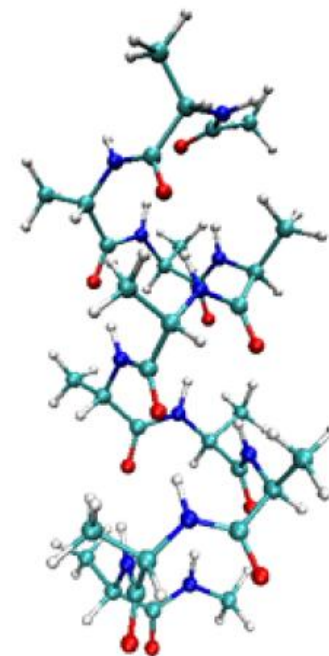
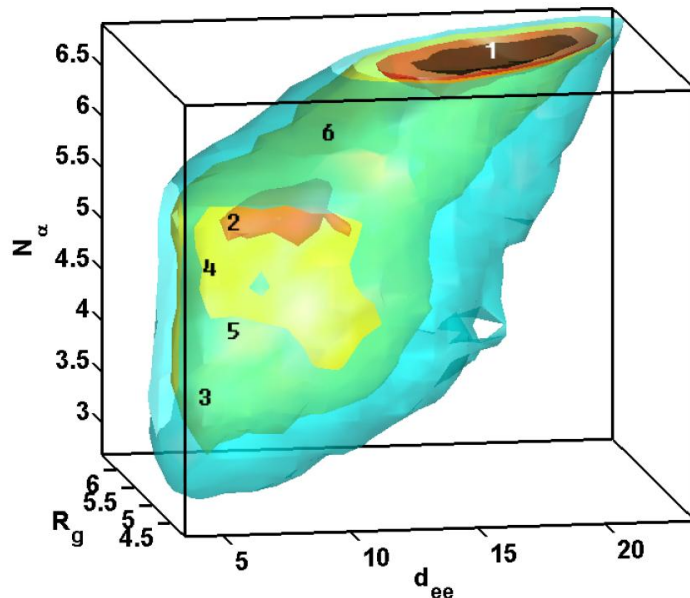
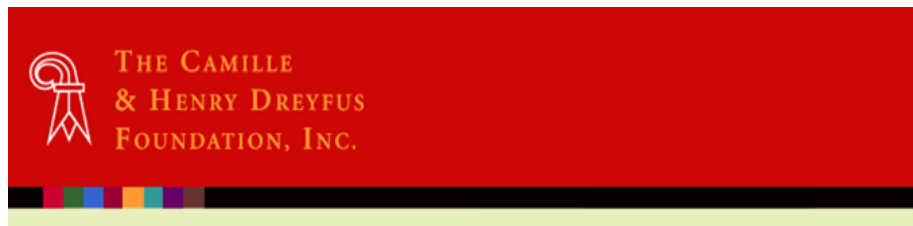


Exploring the free energy landscape of complex molecules and crystalline polymorphs



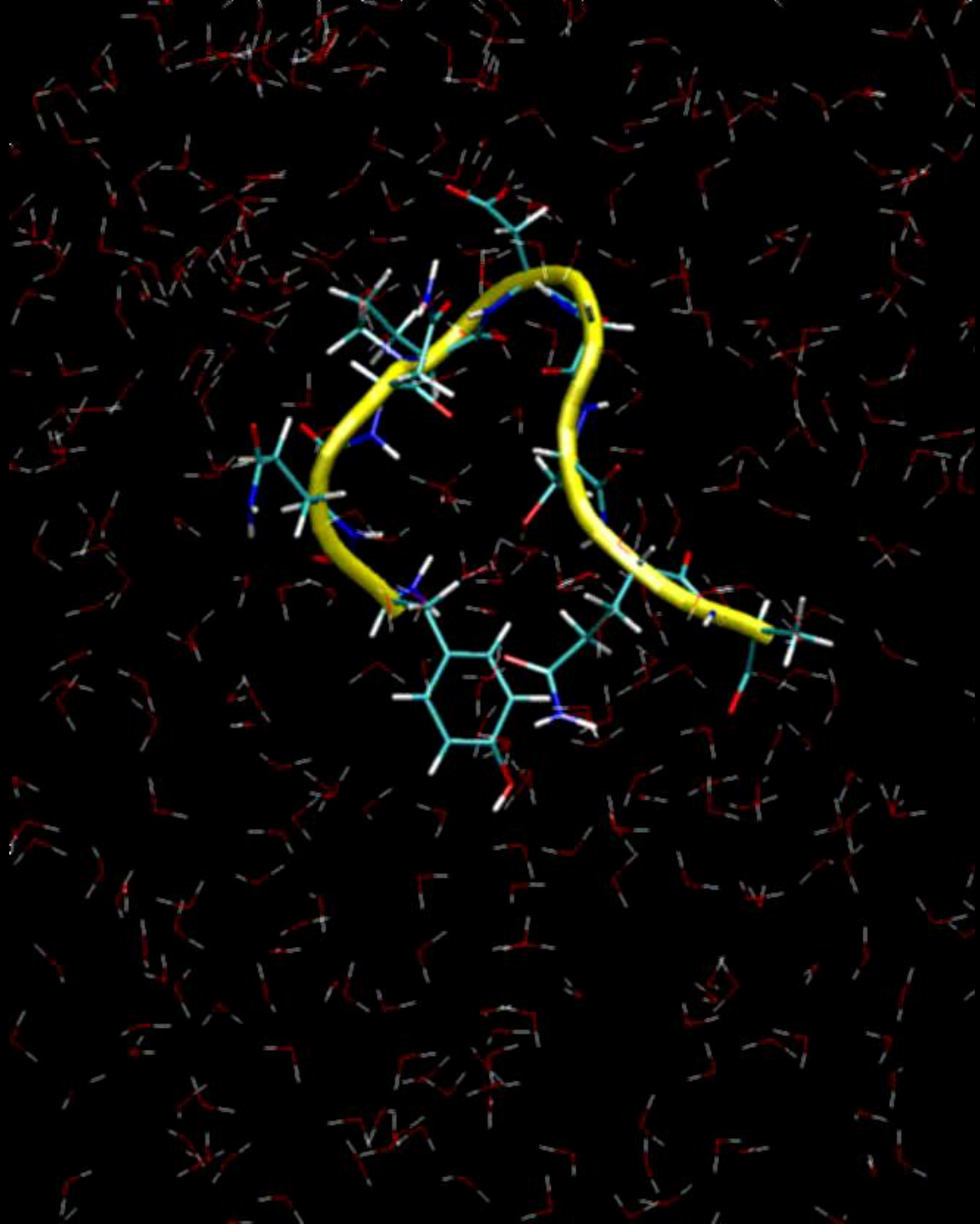
Mark E. Tuckerman

*Dept. of Chemistry and Courant Institute of Mathematical Sciences
New York University, 100 Washington Square East, NY 10003*



Sequence:
YQPDGSQA

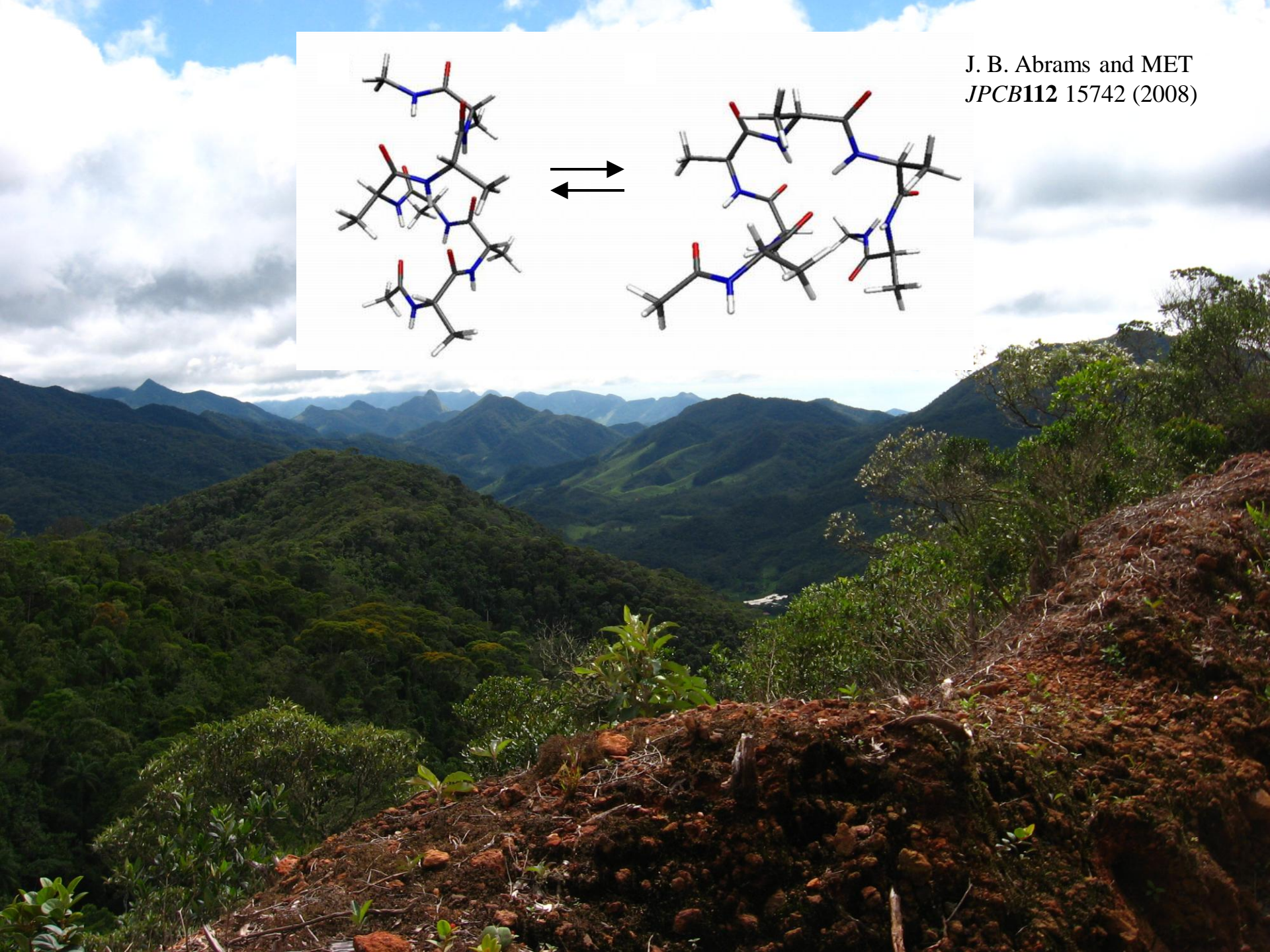
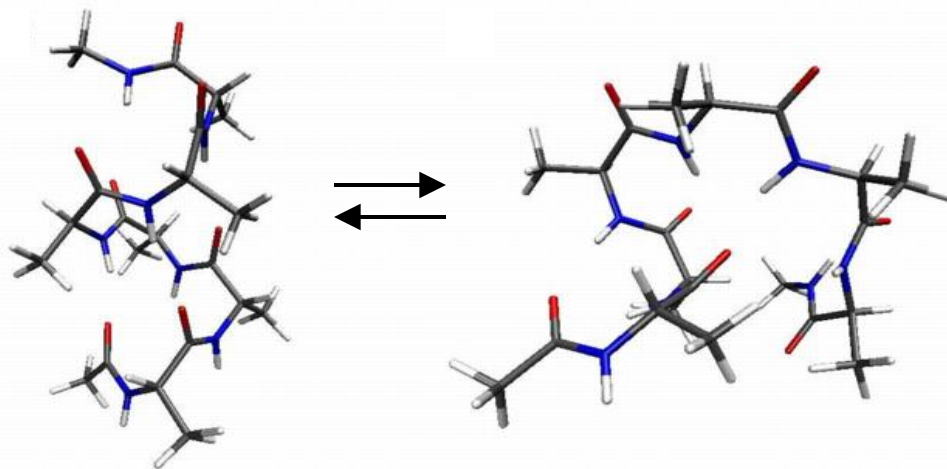
I.-C. Lin and MET
J. Phys. Chem. B
(2010)



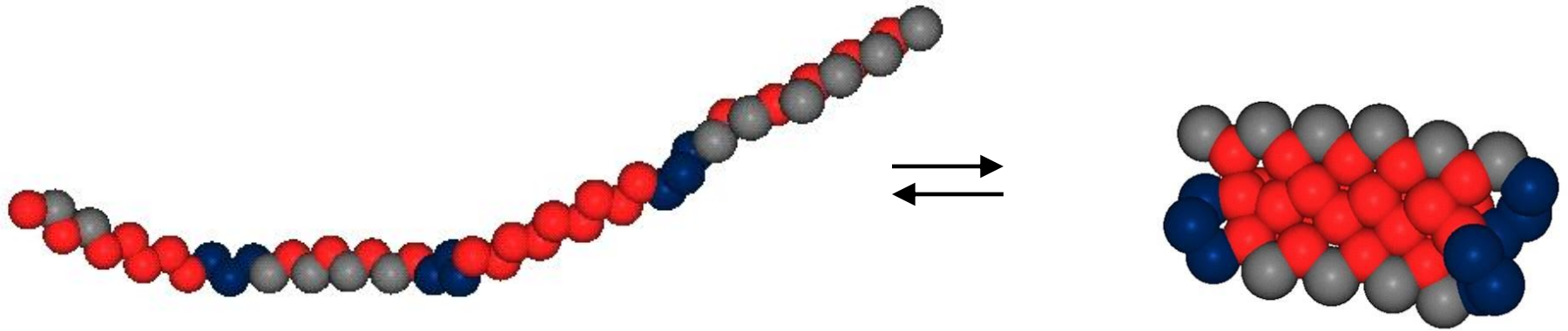
Solvent
mass scaling
0.1

Side-chain
mass scaling
0.6

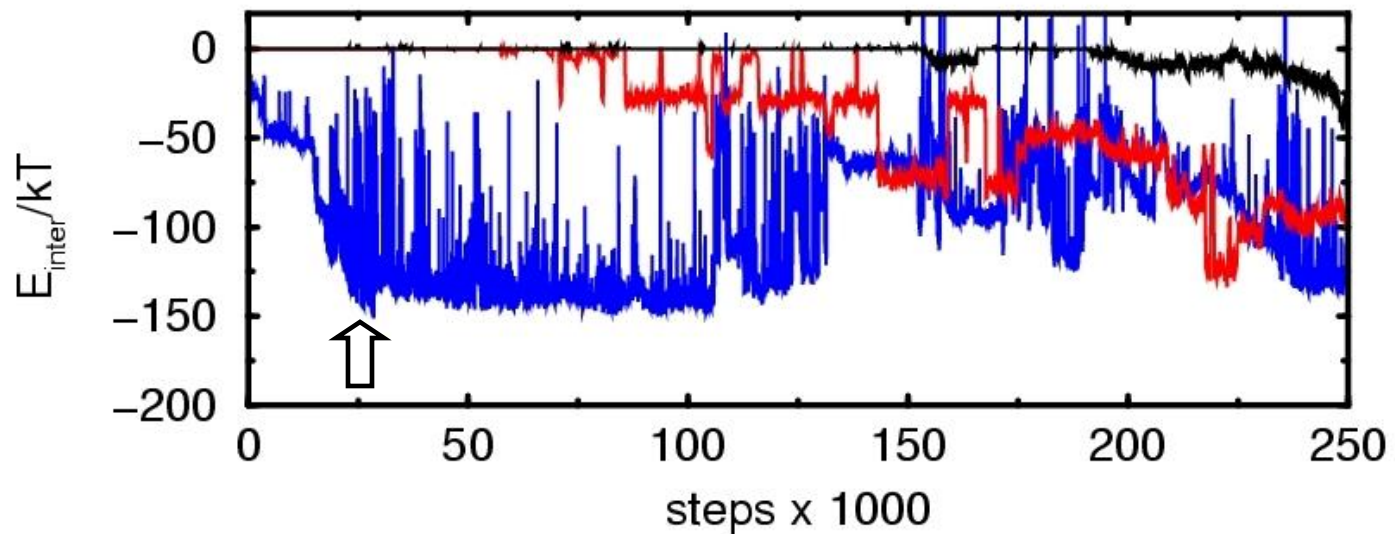
J. B. Abrams and MET
*JPCB*112 15742 (2008)



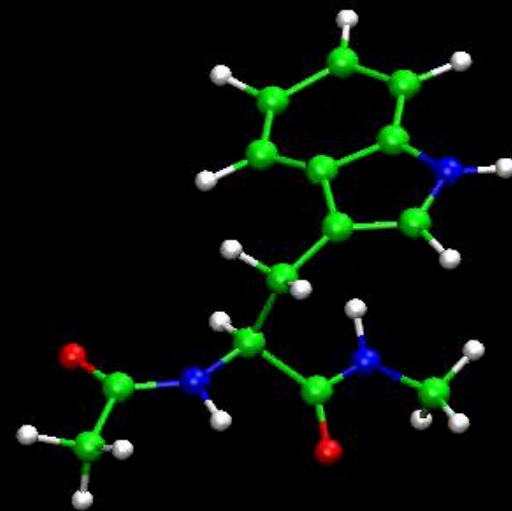
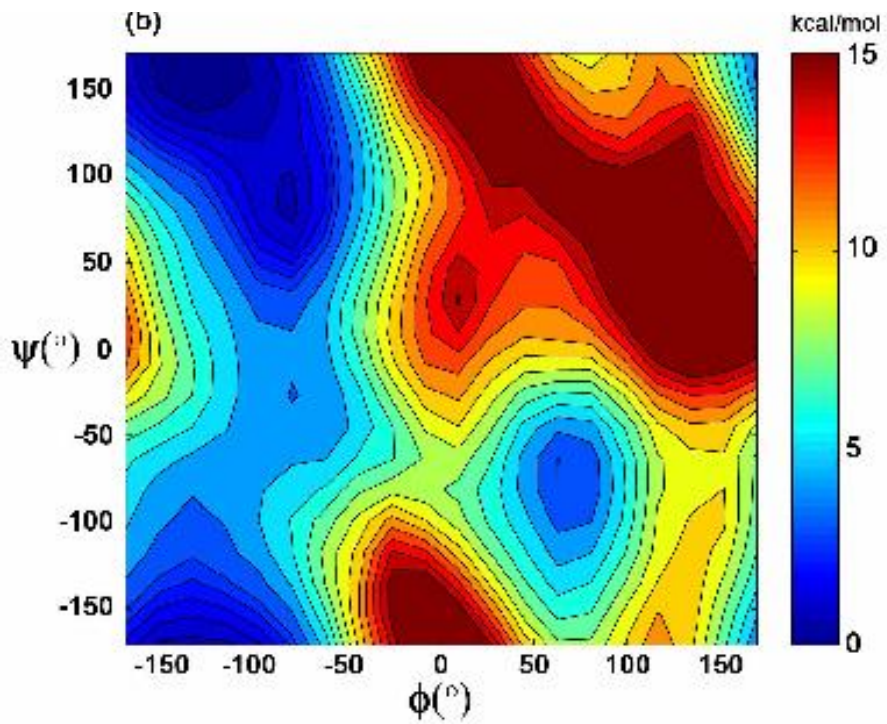
Conformational equilibrium



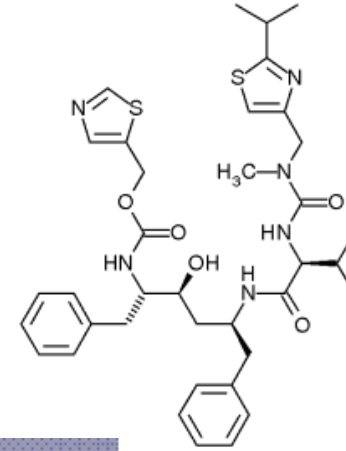
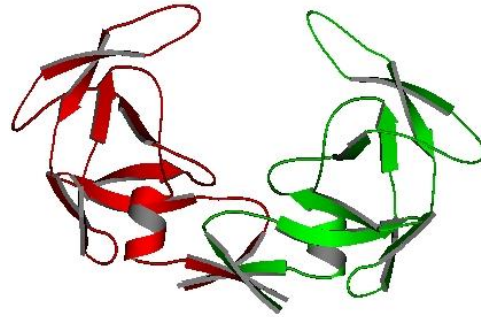
P. Minary, MET, G. J. Martyna *SIAM J. Sci. Comput.* **30**, 2055 (2008)



Ramachandran surface of NATMA

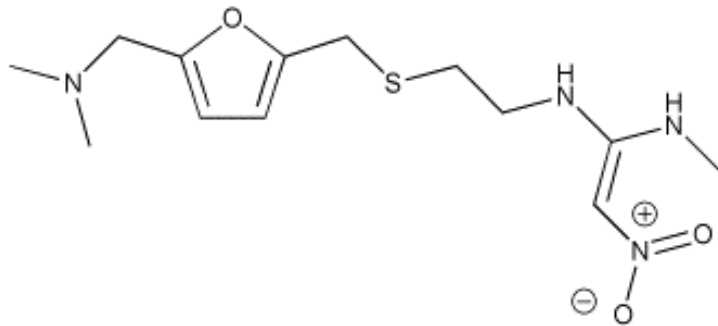
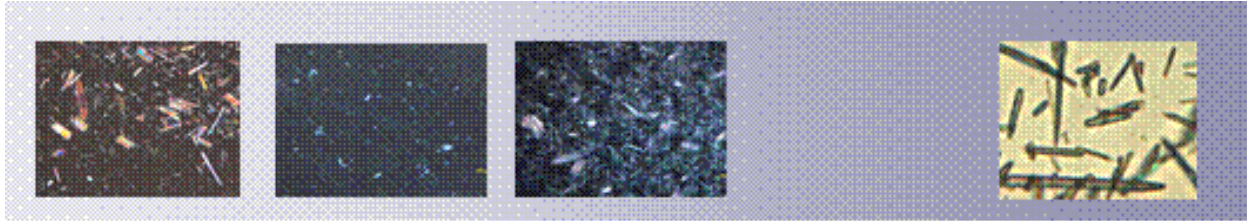


Polymorphic compounds

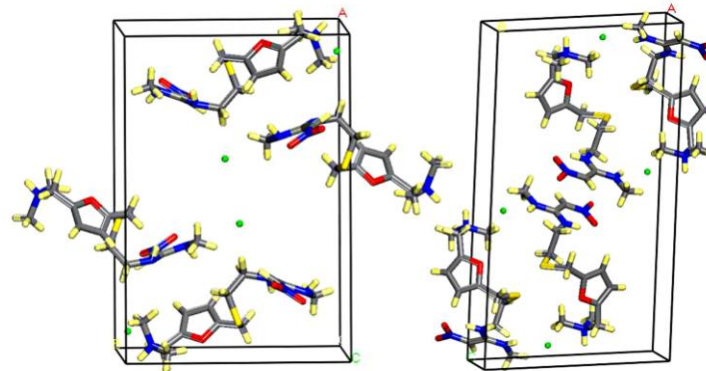


Norvir

1996 (launch) 1998 (recall) 2002 (Trans form)



Chemical Formula: $C_{13}H_{22}N_4O_3S$



Ranitidine
(Zantac)



Adiabatic free-energy dynamics

L. Rosso and MET *Mol. Simulat.* **28**, 91 (2002); L. Rosso, P. Minary, Z. Zhu, MET *J. Chem. Phys.* **116**, 4389 (2002)

In a transformation to generalized coordinates:

$$q_\alpha = q_\alpha(\mathbf{r}_1, \dots, \mathbf{r}_N) \quad \alpha = 1, \dots, 3N \quad ; \quad \mathbf{r}_i = \mathbf{r}_i(q_1, \dots, q_{3N})$$

Suppose first n are of particular interest.

$$\begin{aligned} P(s_1, \dots, s_n) &= \int d^N \mathbf{p} d^N \mathbf{r} e^{-\beta H(\mathbf{p}, \mathbf{r})} \prod_{\alpha=1}^n \delta(q_\alpha(\mathbf{r}) - s_\alpha) \\ &= \int d^N \mathbf{p} d^{3N} q e^{-\beta \tilde{H}(\mathbf{p}, q)} \prod_{\alpha=1}^n \delta(q_\alpha - s_\alpha) \end{aligned}$$

$$H = \frac{1}{2} \mathbf{p}^T M^{-1} \mathbf{p} + U(\mathbf{r}) \quad \tilde{H} = \frac{1}{2} \mathbf{p}^T M^{-1} \mathbf{p} + U(\mathbf{r}(q)) - kT \ln J(q)$$

Adiabatic and temperature conditions: $m_{1, \dots, n} \square m_{n+1, \dots, 3N} \quad T_s \square T$

Free-energy surface:

$$A(s_1, \dots, s_n, T) = -kT_s \ln P_{\text{adb}}(s_1, \dots, s_n)$$

Driven Adiabatic Free Energy Dynamics (d-AFED)

L. Rosso, P. Minary, Z. Zhu, MET *J. Chem. Phys.* **116**, 4389 (2002)

Margliano and Vanden-Eijnden, *Chem. Phys. Lett.* **426**, 168 (2006)

J. B. Abrams and MET, *J. Phys. Chem. B* **112**, 14752 (2008)

Suppose n collective variables characterize a free energy landscape of interest

$$q_\alpha = q_\alpha(\mathbf{r}_1, \dots, \mathbf{r}_N) \quad \alpha = 1, \dots, n$$

Canonical probability distribution:

$$P(s_1, \dots, s_n) = \int d^N \mathbf{p} d^N \mathbf{r} e^{-\beta H(\mathbf{p}, \mathbf{r})} \prod_{\alpha=1}^n \delta(q_\alpha(\mathbf{r}) - s_\alpha)$$

Write δ -functions as product of Gaussians:

$$P(s_1, \dots, s_n) = \lim_{\{\kappa_\alpha \rightarrow \infty\}} \int d^N \mathbf{p} d^N \mathbf{r} e^{-\beta H(\mathbf{p}, \mathbf{r})} \prod_{\alpha=1}^n \left(\frac{\beta \kappa_\alpha}{2\pi} \right)^{1/2} \exp \left[-\frac{\beta \kappa_\alpha}{2} (q_\alpha(\mathbf{r}) - s_\alpha)^2 \right]$$

Introduce uncoupled Gaussian integrations:

$$P_{\{\kappa\}}(s_1, \dots, s_n) = C_{\{\kappa\}} \int d^N \mathbf{p} d^N \mathbf{r} d^n p_s \exp \left\{ -\beta \left[H(\mathbf{p}, \mathbf{r}) + \sum_{\alpha=1}^n \frac{p_{s_\alpha}^2}{2m_\alpha} + \sum_{\alpha=1}^n \frac{\kappa_\alpha}{2} (q_\alpha(\mathbf{r}) - s_\alpha)^2 \right] \right\}$$

Effective Hamiltonian:

$$\mathcal{H}(\mathbf{p}, \mathbf{r}, s, p_s) = H(\mathbf{p}, \mathbf{r}) + \sum_{\alpha=1}^n \frac{p_{s_\alpha}^2}{2m_\alpha} + \sum_{\alpha=1}^n \frac{\kappa_\alpha}{2} (q_\alpha(\mathbf{r}) - s_\alpha)^2$$

Driven Adiabatic Free Energy Dynamics (d-AFED)

Introduce high temperature $T_s \square T$ for extended variables and high masses $m_\alpha \square m_i$

Adiabatically decoupled equations of motion:

$$m_i \ddot{\mathbf{r}}_i = -\frac{\partial H}{\partial \mathbf{r}_i} + \sum_{\alpha} \kappa_{\alpha} (s_{\alpha} - q_{\alpha}(\mathbf{r})) \frac{\partial q_{\alpha}}{\partial \mathbf{r}_i} + \text{heat bath}(T)$$

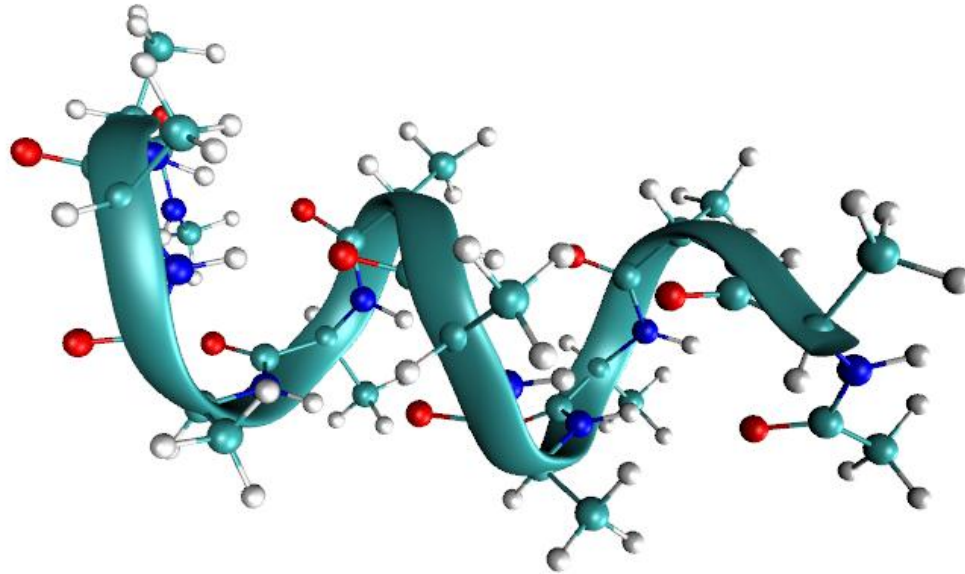
$$m_{\alpha} \ddot{s}_{\alpha} = -\kappa_{\alpha} (s_{\alpha} - q_{\alpha}(\mathbf{r})) + \text{heat bath}(T_s)$$

Under adiabatic conditions, we generate a distribution $P_{\text{adb}}^{(\{\kappa\})}(s_1, \dots, s_n, T_s, T)$

$$P_{\text{adb}}^{(\{\kappa\})}(s_1, \dots, s_n, T_s, T) = \tilde{C}_{\{\kappa\}} \int d^n p e^{-\beta_s \sum_{\alpha} \frac{p_{\alpha}^2}{2m_{\alpha}}} \left[P_{\{\kappa\}}(s_1, \dots, s_n, T) \right]^{T/T_s}$$

$$\lim_{\{\kappa \rightarrow \infty\}} \left[-kT_s \ln P_{\text{adb}}^{(\{\kappa\})}(s_1, \dots, s_n, T_s) \right] = A(s_1, \dots, s_n, T)$$

Alanine Decamer (gas phase)



Force field:

CHARMM22

20 CVs:

All (φ, ψ) pairs

CV Temperature:

$T_s = 900$ K

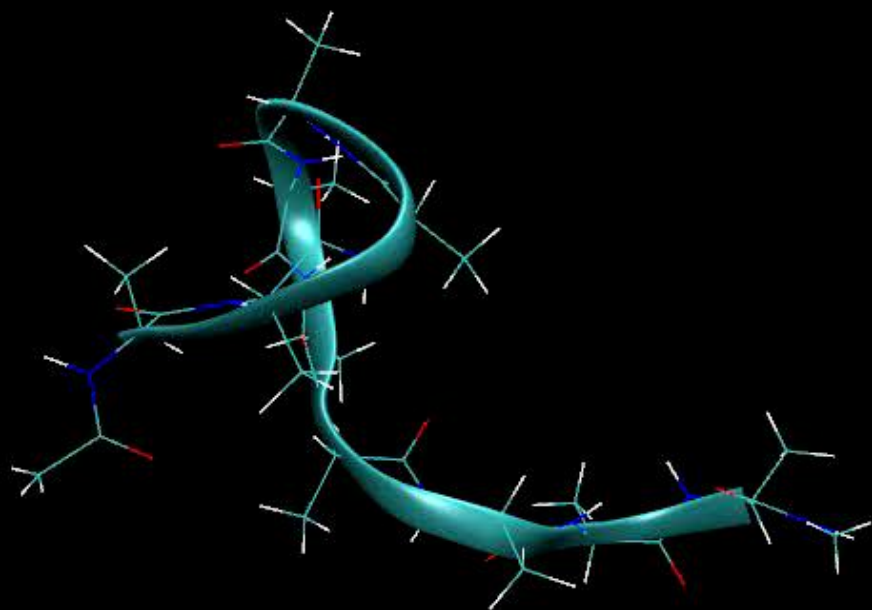
Physical Temp:

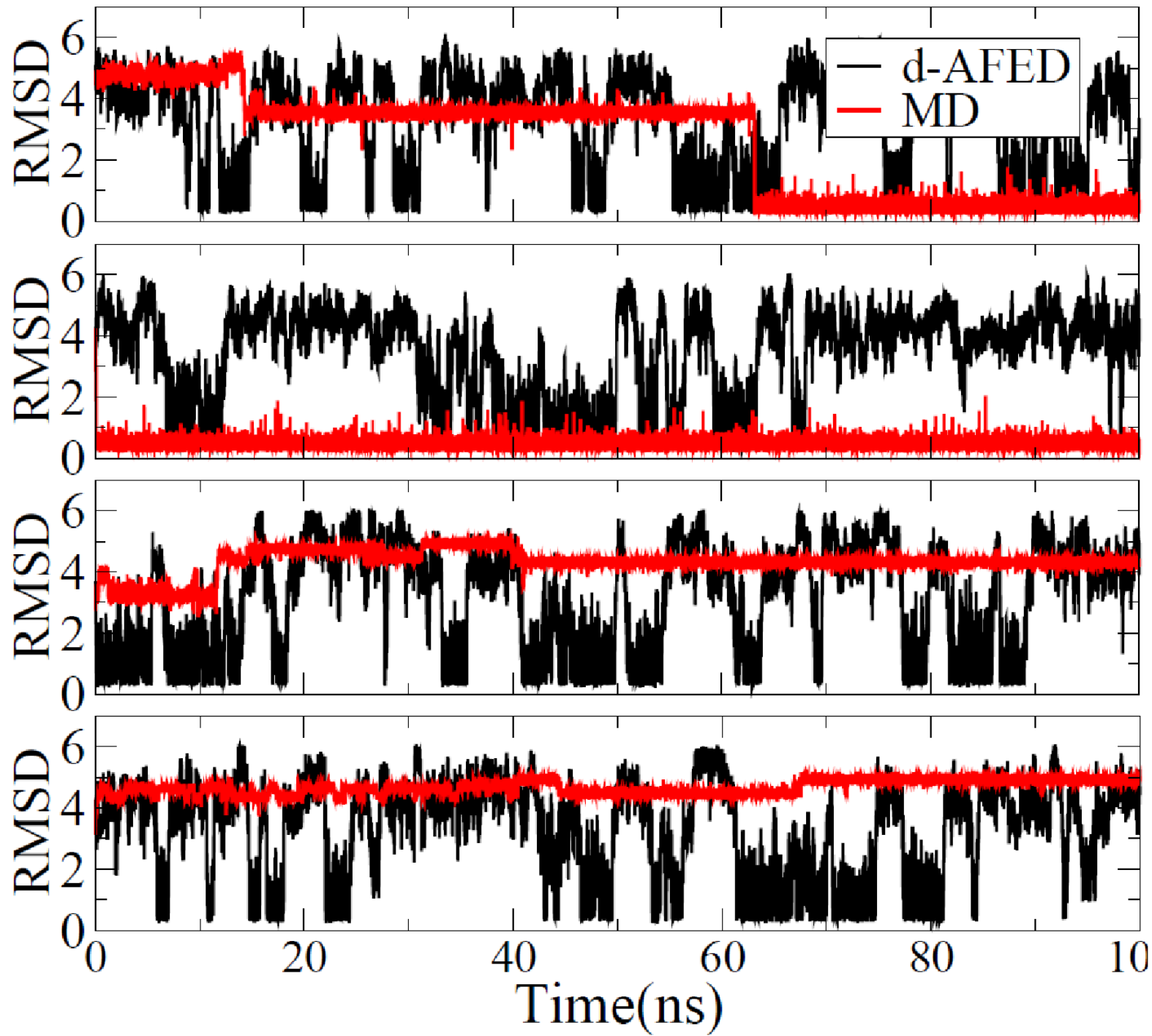
$T = 300$ K

CV mass:

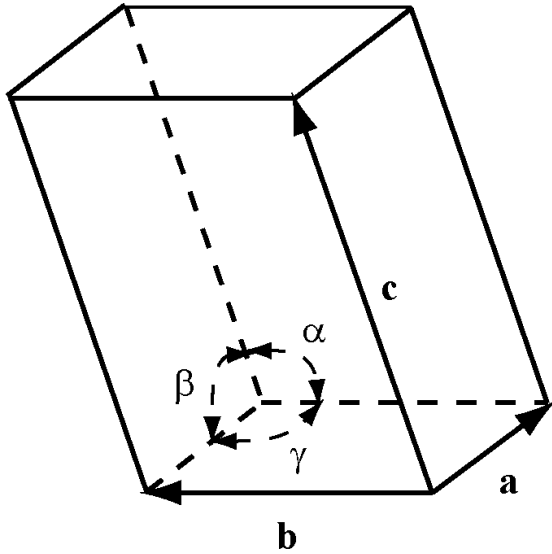
$m_{(\psi, \varphi)} = 600m_H$

Harmonic coupling: 300.5 kcal/mol/rad²





Crystal-AFED [T. Q. Yu and MET *PRL* (2011)]



$$\mathbf{h} = \begin{pmatrix} a_x & b_x & c_x \\ a_y & b_y & c_y \\ a_z & b_z & c_z \end{pmatrix}$$

$$V = \mathbf{a} \cdot (\mathbf{b} \times \mathbf{c}) = \det(\mathbf{h})$$

$$\dot{\mathbf{r}}_i = \frac{\mathbf{p}_i}{m_i} + \frac{\mathbf{p}_g}{W} \mathbf{r}_i, \quad \dot{\mathbf{p}}_i = \mathbf{F}_i - \frac{\mathbf{p}_g}{W} \mathbf{p}_i - \frac{1}{N_f} \frac{\text{Tr}[\mathbf{p}_g]}{W} \mathbf{p}_i + \text{heat bath}(T)$$

$$\dot{\mathbf{h}} = \frac{\mathbf{p}_g \mathbf{h}}{W}, \quad \dot{\mathbf{p}}_g = \det(\mathbf{h}) (\mathbf{P}^{(\text{int})} - P\mathbf{I}) + \frac{1}{N_f} \sum_{i=1}^N \frac{\mathbf{p}_i^2}{m_i} \mathbf{I} + \text{heat bath}(T_h)$$

$$\text{Large } W, \quad T_h \ll T$$

Predicting crystalline polymorphs using adiabatic dynamics

T. Q. Yu and MET *Phys Rev. Lett.* (2011)

Using the AFED framework, we run NPT dynamics using the **h**-matrix as collective variables.

Generates the Gibbs free energy surface from the adiabatic probability distribution:

$$G(\mathbf{h}, T) = -kT_{\mathbf{h}} \ln P_{\text{adb}}(\mathbf{h}, T, T_{\mathbf{h}})$$

Test case: Benzene crystal using the Gromos 96 force field

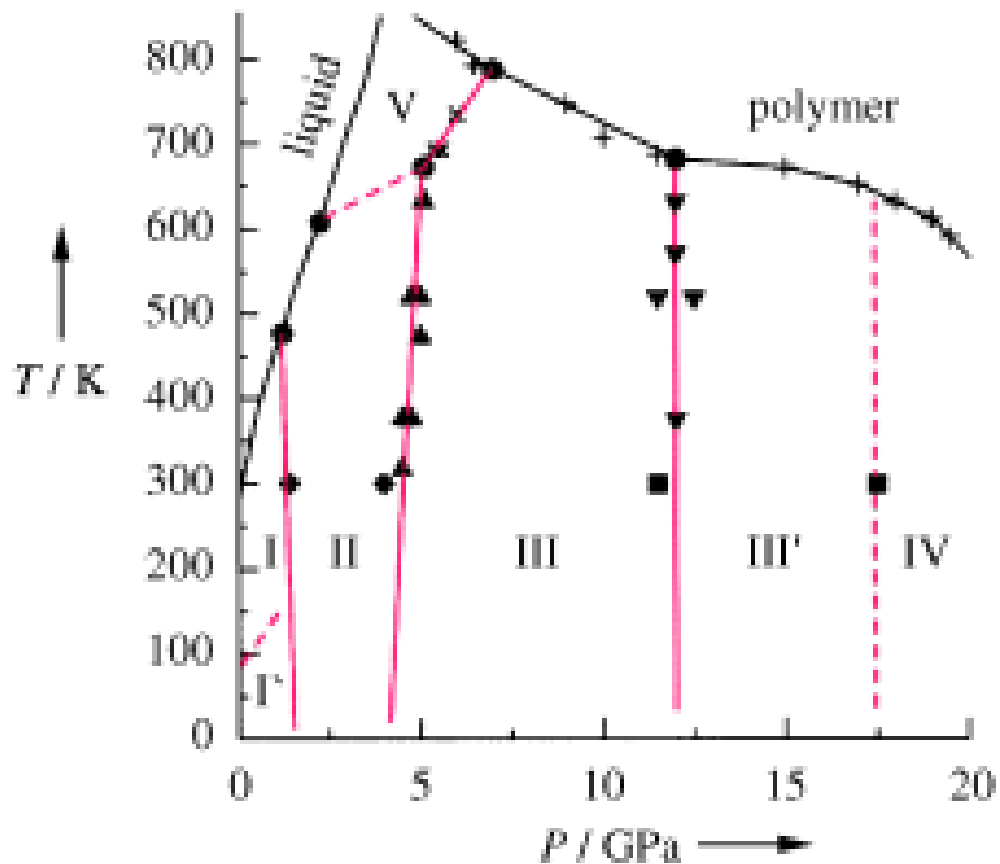
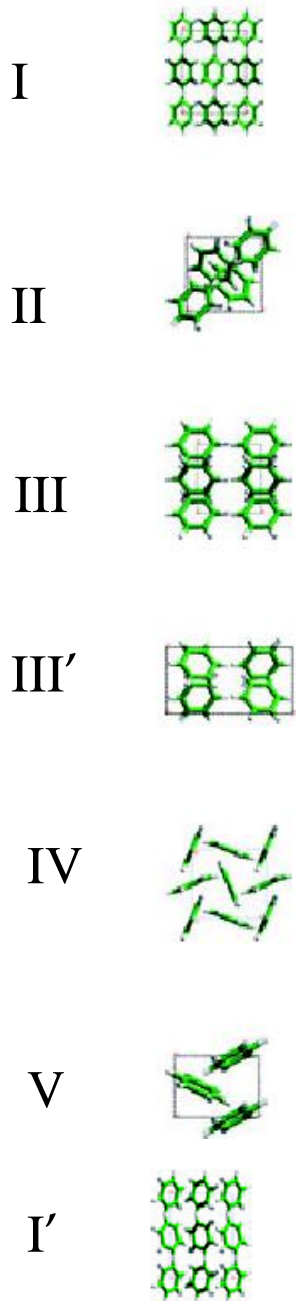
$$T = 300 \text{ K}, \quad T_{\mathbf{h}} = 32,000 - 40,000 \text{ K}$$

$$N = 216 (3 \times 3 \times 3), \quad P = 2 \text{ GPa}$$

Nine different trajectories were initiated from randomly generated initial conditions. Total run time over all nine trajectories is 500 ps.

“Phase diagram” of benzene

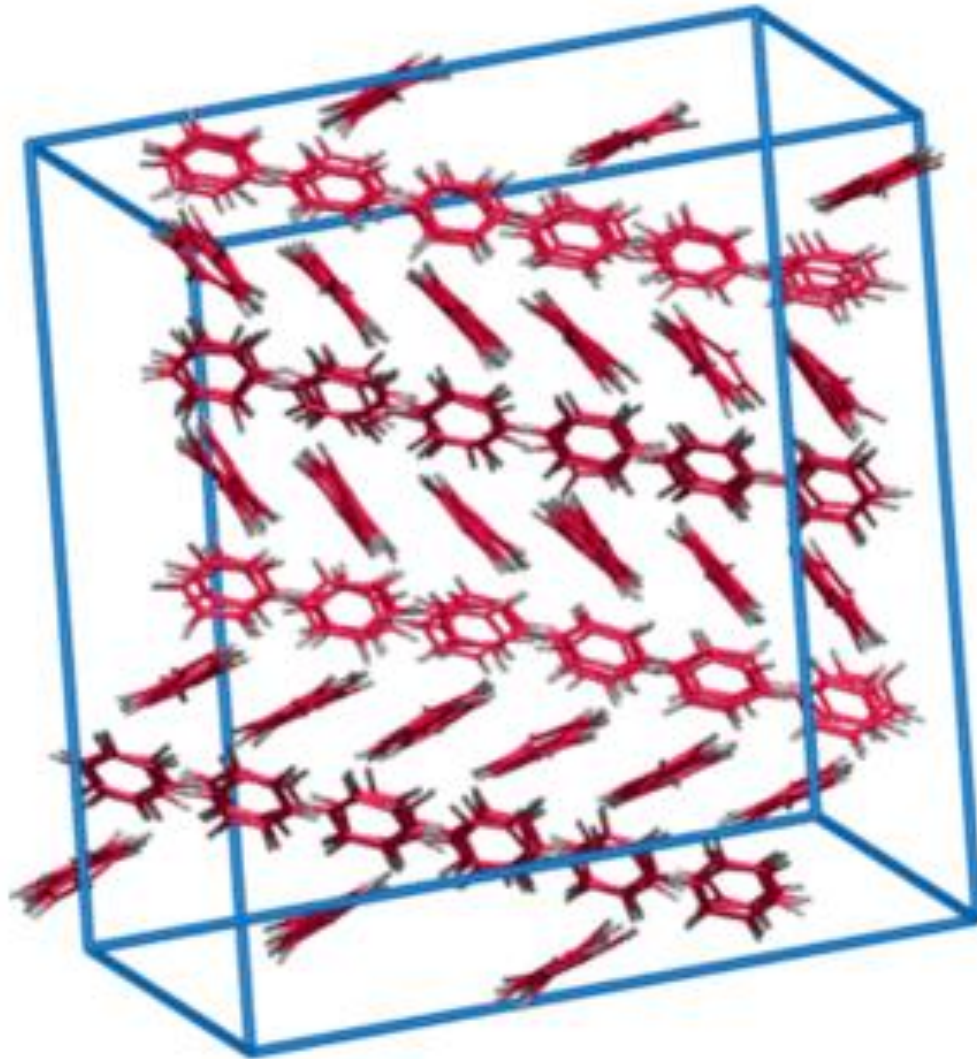
Raiteri, Martonak, Parrinello *Angew. Chem. Intl. Ed.* **44**, 3769 (2005)



$P2_1/c$ (II96) $P4_32_12$ (II98) II01

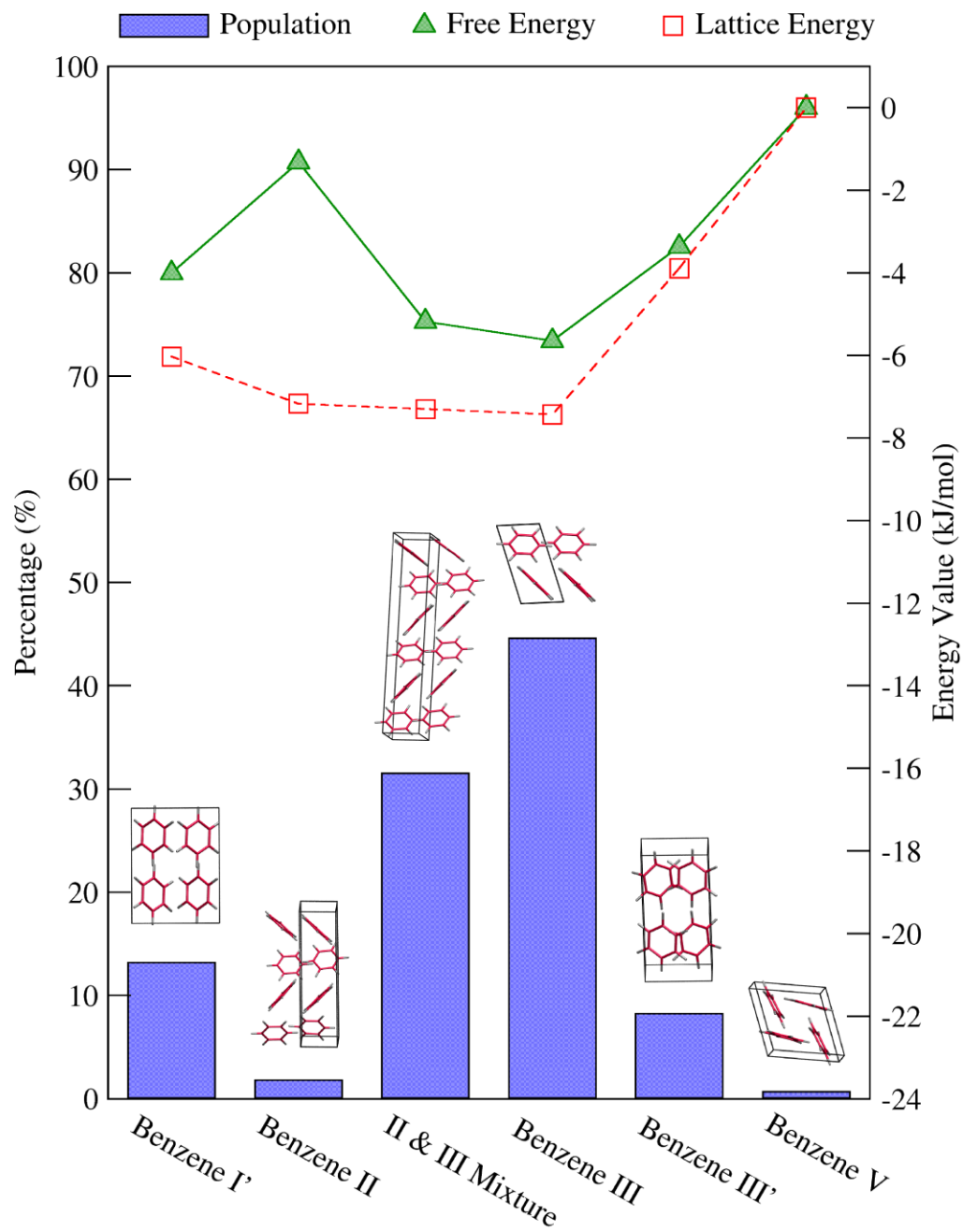
I, II, III, III', IV known from X-ray and Raman scattering

I' and V more speculative



GROMOS FF, $T = 300$ K, $T_h = 32,000 - 40,000$ K

$N = 216$ ($3 \times 3 \times 3$), $P = 2$ GPa, Total run time = 500 ps, 5 ns



Adiabatic conditions on $\{s, p_s\}$: $T_s \square T$, $m_\alpha \gg m$

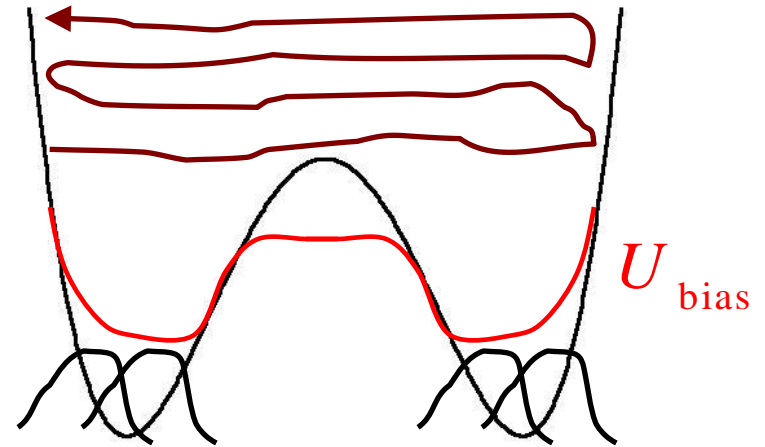
d-AFED with a bias potential

M. Chen, M. Cuendet, and MET *PNAS* (submitted)

Suppose a bias potential $U_{\text{bias}}(s_1, \dots, s_n)$ is applied in the extended phase space.

Apply metadynamics-like bias in the extended-variable space:

$$U_{\text{bias}}(s) = \sum_i A e^{-\sum_{\alpha} \|s_{\alpha} - s_{\alpha}(t_i)\|^2 / 2\sigma^2}$$



Effective Hamiltonian:

$$H(\mathbf{p}, \mathbf{r}, s, p_s) = H(\mathbf{p}, \mathbf{r}) + \sum_{\alpha=1}^n \frac{p_{s_{\alpha}}^2}{2m_{\alpha}} + \frac{1}{2} \sum_{\alpha=1}^n \kappa_{\alpha} (q_{\alpha}(\mathbf{r}) - s_{\alpha})^2 + U_{\text{bias}}(s)$$

Free energy surface:

$$A_{\{\kappa\}}(s_1, \dots, s_n) = -kT_s \ln P_{\text{adb}}^{\{\{\kappa\}\}}(s_1, \dots, s_n) - U_{\text{bias}}(s_1, \dots, s_n)$$

Using the free energy gradient

If the space of CVs is not too large, then we can employ free energy gradients:

$$F_{\alpha}(s_1, \dots, s_n) = -\frac{\partial A}{\partial s_{\alpha}} = \left\langle \kappa_{\alpha} (q_{\alpha}(\mathbf{r}) - s_{\alpha}) \right\rangle$$

Manifestly independent of T_s and can be shown to be independent of bias.

Free energy reconstruction:

1. Expand free energy in a basis set: $A(\{s\}) = \sum_i C_i \psi_i(\{s\})$

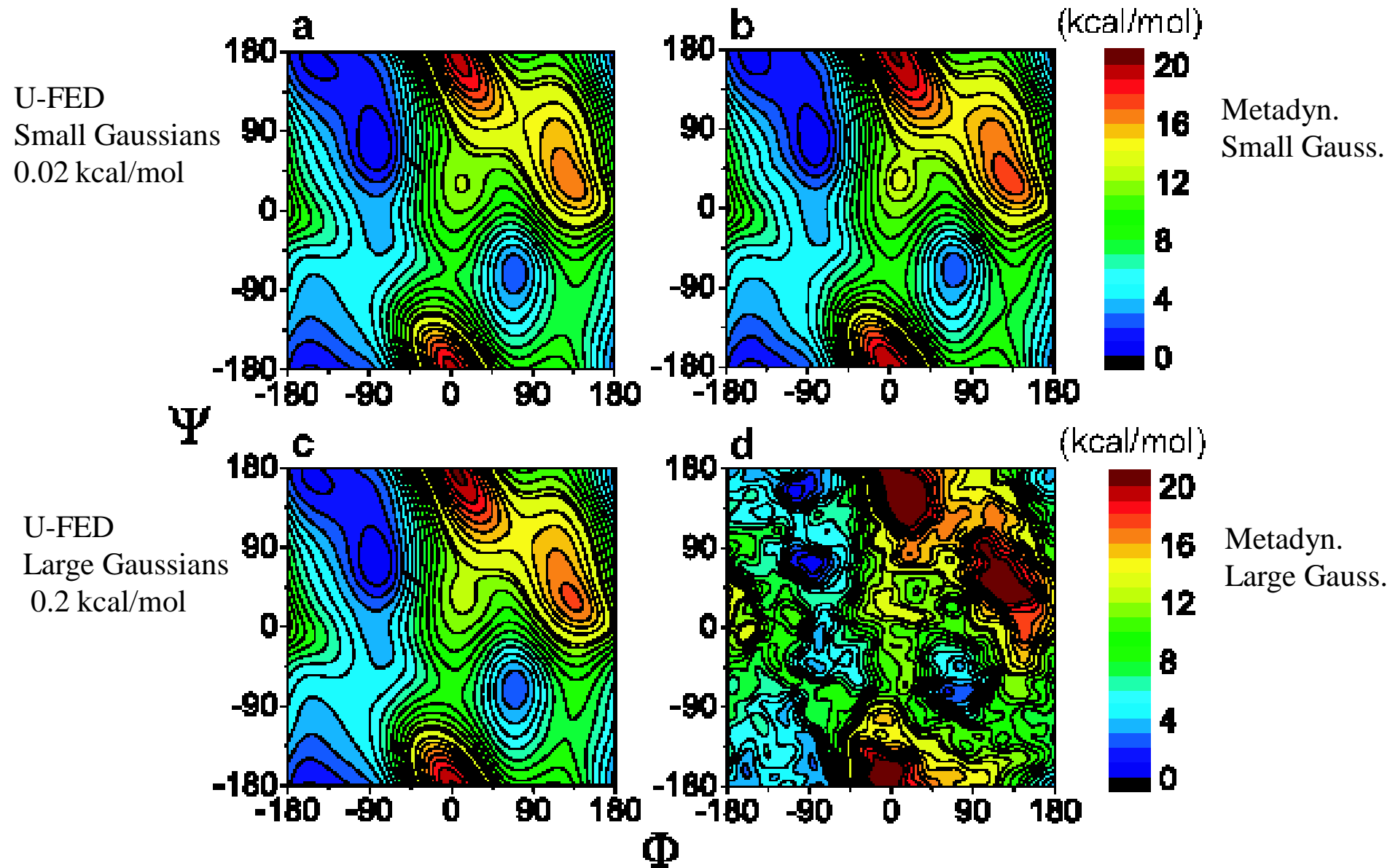
2. Minimize gradient expansion on a grid:

$$f(\{C\}) = \sum_k \sum_{\alpha} \left| \partial_{\alpha} A(\{s^{(k)}\}) + F_{\alpha}(\{s^{(k)}\}) \right|^2$$

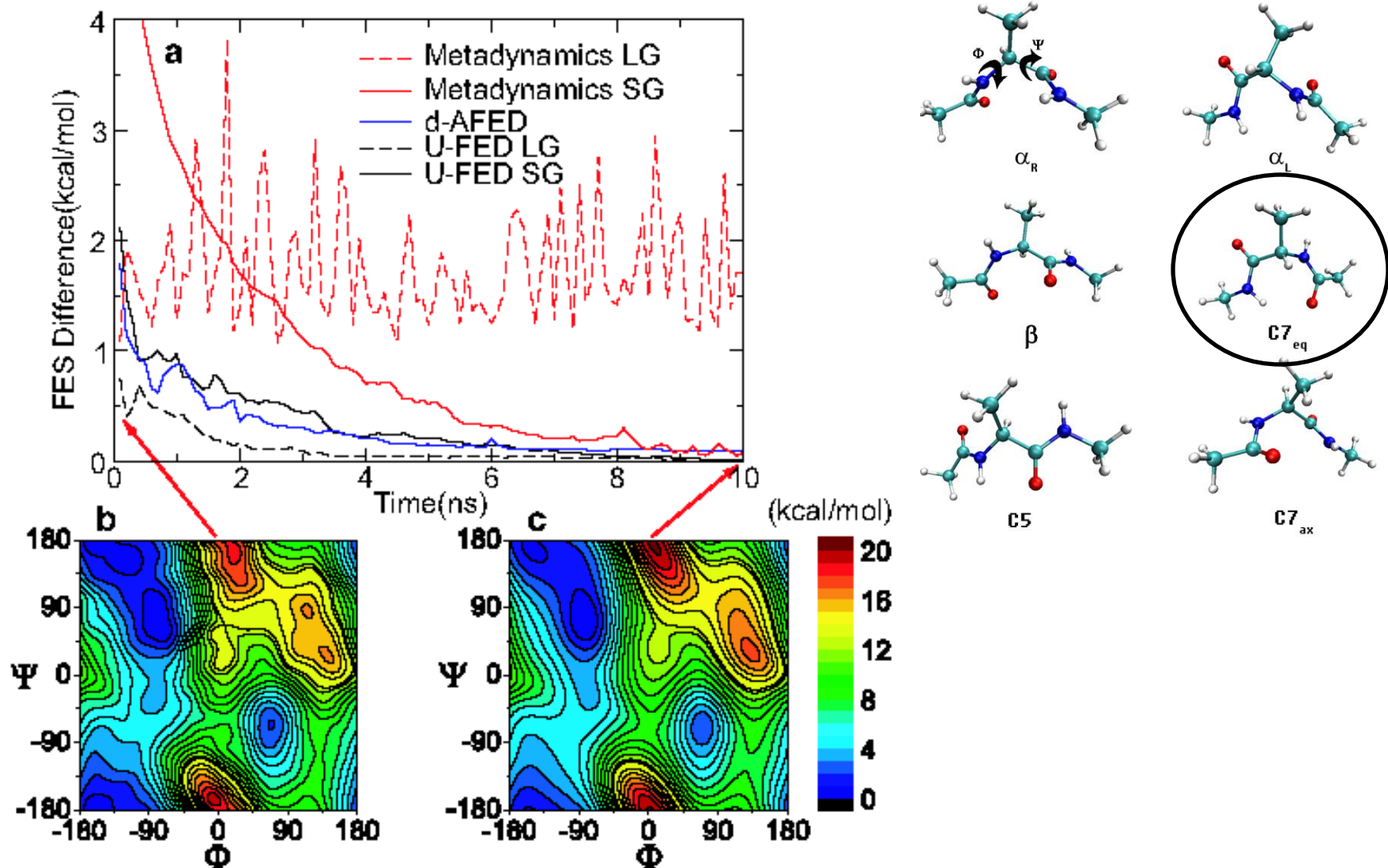
3. Gives a set of linear equations

Overall scheme: Adiabatic dynamics, high T_s , gradients, bias – termed U-FED (unified FED).
M. Chen, M. Cuendet, MET PNAS (submitted).

Ramachandran Surfaces of alanine-dipeptide in gas phase



Convergence of the FES for gas-phase alanine dipeptide

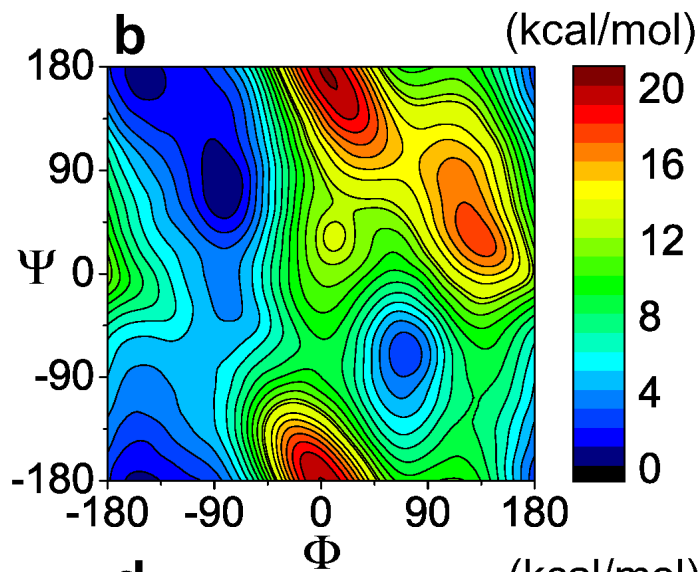
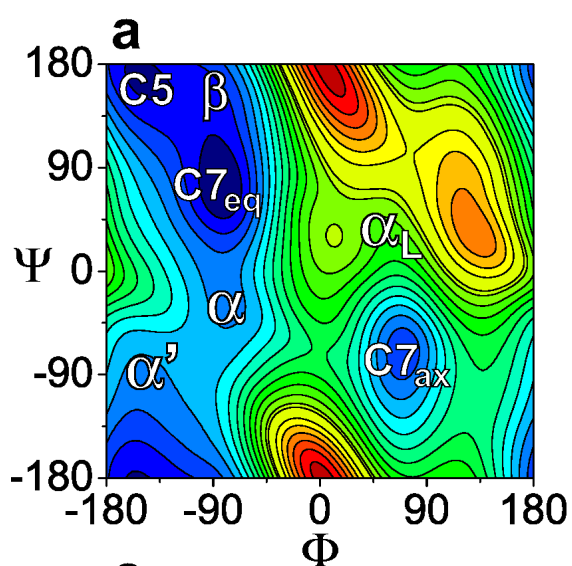


Biased d-AFED 500 ps Large Gauss.

Biased d-AFED 10 ns Large Gauss.

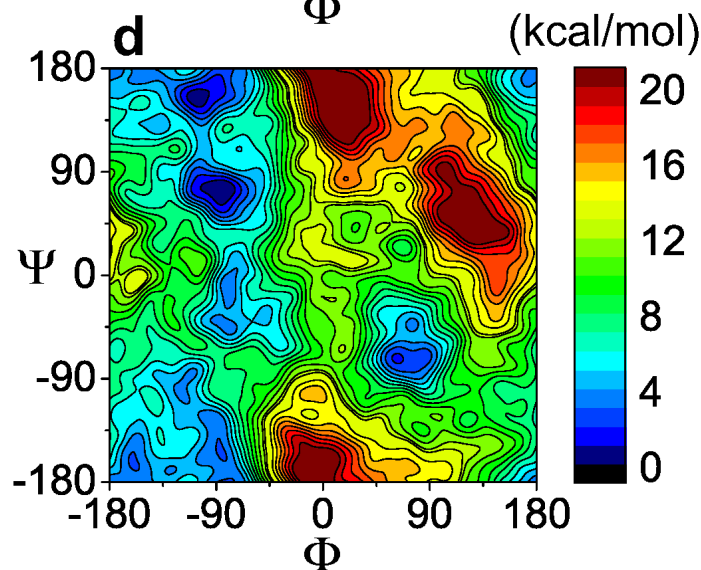
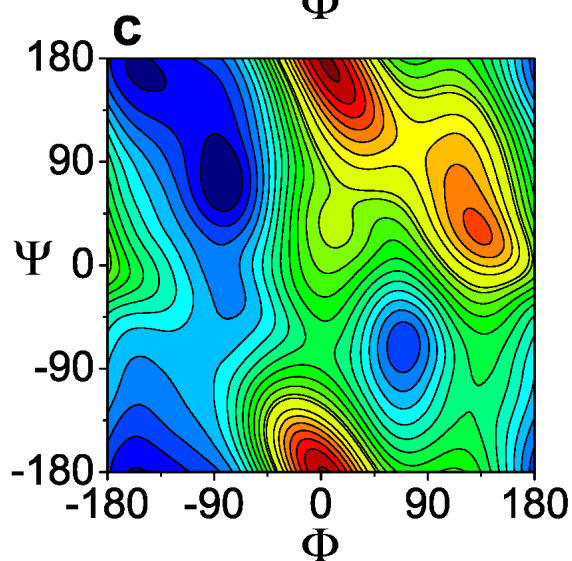
Convergence of FESs for alanine dipeptide in solution

U-FED
Small Gaussians
0.02 kcal/mol

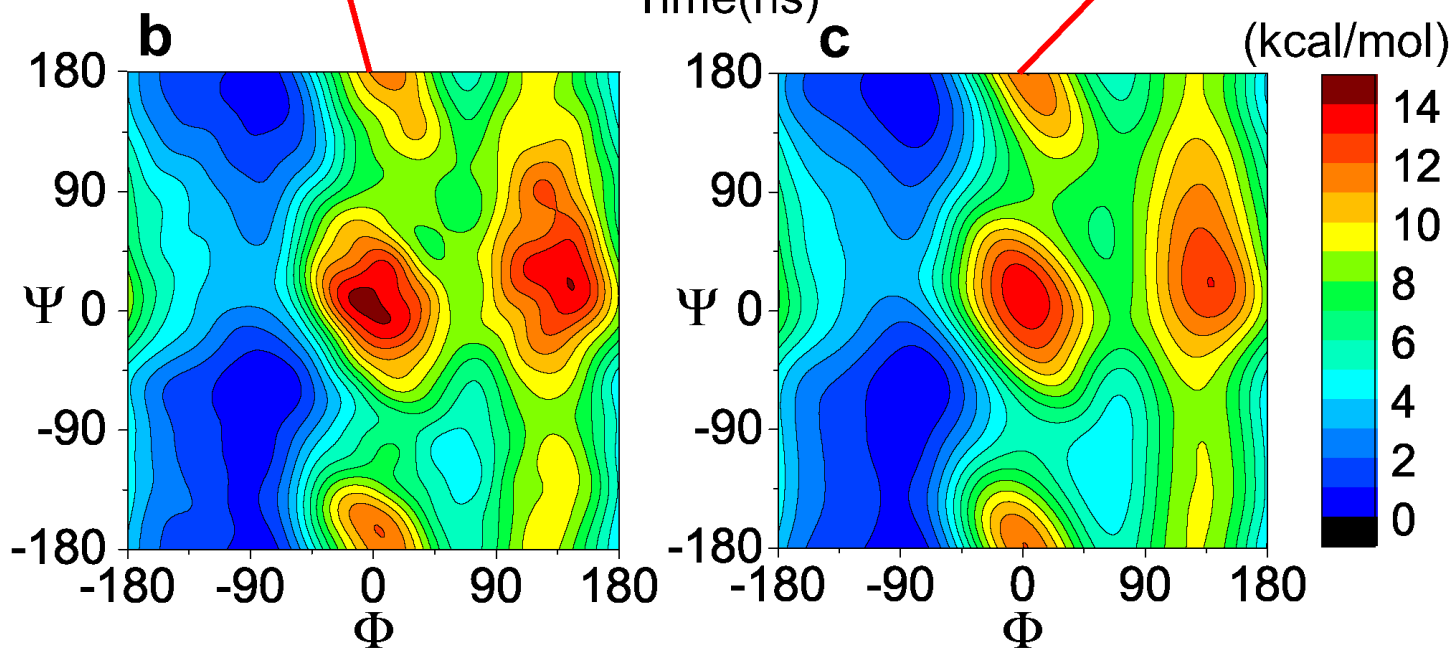
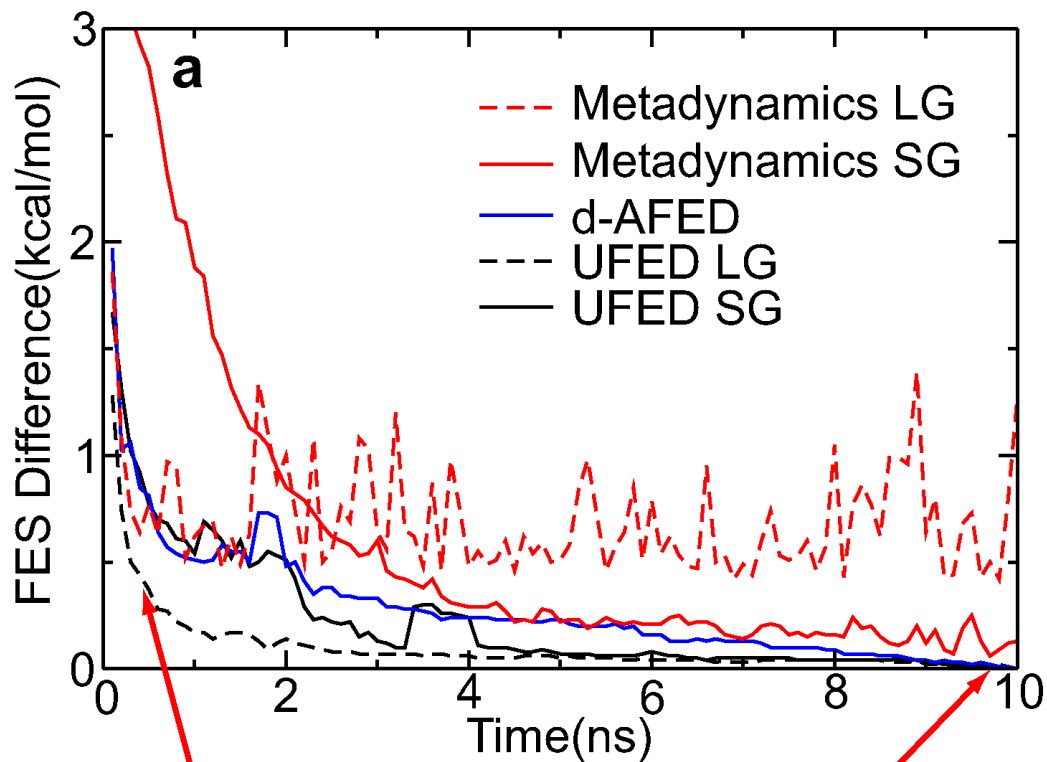


Metadyn.
Small Gauss.

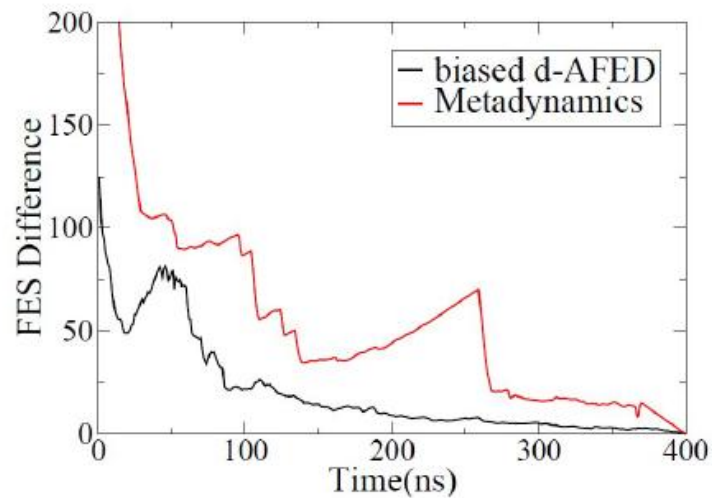
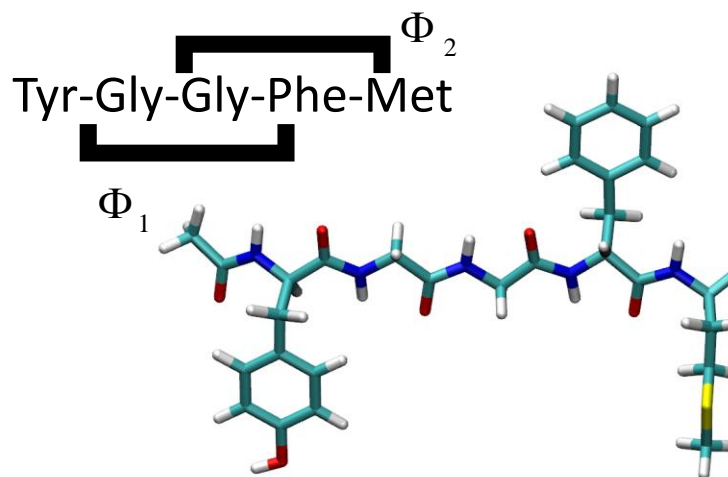
U-FED
Large Gaussians
0.2 kcal/mol



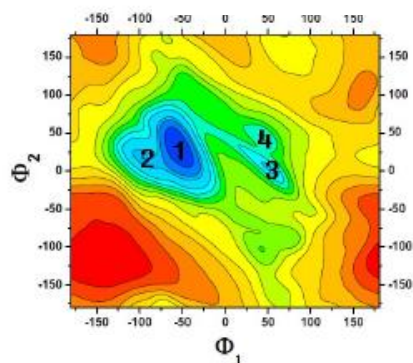
Metadyn.
Large Gauss.



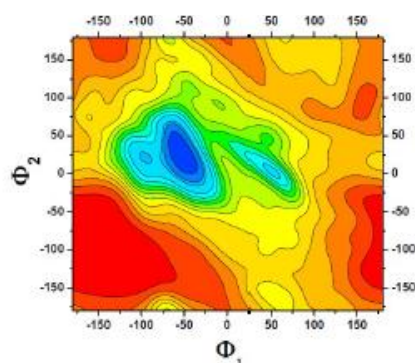
Met-enkephalin



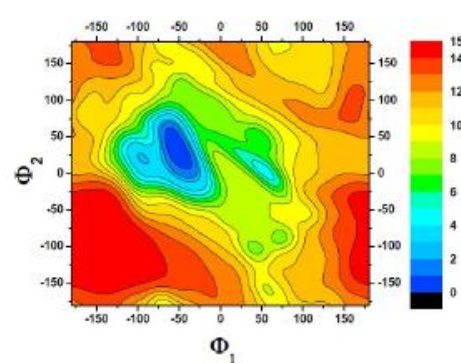
a



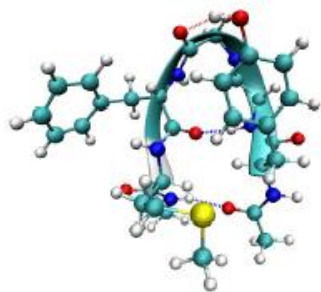
Biased d-AFED



Metadynamics



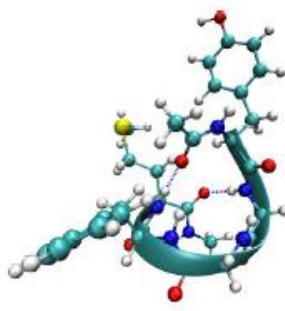
Metadynamics
(GROMACS)



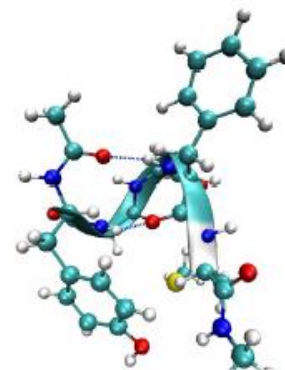
1



2

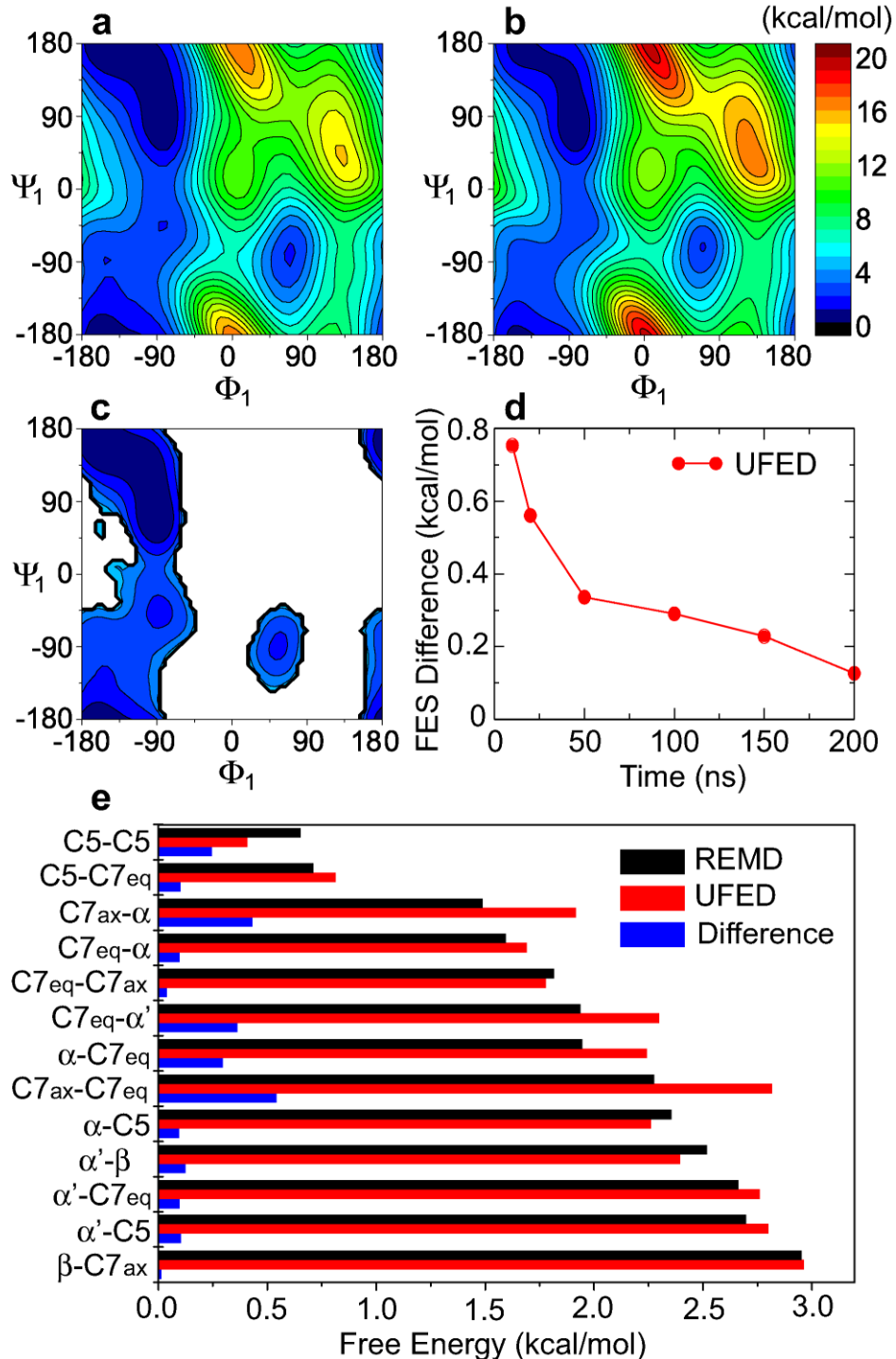
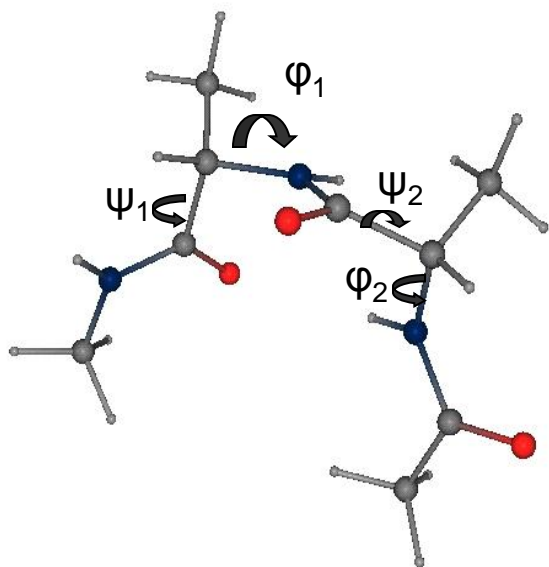


3



4

Conformations of the alanine tripeptide



Conclusions and perspectives

1. Adiabatic dynamics allows free energy surfaces to be explored and generated in a relatively large number of collective variables.
2. Applications include exploration of crystal polymorphs and small oligopeptides.
3. Previously also applied to solvation and binding free energies [J. B. Abrams and MET, *J. Chem. Phys.* (2006)]
4. Combining adiabatic dynamics with bias potentials generated in the spirit of metadynamics further improves the performance. The key is to add the Gaussians only the extended phase space and not to use them to construct the FES.
5. Applications to larger biomolecules and to drug docking currently under way.
6. Combine with Crystal-AFED for studying crystals of strongly interacting molecules.

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- Lula Rosso

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