

Quenched Disorder and Vestigial Nematicity in the Pseudogap Regime of Cuprates



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Abstract

- We carried out a theoretical analysis of the Landau-Ginzburg-Wilson effective field theory of a classical incommensurate CDW in the presence of weak quenched disorder.
- Although the possibility of long-range CDW order is precluded in such systems, any discrete symmetry-breaking aspect of the charge order (nematicity in our case) generically survives up to a nonzero disorder strength.
- Such "vestigial order", which is subject to unambiguous

QCP around doping 0.19?

- Various experiments have indicated a possible quantum critical point in cuprates at around doping 0.19[1,2].
- With disorder taken into account, the QCP (if exists) cannot be • incommensurate-CDW type.
- Discrete symmetry breaking QCP is possible, e.g., nematic or time reversal symmetry breaking.

macroscopic detection, can serve as an avatar of what would be CDW order in the zero disorder limit.

Model

charge density
$$\rho(r) = \overline{\rho} + \left[\psi_x(r)e^{iQ_xr} + \psi_y(r)e^{iQ_yr} + c.c\right] + \dots$$

$$\begin{aligned} H_{\rm CDW} &= -J \sum_{\langle r,r'\rangle,m} \left[\psi^{\dagger}(r,m)\psi(r',m) + h.c. \right] - V_z \sum_{r,m} \left[\psi^{\dagger}(r,m)\psi(r,m+1) + h.c. \right] \\ &+ U \sum_{r,m} [\psi^{\dagger}_x \psi_x + \psi^{\dagger}_y \psi_y - 1]^2 - \Delta \sum_{r,m} [\psi^{\dagger}_x \psi_x - \psi^{\dagger}_y \psi_y]^2 - \sum_{r,m} [h^{\dagger}\psi + h.c.] \end{aligned}$$

 $\overline{h_{\alpha}(r,m)h_{\beta}(r',m)} = \sigma^2 \delta_{\alpha\beta} \delta_{m,m'} \delta(r-r')$ h(r,m): random field

• Apply replica trick and integrate out h; perform Hubbard-Stratonovich transformation and solve self-consistent equations:

$$1 = T \int d^3k \left[\frac{1}{E_x(k,\mathcal{N})} + \frac{1}{E_y(k,-\mathcal{N})} \right] + \sigma^2 \int d^3k \left[\frac{1}{E_x^2(k,\mathcal{N})} + \frac{1}{E_y^2(k,-\mathcal{N})} \right]$$

$$-\frac{\mathcal{N}}{\Delta} = T \int d^3k \left[\frac{1}{E_x(k,\mathcal{N})} - \frac{1}{E_u(k,-\mathcal{N})} \right] + \sigma^2 \int d^3k \left[\frac{1}{E_x^2(k,\mathcal{N})} - \frac{1}{E_u^2(k,-\mathcal{N})} \right]$$

Ongoing work (in collaboration with Prof. Subir Sachdev)

$$\begin{split} H_{\rm \scriptscriptstyle SC-CDW} &= H_{\rm \scriptscriptstyle CDW} + g \sum_{r,m} |\psi(r,m)|^2 + g' \sum_{r,m} |\psi(r,m)|^4 \\ &- K \sum_{\langle r,r'\rangle,m} \left[\psi^{\dagger}_{\rm \scriptscriptstyle sc}(r,m) \psi_{\rm \scriptscriptstyle sc}(r',m) + h.c. \right] - V'_z \sum_{r,m} \left[\psi^{\dagger}_{\rm \scriptscriptstyle sc}(r,m) \psi_{\rm \scriptscriptstyle sc}(r,m+1) + h.c. \right] \\ &\psi_{\rm \scriptscriptstyle sc} : \text{superconductivity order parameter} \quad \text{constraint} : \ |\psi_x|^2 + |\psi_y|^2 + |\psi_{\rm \scriptscriptstyle sc}|^2 = 1 \end{split}$$



Fig. 4 Left: solid and dashed lines are second and first order transitions, respectively. The stripe (SC+stripe) phase becomes nematic (SC+nematic) at finite disorder. **Right**: Different parameter regions. Region 1 is the same as the parameters in ref[3].

CDW structure factor in Region 1

 $\lfloor L_x(n, \mathbf{v}) - L_y(n, \mathbf{v}) \rfloor = J = \lfloor L_x(n, \mathbf{v}) - L_y(n, \mathbf{v}) \rfloor$ where \mathcal{N} is the nematic order parameter (one of the HS fields).

Isotropic	Nematic	Stripe	Checkerboard
$\overline{\langle \psi_x \rangle} = 0,$	$\overline{\langle \psi_x \rangle} = 0,$	$\overline{\langle \psi_x \rangle} \neq 0,$	$\overline{\langle \psi_x \rangle} = \overline{\langle \psi_y \rangle} \neq 0$
$\overline{\langle \psi_y \rangle} = 0,$	$\overline{\langle \psi_y \rangle} = 0,$	$\overline{\langle \psi_y \rangle} = 0,$	$\mathcal{N} = 0$
$\mathcal{N} = 0$	$\mathcal{N} \neq 0$	$\mathcal{N} \neq 0$	

Results





Fig. 5 CDW structure factor as a function of temperature with input parameters from Region 1. Left: under increasing Vz with zero disorder. Right: under increasing disorder with fixed $V_z=0.01J$. Dashed lines mark superconductivity transitions.



Fig. 1 Zero-disorder phase diagram. Solid and dashed lines are second and first order transitions, respectively.

Fig. 2 Finite-disorder phase diagram at *Vz*=0.01*J*. Nematic phase survives up to a critical disorder strength.

CDW correlation length



Fig. 3 T-dependence of CDW correlation length under various disorder strengths, with (thin lines) or without (thick lines) an explicit symmetry breaking field.

- Nematic transition greatly enhances the CDW structure factor.
- SC transition suppresses CDW correlation.
- Nematicity always wins when the two factors compete, resulting structure factor as temperature decreases (no peak structure).

 $\sigma = 1.6, V_z = 0.01$ References $\sigma = 1.9, V_z = 0.01$ [1]K. Fujita et al., Science 344, 612 (2014) [2]B. Ramshaw et al., arXiv 1409.3990 (2014) [3]L. Hayward et al., Science 343, 1336 (2014)