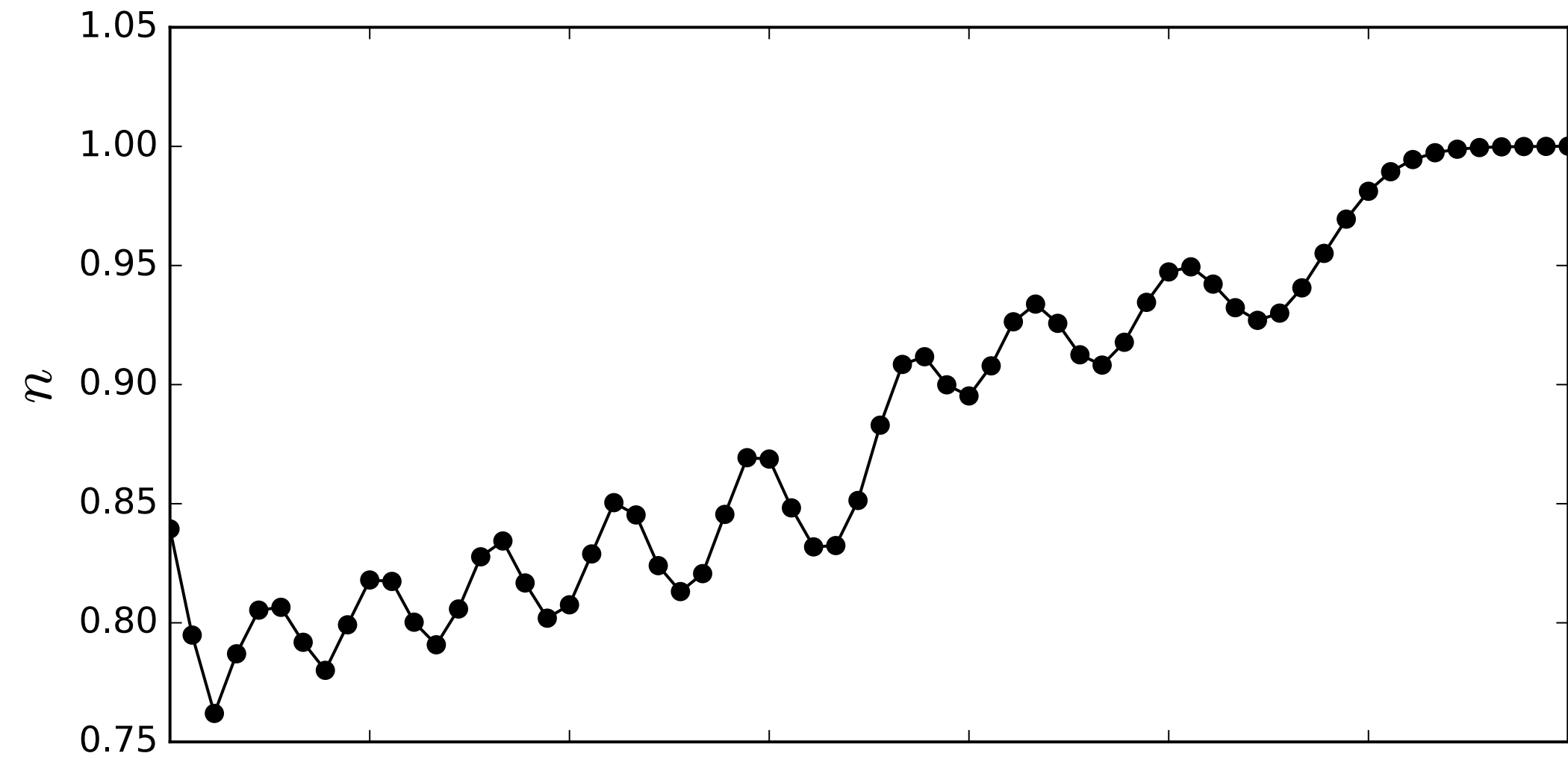
32x4, edge pairing
fields, stripes smeared

For filled stripes,
no d-wave pairing

For partially filled
stripes, d-wave
pairing is seen but
not consistently

Pairing and stripes “intertwined” on Lx4 cylinders (DMRG)

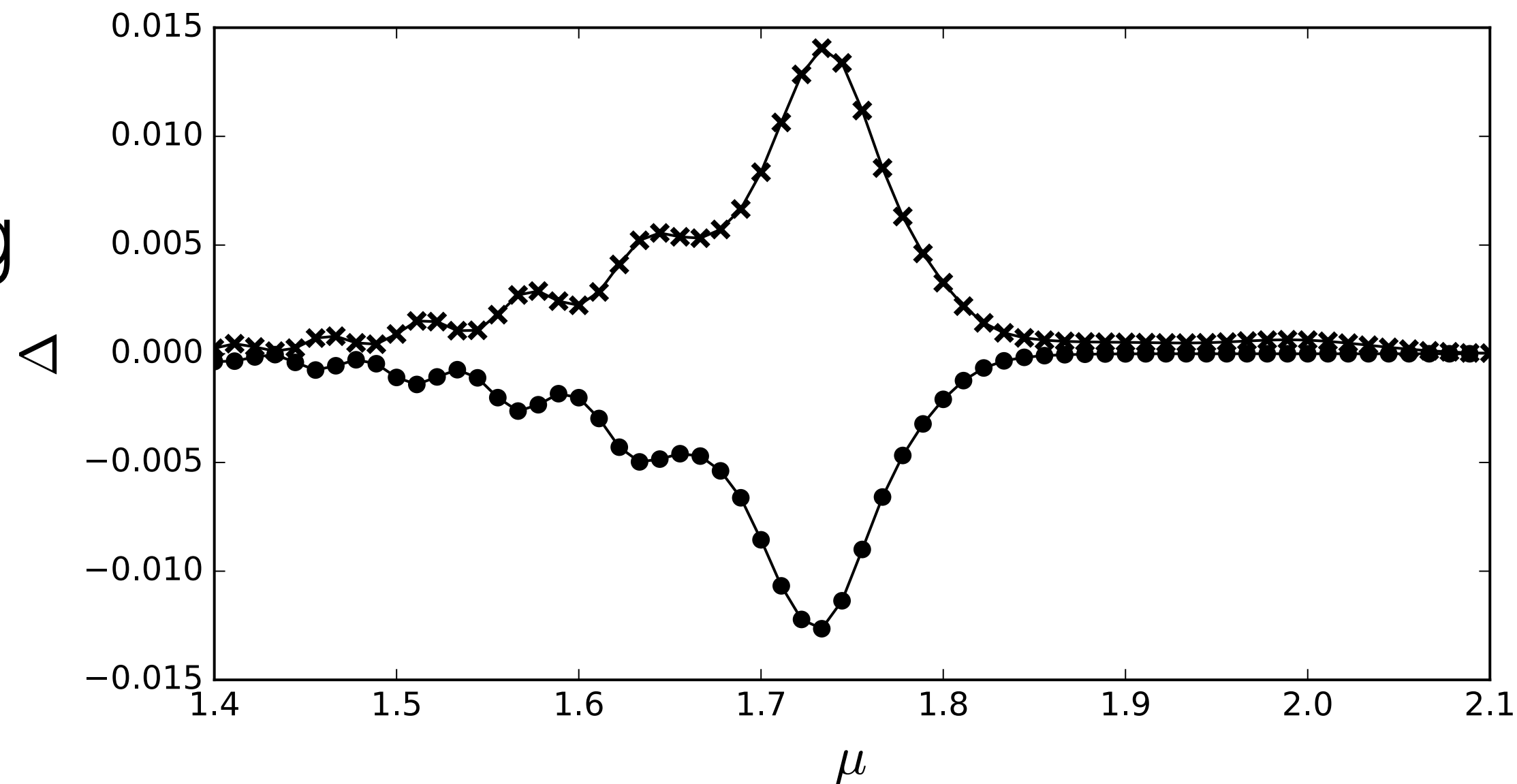
density



Grand canonical simulation so pairing order parameter could be measured locally.

This was a long cylinder with the chemical potential linearly varying with position.

pairing



We see a peak in pairing near optimal doping but coexisting with stripes

← position along cylinder →



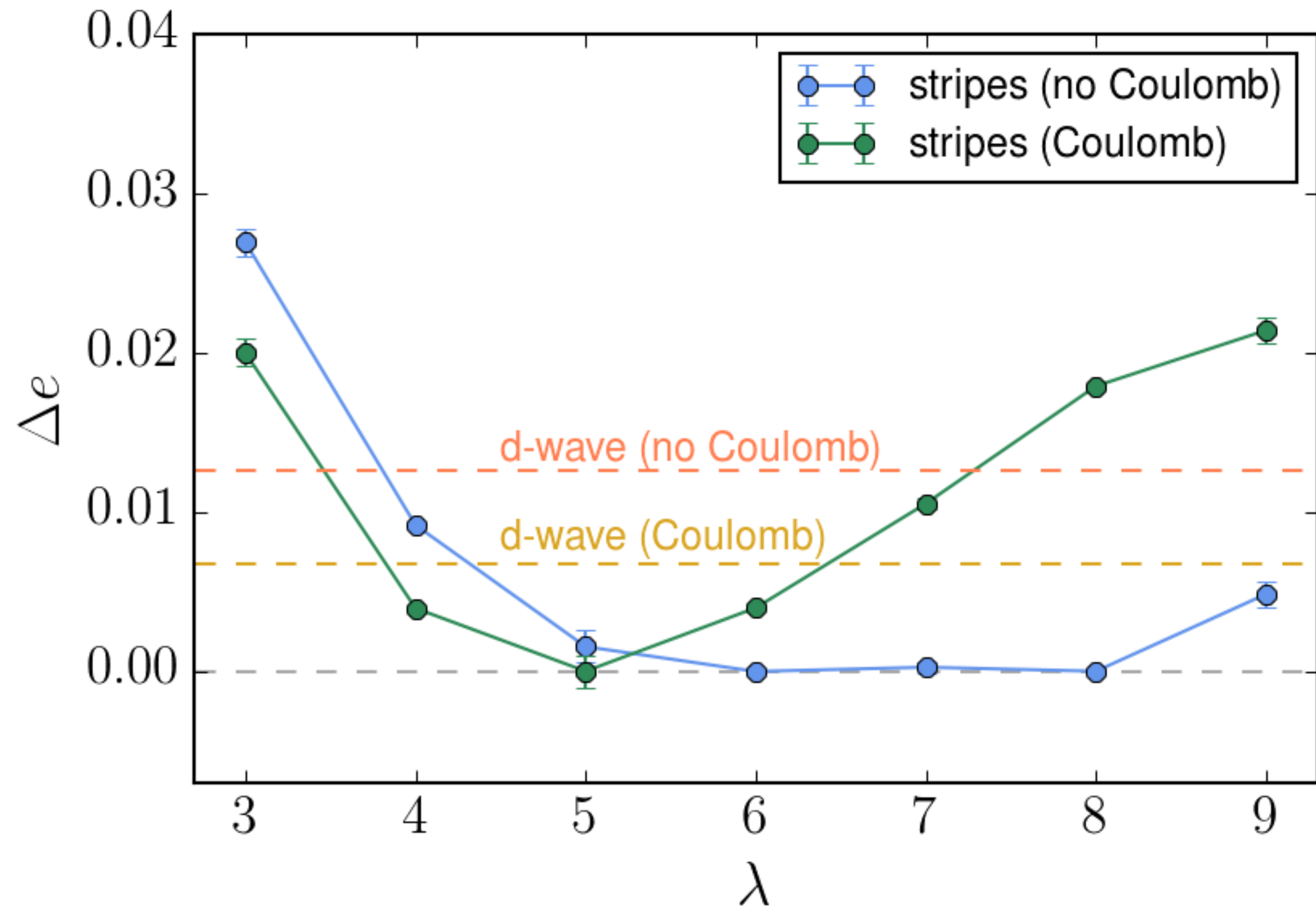
Chia-Min Chung

Solving the Hubbard model: where do we stand?

- Our energy resolution of $0.004t$ corresponds $\pm 10K$ per site. Since superconductivity occurs $\sim 100K$ in the cuprates, this should be enough to understand high- T_c SC!
- Using four different methods with very different uncontrolled errors, we have converged to a consistent general picture of the Hubbard ground state at perhaps the most difficult, important point in the phase diagram
- In applying this to the cuprates, the uncertainty in the Hamiltonian is now central. Small changes to the Hamiltonian from additional terms can change the competition between different phases

Extra Slides

Static corrections from long-range Coulomb



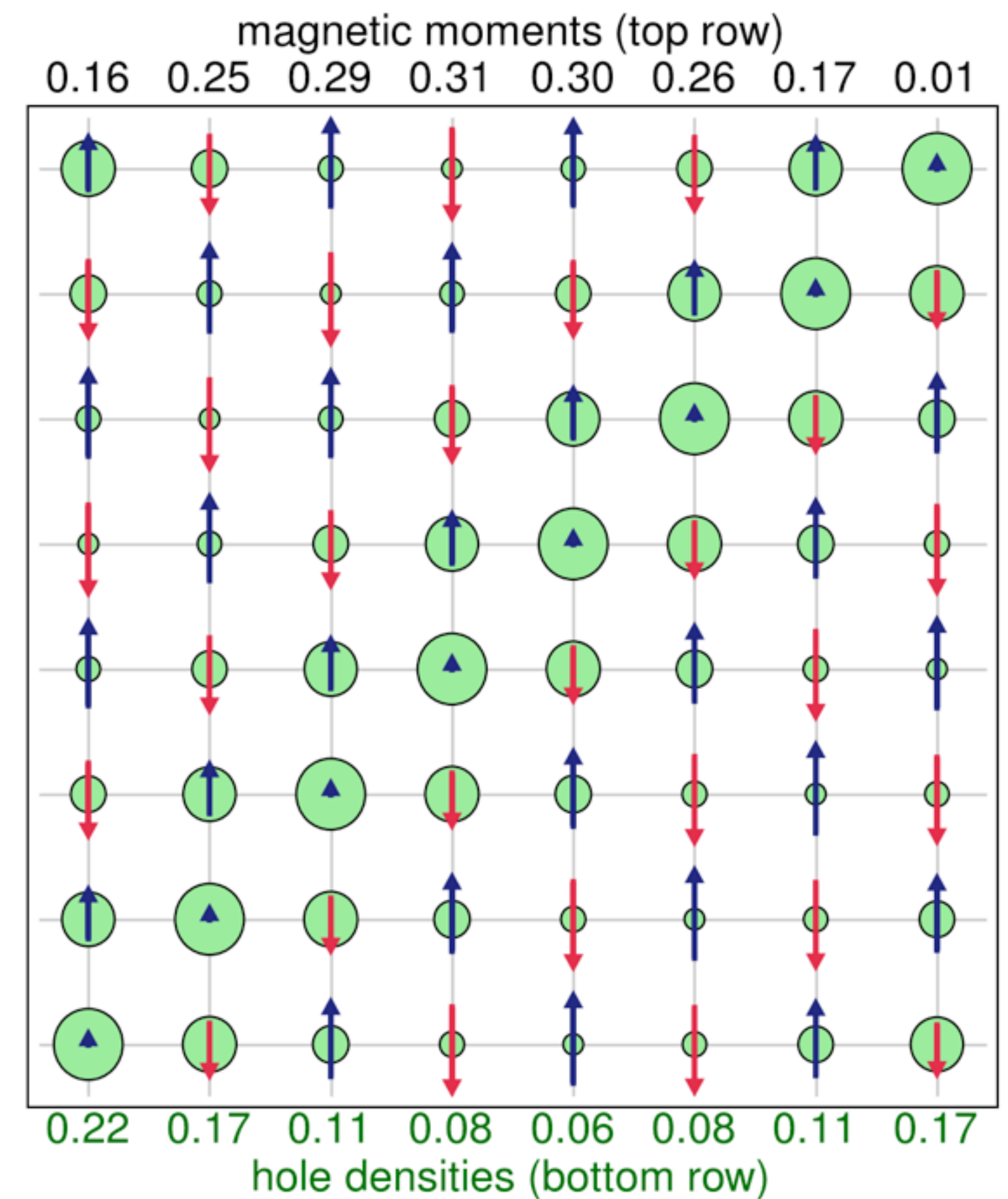
Assume a substantial dielectric constant inducing screening

Calculate energy correction by integrating up the Coulomb contribution from each density pattern. Coulomb favors lower-filled stripes, and favors uniform states

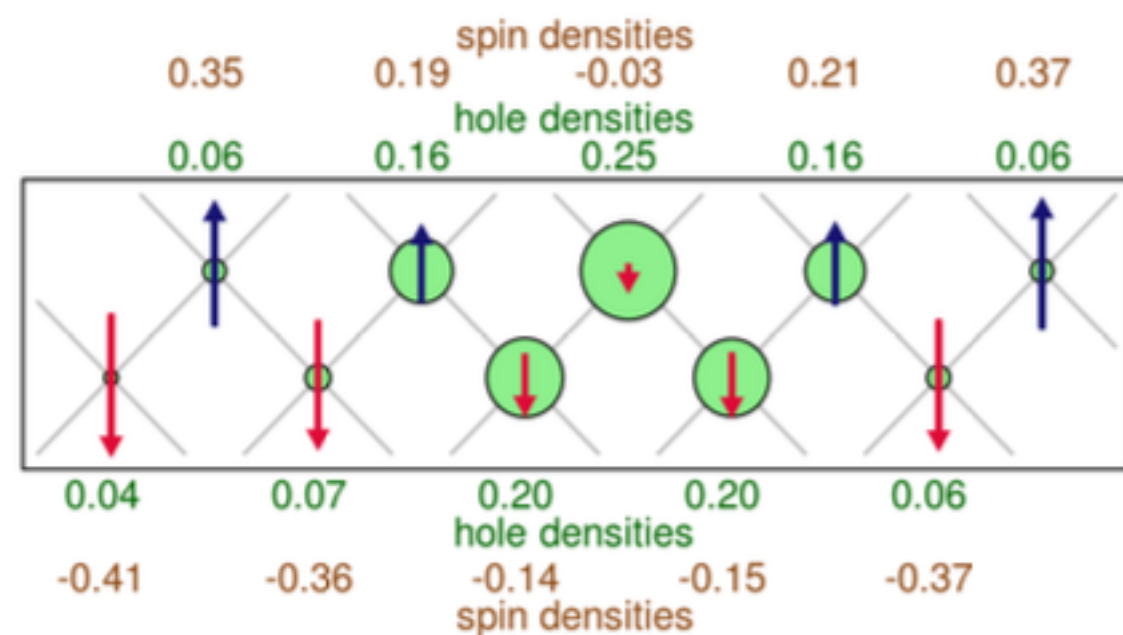
This correction drives our results closer to the cuprates.

Vertical versus diagonal stripes

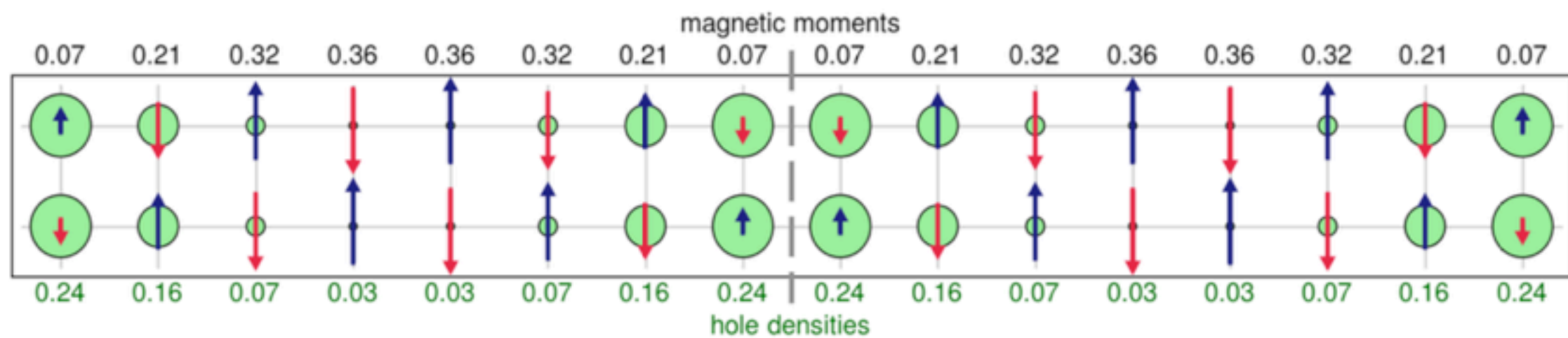
- DMET and iPEPS both can be forced to give diagonal striped states
- Both give higher energies, by $\sim 0.005 t$
- The diagonal stripes are “filled”: one hole per root-2 distance
- Boundary conditions on the cylinders used by DMRG and CPMC frustrate diagonal stripes, and they were not seen.



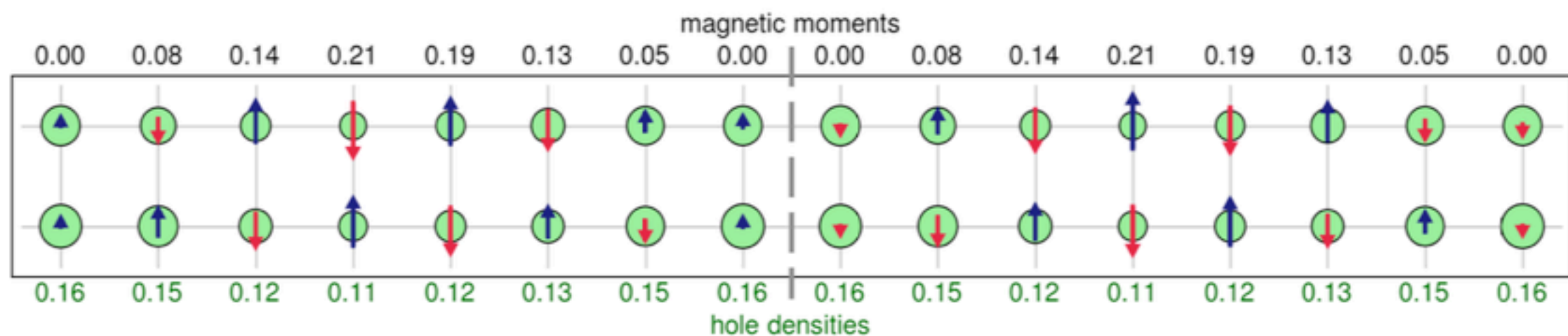
iPEPS



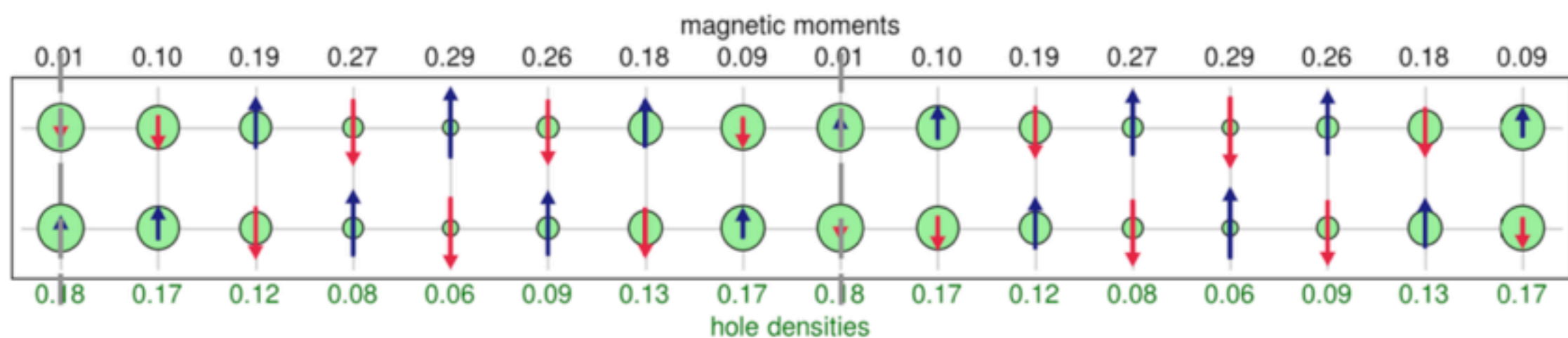
DMET



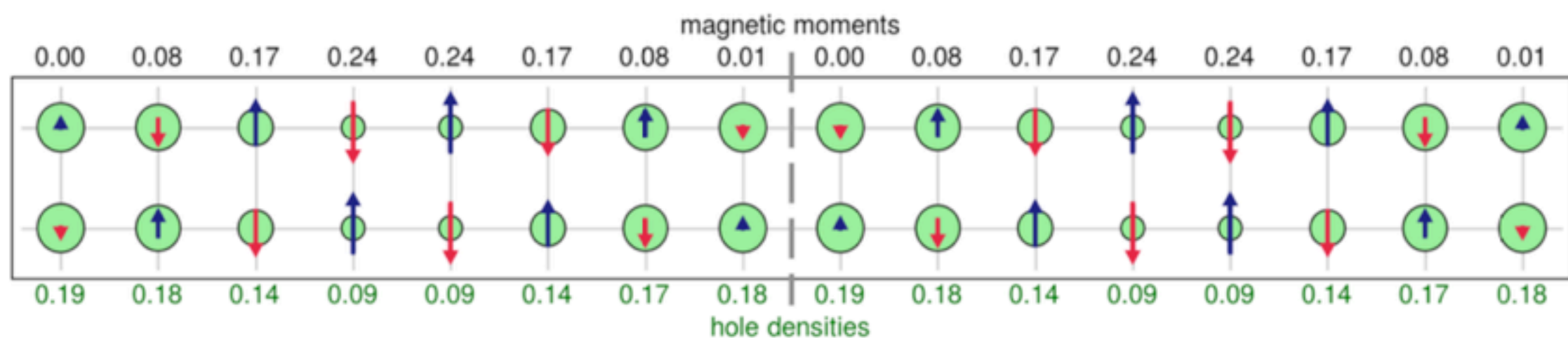
(a) DMET



(b) AFQMC



(c) iPEPS



(d) DMRG

Filled stripes as seen by all four methods

