

1. 3 numbers - ϵ , ϵ' for K

$$\epsilon_K = 2 \times 10^{-3}$$

$$\epsilon_B = \sin 2\beta$$

$$\epsilon' \text{ for } K = 4 \times 10^{-6} \rightarrow$$

evidence against some SWR
 $\sin 2\beta \sim 0.7 \rightarrow$ Qualitative evidence for Stand. Model
2. ϵ'_B is next goal. Now at 2σ from $S_{\pi\pi}$
3. Experimental Future
 - (1) Precision B physics
To limit or find new physics
 - (2) Search for quantities like $d\epsilon$ that vanish in S.M.
 - (3) CP in lepton sector from
 2) oscillations

A. The CKM matrix

In the standard electroweak model, the interactions of the quarks with the charged gauge bosons W are given by

$$g\bar{u}_j V_{ji} \gamma_\lambda (1 - \gamma_5) d_i W^\lambda + \text{H.c.} \quad (3.1)$$

Here $u_j = (u, c, t)$ are the up-type quarks and $d_i = (d, s, b)$ are the down type. V is the unitary CKM (Cabibbo-Kobayashi-Maskawa) matrix, the 3×3 generalization of the Cabibbo mixing matrix. A convenient parametrization of V due to Maiani (1977) is

$$V = \begin{bmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{ts} & V_{tb} \end{bmatrix} = \begin{bmatrix} C_\theta C_\delta & C_\theta S_\theta & S_\theta e^{-i\gamma} \\ -C_\theta S_\theta - C_\delta S_\theta S_\tau e^{i\gamma} & C_\theta C_\delta - S_\theta S_\tau e^{i\gamma} & C_\theta S_\tau \\ S_\theta S_\tau - C_\theta C_\delta S_\theta e^{i\gamma} & -C_\theta S_\tau - C_\delta S_\theta S_\tau e^{i\gamma} & C_\theta C_\delta \end{bmatrix}, \quad (3.2)$$

where $C_\theta = \cos\theta$ and $S_\theta = \sin\theta$. As originally noted by Kobayashi and Maskawa (1973), it is possible by defining the phase of the quark fields to eliminate all but one of the phases in V . Thus all CP violation in this model depends on the phase γ . Experimental data on strange particle and B decay rates can determine the magnitudes V_{us} , V_{cb} , and V_{ub} . Given these magnitudes, there is empirical observation (Wolfenstein, 1983) that the mixing angles have a hierarchical structure allowing expansion in powers of $\lambda = \sin\theta = 0.22$ with

$$\sin\theta = A\lambda^2, \quad (3.3a)$$

$$\sin\theta e^{-i\gamma} = A\lambda^3(\rho - i\eta). \quad (3.3b)$$

where the errors are primarily theoretical.

Expanding V in powers of λ to order λ^3 , we see that the matrix has the simple form

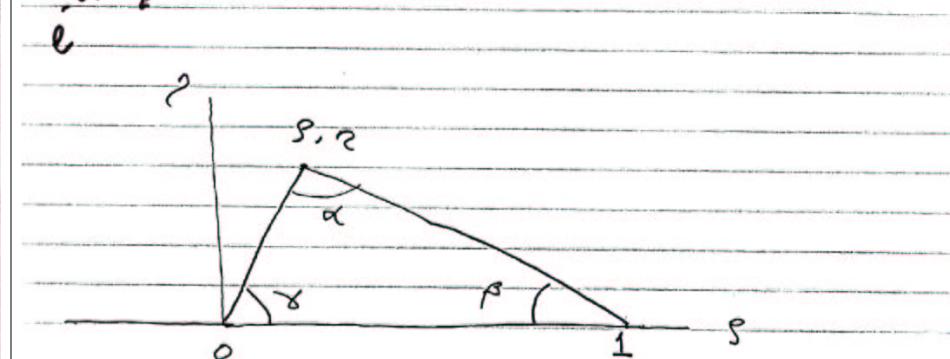
$$V = \begin{bmatrix} 1 - \frac{\lambda^2}{2} & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \frac{\lambda^2}{2} & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{bmatrix}. \quad (3.6)$$

We have chosen a phase convention (that is, a definition of the phases of quark fields) in Eqs. (3.2) and (3.6) such that V is manifestly CP invariant to order λ^2 , and CP

The CP -violating part of the $(K^0 \bar{K}^0)$ mass matrix can be calculated (Ellis *et al.*, 1976) from the second box diagram (Fig. 2). The result of the calculation (Lim, 1981; Buras *et al.*, 1984), including corrections (Gilman and Wise, 1983; Buras *et al.*, 1990), is well represented for $m_t > m_b$ by

$$\epsilon\epsilon^{-i\theta} = 3.4 \times 10^{-3} A^2 \eta B \left[1 + 1.3 A^2 (1 - \rho) \left(\frac{m_t}{m_b} \right)^2 \right]$$

$$V = \begin{bmatrix} 1 - \frac{\lambda^2}{2} & A\lambda^3(\rho - i\eta) & e^{-i\gamma} \\ -\lambda & 1 - \frac{\lambda^2}{2} & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{bmatrix}$$



CP VIOLATION IN B^0

$$P_a(t) \propto e^{-\Gamma t} \left[1 \pm \alpha_a \gamma_a \sin \Delta m t \right]$$

$$\gamma_a = \pm 1 \quad \text{Asymmetry parameter} = \alpha_a$$

$$b \rightarrow c \bar{c} s \quad \alpha_a (\pi K \omega) = \sin 2\beta$$

$$b \rightarrow u \bar{u} d \quad \alpha_a (\pi^+ \pi^-) = \sin 2(\beta + \gamma)$$

$$\tilde{\beta} = \beta$$

LONG TERM FRONTIER FOR CP

1. SIGNALS of PHYSICS BEYOND S.M.

New particle effects on
 Boxes (Mixing)
 Loops (Penguins)
 Electric Dipole Moments

2. CP in LEPTON PHYSICS NEUTRINO MIXING

3. Is CP a SPONTANEOUSLY BROKEN SYMMETRY (θ_{QCD})

4. CP AT A HIGH MASS SCALE TO EXPLAIN BARYON ASYMMETRY

Lepton - Quark Symmetry

$$(u)(c)(t) \\ (d)(s)(b)$$

$$(\nu_e)(\nu_\mu)(\nu_\tau) \\ (e^-)(\mu^-)(\tau^-)$$

1. Three generations

2. $SU(2)_L$ doublets

3. Mass hierarchy e, u, τ

$SO(10)$ 16 of fermions

$$\rightarrow SU(4)_C \times SU(2)_L \times SU(2)_R$$

$$SU(3)_C$$

$$M_R \text{ for } \nu_R$$

See-saw $m_\nu = m_D M_R^{-1} m_D^\top$

m_D = Dirac mass matrix for ν

M_R = Majorana mass matrix for N_R

m_ν defined by 9 numbers

$$U_{PMNS} = \begin{pmatrix} c_\theta & s_\theta & s_{13} e^{-i\delta} \\ -s_\theta & c_\theta & \frac{1}{\sqrt{2}} \\ \frac{s_\theta}{\sqrt{2}} & \frac{c_\theta}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{s_\theta}{\sqrt{2}} & -\frac{c_\theta}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{pmatrix}$$

$$s_\theta^2 \approx 0.3$$

$$m_\nu^{\text{diag}} = \begin{pmatrix} m_3 & & \\ & m_2 e^{i\alpha} & \\ & & m_1 e^{i\beta} \end{pmatrix}$$

of 9 - 4 are measured

$$\begin{aligned} m_3^2 - m_2^2 &\approx 3 \times 10^{-3} \text{ eV}^2 \\ m_2^2 - m_{1,4}^2 &\approx 7 \times 10^{-2} \text{ eV}^2 \end{aligned}$$

$$\nu_x = (\nu_u + \nu_d)/\sqrt{2}$$

$$\nu_y = (\nu_u - \nu_d)/\sqrt{2}$$

$$U_{PMNS} = \begin{pmatrix} c_\theta & s_\theta & s_{13} e^{-i\delta} \\ -s_\theta & c_\theta & e_{12} \\ e_{31} & e_{21} & 1 \end{pmatrix}$$

$s_\theta \approx .55$

$$U_{CKM} = \begin{pmatrix} c_\theta & s_\theta & s_\delta e^{-i\delta} \\ -s_\theta & c_\theta & s_\delta \\ s_\theta c_\delta - s_\theta c_\delta e^{-i\delta} & -s_\delta & 1 \end{pmatrix}$$

$$s_\theta \approx .22$$

$$s_\delta e^{-i\delta} = A \lambda^3 (f - i\eta) \approx .003 e^{-i\delta}$$

$$s_\delta = A \lambda^2 \approx .04$$

Undetermined

1. Mass scale m_1
2. One mixing angle s_{13}
3. One phase δ in U_{PMNS}
4. Discrete ambiguity:
sign of Δm_{32}^2

$$\begin{array}{ccc} \stackrel{3}{\text{---}} & & \stackrel{2}{=}; \\ ; & = & \stackrel{3}{\text{---}} \end{array}$$
5. Majorana or Dirac?
6. If Majorana:
Phases α, β

CP Violation in ν Oscillations (Detecting the phase δ)

$$P(\nu_\mu \rightarrow \nu_e) - P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e)$$

1. Rate is proportional to s_{13}^2
2. Difference depends not only on Δm_{32}^2 , but also Δm_{21}^2

Requires long baseline and/or low E

3. In matter the difference is not a direct sign of CP.
Fit data to s_{13} and δ .

4. T violation ("in" matter effect)
of $\nu_\mu \rightarrow \nu_e$ with $\nu_e \rightarrow \nu_\mu$

CP VIOLATION "PROPOSALS"
 $\nu_\mu \rightarrow \nu_e / \bar{\nu}_\mu \rightarrow \bar{\nu}_e$

1. JHF II \rightarrow Hyper K

0.7 Gev 300 km

Arafune, Sato Phys Rev D 56, 3093
 Y. Itow et al hep-ex/0106019

2. Accelerator Long Baseline

1 Gev, 1200 - 3000 km

Marciano Phys. Rev D

0.5-7 Gev 2000 km BNL-69395

3. "Factory"

20-30 Gev, 3000-7000 km

Freund et al Nucl. Phys. B 578, 27

Barger et al hep-ph/0003184

Cervera et al Nucl. Phys. B 579, 17

4. Low Energy neutrinos

100 Mev, $L = 50$ km

Mishra et al Nucl. Phys. hep-ph/000901

Probing S_{13}^2
 $(\sin^2 2\theta_{13} = 4S_{13}^2)$

CHOOZ $S_{13}^2 < .025 - .03$

Within 10 years ?

JHF \rightarrow SUPER K

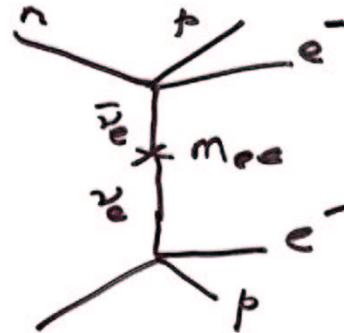
0.6 Gev, 295 km $\rightarrow .003$

Fermilab Off axis beam

"2Gev" 750 km $\rightarrow .003$

CNGS, ICARUS with new beam $\rightarrow .003$

Neutrinoless Double Beta Decay



$$m_{ee} = |U_{e1}|^2 m_1 e^{i\beta} + |U_{e2}|^2 m_2 e^{i\alpha} + |U_{e3}|^2 m_3$$

Example. $m_{ee} = 0.2 \text{ eV}$

$$m_1 \approx m_2 \approx m_3$$

$$(0.2 \text{ eV}) = \cos^2 \theta_0 m_1 e^{i\beta} + \sin^2 \theta_0 m_3 e^{i\alpha}$$

Extremes $e^{i\beta} = e^{i\alpha} = 1$ $m_1 = 0.2 \text{ eV}$

$$e^{i\beta} = \pm 1 \quad e^{i\alpha} = \pm 1 \quad m_1 = \frac{0.2 \text{ eV}}{\cos 2\theta} \approx 0.5 \text{ eV}$$

If Cosmology $\rightarrow m_1 < 0.5 \text{ eV}$
 ${}^3\text{H}$ decay $\rightarrow m_1 > 0.2 \text{ eV}$
 Then 3 Majorana phase

Conclusions

Future for ν Mass + Oscillation

* 1. Probe value of $|U_{e3}|^2$

* 2. Search for $\beta \neq 0$

1 : If $|U_{e3}|^2$ is large enough
 Find CP phase δ
 Hierarchy

2 : If $\langle m_{ee} \rangle$ is large enough

Majorana

Range of m_1 ,
 CP phases α, β ??

3. Physics beyond the standard :

ν_s , FCNC

Electric dipole moments
 of e or n