

# Spatio-temporal chaos in rotating Rayleigh-Bénard convection

Nathan Becker

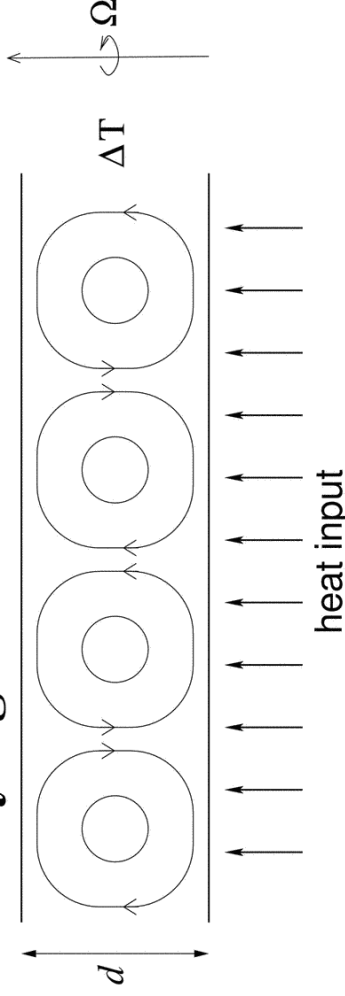
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## Rayleigh-Bénard convection



Convection occurs for

$$\Delta T > \Delta T_c$$

Reduced temperature difference

$$\varepsilon = \frac{\Delta T}{\Delta T_c} - 1$$

Reduced rotation frequency

$$\Omega = \frac{2\pi f d^2}{\nu}$$

Vertical thermal diffusion time

$$\tau_v = \frac{d^2}{\kappa}$$

## Movie of shadowgraph visualization of Küppers-Lortz dynamics



$$\varepsilon = 0.14$$

$$\Omega = 16.2$$

## Model equations

$$\tau \dot{A} = \varepsilon A + \xi^2 \nabla^2 A + N \cdot L.$$

Scaling laws near onset

$$\tau \propto \varepsilon^{-1} \quad \xi \propto \varepsilon^{-1/2}$$

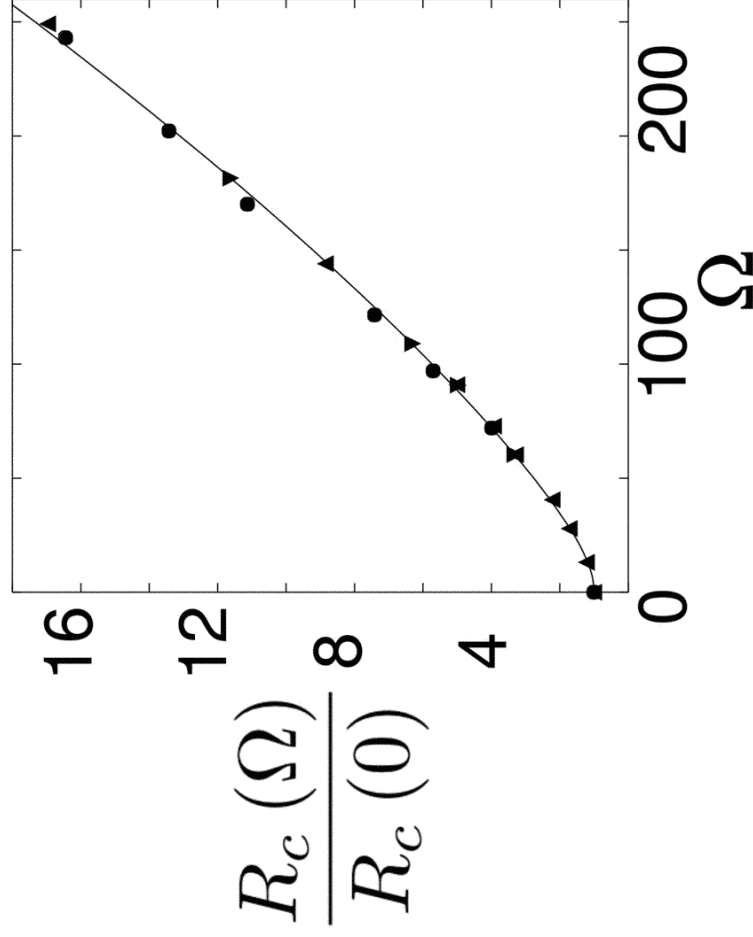
Y. Tu and M.C. Cross, Phys. Rev. Lett. **69** 2515 (1992)

## Rayleigh-Bénard convection

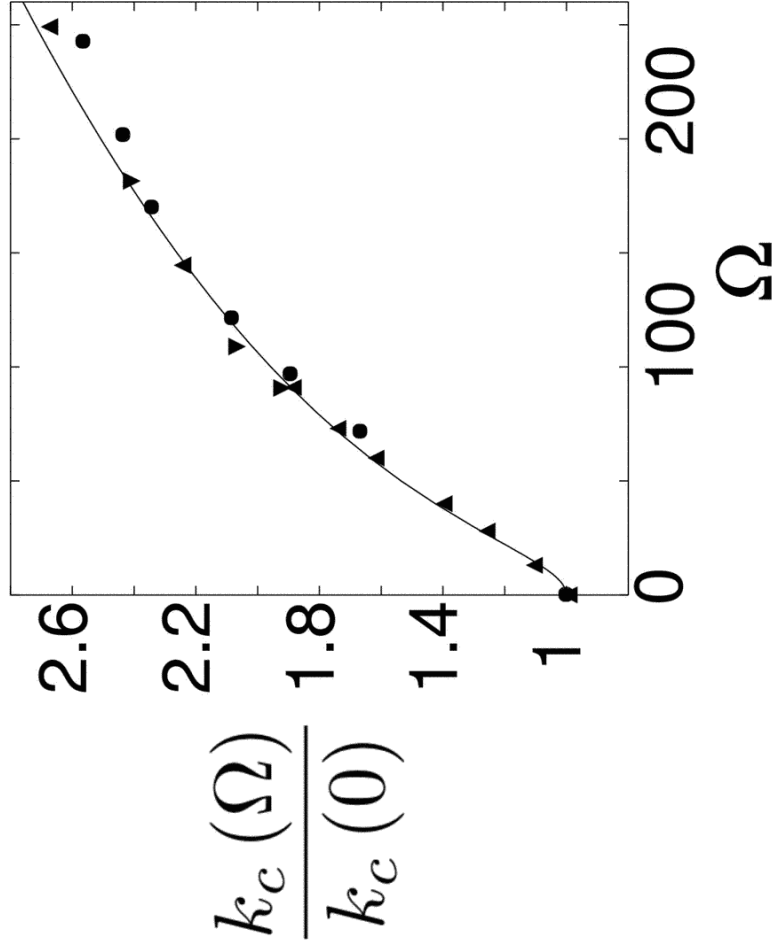
$$\text{Rayleigh number} \quad R = \frac{\alpha g d^3 \Delta T}{\kappa \nu}$$

$$\text{Wave-number} \quad k = \frac{2\pi d}{\lambda}$$

## Onset of convection (from linear stability)

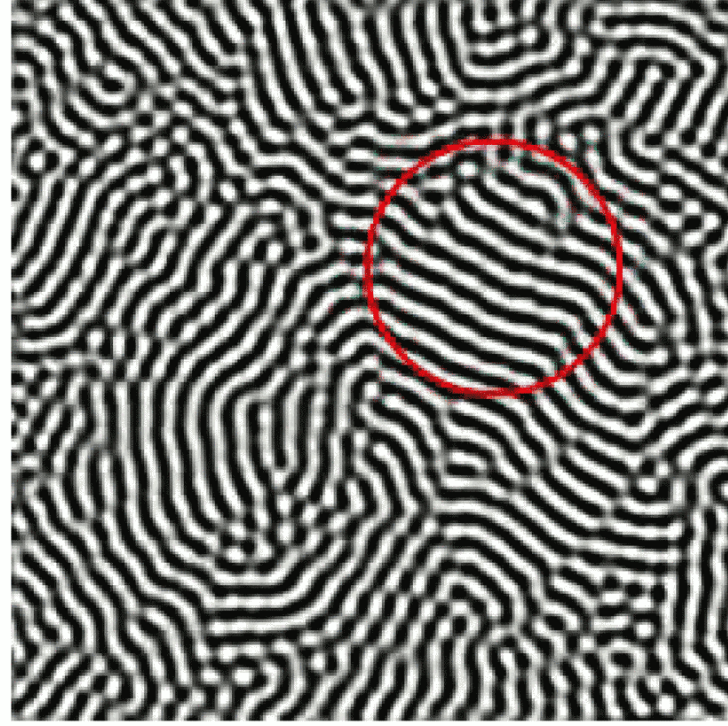


Wave-number at onset (from linear stability)

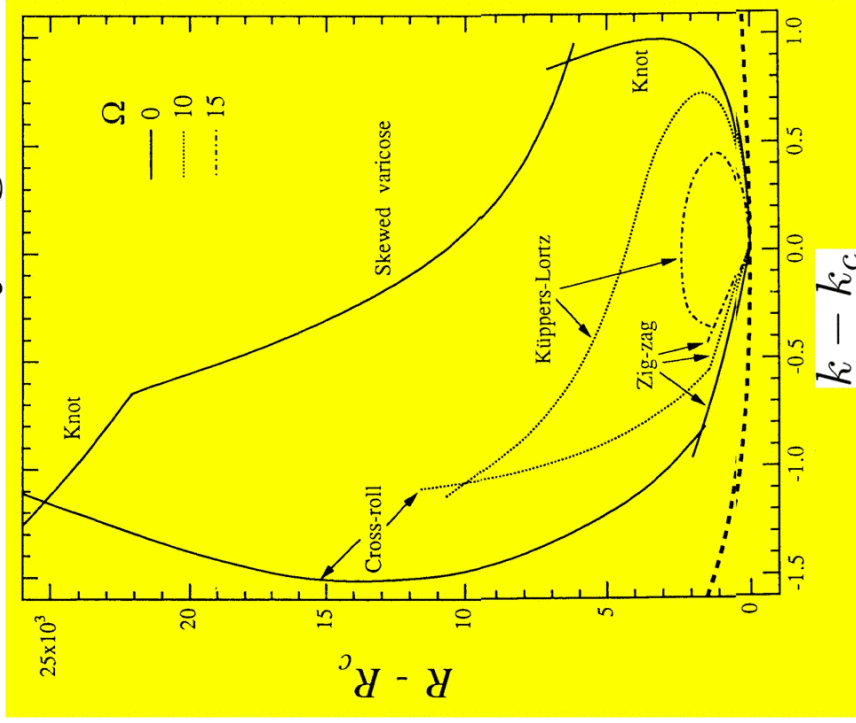


E. Bodenschatz, W. Pesch, and G. Ahlers, *Annu. Rev. Fluid Mech.* **32**, 709 (2000)

Küppers-Lortz bulk instability



## Instability regimes - Busse Balloon



The Küppers-Lortz unstable region shrinks to zero at  $\Omega_c$ . This results in the straight rolls constantly switching because the modes that grow due to instabilities are themselves unstable to the same instabilities.

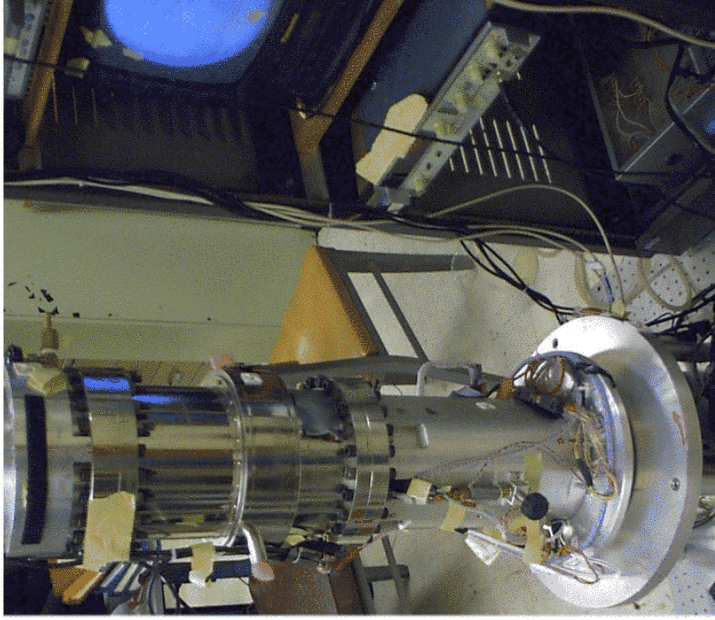
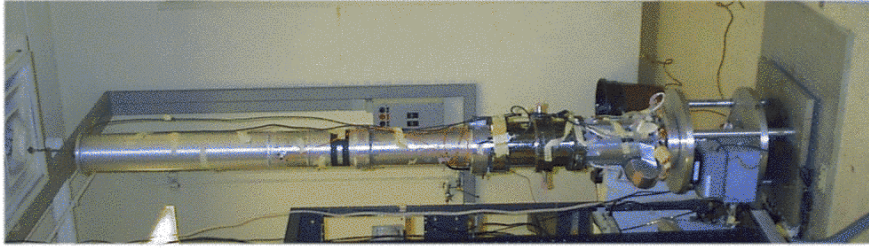
From: F. Zhong, R. Ecke, and V. Steinberg, *Physica D* **51** 596 (1991)

## Domain wall motion



E. Bodenschatz, D.S. Cannell, J. R. de Bruyn, R. Ecke, Y. Hu, K. Lerman, and G. Ahlers, *Physica D* **61**, 77 (1992)

**Photos of the rotating gas convection apparatus**



**Fourier analysis**

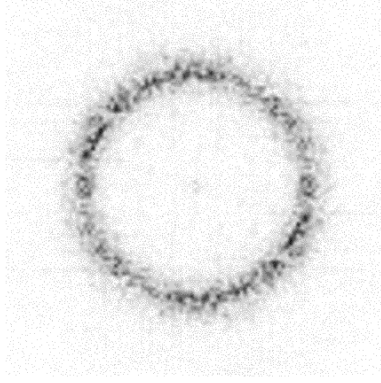


Realspace  
image

**Fourier analysis**



Realspace  
image

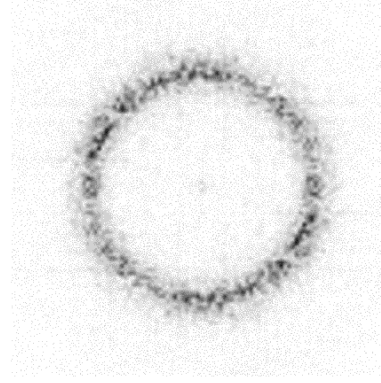


$|FT|^2$

**Fourier analysis**

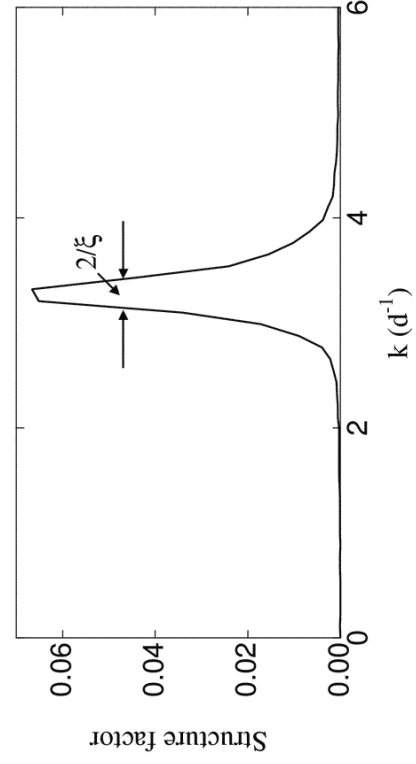


Realspace  
image

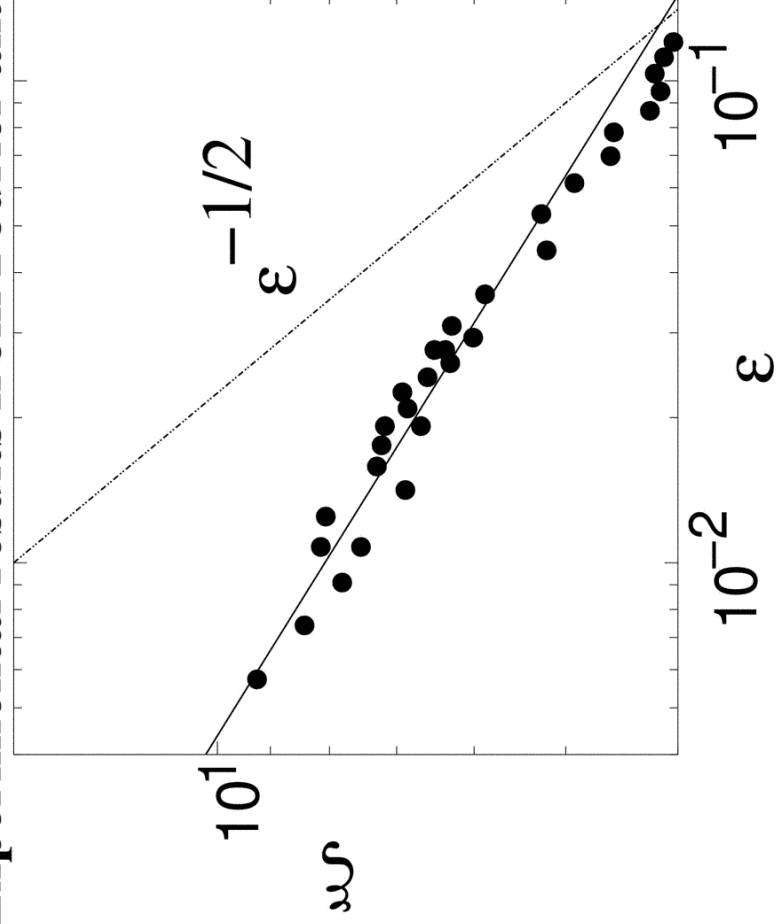


$|FT|^2$

Azimuthal  
average of  
 $|FT|^2$



**Experimental results from Fourier analysis**



Y. Hu, R. Ecke, G. Ahlers, Phys. Rev. Lett. 74, 5040 (1995)

**Küppers-Lortz Domains**



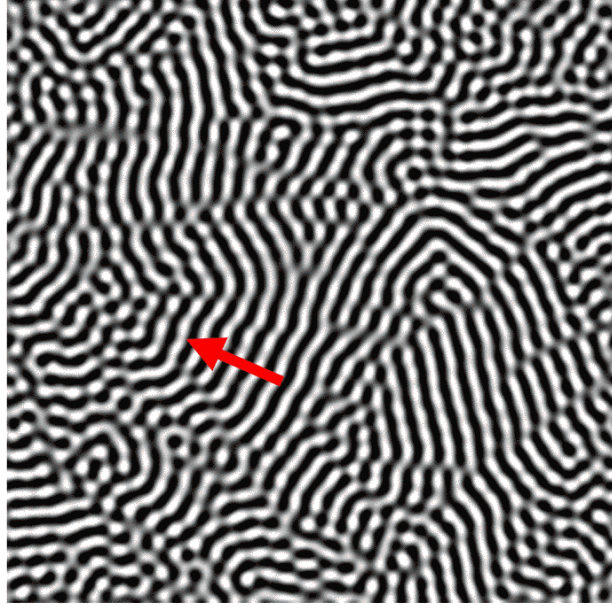
$\epsilon = 0.125$   
 $\Omega = 16.2$



## Küppers-Lortz Domains

$$\varepsilon = 0.125$$

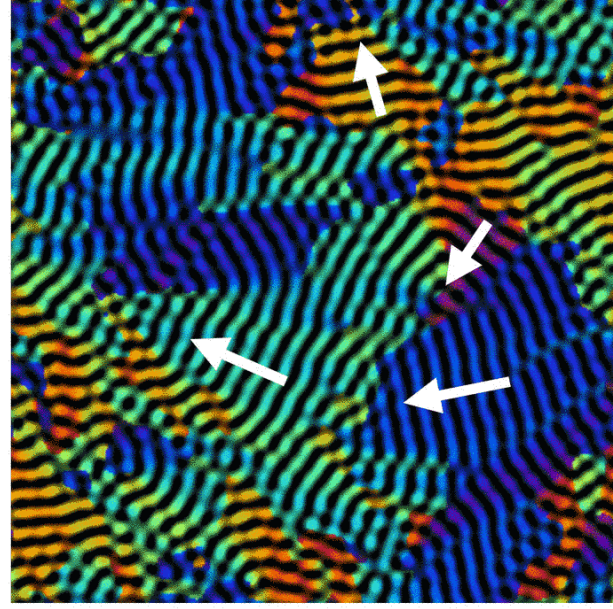
$$\Omega = 16.2$$



## Küppers-Lortz Domains with local wave-director field

$$\varepsilon = 0.125$$

$$\Omega = 16.2$$



M.C. Cross, D. Meiron, and Y. Tu, *Chaos* **4**(4), 607 (1994).

D. Egolf, I. Melnikov, E. Bodenschatz, *Phys. Rev. Lett.* **80**, 3228 (1998).

**Autocorrelation function of the angular component of the local wave-director field**

$$C(\delta \mathbf{x}) = \int \cos 2[\theta(\mathbf{x}) - \theta(\mathbf{x} + \delta \mathbf{x})] d\mathbf{x}$$

**Autocorrelation function of the angular component of the local wave-director field**

$$C(\delta \mathbf{x}) = \int \cos 2[\theta(\mathbf{x}) - \theta(\mathbf{x} + \delta \mathbf{x})] d\mathbf{x}$$

- Use a cosine function so that the correlation varies smoothly from -1 to 1.
- Compare angles by looking at their difference.

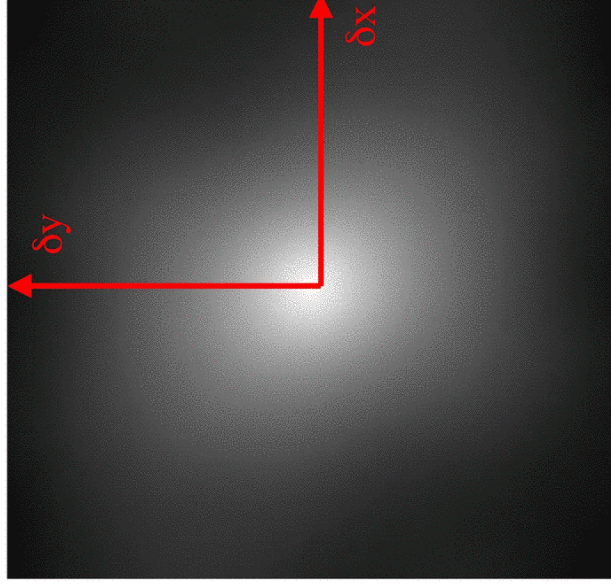
$\cos(2 \times 0) = 1$        $\cos(2 \times \pi) = 1$

$\cos\left(2 \times \frac{\pi}{2}\right) = -1$

## Autocorrelation function of the angular component of the local wave-director field

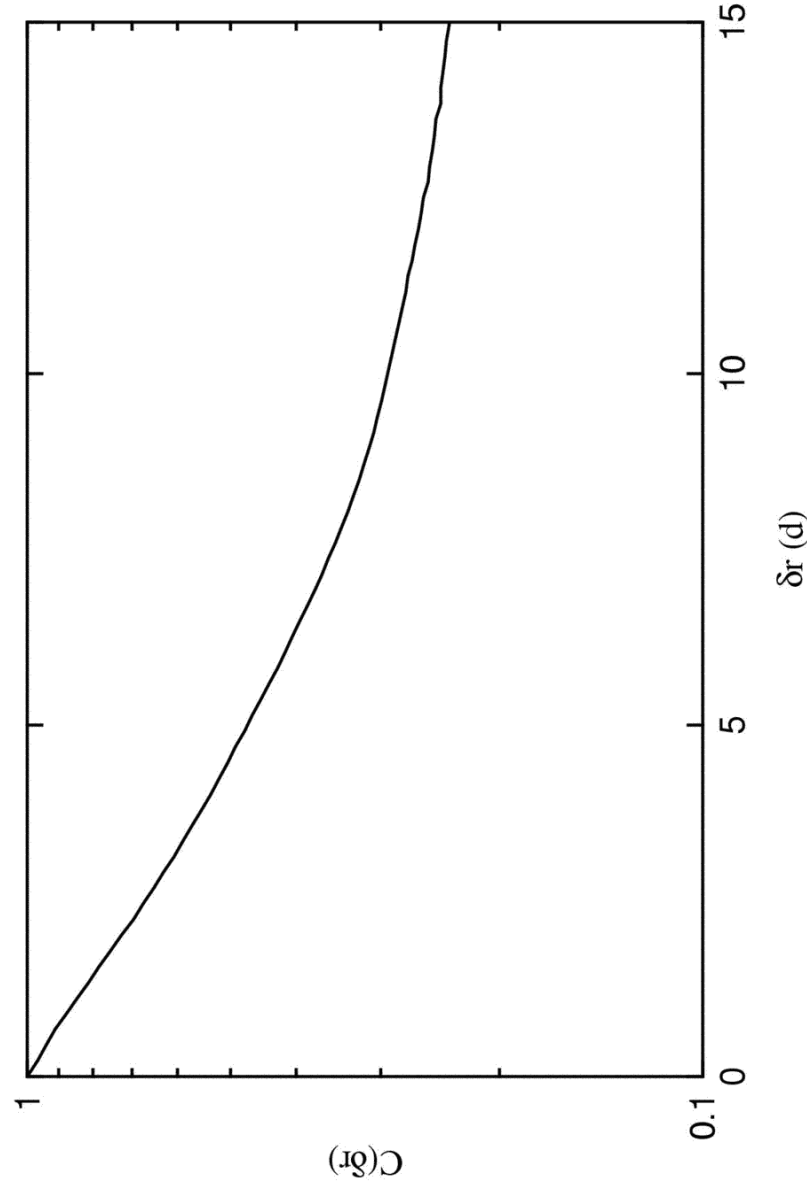
$$C(\delta\mathbf{x}) = \int \cos 2 [\theta(\mathbf{x}) - \theta(\mathbf{x} + \delta\mathbf{x})] d\mathbf{x}$$

- Use a cosine function so that the correlation varies smoothly from -1 to 1.
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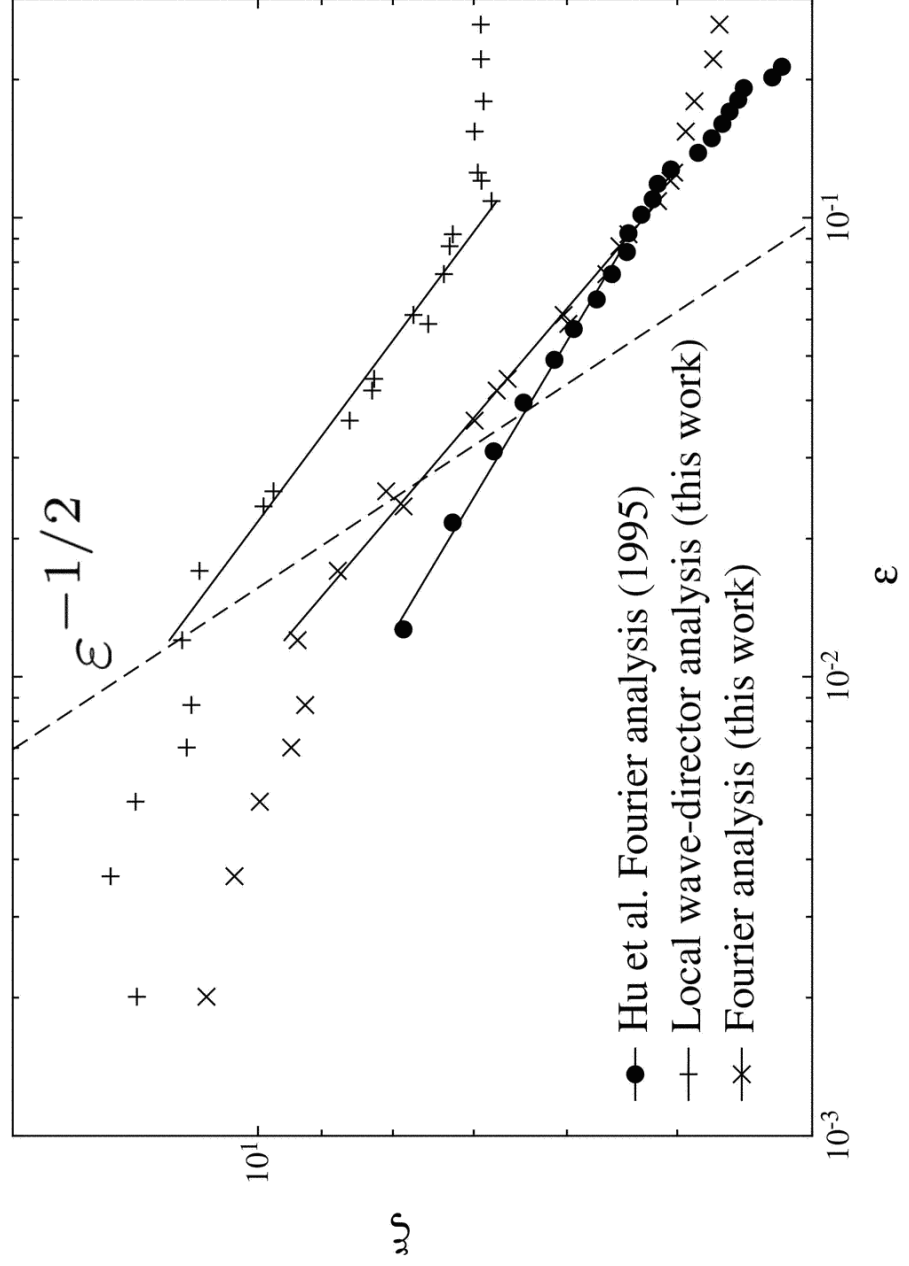
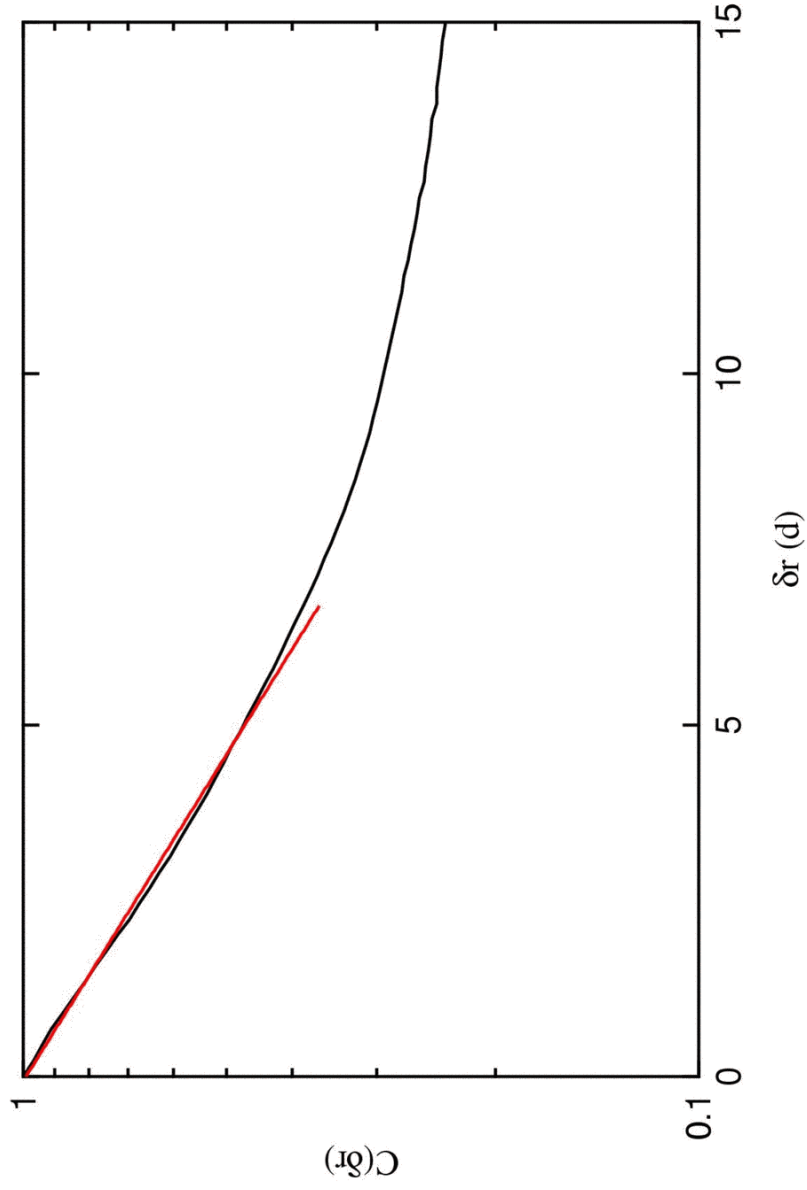


$$\begin{aligned} \cos(2 \times 0) = 1 & \quad \left\| \begin{array}{c} \uparrow \\ \uparrow \end{array} \right\| \cos(2 \times \pi) = 1 \\ \cos\left(2 \times \frac{\pi}{2}\right) = -1 & \quad \left\| \begin{array}{c} \uparrow \\ \rightarrow \end{array} \right\| \end{aligned}$$

## Azimuthal average of autocorrelation



**Azimuthal average of autocorrelation**



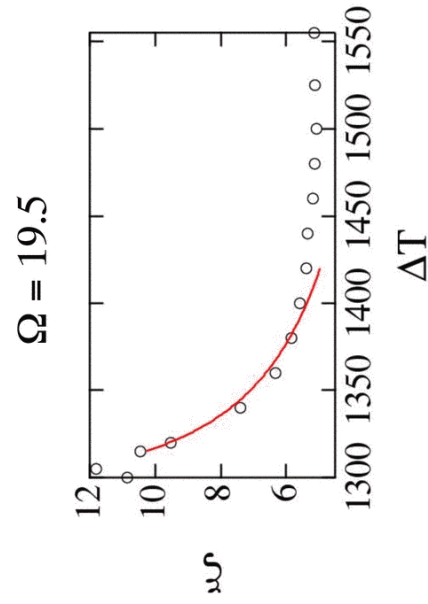
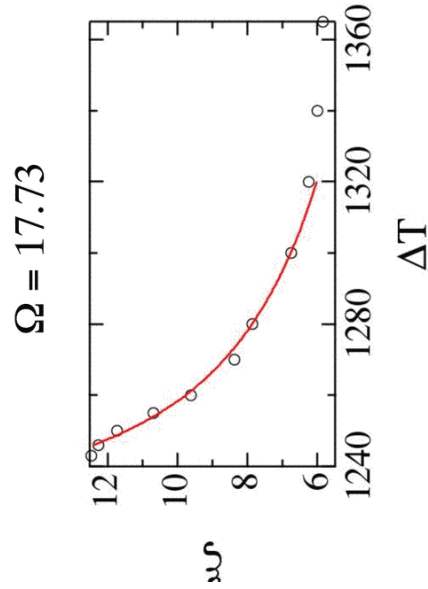
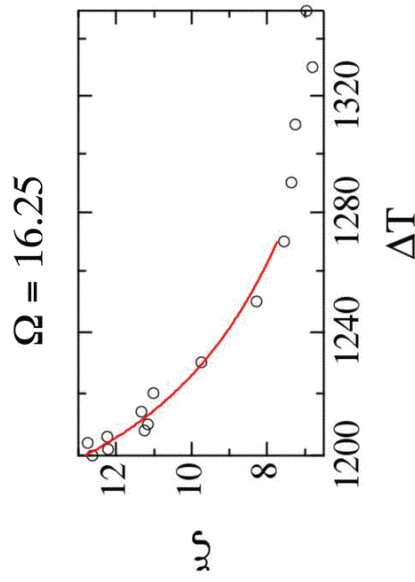
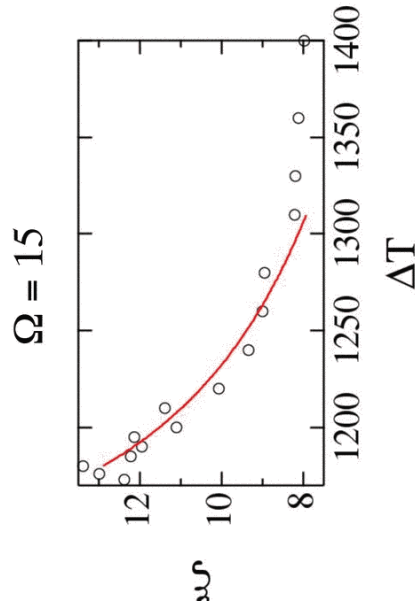
## Theoretical prediction revisited

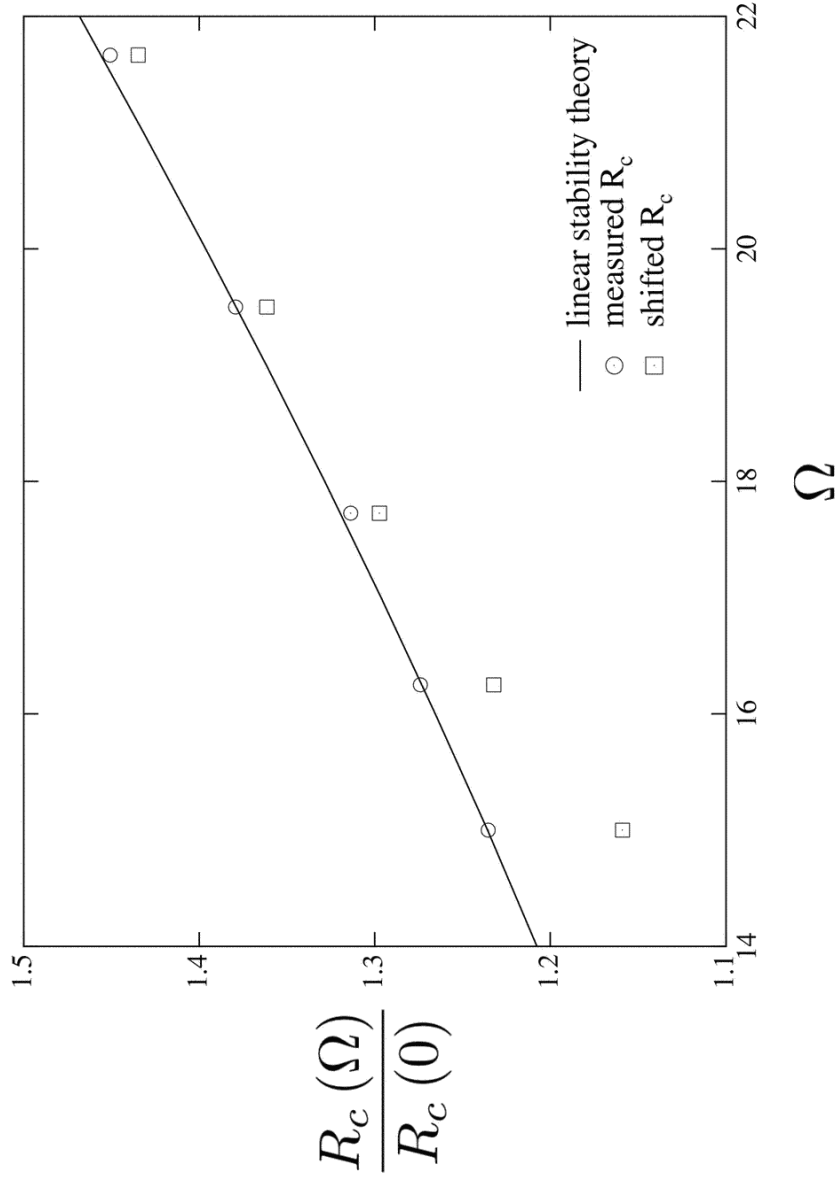
$$\xi \propto \varepsilon^{-1/2}$$

$$\varepsilon \equiv \frac{\Delta T}{\Delta T_c} - 1$$

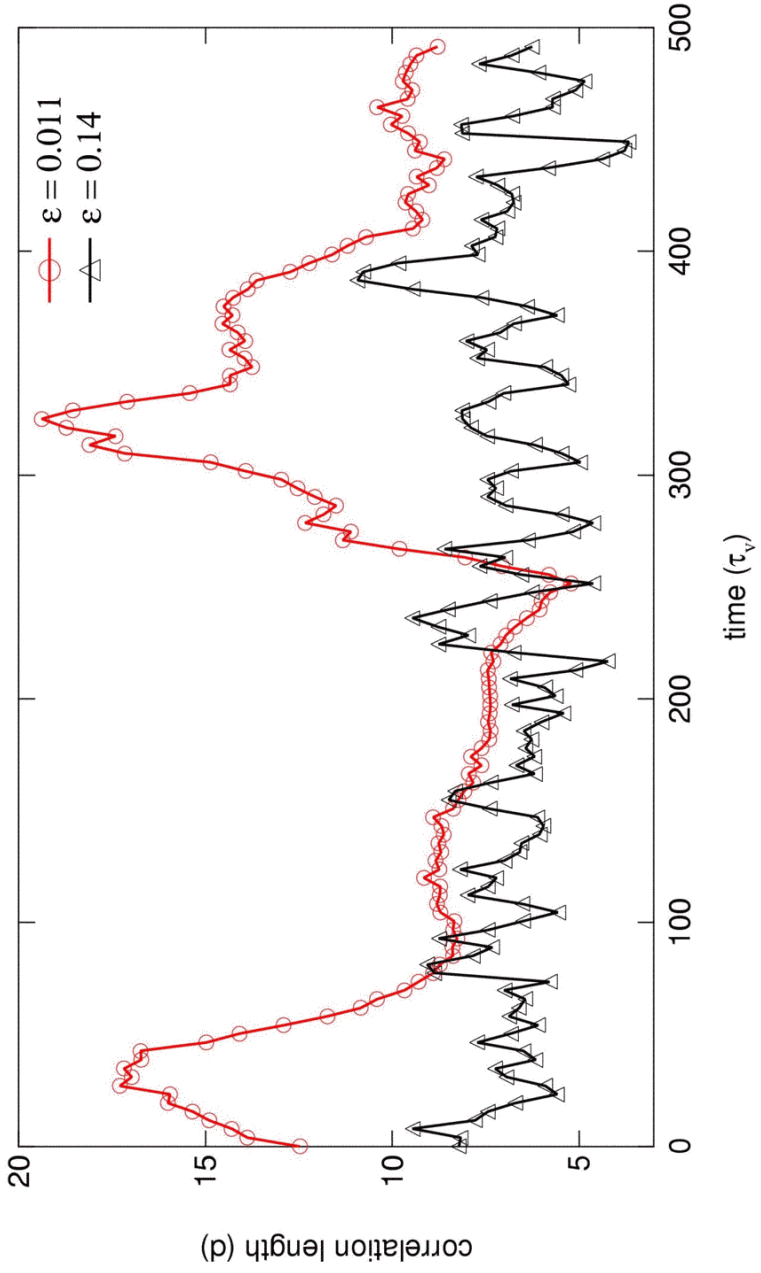
Fitting equation with adjustable onset temperature

$$\xi \propto \left( \frac{\Delta T}{\Delta T_c} - 1 \right)^{-1/2}$$

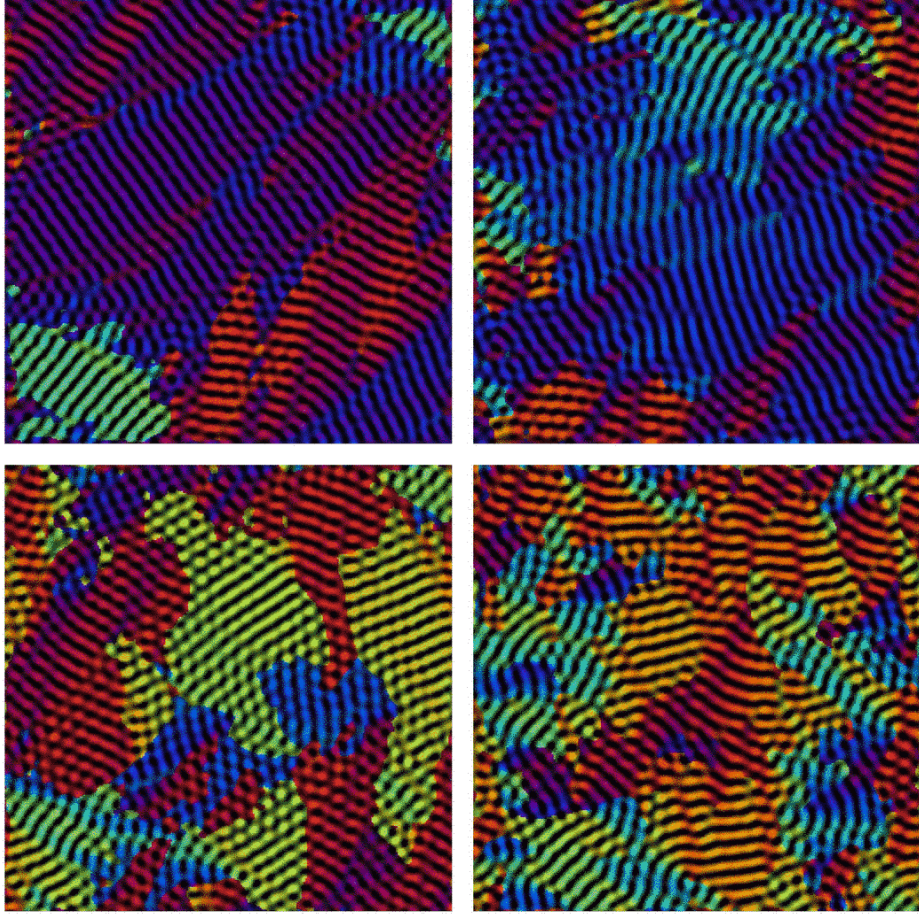




**Time dependence of exponential fit parameter**



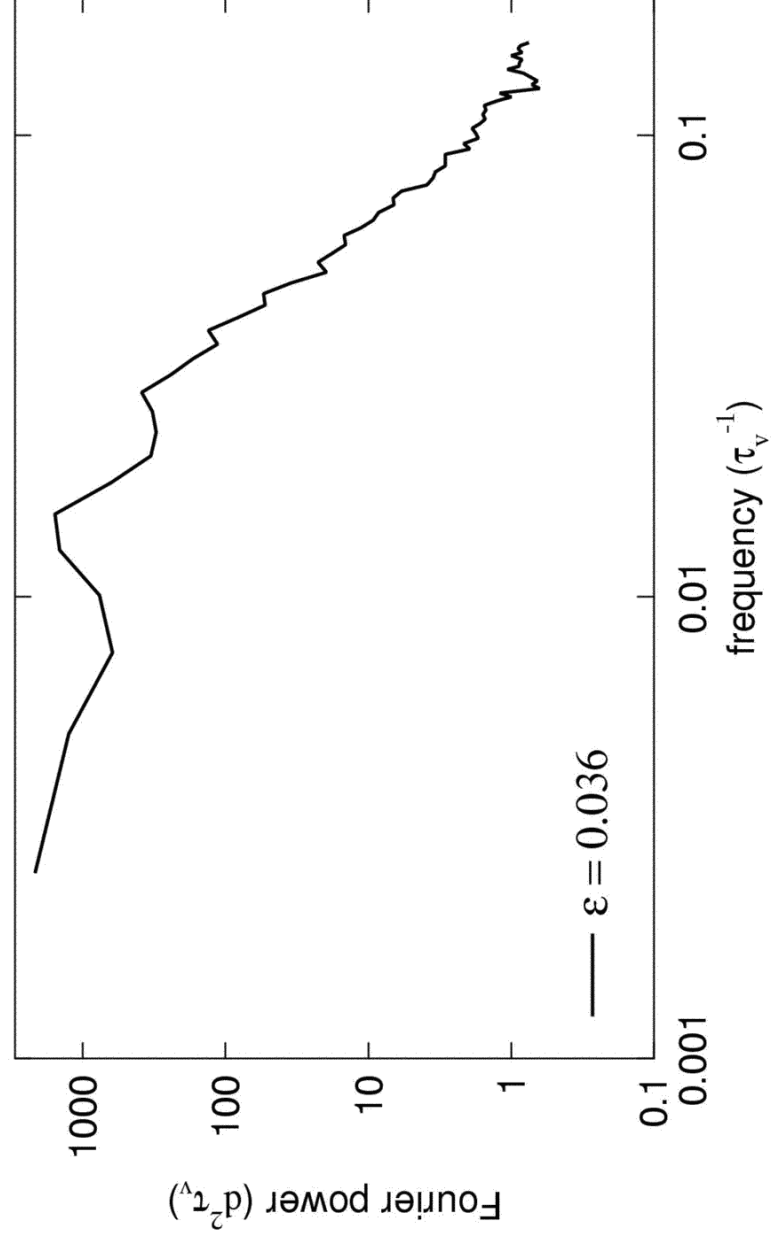
(d)



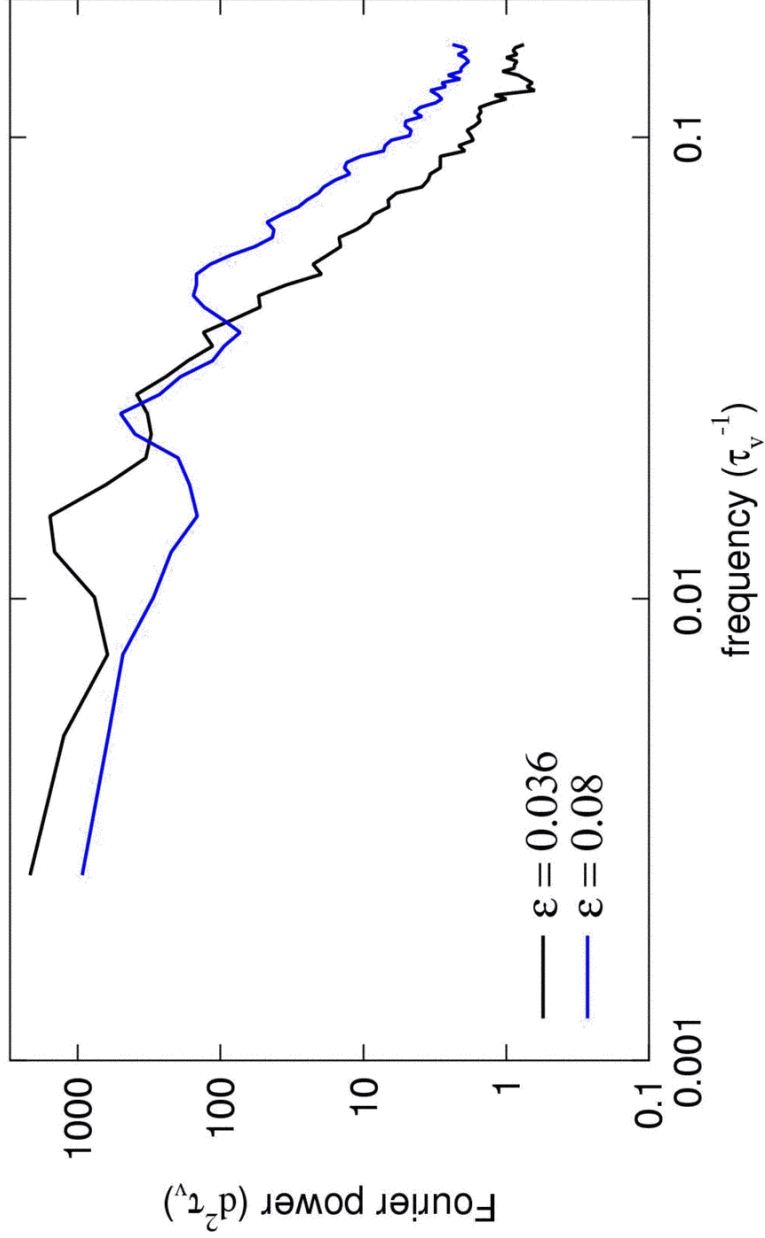
$\epsilon = 0.025$

$\epsilon = 0.125$

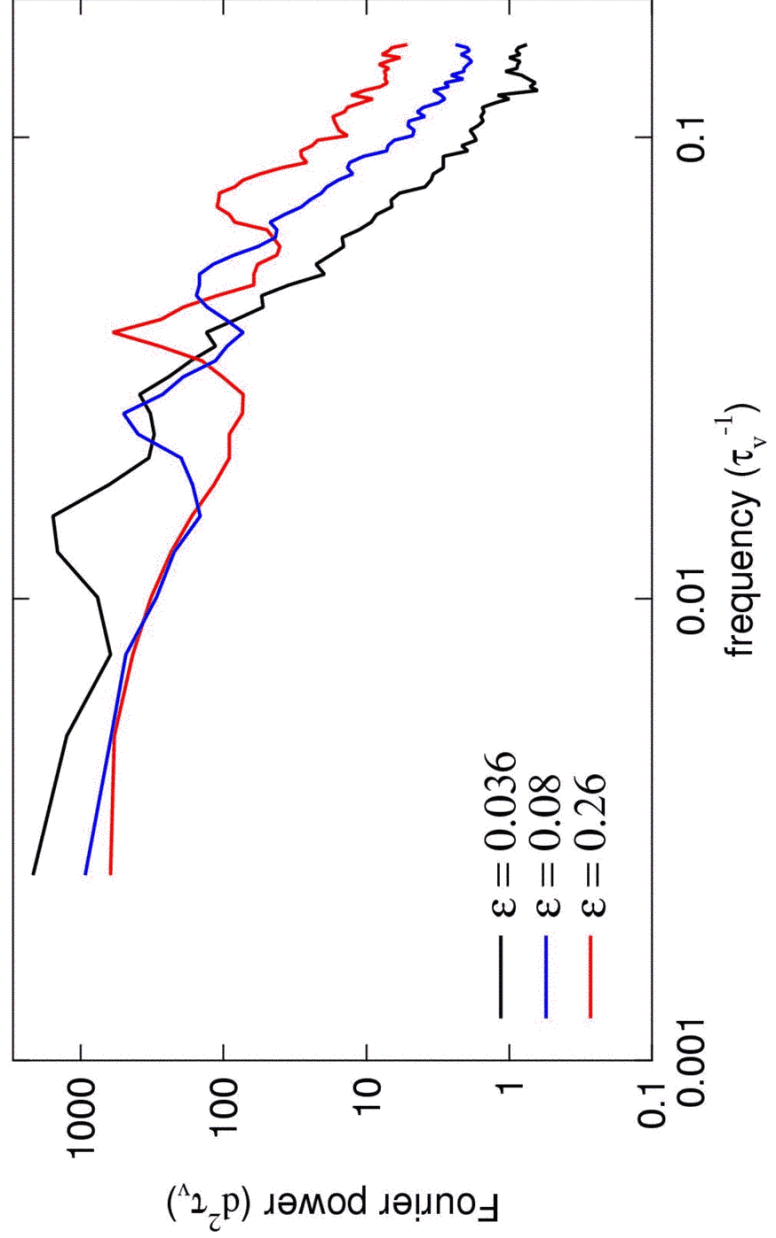
**Power spectrum of correlation length**



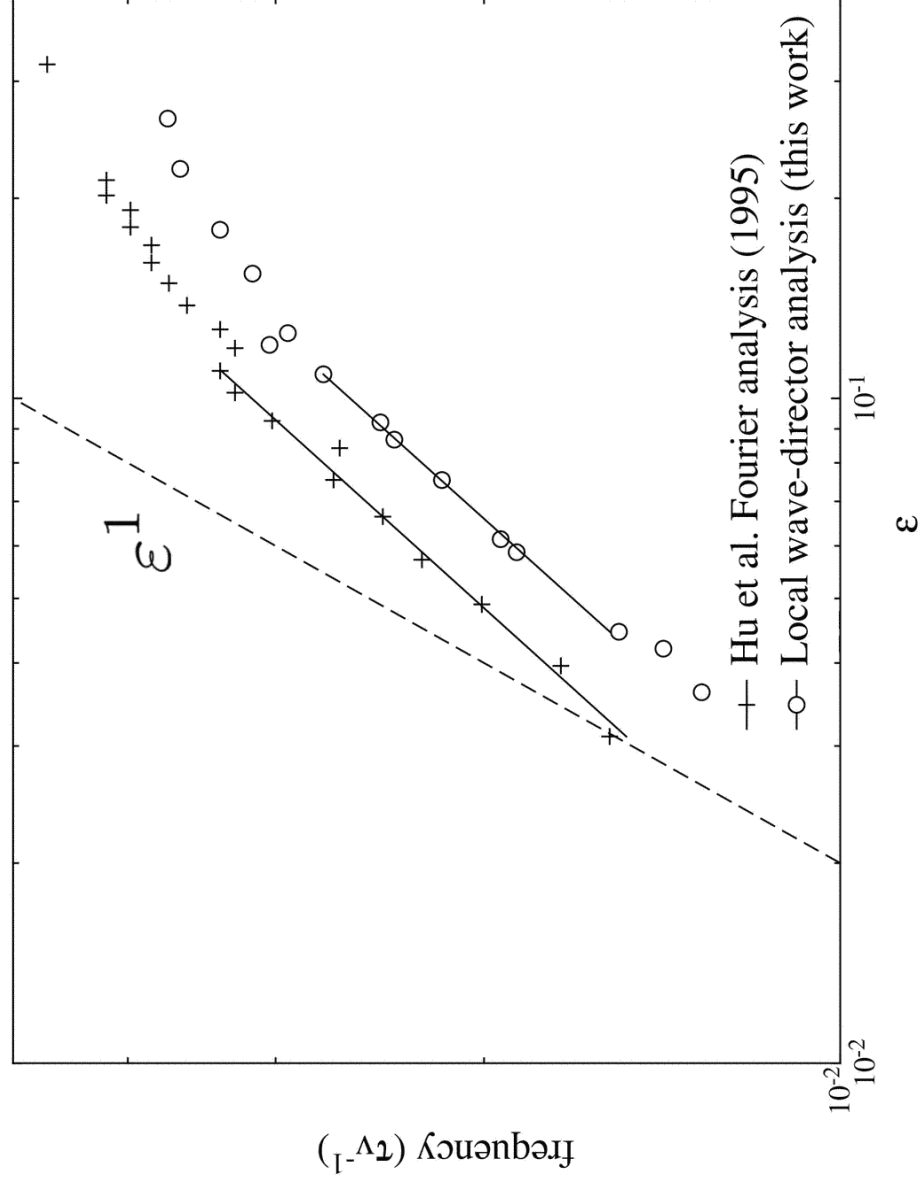
**Power spectrum of correlation length**



**Power spectrum of correlation length**







## Summary

- Experimentally measured time and length scales do not obey the predicted power law dependence. Thus, weakly nonlinear theory in its usual form seems inapplicable.
- The correlation length, measured using local wave-director analysis, exhibits unexplained periodic time dependence. This means that the domain size oscillates. The frequency is roughly the same as the domain switching frequency.