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CW invariants of  $\mathbb{C}P^1$

$\bar{\mathcal{M}}_{g,n,d}$  stable maps  $f: (C_g, z_1, \dots, z_n) \rightarrow \mathbb{C}P^1$   
of degree  $d$

$$\Psi_i = C_i(d_i) \quad (d_i)_f: (C, z_1, \dots, z_n) \rightarrow \mathbb{C}P^1 = T_{z_i}^* \mathbb{C}P^1$$

$$ev_i: (f: (C, z_1, \dots, z_n) \rightarrow \mathbb{C}P^1) = f(z_i)$$

$$\langle \tau_{k_1, a_1} \dots \tau_{k_n, a_n} \rangle_{g,n,d} = \int_{\bar{\mathcal{M}}_{g,n,d}} \Psi_1^{k_1} \dots \Psi_n^{k_n} ev_1^*(\gamma_{a_1}) \dots ev_n^*(\gamma_{a_n})$$

$$\gamma_p = 1 \in H^0(\mathbb{C}P^1)$$

$$\gamma_a = \omega \in H^2(\mathbb{C}P^1)$$

$$t_{n,p} = s_n \quad \left| \quad \partial = \partial / \partial s_0 \right.$$

$$t_{n,a} = t_n \quad \left| \quad \partial_a = \partial / \partial t_0 \right.$$

$$F_g = \sum_{d=0}^{\infty} q^d \sum_{n=0}^{\infty} \frac{1}{n!} \sum_{\substack{k_1, \dots, k_n \\ a_1, \dots, a_n}} t_{k_1, a_1} \dots t_{k_n, a_n} \langle \tau_{k_1, a_1} \dots \tau_{k_n, a_n} \rangle_{g,n,d}$$

$$F = \sum_{g=0}^{\infty} \epsilon^{2g} F_g$$

Large phase space - coordinates  $\{s_u, t_u\}$

Problem: calculate  $F$

Genus 0: Dijkgraaf - Witten

uses topological recursion relations and string equations

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$$\begin{aligned} \sigma &= \partial \partial_q F_0 \\ &= S_0 + \dots \end{aligned}$$

$$\begin{aligned} u &= \partial^2 F_0 \\ &= t_0 + \dots \end{aligned}$$

$$\partial^n \sigma = \begin{cases} 1 + S_1 + \dots & , n=1 \\ S_n + \dots & , n>1 \end{cases} \quad \partial^n u = t_n + \dots, n>0$$

Thus,  $\{\partial^n \sigma, \partial^n u\}$  give coordinates on large phase space.

Toda equation  $\partial_q^2 F_0 = q e^u$

Let  $A = \mathbb{Q}[e^u, \partial^n u, \partial^n \sigma]$

Poisson bracket on  $A[p, p^{-1}]$ :

$$\{f, g\} = p \frac{\partial f}{\partial p} \cdot \partial g - \partial f \cdot p \frac{\partial g}{\partial p}$$

Projection to Lie subalgebra:  $+ : A[p, p^{-1}] \rightarrow A[p]$

$$L = p + \sigma + e^u p^{-1}$$

$$\delta_n L = \{L_+, L\}$$

$$[\partial, \delta_n] = 0$$

$$\delta_n \sigma + e^u \delta_n u \cdot p^{-1}$$

Trace  $[\delta_n, \delta_m] = 0$

Toda hierarchy:  $\delta_n = \frac{1}{2} n! \partial_{n-1, q}$

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$$h_+^R = p + q\sigma$$

$$\{h_+, h\} = \{p + \sigma, p + \sigma + e^u p^-\}$$

$$= \{p + \sigma, e^u p^-\}$$

$$= \partial(e^u) + e^u \partial\sigma p^-$$

$$\delta_1 \sigma = \partial(e^u)$$

$$\delta_1 u = \partial\sigma$$

$$\partial^2(\delta_1 F) = \delta_1^2 u = \partial^2 e^u$$

$\Rightarrow$  Toda equation

Higher genus generalization - Eguchi, Yang  
 (recently proved by Okounkov, Pandharipande -  
 they reduce it to a question about Hodge numbers,  
 i.e. "2D QCD")

$$\nabla = \epsilon^{-1} (e^{\epsilon\partial/2} - e^{-\epsilon\partial/2})$$

$$\sigma = \nabla \partial_a F$$

$$u = \nabla^2 F$$

Algebra  $A[\lambda, \lambda^{-1}]$

$$a \lambda^k \cdot b \lambda^l = ab \lambda^{k+l}$$

$$(e^{-\epsilon l \partial/2} a) (e^{\epsilon k \partial/2} b) \lambda^{k+l}$$

$\dagger : A[\lambda, \lambda^{-1}] \rightarrow A[\lambda]$  as before

$$L = \Lambda + \sigma + e^u \Lambda^{-1}$$

$$\delta_u L = [\varepsilon^{-1} L_+^n, L], \quad [\partial, \delta_u] = 0 \quad \left| \begin{array}{l} \delta_u \sigma = \nabla e^u \\ \delta_u u = \nabla \sigma \end{array} \right.$$

$$\delta_u \sigma + e^u \delta_u u \cdot \Lambda^{-1}$$

$$\nabla^2 (\delta_u^2 F) = \delta_u^2 u = \nabla^2 e^u$$

Toda equation  $\partial_Q^2 F = e^u$

Toda conjecture  $\delta_u = n! \partial_{n-1, Q}$

Another form of the conjecture (G. - soon to appear, KIAS proceedings, Seoul, August 2000)

Let  $D = \sigma \nabla + \ln(e^{\varepsilon \partial/2} + e^{-\varepsilon \partial/2}) \partial_Q$ .

$D \partial_{n-1, Q} F = (n+1) \nabla \partial_{n, Q} F$

This recursion is very powerful.

Example:  $\lim_{g \rightarrow 0} \partial_{n, Q} F =$

$$\sum_{g=0}^{\infty} (-1)^g \varepsilon^{2g} \sum_{n=0}^{\infty} \frac{1}{n!} \sum_{k_1, \dots, k_n} S_{k_1} \dots S_{k_n} \int_{\mathcal{M}_{g, n+1}} \psi_1^{k_1} \dots \psi_n^{k_n} \psi_{n+1}^k \lambda_g$$

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Toda equation is trivial:  $\partial_a^2 F = 0$

$$\begin{aligned} \textcircled{1} \partial_{n-1, a} F &= \nabla \partial_a F \cdot \nabla \partial_{n-1, a} F \\ &= (n+1) \nabla \partial_{n, a} F \end{aligned}$$

$$\nabla \partial_{n, a} F = \frac{(\nabla \partial_a F)^{n+1}}{(n+1)!}$$

Evaluate on line  $\{S_n = 0, n > 0\}$

$$\sum_{g=0}^{\infty} (-1)^g \varepsilon^{2g} \int_{\bar{M}_{g,1}} \psi^{2g-2} \lambda_g = \frac{\varepsilon/2}{\sinh(\varepsilon/2)}$$

(result of Faber, Pandharipande)

What about the flows  $\partial_{n,p}$  ?

Two approaches:

- 1) Virasoro (Eguchi, Hori, Yang; G)
- 2) Equivariant Toda (G, Pandharipande)

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$$\begin{aligned}
 x_n &= \mathbb{D} \partial_{n,p} F - (n+1) \nabla \partial_{n+1,p} F - 2 \nabla \partial_{n,q} F \\
 y_n &= \mathbb{D} \partial_{n,q} F - (n+2) \nabla \partial_{n+1,q} F
 \end{aligned}$$

Let  $Z_n = n^{\text{th}}$  Virasoro constraint

$$= e^{-\epsilon^2 F} \text{Lu} e^{\epsilon^2 F}$$

Formula  $\hat{S}_m = \begin{cases} S_m, & m \neq 1 \\ S_{-1}, & m = 1 \end{cases}$

$$\begin{aligned}
 \mathbb{D} Z_{n+m} - \nabla Z_{n+1} &= \sum_{m=0}^{\infty} \left( (m+1) \dots (m+k+1) t_m y_{m+k} + m(m+1) \dots (m+k) \hat{S}_m x_{m+k} \right. \\
 &\quad \left. + 2m(m+1) \dots (m+k) \sum_{i=0}^k \frac{1}{m+i} \tilde{S}_m y_{m+k-i} \right) \\
 &\quad + \sum_{i+j=k} i! j! \left( \epsilon^2 \partial_{i-1,q} y_{j-1} + \left( e^{\epsilon \partial/\partial t} + e^{-\epsilon \partial/\partial t} \right) \partial_{i-1,q} F \right) y_{j-1}
 \end{aligned}$$

Corollary Toda + Virasoro  $\Rightarrow$  recursive

$$\mathbb{D} \partial_{n-1,p} F = n \nabla \partial_{n,p} F + 2 \nabla \partial_{n-1,q} F$$

Taking  $\lim_{g \rightarrow 0}$ , we get a nice formula for

$$\int_{\bar{\mathcal{M}}_{g,n}} \psi_{k_1}^{k_1} \dots \psi_n^{k_n} \lambda_{g-1}$$

In particular (Faber, Pandharipande)  $\sum (-1)^g \epsilon^{2g} \int_{\bar{\mathcal{M}}_{g,1}} \psi_2^{2g-1} \lambda_{g-1} = \log \left( \frac{\epsilon/2}{\sinh(\epsilon/2)} \right)$

Equivalent Toda

Genus 0 :  $\delta_1 \nearrow \partial_q (\partial_q - \lambda \partial_p) F = e^u$   
 $\delta_1 \nearrow \delta_1$

Toda Field Theory

$$L = p + \sum_{n=0}^{\infty} u_n p^{-n}$$

$$\bar{L} = \sum_{n=1}^{\infty} v_n p^n$$

$$\delta_n L = \{ L_+^n, L \}$$

$$\delta_n \bar{L} = \{ \bar{L}_+^n, \bar{L} \}$$

$$[\partial, \delta_n] = 0$$

Commuting flows ( $\approx$  2 component KP)

$$\delta_n L = - \{ \bar{L}_-^{-n}, L \}$$

$$\delta_n \bar{L} = - \{ L_-^{-n}, \bar{L} \}$$

$$[\partial, \delta_n] = 0$$

But  $\{ L_+ - \bar{L}_-^{-1}, L \} = \lambda \partial L$

$$\{ L_+ - \bar{L}_-^{-1}, \bar{L} \} = \lambda \partial \bar{L}$$

Hypothesis :  $u_0 = v$        $v_1 = e^{-u}$

$$\Leftrightarrow L_+ - \bar{L}_-^{-1} = p + v + e^u p^{-1}$$

Idea : Basis of  $H_{S'}^*(\mathbb{C}P^1)$  associated to two fixed points is better than homogeneous basis of  $H^*(\mathbb{C}P^1)$

Promote to higher genus:

$$\text{Let } L(0) = \lambda + \sigma + e^u \lambda^{-1}$$

$$L(\lambda) = L(0) + O(\lambda)$$

$$\bar{L}(\lambda) = L(0)^{-1} + O(\lambda)$$

$$\begin{cases} [\bar{\epsilon}^{-1} L(0), L(\lambda)] = \lambda \partial L(\lambda) \\ [\bar{\epsilon}^{-1} L(0), \bar{L}(\lambda)] = \lambda \partial \bar{L}(\lambda) \end{cases}$$

Unique solution - determines  $\partial_{n,q}, \partial_{n,p}$  by

$$\delta_n = n! \partial_{n-1,q}$$

$$\bar{\delta}_n = n! (\partial_{n-1,q} - \lambda \partial_{n-1,p})$$

Puzzle: recover Virasoro in limit  $\lambda \rightarrow 0$ .

Very powerful - determines

$$\sum_{g=0}^{\infty} (-1)^g \epsilon^{2g} \int_{\bar{\mathcal{M}}_{g,1}} \frac{1 + \lambda_1 + \lambda_2 + \dots + \lambda_g}{1 - t \Psi}$$

$$= t^{-2} \frac{\epsilon/2}{\sinh(\epsilon/2)} + t^{-1} \log \left( \frac{\epsilon/2}{\sinh(\epsilon/2)} \right) + \dots$$

Matrix model?