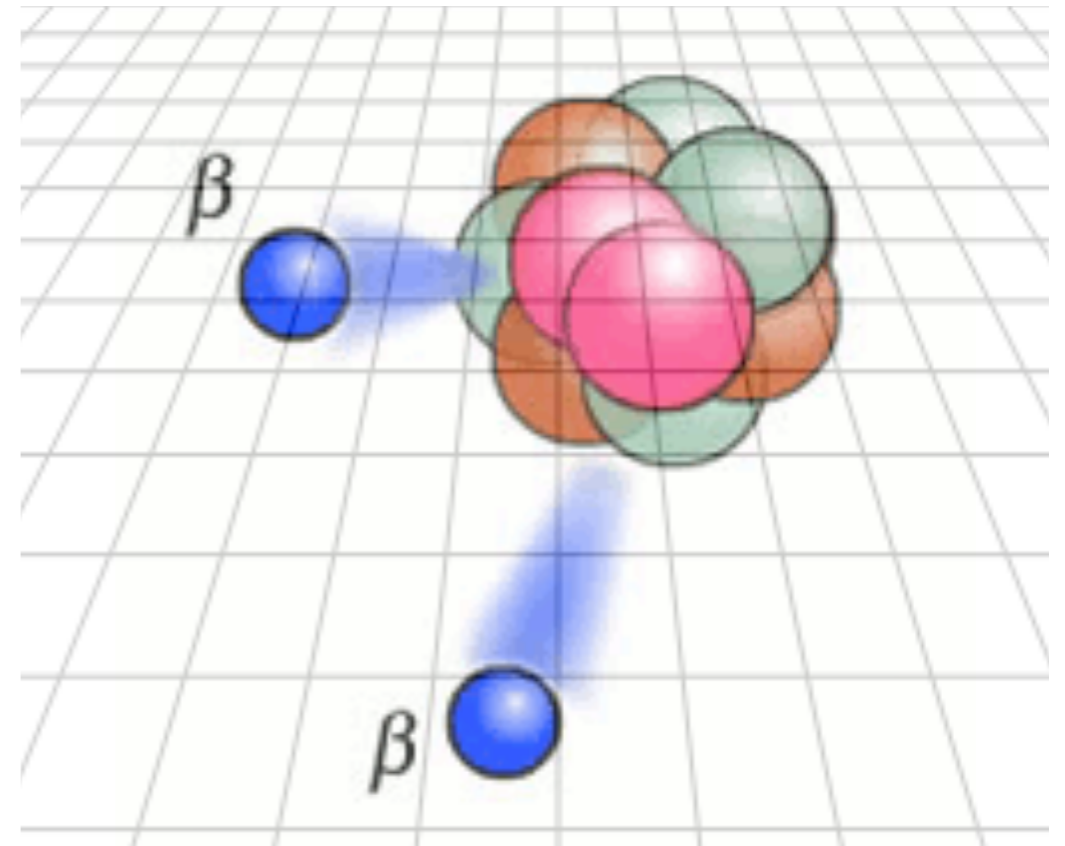


# Neutrinoless Double Beta Decay from Lattice QCD

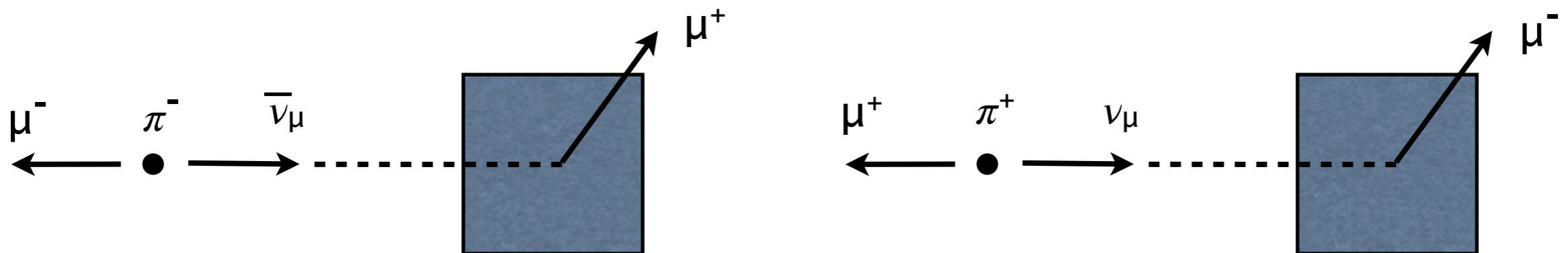
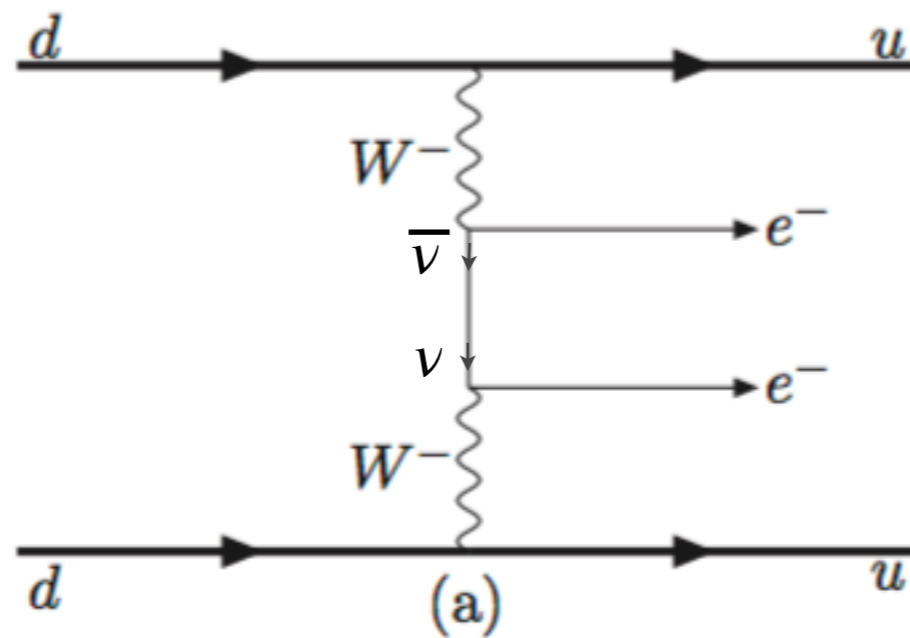


Amy Nicholson  
UC Berkeley  
*Symmetry Tests in Nuclei and Atoms*  
KITP, Sept. 22, 2016



# $0\nu\beta\beta$ and Lepton Number

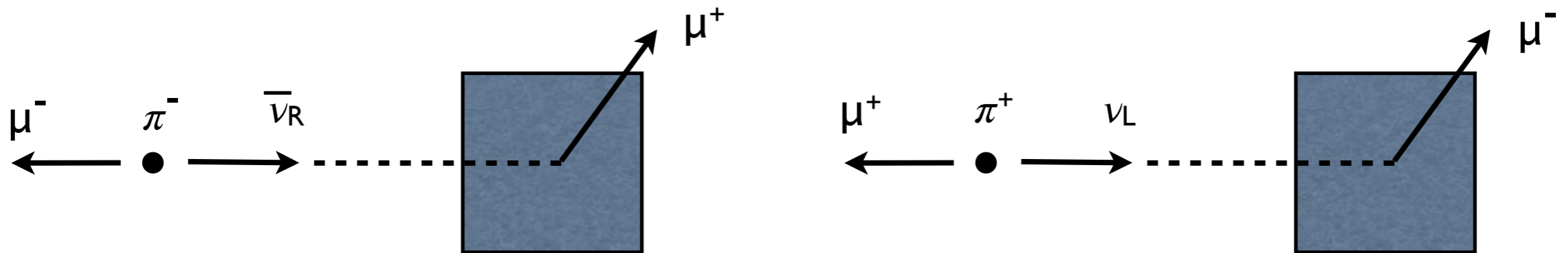
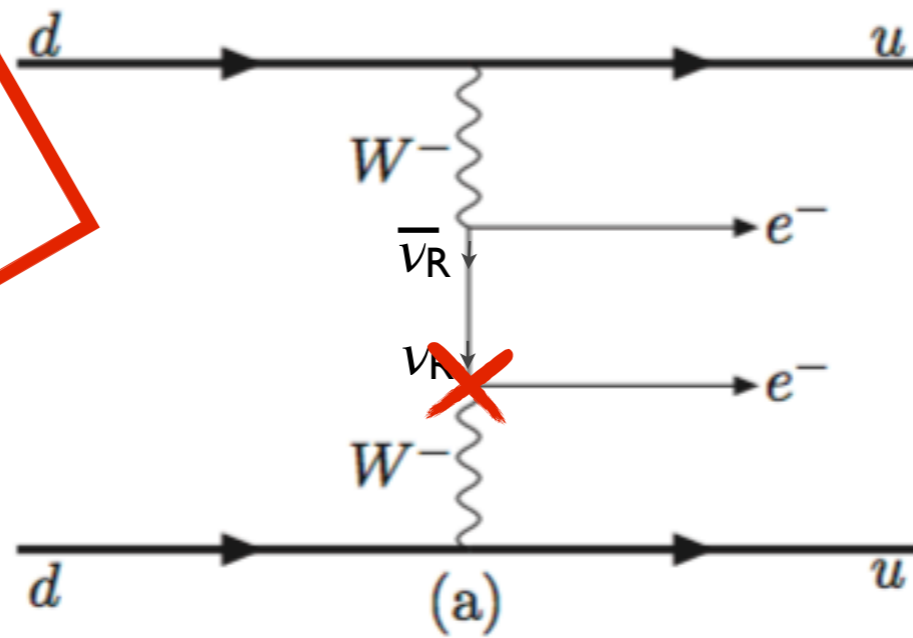
Neutrinos have no known charge or other additively conserved quantum number



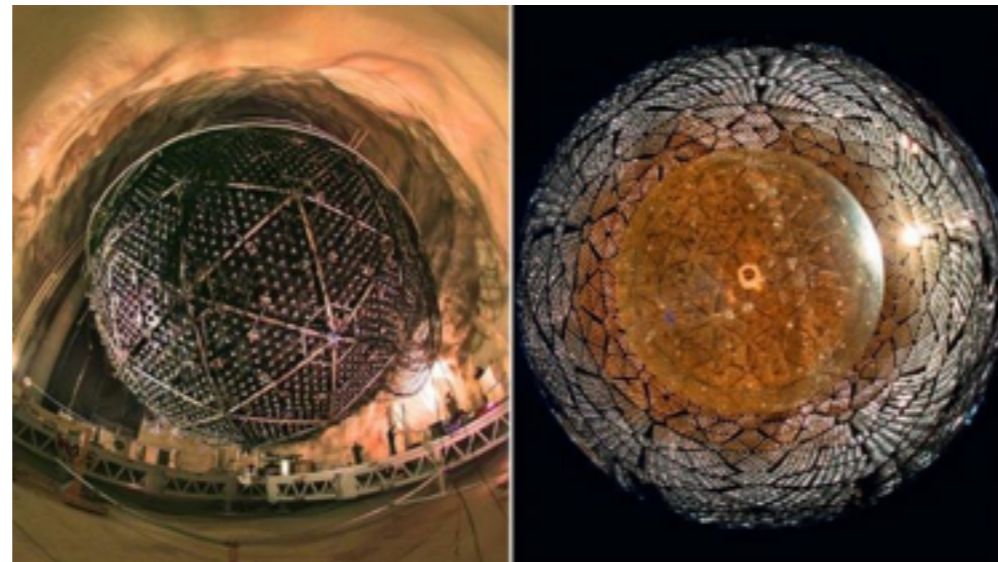
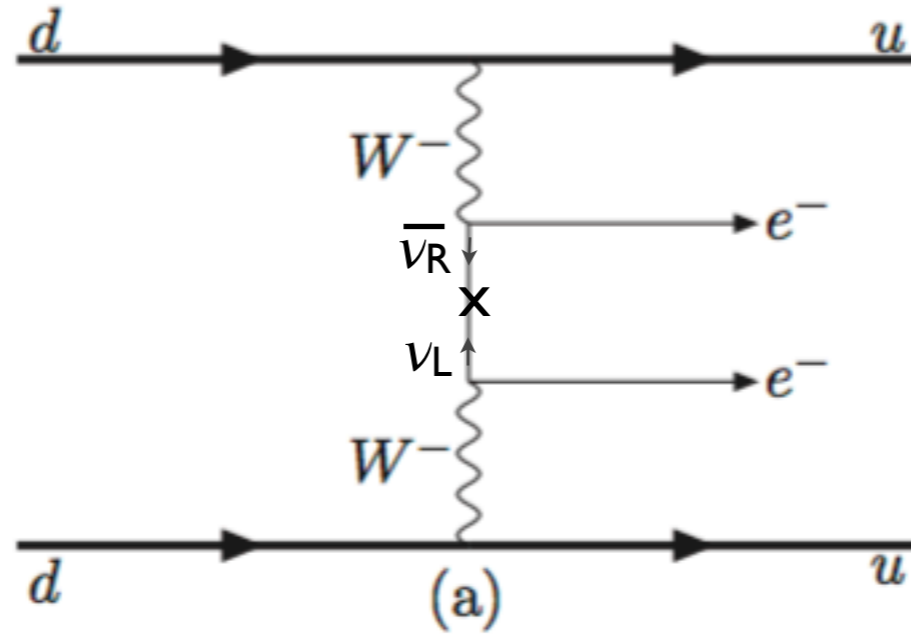
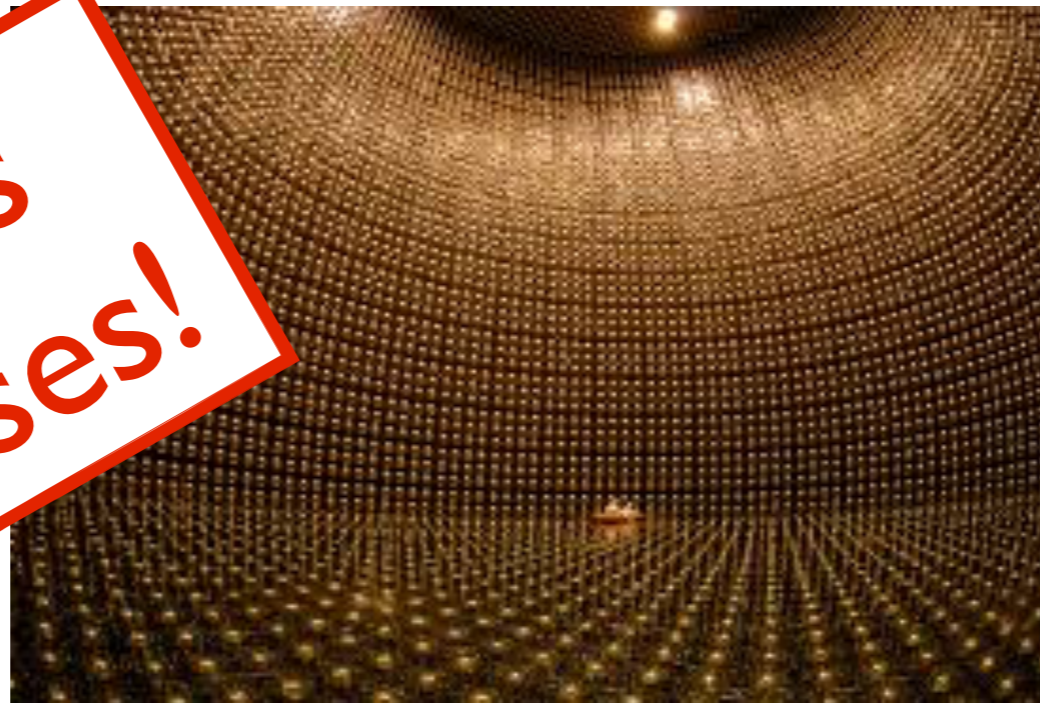
# $0\nu\beta\beta$ and Lepton Number

Neutrinos have no known charge or other additively conserved quantum number

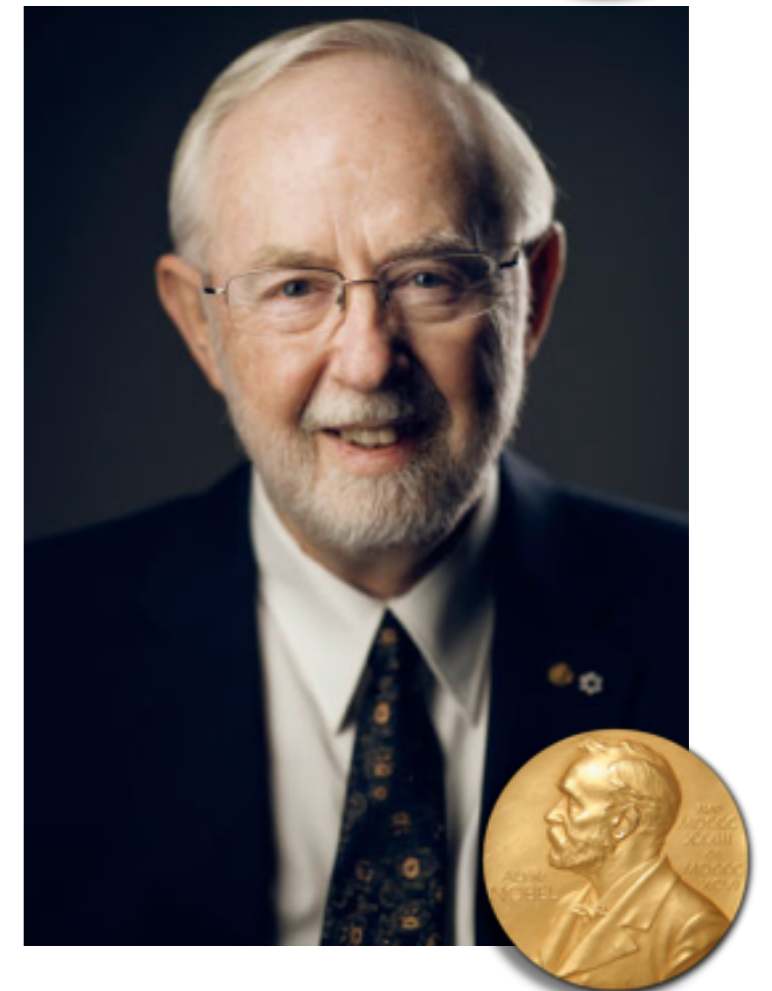
Forbidden by helicity?



Neutrinos  
have masses!



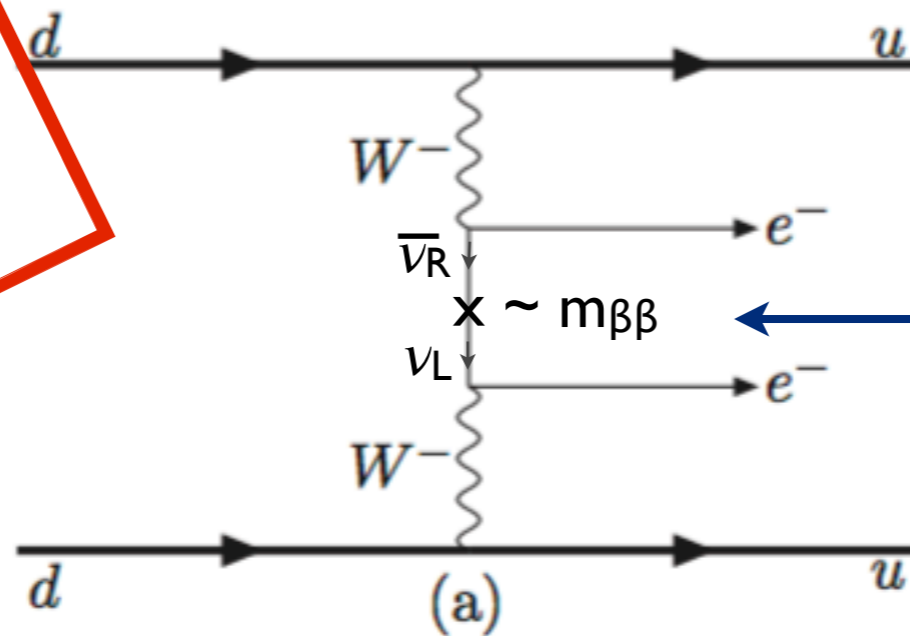
Takaaki Kajita  
(Super-K)  
Arthur B.  
McDonald  
(SNO)  
Nobel Prize,  
2015



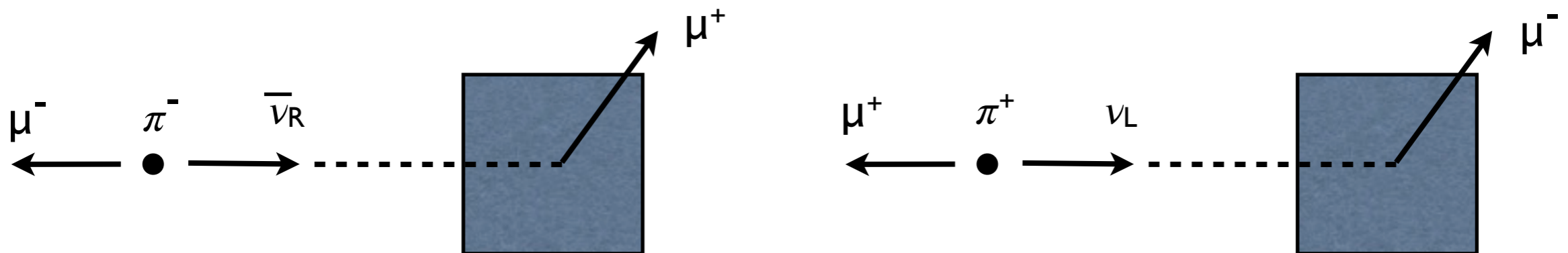
# $0\nu\beta\beta$ and Lepton Number

Neutrinos have no known charge or other additively conserved quantum number

But they're tiny!



oscillation experiments don't tell us absolute mass scale  
 $0\nu\beta\beta$  will!



# Majorana or Dirac?

- Anything not forbidden by symmetry should occur in nature

$$\mathcal{L}_5 = -m \left( \bar{L} \tilde{H} \right) \left( \tilde{H} L \right)^\dagger$$

- Why are neutrinos so light?
  - Dirac mass on its own requires fine-tuning



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$$\begin{pmatrix} M_L & M_D \\ M_D & M_R \end{pmatrix}$$

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$$m_l \sim M_D^2 / M_R \quad m_h \sim M_R$$

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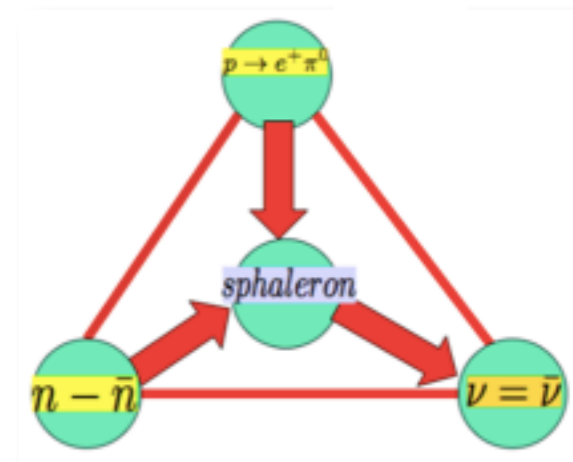
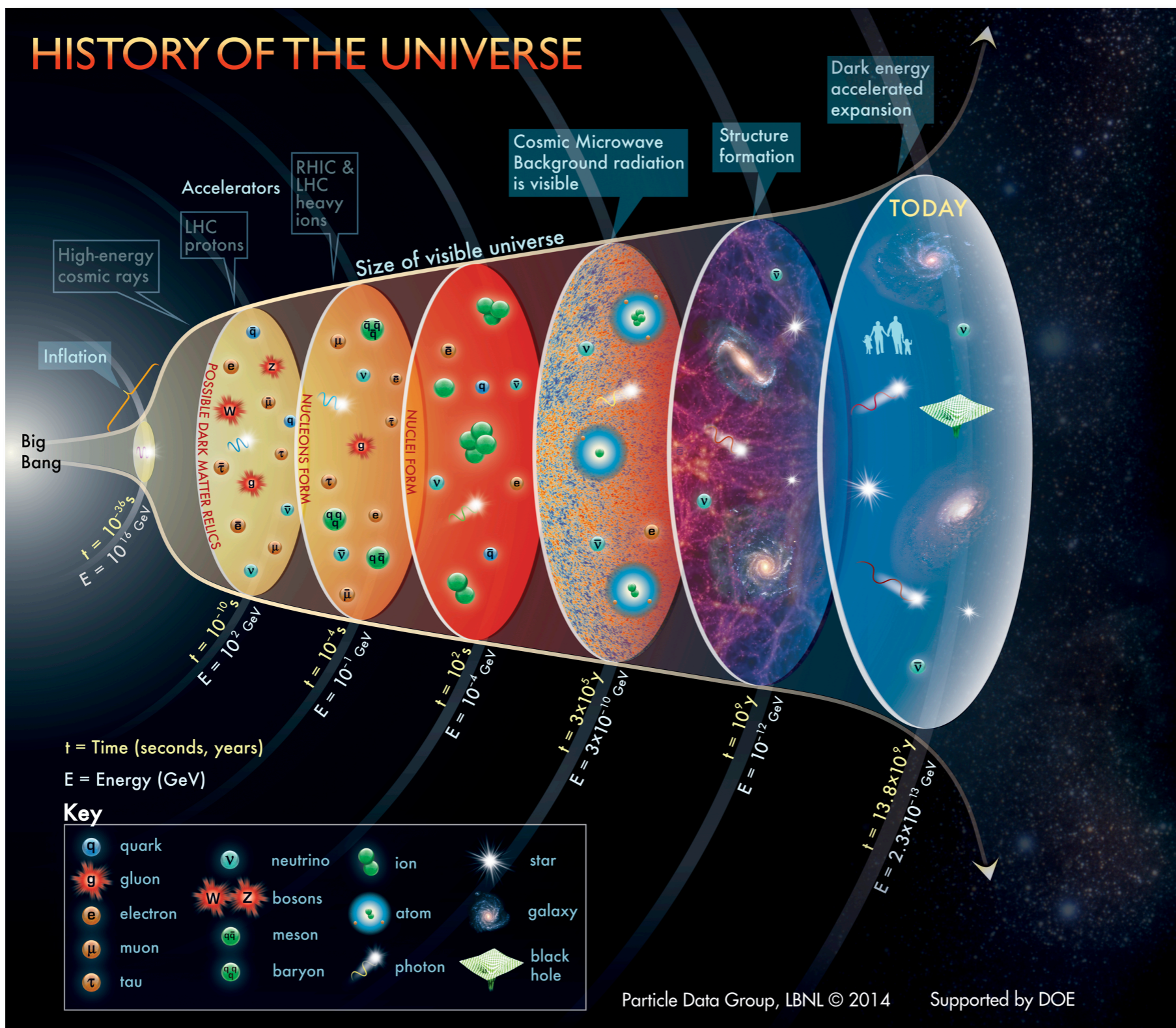
$$\begin{pmatrix} 0 & M_D \\ M_D & M_R \end{pmatrix}$$

$$m_l \sim M_D^2 / M_R \quad m_h \sim M_R$$

$$M_D \sim 200 \text{ GeV} \quad m_l \sim 0.05 \text{ eV}$$

$$M_R \sim 10^{15} \text{ GeV}$$

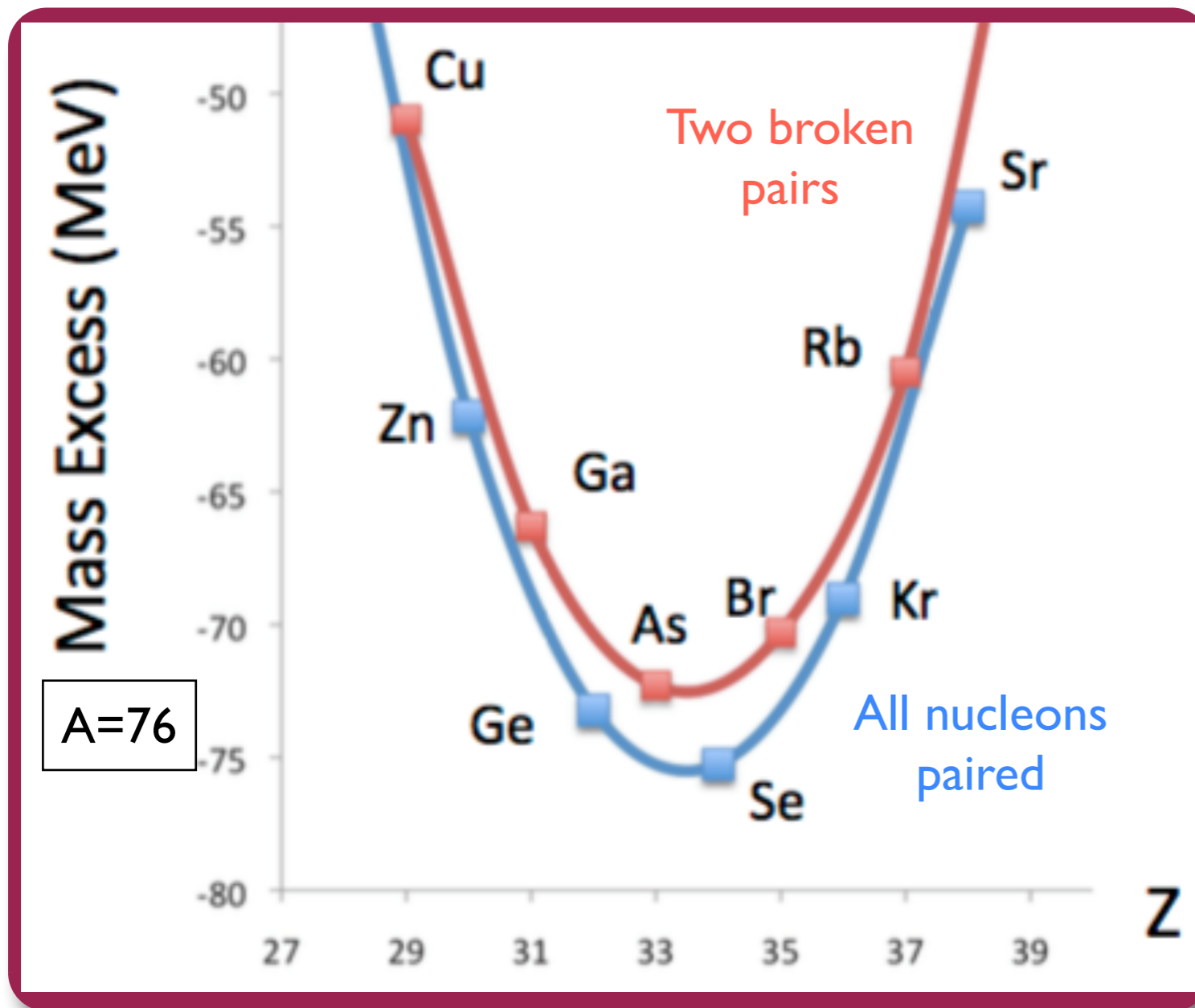
# If observed, could help explain matter/anti-matter asymmetry in the universe!



Jansen (1996)  
 Bödeker,  
 Moore,  
 Rummukainen  
 (2000)  
 Fodor (2000)

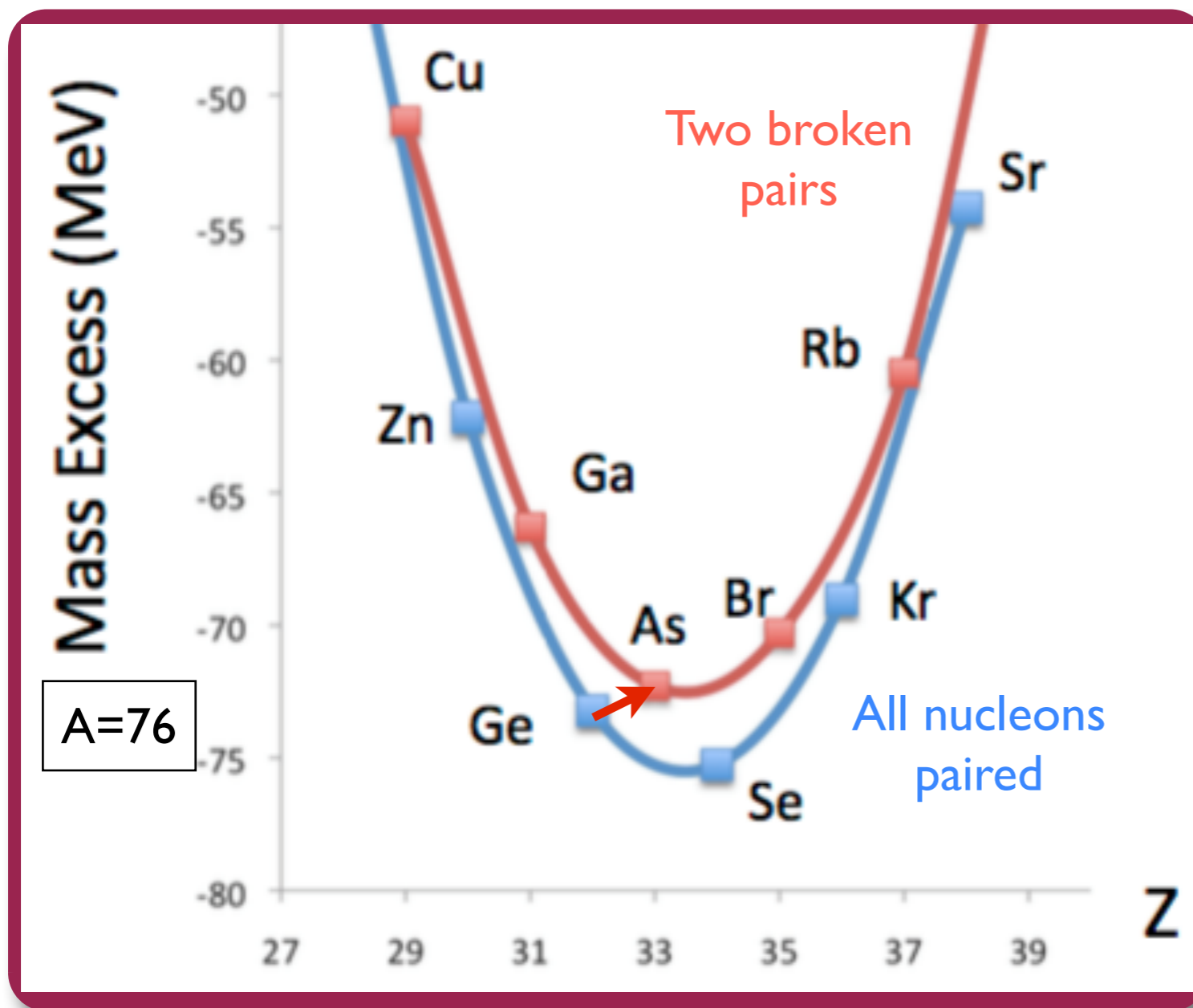
# Experiment

Nuclear physics gives us a natural filter for the process

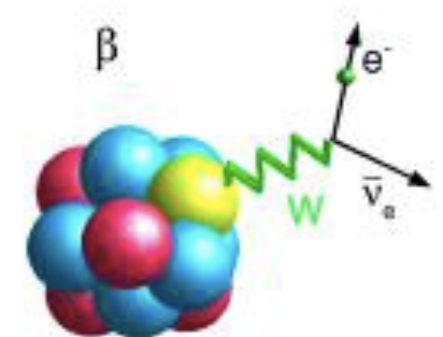


# Experiment

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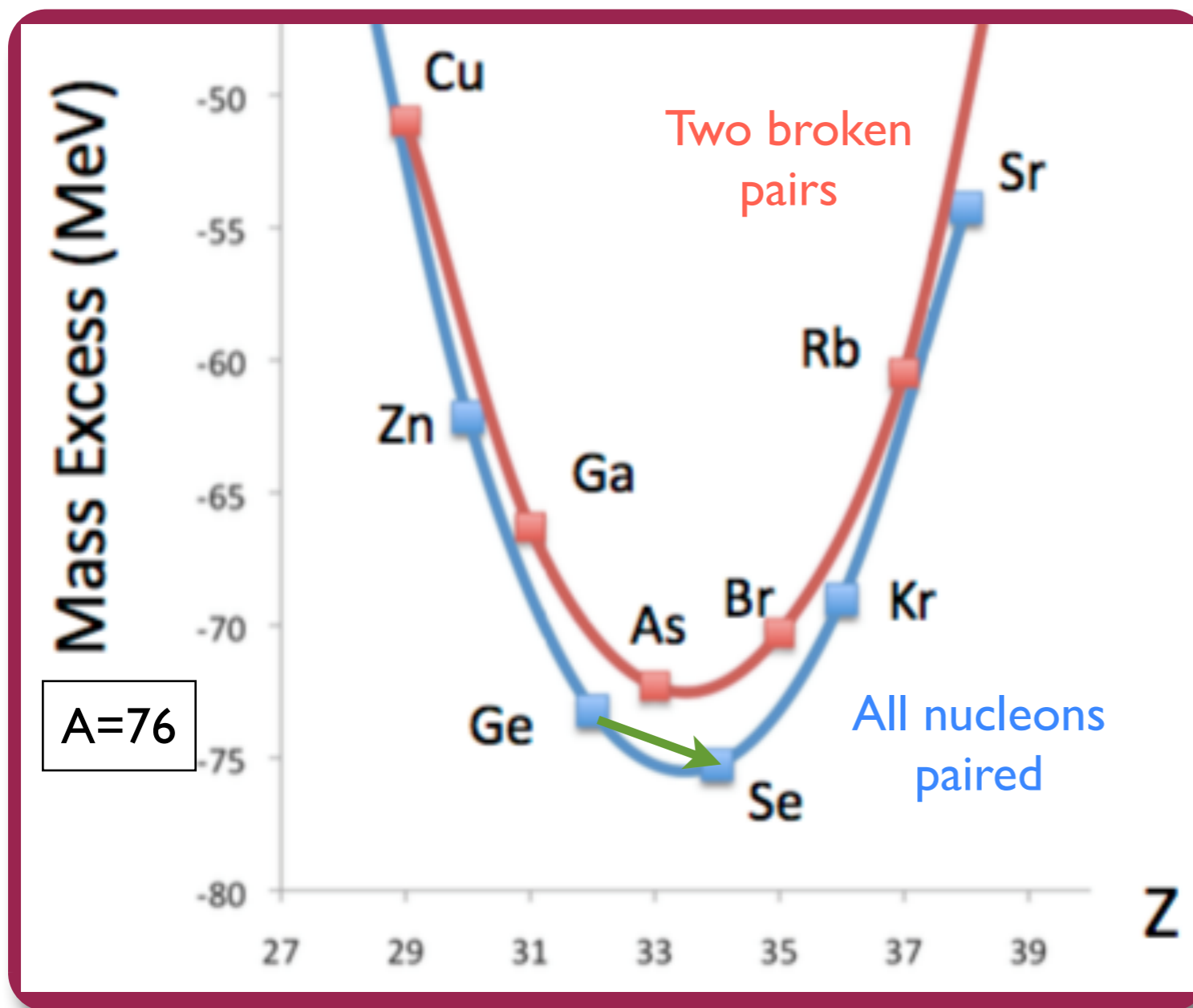


Energetically forbidden

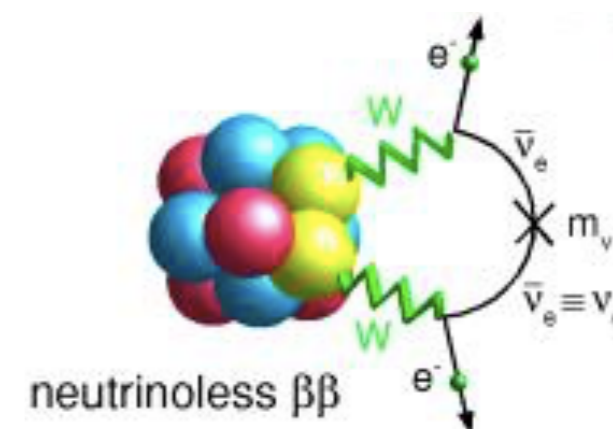
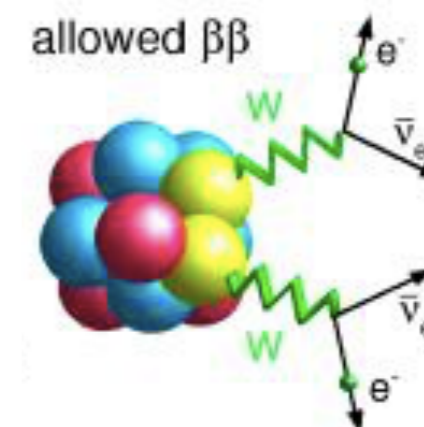


# Experiment

Nuclear physics gives us a natural filter for the process

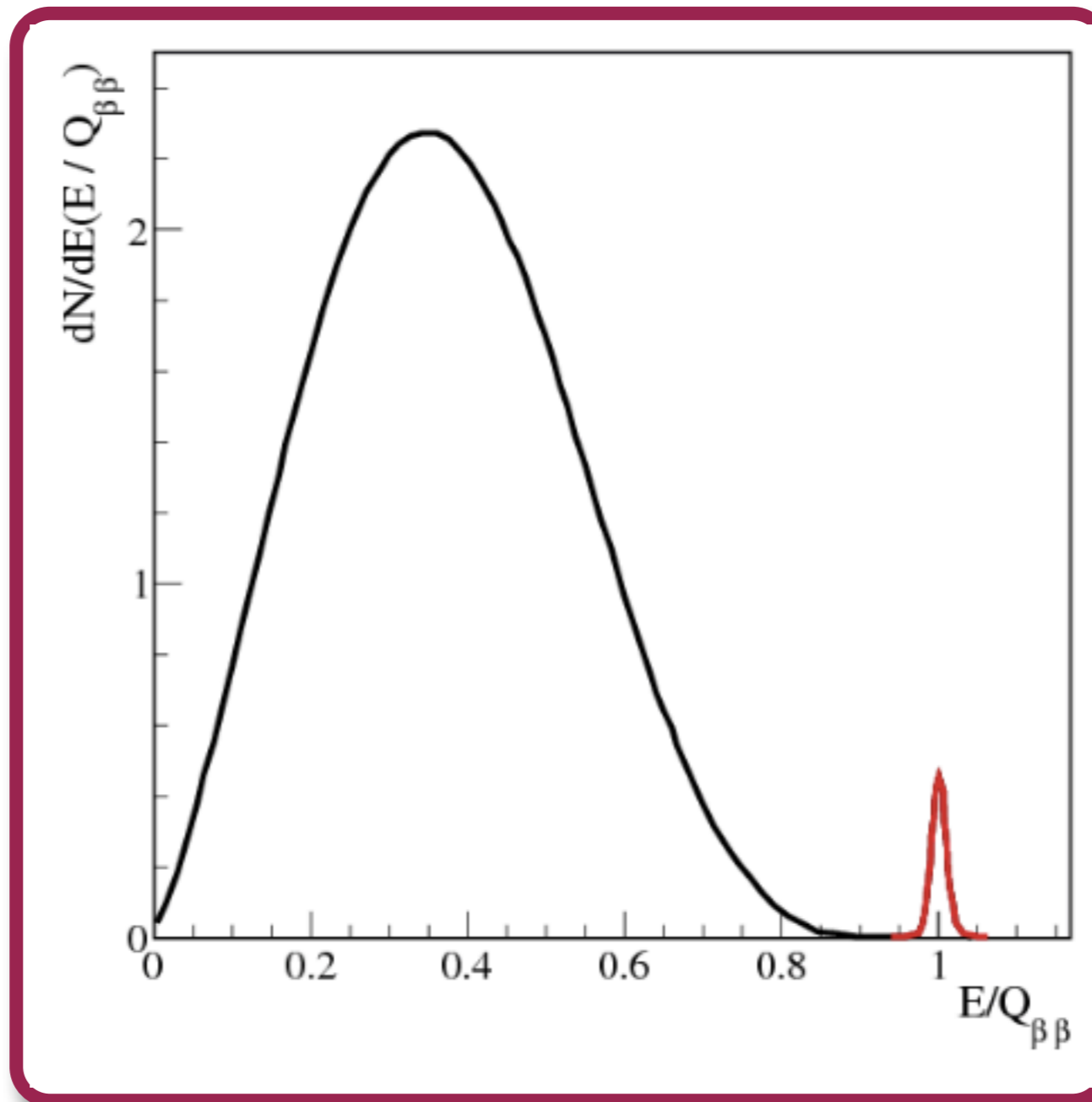


Second order, allowed



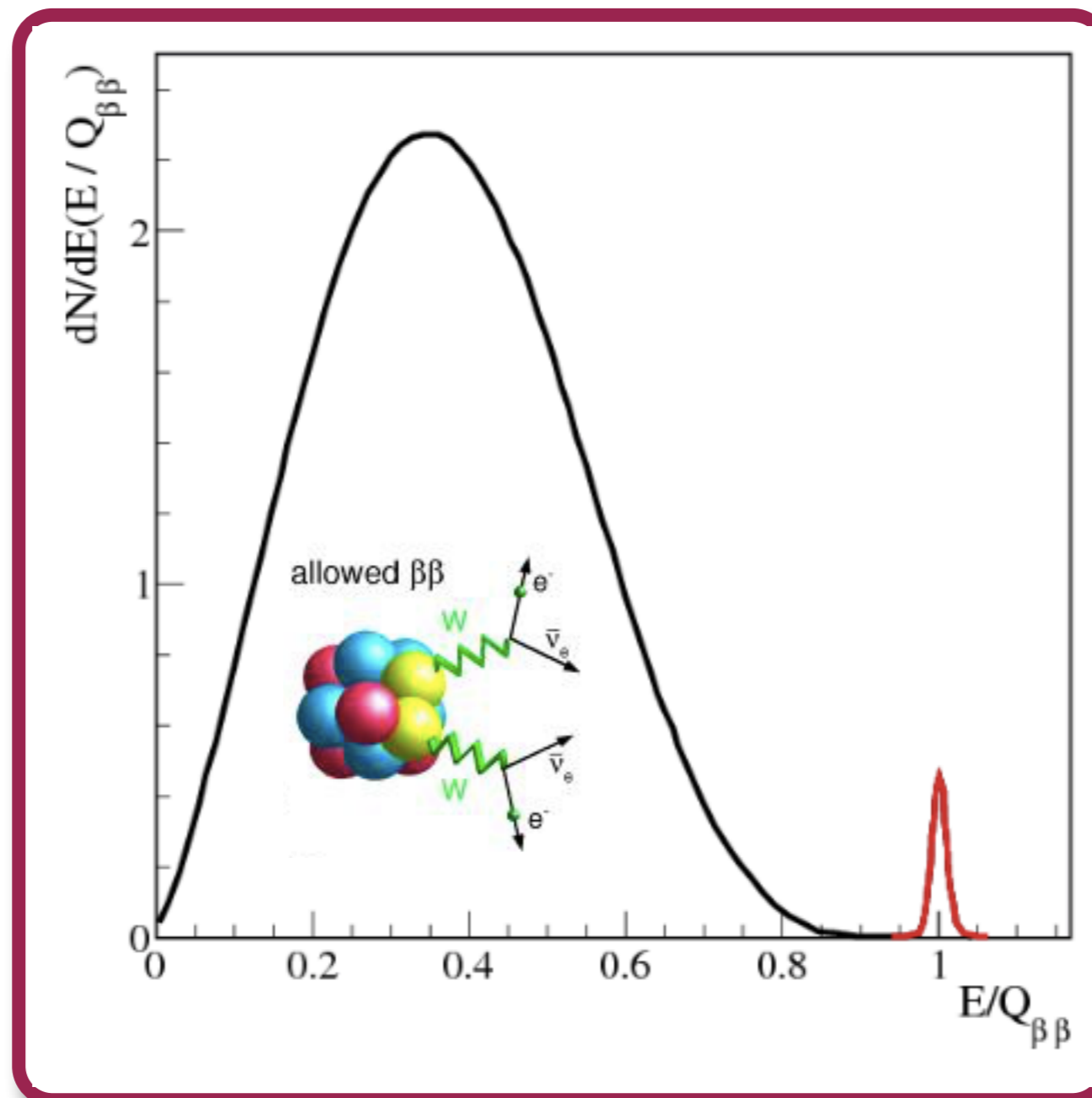
# Experiment

Neutrinoless mode can be isolated using spectroscopic methods



# Experiment

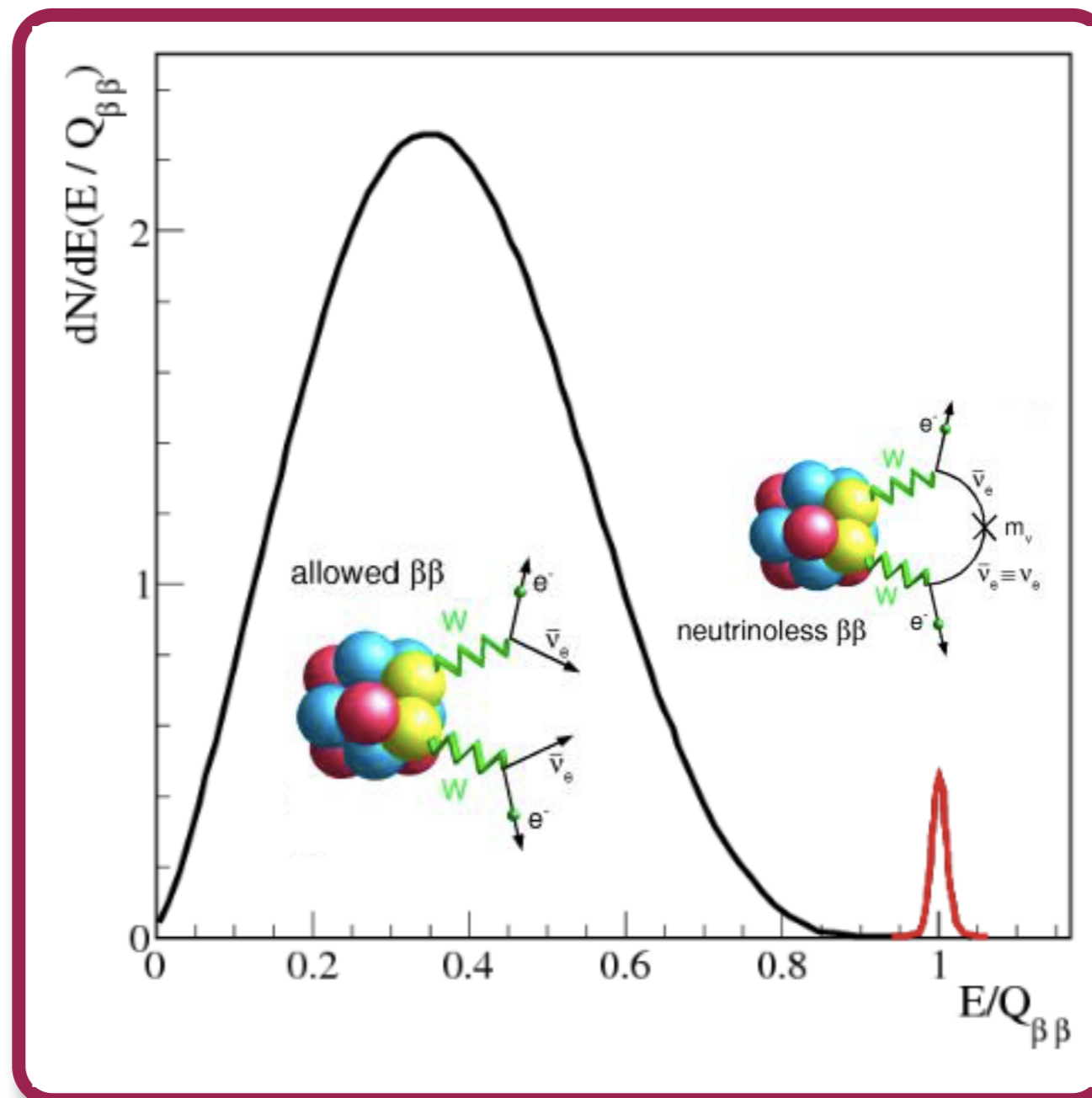
Neutrinoless mode can be isolated using spectroscopic methods





# Experiment

Neutrinoless mode can be isolated using spectroscopic methods



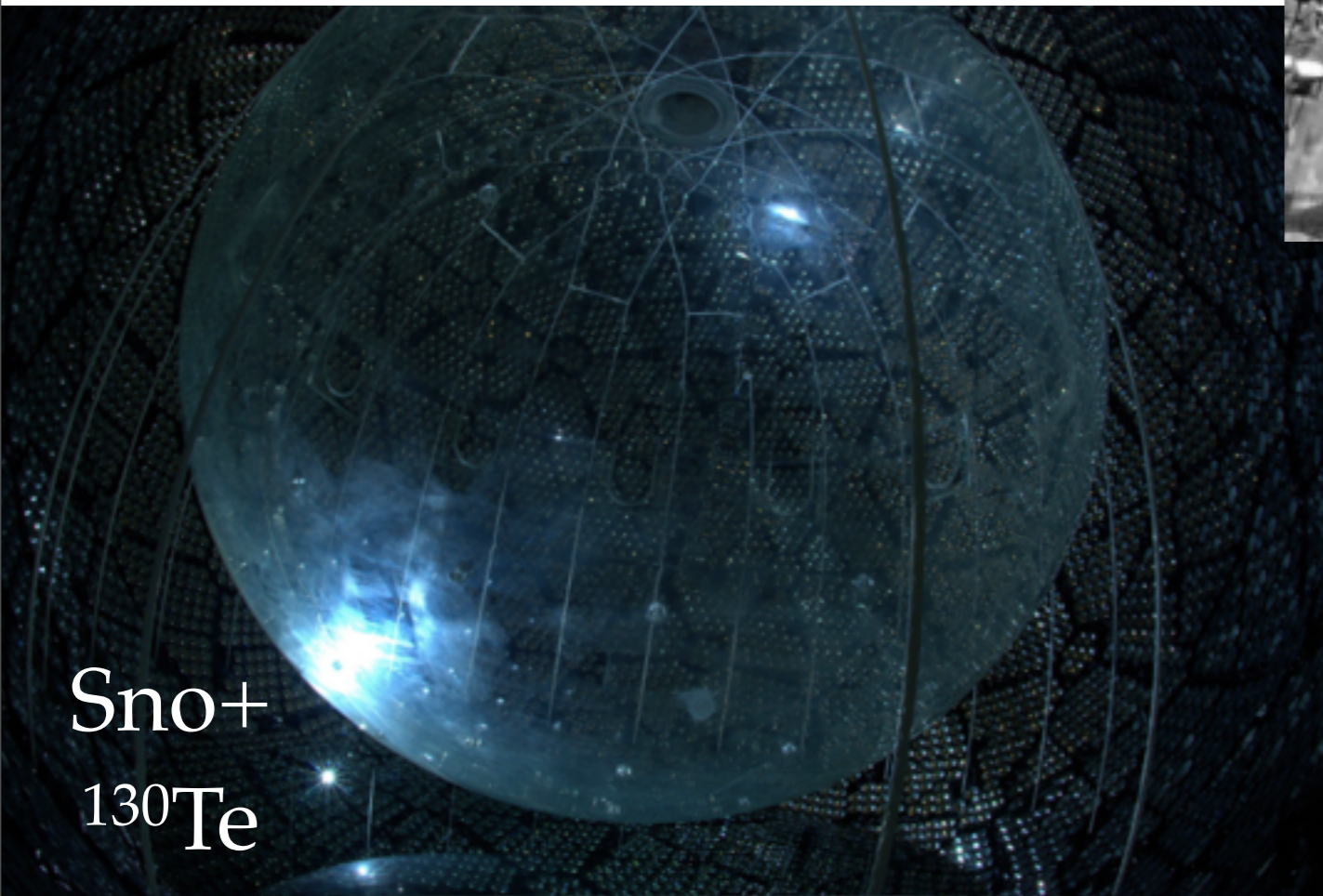
# Experiment



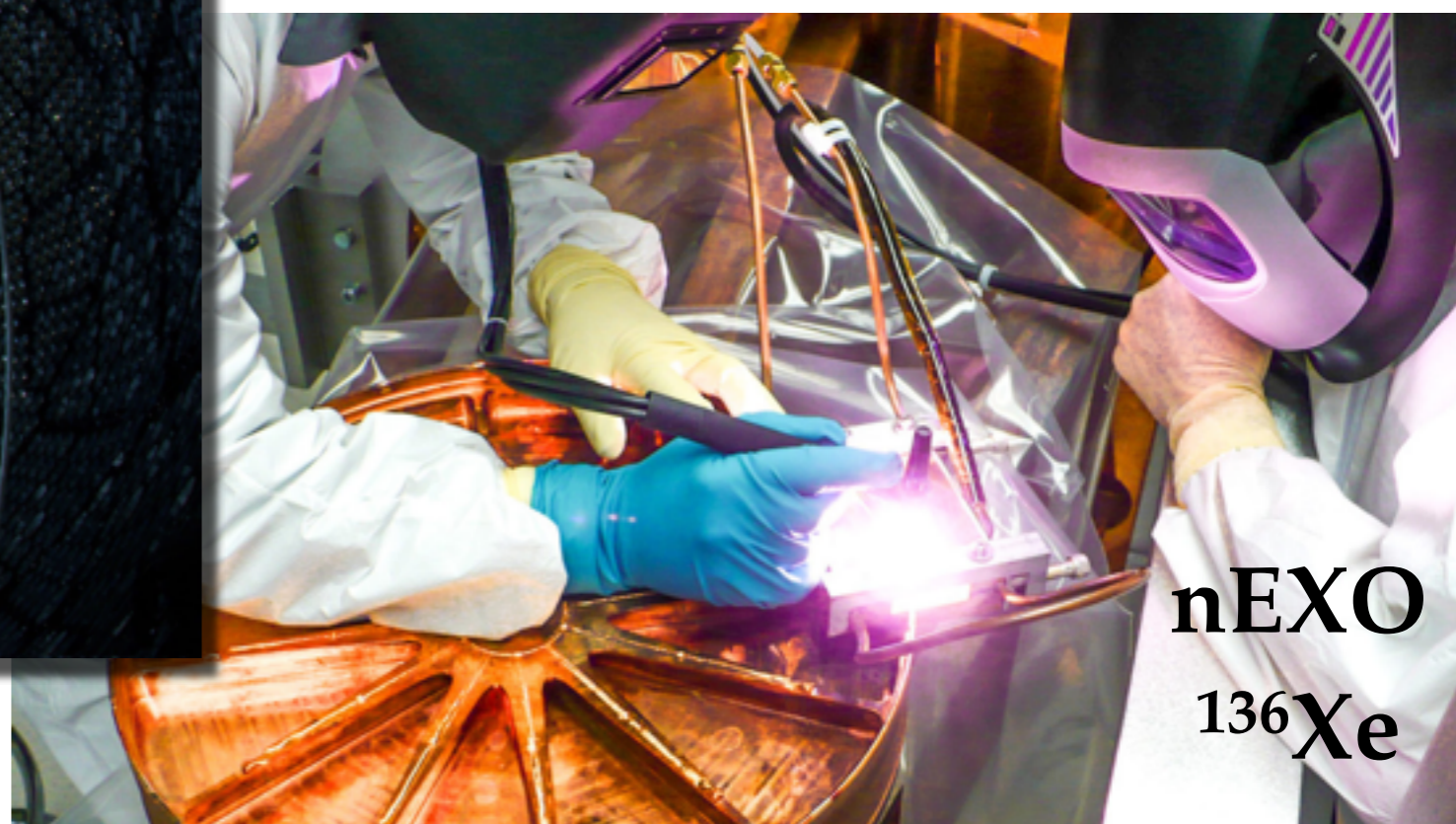
Cuore  
 $^{130}\text{Te}$



Gerda  
 $^{76}\text{Ge}$



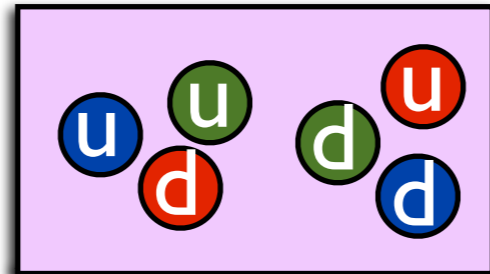
Sno+  
 $^{130}\text{Te}$



nEXO  
 $^{136}\text{Xe}$

How can LQCD  
contribute?

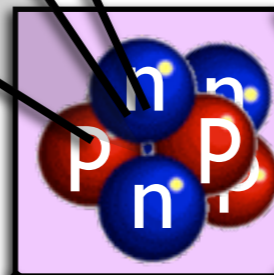
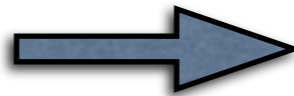
LQCD



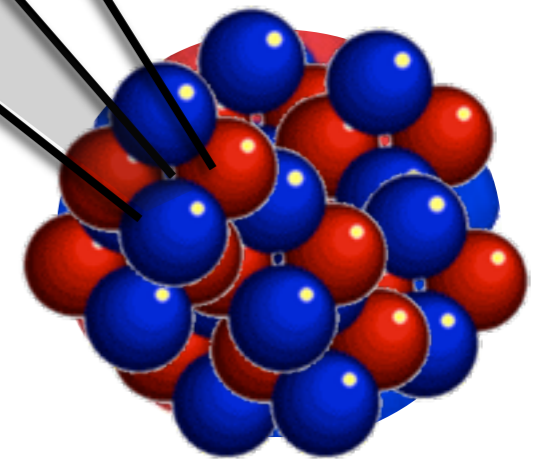
Paradigm



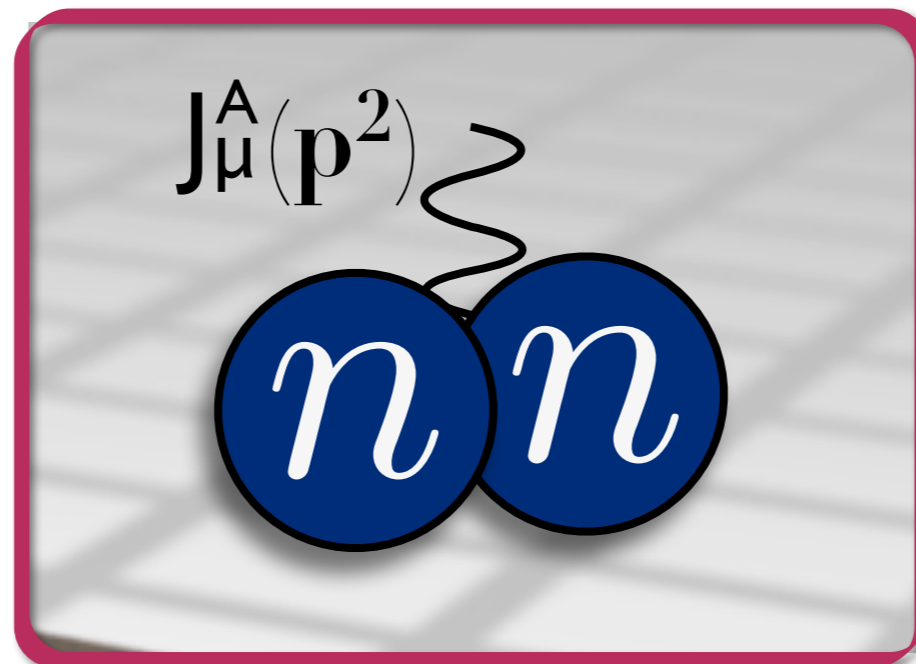
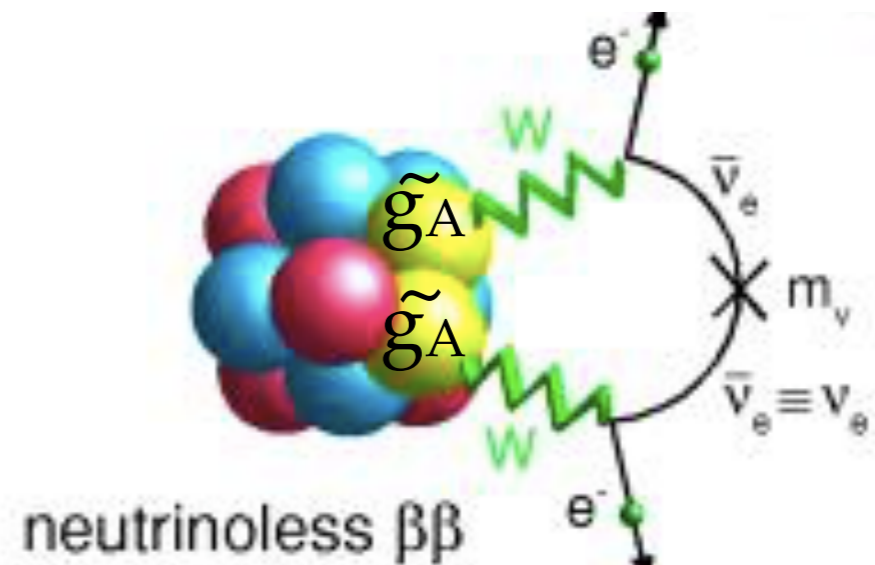
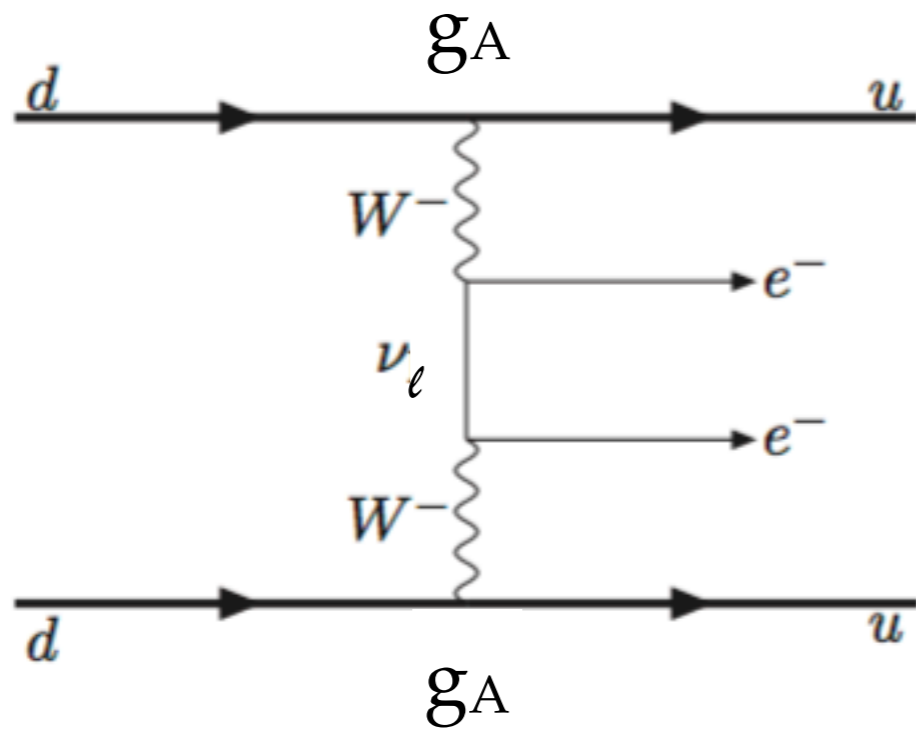
*ab-initio* methods



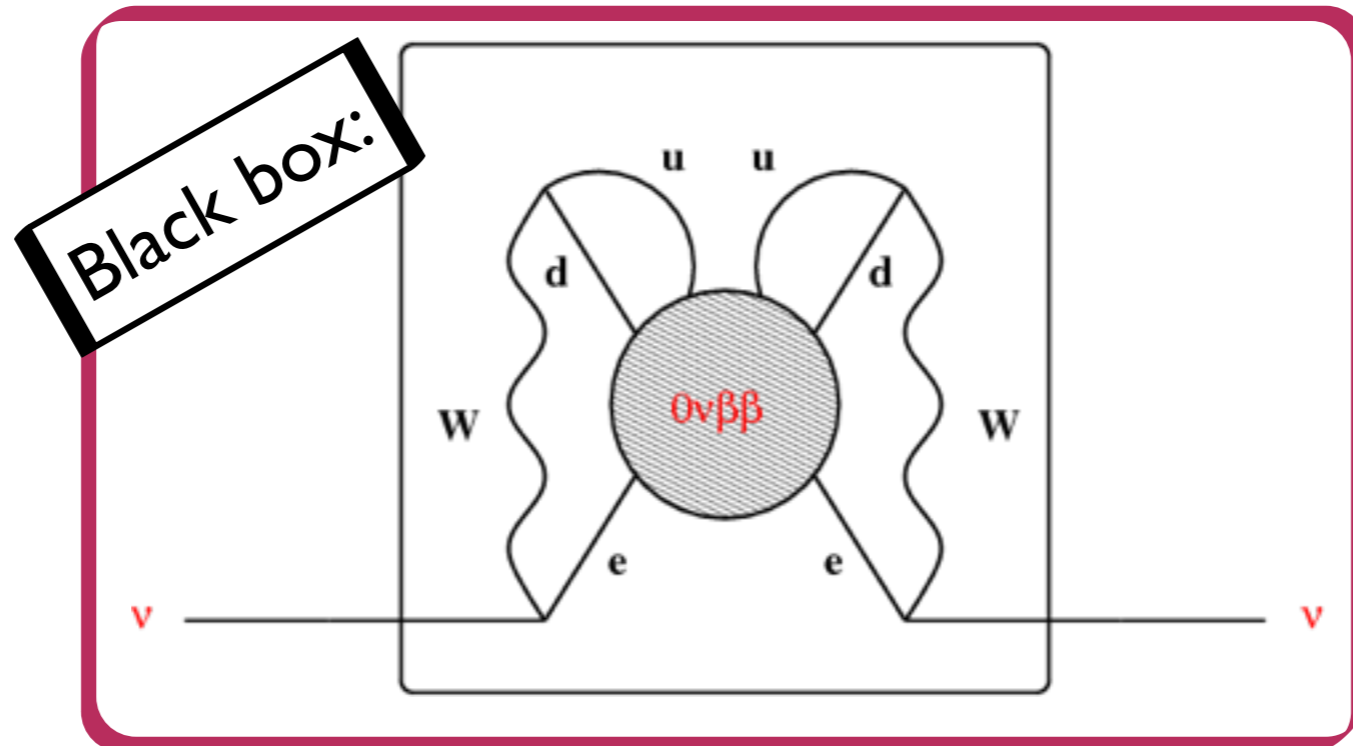
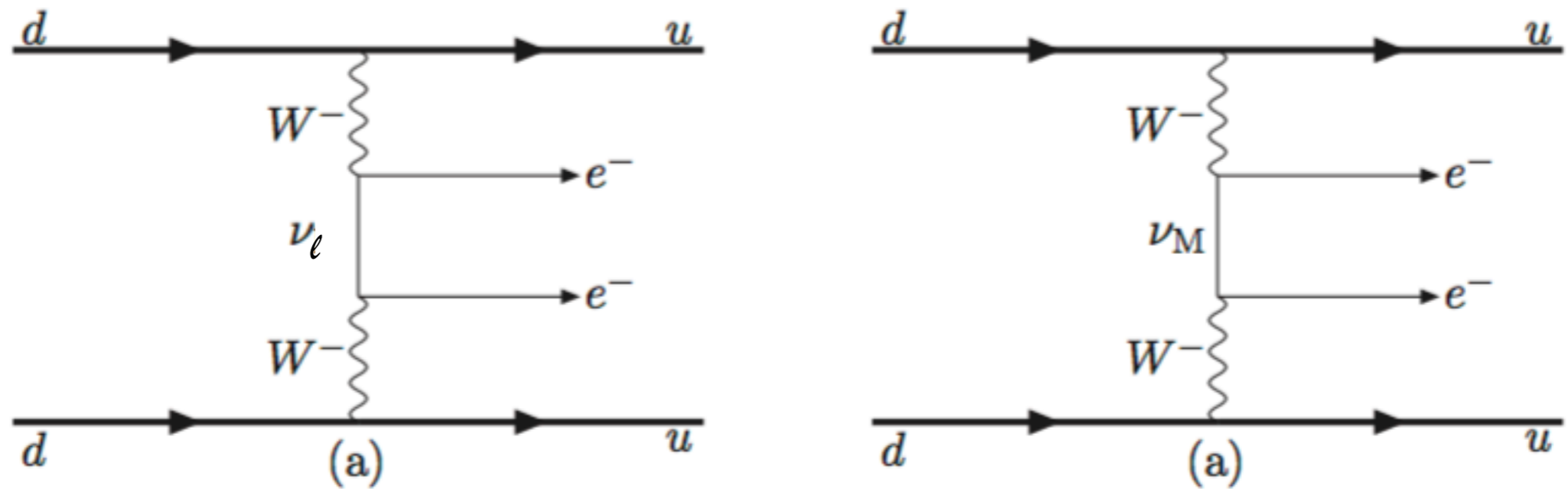
Many-body methods



# Standard picture: long-range contribution

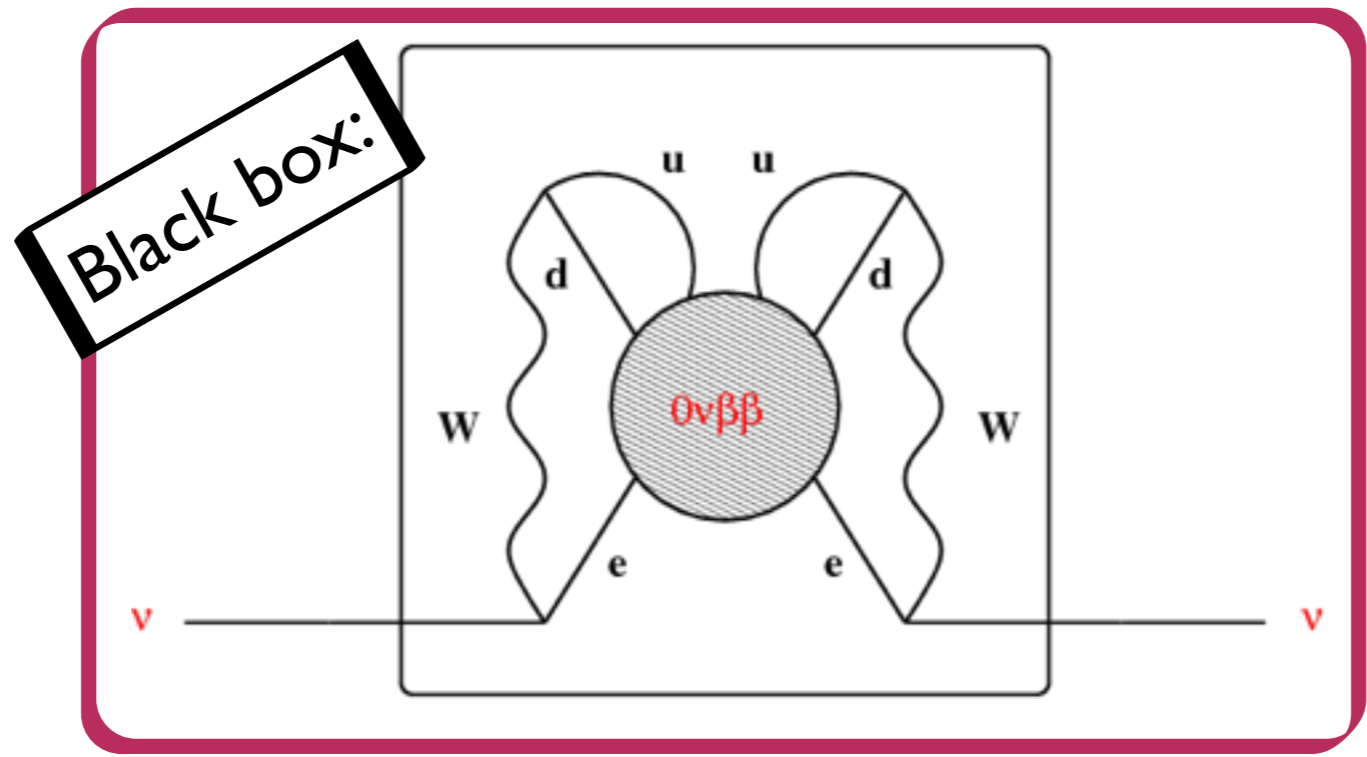
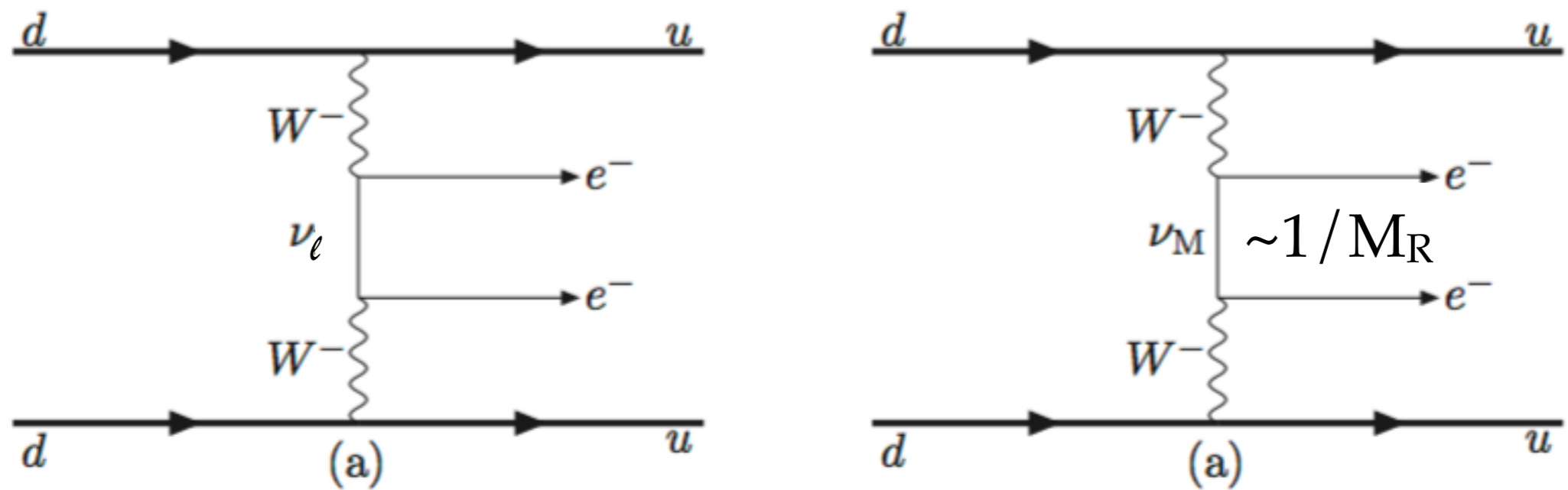


# Short-range contribution: probe for heavy physics



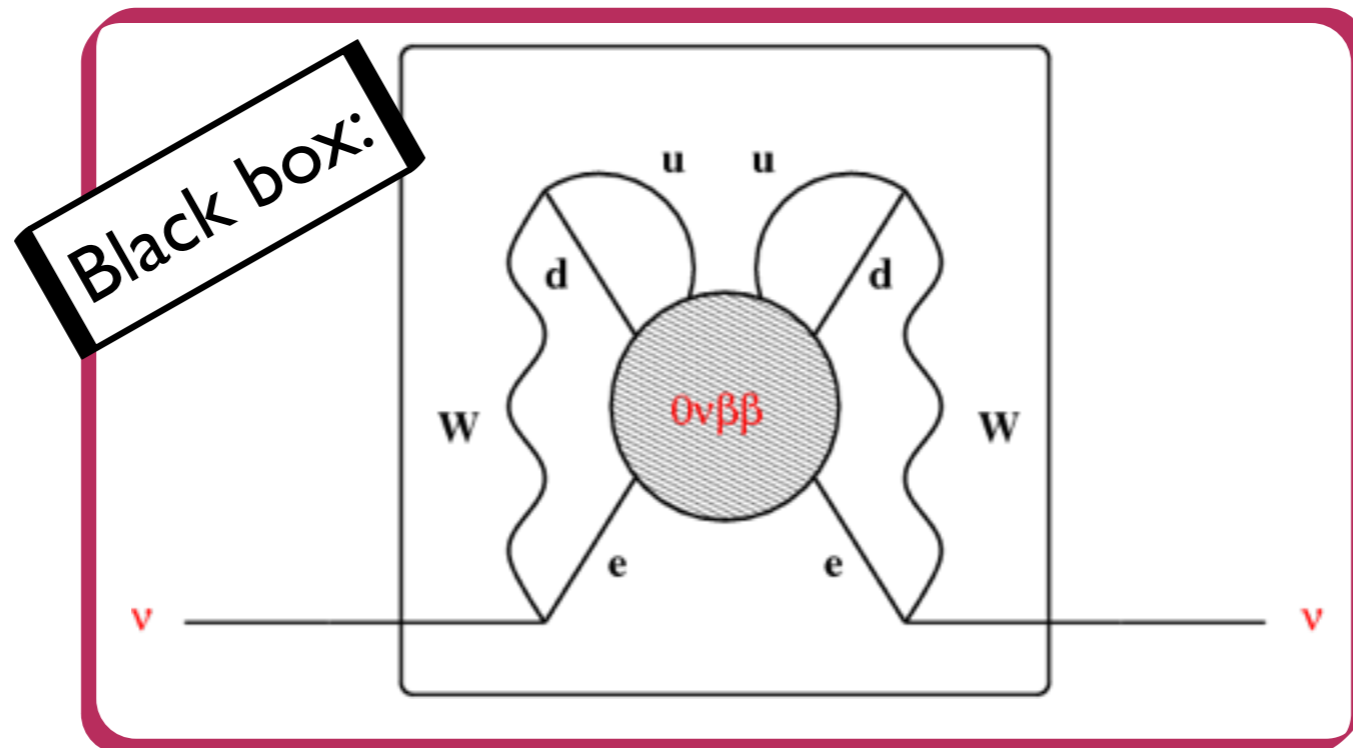
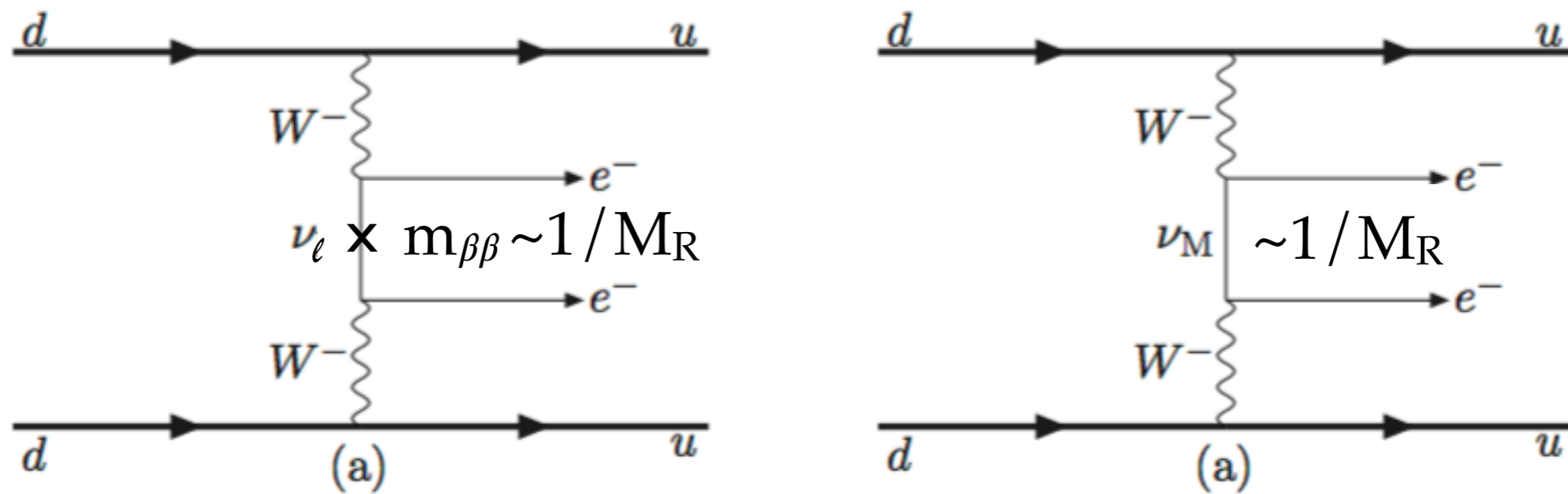
Valle & Schechter, Fig.: H. Päs, W. Rodejohann New J.Phys. 17 (2015) no.11, 115010

# Short-range contribution: probe for heavy physics



Valle & Schechter, Fig.: H. Päs, W. Rodejohann New J.Phys. 17 (2015) no.11, 115010

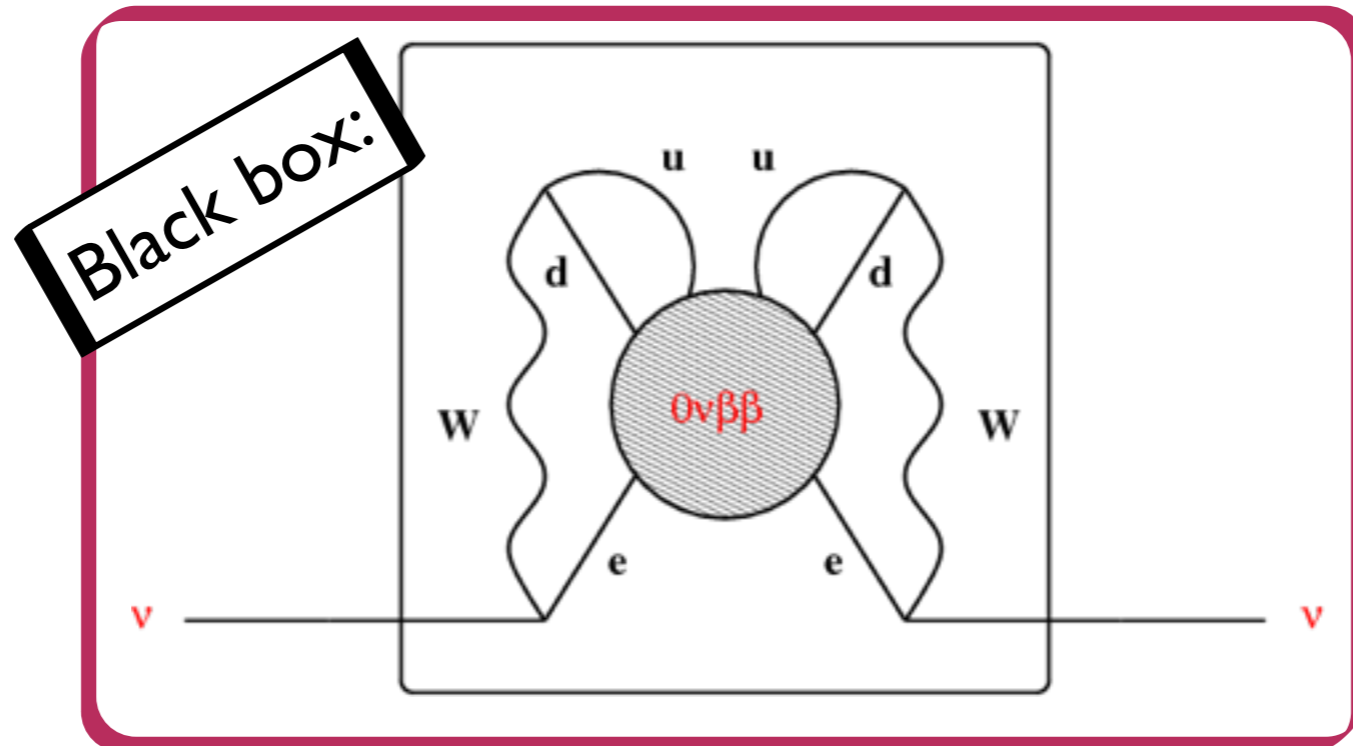
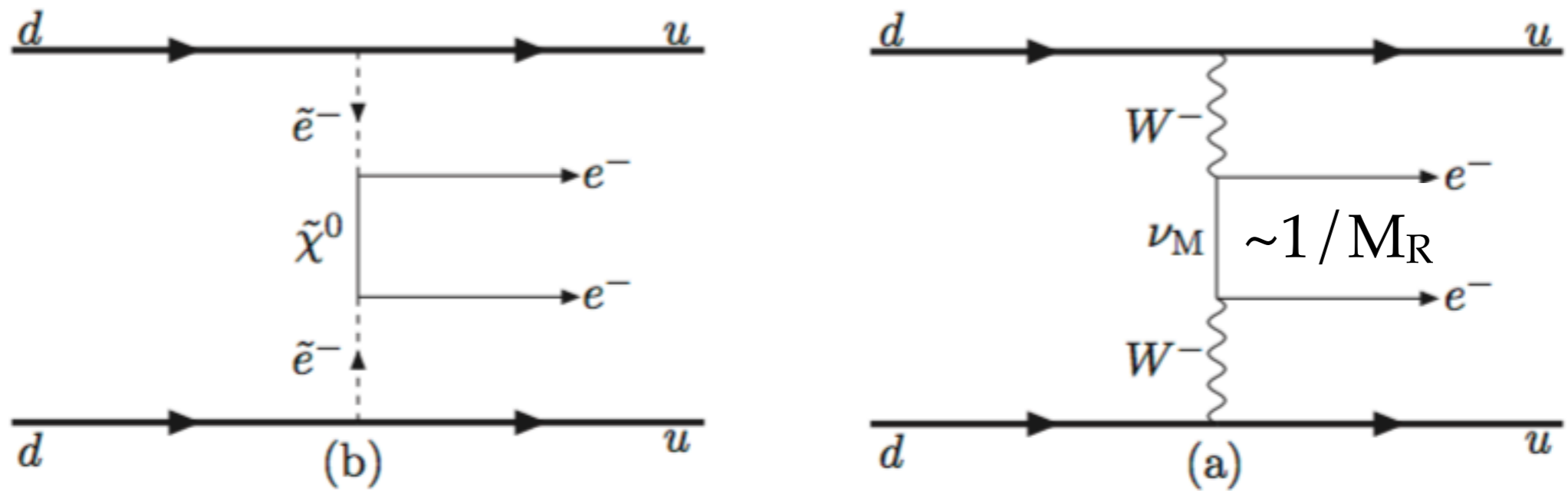
# Short-range contribution: probe for heavy physics



Valle & Schechter, Fig.: H. Päs, W. Rodejohann New J.Phys. 17 (2015) no.11, 115010



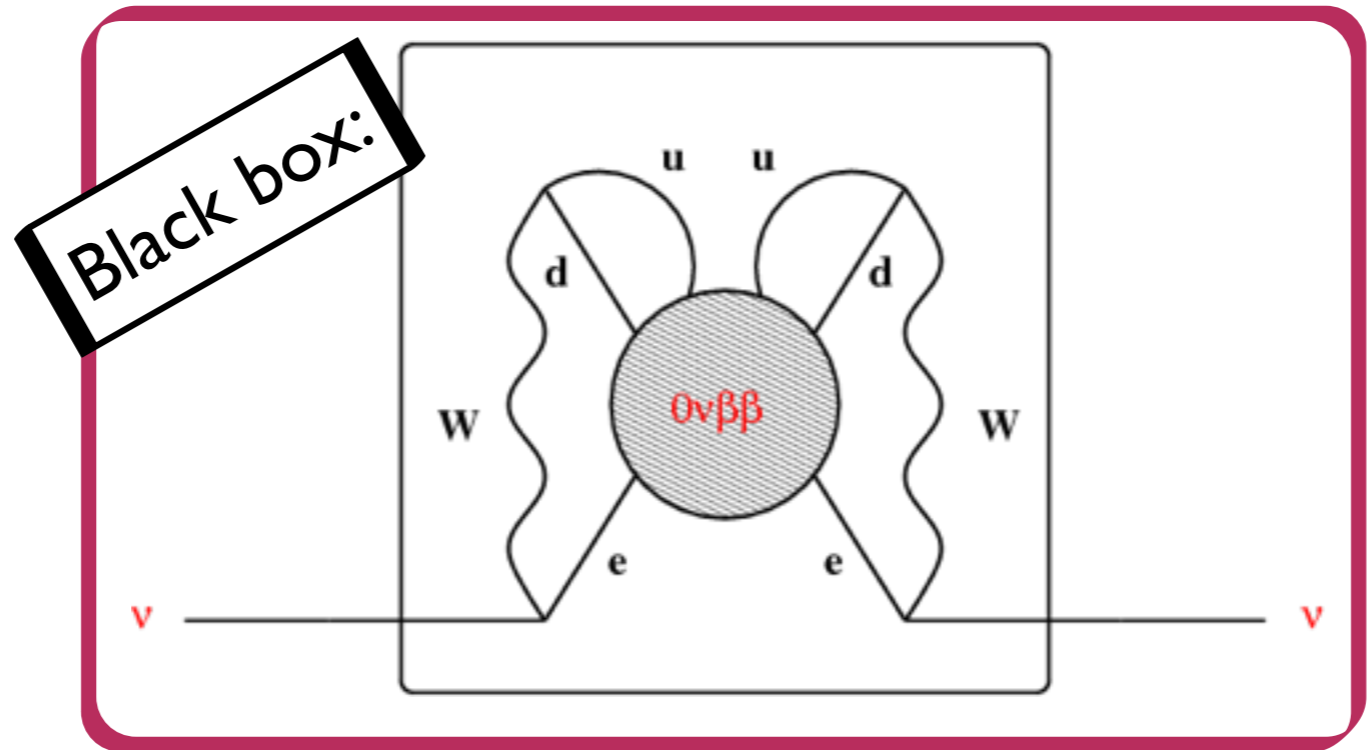
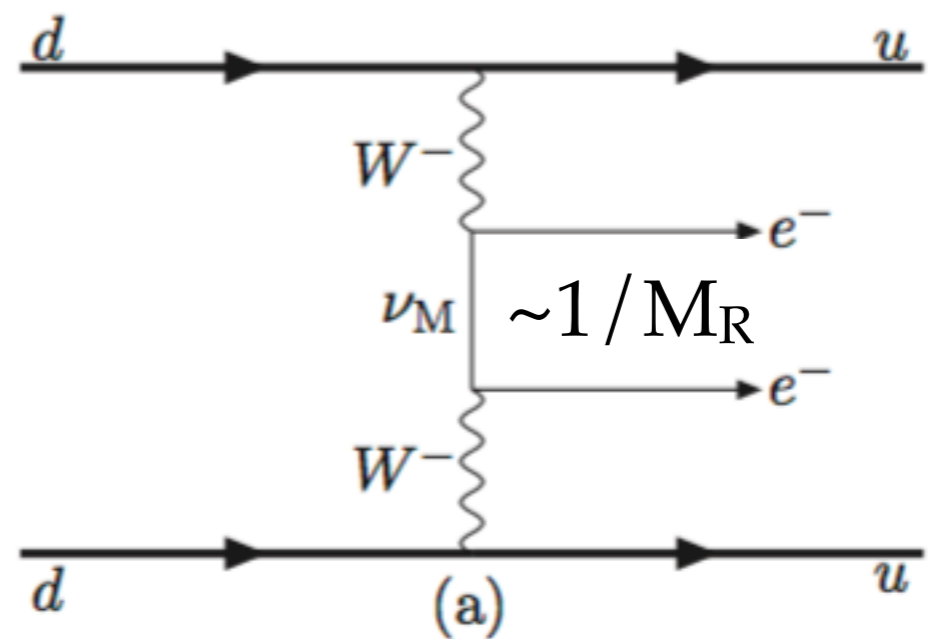
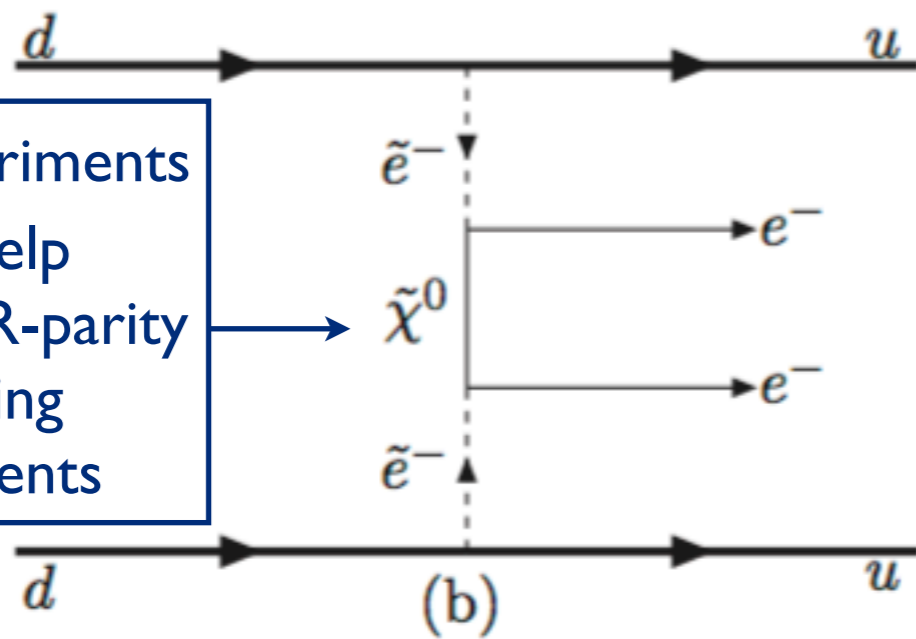
# Short-range contribution: probe for heavy physics



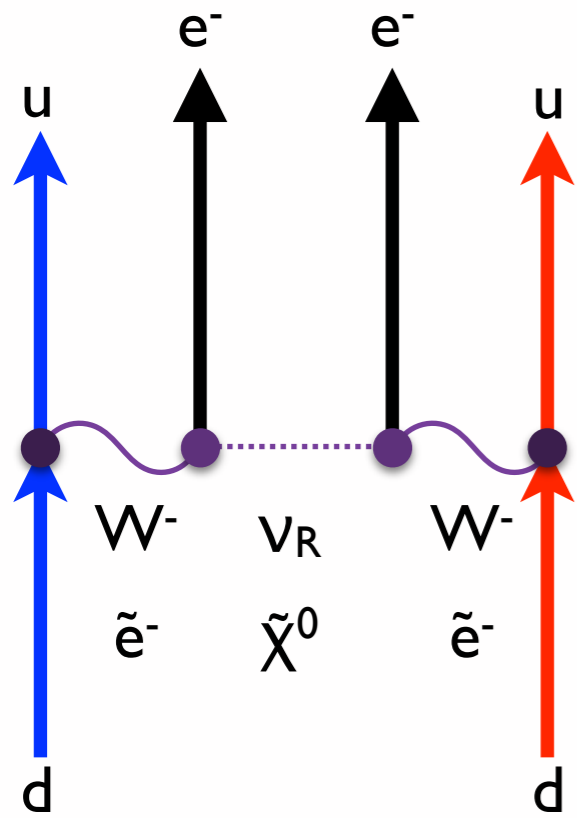
Valle & Schechter, Fig.: H. Päs, W. Rodejohann New J.Phys. 17 (2015) no.11, 115010

# Short-range contribution: probe for heavy physics

$0\nu\beta\beta$  experiments may help constrain R-parity violating coefficients

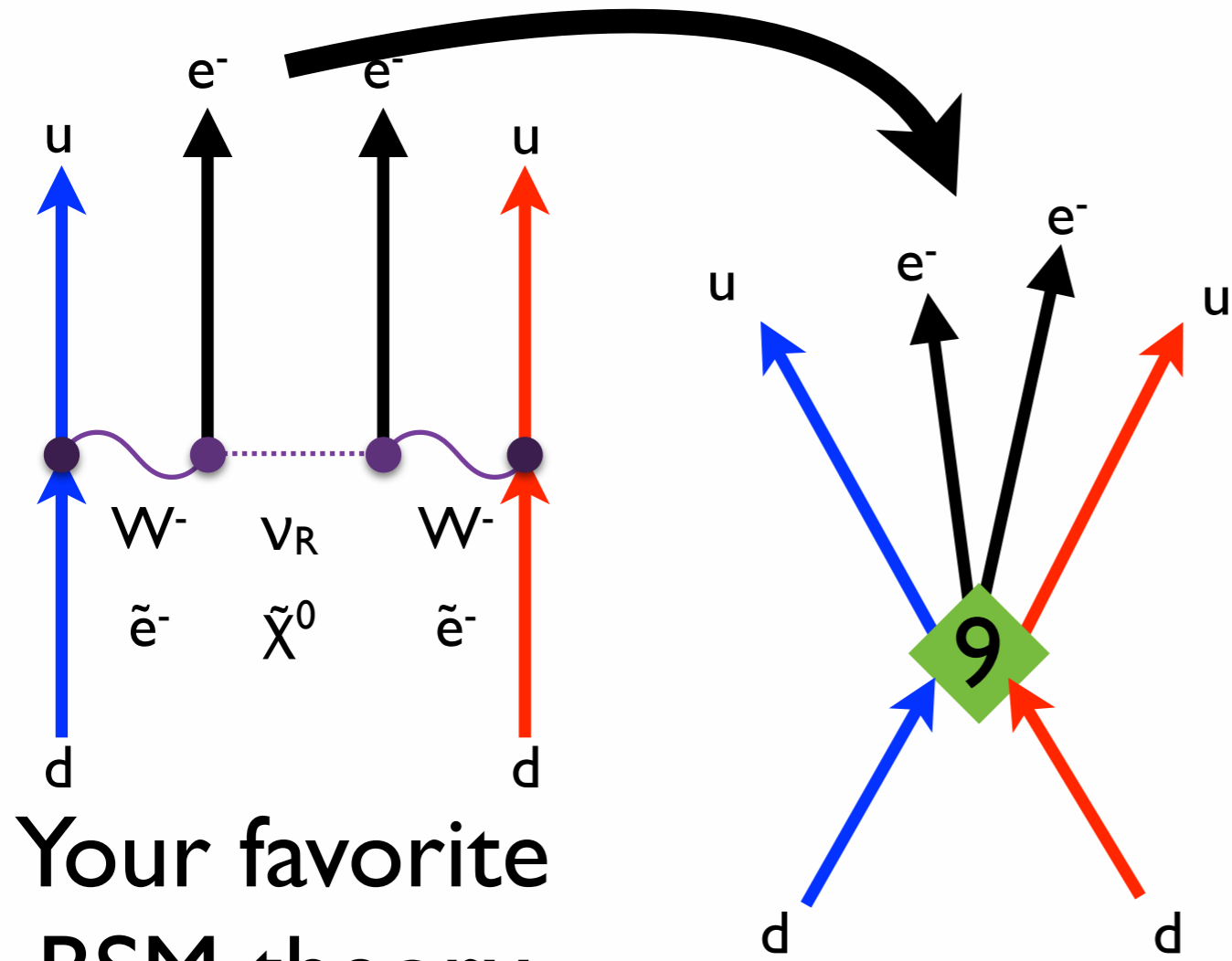


Valle & Schechter, Fig.: H. Päs, W. Rodejohann New J.Phys. 17 (2015) no.11, 115010



Your favorite  
BSM theory

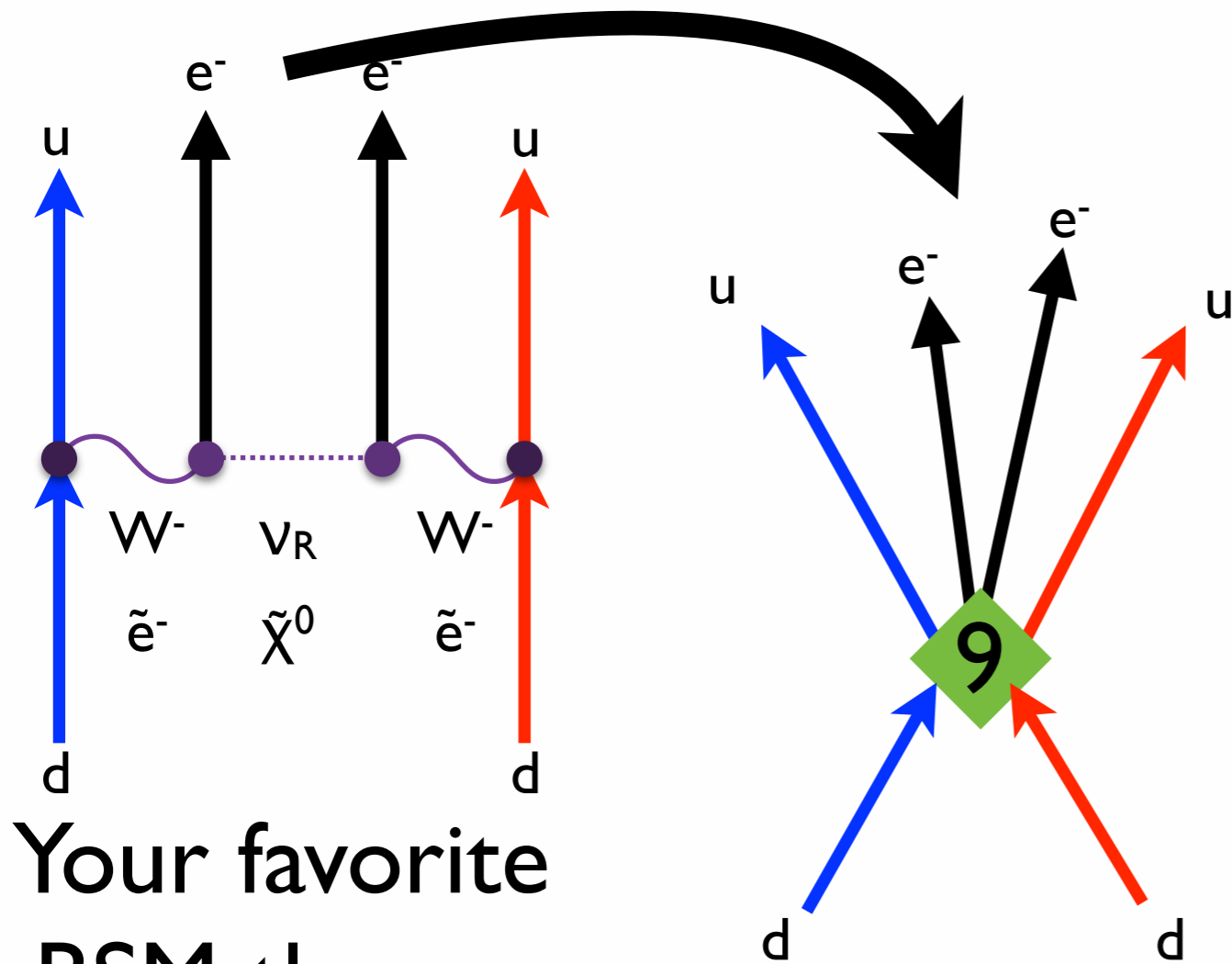
# Pen & Paper



Your favorite  
BSM theory

effective 4-quark  
operators in QCD

# Pen & Paper



Your favorite  
BSM theory

op

$$\mathcal{O}_{1+}^{ab} = (\bar{q}_L \tau^a \gamma^\mu q_L) (\bar{q}_R \tau^b \gamma_\mu q_R),$$

$$\mathcal{O}_{2\pm}^{ab} = (\bar{q}_R \tau^a q_L) (\bar{q}_R \tau^b q_L) \pm (\bar{q}_L \tau^a q_R) (\bar{q}_L \tau^b q_R),$$

$$\mathcal{O}_{3\pm}^{ab} = (\bar{q}_L \tau^a \gamma^\mu q_L) (\bar{q}_L \tau^b \gamma_\mu q_L) \pm (\bar{q}_R \tau^a \gamma^\mu q_R) (\bar{q}_R \tau^b \gamma_\mu q_R),$$

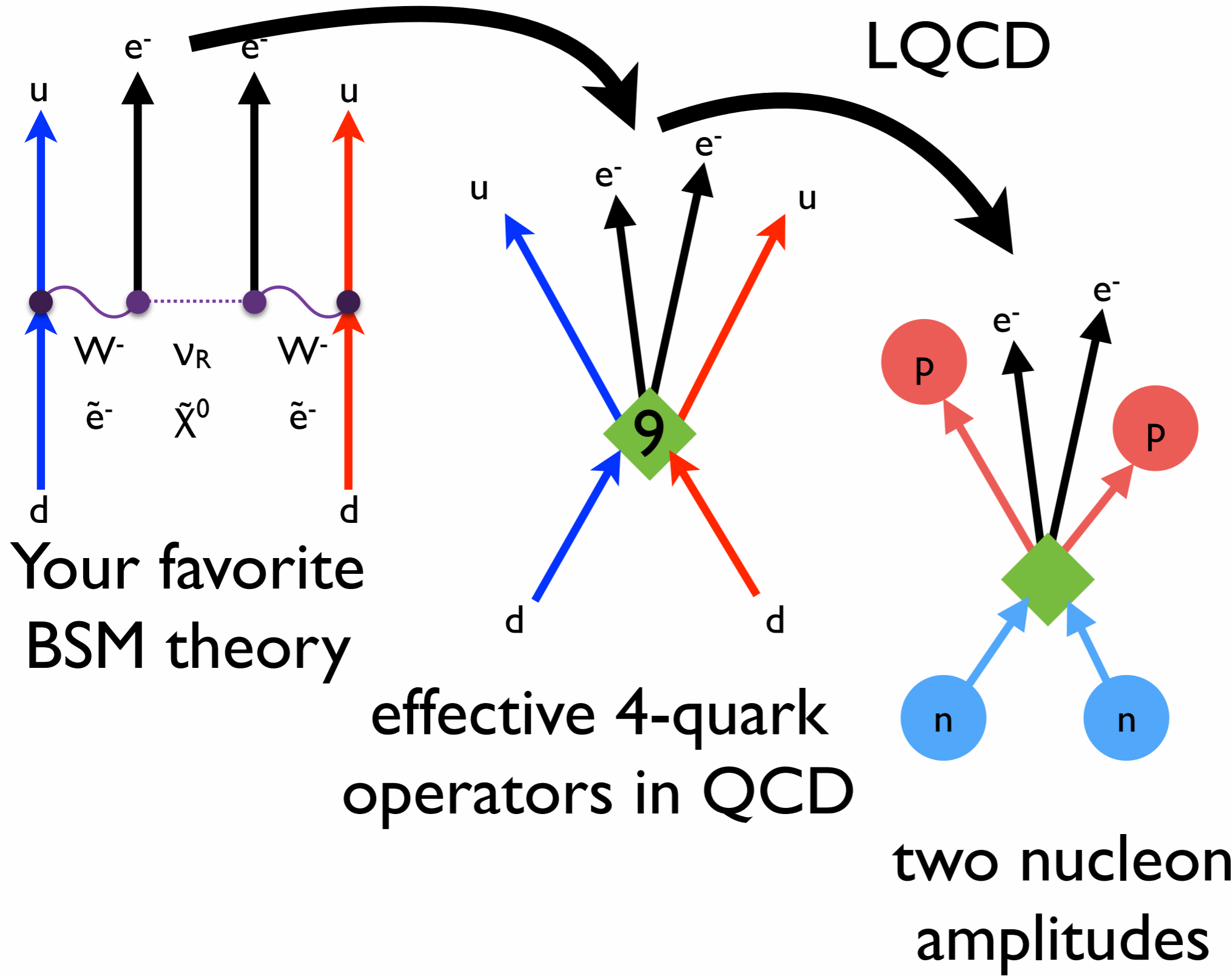
$$\mathcal{O}_{4\pm}^{ab,\mu} = (\bar{q}_L \tau^a \gamma^\mu q_L \mp \bar{q}_R \tau^a \gamma^\mu q_R) (\bar{q}_L \tau^b q_R - \bar{q}_R \tau^b q_L),$$

$$\mathcal{O}_{5\pm}^{ab,\mu} = (\bar{q}_L \tau^a \gamma^\mu q_L \pm \bar{q}_R \tau^a \gamma^\mu q_R) (\bar{q}_L \tau^b q_R + \bar{q}_R \tau^b q_L).$$

Prezeau, Ramsey-Musolf,  
Vogel (2003)

# Pen & Paper

# LQCD

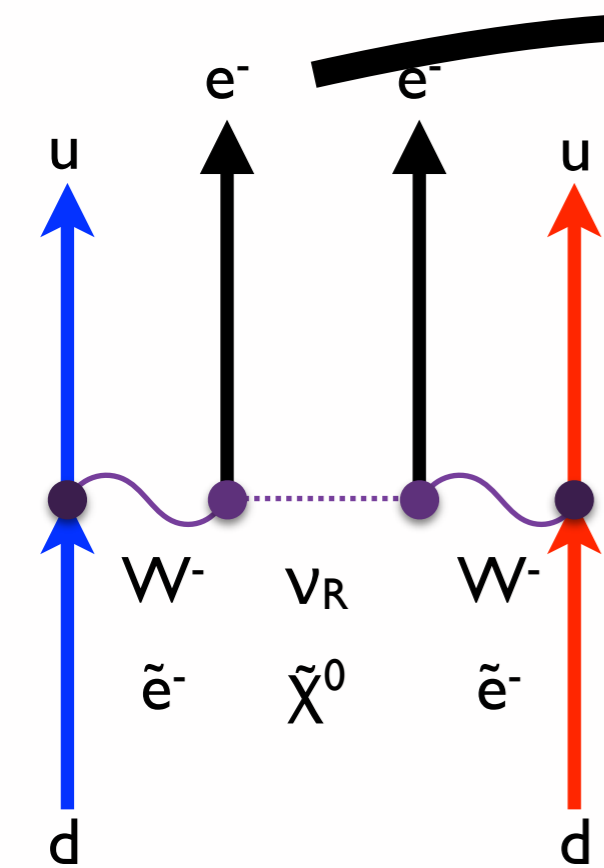


Your favorite  
BSM theory

effective 4-quark  
operators in QCD

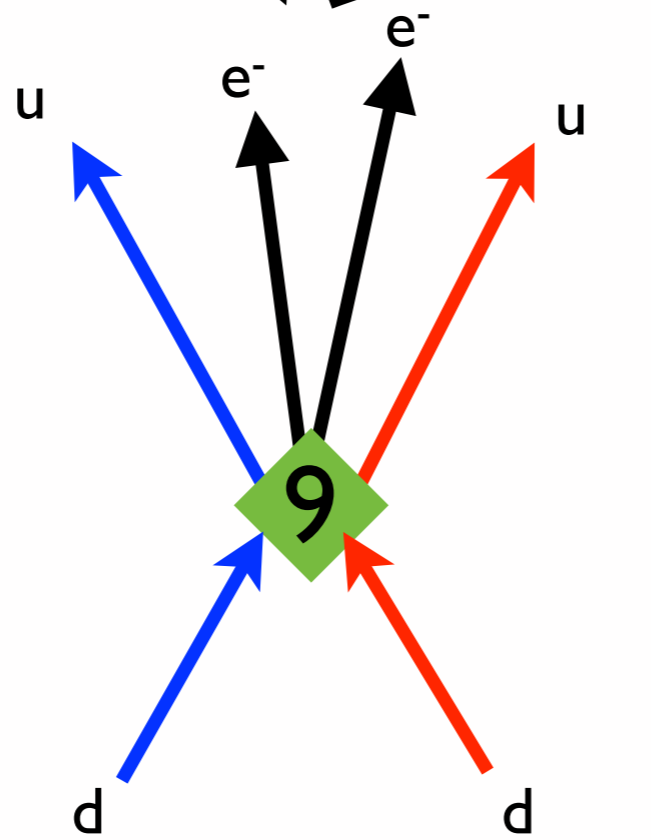
two nucleon  
amplitudes

# Pen & Paper



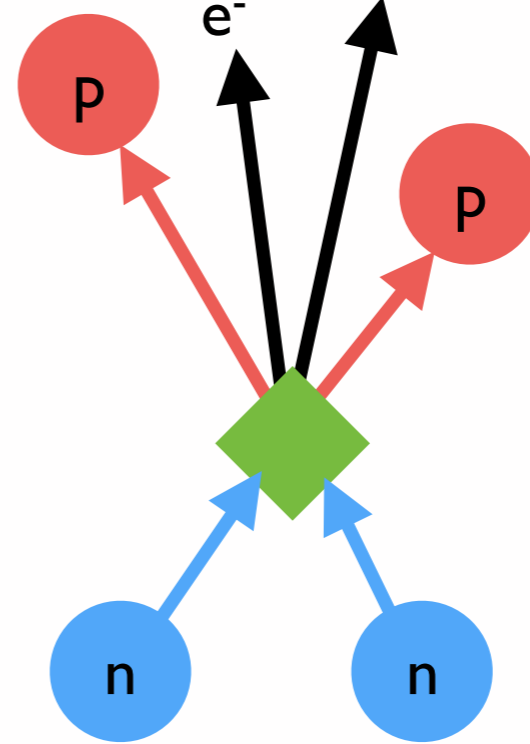
Your favorite  
BSM theory

# LQCD

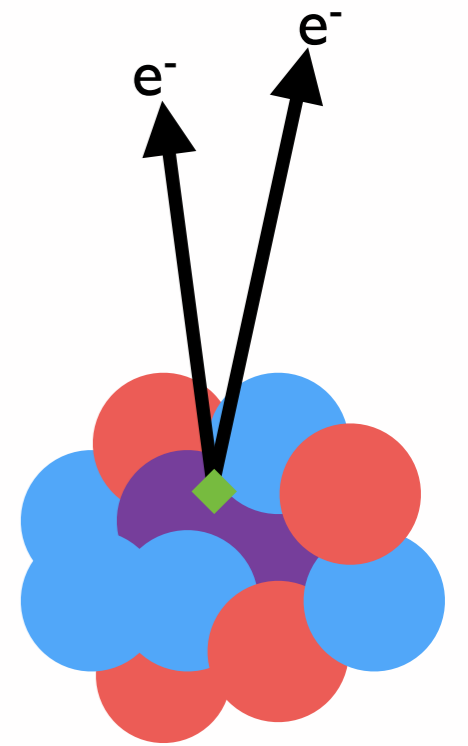


effective 4-quark  
operators in QCD

# NP

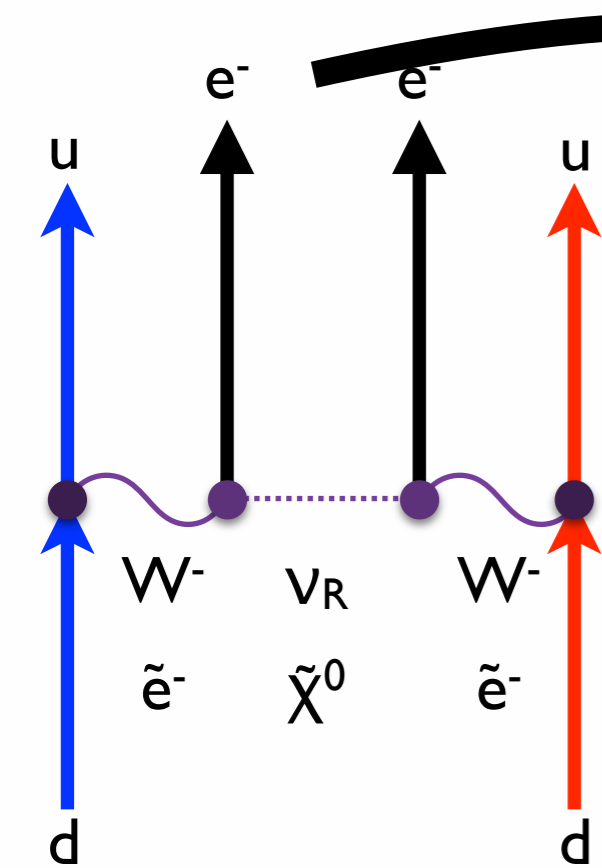


two nucleon  
amplitudes



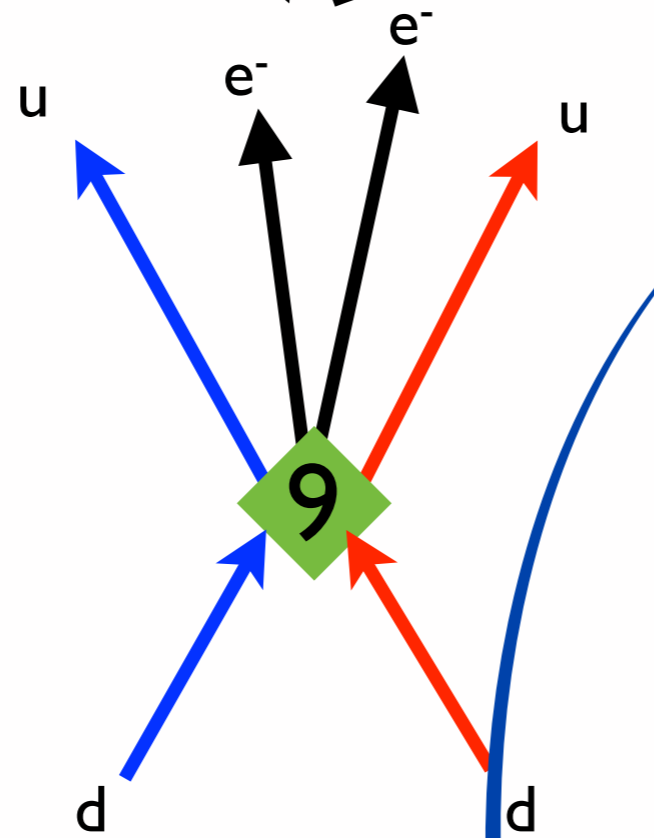
decay rates  
in big nuclei

# Pen & Paper



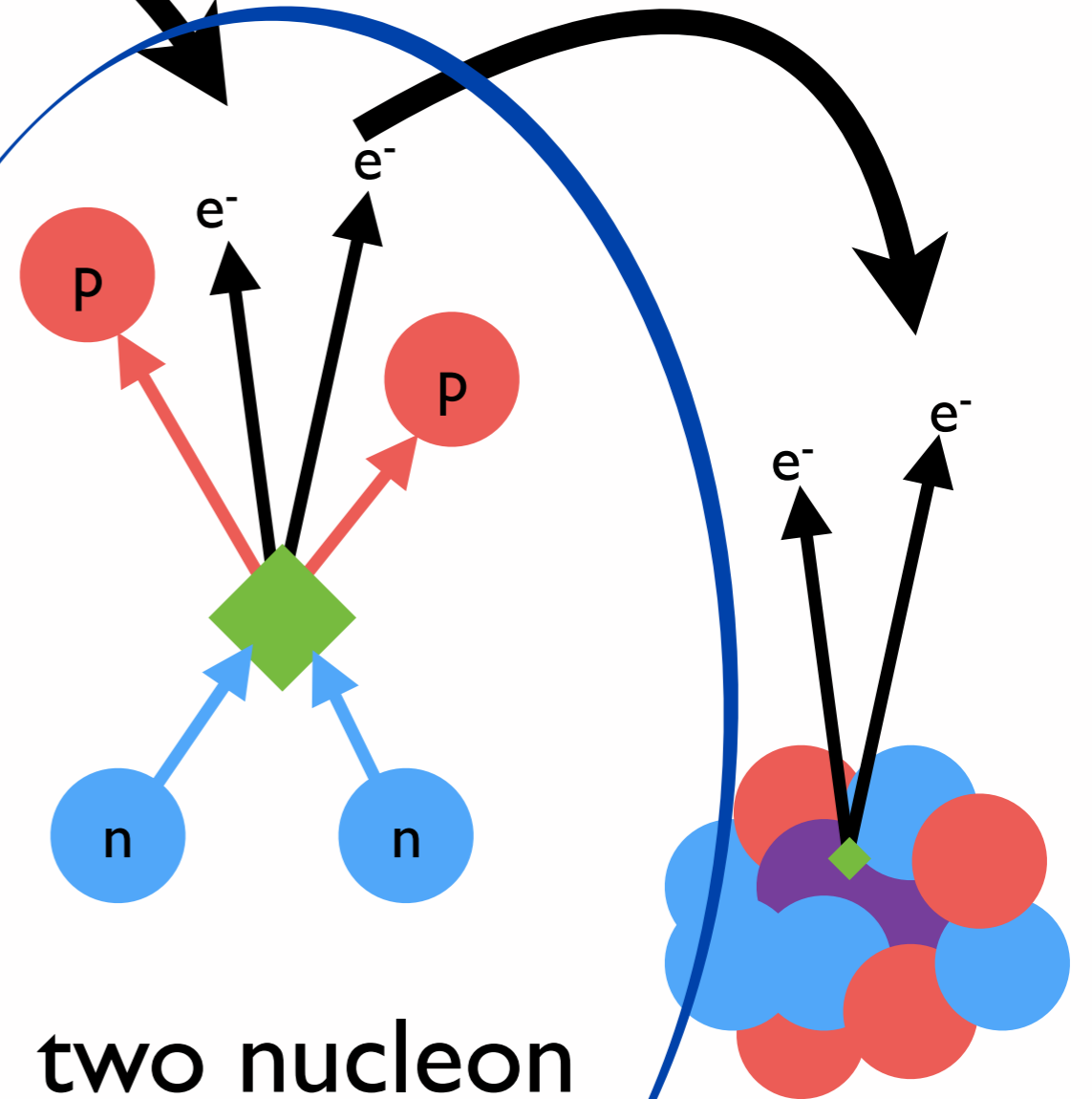
Your favorite  
BSM theory

# LQCD



effective 4-quark  
operators in QCD

# NP



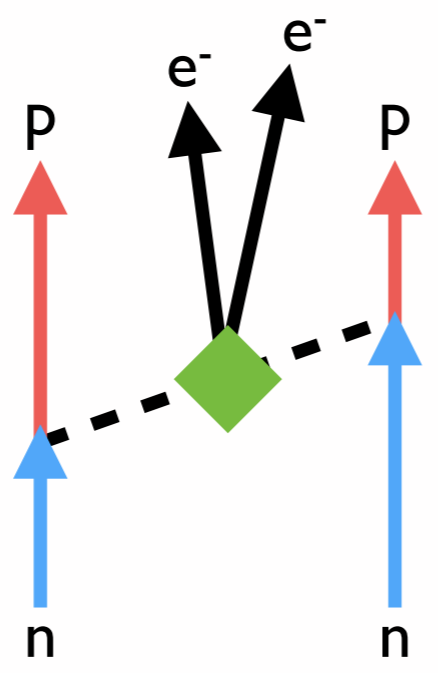
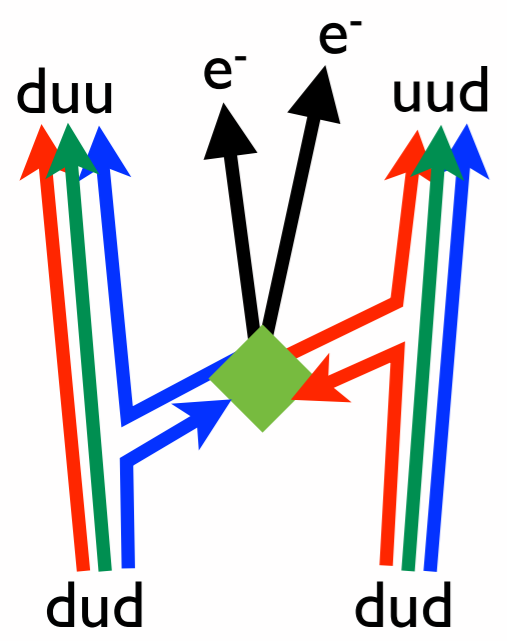
two nucleon  
amplitudes

decay rates  
in big nuclei

Can determine power-counting  
of various contributions using  
chiral perturbation theory

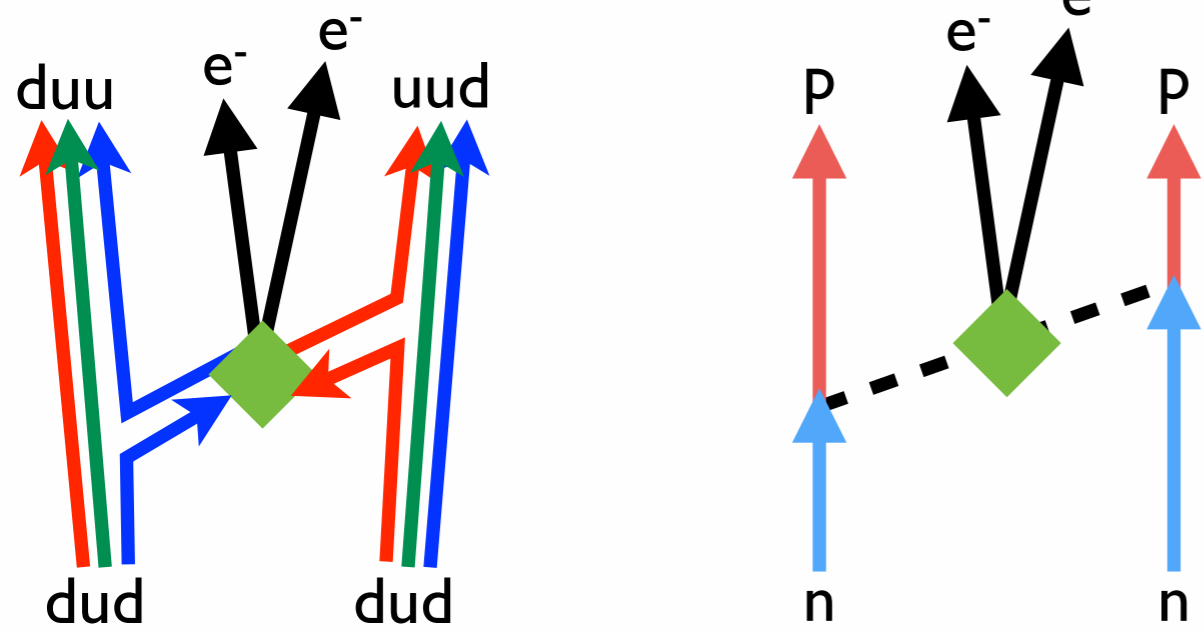


# $\chi$ PT power-counting of operators



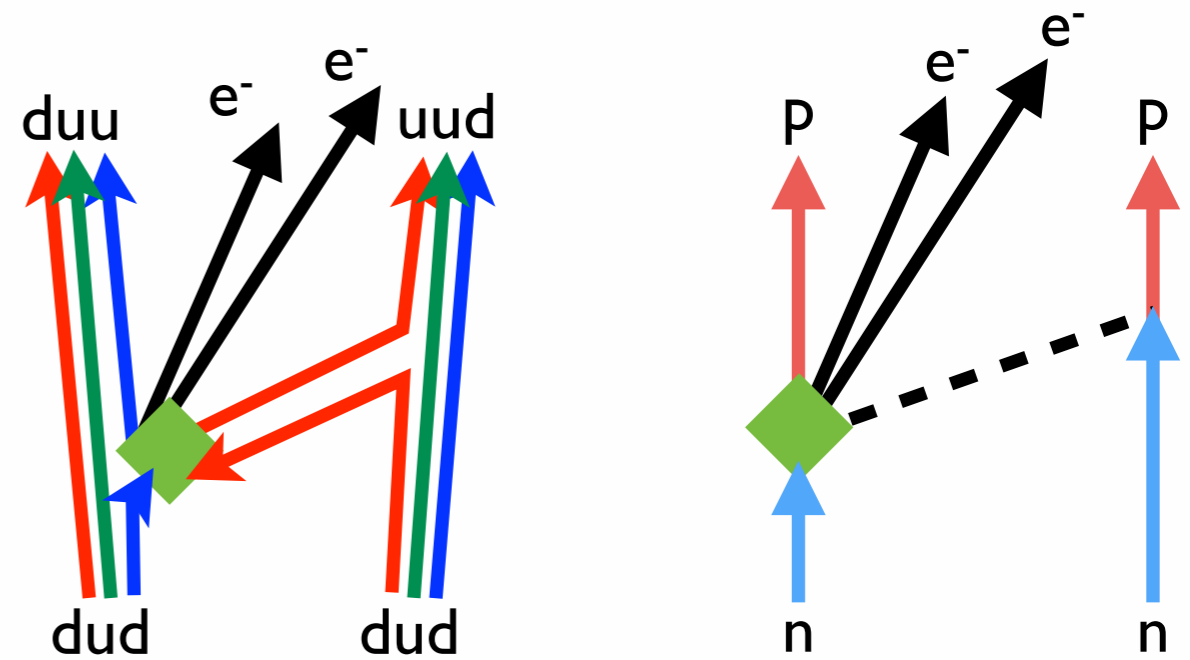
$\mathcal{O}(p^{-2})$  long-range  $\pi$  exchange

# $\chi$ PT power-counting of operators



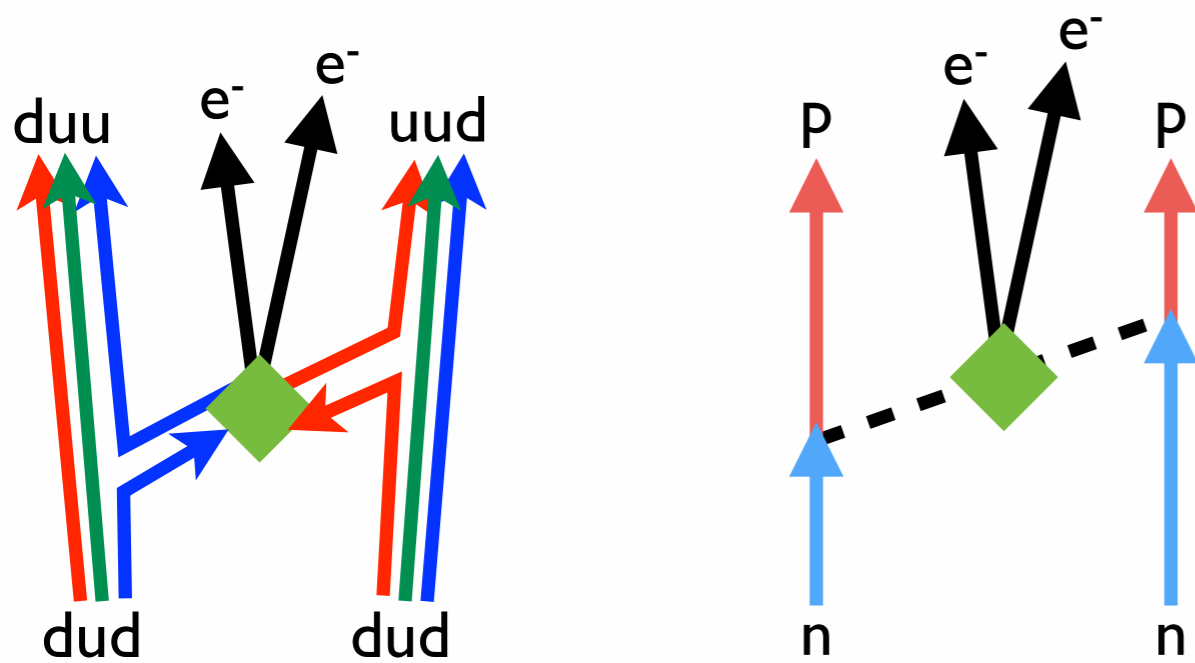
$\mathcal{O}(p^{-2})$  long-range  $\pi$  exchange

$\mathcal{O}(p^{-1})$   $\pi N$  vertex

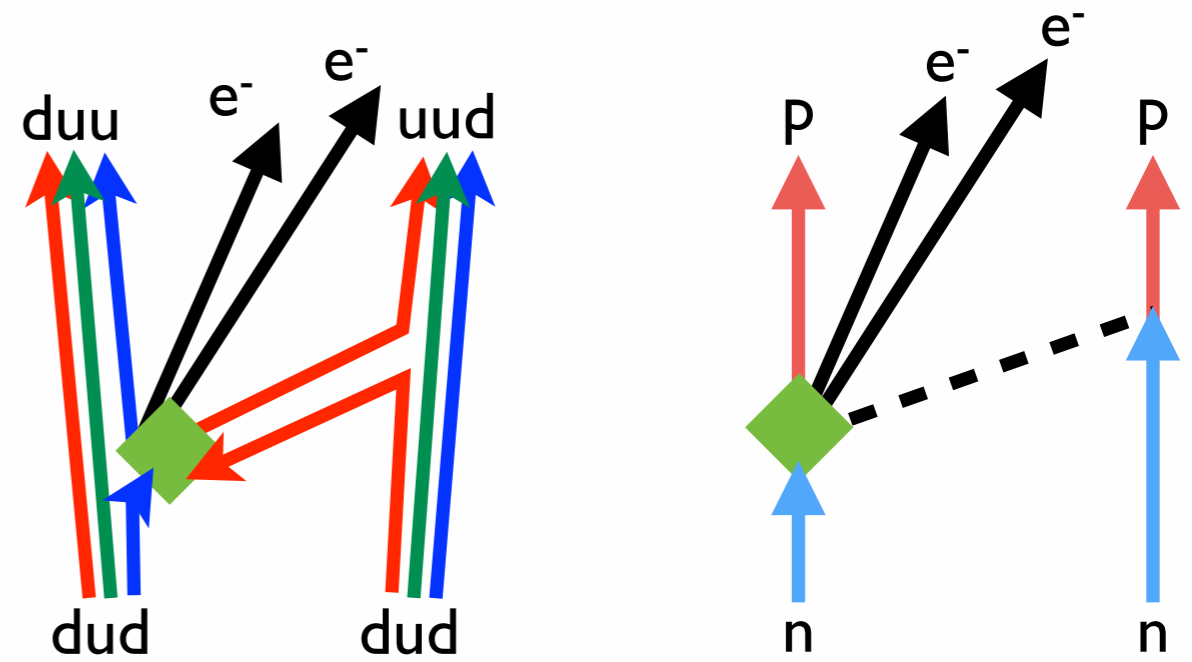


# $\chi$ PT power-counting of operators

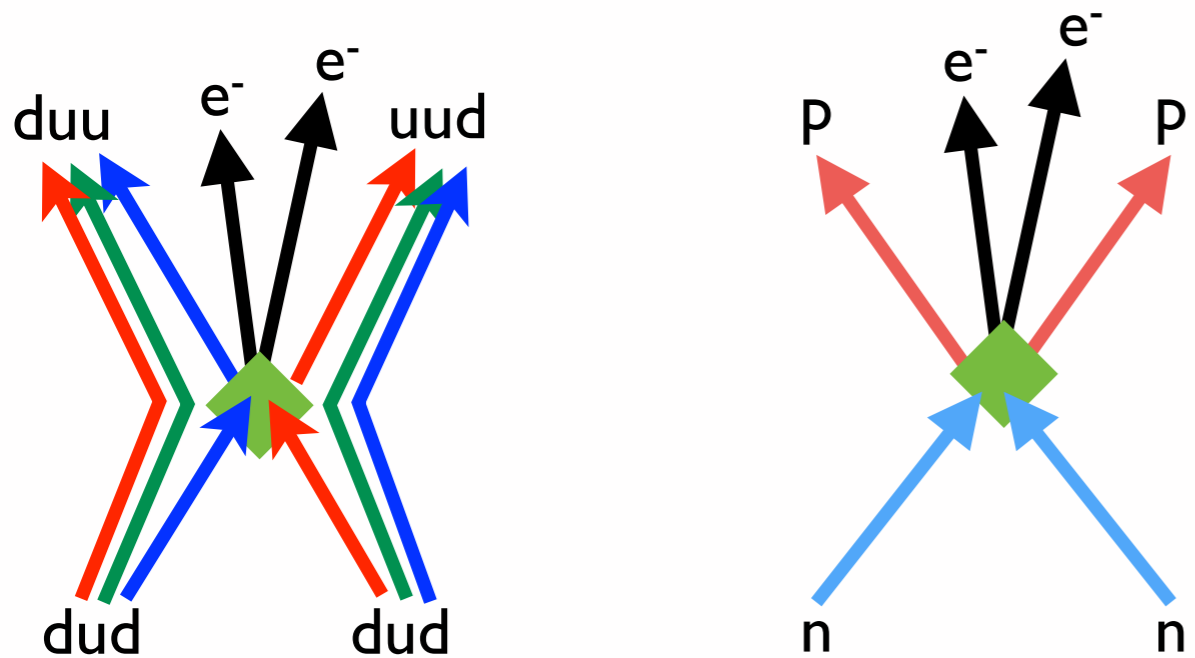
$\mathcal{O}(p^{-2})$  long-range  $\pi$  exchange



$\mathcal{O}(p^{-1})$   $\pi$ N vertex

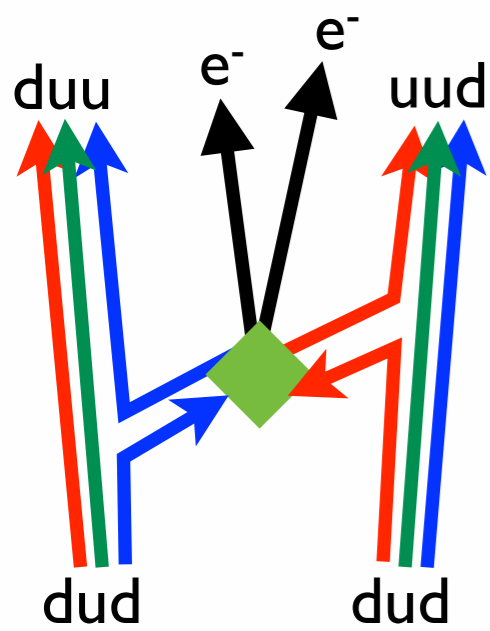


$\mathcal{O}(p^0)$  NN contact operator

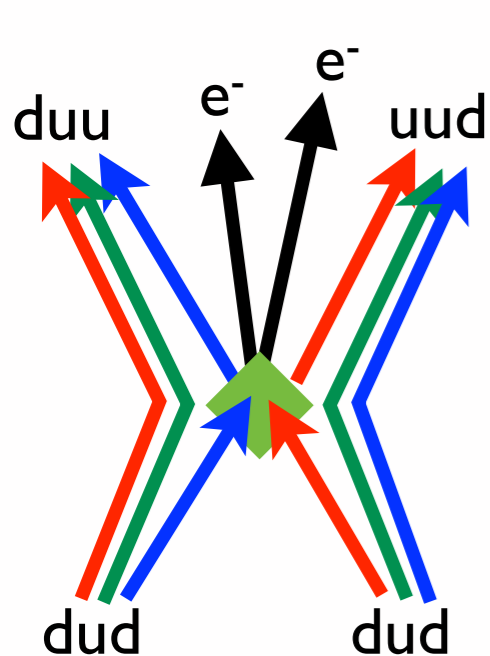


# $\chi$ PT power-counting of operators

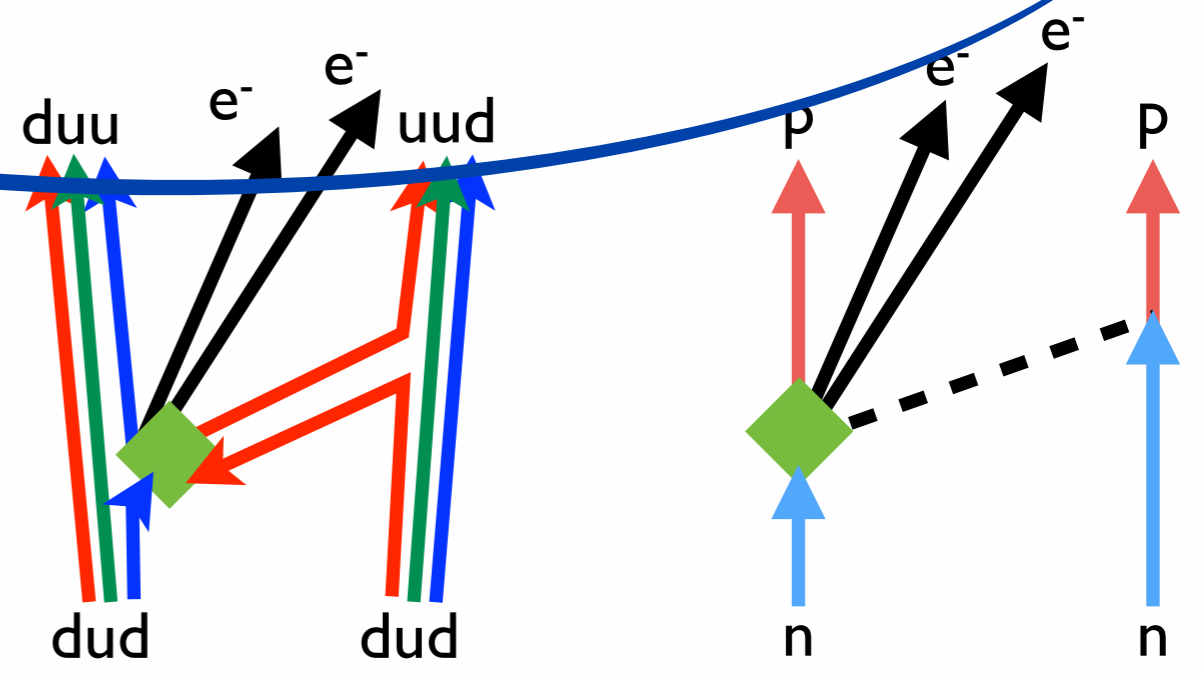
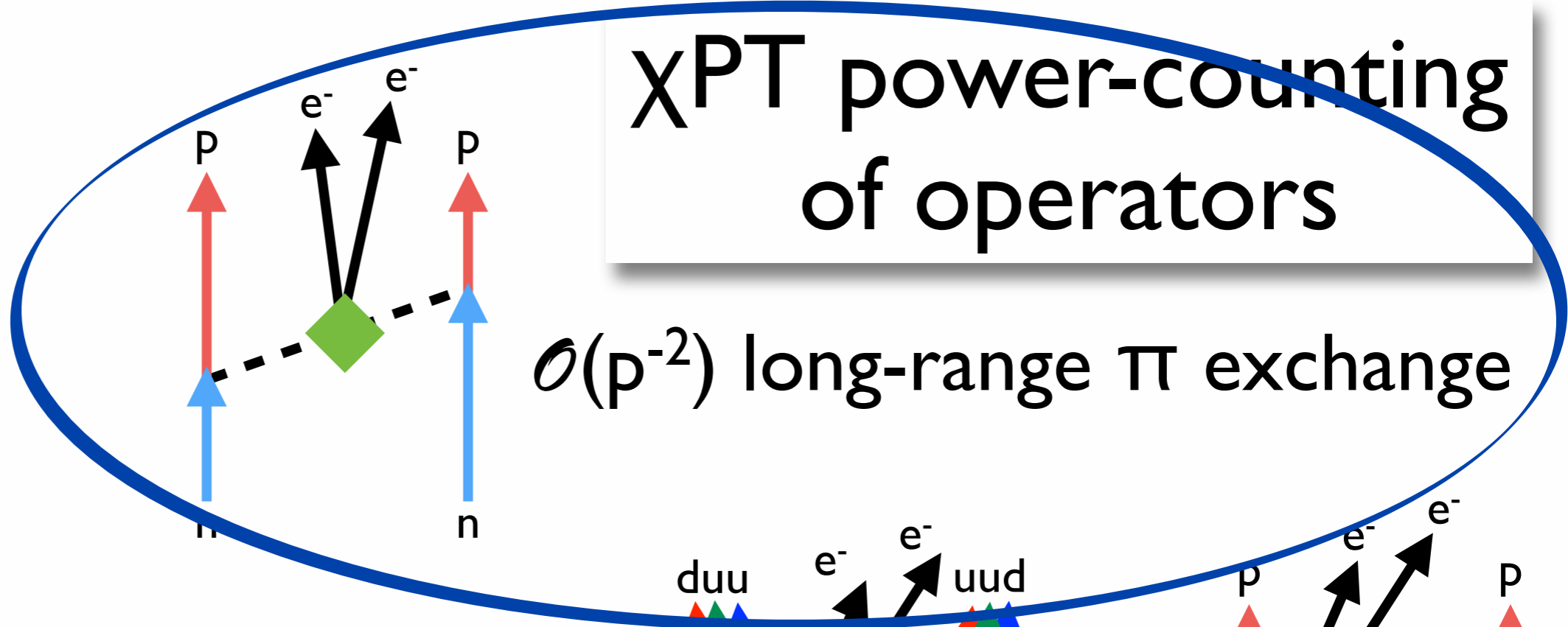
$\mathcal{O}(p^{-2})$  long-range  $\pi$  exchange



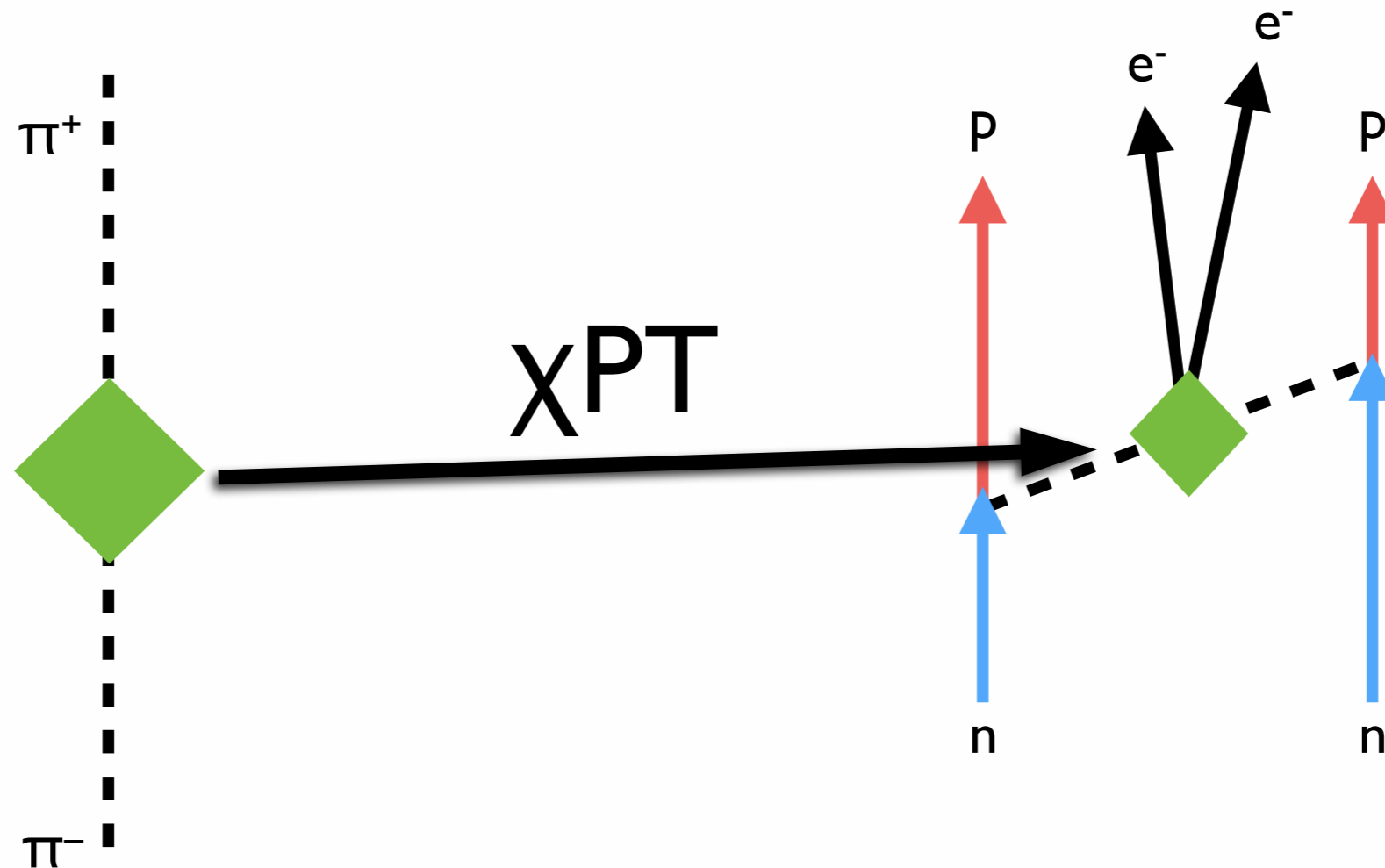
$\mathcal{O}(p^{-1})$   $\pi$ N vertex



$\mathcal{O}(p^0)$  NN contact operator

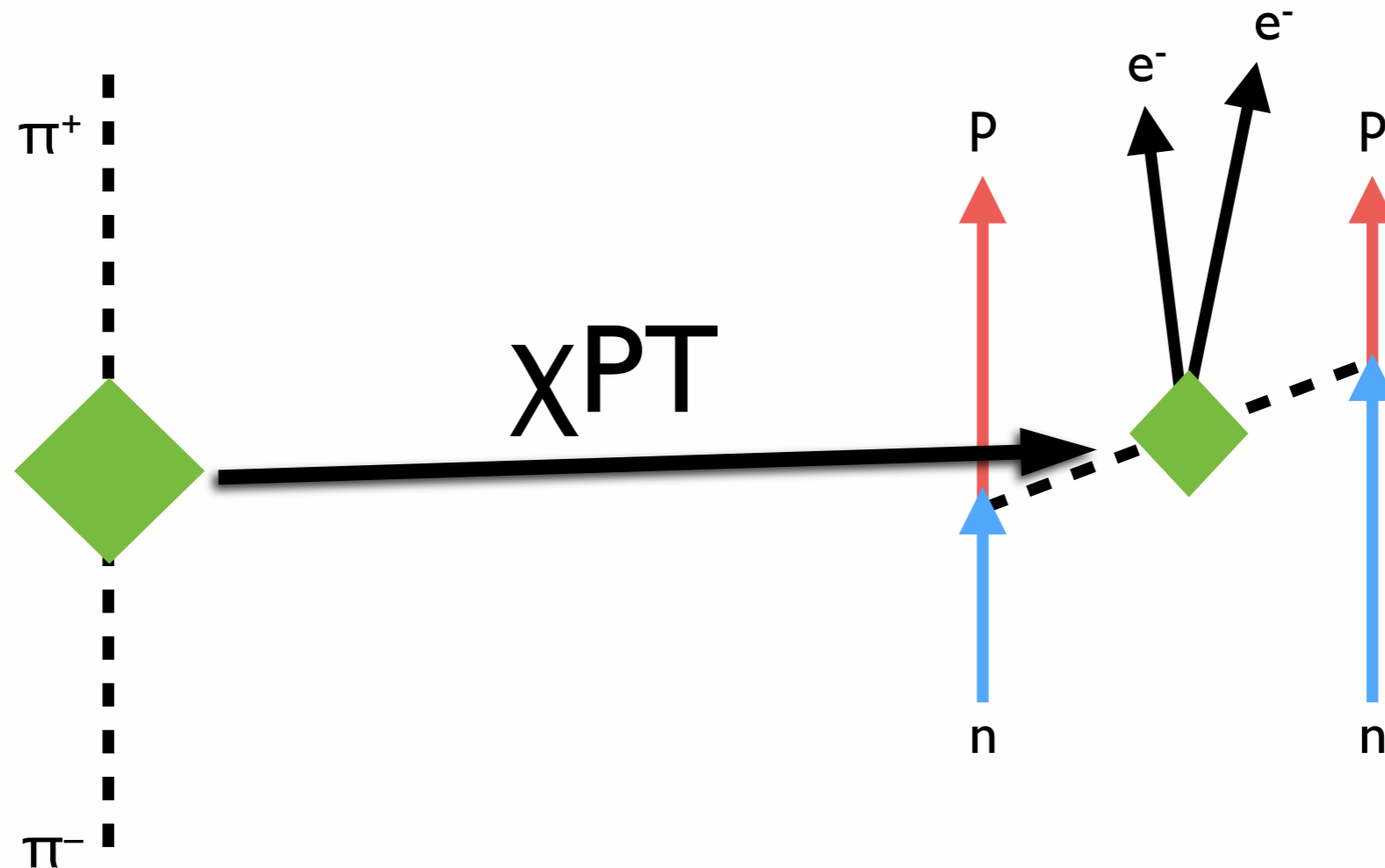


# Long Range pion contribution



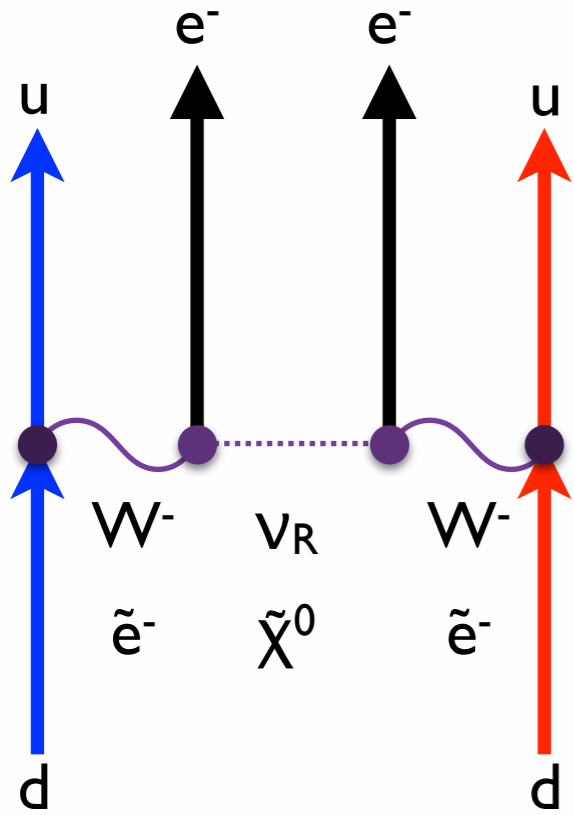
Unknown LEC is the same as for  $\pi^-$  to  $\pi^+$  transition: don't have to explicitly calculate two nucleon amplitude in LQCD

# Long Range pion contribution

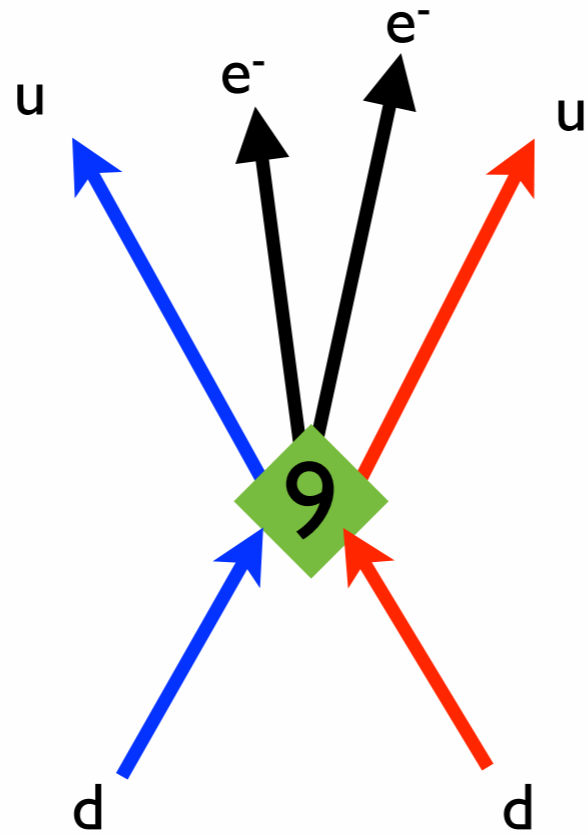


Easy to compute pion physics on the lattice!

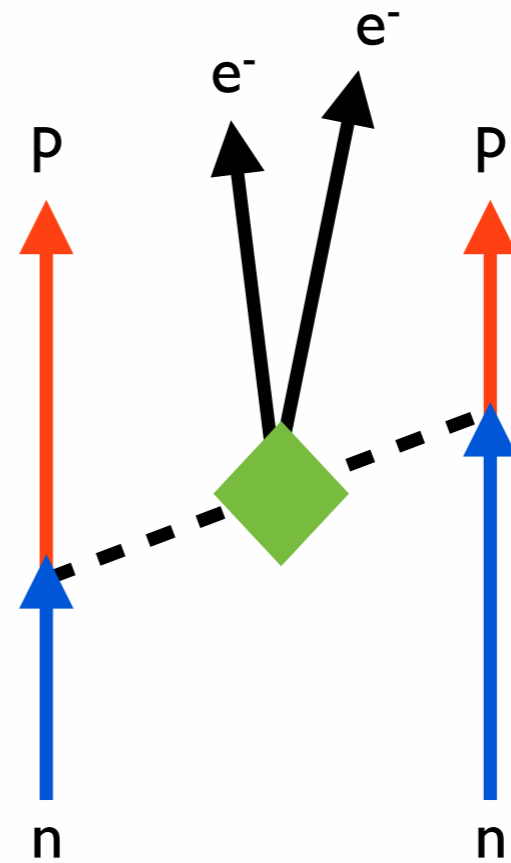
- Cheap
- Clean signals
- $I=2$ : no disconnected pieces



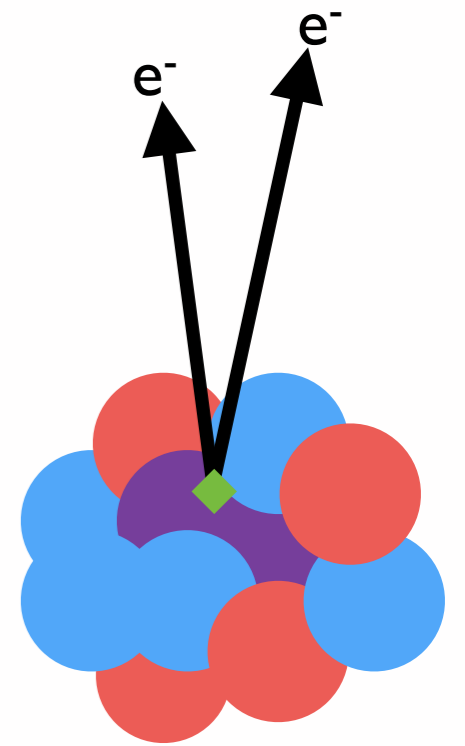
Your favorite  
BSM theory



effective 4-quark  
operators in QCD



two nucleon  
amplitudes

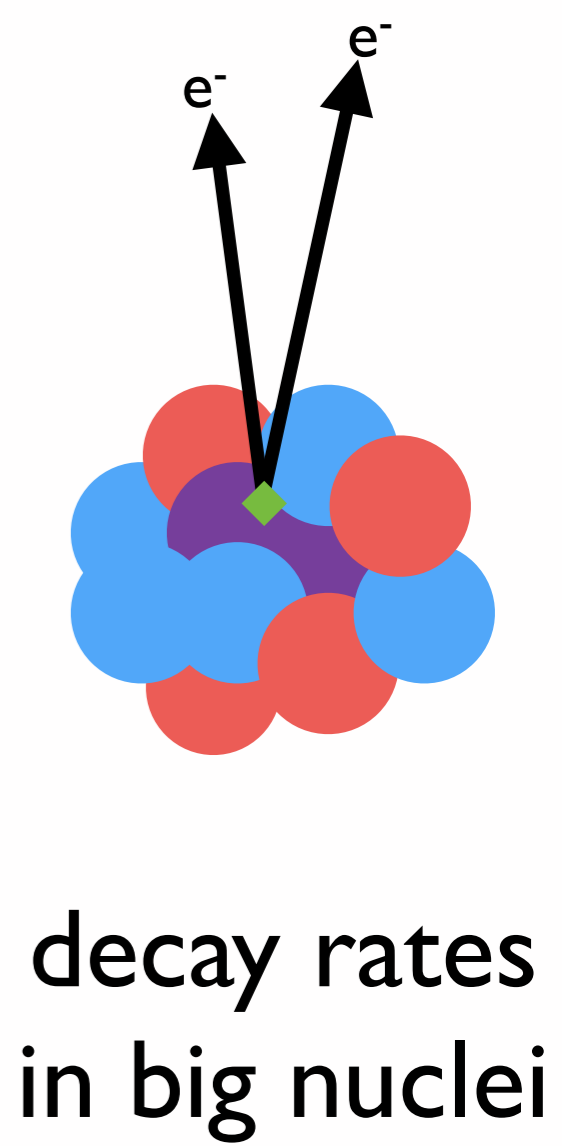
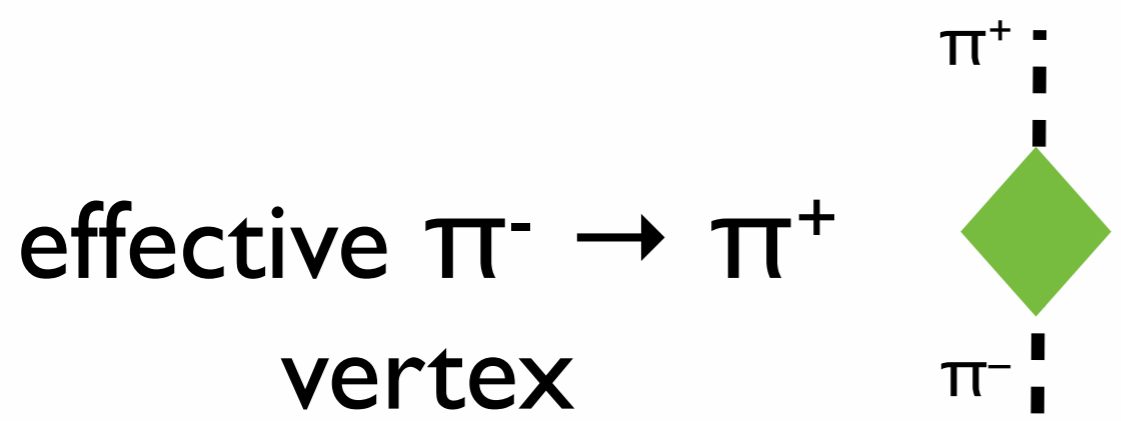
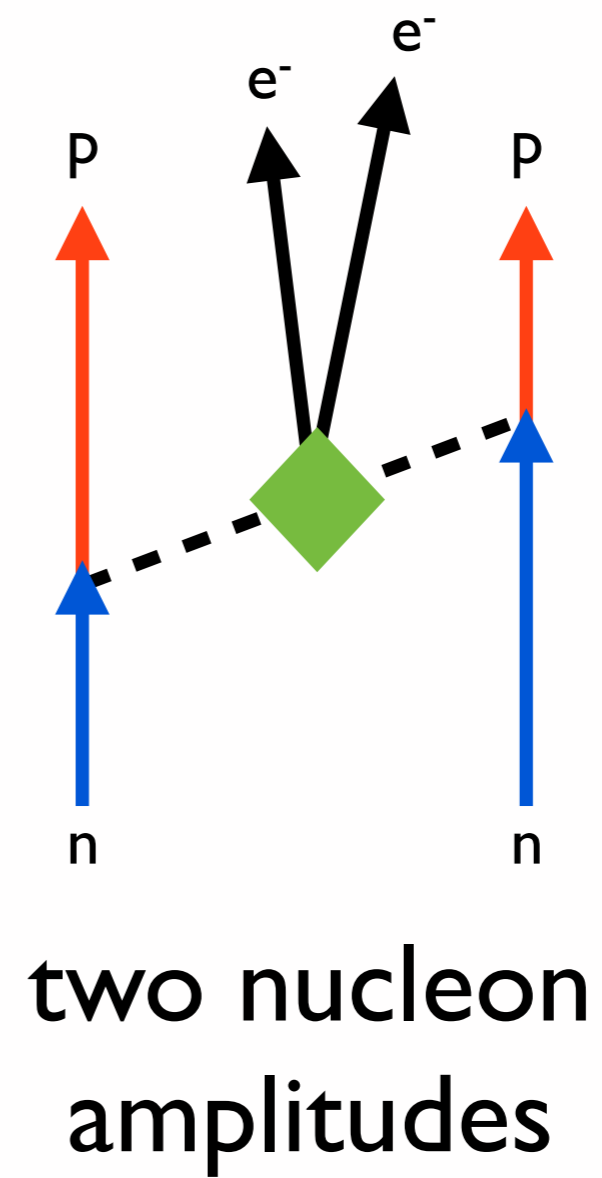
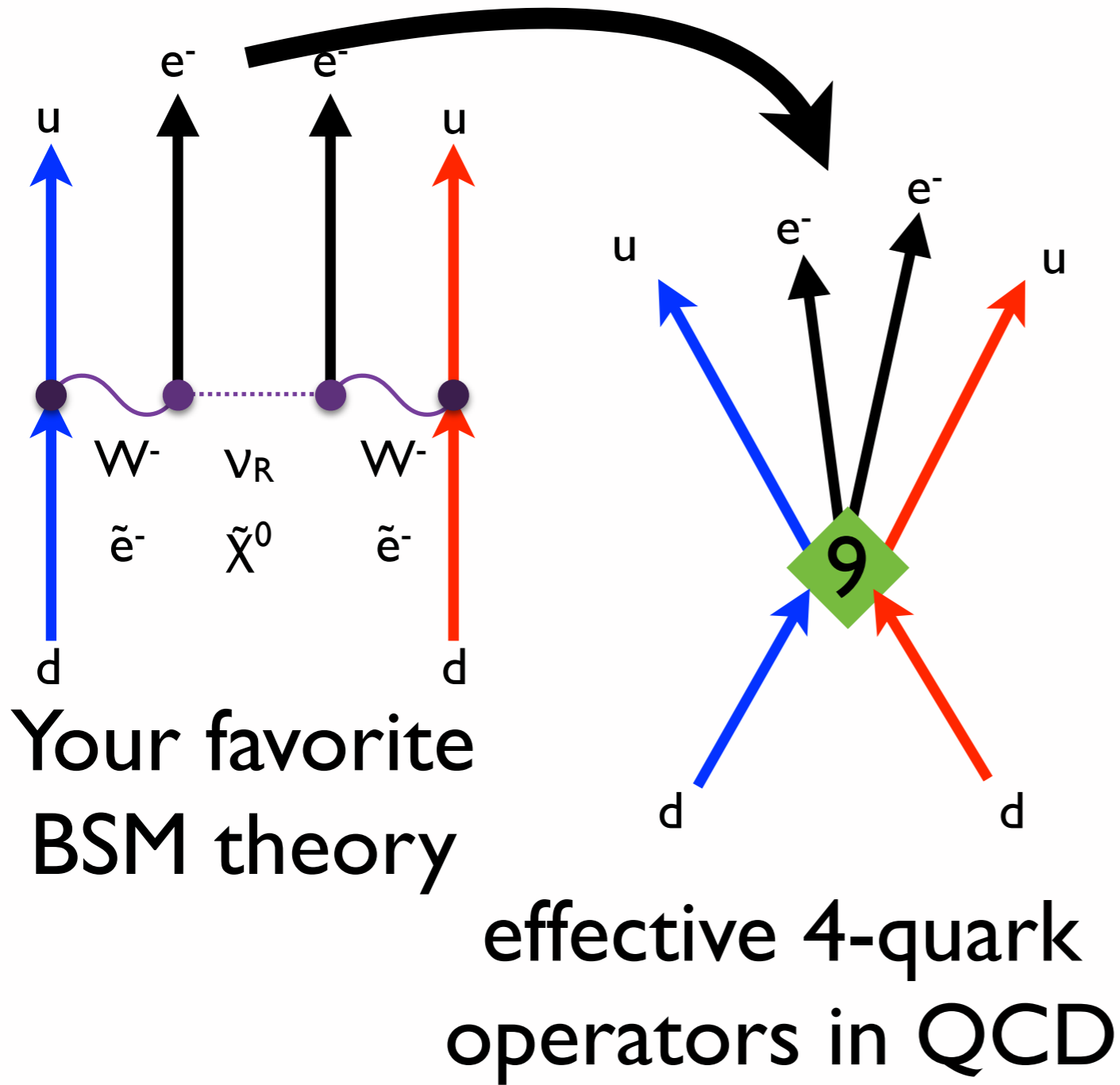


decay rates  
in big nuclei

effective  $\pi^- \rightarrow \pi^+$   
vertex

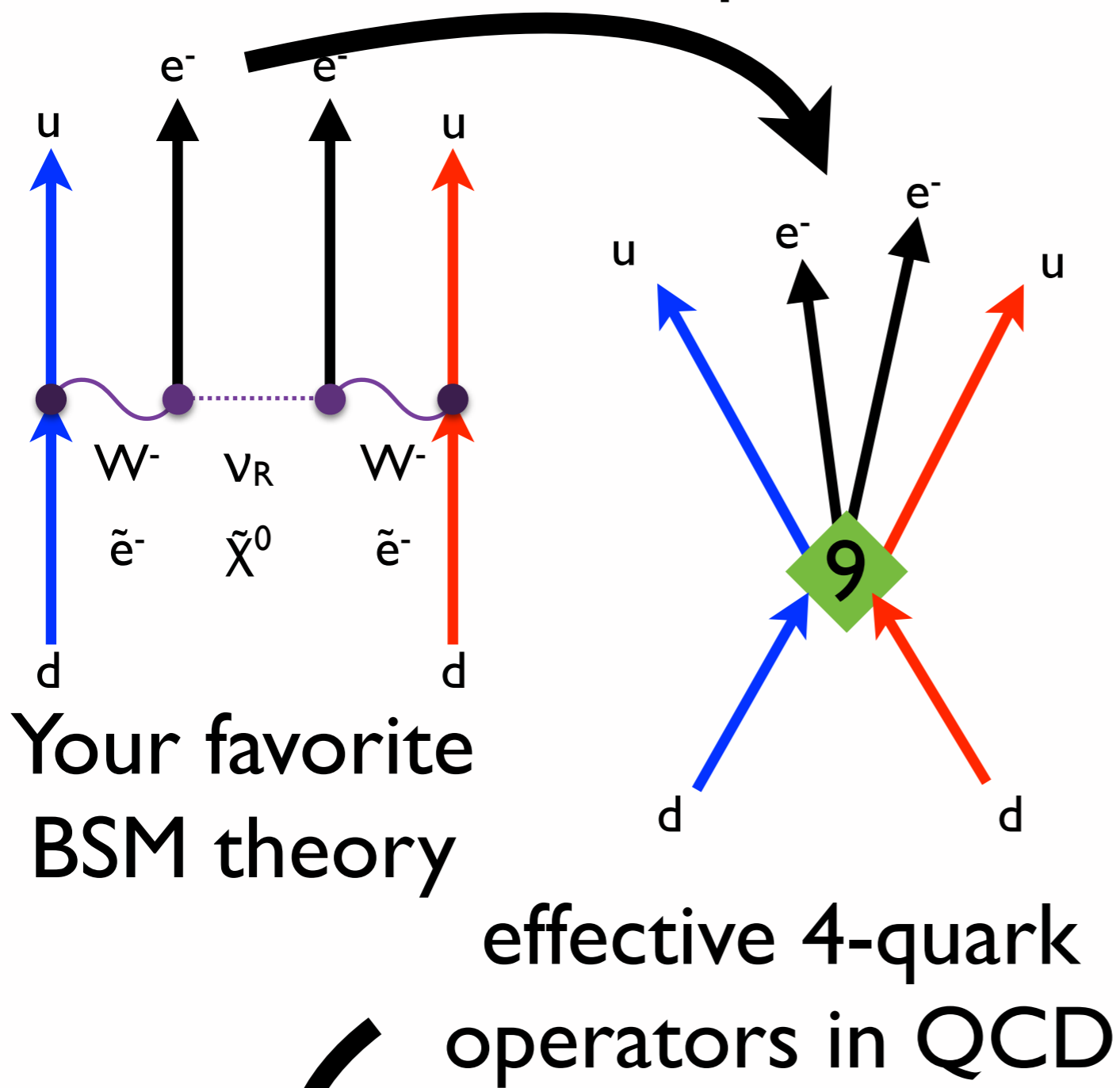


# Pen & Paper





# Pen & Paper

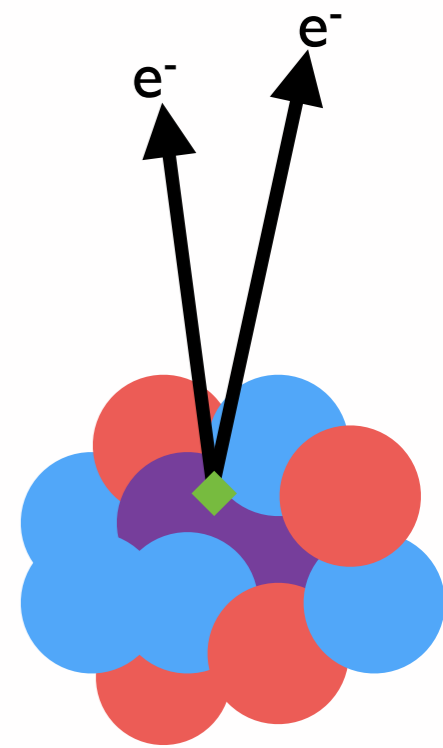


LQCD

effective  $\pi^- \rightarrow \pi^+$  vertex

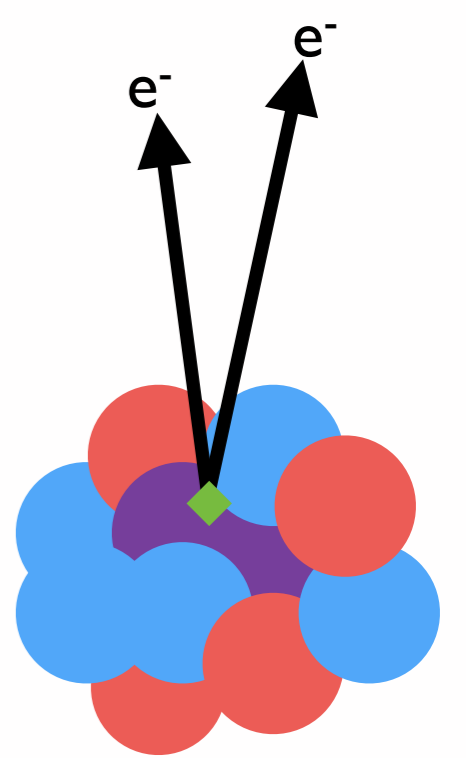
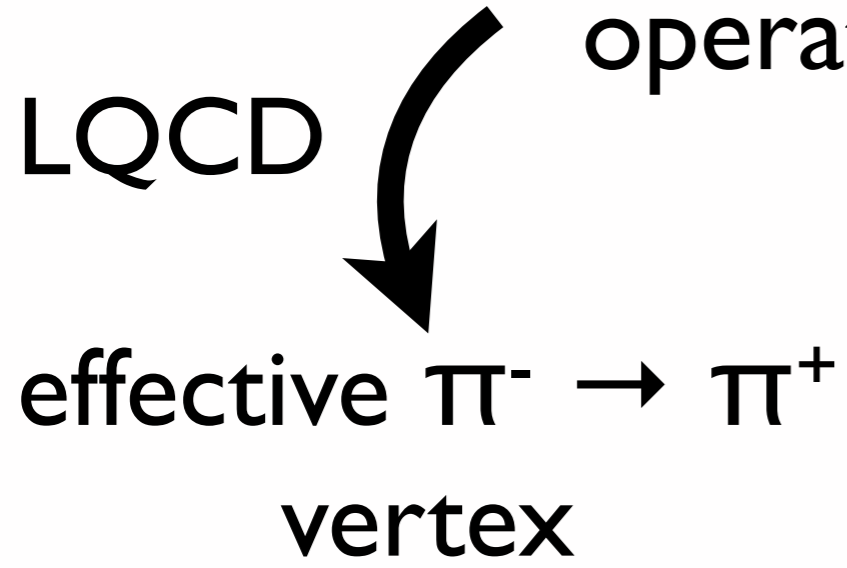
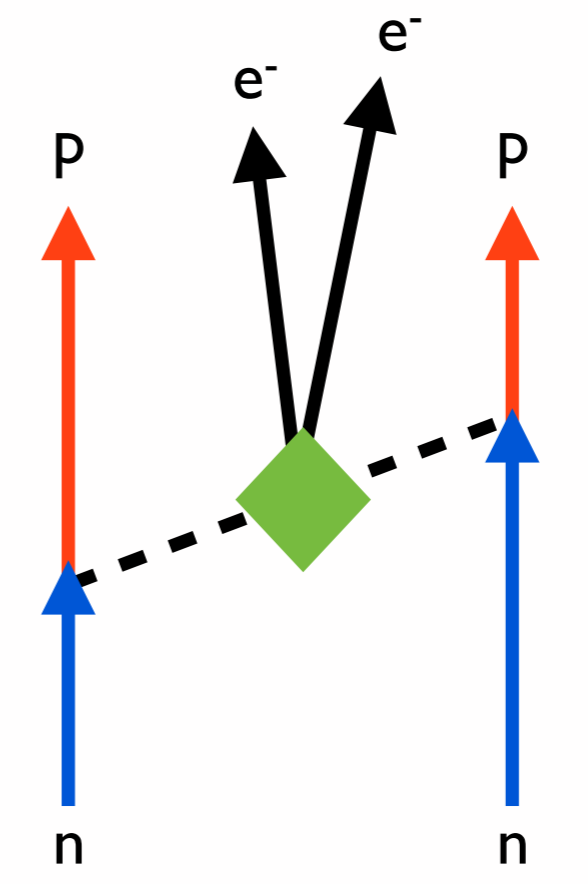
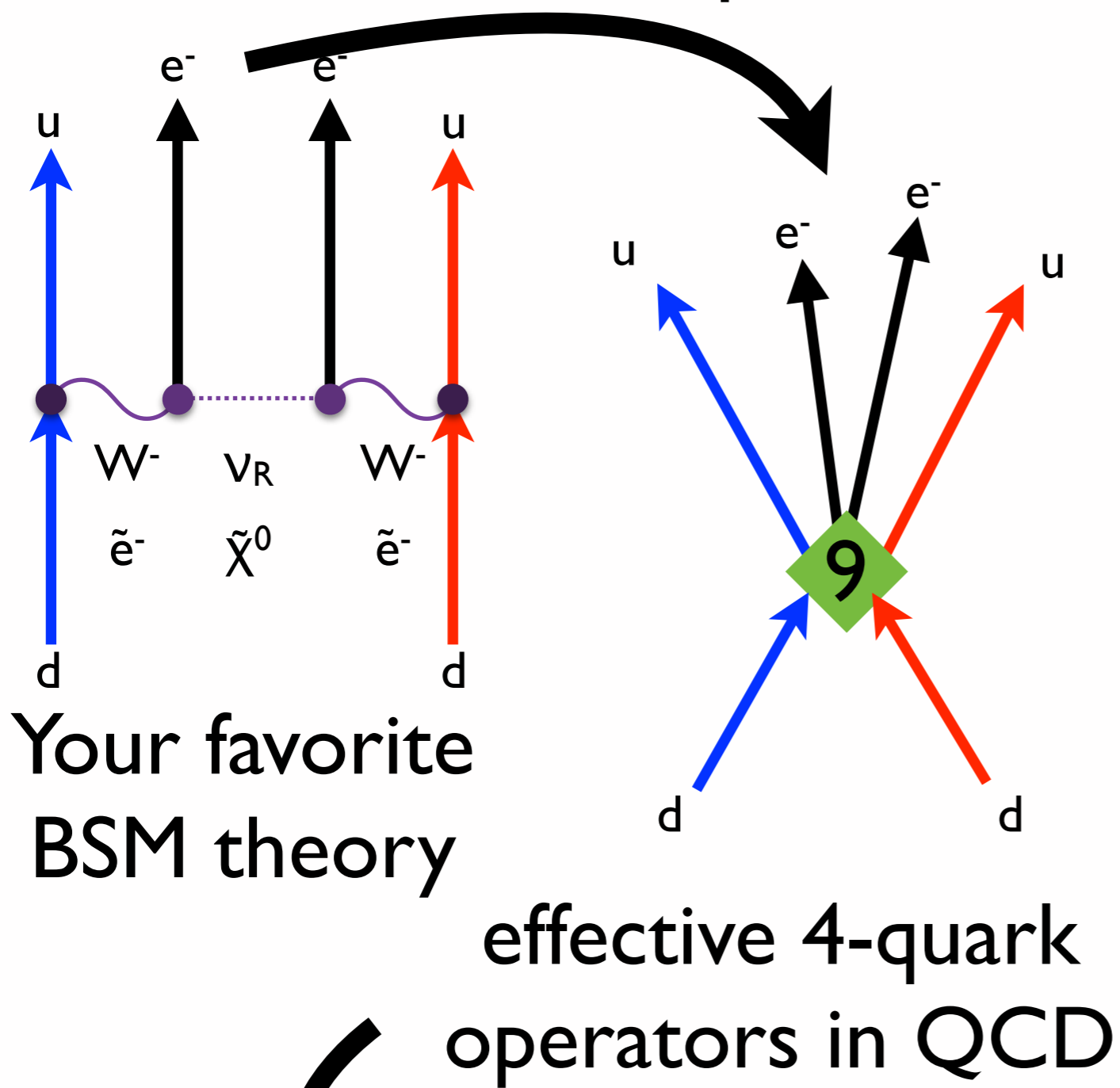


two nucleon amplitudes

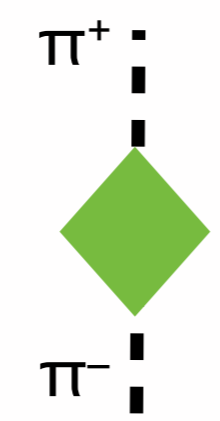
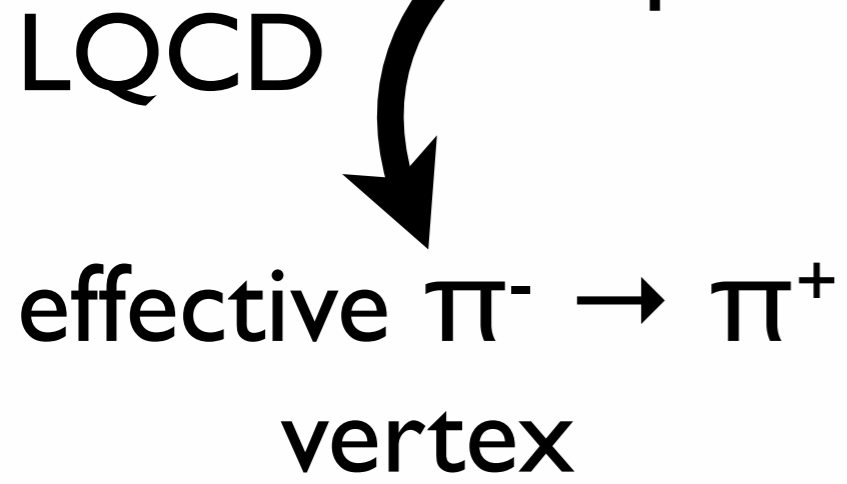
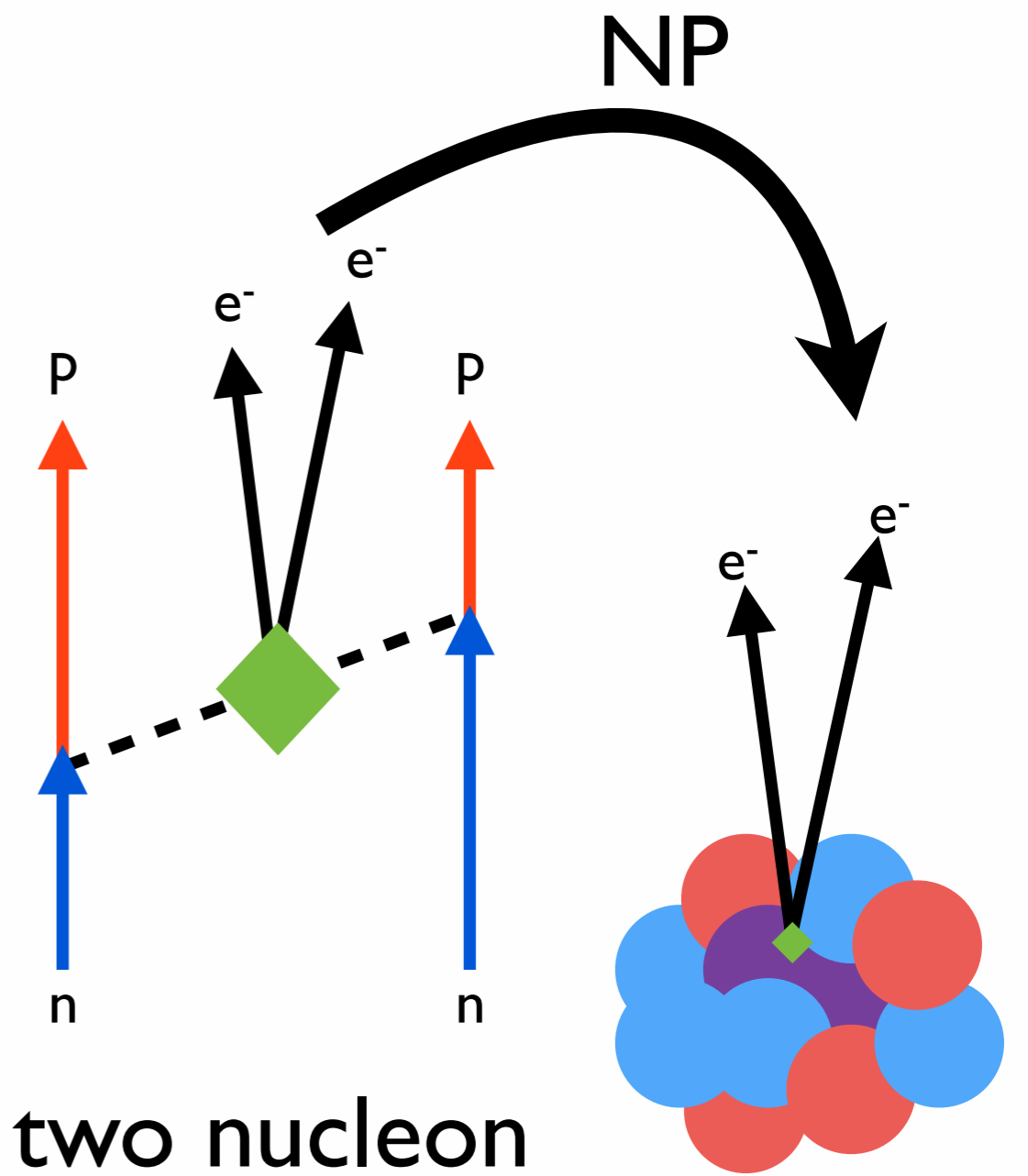
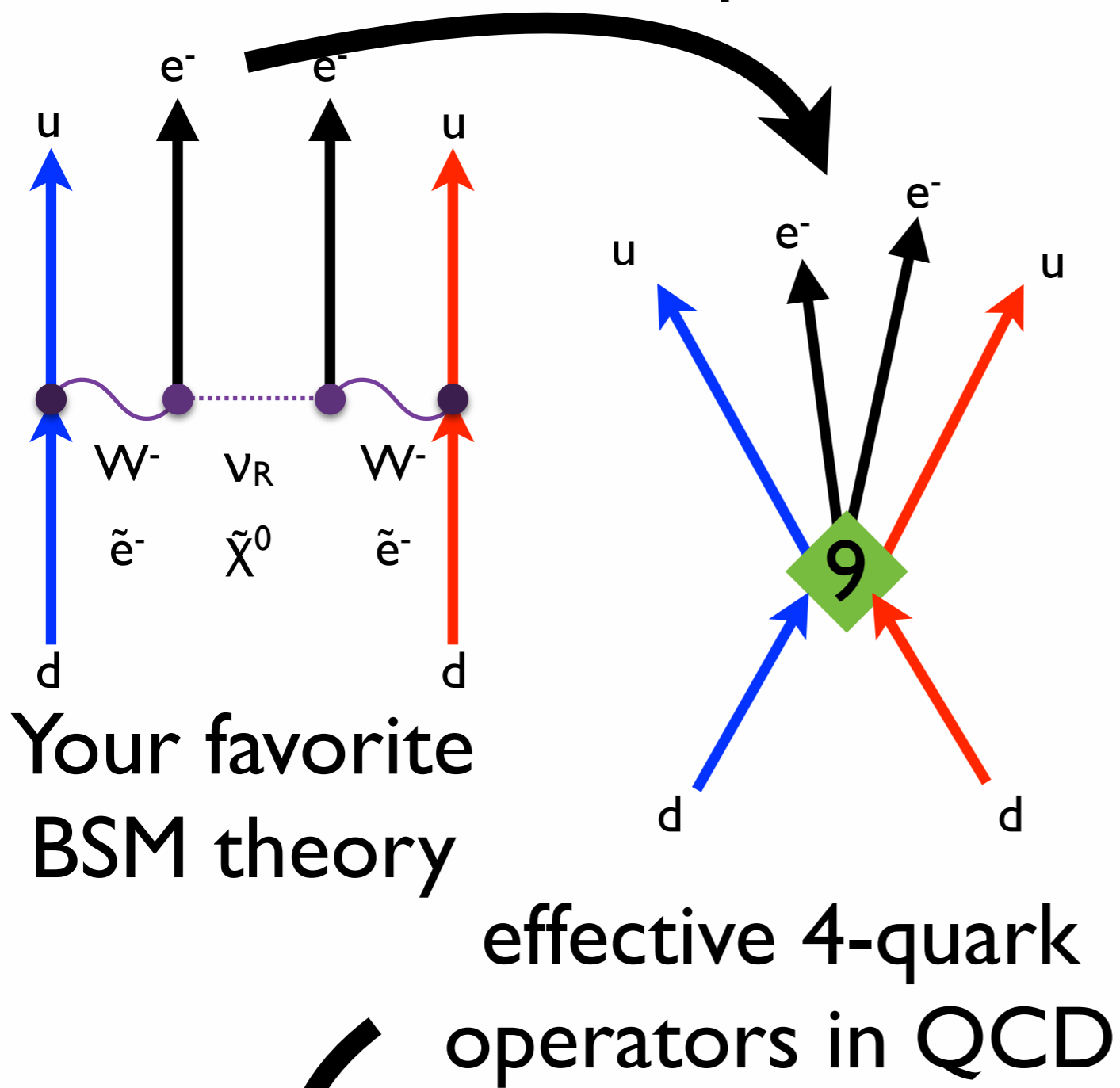


decay rates in big nuclei

# Pen & Paper



# Pen & Paper

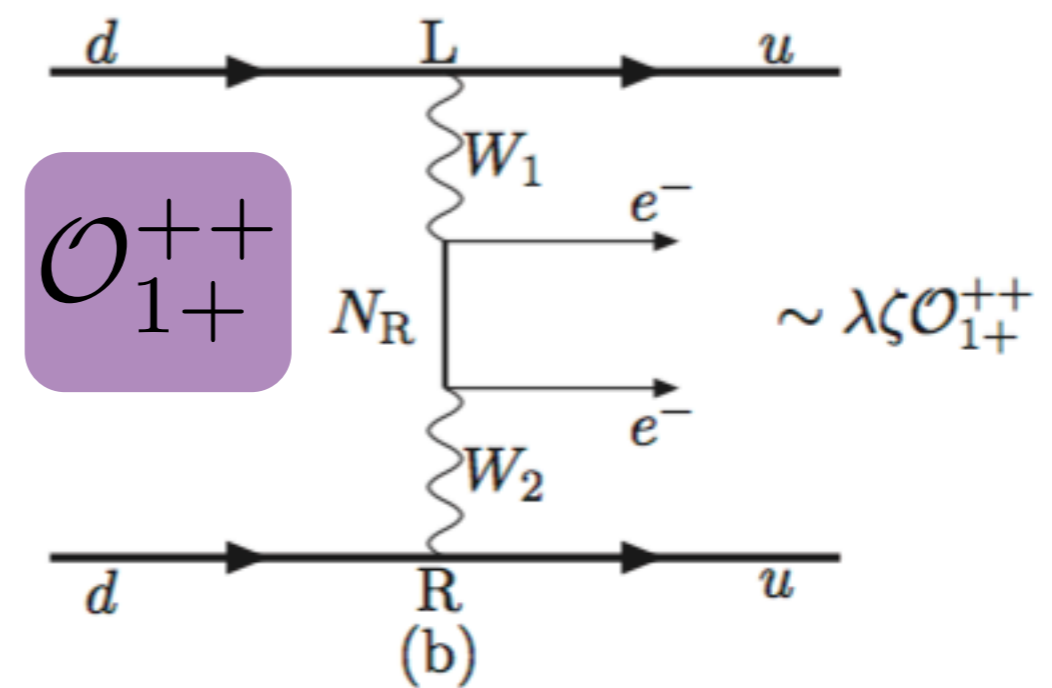
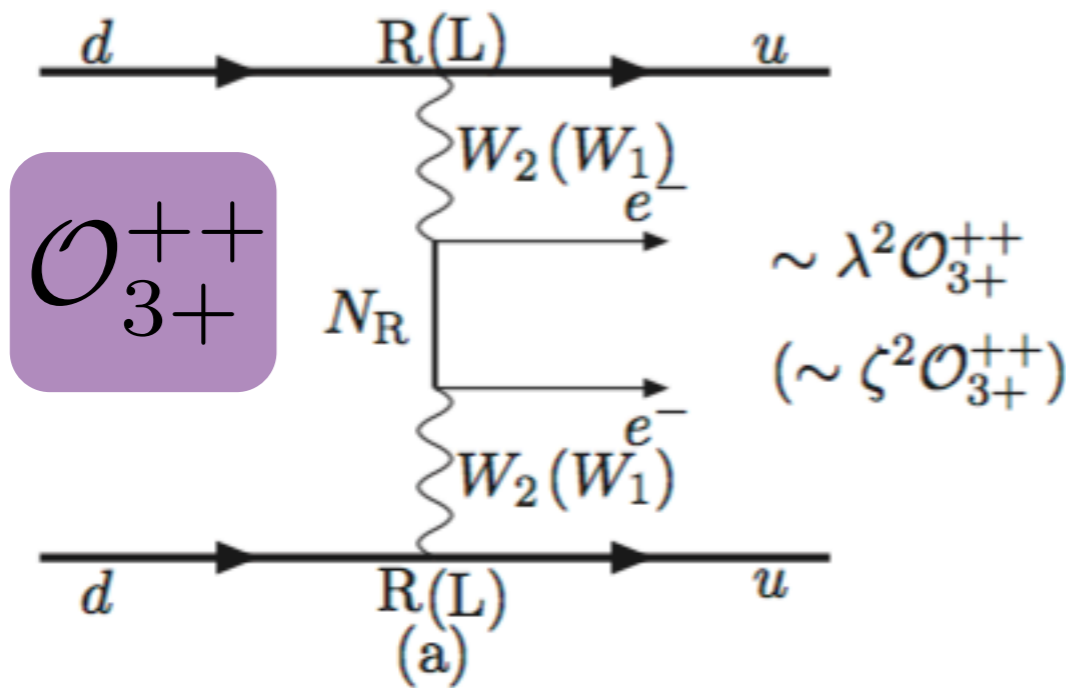


# LQCD calculation: details/results

XPT:

$0\nu\beta\beta$ -decay ops.	$\mathcal{O}_{1+}^{\pm\pm}$	$\mathcal{O}_{2+}^{\pm\pm}$	$\mathcal{O}_{2-}^{\pm\pm}$	$\mathcal{O}_{3+}^{\pm\pm}$	$\mathcal{O}_{3-}^{\pm\pm}$	$\mathcal{O}_{4+}^{\pm\pm,\mu}$	$\mathcal{O}_{4-}^{\pm\pm,\mu}$	$\mathcal{O}_{5+}^{\pm\pm,\mu}$	$\mathcal{O}_{5-}^{\pm\pm,\mu}$
$\pi\pi ee$ LO	✓	✓	X	X	X	X	X	X	X
$\pi\pi ee$ NNLO	✓	✓	X	✓	X	X	X	X	X
$NN\pi ee$ LO	X	X	✓	X	X	✓	✓	✓	✓
$NN\pi ee$ NLO	X	✓	X	✓	X	✓	✓	✓	✓
$NNNNee$ LO	✓	✓	X	✓	X	✓	✓	✓	✓

Left-right symmetric models



Prezeau, Ramsey-Musolf, Vogel (2003), Savage (1999)

# Contractions

- Can perform exact momentum projection at source and sink
- Must add color mixed versions of Prezeau, Ramsey-Musolf, Vogel ops 1&2

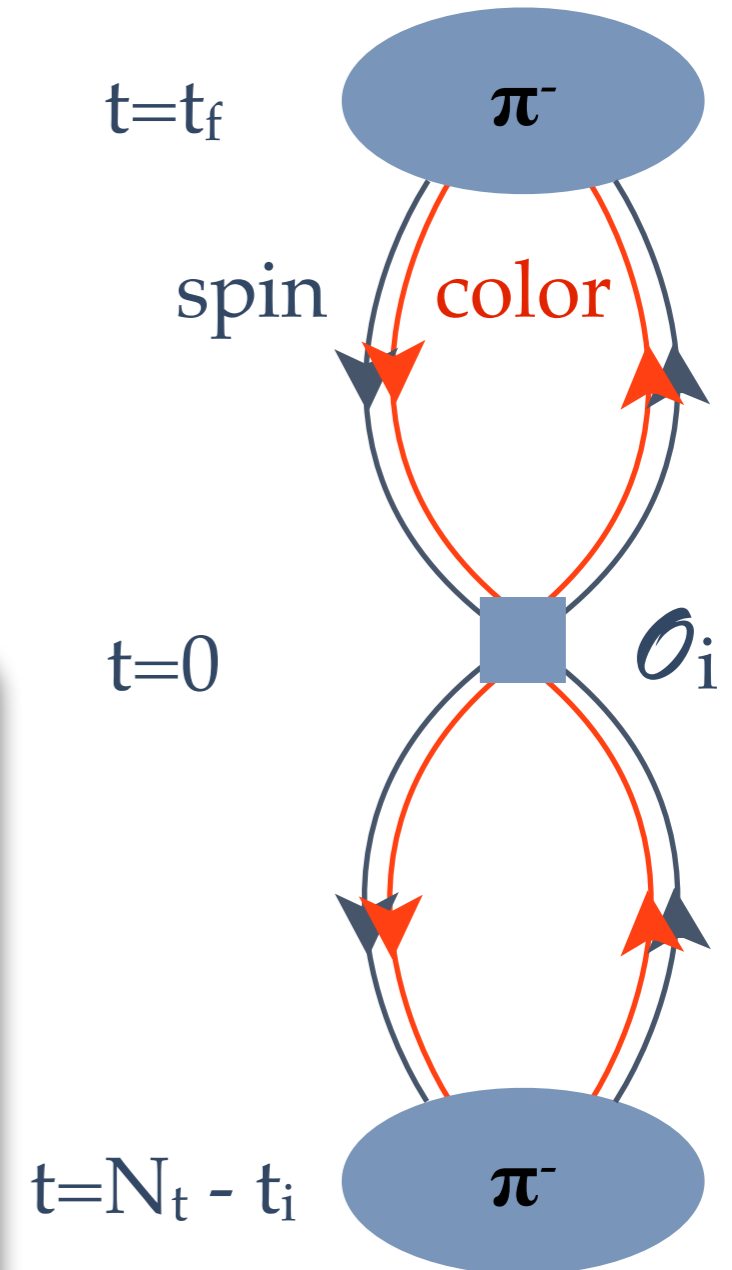
$$\mathcal{O}_{1+}^{++} = (\bar{q}_L \tau^- \gamma^\mu q_L) [\bar{q}_R \tau^- \gamma_\mu q_R]$$

$$\mathcal{O}'_{1+}^{++} = (\bar{q}_L \tau^- \gamma^\mu q_L) [\bar{q}_R \tau^- \gamma_\mu q_R]$$

$$\mathcal{O}_{2+}^{++} = (\bar{q}_R \tau^- q_L) [\bar{q}_R \tau^- q_L] + (\bar{q}_L \tau^- q_R) [\bar{q}_L \tau^- q_R]$$

$$\mathcal{O}'_{2+}^{++} = (\bar{q}_R \tau^- q_L) [\bar{q}_R \tau^- q_L] + (\bar{q}_L \tau^- q_R) [\bar{q}_L \tau^- q_R]$$

$$\mathcal{O}_{3+}^{++} = (\bar{q}_L \tau^- \gamma^\mu q_L) [\bar{q}_L \tau^- \gamma_\mu q_L] + (\bar{q}_R \tau^- \gamma^\mu q_R) [\bar{q}_R \tau^- \gamma_\mu q_R]$$



# Contractions

- Can perform exact momentum projection at source and sink
- Must add color mixed versions of Prezeau, Ramsey-Musolf, Vogel ops 1&2

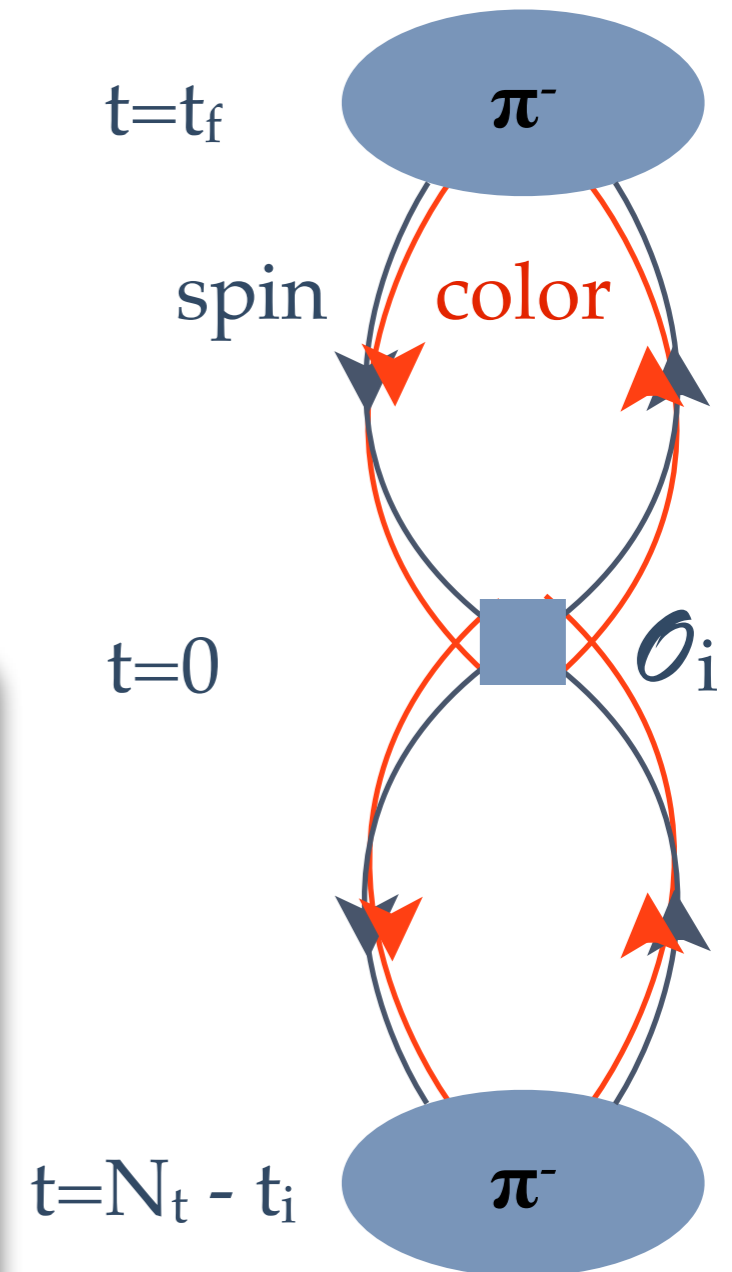
$$\mathcal{O}_{1+}^{++} = (\bar{q}_L \tau^- \gamma^\mu q_L) [\bar{q}_R \tau^- \gamma_\mu q_R]$$

$$\mathcal{O}'_{1+}^{++} = (\bar{q}_L \tau^- \gamma^\mu q_L) [\bar{q}_R \tau^- \gamma_\mu q_R]$$

$$\mathcal{O}_{2+}^{++} = (\bar{q}_R \tau^- q_L) [\bar{q}_R \tau^- q_L] + (\bar{q}_L \tau^- q_R) [\bar{q}_L \tau^- q_R]$$

$$\mathcal{O}'_{2+}^{++} = (\bar{q}_R \tau^- q_L) [\bar{q}_R \tau^- q_L] + (\bar{q}_L \tau^- q_R) [\bar{q}_L \tau^- q_R]$$

$$\mathcal{O}_{3+}^{++} = (\bar{q}_L \tau^- \gamma^\mu q_L) [\bar{q}_L \tau^- \gamma_\mu q_L] + (\bar{q}_R \tau^- \gamma^\mu q_R) [\bar{q}_R \tau^- \gamma_\mu q_R]$$



## HISQ ensembles

$a[fm] : m_\pi[MeV]$	310	220	135
0.15	$16^3 \times 48, m_\pi L \sim 3.78$	$24^3 \times 48, m_\pi L \sim 3.99$	$32^3 \times 48, m_\pi L \sim 3.25$
0.12		$24^3 \times 64, m_\pi L \sim 3.22$	
0.12	$24^3 \times 64, m_\pi L \sim 4.54$	$32^3 \times 64, m_\pi L \sim 4.29$	$48^3 \times 64, m_\pi L \sim 3.91$
0.12		$40^3 \times 64, m_\pi L \sim 5.36$	
0.09	$32^3 \times 96, m_\pi L \sim 4.50$	$48^3 \times 96, m_\pi L \sim 4.73$	

- Möbius DWF on HISQ
- Gradient flow method for smearing configs
  - $m_{\text{res}} < 0.1 m_\ell$  for moderate  $L_5$
- Wall + point sources for pions
- $\sim 1000$  cfigs, 1 source/cfg

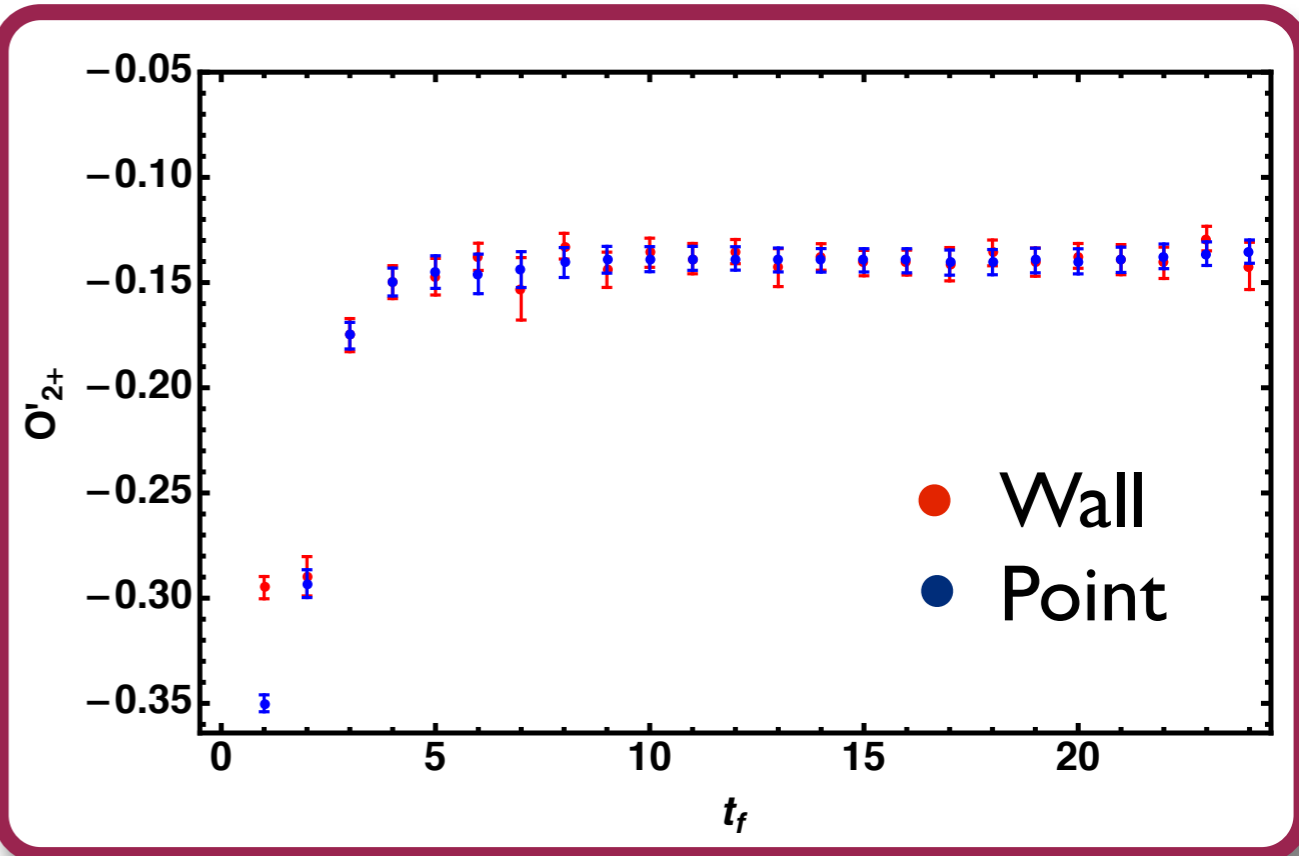
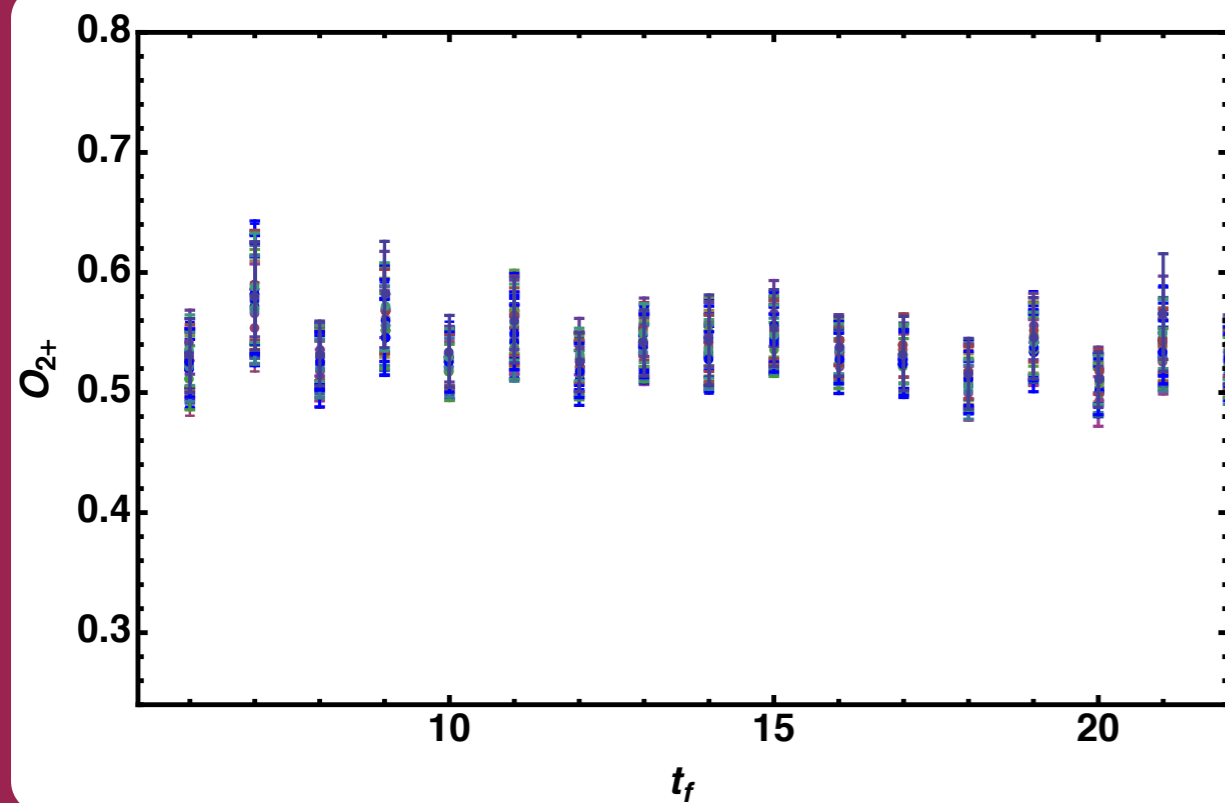
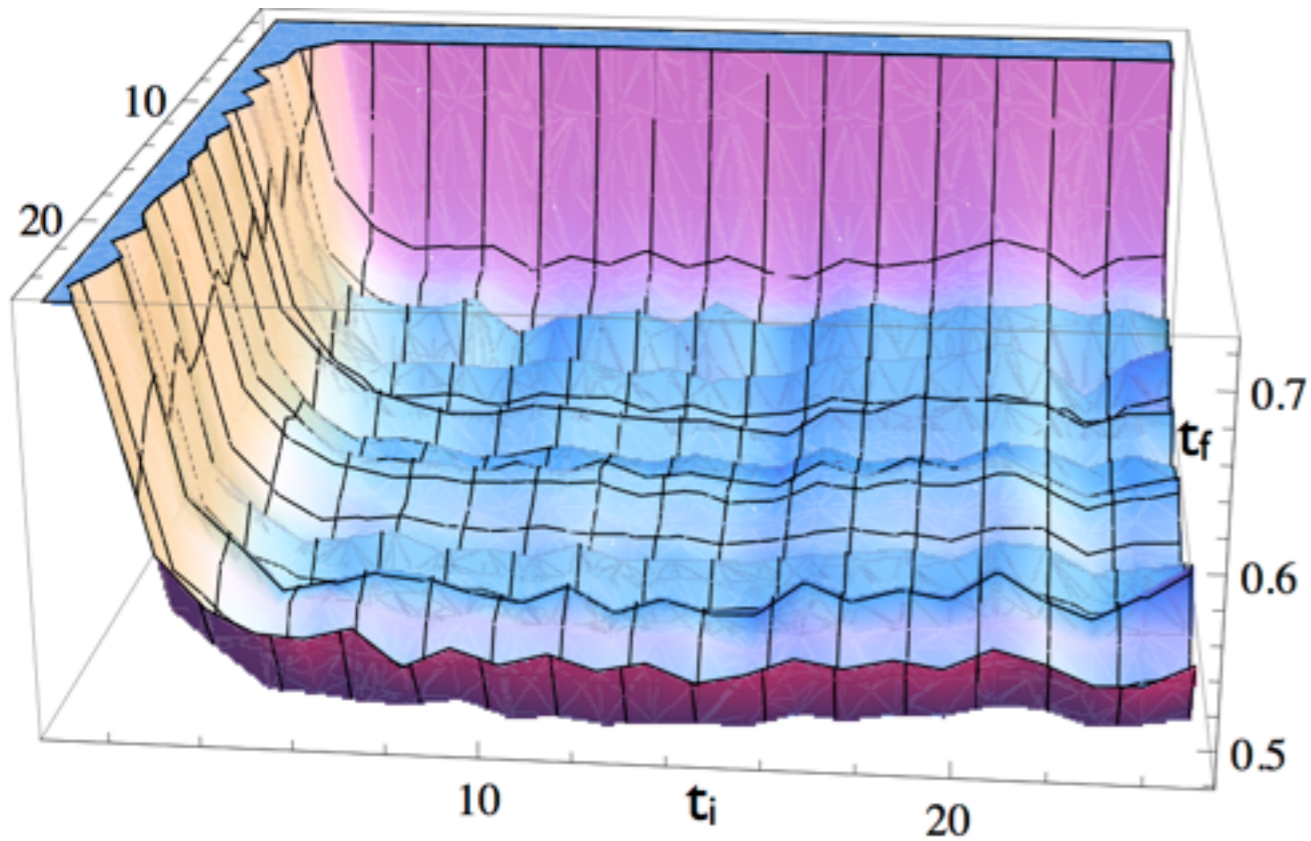
MILC Collaboration Phys. Rev. D87 (2013) 054505

Narayanan, Neuberger (2006), Luscher (2010)

K. Orginos, C. Monahan (private communication)

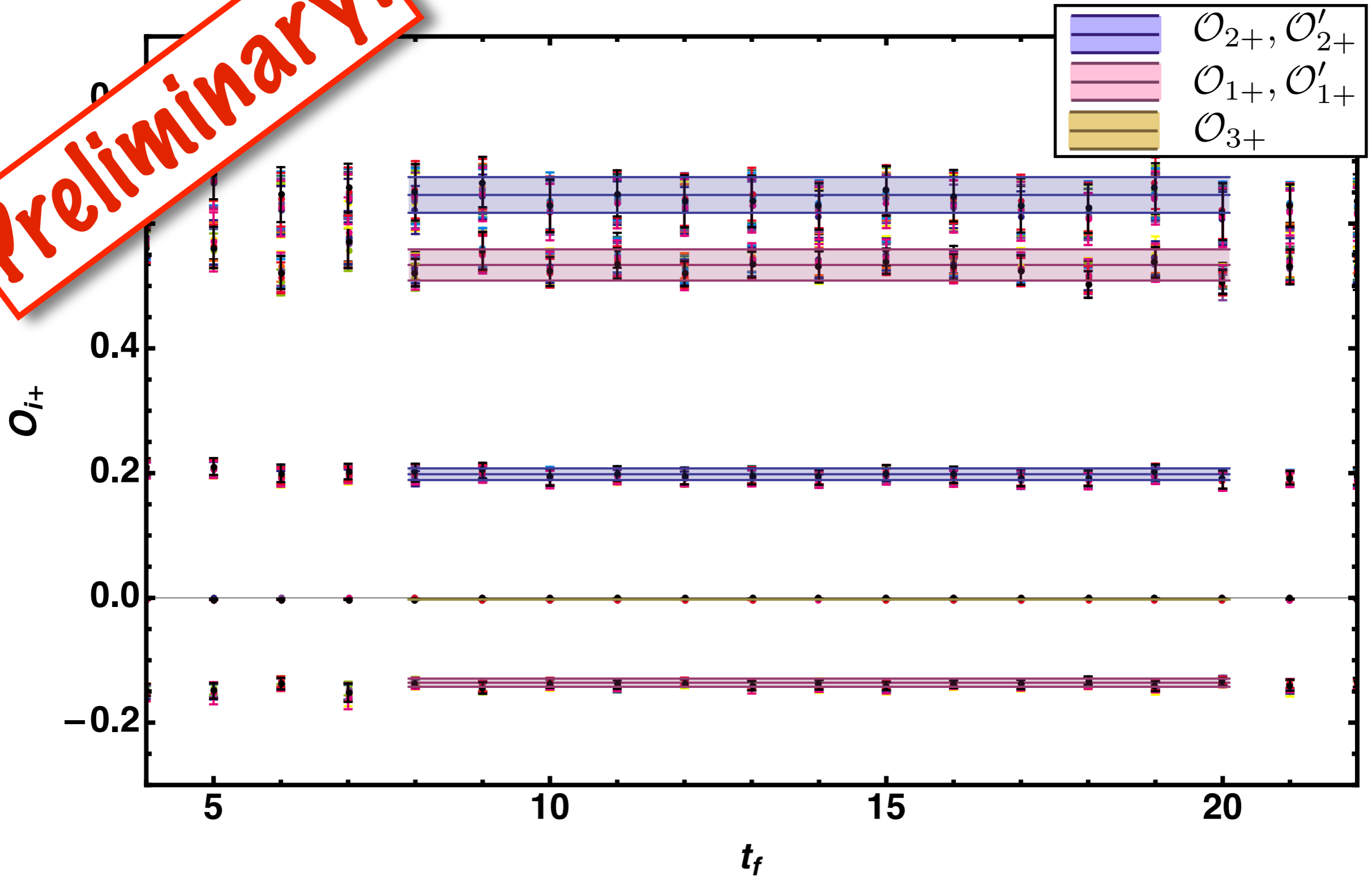


# Signals

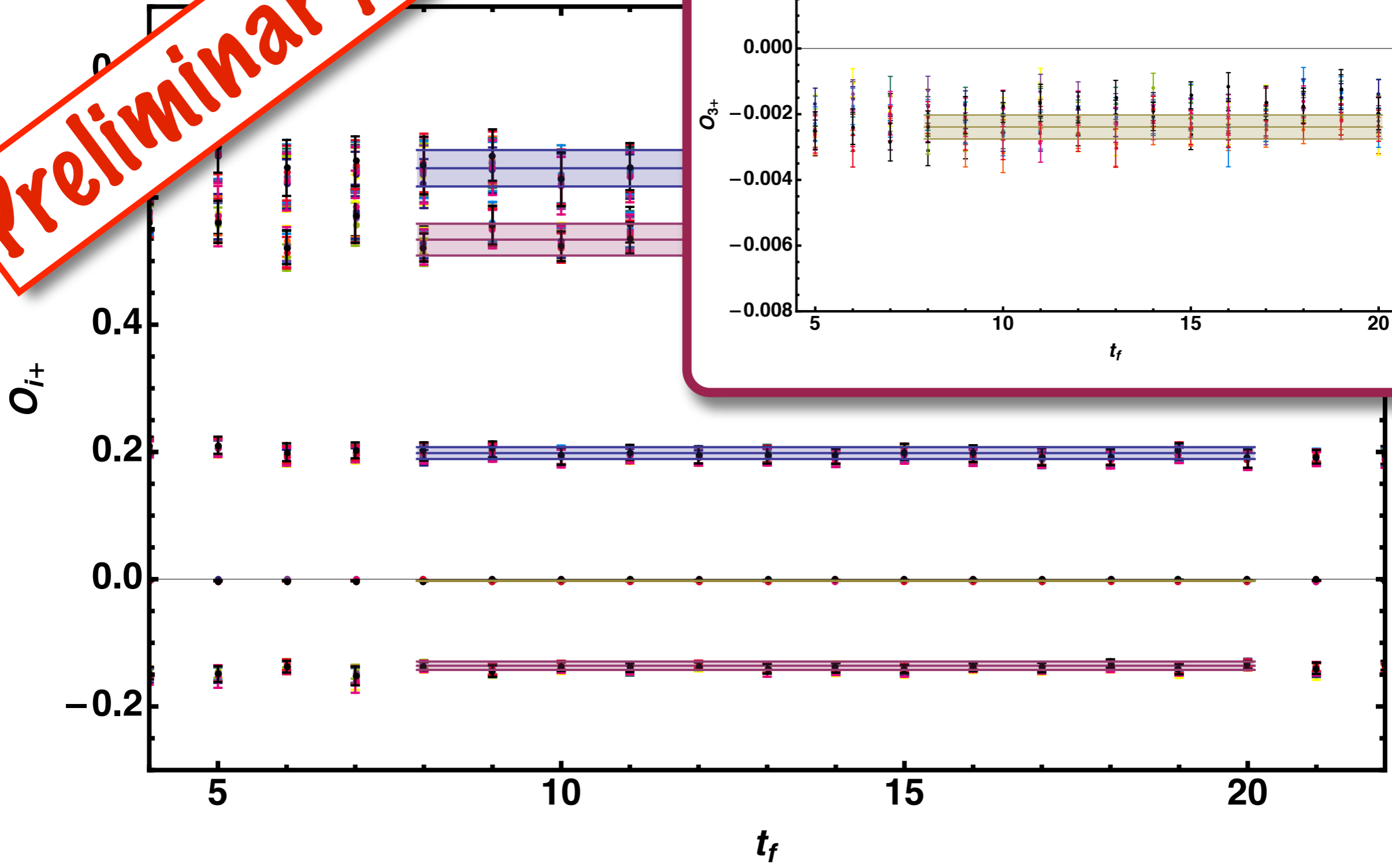


- $m_\pi \sim 135$  MeV
- $L = 5.76$  fm
- $a = 0.12$  fm

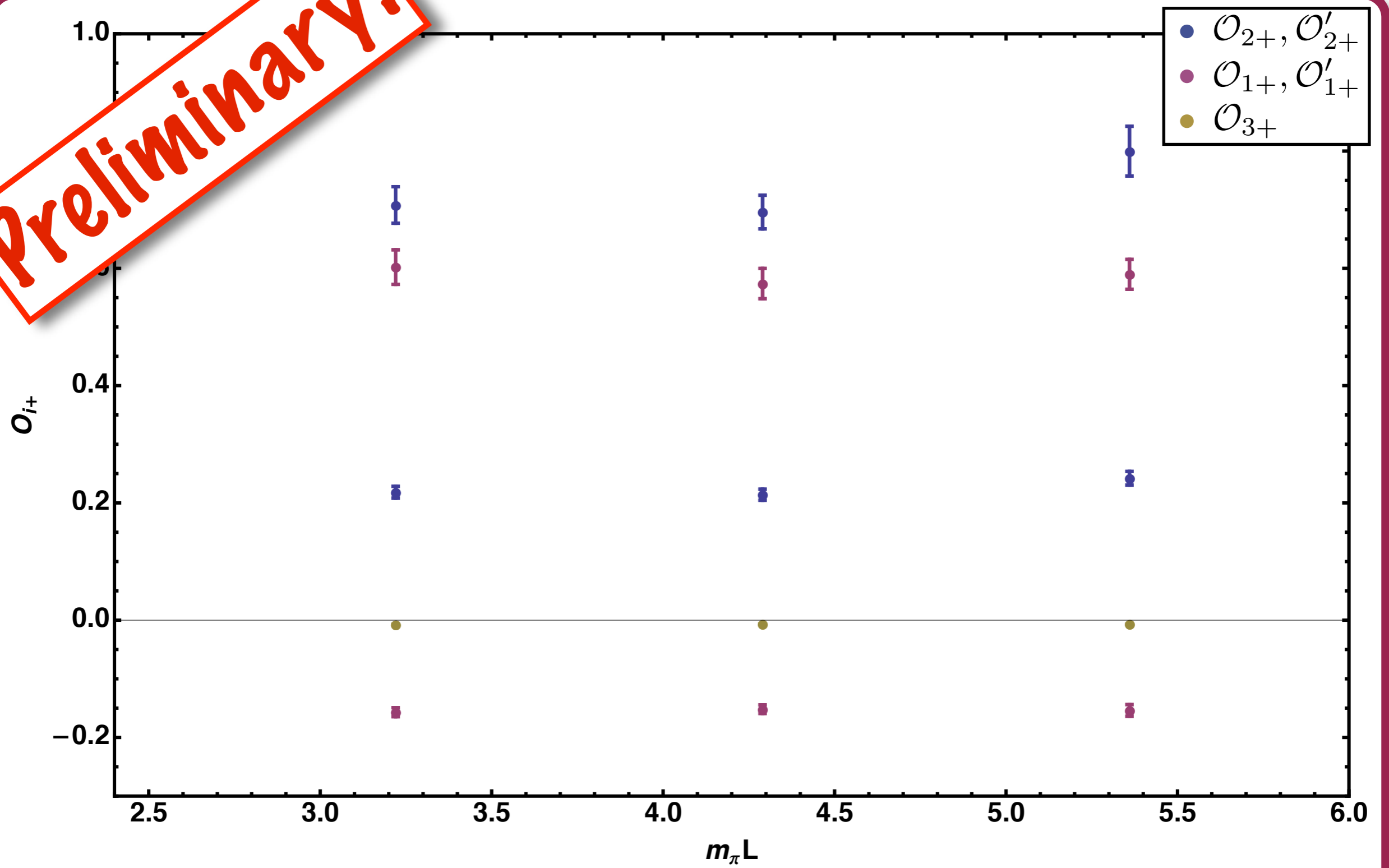
**Preliminary!**



**Preliminary!**

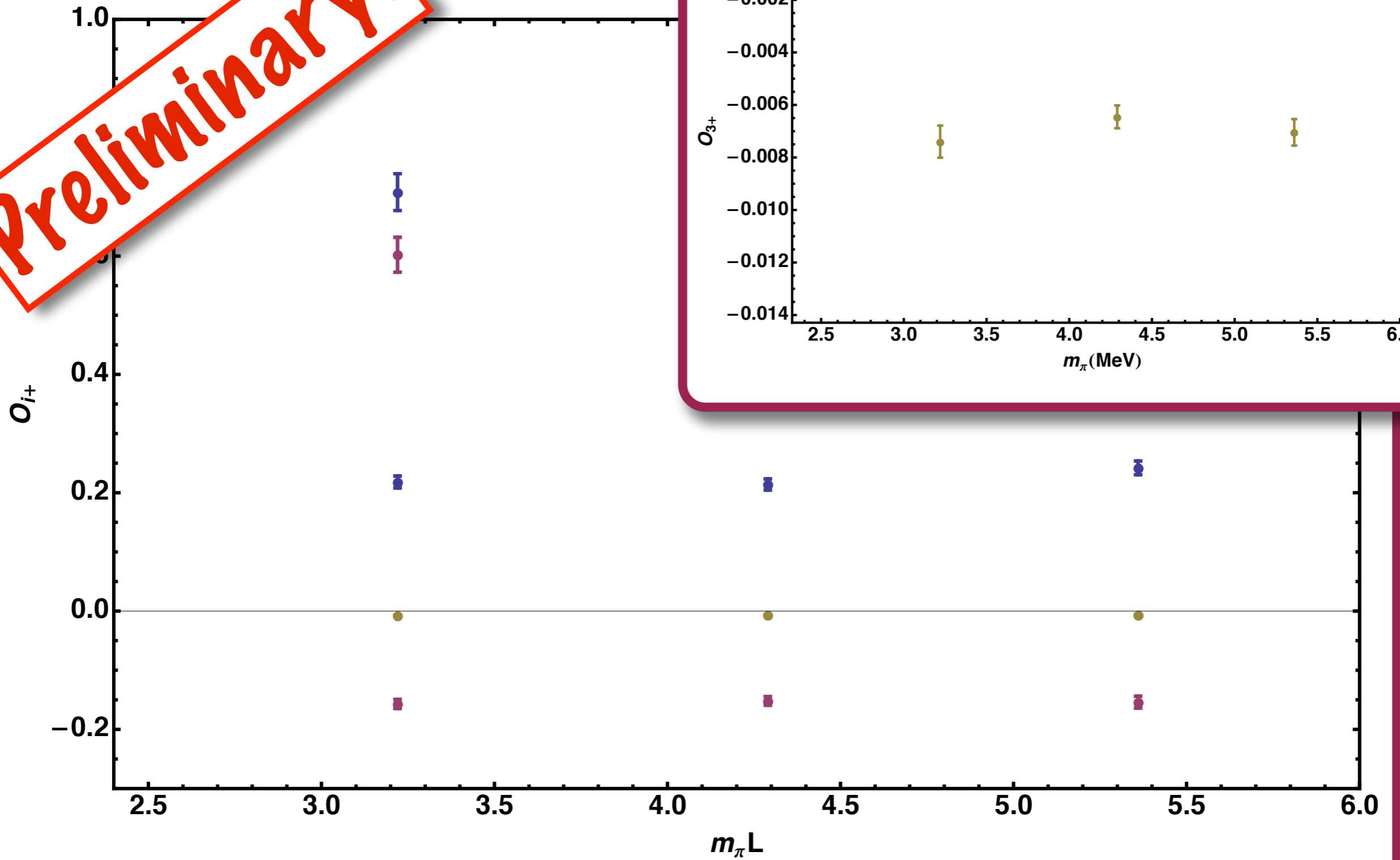


Preliminary!



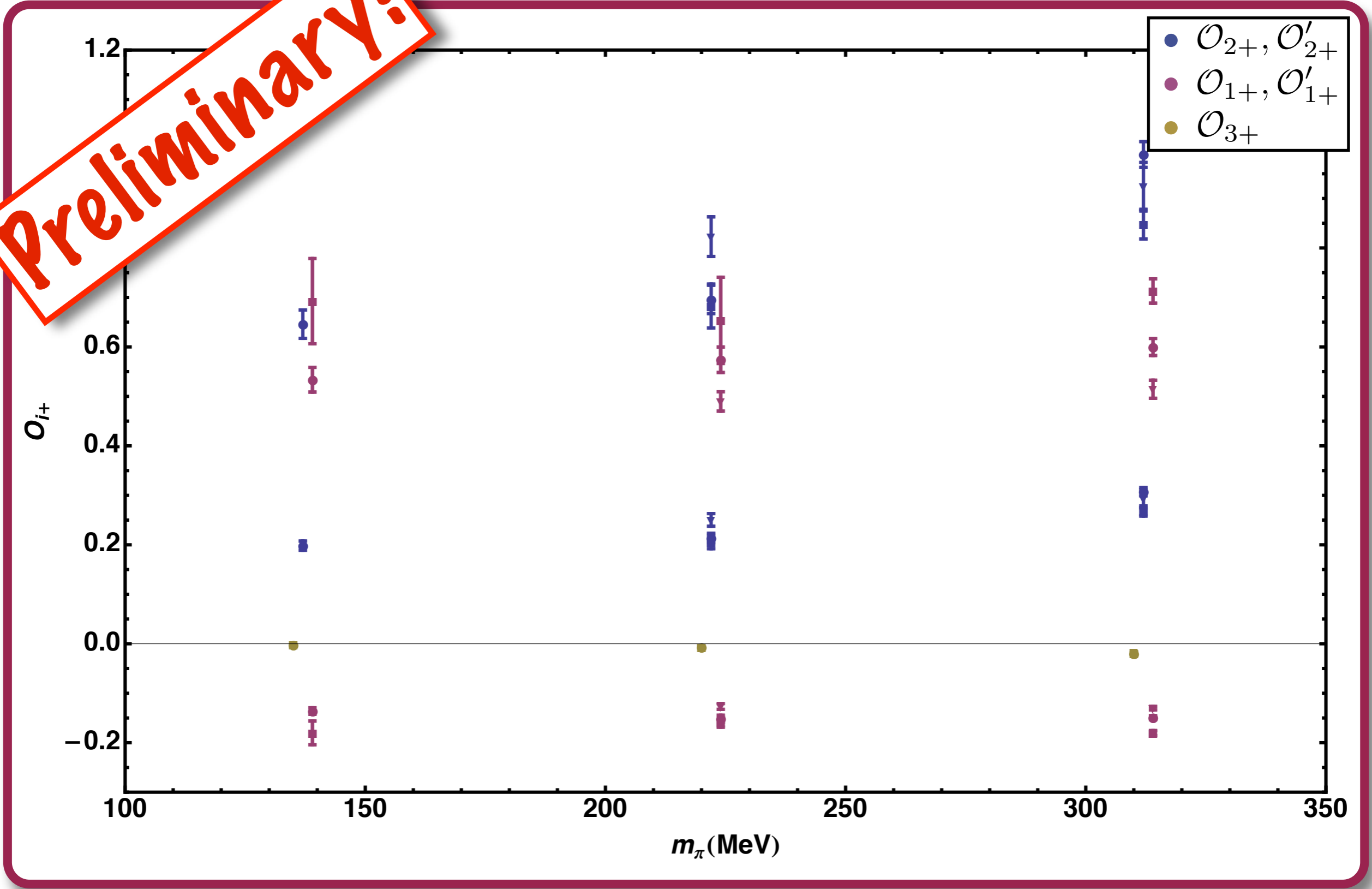
$m_{\pi} \sim 220$  MeV,  $a = 0.12$  fm

**Preliminary!**

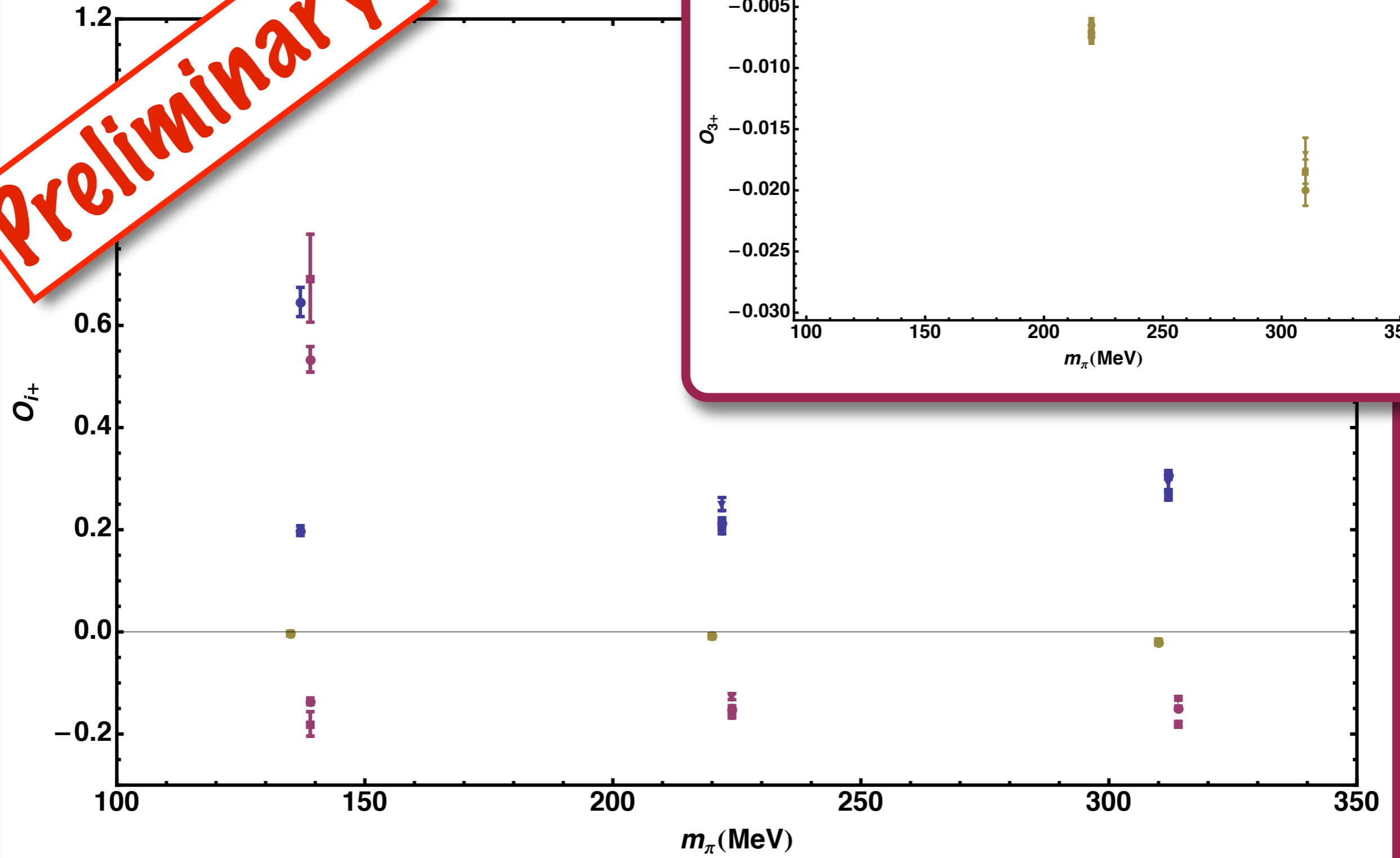


$m_\pi \sim 220$  MeV,  $a = 0.12$  fm

**Preliminary!**



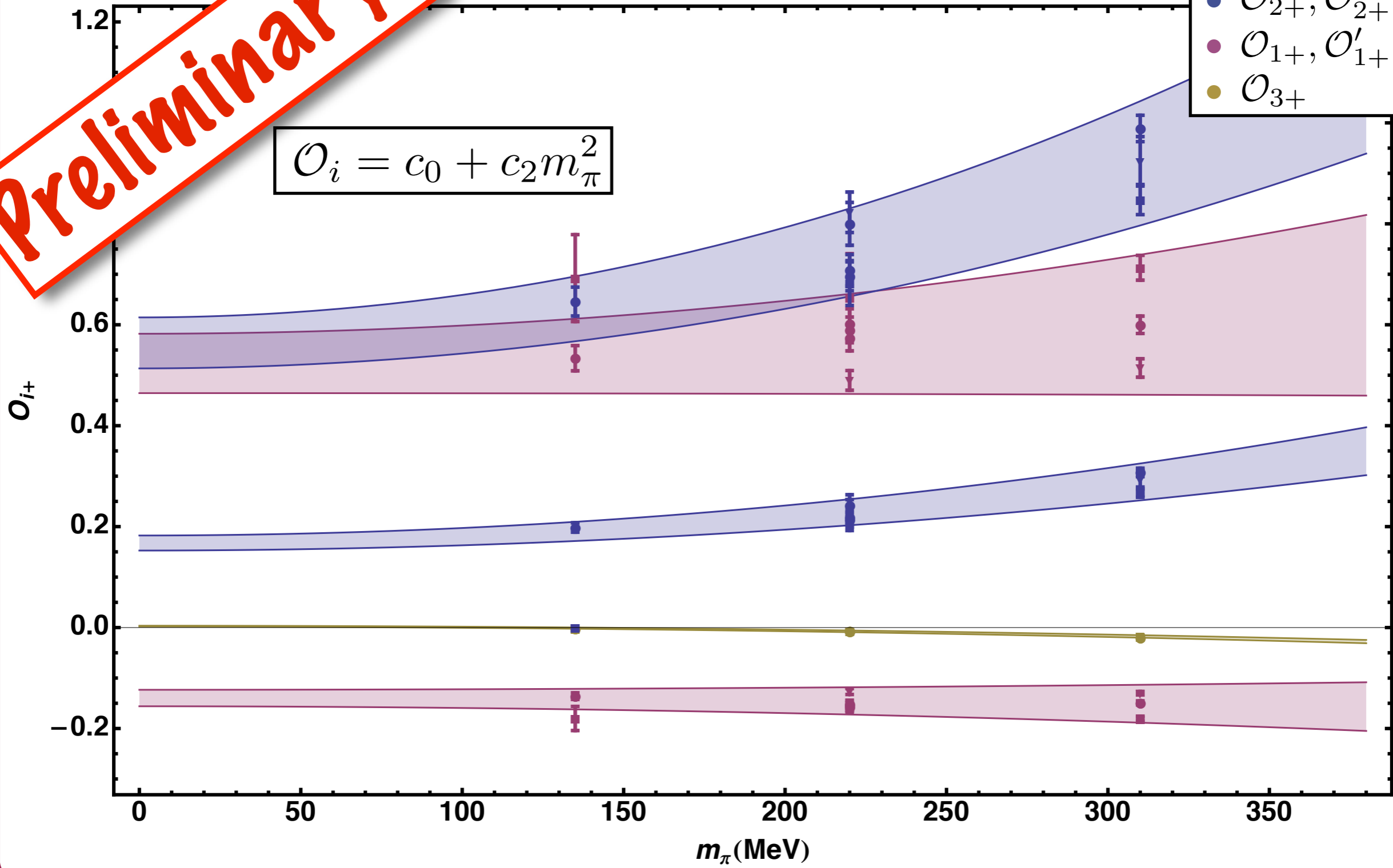
**Preliminary!**



**Preliminary!**

$$\mathcal{O}_i = c_0 + c_2 m_\pi^2$$

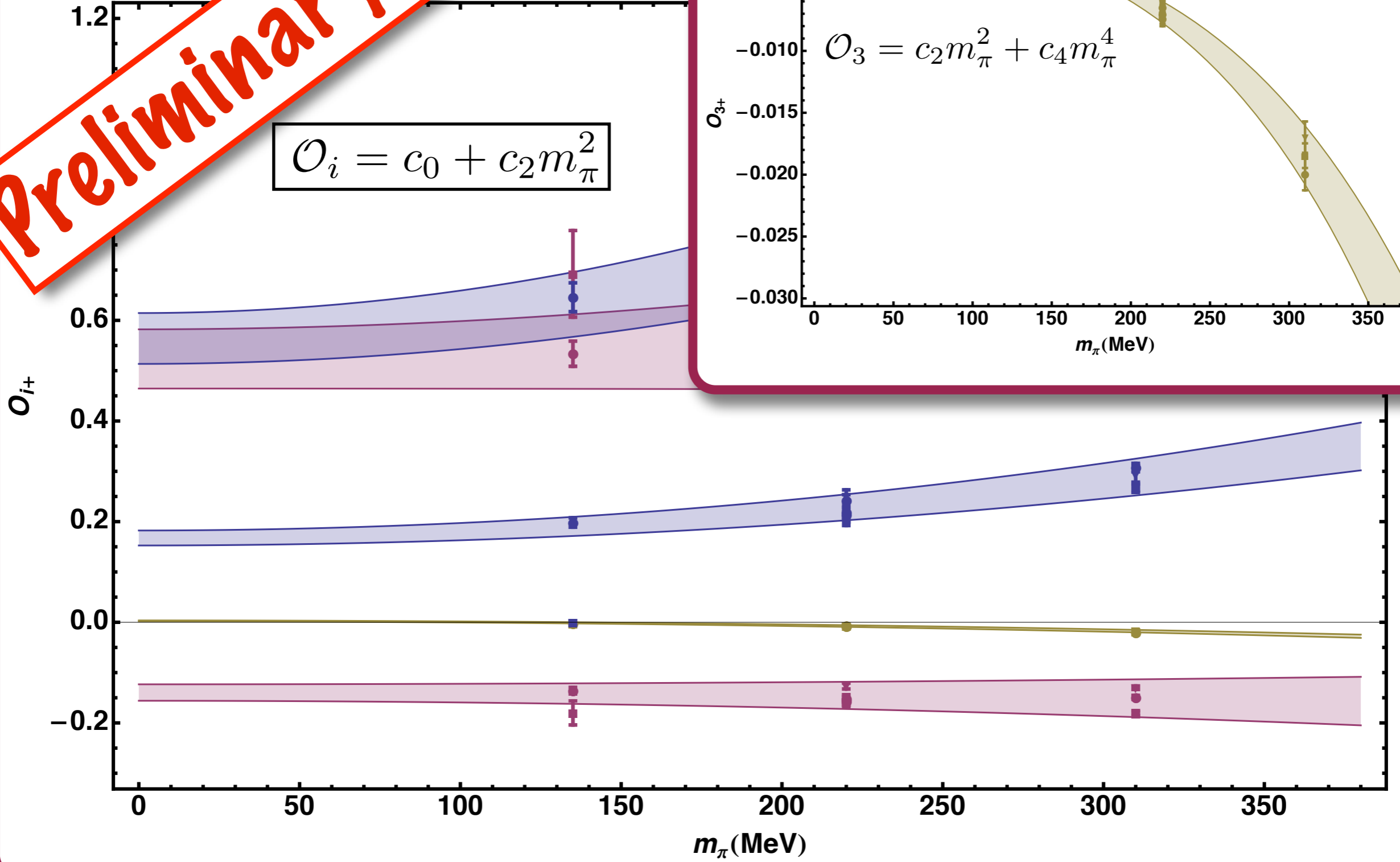
- $\mathcal{O}_{2+}, \mathcal{O}'_{2+}$
- $\mathcal{O}_{1+}, \mathcal{O}'_{1+}$
- $\mathcal{O}_{3+}$





**Preliminary!**

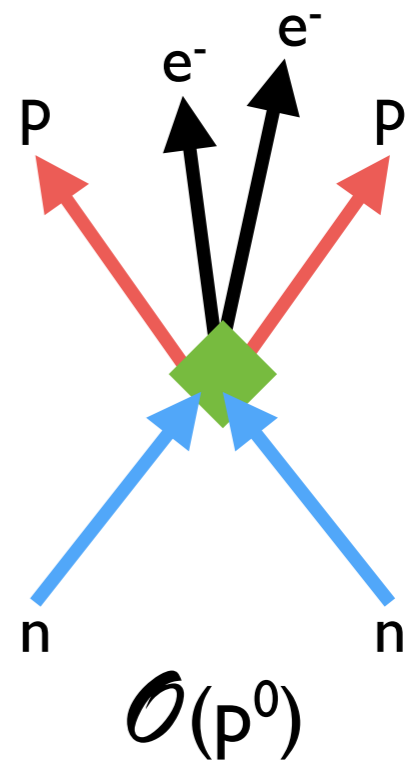
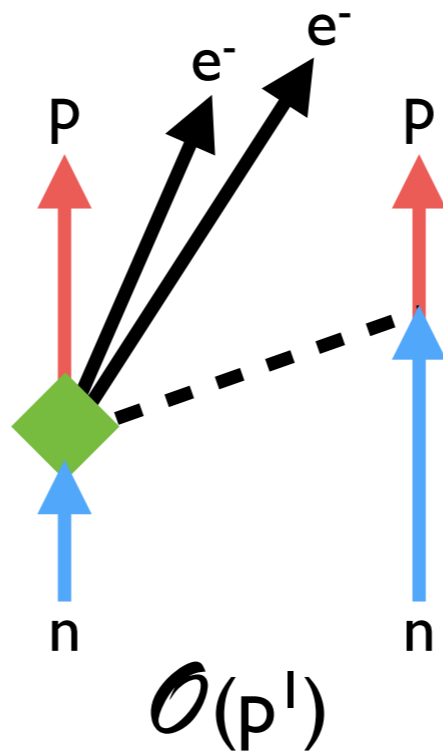
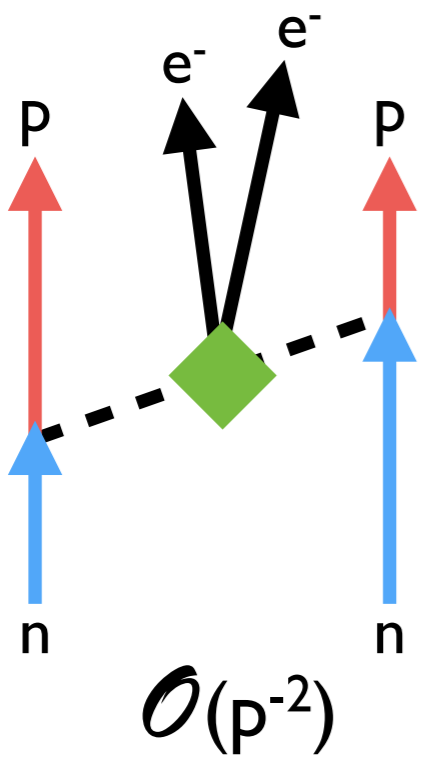
$$\mathcal{O}_i = c_0 + c_2 m_\pi^2$$



# Summary

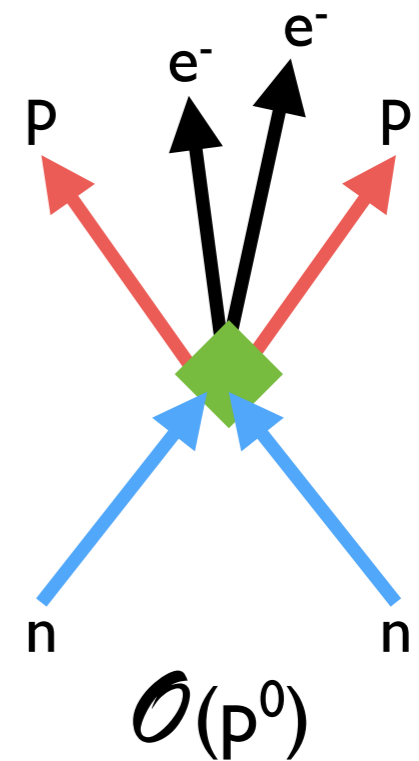
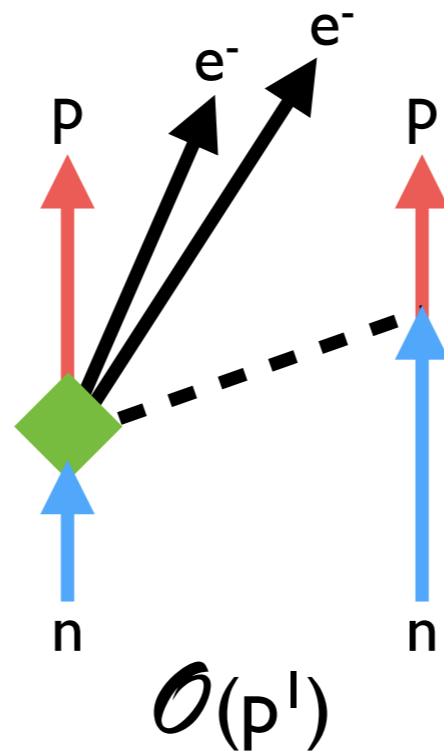
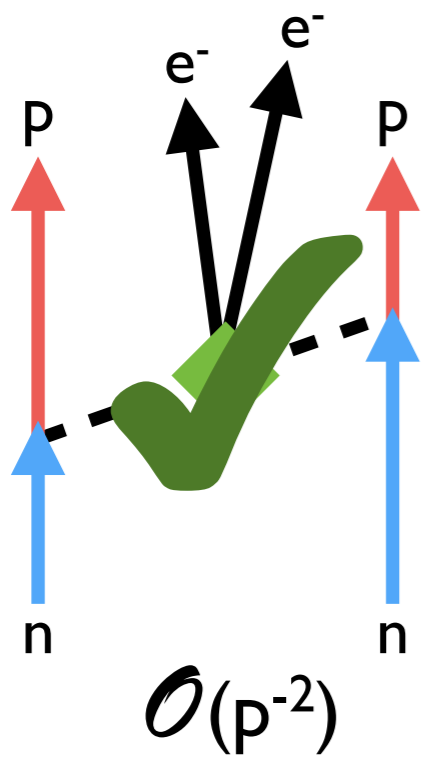
- $0\nu\beta\beta$ : search for Majorana mass signature
  - Lepton number violation could be source of matter / anti-matter asymmetry
  - Huge experimental efforts planned / underway
  - LQCD can make major impact on understanding of short-range operators
- Preliminary results for  $\pi^- \rightarrow \pi^+$  matrix element
  - Multiple pion masses, lattice spacings, volumes
  - Pion mass dependence as expected from chiral EFT counting
- To do:
  - Renormalization Buras, Misiak, Urban (2000), Tiburzi (2012)
  - Extrapolations in pion mass / lattice spacing
  - Other contact operators....

# Contact operators



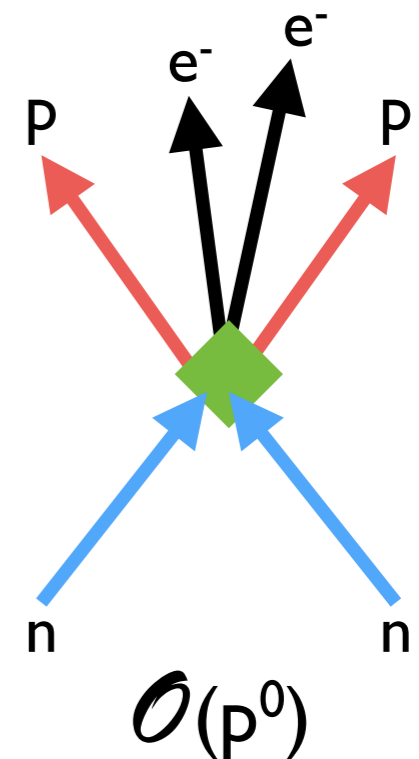
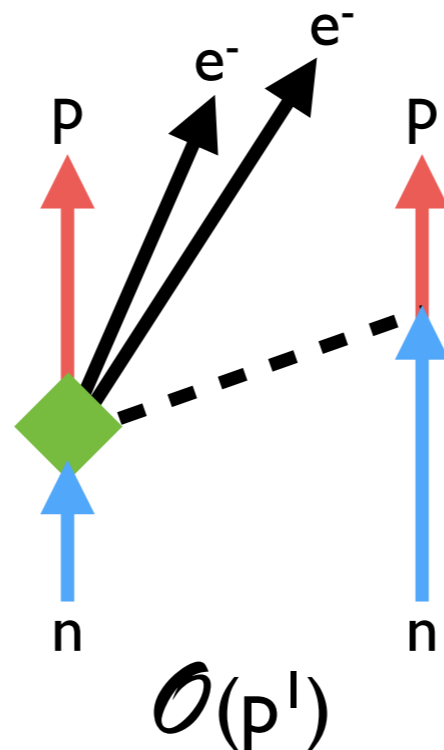
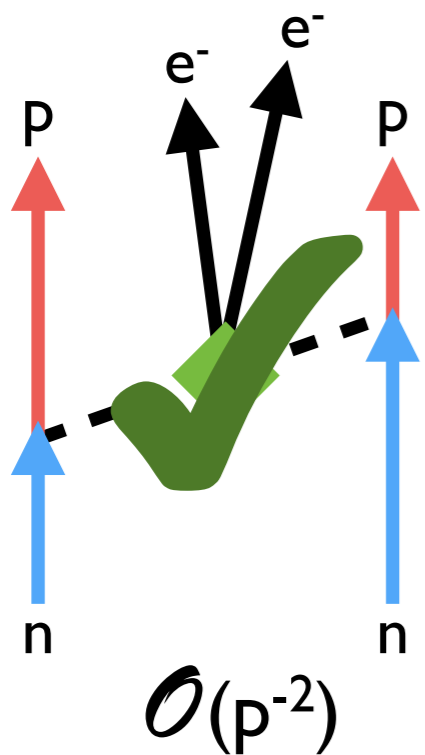
# Contact operators

- LO almost complete!



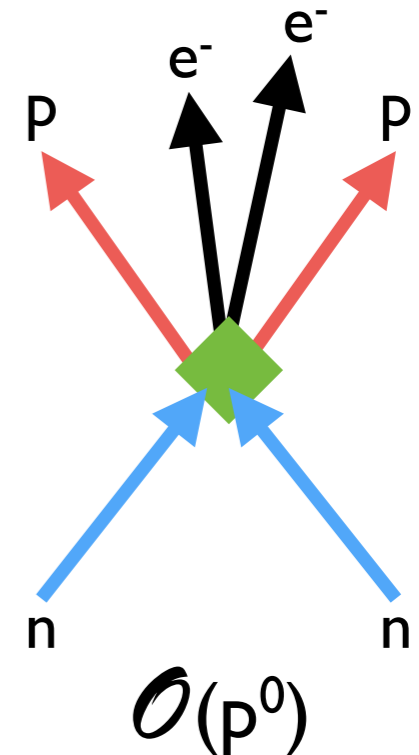
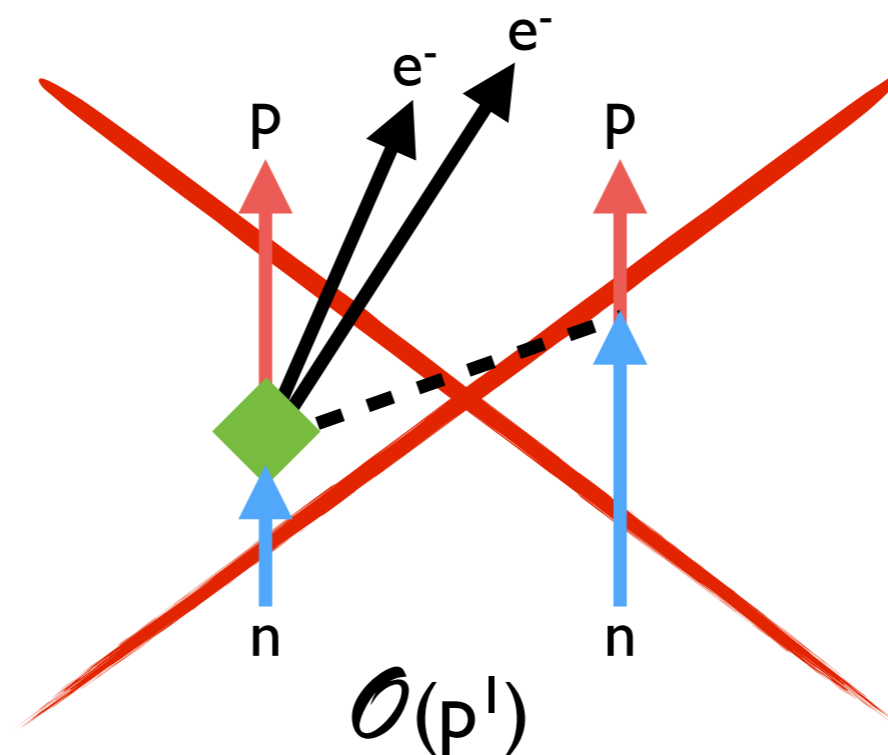
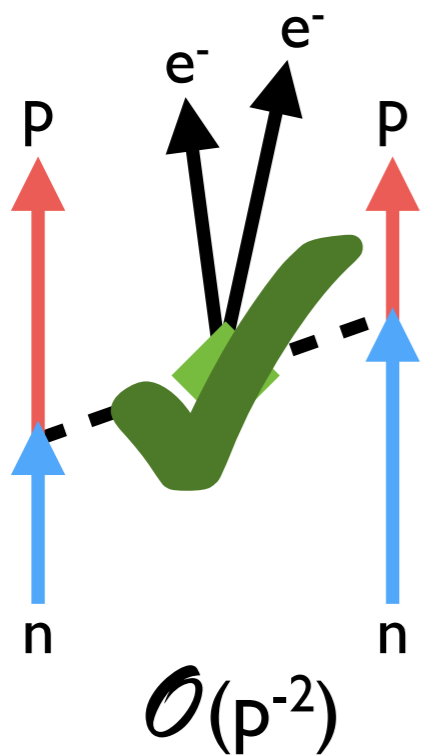
# Contact operators

- LO almost complete!
- NLO: **disconnected diagrams**



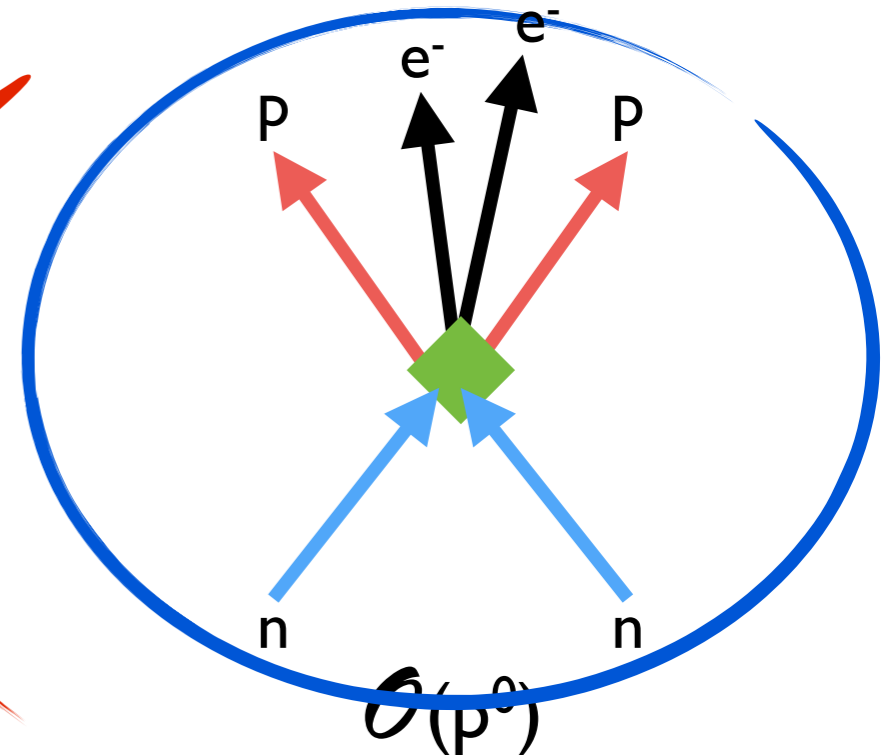
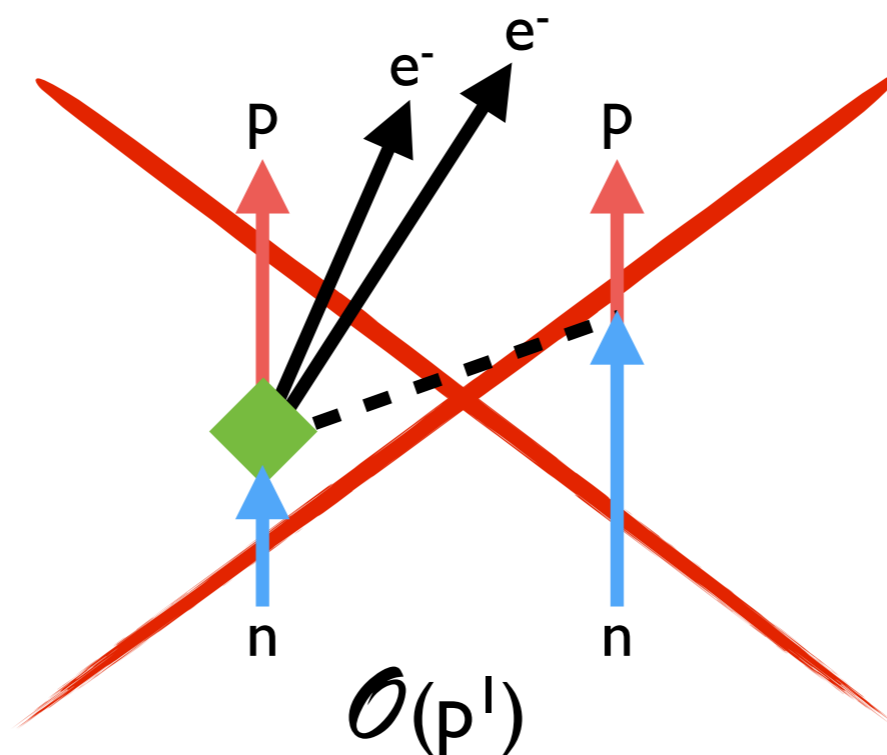
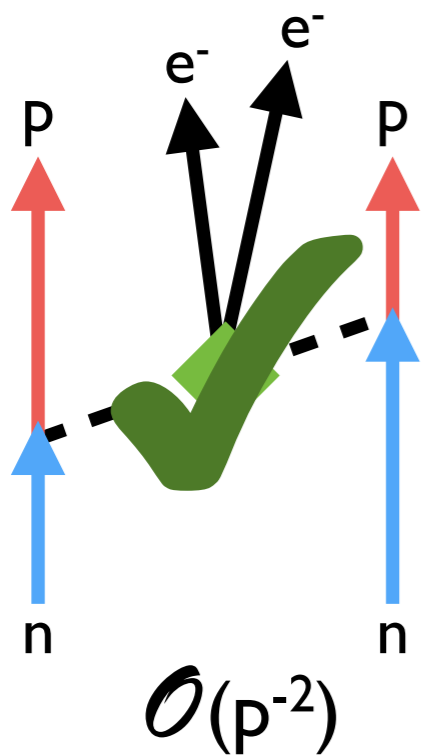
# Contact operators

- LO almost complete!
- NLO: **disconnected diagrams**
  - Don't contribute to  $0^+ \rightarrow 0^+$  nuclear transitions



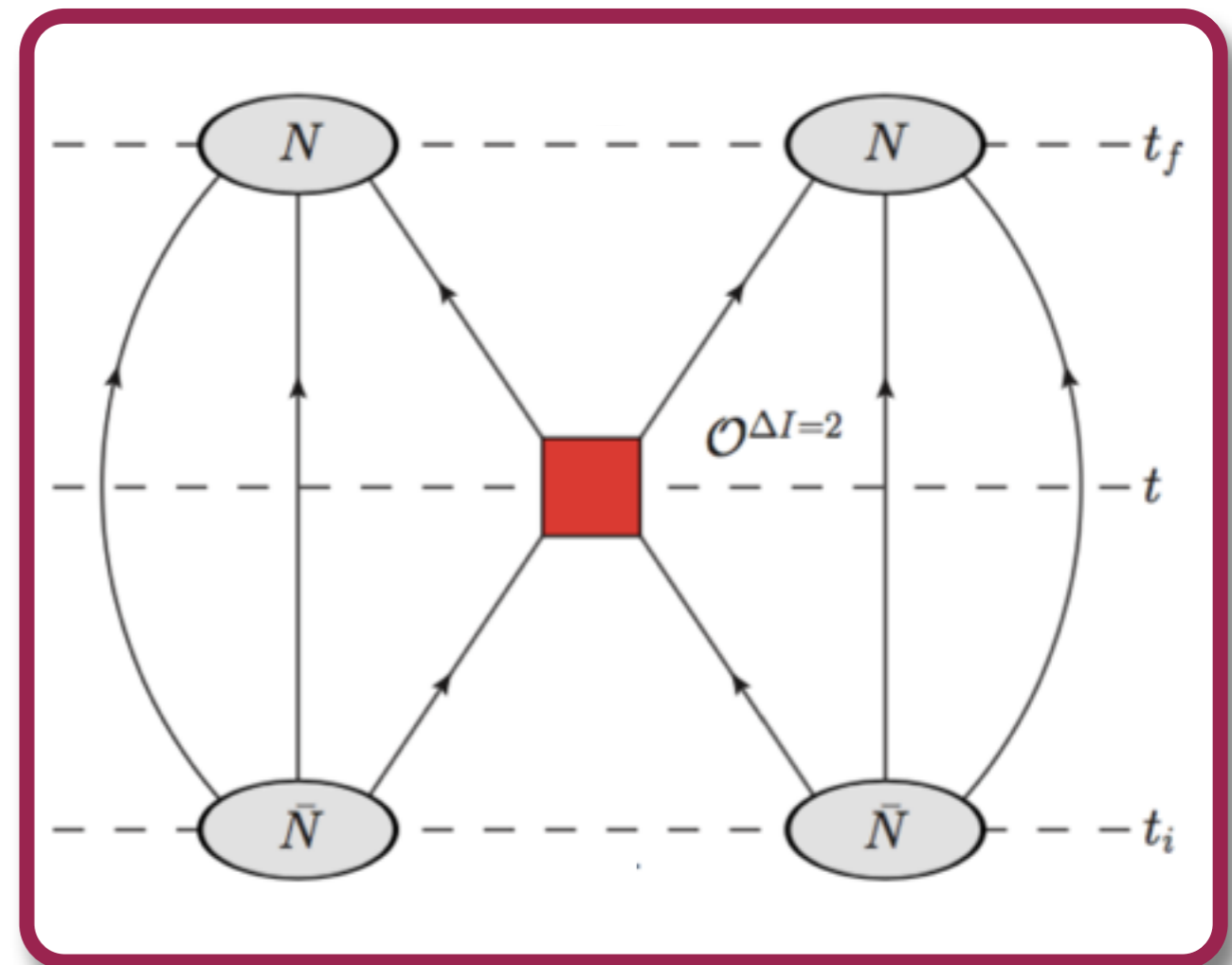
# Contact operators

- LO almost complete!
- NLO: **disconnected diagrams**
  - Don't contribute to  $0^+ \rightarrow 0^+$  nuclear transitions
- $nn \rightarrow pp$  contact operators

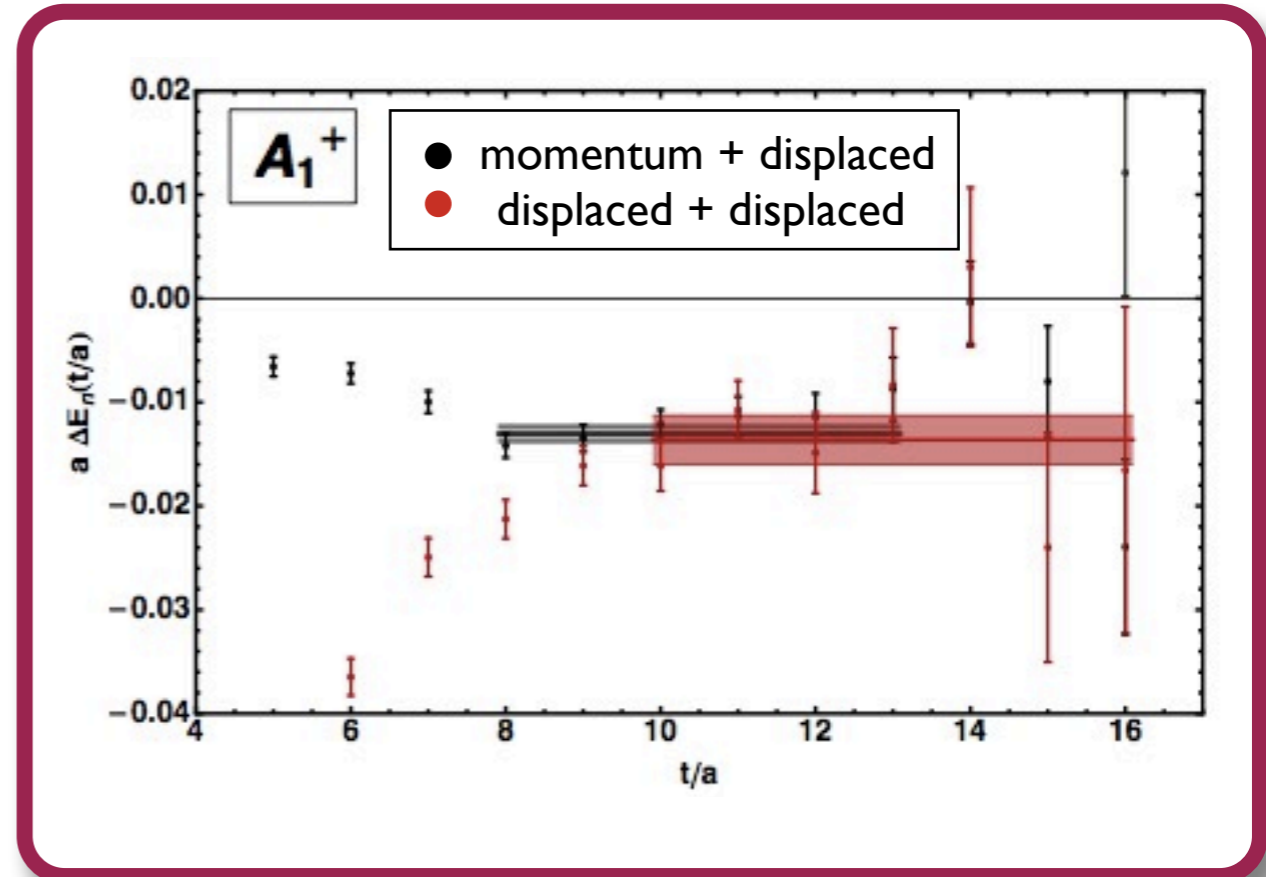
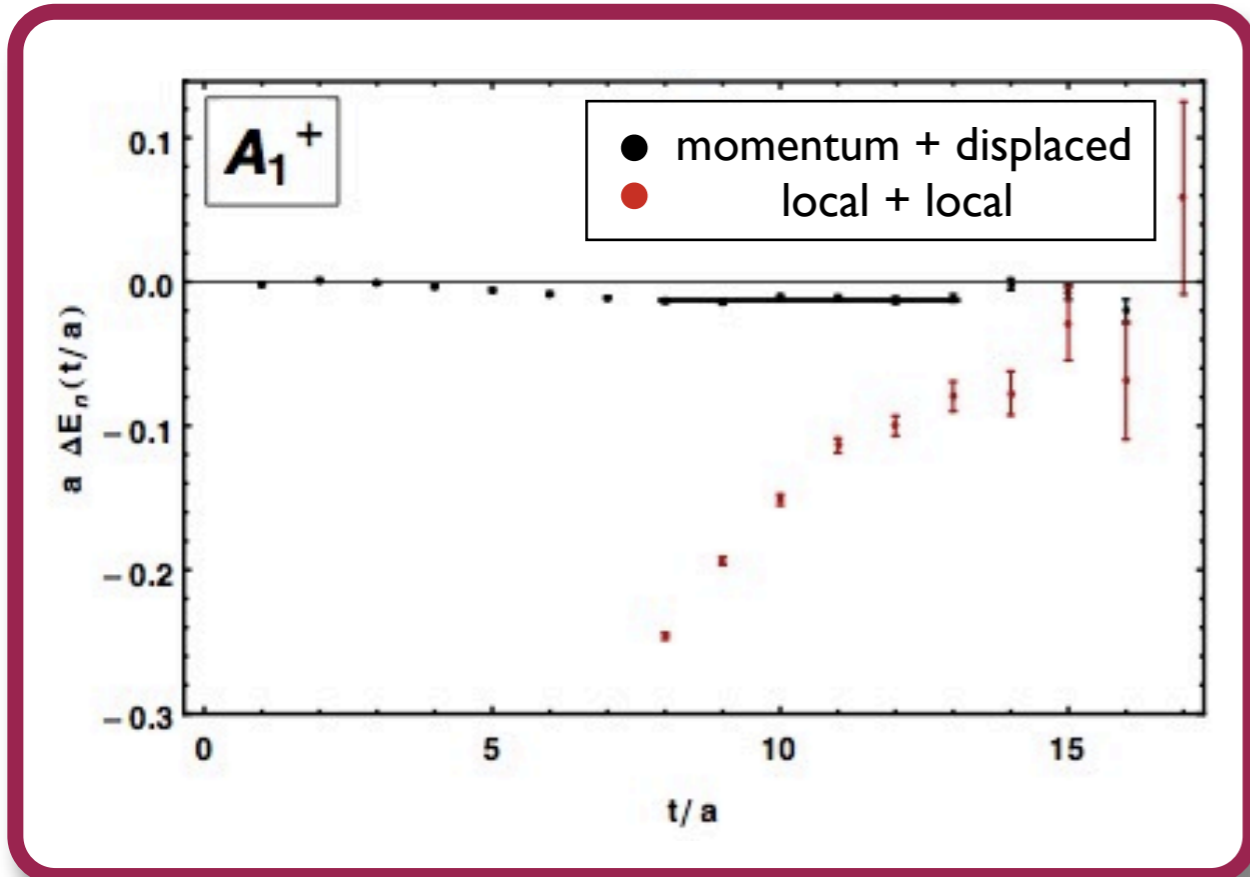
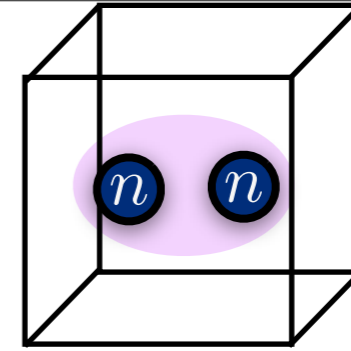
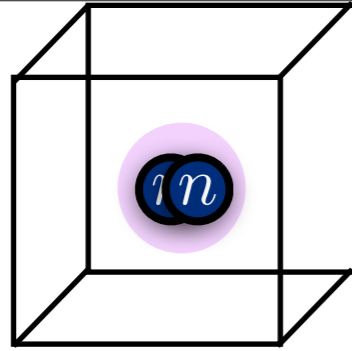


# Contractions

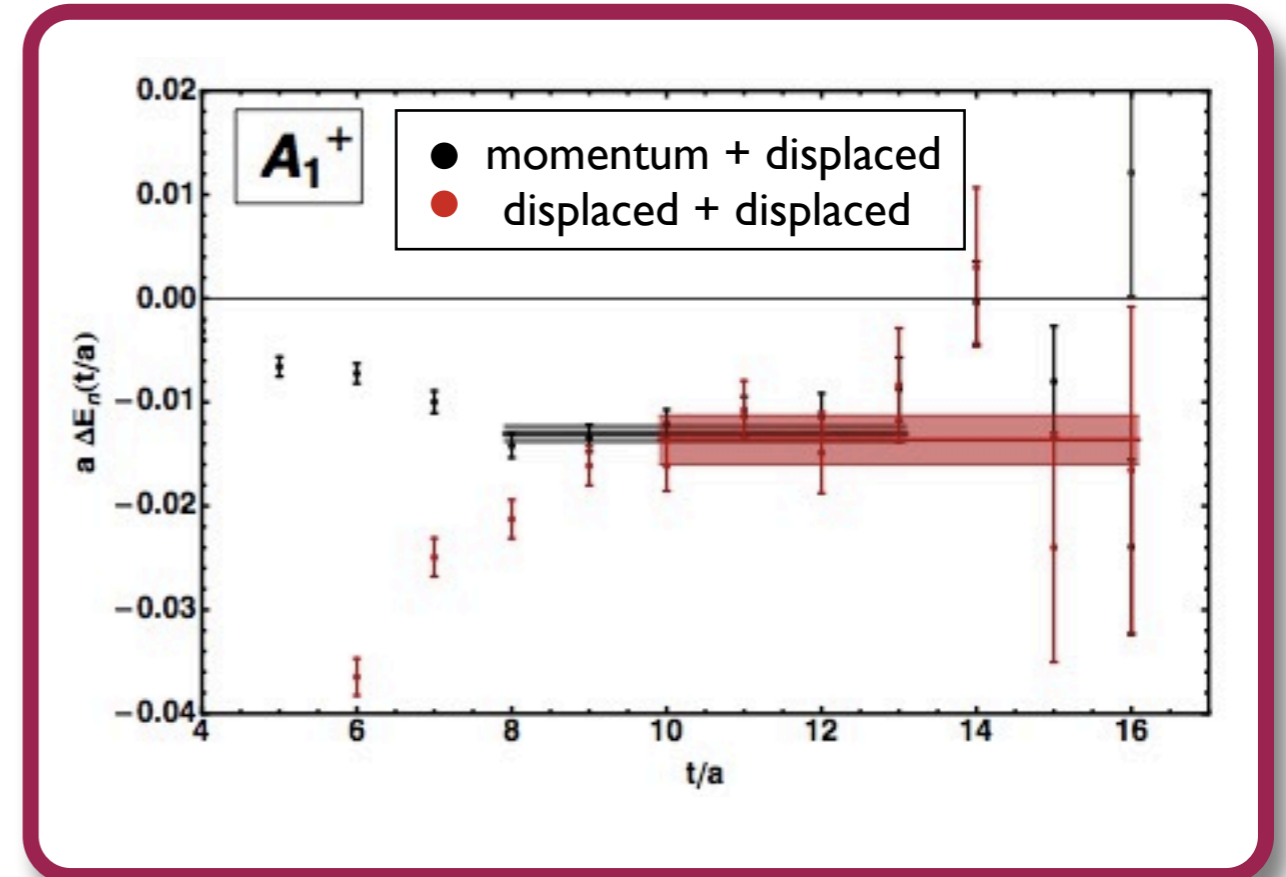
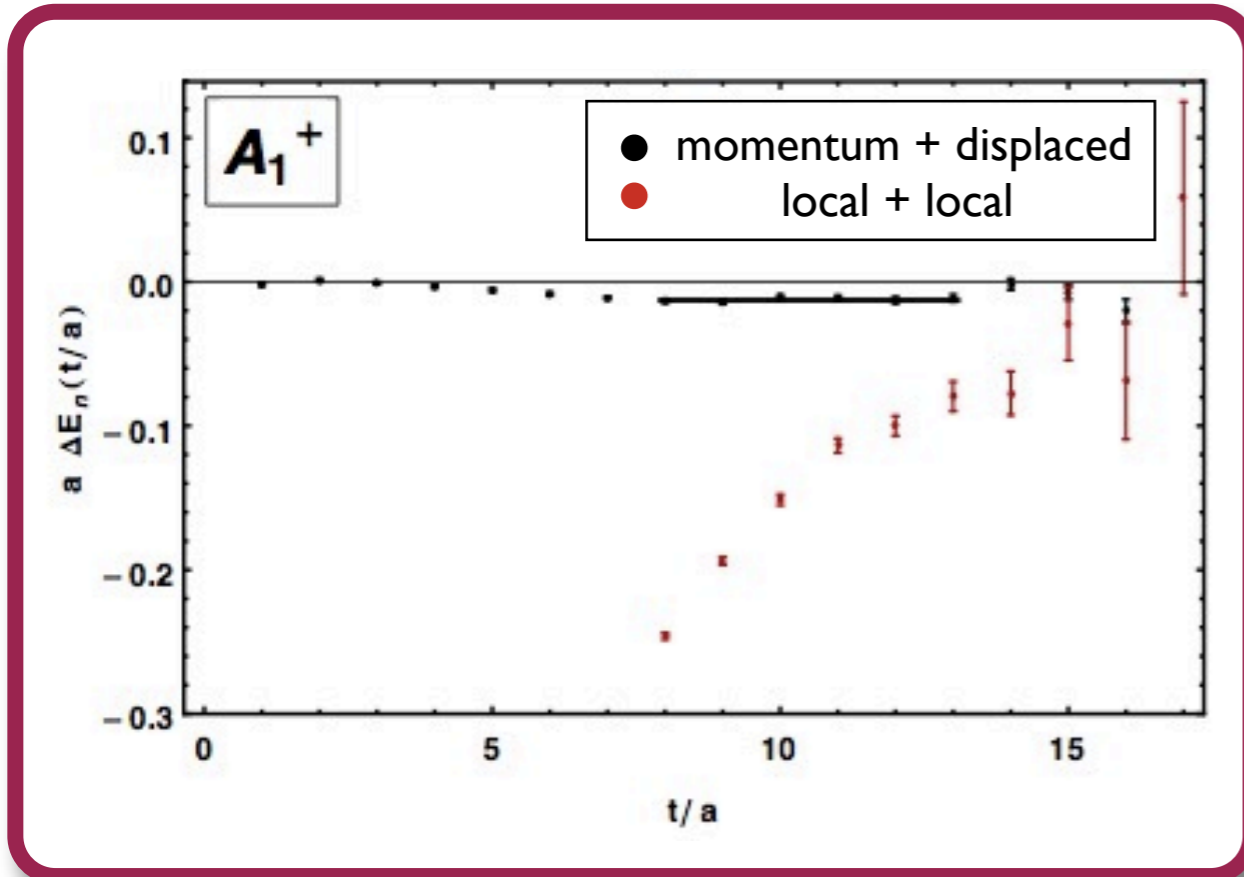
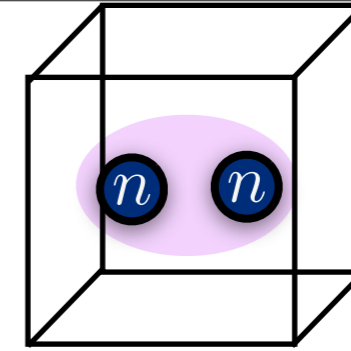
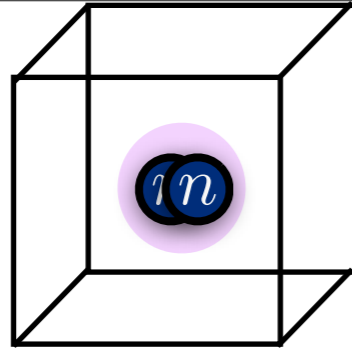
- Isospin limit: 576 contractions\*
- Baryon signal-to-noise problem, small excited state energy splittings, ....
- Need position space source & sink
  - otherwise all-to-all propagators connect to 4-quark operator
  - stochastically project onto zero total momentum







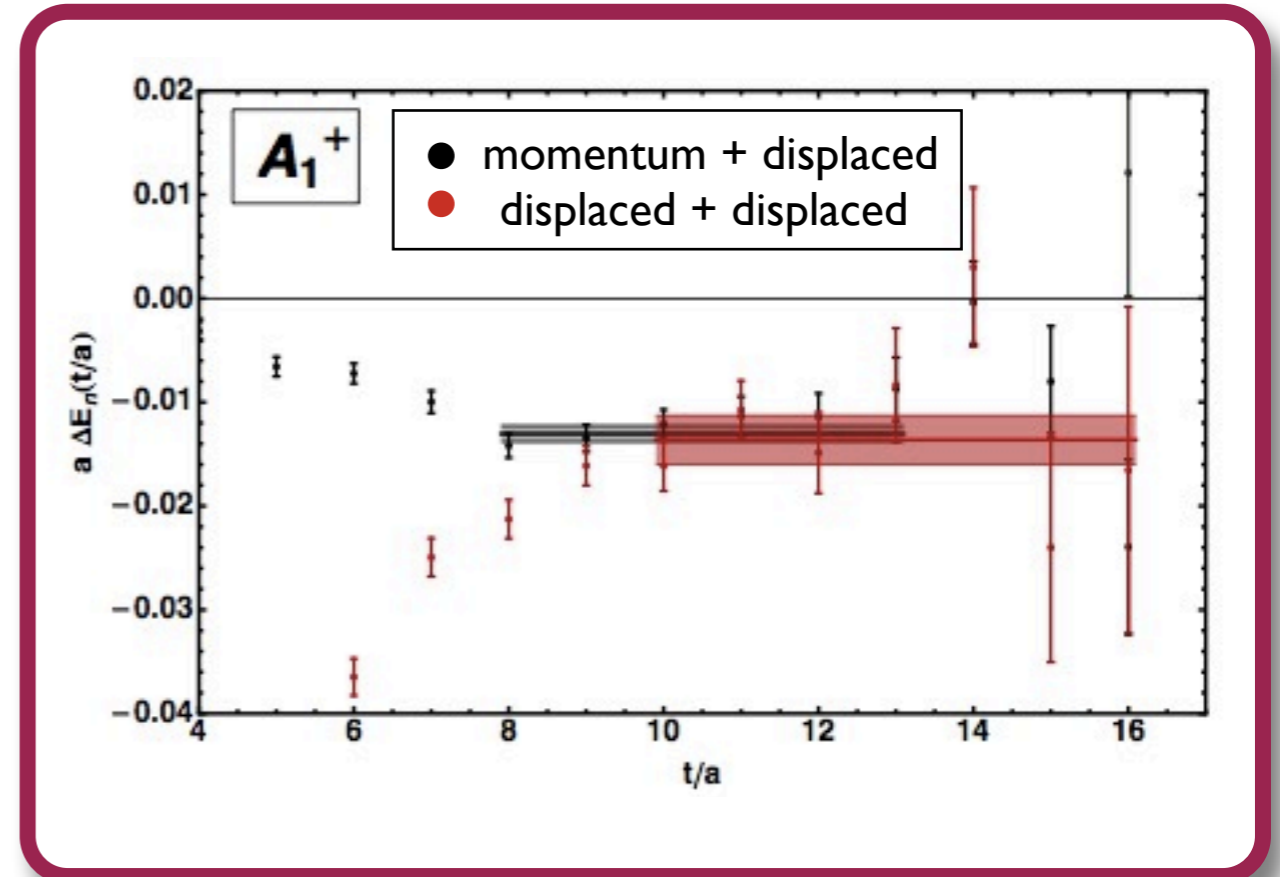
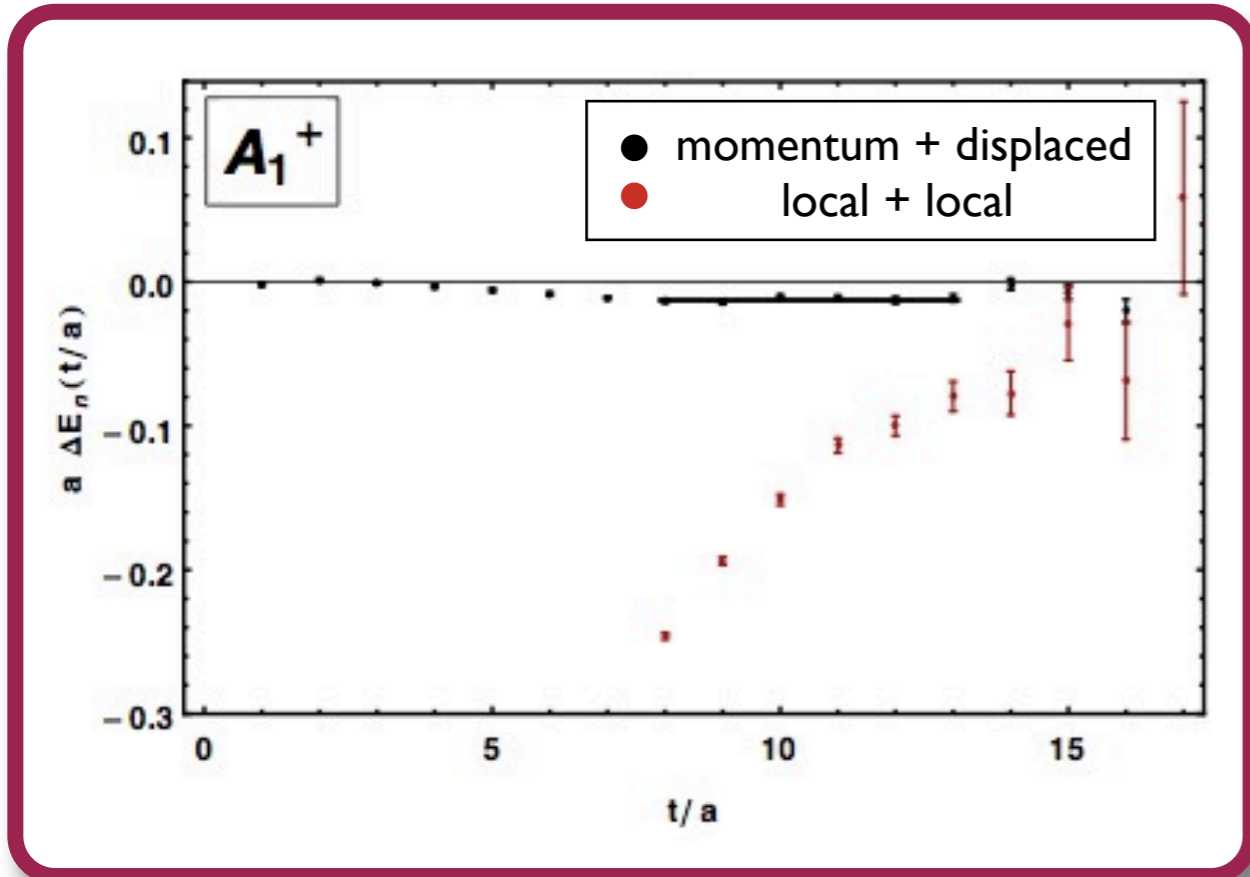
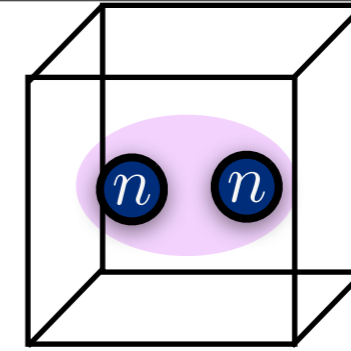
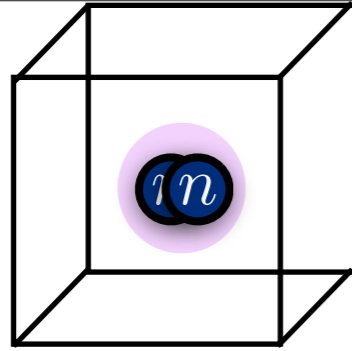
Iso-clover cfgs (W. Detmold,  
R. Edwards, D. Richards, K. Orginos)



Need displaced operators!

E. Berkowitz, T. Kurth, A.N., B. Joo, E. Rinaldi, M. Strother, P.Vranas, A. Walker-Loud arXiv:1508.00886 (2015)

Iso-clover cfgs (W. Detmold, R.Edwards, D. Richards, K. Orginos)

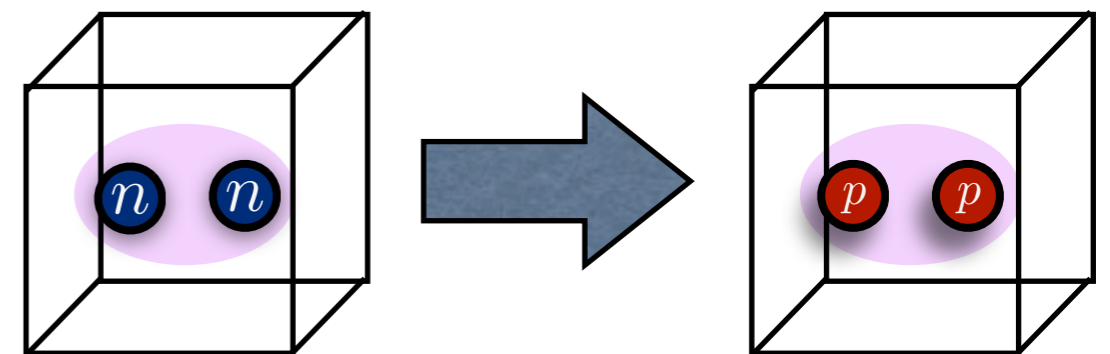


Finite volume formalism for  $2 \rightarrow 2$   
matrix elements completed:

R. Briceño, M. Hansen Phys.Rev. D94  
(2016) no.1, 013008

Renormalization known in  $\overline{MS}$ :

B. Tiburzi Phys.Rev. D86 (2012) 097501



*Stay tuned!*



- LBL/UCB: Chia Cheng Chang, AN, André Walker-Loud,
- LLNL: Evan Berkowitz, Enrico Rinaldi, Pavlos Vranas
- NERSC: Thorsten Kurth
- JLab: Balint Joo
- CCNY: Brian Tiburzi
- nVidia: Kate Clark

