

# High-energy scattering and locality: gravity and strings

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Based on:

[hep-th/0604072](#);

[hep-th/0703116](#);

[arXiv:0705.1816](#), w. Gross and Maharana;

[arXiv:0705.2197](#);

WIP w/ M. Gary; D. Marolf; M. Srednicki

One of the cornerstones of Quantum  
Field Theory is locality

Yet, we suspect not strictly valid.

Some reasons:

BH information paradox (most acceptable resolution)  
(e.g. holography - rather extreme form)

Hints from strings: extendedness; AdS/CFT

Expected breakdown of classical spacetime

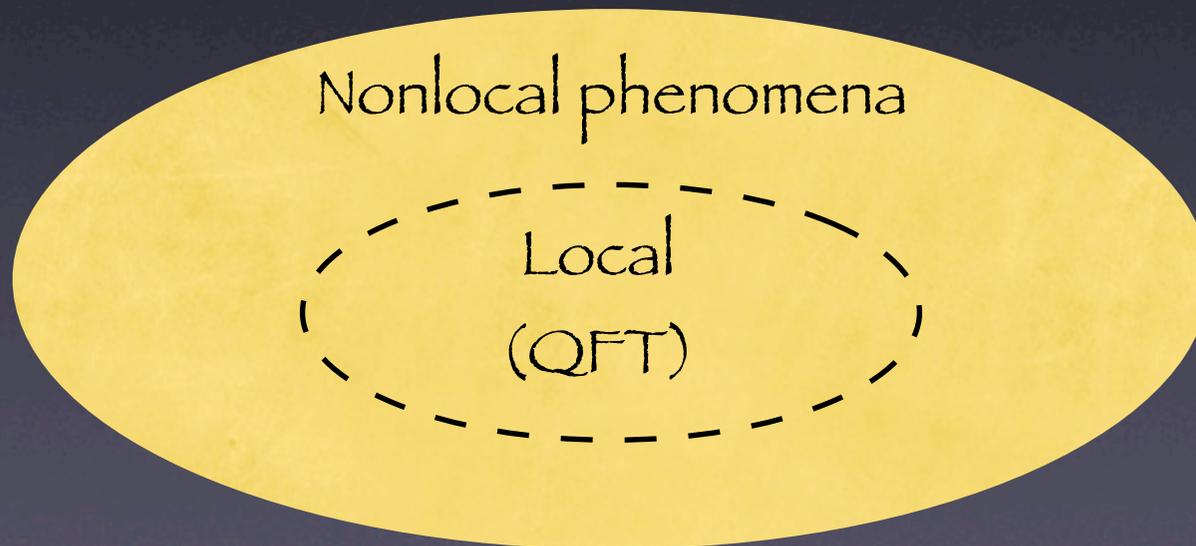
...

## Some questions:

1) What is the mechanism for breakdown?

extendedness of strings (branes ...)?

2) What is the “correspondence boundary”?



These questions are clearly related.

One way to address them is via high-energy scattering

(There's a related story more closely based on observables, etc.)

E.g. is the basic mechanism string extendedness?

Possible picture in HE scatt:



Long strings  
 $L \sim E/M_s^2$

$\longleftrightarrow$  String uncertainty principle  
(Veneziano, Gross)  $\Delta X \geq \frac{1}{\Delta p} + \alpha' \Delta p$   
( $\longleftrightarrow$  nonlocality)

(Proposed app. to BH info: LPSTU)

What does this have to do w/BH formation?

Does it prevent? (Strominger, Gross)

Or is this BH formation? (~Susskind)

Let's investigate ...

Begin w/tree-level amplitude: high  $E$

$$\mathcal{A}_0^{\text{string}}(s, t) \propto g_s^2 \frac{\Gamma(-t/8)}{\Gamma(1+t/8)} s^{2+t/4} e^{2-t/4}$$

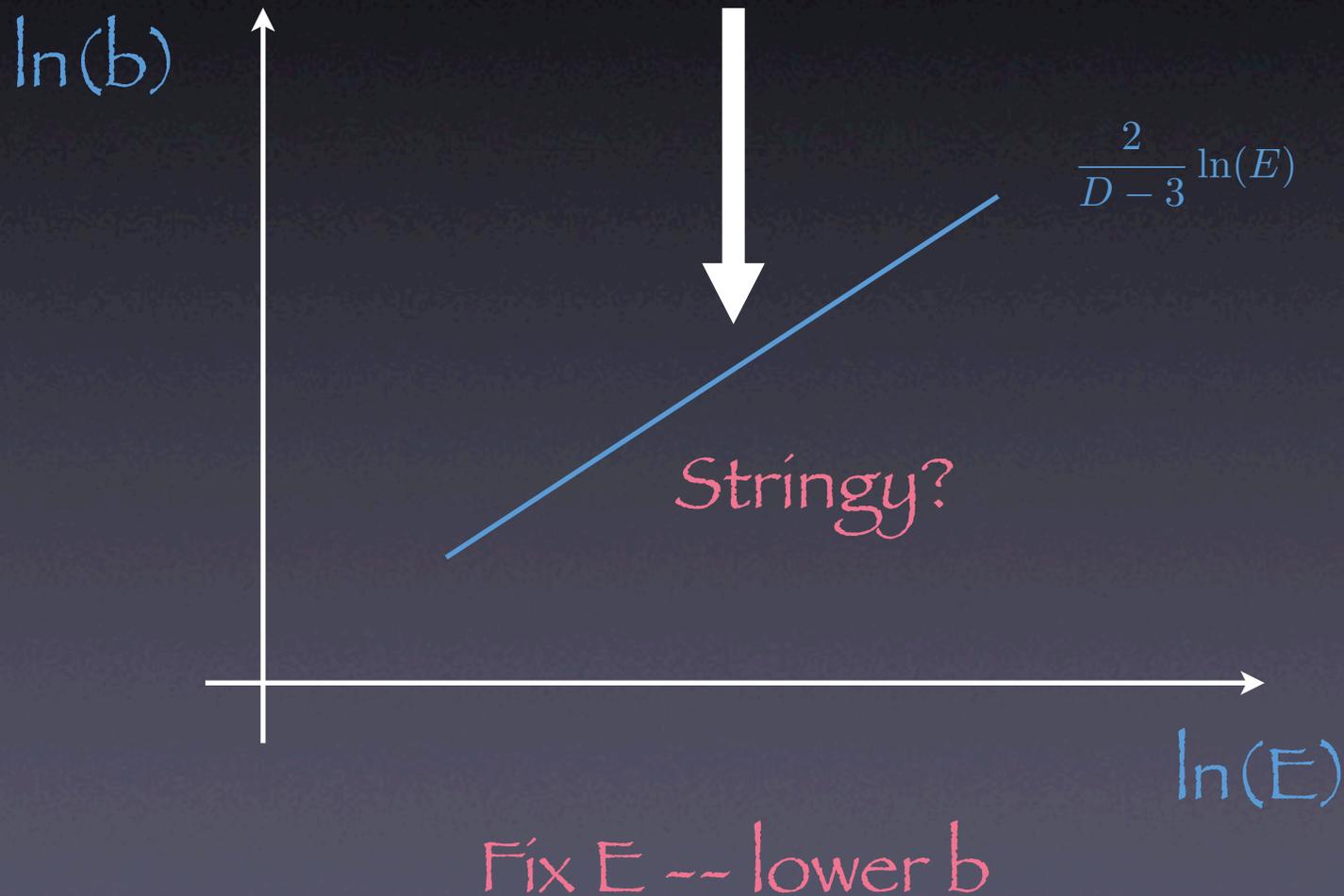
vs.

$$\mathcal{A}_0^{\text{grav}}(s, t) \propto G_D \frac{s^2}{t}$$

- No evidence for long string effects;
- But significant modifications for  $t \sim -1$

To investigate:  $(s,t) \longrightarrow (E,b) \quad E \gg M_P$

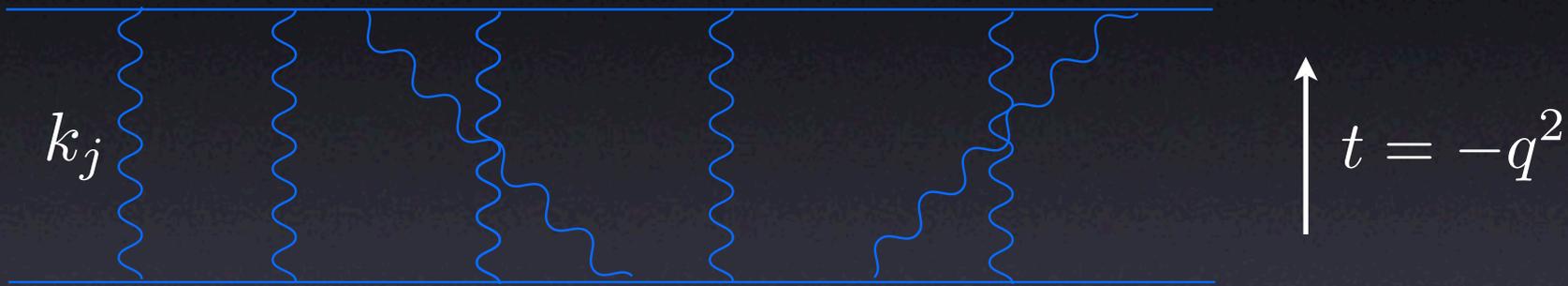
e.g.  $t \sim -1 \Leftrightarrow b \sim E^{2/D-3} \quad (D \text{ noncompact dims})$



To check, include loops:

(Following Amati, Ciafaloni, Veneziano; Muzinich-Soldate; SBG, Gross, Maharana)

Ultrahigh-E: **Eikonal**



$$i\mathcal{A}_N^{\text{string}} = \frac{2s}{(N+1)!} \int \left[ \prod_{j=1}^{N+1} \frac{d^{D-2}k_j}{(2\pi)^{D-2}} \frac{i\mathcal{A}_0^{\text{string}}(s, -k_j^2)}{2s} \right] (2\pi)^{D-2} \delta^{D-2} \left( \sum_j k_j - q_{\perp} \right)$$

$$\prod_{j=1}^{N+1} \frac{E^{2-\alpha'k_j^2}}{k_j^2}$$

$$1) k_j \approx q/(N+1)$$

$$2) E^{-\alpha'q^2}/(N+1)$$

Thus at large  $N$ , string corrections get smaller

Which  $N$  dominates?

Can sum eikonal series:

$$i\mathcal{A}_{\text{eik}}(s, t) = 2s \int d^{D-2} \mathbf{b} e^{-iq_{\perp} \cdot \mathbf{b}} (e^{i\chi(b)} - 1)$$

with 
$$\chi(b) \sim G_D \frac{E^2}{b^{D-4}}$$

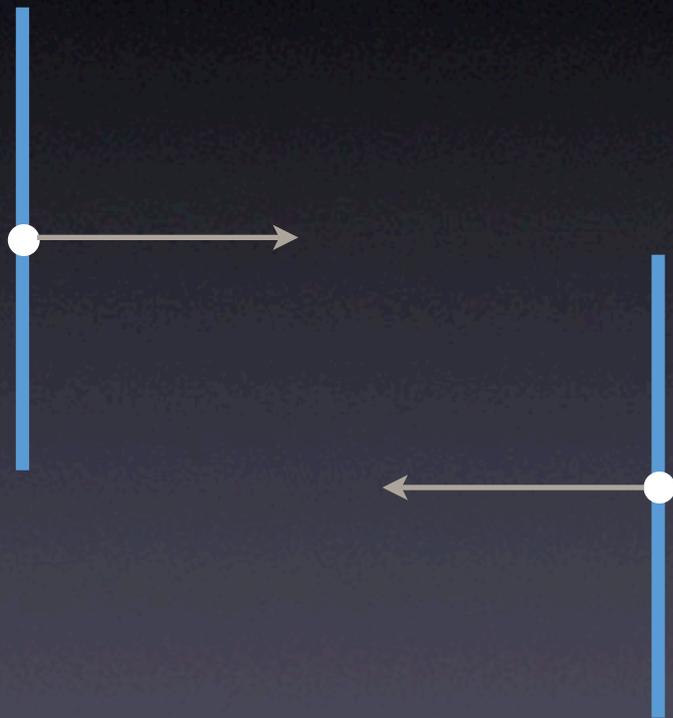
$\Leftrightarrow$  Dominant  $N$ : 
$$N \sim \frac{G_D E^2}{b^{D-4}} ;$$

At  $t \sim -1$  : 
$$N \sim (G_D E^2)^{\frac{1}{D-3}}$$

$\therefore$  Large loop order dominates.

Eikonal  $\longleftrightarrow$  classical scattering

Two Aichelburg-Sexl shocks (ACV: checks)



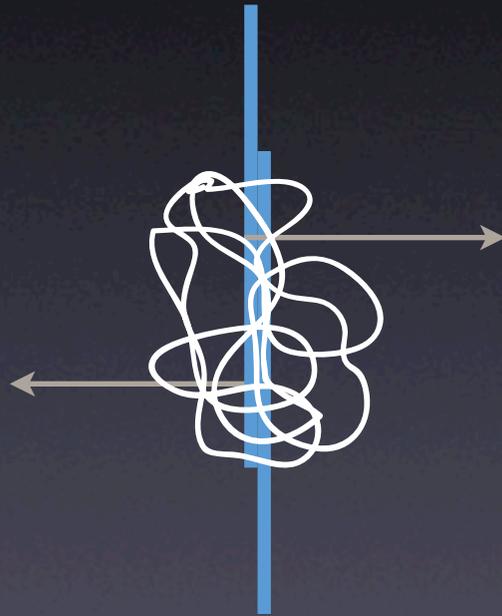
Black hole formation?

But - can excite strings: “diffractive excitation” (ACV)

Indeed, unexcited (elastic) amplitude, near Schwarzschild radius:

$$A_{el} \sim \exp \left\{ -E^{(D-4)/(D-3)} \right\} \quad !!$$

So:



??

No black hole??

Info carried away?

(Veneziano, 2004)

But there is a contrary intuition: string only  
“spreads out” “after” collision??

However, string spreading is a notoriously  
fuzzy concept...

Where is the string?

Karliner, Klebanov, Susskind: it depends



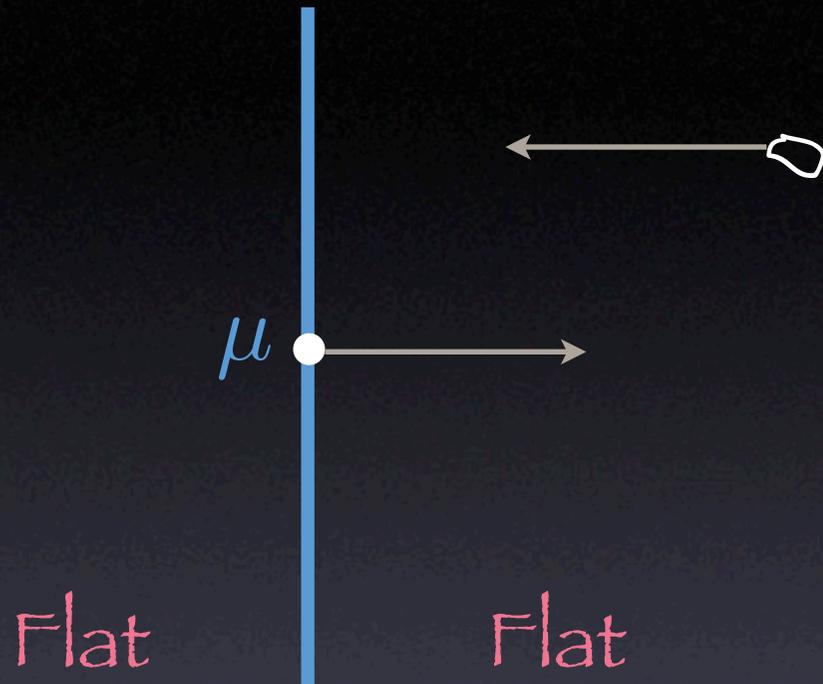
“low resolution”



“high resolution”

So: need to check for process in question ...

A test:



$$ds^2 = -dudv + dx^i dx^i + \Phi(\rho)\delta(u)du^2$$

$$\Phi(\rho) = -8G\mu \ln \rho \quad , \quad D = 4$$

$$\Phi(\rho) = \frac{16\pi G\mu}{\Omega_{D-3}(D-4)\rho^{D-4}} \quad , \quad D > 4$$

Scattering in a plane-wave metric:

de Vega and Sanchez; Horowitz and Steif

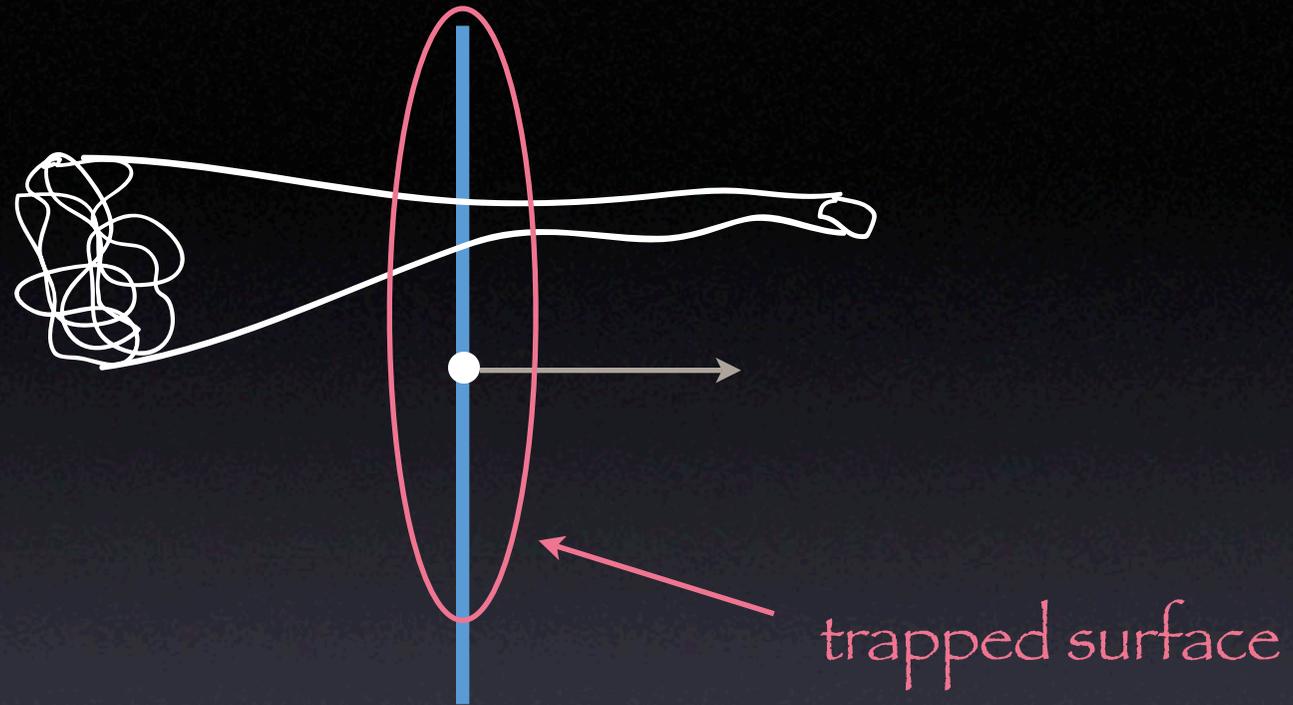
\*Light cone quantization\*

Compute for incoming unexcited string:

$$\langle \hat{X}_\epsilon^i(\tau, \sigma) \hat{X}_\epsilon^i(\tau, \sigma) \rangle$$

Where  $\hat{X}_\epsilon^i(\tau, \sigma)$  is deviation from CM of string,  
w/world sheet regulator  $\epsilon$

Find:



Indeed, origin of effect is “tidal string excitation”

$$(\Delta X)^2 \sim |\ln \epsilon| + \left[ \frac{G_D E^2}{b^{D-2}} \tau \right]^2 |\ln \tau| \quad \epsilon \ll \tau$$

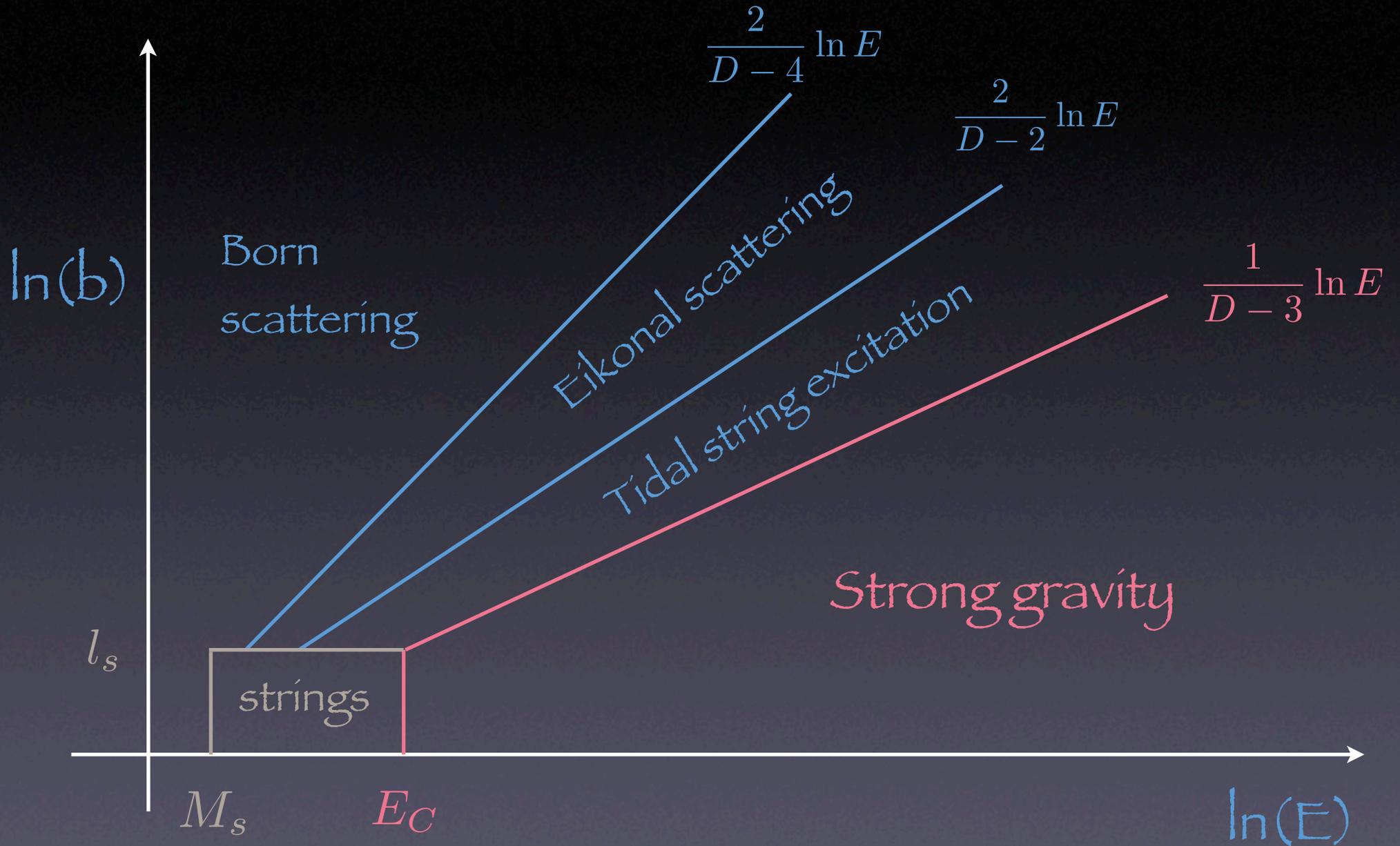
For small tau: inside trapped surface

Thus:

- String appears to behave ~locally during collision
- Trapped surface (aka black hole) appears to form

What conclusions can we draw?

# 1. A suggested "phase diagram:"



## 2. Perturbation theory apparently breaks down



$$1 + \mathcal{O} \left[ \left( R_S(E)/b \right)^{D-3} \right]$$

...not short distance

This is a challenge to string calculability ...  
(AdS/CFT)

3. Mechanism: no apparent fundamental role for string extendedness

A. Assuming scattering is finite and unitary:

unitarization apparently via intrinsically nonperturbative gravitational effects?

B. This dynamics is probably not local.

A possible “principle:” the nonperturbative physics that unitarizes gravity in domains where gravitational perturbation theory fails is nonlocal

4. Suggested correspondence boundary:

where does GR+LQFT break down?

2 part Fock sp.:

$$\phi_{x,p} \phi_{y,q} |0\rangle$$

(min uncertainty wavepackets)

good description for  $|x - y|^{D-3} > G|p + q|$

where  $G \sim G_{Newton}$

“the locality bound”

(extends off shell?)

How does this relate to BH info?

Criteria/evidence for locality vs. nonlocality?

What can one say about cosmology?

What does AdS/CFT say?

...

# Criteria for locality, and breakdown

1. Derivable from local QFT

2.  $[\mathcal{O}_1(x), \mathcal{O}_2(y)] = 0$  ,  $(x - y)^2 > 0$

3. Bounds/analyticity: Froissart; Cerulus-Martin; polynomial boundedness ...

$$2. [\mathcal{O}_1(x), \mathcal{O}_2(y)] = 0, \quad (x - y)^2 > 0$$

- There are no local gauge invariant observables in gravity. (Diffeos!)
- One can construct gauge invariant “proto-local” observables that approximately reduce to local observables in certain states (SBG, Marolf, and Hartle hep-th/0512200; Gary and SBG, hep-th/0612191).
- However, in situations characterized by the locality bound (and generalizations), one encounters obstacles to such a reduction.
- Thus this criterion for locality apparently is only approximate and appears to break down in situations of interest.

3. Bounds/analyticity: Froissart; Cerulus-Martin;  
polynomial boundedness ...

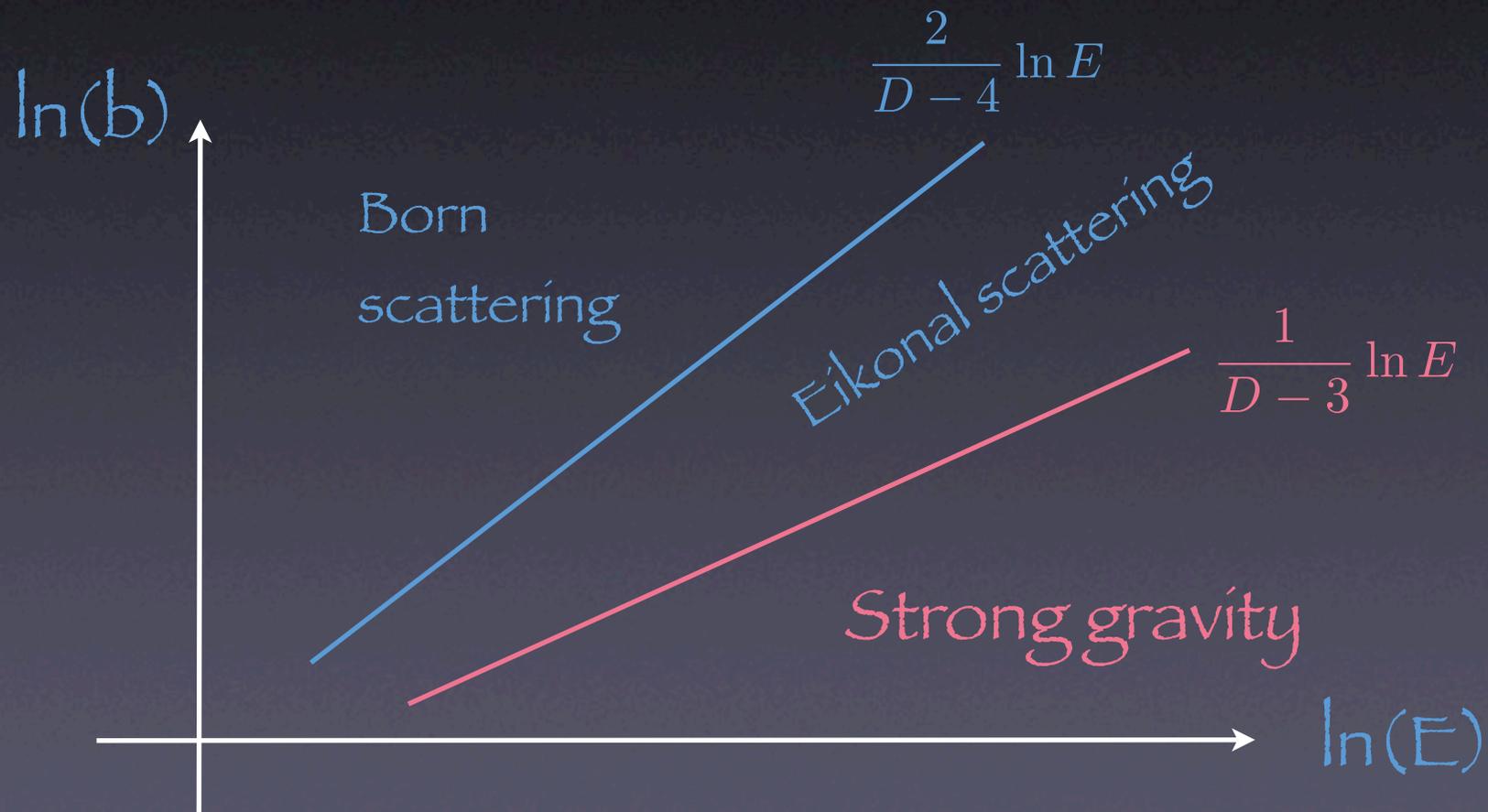
What general properties of gravitational  
amplitudes can we infer?

(WIP w/ M. Srednicki)

E.g. study partial wave expansion of 2-2 scattering  
(IR:  $D > 6$ )

$$T(s, t) = (\text{const}) E^{4-D} \sum_{l=0}^{\infty} (l + \nu) C_l^\nu(\cos \theta) \left[ e^{2i\delta_l(s) - 2\beta_l(s)} - 1 \right]$$

$$\nu = \frac{D-3}{2}$$



Some features:

A. Understand Born, eikonal regions

$$\text{e.g. } \delta_l \approx [ER_S(E)]^{D-3} / l^{D-4} \quad , \quad \beta_l \approx 0$$

B. Ansatz for BH region  $\beta_l \approx \frac{S(E, l)}{4}$

consequences (tentative)

- amplitudes apparently obey Cerulus Martin  
(contrary to earlier expectations)

- absorptive amplitude violates Froissart

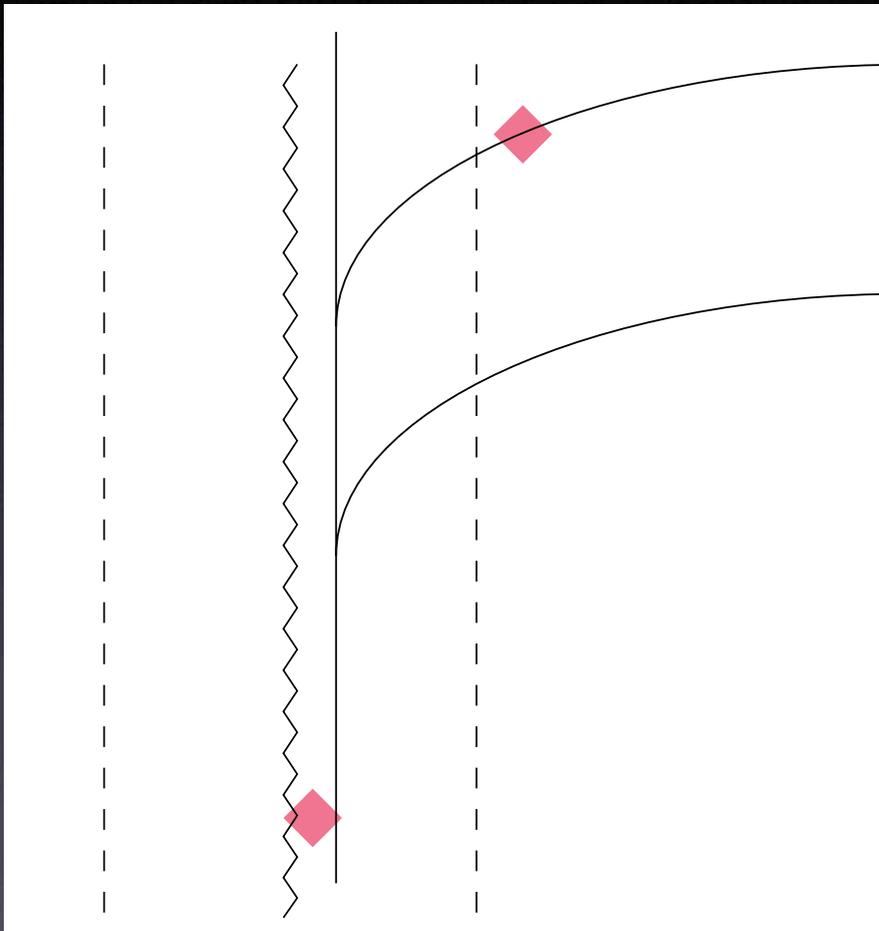
$$\sigma_{abs} \sim [R_S(E)]^{D-2}$$

- correspondingly, amplitudes not poly bdd:

$$T(s, t) \sim e^{R_S(E)\sqrt{t}}$$

(Not “local” by usual criterion)

What is missing in Hawking's argument for information loss?



proposal:

non-perturbative,  
non-local effects that  
become important by

$$T \sim R_S S_{BH}$$

(strong complementarity/  
strong holography not  
needed?)

## Some comments:

- there exists motivation for such effects, based on a) apparent breakdown of perturbation theory in nice slice quantization and b) limitations on observation of state. See [hep-th/0703116](https://arxiv.org/abs/hep-th/0703116)
- a complete picture of how/why such physics enters and how it relays the information would require knowledge of this nonpert. dynamics ...

What does AdS/CFT (or matrix thy) say?

Commonly believed that one has complete nonpert descriptions of string theory

- AdS/CFT: need to extract flat space limit.

Subtle. Concrete First test: can we see  $1/q^2$  of Born regime? (WIP w/ M. Gary)

- Matrix thy: does it avoid divergences of grav. pert thy?

## Cosmology: de Sitter, etc.

- if complementarity inessential in BHs, there should also be a global picture for dS.
- steps towards the formulation of such a picture:  
SBG and Marolf, arXiv:0705.1178 and WIP

- some features: finite number of pert. dS states  
not violating loc. bd.; no recurrences; relational  
observation

- apparent limitations to local QFT description of  
global picture by timescale  $R_{dS} S_{dS}$ : Boltzmann brain  
observers; large perturbative corrections  $\sim$  BH case ...  
(longer times likely allowed in static patch picture)

- related constraints found in Arkani-Hamed et al  
[arXiv:0704.1814] picture of regulating dS: large  
fluctuations at time  $\sim R_{dS} S_{dS}$

## Summary

- several considerations (HE scattering; observables; BH information, ...) support breakdown of conventional locality
- mechanism: no apparent role for string extendedness; rather nonperturbative gravity
- not clear how string theory addresses these issues
- correspondence boundary for such a “nonlocal mechanics:” locality bound, etc.

- such nonlocality could explain how info escapes black holes (“unitarity restored at the price of locality”)

- related story for inflationary cosmology; potentially places limitations on regime of local QFT description

(might we expect corresponding limitations on eternally inflating landscape?)