

Pattern Formation and Symmetry in the Visual Cortex

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Klüver: We wish to stress merely one point, namely, that under diverse conditions **the visual system** responds in terms of a **limited number of form constants**.

Outline

1. Visual Hallucinations
2. Structure of Visual Cortex
 - (a) Hubel and Wiesel hypercolumns
 - (b) local and lateral connections
 - (c) isotropy versus anisotropy
3. Pattern Formation in Planar Systems
 - (a) Symmetry
 - (b) Four models
4. Interpretation of Patterns in Retinal Coordinates
 - (a) threshold patterns
 - (b) thin line contour patterns
 - (c) time-periodic patterns

Visual Hallucinations

- Drug **uniformly** forces activation of cortical cells
- Leads to **spontaneous** pattern formation on cortex
- Map from retina to primary visual cortex;
translates pattern on cortex to visual image
- Patterns fall into four *form constants* (Klüver, 1928):
 - tunnels and funnels
 - spirals
 - lattices includes honeycombs and triangles
 - cobwebs

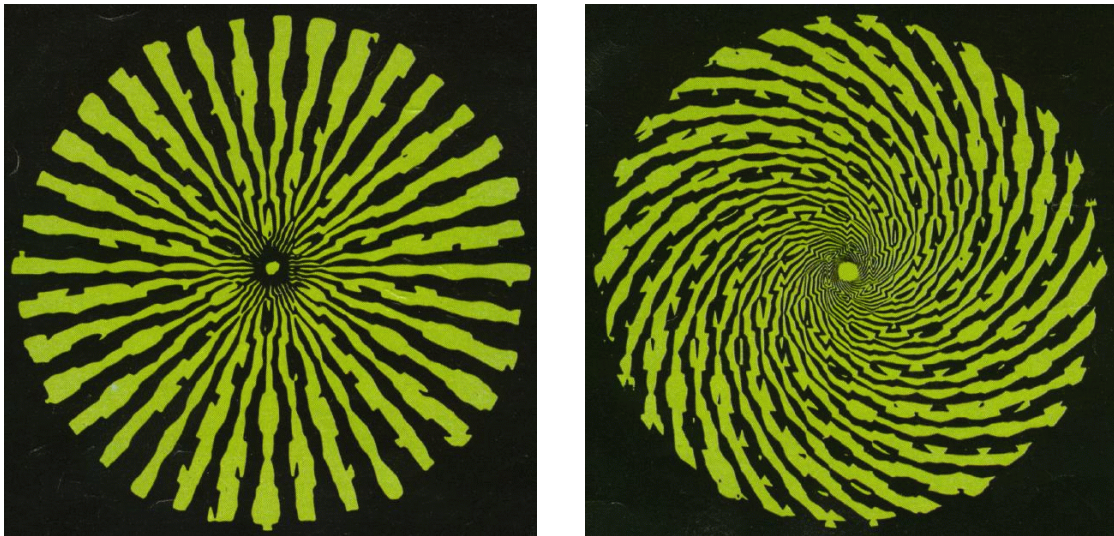


Figure 1: **Funnels** and **spirals** (G. Oster, *Scientific American*, 1970)

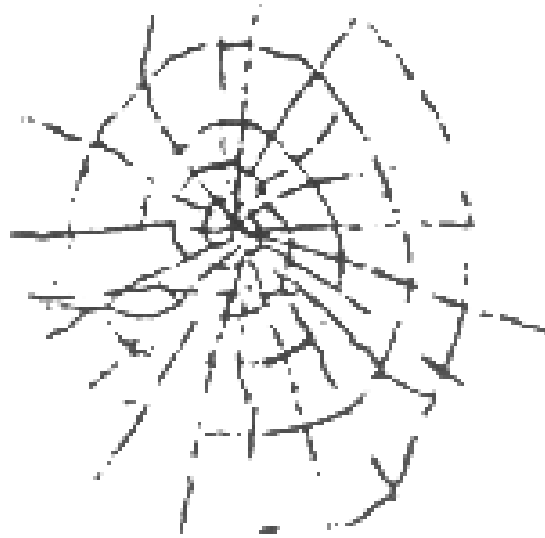


Figure 2: Cobweb (Patterson, 1992).

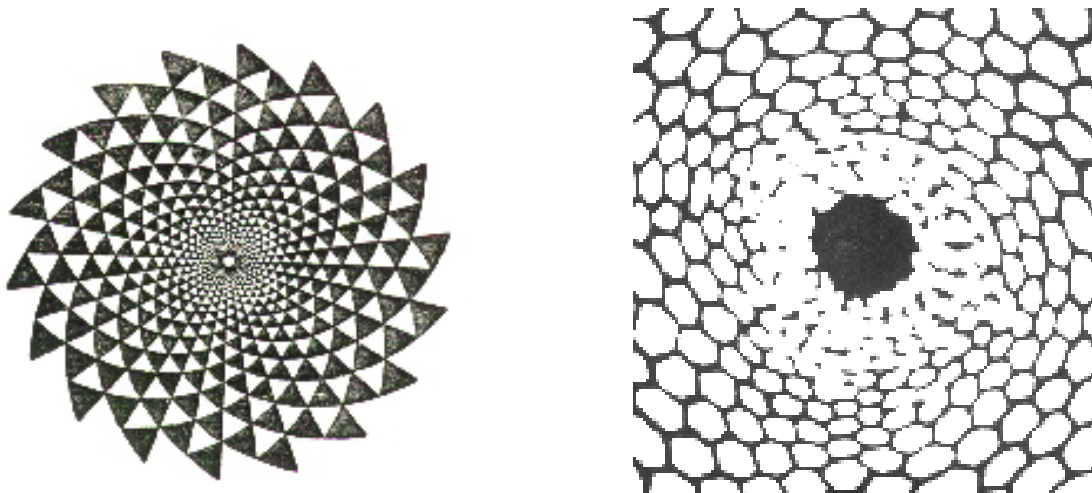


Figure 3: (Left) Phosphene produced by deep binocular pressure on eyeballs; (Right) Honeycomb generated by marihuana



Figure 4: **Lattice-tunnel** generated by *marihuana* (Hall)

Orientation Sensitivity of Cells in V1

- Most V1 cells sensitive to *orientation* of contrast edge

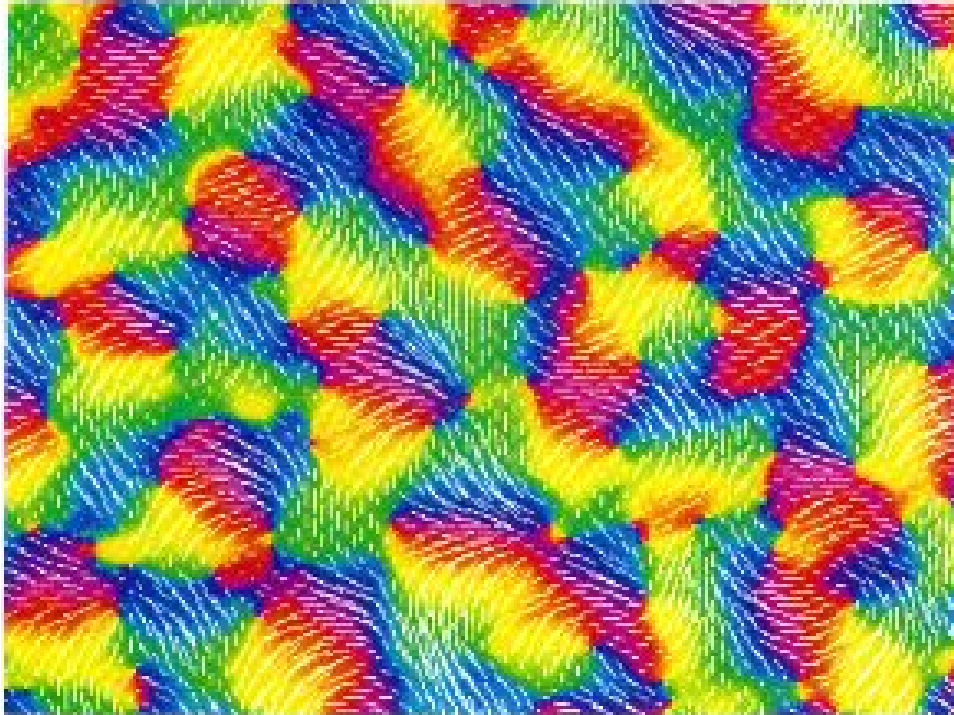


Figure 5: Distribution of orientation preferences in Macaque V1 (Blasdel)

- Hubel and Wiesel, 1974

Each millimeter there is a *hypercolumn* consisting of orientation sensitive cells in every direction preference

Structure of Primary Visual Cortex (V1)

- Optical imaging exhibits pattern of connection

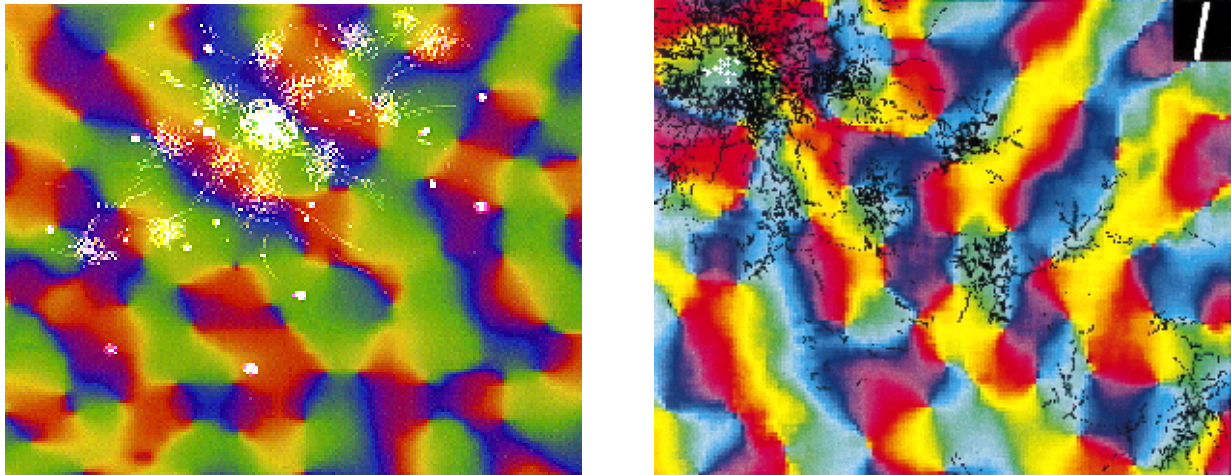


Figure 6: V1 lateral connections: Macaque (left, Blasdel) and Tree Shrew (right, Fitzpatrick)

- Two kinds of coupling: **local** and **lateral**
 - (a) **local**: cells $< 1mm$ connect with most neighbors
 - (b) **lateral**: cells make contact each mm along axons; connections in direction of cell's preference
 - Lateral coupling **small** compared to local coupling
- Anisotropy in lateral coupling small**

Optical imaging suggests **spatial anisotropy**.

Tree shrew: anisotropy pronounced

Macaque: most anisotropy due to stretching in direction orthogonal to ocular dominance columns

Action of Euclidean Group

- Euclidean group: **rotations, reflections, translations**
- Many differential equations are **Euclidean invariant**
Similarity of **pattern formation** due to symmetry
- Abstract **physical space** of V1 is $\mathbf{R}^2 \times \mathbf{S}^1$ — not \mathbf{R}^2
Hypercolumn becomes **circle** measuring orientation

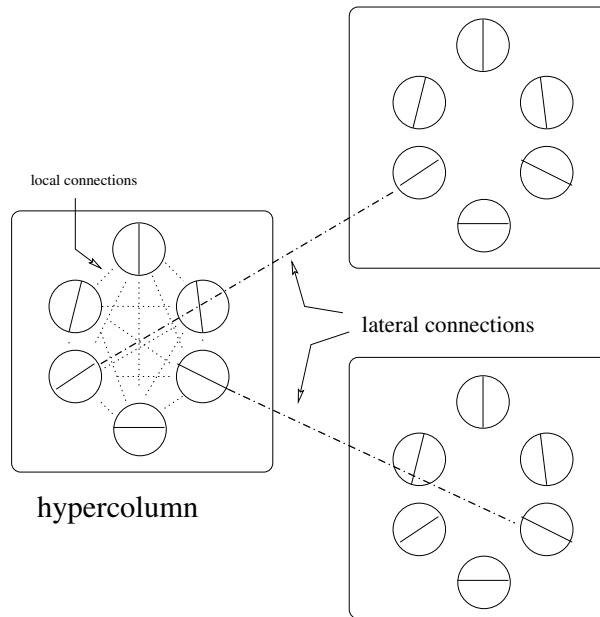


Figure 7: Abstraction of **local** and **anisotropic lateral** connections in V1

- Euclidean groups acts on $\mathbf{R}^2 \times \mathbf{S}^1$ by

$$R_\theta(x, \varphi) = (R_\theta x, \varphi + \theta) \quad \kappa(x, \varphi) = (\kappa x, -\varphi)$$

$$T_y(x, \varphi) = (T_y x, \varphi)$$

Isotropic Lateral Connections

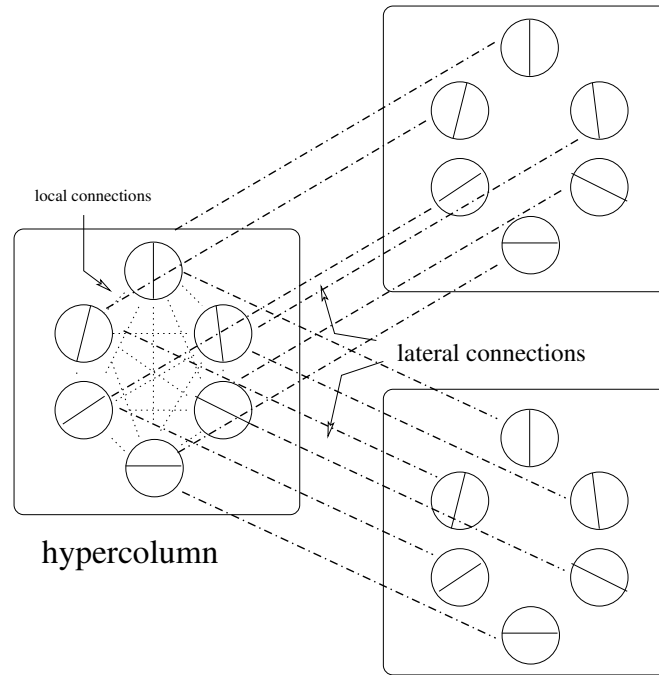


Figure 8: Abstraction of **local** and **isotropic lateral** connections in V1

- **Isotropic** lateral connections introduce **new** $\mathbf{O}(2)$ symmetry

$$\hat{\phi}(x, \varphi) = (x, \varphi + \hat{\phi})$$

- Weak anisotropy is **forced symmetry breaking** of

$$\mathbf{E}(2) \dot{+} \mathbf{O}(2) \rightarrow \mathbf{E}(2)$$

Four Models

1. $\mathbf{E}(2)$ acting on \mathbf{R}^2 (Ermentrout-Cowan)
neurons located at each point x
2. Shift-twist action of $\mathbf{E}(2)$ on $\mathbf{R}^2 \times \mathbf{S}^1$ (Bressloff-Cowan)
hypercolumns located at x ; neurons tuned to φ
anisotropic lateral connections
3. $\mathbf{E}(2) \dot{+} \mathbf{O}(2)$ acting on $\mathbf{R}^2 \times \mathbf{S}^1$ (Wolf)
isotropic lateral coupling
4. Symmetry breaking: $\mathbf{E}(2) \dot{+} \mathbf{O}(2) \rightarrow \mathbf{E}(2)$
weakly anisotropic lateral coupling

Pattern Formation Outline

1. Double-Periodicity and Planar Lattices

- **Translations**: **plane waves**
- **Reflections**: **even** and **odd** representations
- **Rotations**: **infinite-dimensional** eigenspaces
- **Lattices**: back to **finite** dimensions

2. Bifurcation Theory with Symmetry

- Equivariant Branching Lemma
- **Scalar** and **pseudoscalar** bifurcations

3. Planforms

- Adaptation to **Visual Cortex**
Line Fields, contours, and thresholding
- **Winner-take-all** strategy
- **Cortex to Retina** transformation

Observations Using Symmetry

Bosch Vivancos, Chossat, Melbourne

- Assume system of differential equations on $\mathbf{R}^2 \times \mathbf{S}^1$ with **Euclidean equivariant** linearization \mathbf{L}

$$\mathbf{L}\gamma = \gamma\mathbf{L} \quad \forall \gamma \in \mathbf{E}(2)$$

- **Planforms** are approximated by eigenfunctions of \mathbf{L}
Symmetry dictates eigenfunctions
- **TRANSLATIONS** on $\mathbf{R}^2 \times \mathbf{S}^1$ imply

$$W_{\mathbf{k}} = \{u(\varphi)e^{i\mathbf{k}\cdot\mathbf{x}} + \text{c.c.} : u : \mathbf{S}^1 \rightarrow \mathbf{C}\}$$

is \mathbf{L} -invariant subspace for every **dual wave vector**

$$\mathbf{k} \in \mathbf{R}^2$$

- Eigenfunctions have **plane wave factors**

$$u(\varphi)e^{i\mathbf{k}\cdot\mathbf{x}} + \text{c.c.}$$

Action of Reflections

- Choose **REFLECTION** ρ so that $\rho\mathbf{k} = \mathbf{k}$

$$\rho(u(\varphi)e^{i\mathbf{k}\cdot\mathbf{x}}) = \rho(u(\varphi))e^{i\mathbf{k}\cdot\mathbf{x}}$$

So $\rho : W_{\mathbf{k}} \rightarrow W_{\mathbf{k}}$

- $\rho^2 = 1$ implies $W_{\mathbf{k}} = W_{\mathbf{k}}^+ \oplus W_{\mathbf{k}}^-$
where ρ acts as $+1$ on $W_{\mathbf{k}}^+$ and -1 on $W_{\mathbf{k}}^-$

- Eigenfunctions are **even** or **odd**. When $\mathbf{k} = (1, 0)$

$$\begin{aligned} u(-\varphi) &= u(\varphi) & u \in W_{\mathbf{k}}^+ \\ u(-\varphi) &= -u(\varphi) & u \in W_{\mathbf{k}}^- \end{aligned}$$

- Both kinds of eigenfunctions occur in models
- Study **nonoriented** directions: $u(\mathbf{x}, \varphi + \pi) = u(\mathbf{x}, \varphi)$

Action of Rotations

$$R_\theta (u(\varphi)e^{i\mathbf{k}\cdot\mathbf{x}}) = R_\theta(u(\varphi))e^{iR_\theta(\mathbf{k})\cdot\mathbf{x}}$$

Therefore

$$R_\theta(W_{\mathbf{k}}) = W_{R_\theta(\mathbf{k})}$$

- Rotation symmetry implies $\ker L$ is ∞ -dimensional

Planar Lattices

- **Double-periodicity**: Look for solns on lattice
- The space of *doubly periodic* functions w.r.t \mathcal{L} is

$$\mathcal{F}_{\mathcal{L}} = \{f \in \mathcal{F} : f(\mathbf{x} + \boldsymbol{\ell}) = f(\mathbf{x}) \quad \forall \boldsymbol{\ell} \in \mathcal{L}\}$$

- Finite number of rotations: $\ker L$ is finite-dimensional
- Choose lattice size so **shortest** dual vectors are **critical**

Equivariant Bifurcation Theory

- Symmetry group Γ : $f(\gamma x) = \gamma f(x)$
- $\text{Fix}(\Sigma) = \{x \in \mathbf{R}^n : \sigma x = x \quad \forall \sigma \in \Sigma\}$
- $\text{Fix}(\Sigma)$ is **flow invariant**
Proof: $f(x) = f(\sigma x) = \sigma f(x)$

The Equivariant Branching Lemma

- Isotropy subgroup $\Sigma \subset \Gamma$ is *axial* if

$$\dim \text{Fix}(\Sigma) = 1$$

on critical eigenspace

- **Generically, there exists a branch of solutions with Σ symmetry for every axial subgroup Σ**

Planforms For Ermentrout-Cowan

Square lattice: Two axial subgroups of $\mathbf{T}^2 + \mathbf{D}_4$
 $\mathbf{O}(2) \oplus \mathbf{Z}_2$ stripes and \mathbf{D}_4 squares

Hexagonal lattice: Two axial subgroups of $\mathbf{T}^2 + \mathbf{D}_6$
 $\mathbf{O}(2) \oplus \mathbf{Z}_2$ stripes and \mathbf{D}_6 hexagons

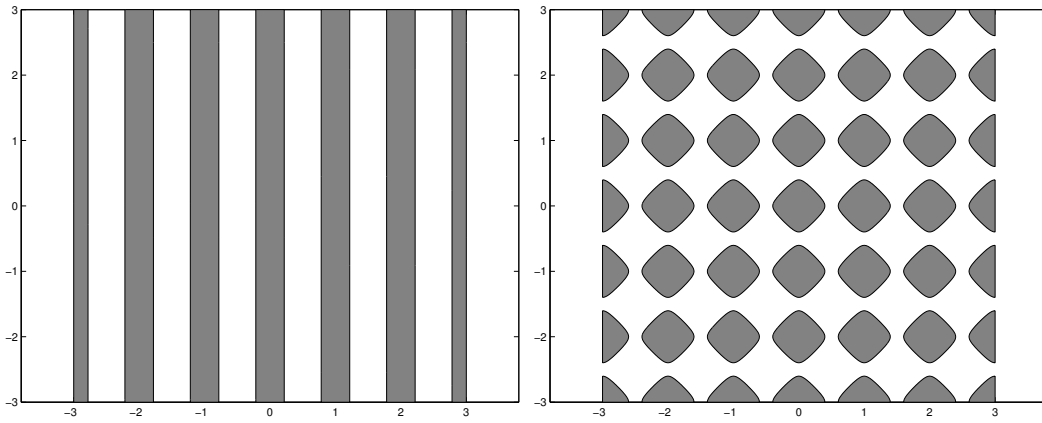


Figure 9: **Thresholding** of axial eigenfunctions: (left) *stripes*; (right) *squares*

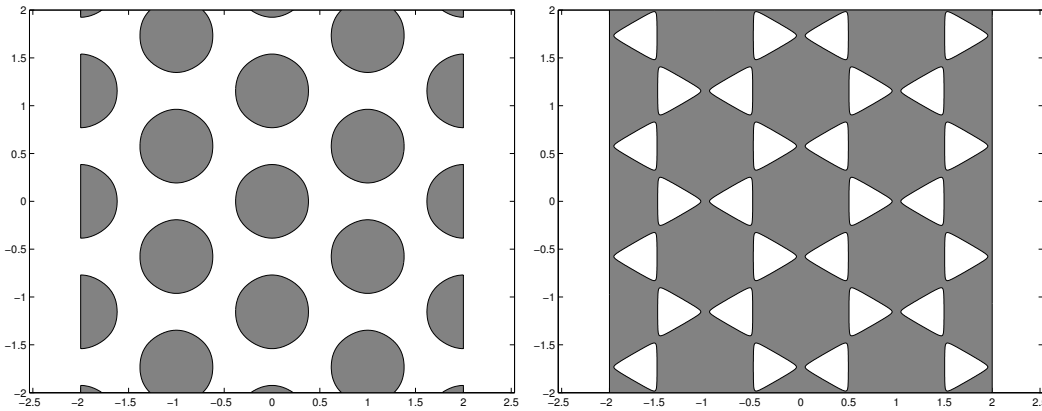


Figure 10: **Thresholding** of axial eigenfunction *hexagons*

Axial Subgroups in Orientation Tuned Models

Name	Axial	Planform Eigenfunction
squares	\mathbf{D}_4	$u(\varphi) \cos x + u\left(\varphi - \frac{\pi}{2}\right) \cos y$
stripes	$\mathbf{O}(2) \oplus \mathbf{D}_1$	$u(\varphi) \cos x$
hexagons	\mathbf{D}_6	$\sum_{j=0}^2 u(\varphi - j\pi/3) \cos(\mathbf{k}_j \cdot \mathbf{x})$
stripes	$\mathbf{O}(2) \oplus \mathbf{D}_1$	$u(\varphi) \cos(\mathbf{k}_1 \cdot \mathbf{x})$

Table 1: Axial planforms when $u(\varphi) = u(-\varphi)$ is **even**.

Name	Axial	Planform Eigenfunction
square	\mathbf{D}_4^*	$u(\varphi) \cos x - u\left(\varphi - \frac{\pi}{2}\right) \cos y$
stripes	$\mathbf{O}(2) \oplus \mathbf{D}_1^*$	$u(\varphi) \cos x$
hexagons	\mathbf{Z}_6	$\sum_{j=0}^2 u(\varphi - j\pi/3) \cos(\mathbf{k}_j \cdot \mathbf{x})$
triangles	\mathbf{D}_3	$\sum_{j=0}^2 u(\varphi - j\pi/3) \sin(\mathbf{k}_j \cdot \mathbf{x})$
rectangles	\mathbf{D}_2	$u\left(\varphi - \frac{\pi}{3}\right) \cos(\mathbf{k}_2 \cdot \mathbf{x}) - u\left(\varphi + \frac{\pi}{3}\right) \cos(\mathbf{k}_3 \cdot \mathbf{x})$
stripes	$\mathbf{O}(2) \oplus \mathbf{D}_1^*$	$u(\varphi) \cos(\mathbf{k}_1 \cdot \mathbf{x})$

Table 2: Axial planforms when $u(\varphi) = -u(-\varphi)$ is **odd**. * = glide reflection

How to Find Amplitude Function $u(\varphi)$

- **Isotropic connections** imply EXTRA $\mathbf{O}(2)$ symmetry
- $\mathbf{O}(2)$ decomposes $W_{\mathbf{k}}$ into sum of irreducible subspaces

$$W_{\mathbf{k},p} = \{ze^{p\varphi i} e^{i\mathbf{k}\cdot x} + \text{c.c.} : z \in \mathbf{C}\} \cong \mathbf{R}^2$$

Generically, eigenfunctions of \mathbf{L} lie in $W_{\mathbf{k},p}$ for some p

- $W_{\mathbf{k},p}^+ = \{\cos(p\varphi)e^{i\mathbf{k}\cdot x}\}$ even case
 $W_{\mathbf{k},p}^- = \{\sin(p\varphi)e^{i\mathbf{k}\cdot x}\}$ odd case
- **Wilson-Cowan models** lead to
 $p = 0$ or $p = 1$ bifurcations in even case
 $p = 1$ bifurcations in odd case
- Compute pictures in $p = 1$ cases

$$u(\varphi) \approx \cos(\varphi) \text{ and } u(\varphi) \approx \sin(\varphi)$$

Winner-Take-All Strategy

Creation of Line Fields

- **Given:** Activity of neuron in hypercolumn at \mathbf{x} sensitive to direction φ
- **Assumption:** Most active neuron in hypercolumn suppresses other neurons in hypercolumn
- **Consequence:** For all \mathbf{x} find $\varphi_{\mathbf{x}} \in \mathbf{S}^1$ where activity is maximum
- **Planform:** Draw small line segment at \mathbf{x} oriented at angle $\varphi_{\mathbf{x}}$

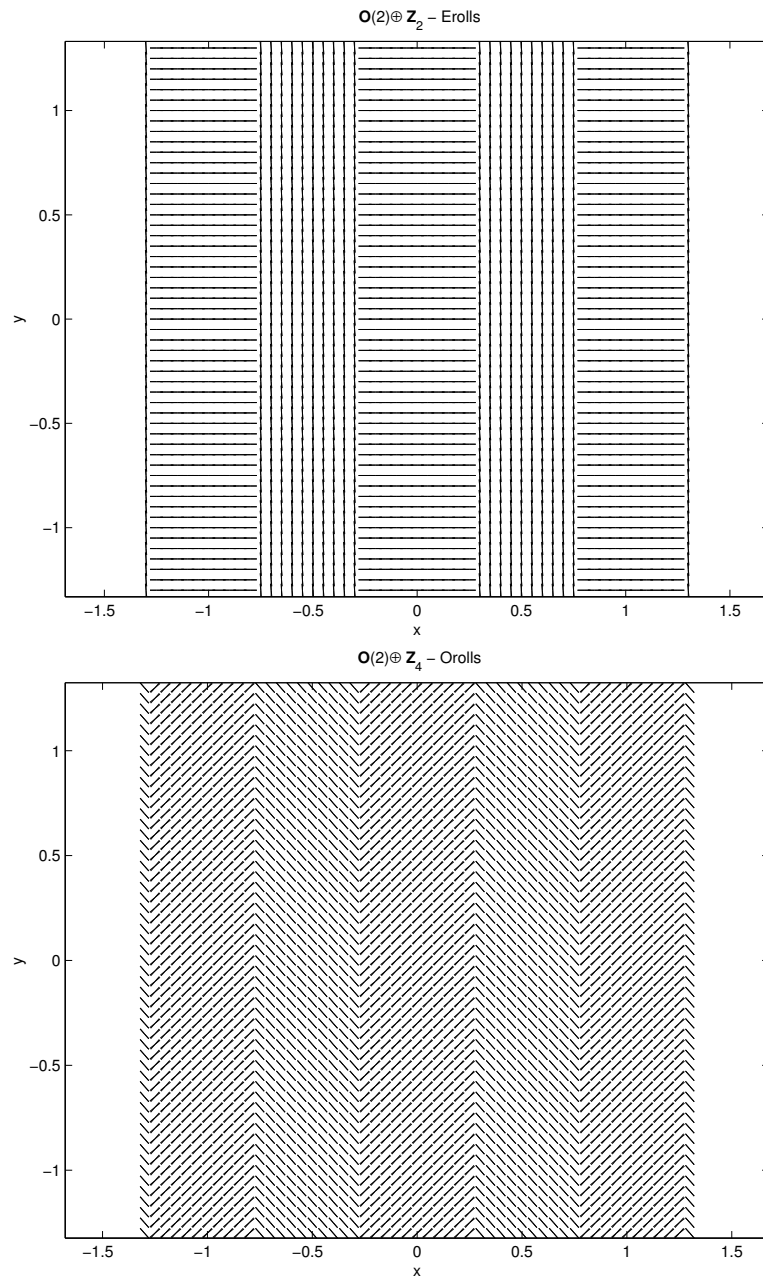


Figure 11: Even Stripes (upper); odd Stripes (lower)

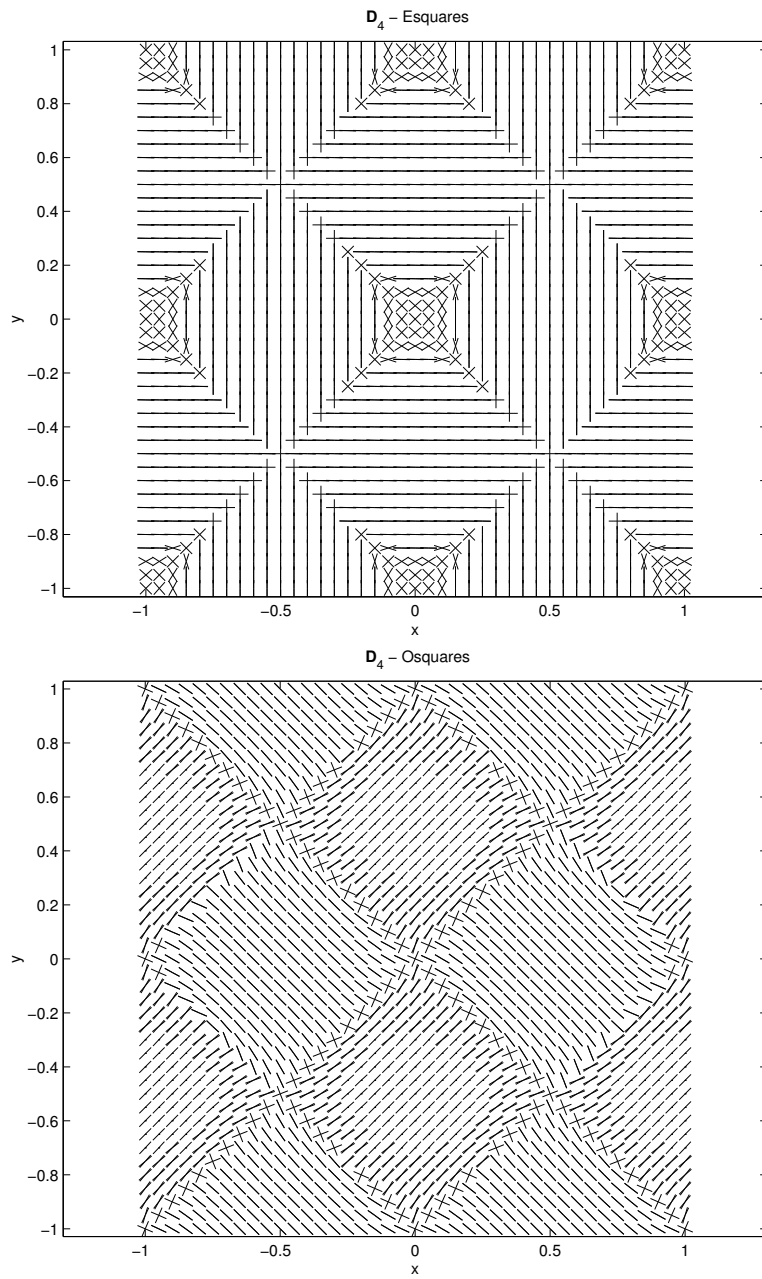


Figure 12: Even Squares (upper); odd Squares (lower)

Cortex to Retina

- Neurons on cortex are **uniformly** distributed
- Neurons in retina fall off by $1/r^2$ from fovea
- Unique conformal map takes uniform density square to $1/r^2$ density disk: **complex exponential**
- Cortex to retinal map is

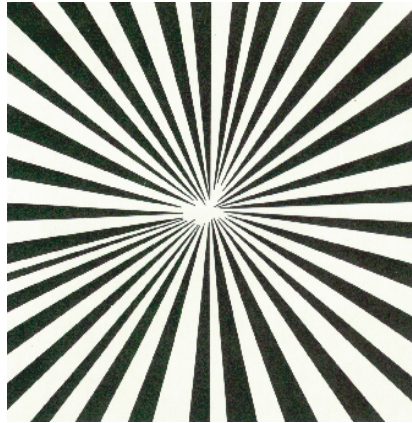
$$\begin{aligned}r &= \omega \exp(\epsilon x) \\ \theta &= \epsilon y\end{aligned}$$

In retinal images we take

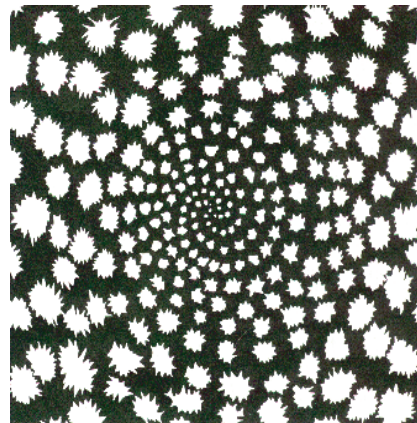
$$\omega = 30/e^{2\pi} \quad \text{and} \quad \epsilon = 2\pi/n_h$$

where $n_h = 36 = \#$ hypercolumn widths in cortex

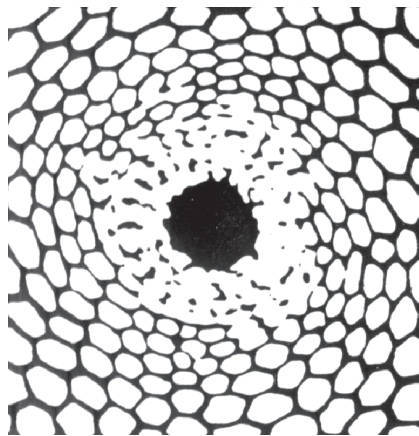
- Straight lines on cortex \mapsto
circles, logarithmic spirals, and rays in retina



(I)



(II)



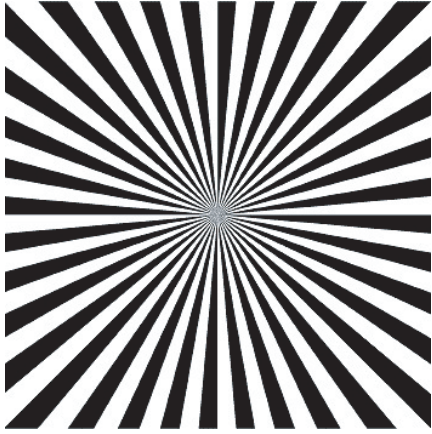
(III)



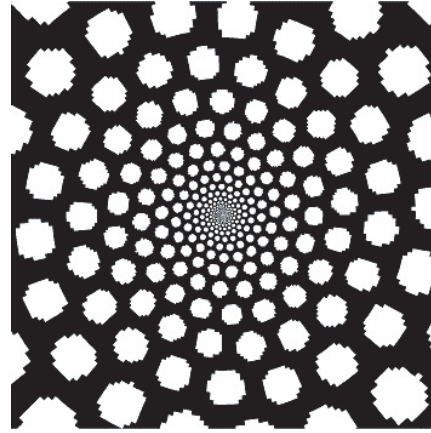
(IV)

Figure 13: Hallucinatory form constants. (I) **funnel** and (II) **spiral** images seen following ingestion of LSD [Siegel & Jarvik, 1975], (III) **honeycomb** generated by marihuana [Clottes & Lewis-Williams (1998)], (IV) **cobweb** petroglyph [Patterson, 1992].

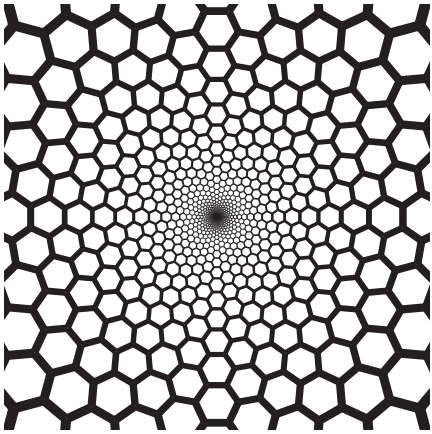
Planforms in the Visual Field



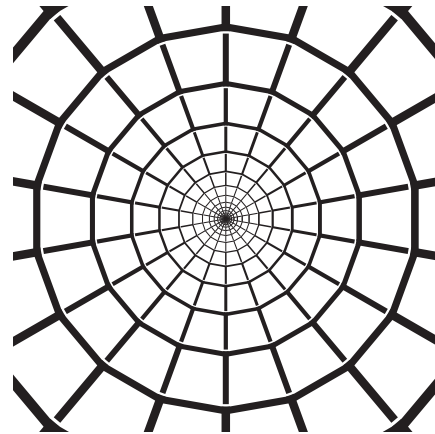
(a)



(b)



(c)



(d)

Visual field planforms

Isotropic Coupling: Extra $O(2)$ symmetry

- $\widehat{\varphi}(\mathbf{x}, \varphi) = (\mathbf{x}, \varphi + \widehat{\varphi})$
- Eigenspaces: sum of even and odd
- Square lattice:
 - four axials
 - one maximal subgroup with 2D fixed-point subspace

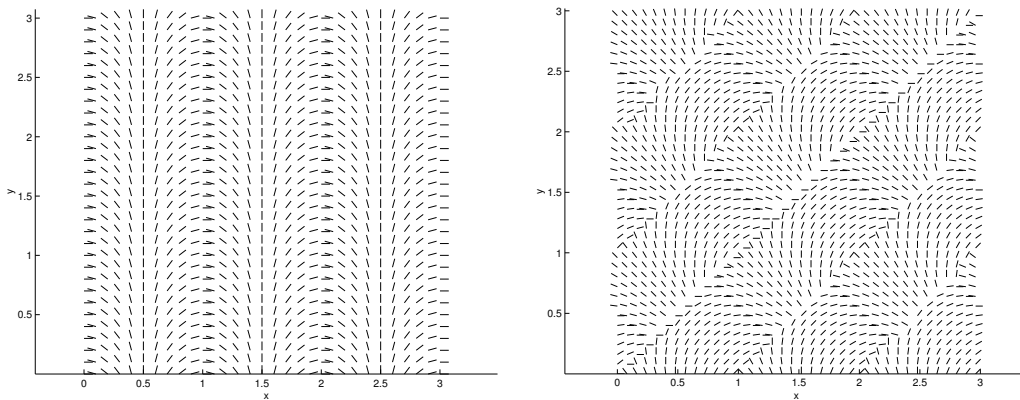


Figure 14: Direction fields of new planforms in isotropic model.

- Hexagonal lattice:
 - Nine axials
 - three maximal subgroups with 2D fixed-pt subspaces

Hallucinations in Isotropic Coupling Model

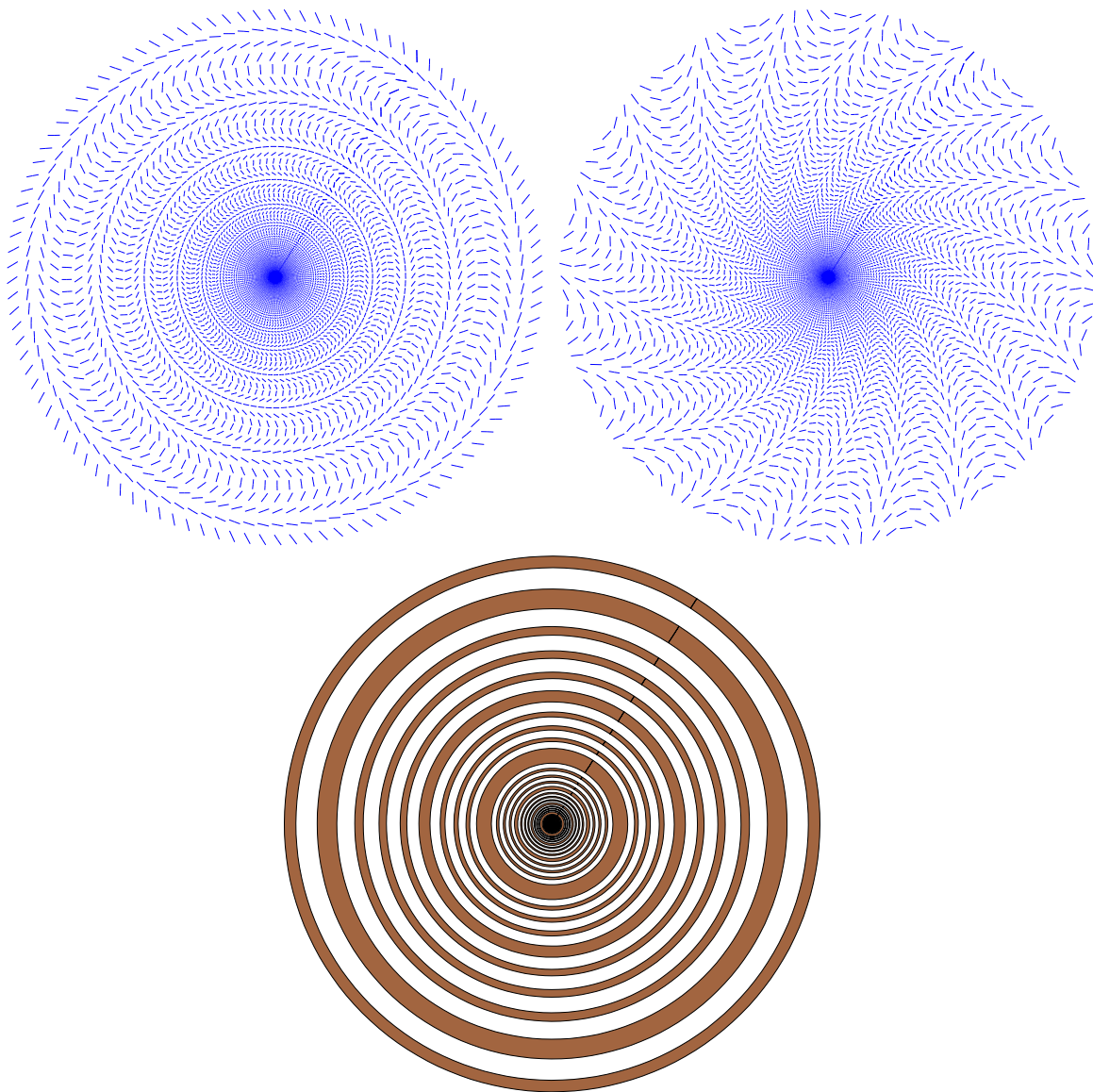


Figure 15: (Top) Conjugate solutions (7); (bottom) threshold.

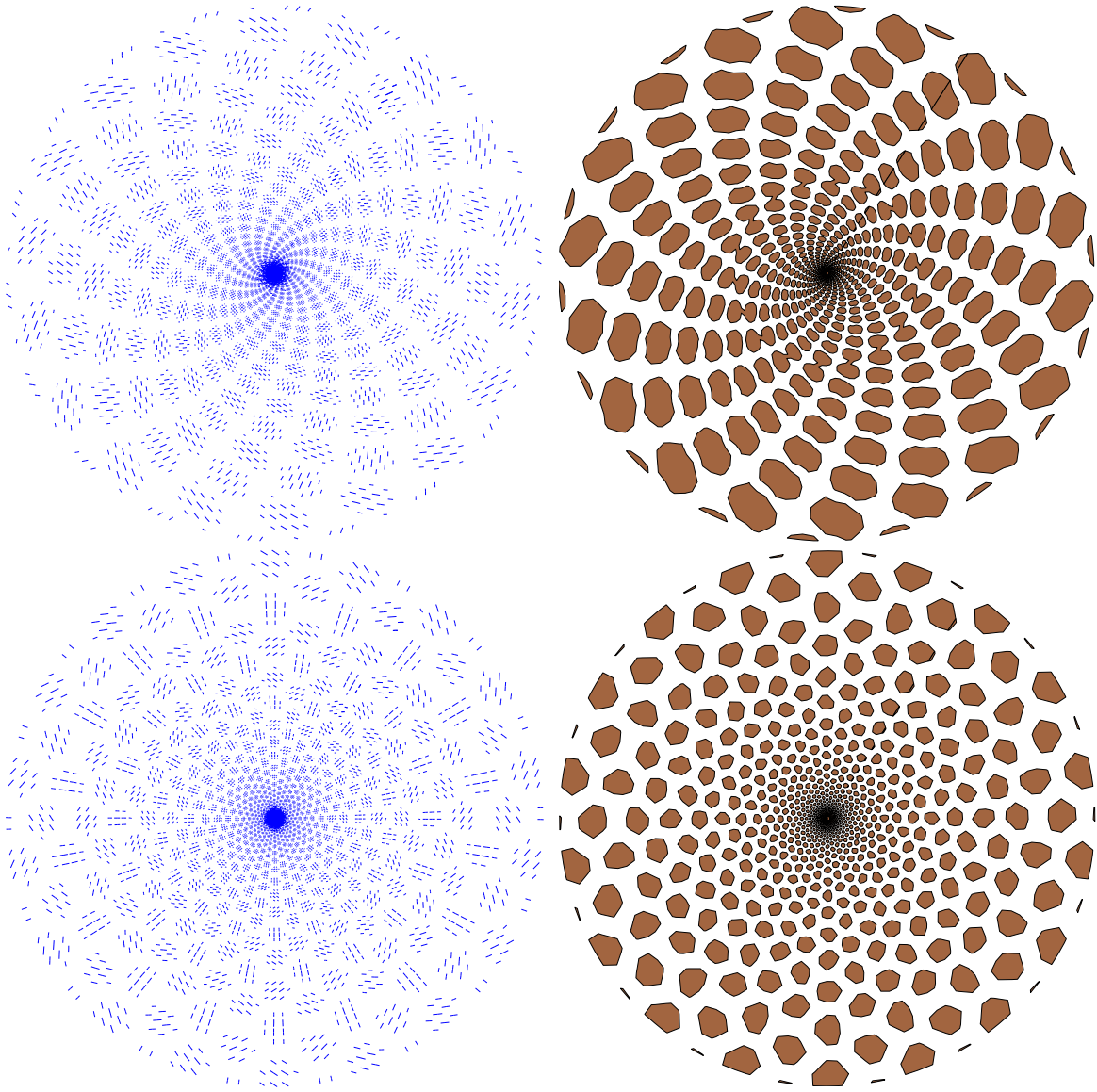


Figure 16: **Phosphene-like** planforms: (top) planform (12); (bottom) planform (9)

Weakly Anisotropic Coupling

- Square lattice: Forced symmetry-breaking to
 1. even and odd stripes
 2. even and odd squares
 3. two new equilibrium planforms
 4. a time-periodic rotating wave

- Hexagonal lattice: Forced symmetry-breaking to
 1. seven types of equilibria
 2. two contracting or expanding periodic states
 3. two rotating waves
 4. state that is an equilibrium or time-periodic

Landau Theory of Phase Transitions for a Liquid Crystal

- **nematic phase**

preferred direction along which molecules align

- Alignment of molecules represented by

3×3 symmetric trace zero matrices $Q(x)$

- Molecule at x aligns along **eigendirection** of $Q(x)$
corresponding to **largest** eigenvalue

- Q is **second moment** of probability distribution
for alignment of rod-like molecule

- Action of $\mathbf{E}(3)$: Let $\gamma \in \mathbf{O}(3)$ and $y \in \mathbf{R}^3$

$$\begin{aligned}(T_y Q)(x) &= Q(x - y) \\ (\gamma \cdot Q)(x) &= \gamma Q(\gamma^{-1}x)\gamma^{-1}\end{aligned}$$

Scalar and Pseudoscalar Eigenspaces

Chillingworth and Golubitsky

- Spatially uniform liquid crystal: Q independent of x
homeotropic — all molecules aligned in one direction
isotropic — equally likely to align in any direction

- Bifurcation from Euclidean invariant Q_0

Translation invariance implies Q_0 is constant

Rotation invariance implies: $Q_0 = \alpha \begin{bmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 2 \end{bmatrix}$

- $\alpha > 0$: **homeotropic** — directors point vertically
- $\alpha < 0$: **isotropic** — directors equally likely to point in any horizontal direction

- Q_0 is invariant under up-down reflection: $\tau = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -1 \end{bmatrix}$

- Symmetry group is $\Gamma = \mathbf{E}(2) \oplus \mathbf{Z}_2(\tau)$

Just like Bénard convection with midplane reflection

Linear Theory

- **Translations**: eigenfunctions have form $e^{2\pi i \mathbf{k} \cdot x} Q + \text{c.c.}$

$$W_{\mathbf{k}} = \{e^{2\pi i \mathbf{k} \cdot x} Q + \text{c.c.} : Q \text{ is complex-valued}\}$$

where $\dim W_{\mathbf{k}} = 10$

- **Four** possible bifurcations

scalar/pseudoscalar
 τ acts trivially/nontrivially

- Since L commutes with τ , we can subdivide

$$W_{\mathbf{k}} = W_{\mathbf{k}}^{++} \oplus W_{\mathbf{k}}^{+-} \oplus W_{\mathbf{k}}^{-+} \oplus W_{\mathbf{k}}^{--}$$

where each $W_{\mathbf{k}}^{\pm\pm}$ is **L-invariant**

-

$$Q^{++} = \begin{bmatrix} a & 0 & 0 \\ 0 & b & 0 \\ 0 & 0 & -a - b \end{bmatrix} \quad Q^{+-} = \begin{bmatrix} 0 & 0 & i \\ 0 & 0 & 0 \\ i & 0 & 0 \end{bmatrix}$$

and

$$Q^{-+} = \begin{bmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \quad Q^{--} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & i \\ 0 & i & 0 \end{bmatrix}$$

Pattern Formation from the Isotropic State

Rolls solutions on bifurcation from αQ_0 when $\alpha < 0$

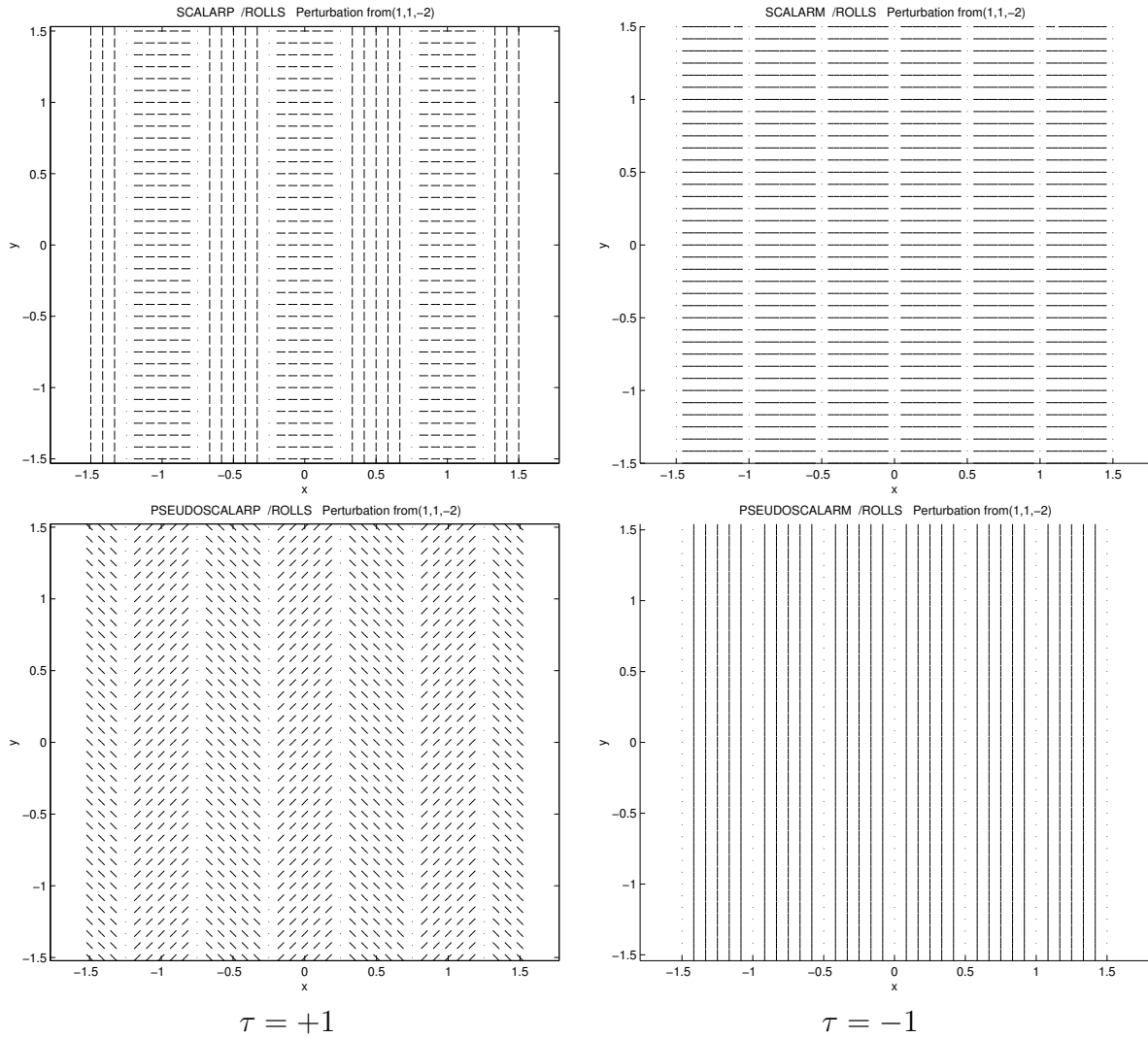


Figure 17: Rolls perturbation from $-Q_0$: (upper) scalar; (lower) pseudoscalar $\tau = +1$.

Pattern Formation from the Homeotropic State

Rolls solutions on bifurcation from αQ_0 when $\alpha > 0$

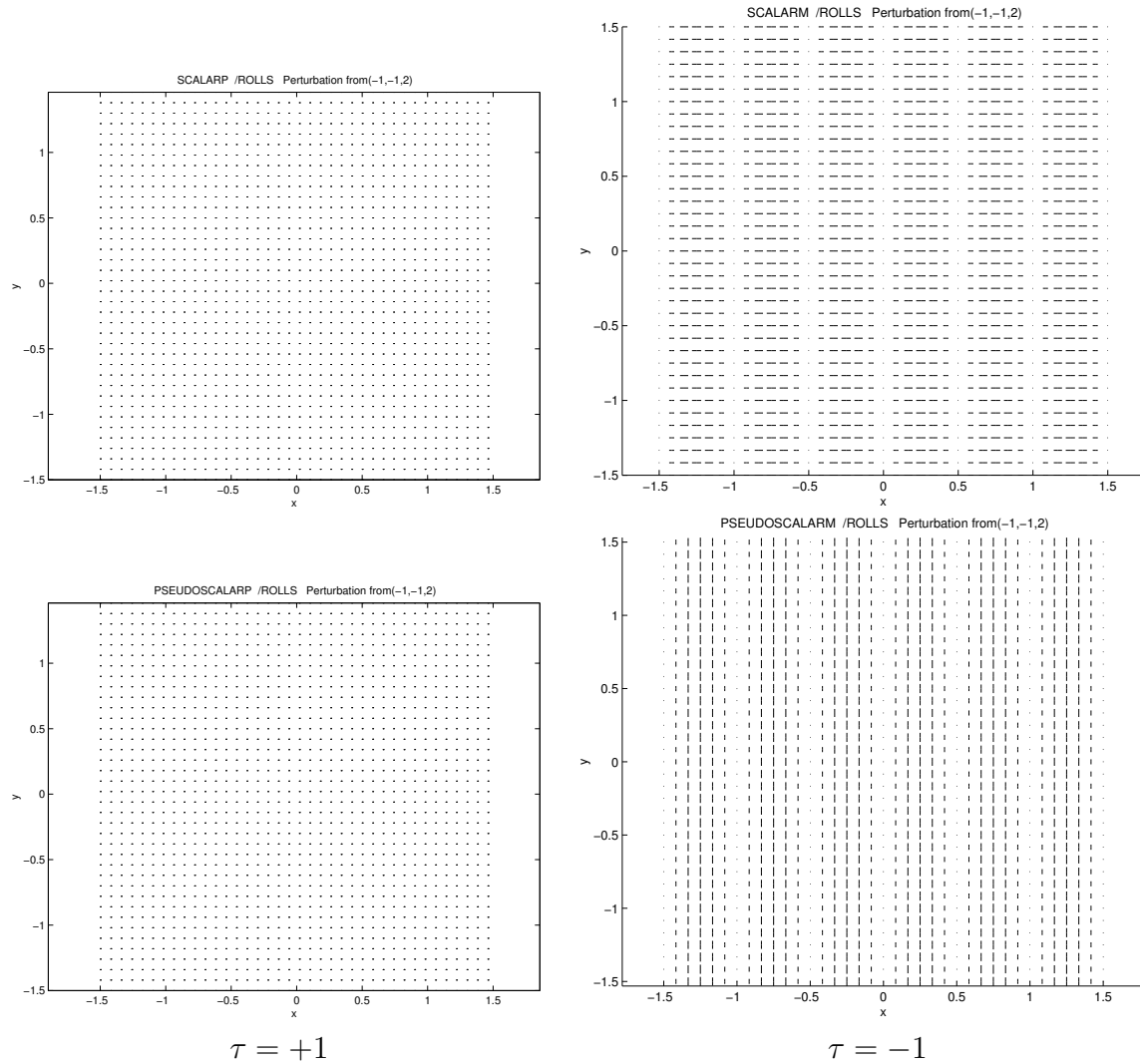


Figure 18: Rolls perturbation from Q_0 : (upper) scalar; (lower) pseudoscalar $\tau = +1$.