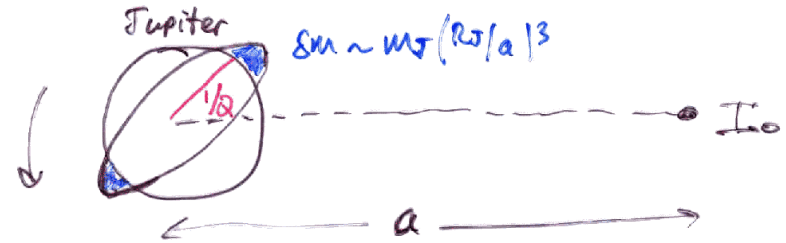


Tides in Gas Giant Planets

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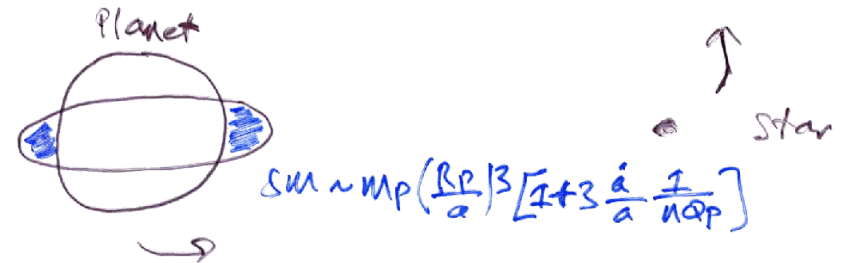
Orbit Evolution by Tides

- Expansion of Io's orbit



$$\text{torque} \sim \frac{GM_J}{R_J^2} \cdot R_J \cdot m_J \left[\frac{m_{Io}}{m_J} \left(\frac{R_J}{a} \right)^3 \right]^2 \frac{1}{Q_J}$$

- Circularization of extrasolar planet orbits



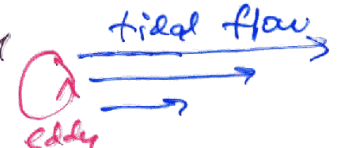
$$\frac{da}{dt} \sim - \frac{M_*}{m_p} \left(\frac{R_p}{a} \right)^5 \sqrt{\frac{GM_*}{a^3}} \frac{1}{Q_p}$$

Q for giant planets

$$Q \sim \frac{W_{\text{tide}} E_{\text{tide}}}{\dot{E}_{\text{tide}}} \quad \text{measures dissipation rate}$$

- tides push Io into resonance with Europa & Ganymede $\Rightarrow Q_J \lesssim 10^6$ (Goldreich & Soter)
- circularization of extrasolar planet orbits at 3 day orbits $\Rightarrow Q_p \lesssim 10^6$ (e.g. Marcy et al 1997)

Theory of Q

Hubbard: turbulent viscosity 

$$\dot{E}_{\text{viscous}} \sim E \frac{V}{R^2}$$

$$V \sim V_{\text{conv}} H \sim \left(\frac{E}{\rho}\right)^{1/3} \frac{\rho}{\rho_0} \sim \left(\frac{LR^4}{M}\right)^{1/3}$$

$$\Rightarrow Q \sim \left[\frac{\frac{GM^2}{R} \sqrt{\frac{GM}{R^3}}}{L} \right]^{1/3} \sim \underline{\underline{10^5}}$$

Goldreich & Nicholson: reduction factor

$$\tau_{\text{eddy}} \sim H/V_{\text{conv}} \sim \gamma R \quad \text{for large scale eddy}$$

$$P_{\text{tide}} \sim 5 \text{ hr}$$

\Rightarrow energy bearing eddies too slow!
Can't transport momentum.
Need small scale eddies ($V_e \sim e^{1/3}$).

$$V \sim V_{\text{conv}} H (\omega \tau_{\text{eddy}})^{-2} \sim \frac{L}{GM^2/R^3}$$

$$\Rightarrow Q \sim \frac{\frac{GM^2}{R} \sqrt{\frac{GM}{R}}}{L} \sim \underline{\underline{10^{15}}} \quad \text{!!}$$

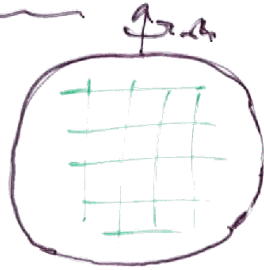
Disagrees with observations by $\sim 10^7$

Waves in giant planets

Convective core: inertial modes

$$\omega \approx 2\Omega_{\text{spin}} \frac{k_z}{k}$$

$\delta U \propto \rho^{-1/2}$ amplitude



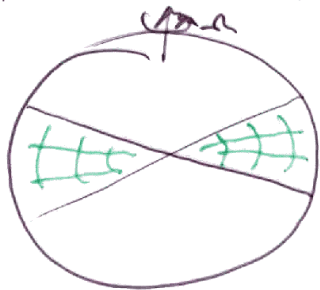
Radiative envelope: rotation modified g-modes

$$\omega^2 \approx 4\Omega_{\text{spin}}^2 \frac{k_z^2}{k^2} + N^2 \frac{k_{\text{hor}}^2}{k^2}$$

$$k_r \gg k_\theta, k_\phi, k_z \approx k_r \cos \theta \Rightarrow$$

$$k_{\text{hor}}^2 \approx k_r^2 [\omega^2 - 4\Omega_{\text{spin}}^2 \cos^2 \theta]$$

\Rightarrow horizontal propagation for $\cos \theta \leq \frac{\omega}{2\Omega_{\text{spin}}}$.



Amplitude $\delta U \propto \left(\frac{M}{\rho}\right)^{1/2}$.

How are tides different in a Rotating Star/Planet?

$$1) \delta P_m(r, \theta, \phi) = \sum_{l, m} \delta P_{lm}(r) Y_{lm}(\theta, \phi)$$

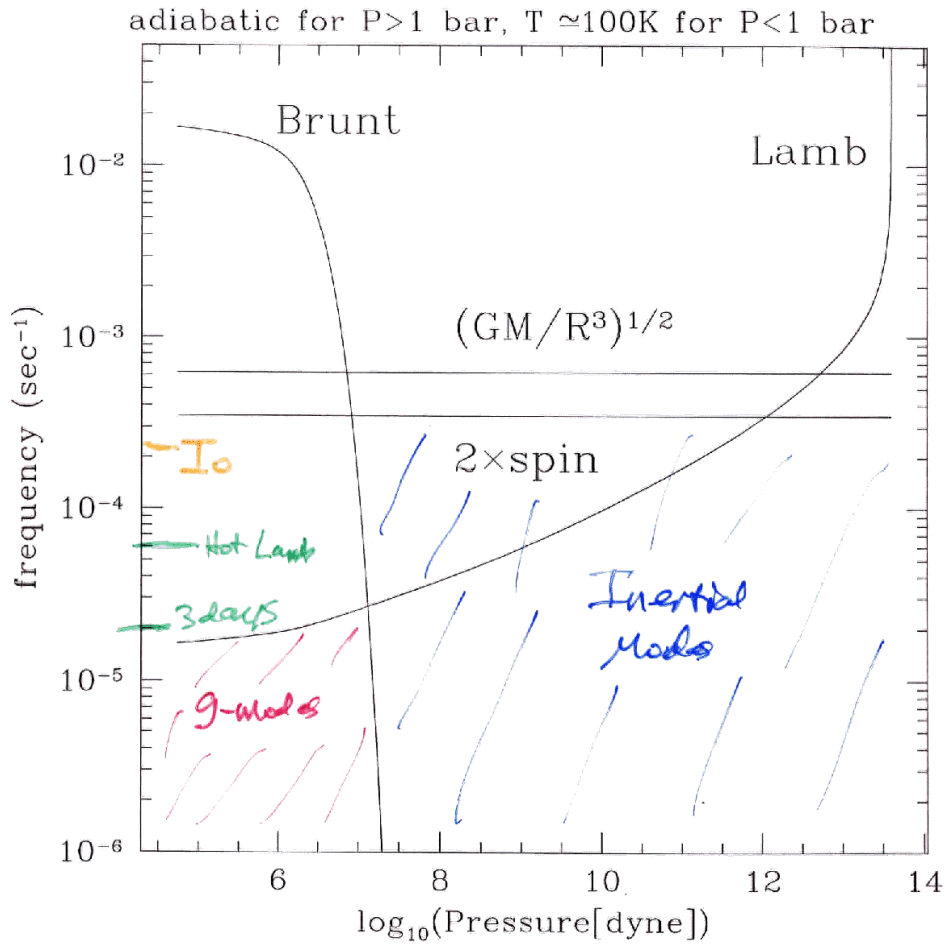
Coriolis force couples all spherical harmonics.

Tidal potential $\sim Y_{2m}$ overlaps with modes better (only $l=2$ in non-rotating star).

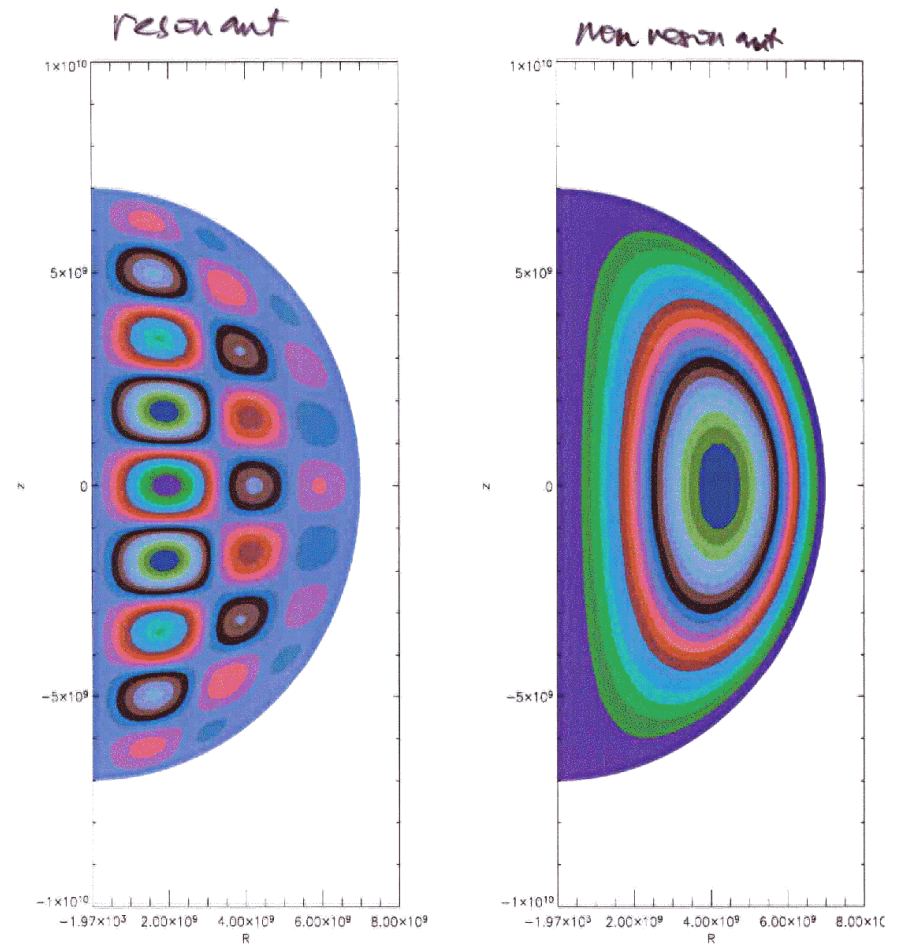
2) Can adjust two quantum #s to find good resonance (in non-rotating case, l & m fixed; only n can vary).

EX: $\omega = 2\Omega_{\text{spin}} \frac{k_z}{k_r}$

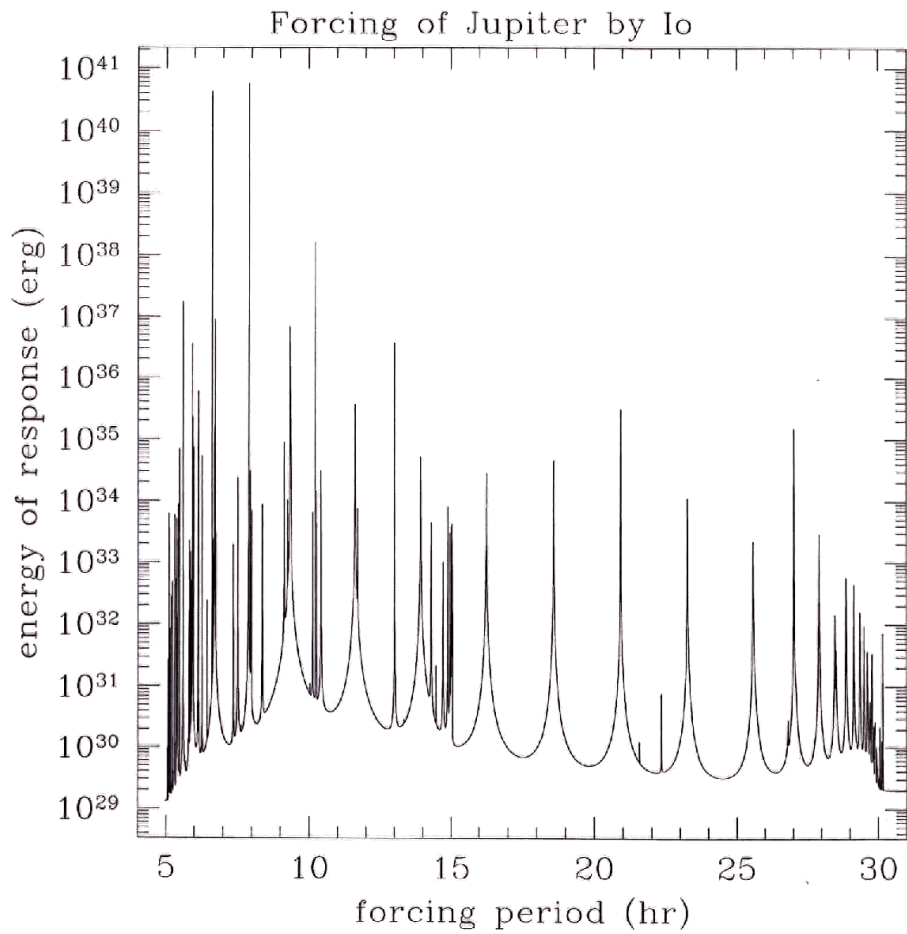
Where do waves propagate?



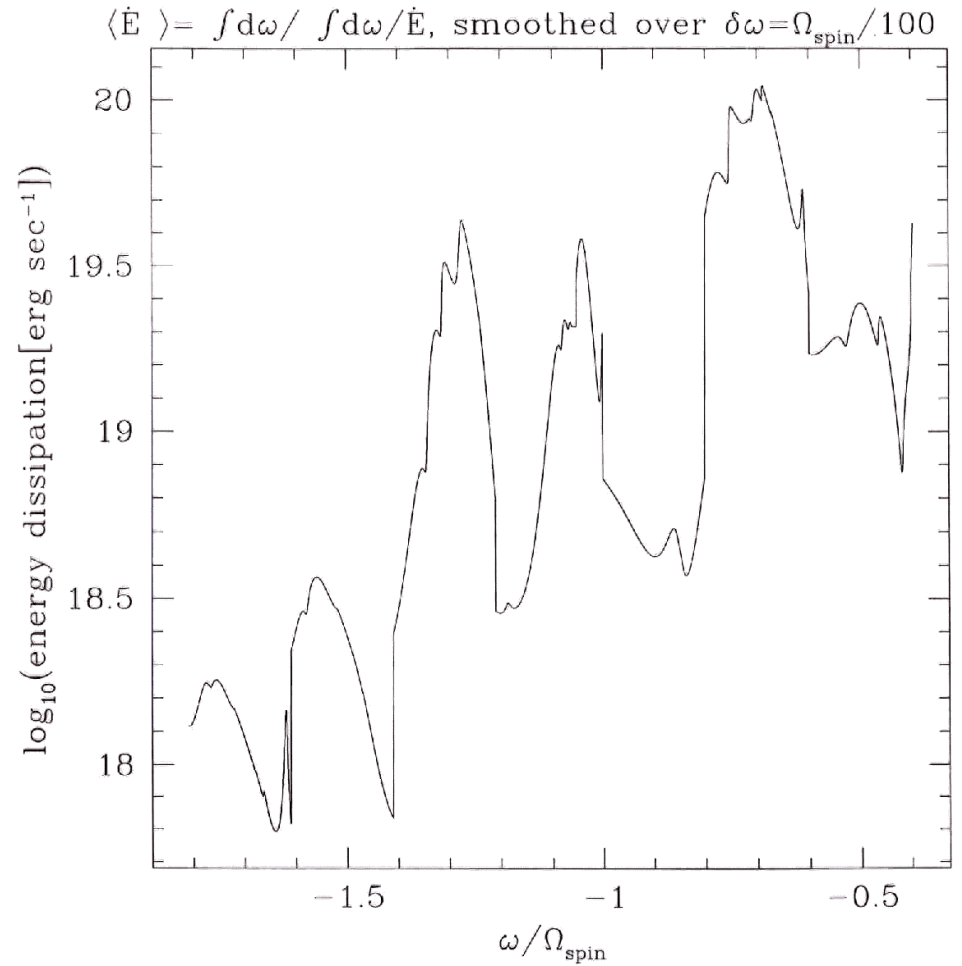
Tidal Response of Jupiter to Io



Plotted: Pressure perturbation



$N_r \approx 164, N_e = 8$



Estimate Q for turbulent viscosity

$$\text{torque} \sim \frac{\dot{E}}{\omega} \sim \left(\frac{m v_0}{m_J}\right)^2 \left(\frac{R_J}{a}\right)^6 \frac{GM_J^2}{R_J} Q^{-1}$$

$$\text{time spent in } (\omega, \omega + d\omega) \propto \dot{E}^{-1}$$

$$\Rightarrow \langle \dot{E} \rangle_{\text{time}} = \frac{\int \dot{E} d\omega}{\int d\omega / \dot{E}}$$

\Rightarrow off-resonant torque dominates.

Plug in #s from plot

$$\langle Q \rangle_{\text{turb}} \approx 10^{7-9} \frac{\langle \dot{E} \rangle}{10^{19-20} \text{ erg/sec}}$$

Needs more work!

Resolution, better Jup. models, etc

Apply to hot Jup. next.