

Migration type III





Migration Type III

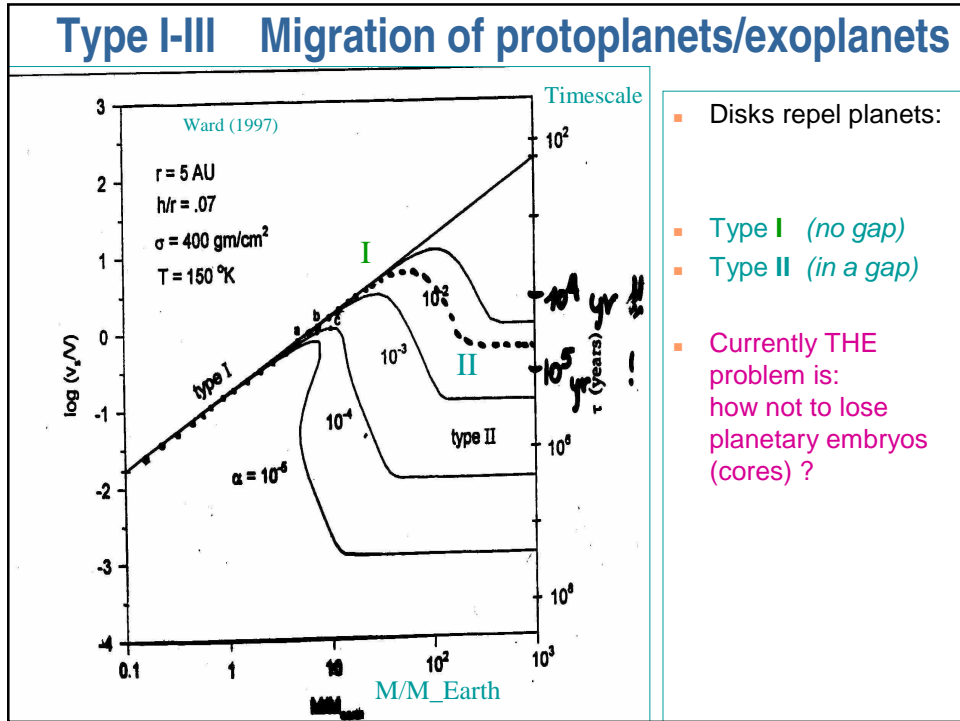
Pawel Artymowicz, Stockholm Obs.

Corotational vs. Lindblad torques
Slow migration, fast migration, ...
Is there type I and II migration?
What stops type III motion and where?

KITP, UCSE, 15 Mar 04

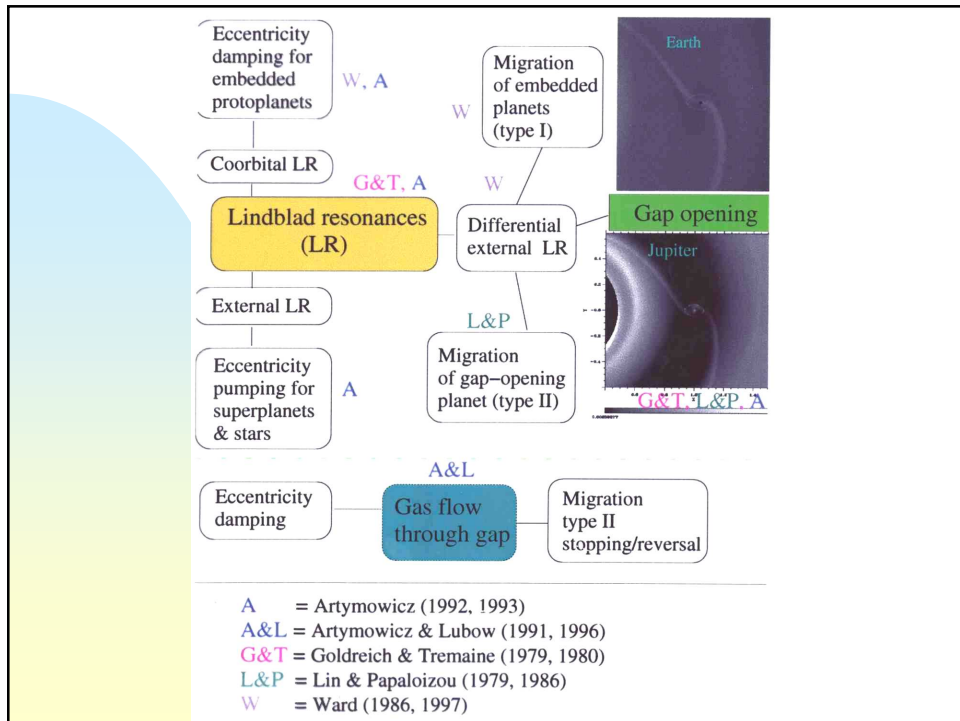
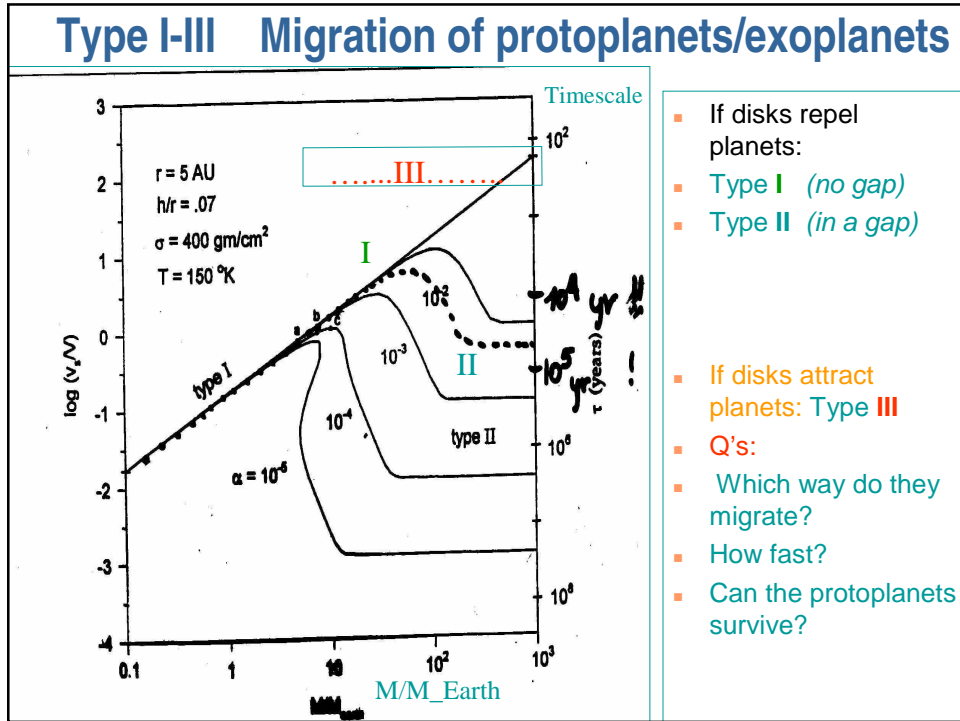
<p>Migration Type I : embedded</p>	<p>Migration Type II : more in the open</p>
	

Migration type III

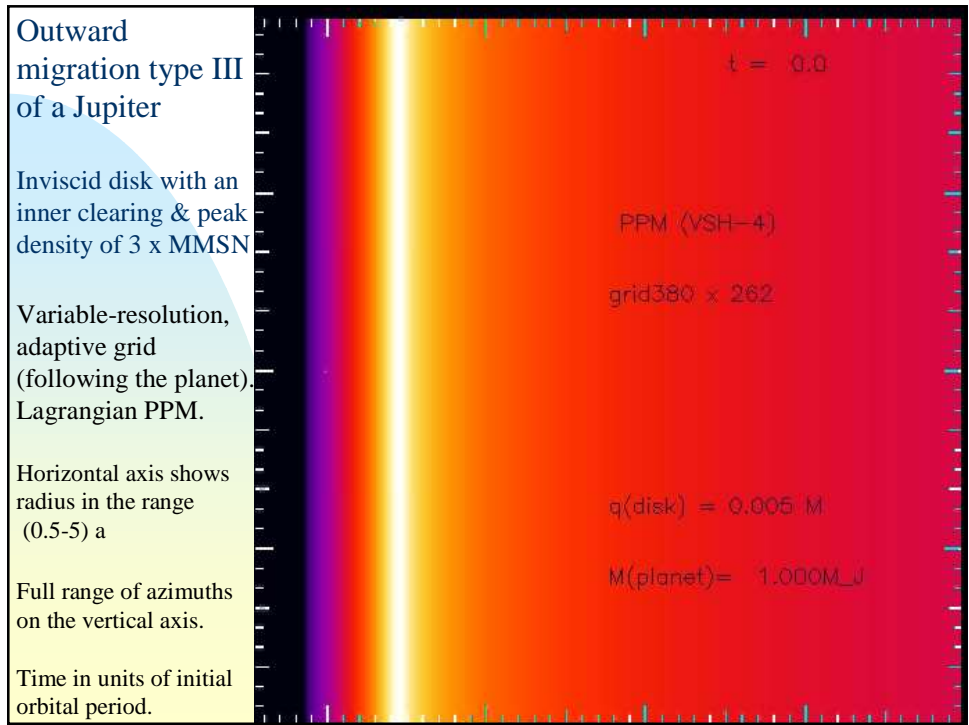
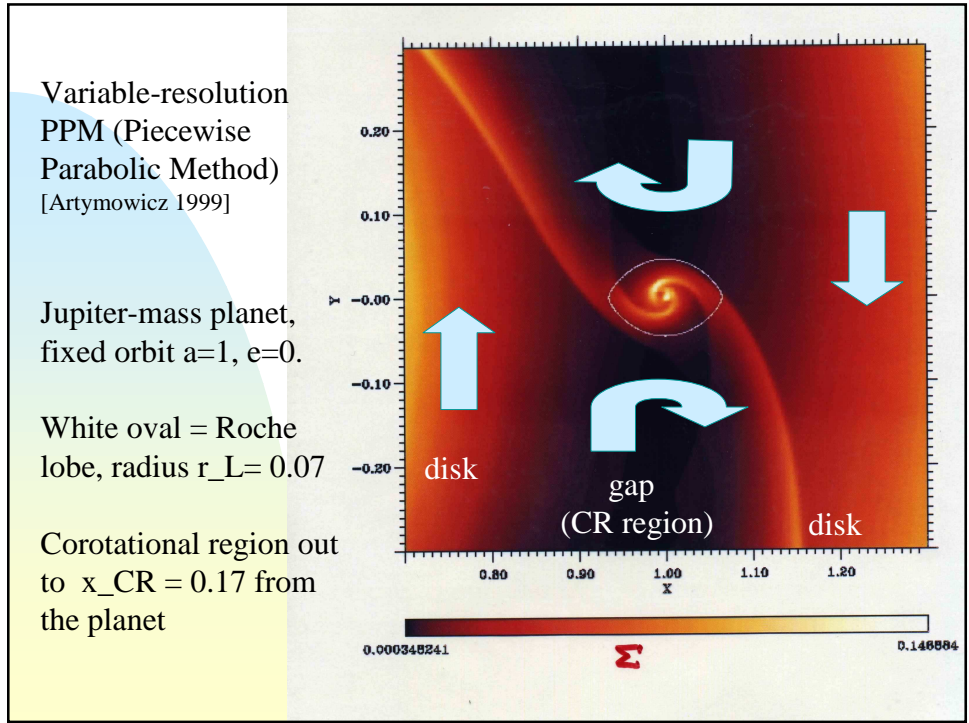


Migration:	Interaction:	Timescale of migration:
type 0	Gas drag + Radiation press.	from $\sim 1e2$ yr to disk lifetime (up to $1e7$ yr)
type I	Resonant excitation of waves (LR)	$> 1e4$ yr
type II & IIb	Tidal excitation of waves (LR)	$> 1e5$ yr
type III	Corotational flows (CR)	$> 1e2 - 1e3$ yr
N-body	Gravity	$> 1e5$ yr (?)

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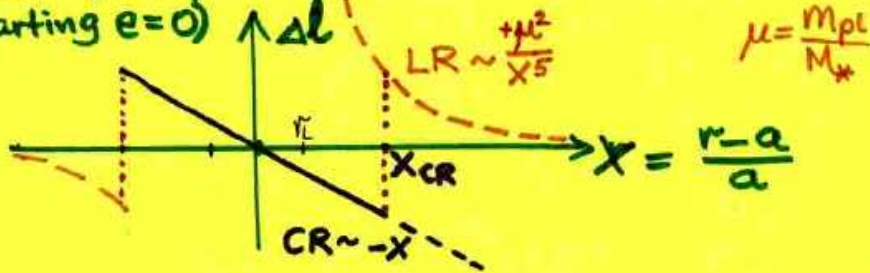
What is more important: Lindblad Resonances & waves or Corotation?

$$X_{CR} =: \zeta r_L \quad X_{CR} = r_L = \left(\frac{\mu}{3}\right)^{\frac{1}{3}}$$

Why is ζ important?

Competition between CR & LRs.

Angular momentum gain in one passage near the planet (starting $e=0$)



Shearing sheet (disk) provides flux of incoming matter $\sim \Sigma \left| \frac{3}{2} \Omega x \right|$, therefore a one-sided torque estimation gives

$$T_{CR}^{(1)} \sim \int_0^{X_{CR}} (-x) |x| \Sigma dx \sim -\Sigma X_{CR}^3$$

$$T_{LR}^{(1)} \sim \int_{X_{CR}}^{+\infty} \frac{\mu^2}{X^5} |x| \Sigma dx \sim \mu^2 \Sigma X_{CR}^{-3} \sim$$

$$T_{CR}^{(1)} : T_{LR}^{(1)} \sim 10 : 1 \quad \sim \Sigma X_{CR}^3 \zeta^{-6}$$

(cf. R3B results of Ida & Lin 2001)

Levels of description :

symmetric switch model

- at conjunction: $a \pm x \mapsto a \mp x$ if $|x| \leq x_{CR}$
- $x_{CR} = ?$
- fairly good kinem. & timing of global flow

guiding center approximation to Hill's or R3B problem

- neglect epicyclic motion & pressure, viscosity, ...
- $x_{CR} \approx 2.5 r_L$
- improved kinem.

numerical hydrodynamics

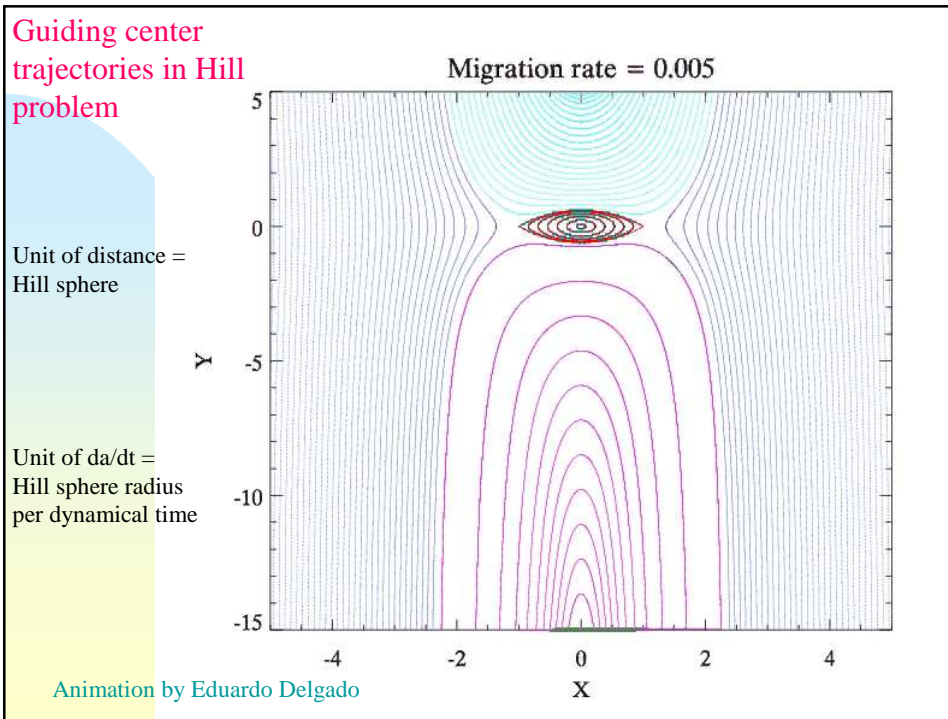
- maximum physics, but
- many difficult issues (what's real? what's numerical? why differences?)
- $x_{CR} \approx 2.5 r_L$
- $\Sigma(r, \theta)$

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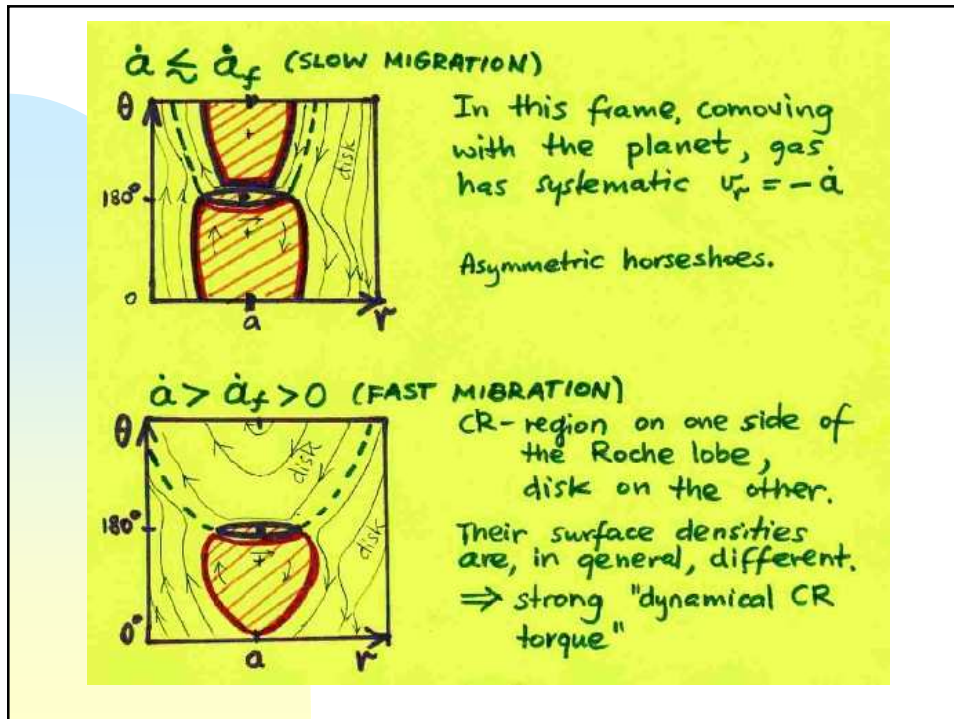
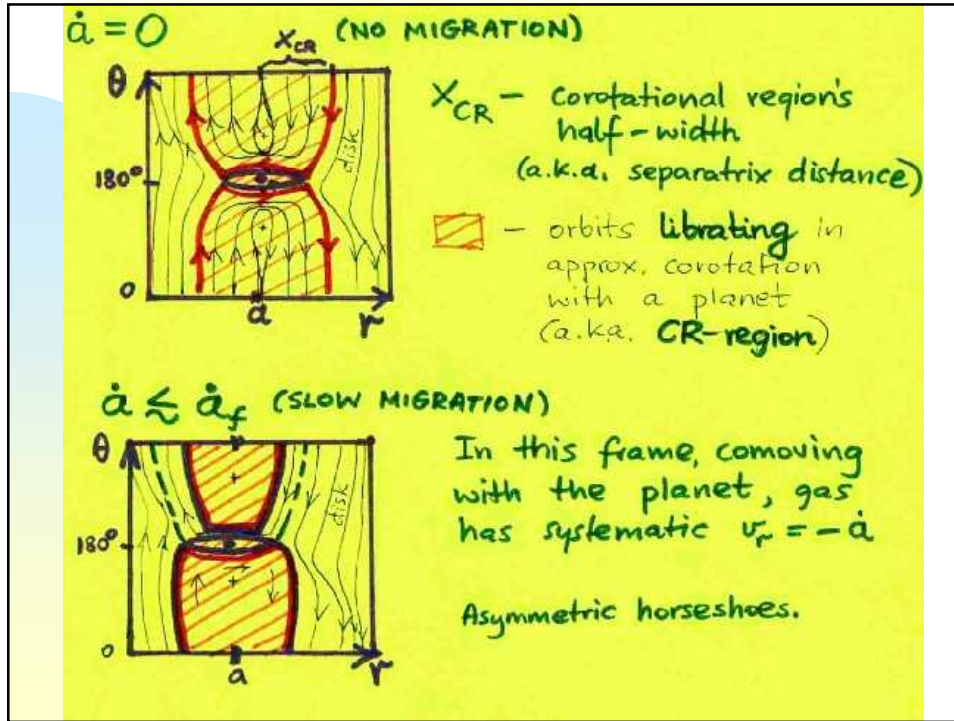
Toy models, regularized theories

looking at orbits and separatrices rather than doing a questionable Fourier decomposition + infinitesimal perturbations theory

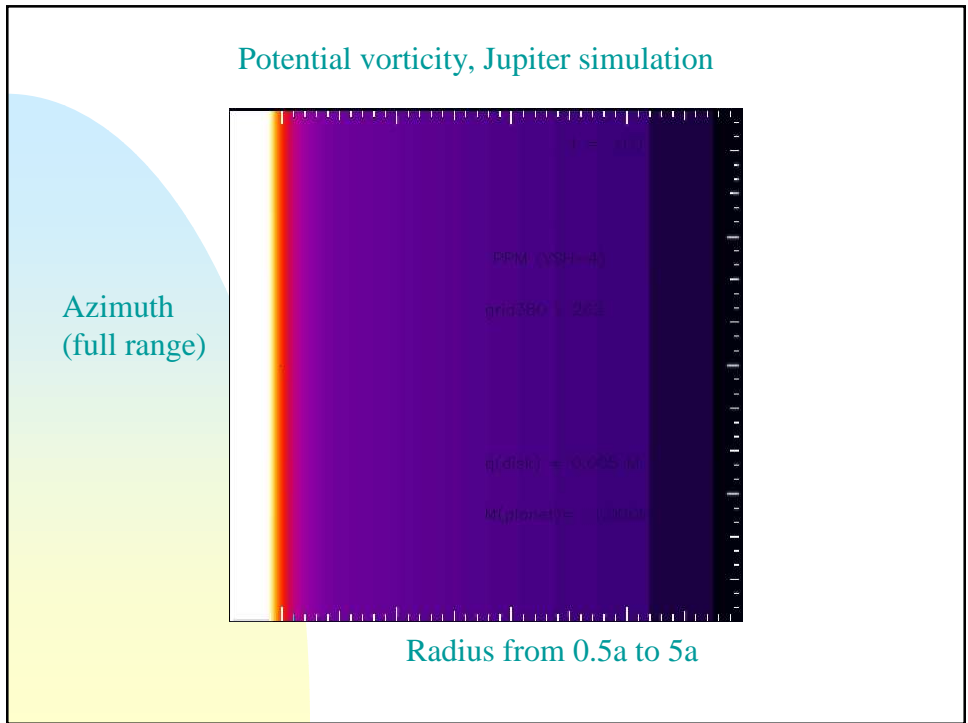
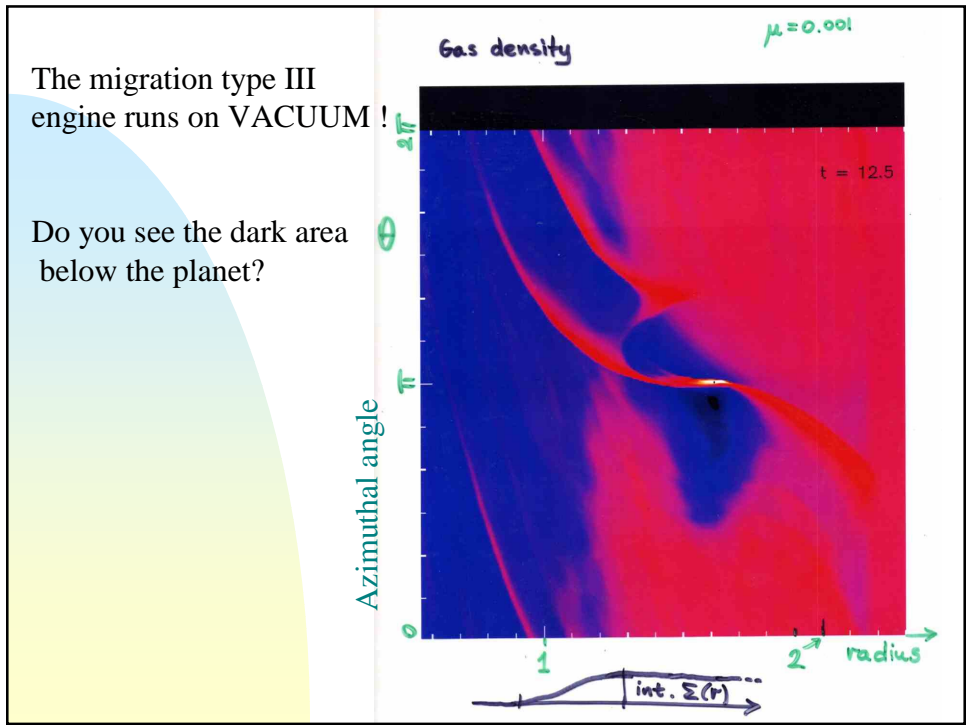
- Capture essentials!
- Give results resembling full 2-d hydrodynamics



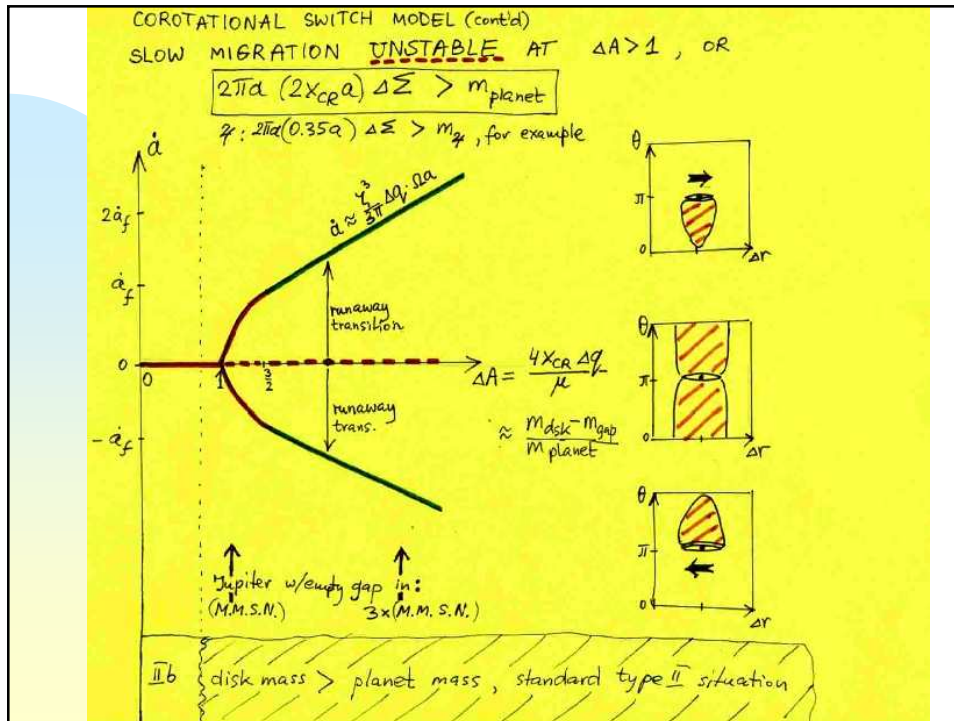
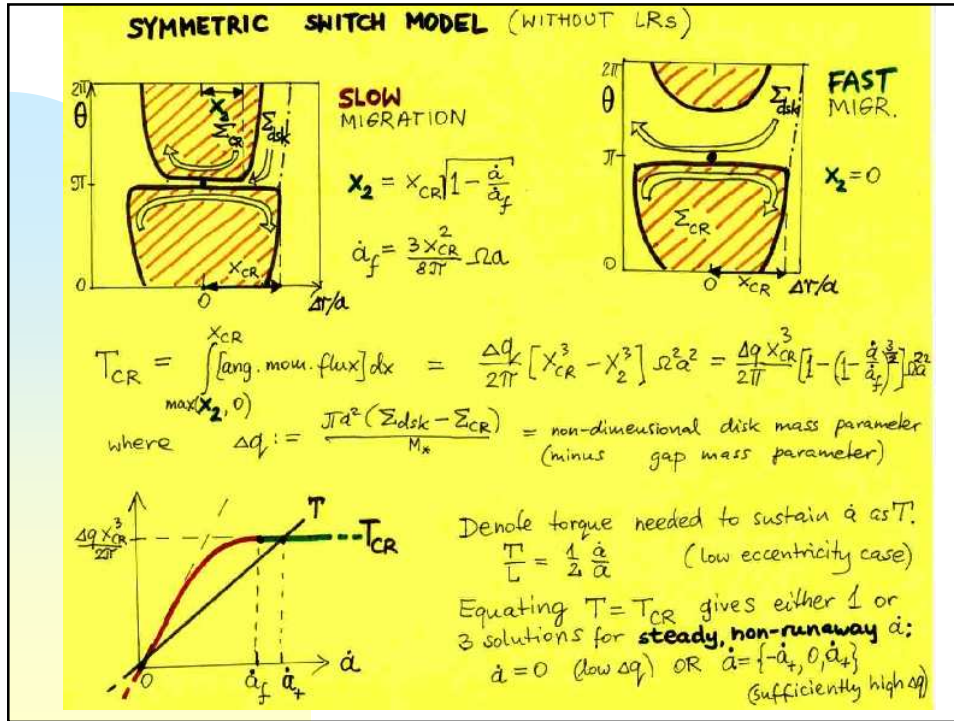
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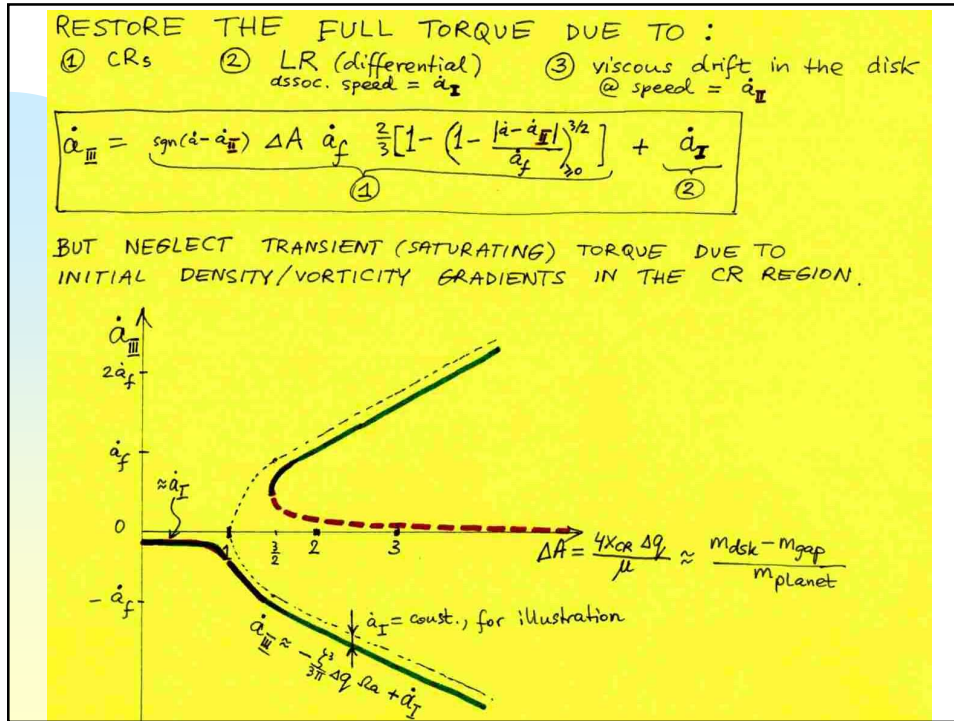
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TYPE III vs. TYPE II

- TYPE III wins if a planet catches an underdense bubble (balloon?) via initial conditions, or carves a gap.
- TYPE II \rightarrow TYPE III seen in numerical work.
- THE ONLY WAY TO RESTORE STANDARD TYPE II (II b ?) : ENSURE A VERY WIDE GAP OPENS, VERY LOW Σ_{disk} (near x_{CR})

TYPE III vs. TYPE I

- For $\mu \lesssim 20 \mu_{\oplus}$, the issue is vertical averaging.
 $x_{CR} < h = \frac{c}{\alpha a}$ (disk scale height). Also gap filling: $\Delta \rho \sim x_{CR} \rho_{\oplus}^{\downarrow}$
- But then, in 2-D, $\frac{\dot{a}_{III}}{\dot{a}_{LR(I)}} \approx \frac{T_{CR}}{T_{LR}} \sim 2.5 \left(\frac{x_{CR}}{h} \right)^2 \xrightarrow{\mu \rightarrow 0} \infty$
- 3-D, expect $\sim \left(\frac{x_{CR}}{h} \right)^{-1} \rightarrow \infty$, STILL!
- For $\mu \gtrsim 10 \mu_{\oplus}$, numerics shows type III, but more fragile than for Jupiters (little gap-opening).

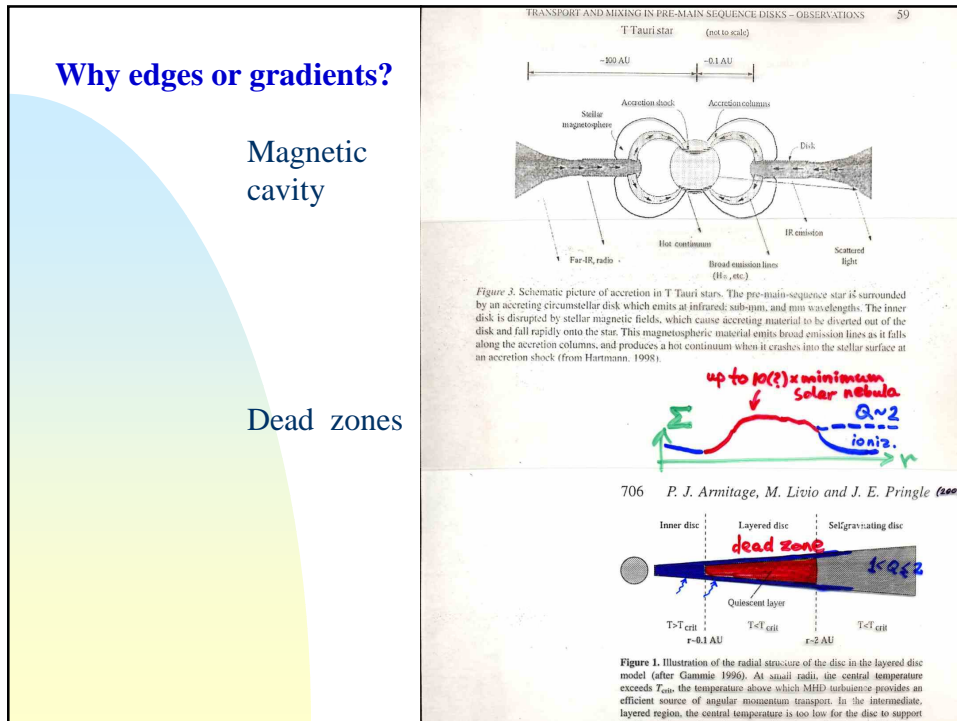
Next Steps: Toward a better LR/CR perturbation theory

- Previous perturbation theories started from circular unperturbed orbits [do not exist]
- and assumed infinitesimal perturbations (Fourier decomposition allowed) [not always!]
- Alternative way: unperturbed state adjusted for perturbation. Trajectories of all essential types (disk orbits, corotational hairpin/horeseshoes, closed orbits around planet)
- On that set of unperturbed flow lines, 1st order perturbation should give a better approximation
- Migration and additional effects can be incorporated

Are there ANY SURVIVORS
of type III migration?!

Of course...

Migration type III

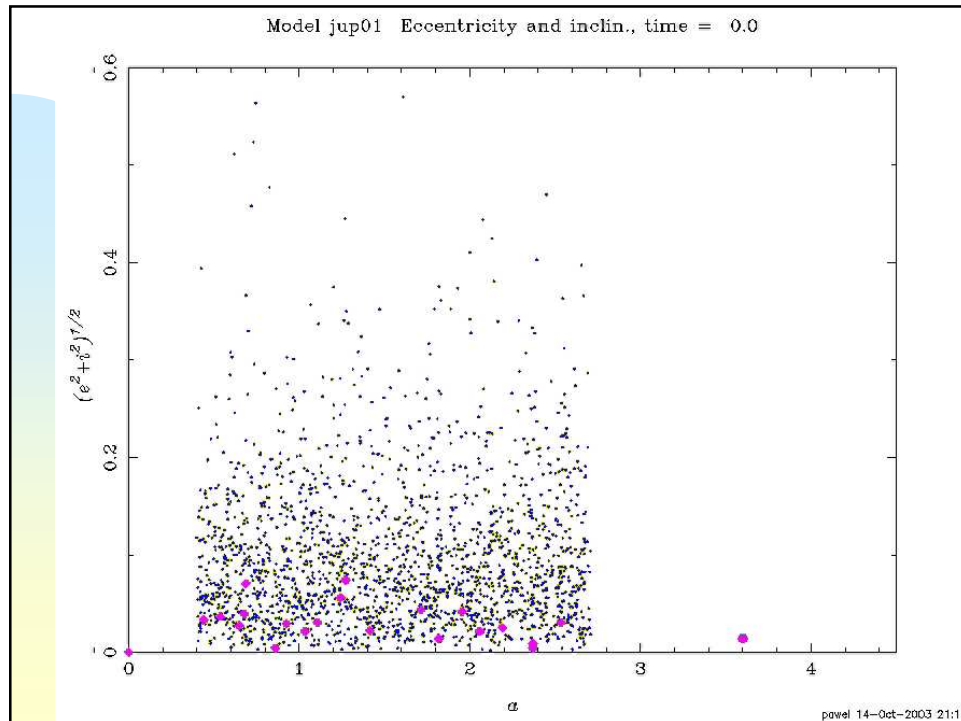


Unsolved problem of the Last Mohican scenario of planet survival in the solar system:

Can the terrestrial zone survive a passage of a giant planet?

- N-body simulations, $N \sim 1000$ (Edgar & Artymowicz 2004)
- A quiet disk of sub-Earth mass bodies reacts to the rapid passage of a much larger protoplanet (migration speed = input parameter).
- Results show increase of velocity dispersion/inclinations and limited reshuffling of material in the terrestrial zone.
- Migration type III too fast to trap bodies in mean-motion resonances and push them toward the star
- Evidence of the passage can be obliterated by gas drag on the time scale \ll Myr --- passage of a pre-jupiter planet(s) not excluded dynamically.

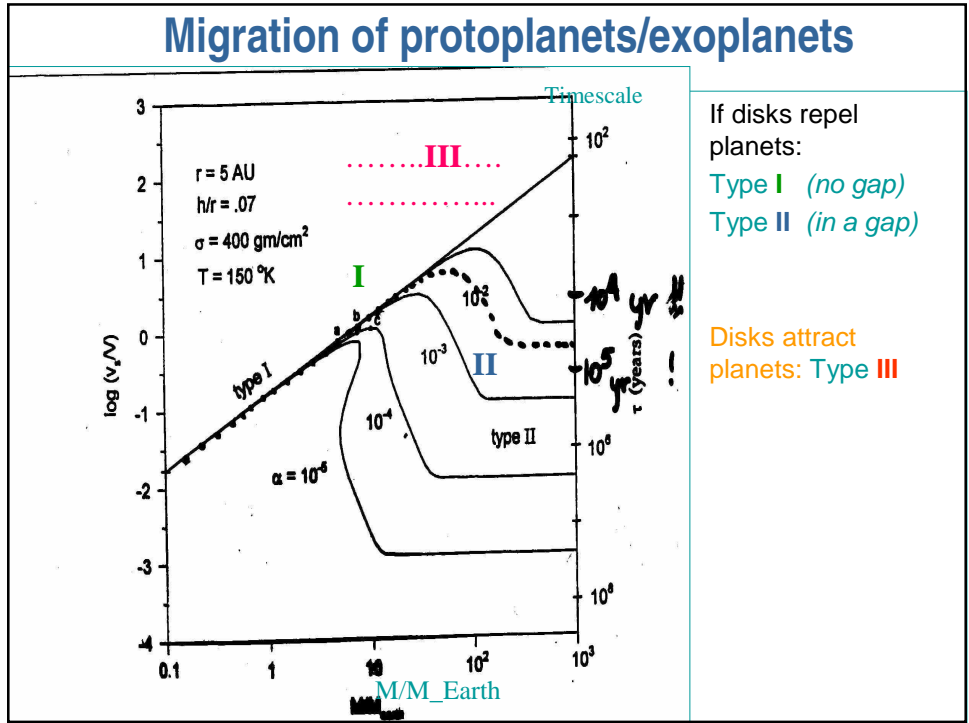
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Summary of type-III migration

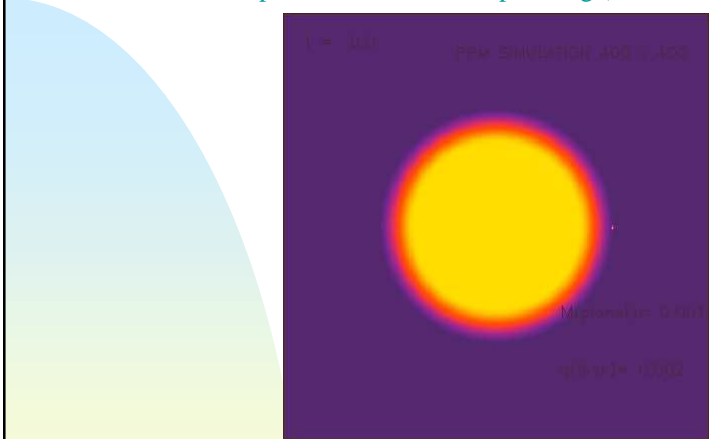
- New type, sometimes extremely rapid (timescale < 1000 years). CRs \gg LRs
- Direction depends on prior history, not just on disk properties.
- Supersedes a much slower, standard type-II migration in disks more massive than planets
- Preliminary results: modifies or replaces type-I migration
- Very sensitive to disk density (or vortensity) gradients.
- Migration stops on disk features (rings, edges and/or substantial density gradients.) Such edges seem natural (dead zone boundaries, magnetospheric inner disk cavities, formation-caused radial disk structure)
- Offers possibility of survival of giant planets at intermediate distances (0.1 - 1 AU),
- ...and of terrestrial planets during the passage of a giant planet on its way to the star.
- If type I superseded by type III then these conclusions apply to cores as well, not only giant protoplanets.

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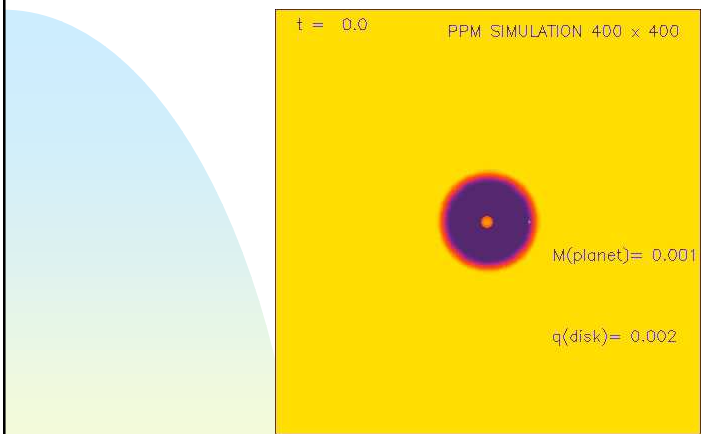
Consider a one-sided disk (**inner disk only**). The rapid inward migration is **OPPOSITE** to the expectation based on shepherding (**Lindblad resonances**).



The diagram on the left shows a cross-section of a disk with a blue outer region and a yellow-green inner region. The PPM simulation on the right shows a central yellow-orange disk with a small white planet. Text in the simulation includes: $t = 0.0$, PPM SIMULATION 400 x 400, Migration (cm), and $q(\text{disk}) = 0.002$.

Like in the well-known problem of “sinking satellites” (small satellite galaxies merging with the target disk galaxies), **Corotational torques** cause rapid inward sinking. (Gas is transferred from orbits inside the perturber to the outside. To conserve angular momentum, satellite moves in.)

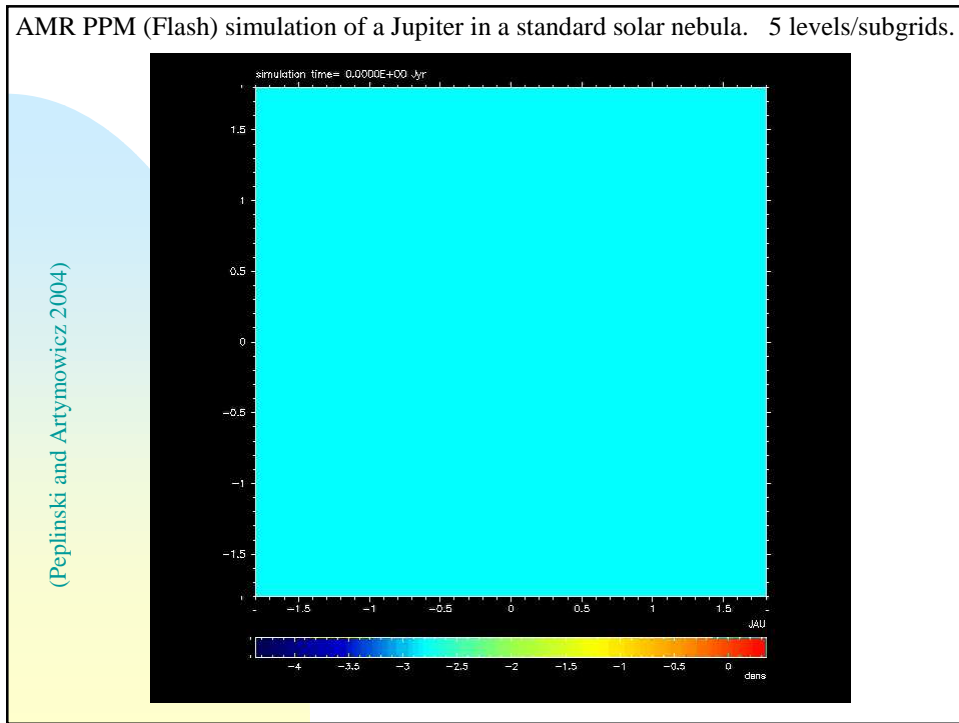
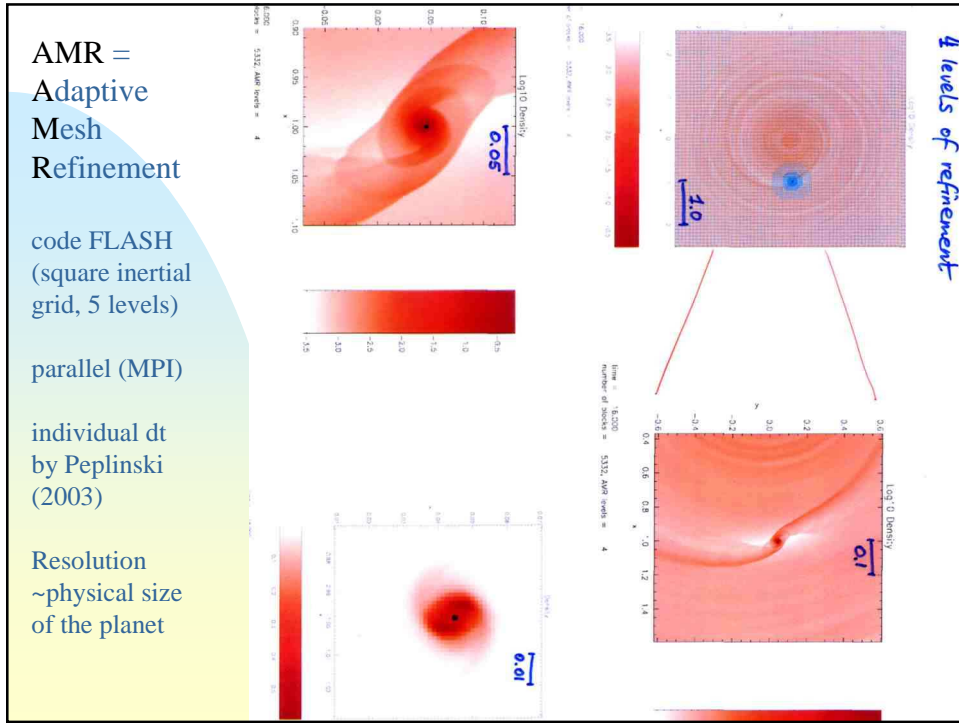
Now consider the opposite case of an **inner hole** in the disk. Unlike in the shepherding case, the planet rapidly migrates **outwards**.



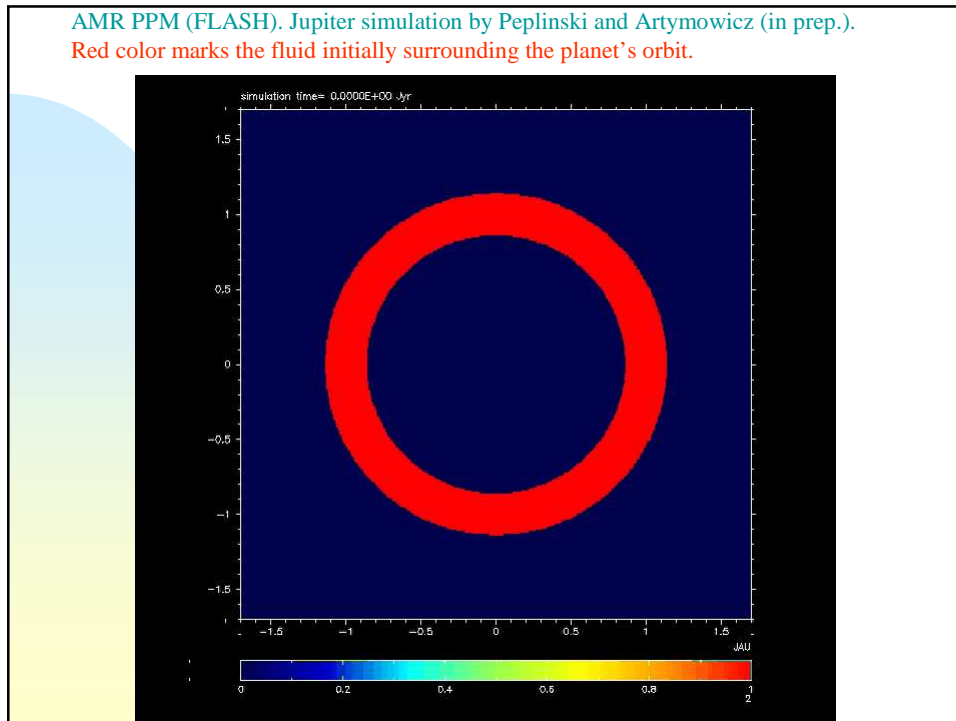
The diagram on the left shows a cross-section of a disk with a blue outer region and a yellow-green inner region. The PPM simulation on the right shows a central yellow-orange disk with a small white planet. Text in the simulation includes: $t = 0.0$, PPM SIMULATION 400 x 400, $M(\text{planet}) = 0.001$, and $q(\text{disk}) = 0.002$.

Here, the situation is an inward-outward reflection of the sinking satellite problem. Disk gas traveling on hairpin (half-horseshoe) orbits fills the inner void and moves the planet out rapidly (type III outward migration). **Lindblad resonances** produce spiral waves and try to move the planet in, but lose with **CR torques**.

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Numerical Troubles:

resolution

grav. Softening and zones where torques are ignored

self-gravity of gas (missing)

gas heating (and other effects)

the usual troubles: boundary conditions, instabilities, unexplained crashes,

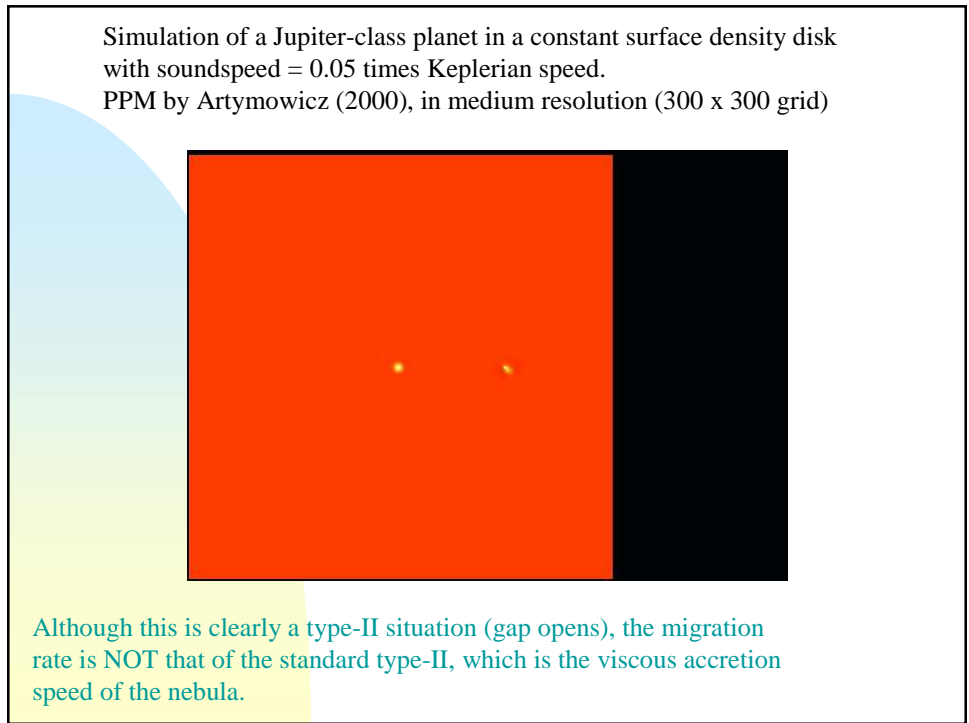
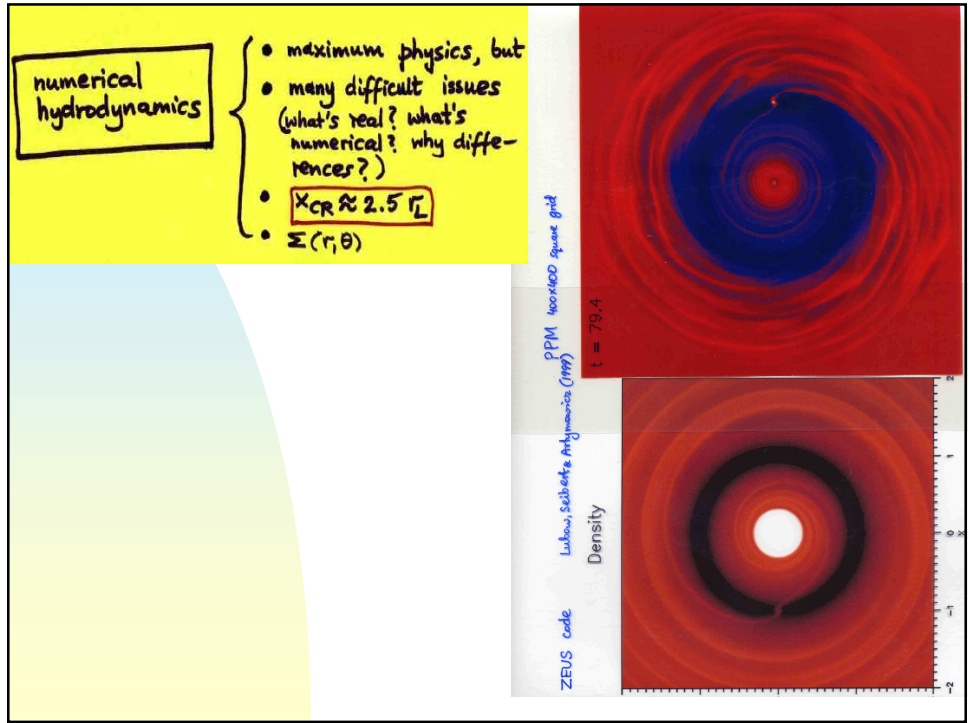
the unusual troubles: extreme vortex production and/or variability of flow in some codes

---> **Comparison or Test Problem**

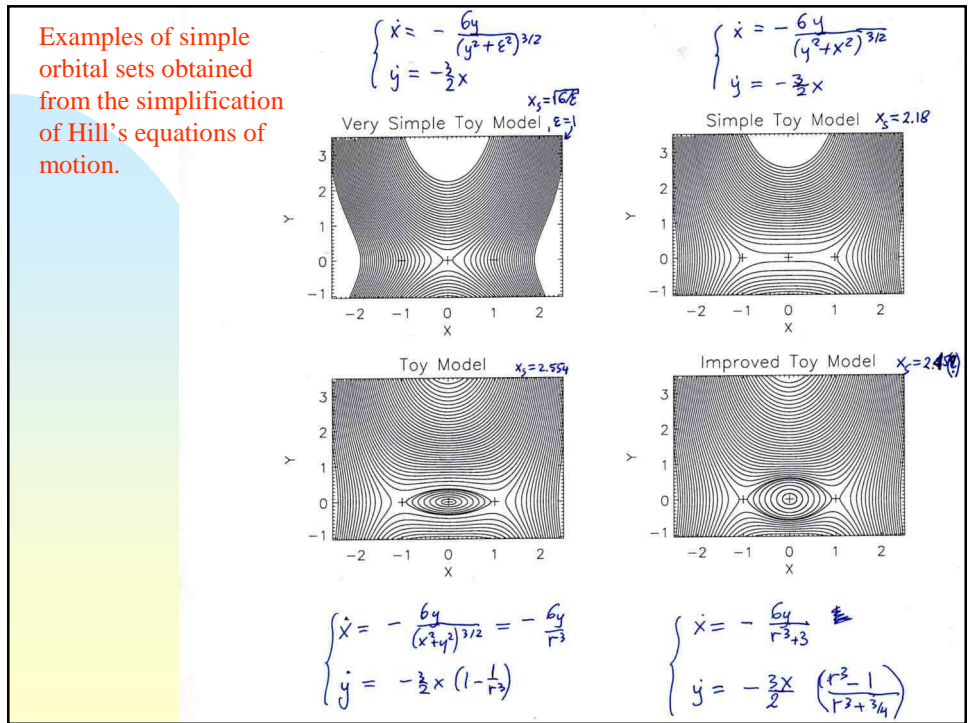
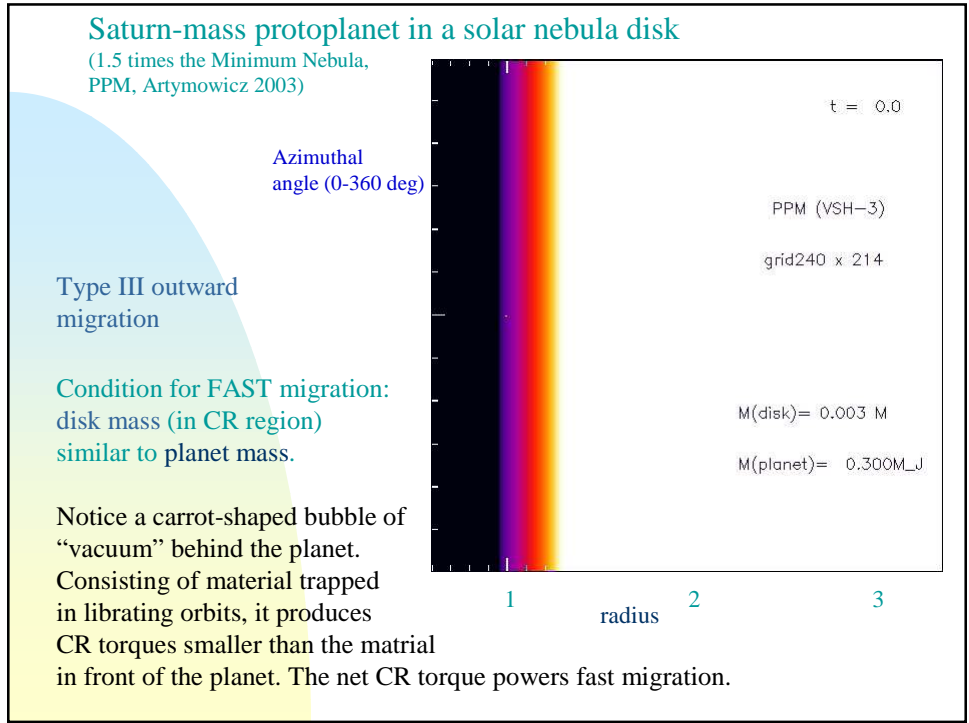
mini-workshop scheduled for May 2004 in Stockholm (EU Network on Planets)

(www.astro.su.se/~pawel/planets/test.hydro.html)

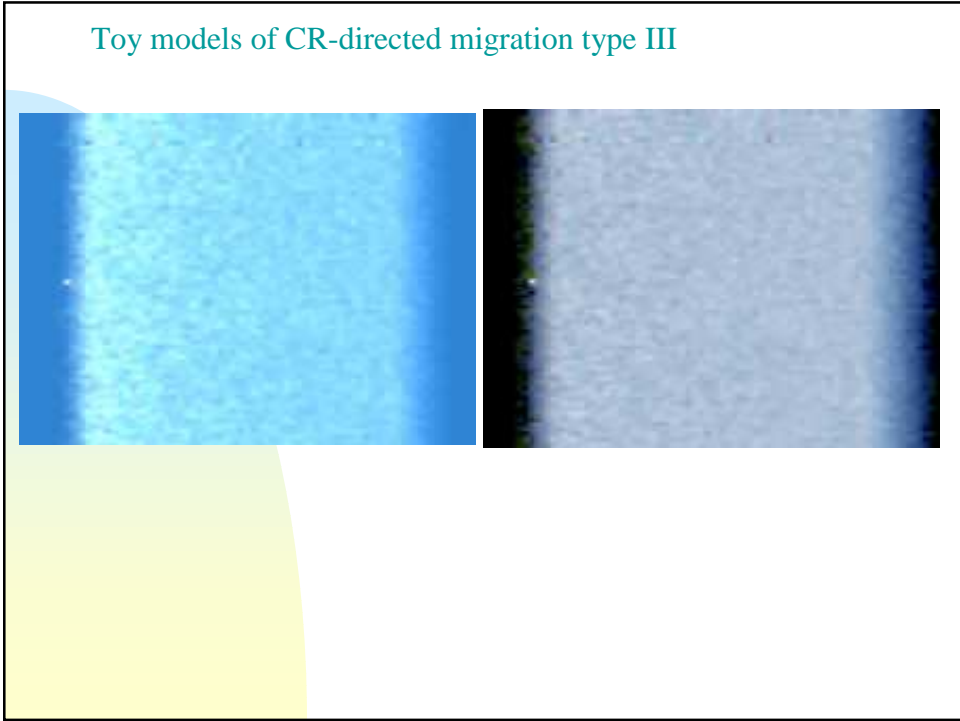
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**WHAT IS MORE IMPORTANT:
LR or CR?**

ESTIMATE the so-called ONE-SIDED TORQUES:

$$\frac{T_{LR}}{T_{CR}} \Big|_{\text{separatrix}} \approx -\frac{\delta x}{2x} \approx -\frac{1}{10}$$

⇒ Corotational torques potentially much more important!

MORE ELABORATE THEORY INTEGRATES $\frac{dT}{dr}$ functions, which look like this:

$x_m \approx \% \cdot r_c, x \approx 2.2$

$$\frac{\dot{a}}{\alpha} = \frac{2(T_{CR} - T_{LR})}{L_{pl}} = \frac{c_s}{3\pi} \left(\frac{\rho_0 \Sigma}{M} \right) \left(x^2 - \frac{6}{3x^3} \right)$$

$+10.6 - 1.4 = +9.2$

~ reproduces Ida & Lin (2001) cold disk limit of +8.8