

**KITP** 

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# Spin Excitations in Stoichiometric Yb<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub>

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Yb<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub> 100 mK

## Outline

- Anisotropic exchange in rare-earth pyrochlores
- "Stoichiometric" vs. "Stuffed" Yb<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub>
- Elastic and inelastic neutron scattering on stoichiometric Yb<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub> powders
  - development of 50% long range ordered moment
  - Unconventional, gapless spin excitations
- Comparison to excitations in crystals and Yb<sub>2</sub>Sn<sub>2</sub>O<sub>7</sub>





### Anisotropic Exchange in Pyrochlores

#### $R_2 Ti_2 O_7 (R = Yb^{3+}, Er^{3+}, Ho^{3+}, Dy^{3+} etc.)$

- Rare-Earth Oxides: dominant spin orbit coupling
- Moments described by total angular momentum, J
- low temperatures: anisotropic g-tensors and exchange interactions
- Can get anything from long range AFM order to spin ice depending on details





### Anisotropic Exchange in Pyrochlores

- Pyrochlore: 4 Symmetryallowed exchange parameters (J<sub>zz</sub>, J<sub>±</sub>, J<sub>±±</sub>, J<sub>z±</sub>)
- Yb<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub>: XY-like g-tensor, but strong exchange coupling with Ising components



$$H = \sum_{\langle ij \rangle} \left\{ J_{zz} S_{i}^{z} S_{j}^{z} - J_{\pm} (S_{i}^{+} S_{j}^{-} + S_{i}^{-} S_{j}^{+}) + J_{++} \left[ \gamma_{ij} S_{i}^{+} S_{j}^{+} + \gamma_{ij}^{*} S_{i}^{-} S_{j}^{-} \right] \right.$$

$$+ J_{z\pm} \left[ S_{i}^{z} (\zeta_{ij} S_{j}^{+} + \zeta_{ij}^{*} S_{j}^{-}) + i \leftrightarrow j \right] \right\}$$

K.A. Ross, L. Savary, B. D. Gaulin, and L. Balents, Phys. Rev. X 1, 021002 (2011)

### Yb<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub> Specific Heat



Powder samples

usually have sharper, higher temperature anomalies

- Some samples show signs of ferromagnetic order below T<sub>c</sub>, others do not
  - Ordered moment size quoted from 0.8 to 1.1 µB (47 - 64% of total moment)
- Evidence for hysteresis at the sharper transitions: **first order**

### Some single crystals do not order



K. A. Ross et al, PRL **103** 227202 (2009)

• Diffuse continuum-like scattering at low temperatures and low fields

•No sharp magnetic Bragg scattering below Tc

### Evolution of $S(Q, \omega)$ in Single Crystals



Broad scattering develops into sharp magnons with increasing field (H || [1,-1,0]) K. A. Ross et al, PRL **103** 227202 (2009)

#### "Quantum Spin Ice" Exchange Parameters for Yb<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub>



K.A. Ross, L. Savary, B. D. Gaulin, and L. Balents, Phys. Rev. X 1, 021002 (2011)

### What is the ground state?

Exchange parameters for Yb<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub> compared to "Gauge Mean Field" phase diagram



 $L = (-0.5, -0.5, -0.5) \Gamma = (0, 0, 0)$ 

 $\Gamma = (2, 2, 2)$ 

 $\Gamma = (2, 2, 0)$ 

(1,1,0)

 $\Gamma = (1$ 

L. Savary, L. Balents, Phys. Rev. Lett. 108, 037202 (2012)

### What is the ground state?

Yb<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub> also lies "close" to a phase boundary between AFM and FM states



- New proposed parameters from other groups<sup>[1,2]</sup> suggest Yb<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub> is right on the edge of AFM order
- Do quantum fluctuations arise from proximity to AFM state?
- What role does the known sample dependence play?

[1] J. Robert, arXiv:1506.01729 [cond-mat.str-el] (2015)
[2] R. Coldea, KITP talk (<u>http://online.kitp.ucsb.edu/online/lsmatter15/coldea/</u>)

#### Powder vs.crystal: Log scale



#### Powder vs.crystal: Linear scale



K. A. Ross et al, PRB 84 174442 (2011)

#### Crushed Crystal vs. Sintered Powder

Structural difference between crushed crystal and sintered powder can be attributed to "stuffing" 2.3% excess Yb<sup>3+</sup> on Ti<sup>4+</sup> sublattice



### Powder Diffraction data at 15 K



## Stoichiometric powder

#### Crushed Crystal 2.3% stuffing

Ross et al, PRB **86**, 174424 (2012)

#### Crushed Crystal vs. Sintered Powder

# What are the ground state magnetic properties of the stoichiometric powder?



#### Neutron Scattering Measurements at the NCNR

10 grams of the **previously studied stoichiometric powder sample**, sealed with 10 atm helium gas at room temperature

SPINS Triple Axis Spectrometer (Elastic)





Disk Chopper Spectrometer (inelastic)



#### "Elastic" Scattering in a Stoichiometric Powder

- Elastic scattering,  $E = 0.0 \pm 0.25$  meV, at (111) Bragg Peak
- Increase of intensity visible in raw data



### Stoichiometric Powder of Yb<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub>

- (111) peak: 3.2 % increase from 8 K to 100 mK
- similarly, 3.3(5) % increase on (113) peak, 3.6(6)% increase on (222)



#### Fit to Stoichiometric Model at T = 8K3 × 10<sup>4</sup> Diffraction at 15 K Counts (3.49 × 10<sup>5</sup> monitor units) (Ross et al, PRB 2012) 2.5 (311)(222)2.5 2 S(Q) arp. nuits 1.5 0.5 2.05 2.1 2.15 2.2 Q(Å-1) background from \* sample environment 0.5 \* \* $\left( \right)$ (111)(002)(220)(113)(222)

![](_page_18_Figure_1.jpeg)

#### Resolution-Limited Elastic Scattering (100mK - 8K)

![](_page_19_Figure_1.jpeg)

### Temperature dependence of (111)

- gradual decline of intensity at (111) on warming (1 mK / min)
- Same temperature dependence as the diffuse scattering near (111) in our crystal
- No sharp changes at  $T_c$

![](_page_20_Figure_4.jpeg)

[1] K. Ross et al, PRB 84, 174442 (2011)

### Splayed Ferromagnetism

![](_page_21_Figure_1.jpeg)

- Stoichiometric powder shows splayed "ice" ferromagnetism
- **0.90(6)**  $\mu_B$  at 100 mK (compare to  $\mu_{sat} = 1.7 \ \mu$ B, 53% ordered)
- Large splaying angle of the moments, 14° (compare to spin ice, 54°)
- No sharp onset at transition

#### Temperature Dependence of $S(Q,\omega)$

![](_page_22_Figure_1.jpeg)

![](_page_23_Figure_1.jpeg)

![](_page_24_Figure_1.jpeg)

![](_page_25_Figure_1.jpeg)

![](_page_26_Figure_1.jpeg)

#### Temperature independence of spectrum

![](_page_27_Figure_1.jpeg)

- Continuum of scattering extending to ~1.5 meV
  - No change in S(Q,w)
     up to 2.5 K:
     insensitive to the
     transition seen in
     the specific heat

#### Gapless Excitations Not Expected

## Calculated zero-field spin waves

#### Measured Yb<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub> 100 mK

![](_page_28_Figure_3.jpeg)

Using Exchange parameters from Ross *et al*, Phys. Rev. X **1**, 021002 (2011)

Time-of-Flight Spectrometer (DCS)

#### Comparison to Crystals

![](_page_29_Figure_1.jpeg)

- Lineshapes are essentially the same as for crystals
- •Is the spectrum insensitive to details of the transition seen in heat capacity?
- •Why is the spectrum insensitive to the presence or absence of order?

![](_page_29_Figure_5.jpeg)

Powder

#### Inelastic Spectrum Compared to Yb<sub>2</sub>Sn<sub>2</sub>O<sub>7</sub>

![](_page_30_Figure_1.jpeg)

Thanks to C. Wiebe, H. Zhou, Z. Dun

### Compare $Yb_2Sn_2O_7$ and $Yb_2Ti_2O_7$

#### Yb<sub>2</sub>Sn<sub>2</sub>O<sub>7</sub> 100 mK 20K background

#### Yb<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub> 100 mK 8K background

![](_page_31_Figure_3.jpeg)

See: Dun et al, PRB **87**, 134408 (2013)

Thanks to C. Weibe, H. Zhou, Z. Dun

## Summary

- Yb<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub>'s true ground state has seemingly been obscured by sample dependence issues
- We showed elastic and inelastic neutron scattering studies of a powder sample known to be stoichiometric
- Partially ordered moment of ~0.9 μ<sub>B</sub> (53 % ordered), consistent with splayed "Ice" ferromagnetism, Temp dependence does not correlate with anomaly in specific heat
- **Gapless, continuum-like** spectrum at 100 mK is not conventional magnons, is insensitive to details of transition, and strongly resembles Yb<sub>2</sub>Sn<sub>2</sub>O<sub>7</sub>

![](_page_32_Figure_5.jpeg)

#### Thanks to...

- Bruce Gaulin
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#### Resolution-Limited Elastic Scattering (100mK - 8K)

![](_page_34_Figure_1.jpeg)

#### Collinear model

![](_page_35_Figure_1.jpeg)

### AFM (psi2/psi3) model

![](_page_36_Figure_1.jpeg)

#### Ordered Ice

![](_page_37_Figure_1.jpeg)