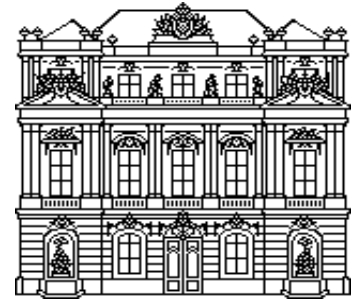


# Polar Molecules: QIPC & Condensed Matter Physics



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G. Pupillo

In collaboration with:

Harvard: J. Doyle, M Lukin  
Yale: D. DeMille, R. Schoelkopf

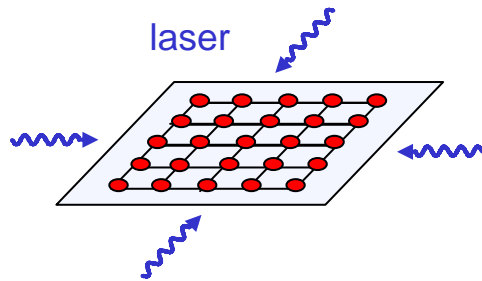
**SFB**

*Coherent Control of  
Quantum Systems*

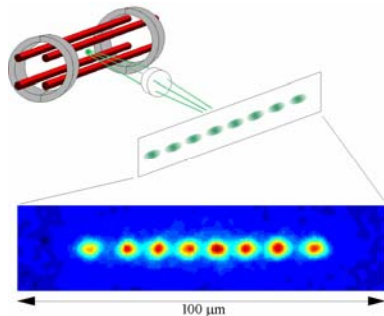
**EU networks**

# Atoms & Ions

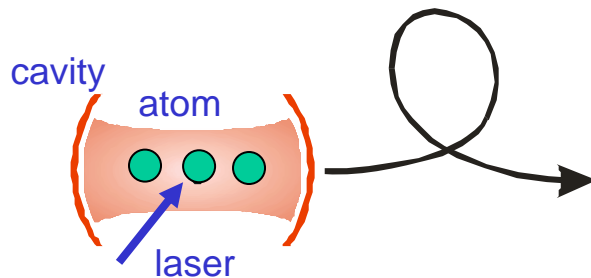
- cold atoms in optical lattices



- trapped ions / crystals of ...

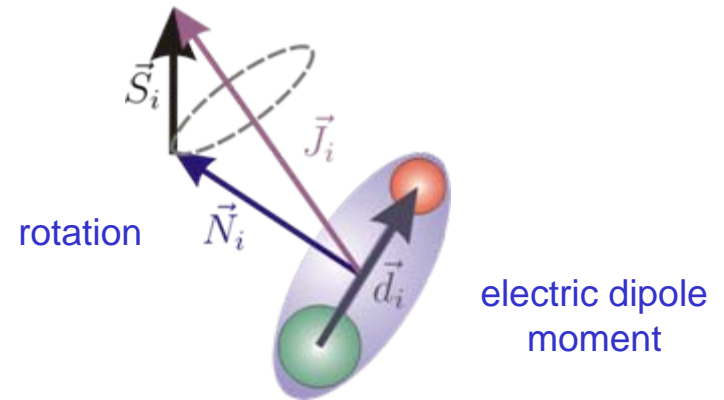


- CQED



- atomic ensembles

# Polar Molecules



- single molecules / molecular ensembles
- coupling to optical & microwave fields
  - trapping / cooling
  - CQED (strong coupling)
  - spontaneous emission / engineered dissipation
- interfacing solid state / AMO & microwave / optical
  - strong coupling / dissipation
- superfluid – crystalline phases
  - cond mat / quantum info
  - dephasing

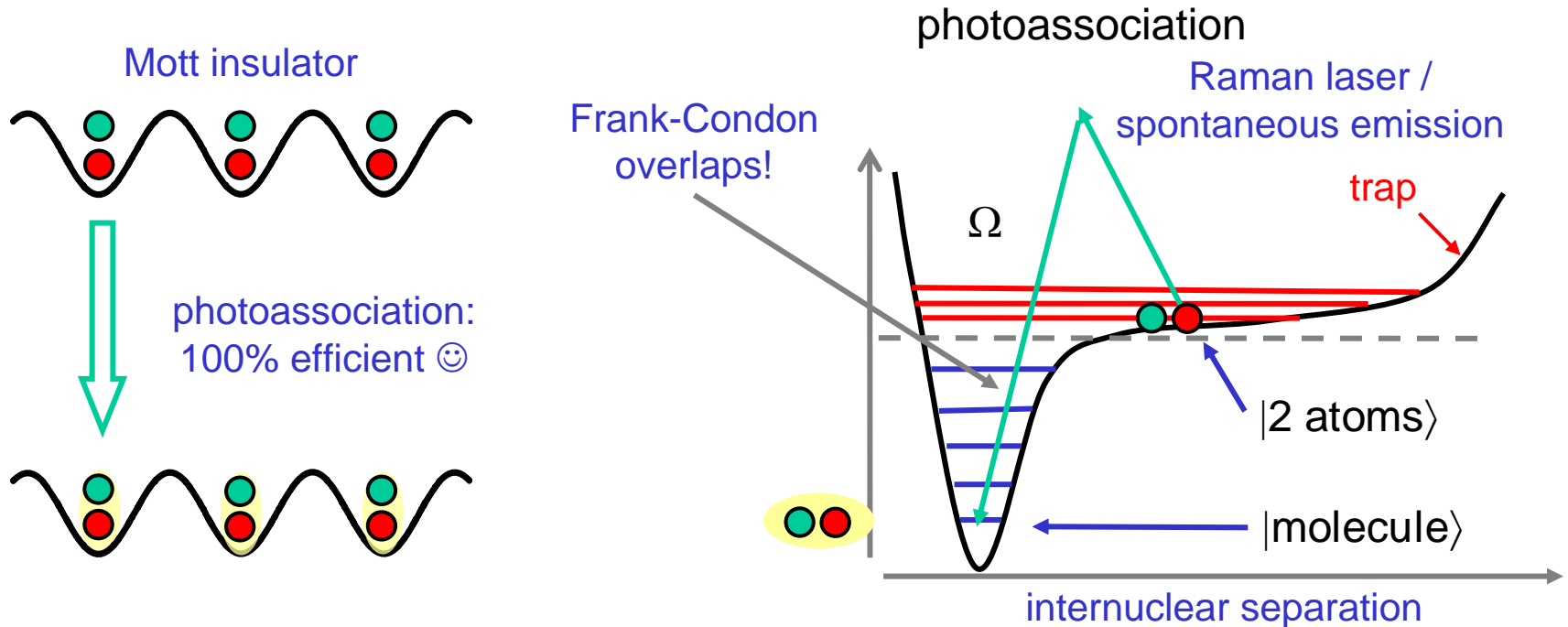
## Polar Molecules I (basic features)

- rotational spectra of single molecules
  - ... spins, and how to encode qubits
- dipole-dipole interactions

# Polar Molecules

- We consider polar molecules *in electronic and vibrational ground state*
  - cooling and trapping techniques being developed
  - preparation via ...
    - photoassociation (e.g. from two-species BEC)
    - buffer gas cooling

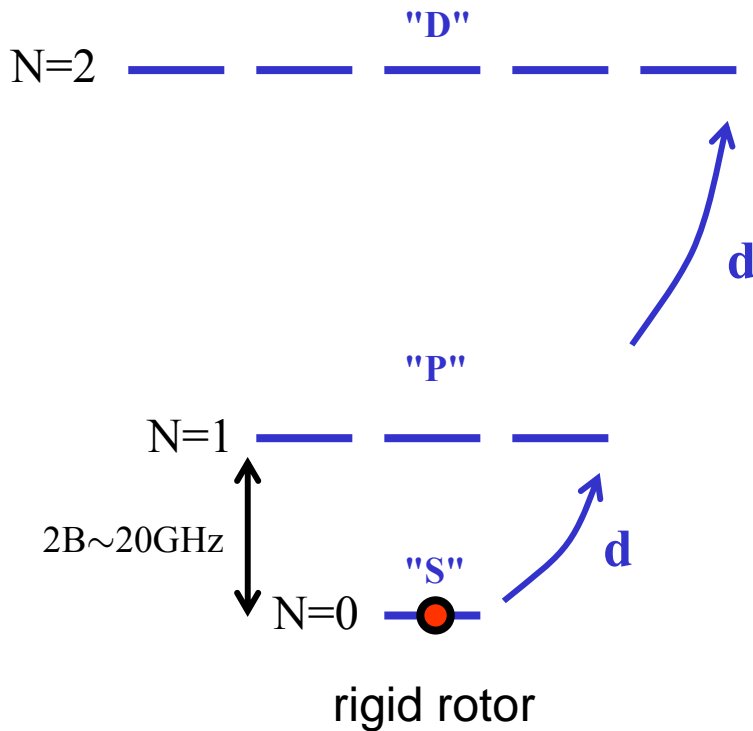
exp: all cold atom labs  
exp: Demille, Doyle,  
Mejer, Rempe, Ye ...



- Cold polar molecules will exist in the lab in a few years.

# 1a. Single Polar Molecule: rigid rotor

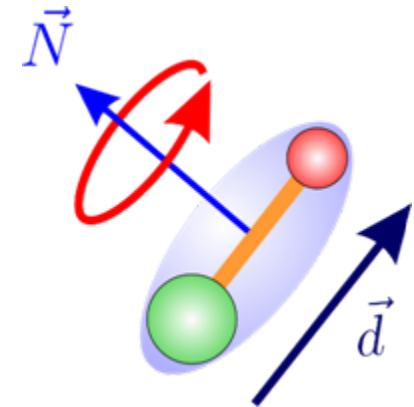
- single heteronuclear molecule



$$H = B N^2$$

$$|N, M_N\rangle$$

$$E_N = B N(N+1)$$



- ✓ dipole-moment  $d \sim 10$  Debye ☺
- ✓ rotation  $B \sim 10$  GHz ... anharmonic ☺
- ✓ essentially no spontaneous emission  $10^{-3}$  Hz ☺  
black-body scatt.rate  $\sim 1$  Hz ☺  
i.e. excited states "useable" ☺

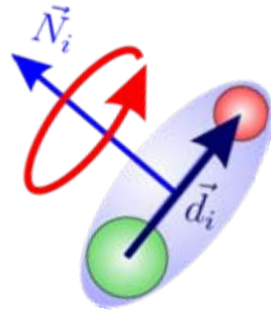
# 1b. Single Polar molecule: adding spin rotation coupling

- rigid rotor

$$H = B N^2$$

$$|N, M_N\rangle$$

$$E_N = B N(N+1)$$

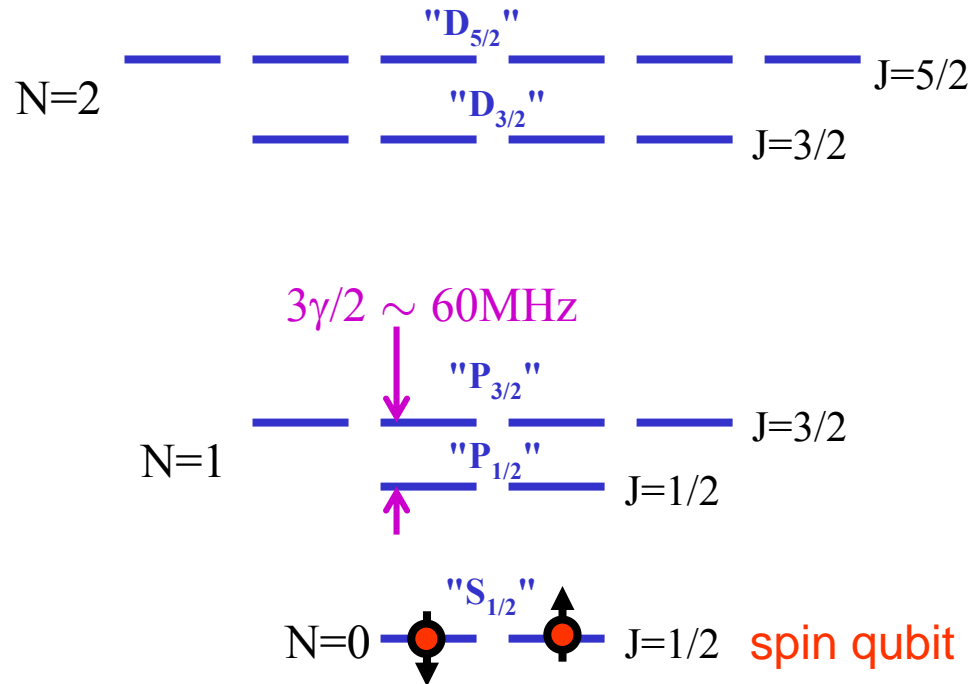
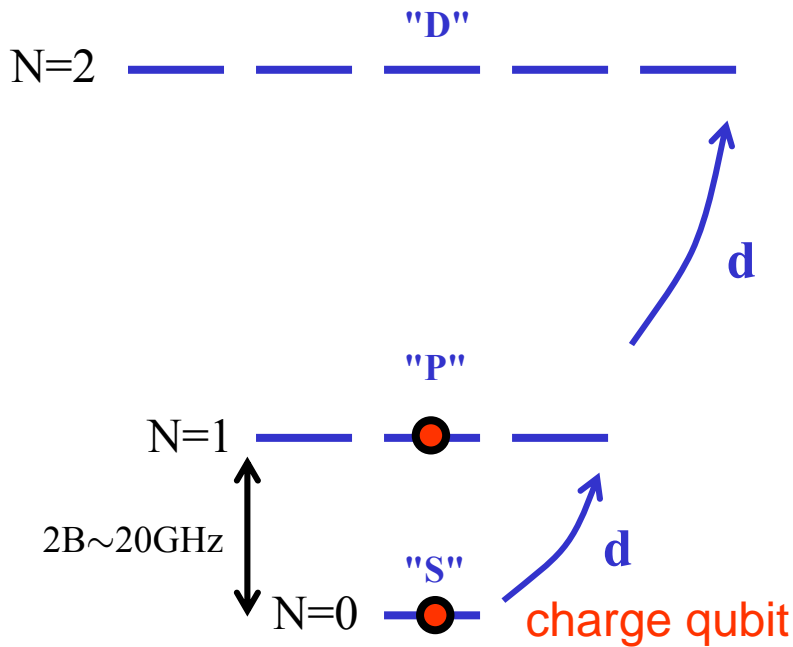
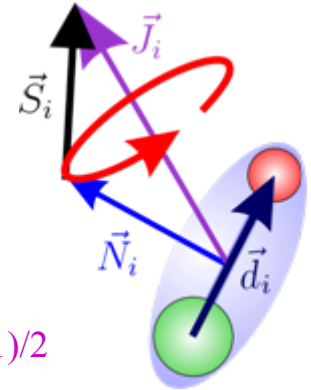


- adding spin-rotation coupling ( $S=1/2$ )

$$H = B N^2 + \gamma \mathbf{N} \cdot \mathbf{S}$$

$$|N, J, M_J\rangle \quad (J=|N \pm 1/2|)$$

$$E_{N,J=N \pm 1/2} = B N(N+1) + \begin{cases} +\gamma N/2 \\ -\gamma(N+1)/2 \end{cases}$$



- How to encode qubits?

“looks like an Alkali atom on GHz scale”  
(we adopt this below as our model molecule)

## 2. Two Polar Molecules: dipole – dipole interaction

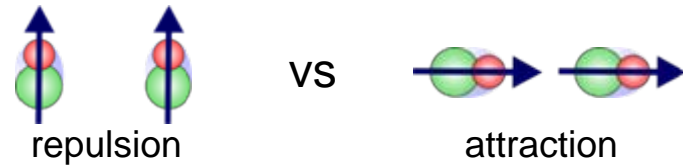
- interaction of two molecules



$$V_{dd} = \frac{\vec{d}_1 \cdot \vec{d}_2 - 3(\vec{d}_1 \cdot \vec{e}_b)(\vec{e}_b \cdot \vec{d}_2)}{r^3}$$

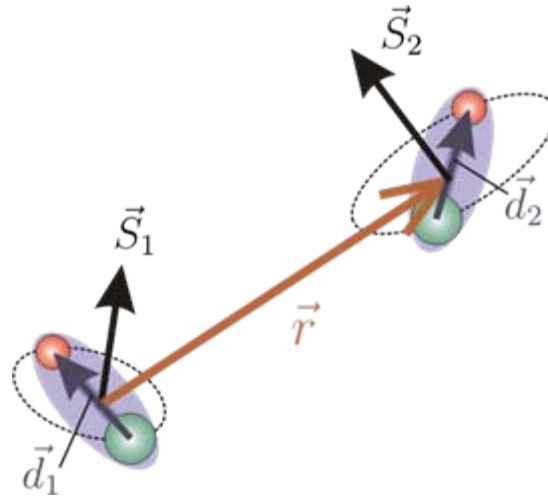
features of dipole-dipole interaction

- ✓ long range  $\sim 1/R^3$
- ✓ angular dependence

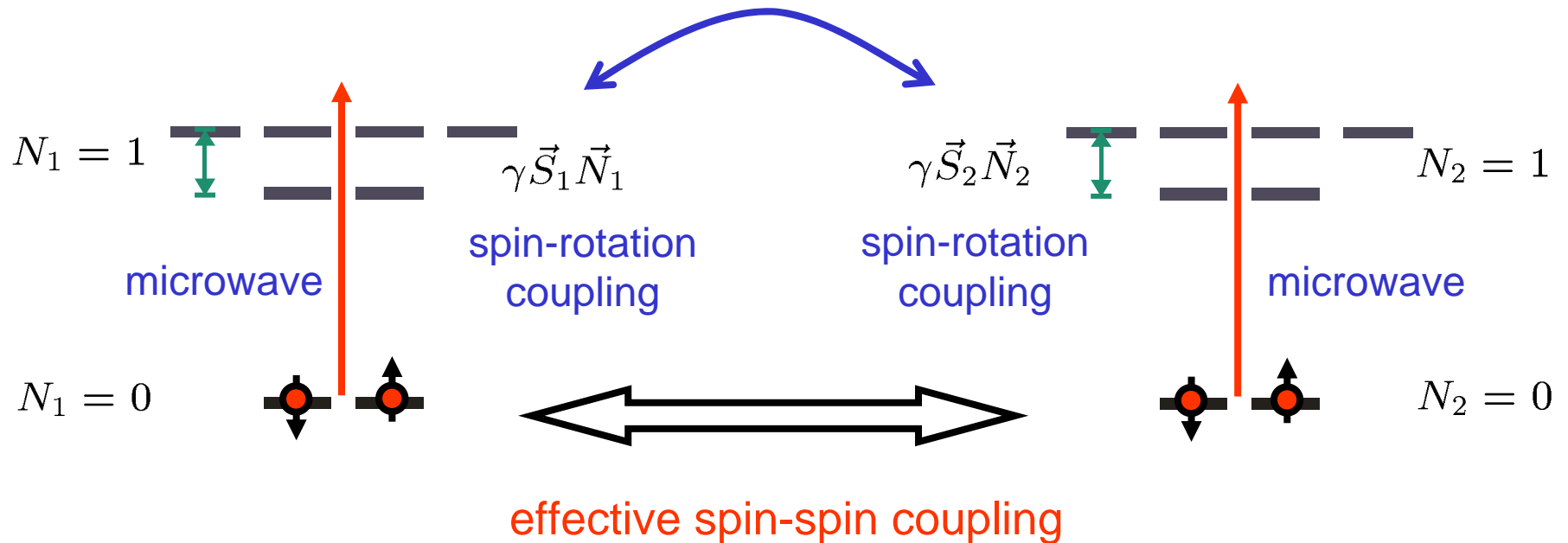


- ✓ strong! (temperature requirements)

... combining 1+2: engineering spin – spin interactions



dipole-dipole:  
anisotropic + long range



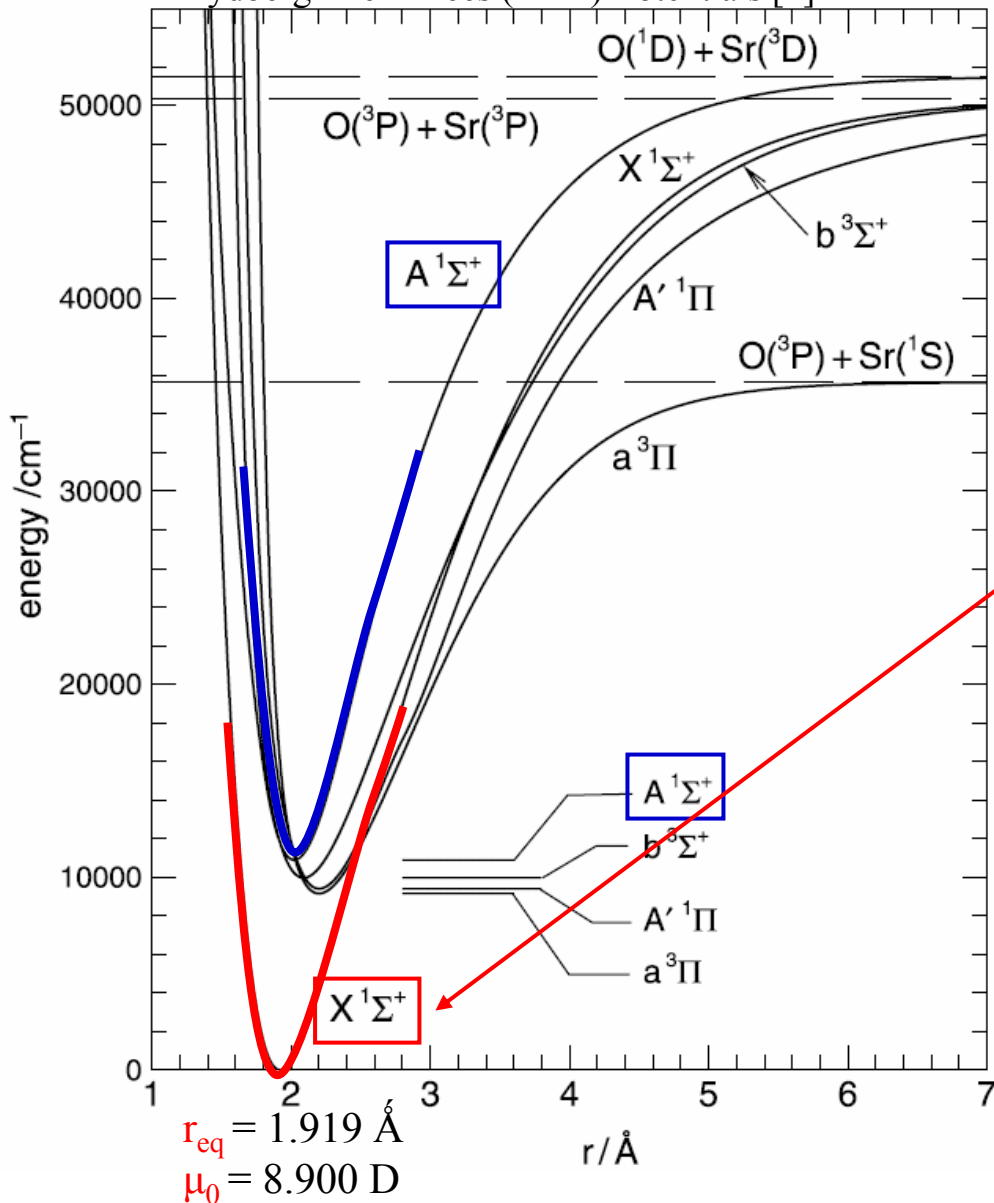


## Polar Molecules II (for the Brave)

- how real molecules look like
  - SrO
  - CaF

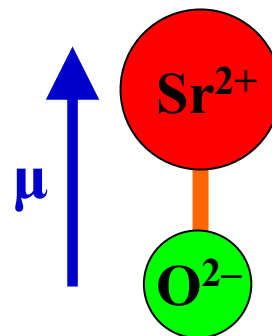
# Polar Molecules I : Strontium Oxide (SrO)

Rydberg-Klein-Rees (RKR)-Potentials [1]



heteronuclear molecule with strong **persistent** dipole moment in electronic groundstate.

**Sr<sup>2+</sup>O<sup>2-</sup>** ... ionic binding



<sup>38</sup> Sr	[Kr]5s <sup>2</sup>	
<sup>88</sup> Sr	I <sup>p</sup> = 0 <sup>+</sup>	(83%)
<sup>86</sup> Sr	I <sup>p</sup> = 0 <sup>+</sup>	(10%)
<sup>87</sup> Sr	I <sup>p</sup> = 3/2 <sup>+</sup>	(7%)
<sup>8</sup> O	1s2s <sup>2</sup> p <sup>4</sup>	
<sup>16</sup> O	I <sup>p</sup> = 0 <sup>+</sup>	(99.76%)
<sup>18</sup> O	I <sup>p</sup> = 0 <sup>+</sup>	(0.20%)

**X ¹Σ<sup>+</sup>** ... electronic groundstate:

S=0 ... closed shell ( $\cdot 9\sigma^2 10\sigma^2 4\pi^4$ )

$r_{\text{eq}} = 1.919 \text{ Å}$  ... equilibrium distance

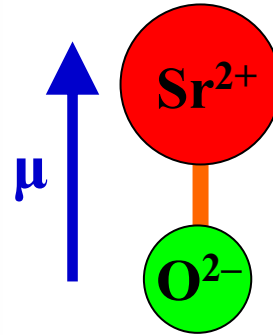
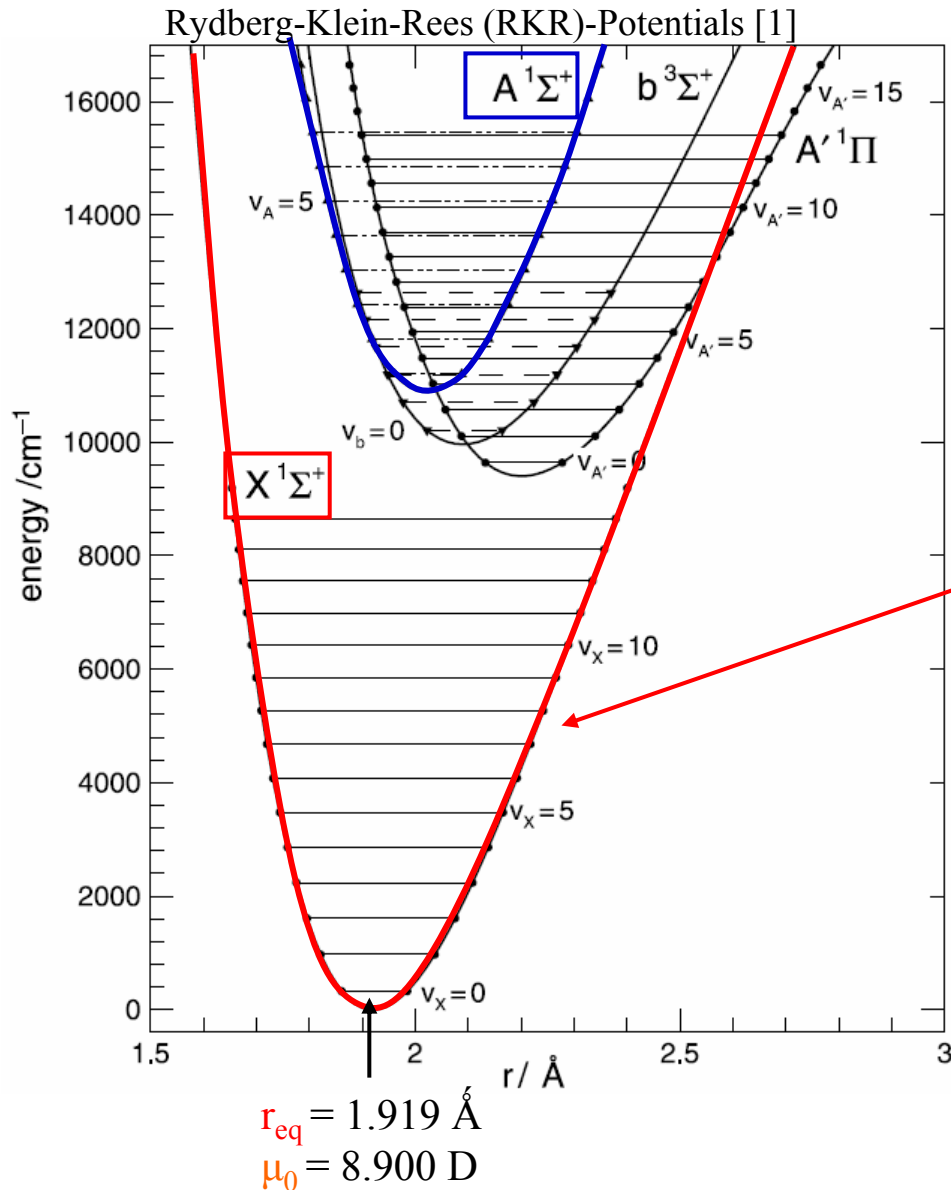
$\mu_0 = 8.900 \text{ D}$  ... dipole-moment

$\omega_{\text{eq}} = 19.586 \text{ THz}$  ... vibrational const.

$B_{\text{eq}} = 10.145 \text{ GHz}$  ... rotational ( $h/8\pi^2 M_{\text{red}} r_{\text{eq}}^2$ )

I=0 ... no nuclear momenta for <sup>88</sup>SrO, <sup>86</sup>SrO  
 [<sup>87</sup>SrO has I=3/2 and thereby nuclear quadrupole momenta ...  $eQq_n \approx -42\text{MHz}$ ]

# Polar Molecules I : Strontium Oxide (SrO)



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## X 1Σ<sup>+</sup> ... el. groundstate:

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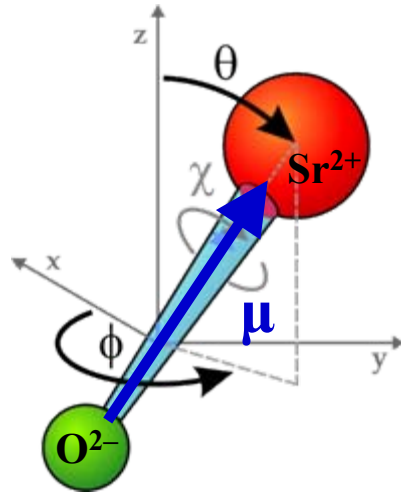
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# SrO as an "ideal" rigid Rotor

## $^{88}\text{SrO}$ in vibrational gs $X^1\Sigma^+$ ( $v=0$ ):

[ $v=1$  at  $\omega_{\text{eq}} = 19.586$  THz, no fine/hfs:  $S=I=0$ ]

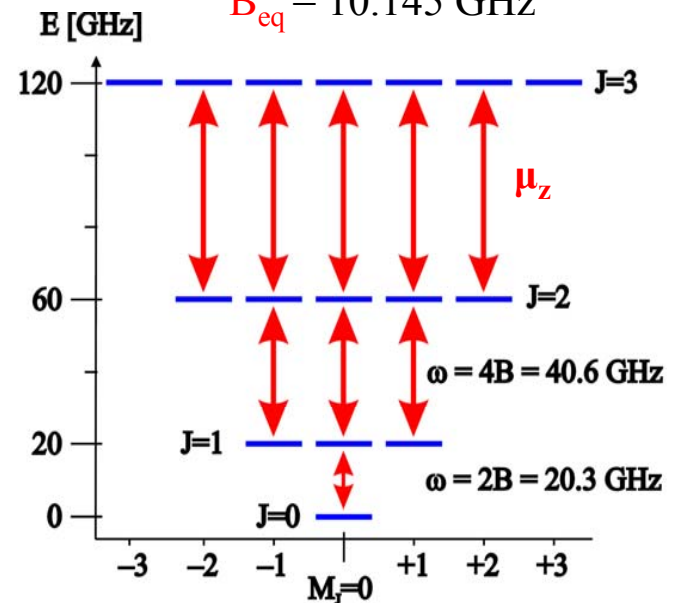


### • Rigid Rotor

$$\hat{H} = \frac{\hat{\mathbf{J}}^2}{2I}$$

$$E_{J,M_J} = BJ(J+1)$$

$$B_{\text{eq}} = 10.145 \text{ GHz}$$



- strong anharmonicity  $\sim J^2$
- unpolarized
- electric dipole transitions  
 $\Delta J = 1$  and  $\Delta M_J = 0, 1$

# CaF - rotational, fine and hyperfine structure

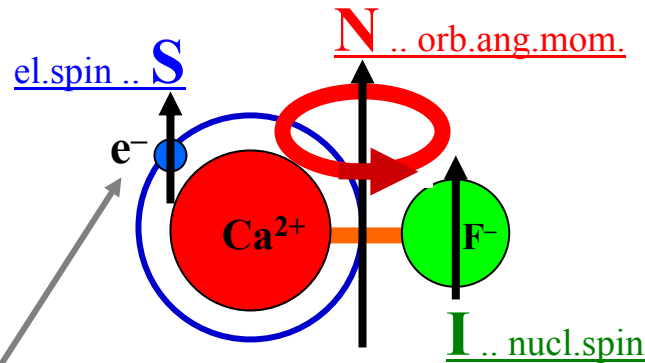
$X^2\Sigma_{1/2}$  ... el. groundstate:

$S=1/2$  ... from open (Ca-)shell

$I=1/2$  ... nuclear momentum (of F)

$r_{eq} = 1.951 \text{ \AA}$  ... eq. distance

$\mu_0 = 3.077 \text{ D}$  ... dipole-moment



optical excitation  
Alkali-like

talks to optical radiation,  
like an alkali atom

talks to microwave radiation  
(alkali atom: hyperfine / magnetic)

# CaF - rotational, fine and hyperfine structure

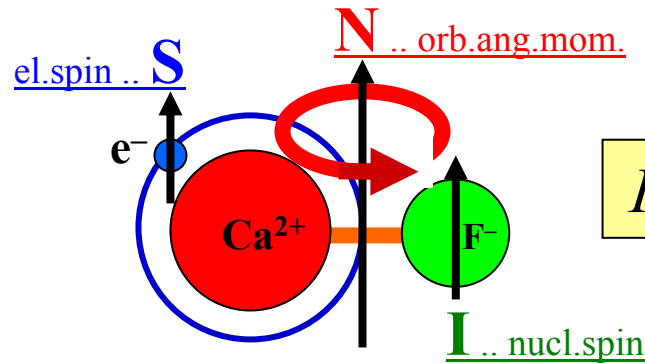
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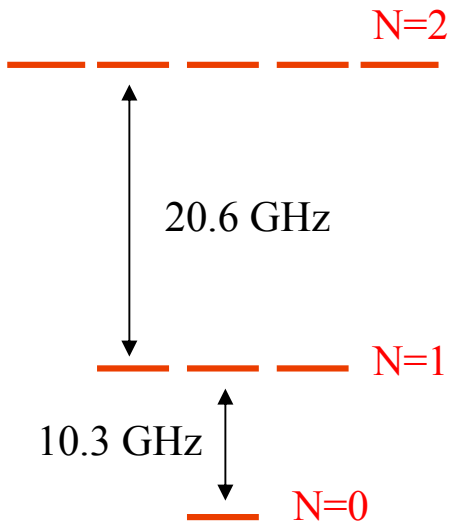
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$$H_0 = H_{rot} + H_{spin-rot} + H_{hfs}$$

$$H_{rot} = B N^2$$

$$B = 10.3 \text{ GHz}$$



( rotor )

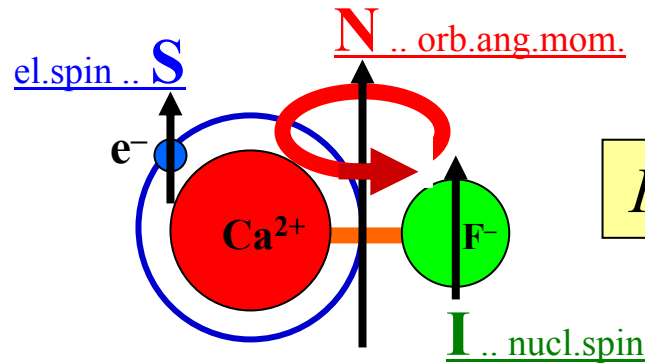
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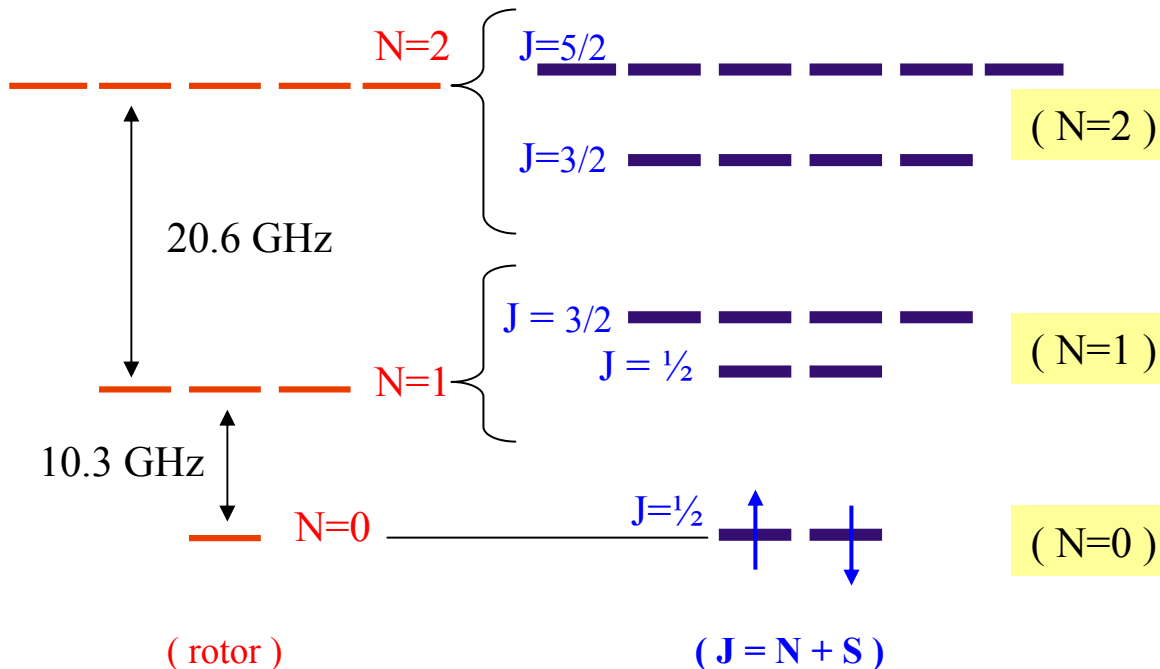
$$H_0 = H_{rot} + H_{spin-rot} + H_{hfs}$$

$$H_{rot} = B N^2$$

$$B = 10.3 \text{ GHz}$$

$$H_{spin-rot} = \gamma N \cdot S$$

$$\gamma = 39.66 \text{ MHz}$$

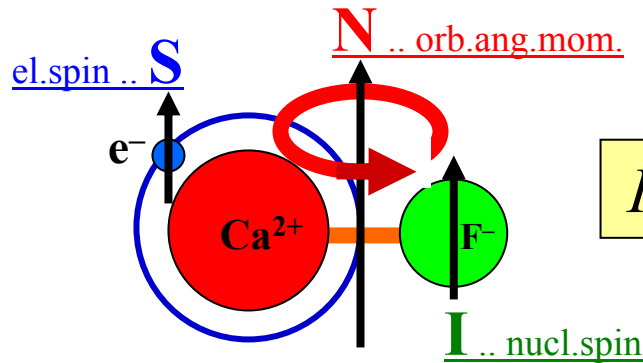


# CaF - rotational, fine and hyperfine structure

$X^2\Sigma_{1/2}$  ... el. groundstate:

$S=1/2$  ... from open (Ca-)shell  
 $I=1/2$  ... nuclear momentum (of F)

$r_{eq} = 1.951 \text{ \AA}$  ... eq. distance  
 $\mu_0 = 3.077 \text{ D}$  ... dipole-moment



$$H_0 = H_{rot} + H_{spin-rot} + H_{hfs}$$

...  $S_\zeta$  component along figure axis

$$H_{rot} = B N^2$$

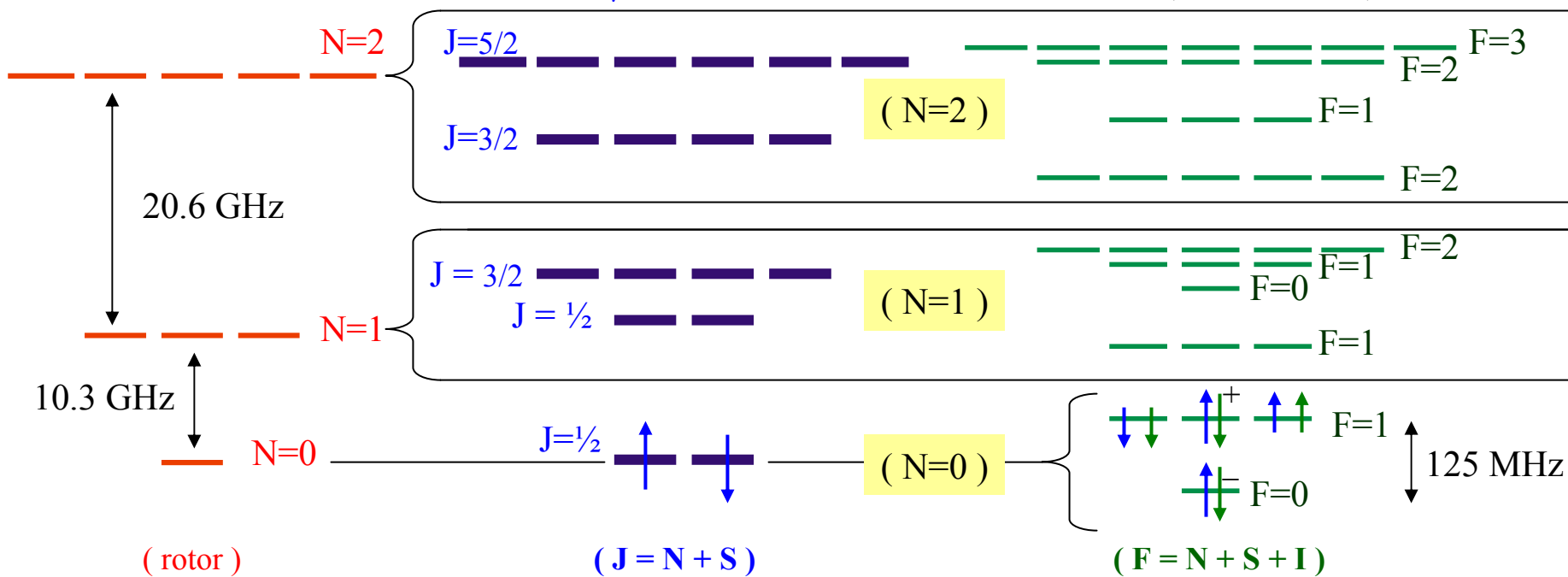
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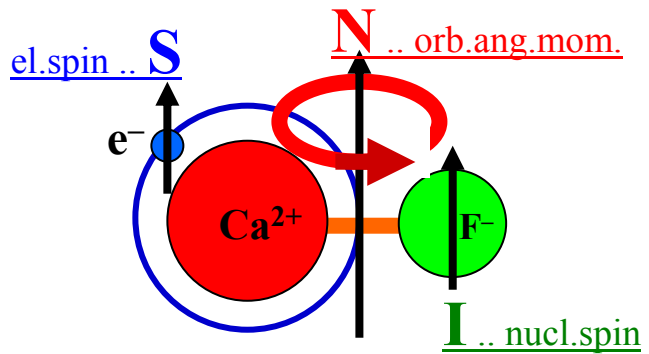
$$H_{hfs} = (b S \cdot I + c S_\zeta I_\zeta) + C N \cdot I$$

$b=109.18 \text{ MHz}, c=40.12 \text{ MHz}, C=30 \text{ KHz}$

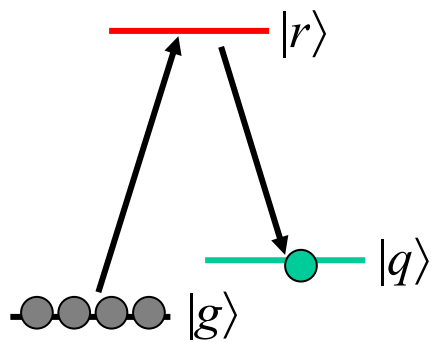




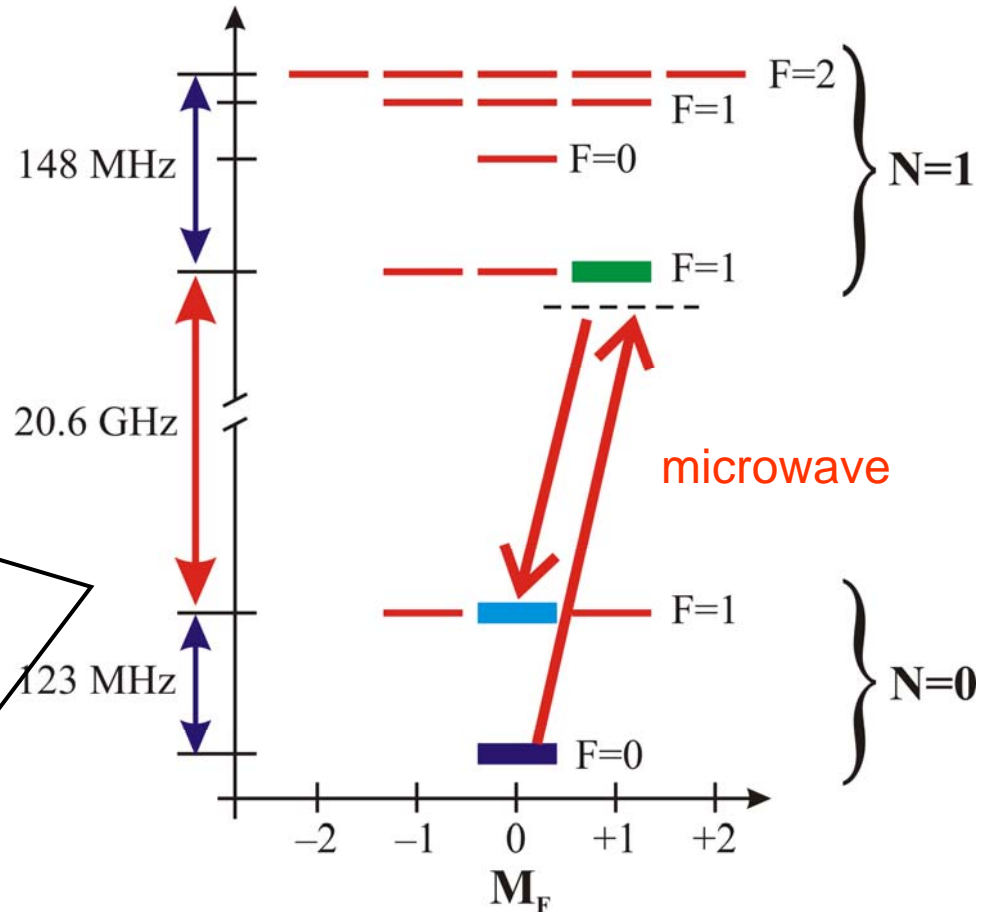
# Calcium Monofluoride (CaF): in external fields III



Summary for Quantum Optics Theorists



two & three level systems  
spontaneous emission from excited states negligible



clock transition

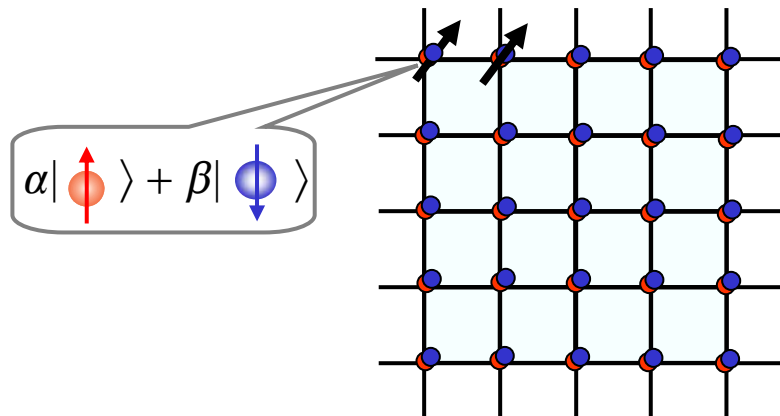
Sneak Preview:

What can we do with Polar Molecules?

- a few examples & ideas

# 1. Polar Molecules in Optical Lattices: “Spin Lattice Toolbox“

- Mott insulator state of molecules in optical lattices



engineer spin models:

$$H = g \sum_{i \neq j} \sigma_{\alpha}^{(i)} A^{\alpha\beta}(\vec{x}_i, \vec{x}_j) \sigma_{\beta}^{(j)}$$

- ✓ range
- ✓ spin / space anisotropy

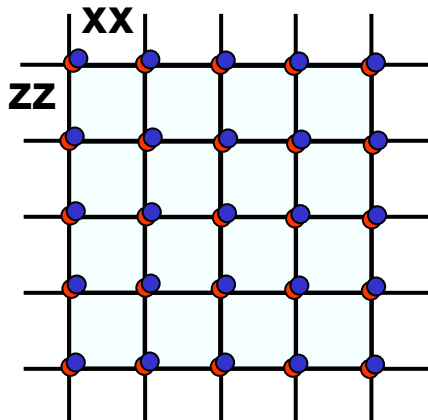
dipoles: large offsite couplings!

- [... and Hubbard models with large offsite interactions]

# Examples ...

- Protected Quantum Memory

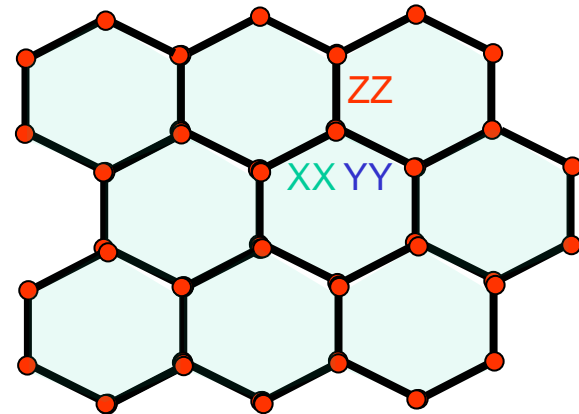
B.Douçot, M.V.Feigel'man, L.B.Ioffe,  
A.S.Ioselevich, Phys. Rev. B **71**, 024505 (2005).



$$H_{\text{spin}}^{(I)} = \sum_{i=1}^{\ell-1} \sum_{j=1}^{\ell-1} J(\sigma_{i,j}^z \sigma_{i,j+1}^z + \cos \zeta \sigma_{i,j}^x \sigma_{i+1,j}^x)$$

- Kitaev 2005

E.Dennis, A.Y.Kitaev, J.Preskill, JMP **43**, 4452;  
A.Y.Kitaev, cond-mat/0506438 (2005);

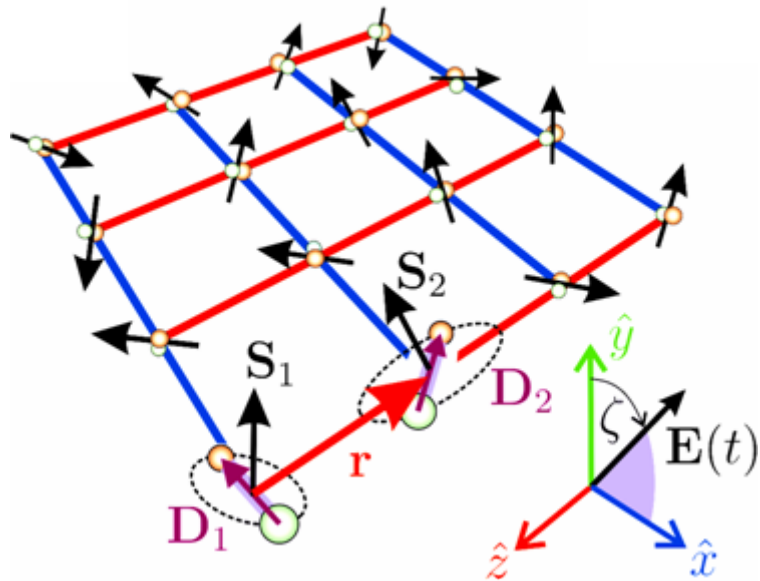


$$H_{\text{spin}}^{(II)} = J_{\perp} \sum_{x\text{-links}} \sigma_j^x \sigma_k^x + J_{\perp} \sum_{y\text{-links}} \sigma_j^y \sigma_k^y + J_z \sum_{z\text{-links}} \sigma_j^z \sigma_k^z$$

# Examples ...

- Protected Quantum Memory

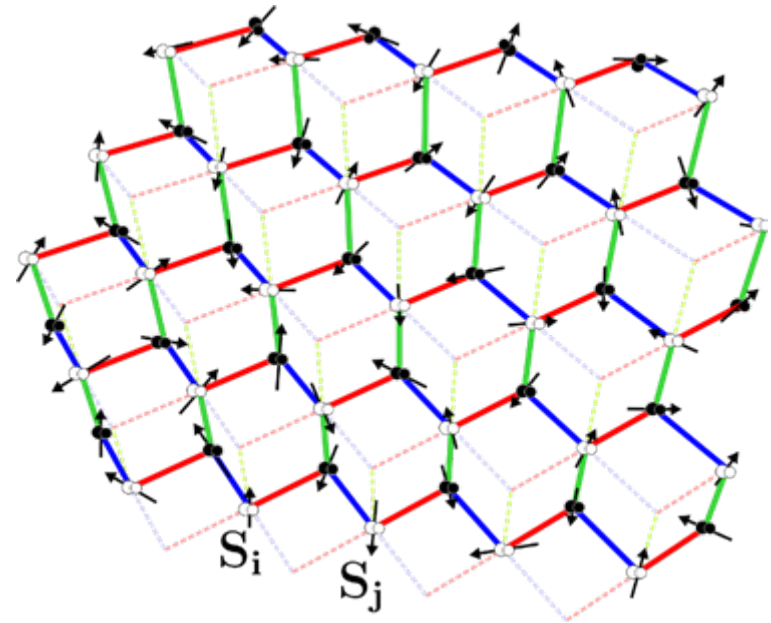
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- Kitaev 2005

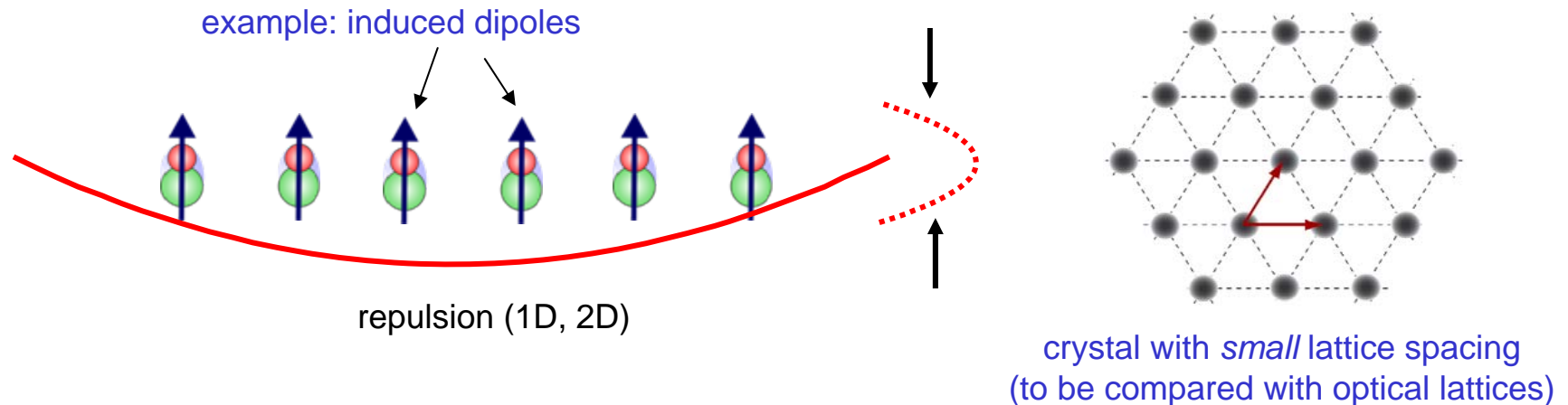
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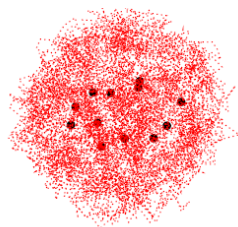
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## 2. Self-assembled Lattices / Crystals with Atoms / Molecules

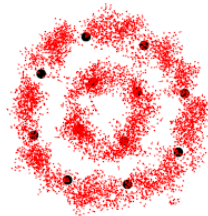
- engineer *repulsive* interactions between atoms / molecules so that ...



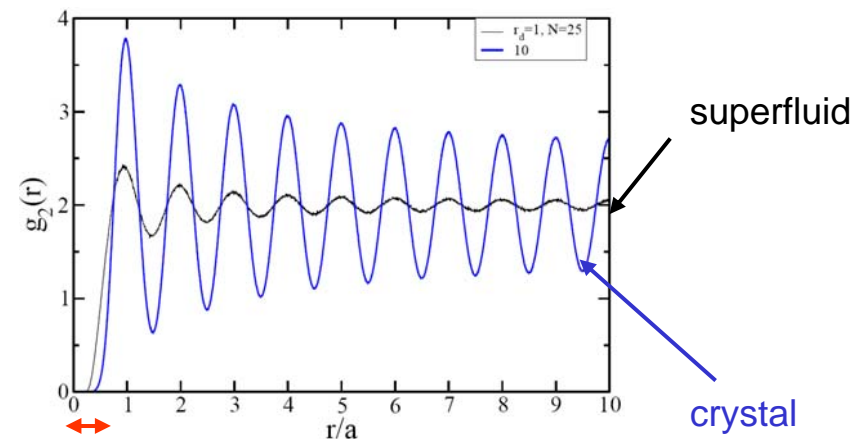
- Liquid – Crystal phase transitions:



superfluid

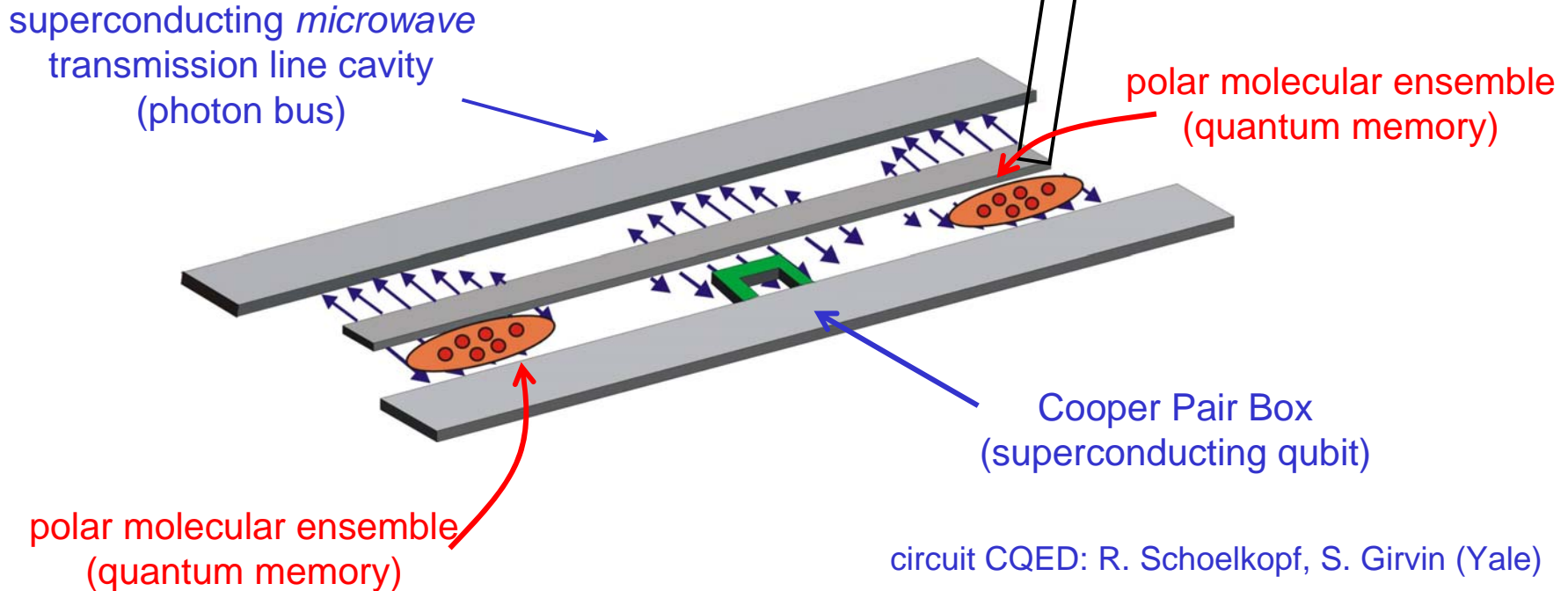


crystal: shell structure  
(compare: ions)



### 3. Solid State / AMO Hybrid Quantum Processor

- solid state processor
- AMO quantum memory here: **molecular ensembles**
- interface to ‚flying‘ optical qubits



P.Rabl, D. DeMille, J. Doyle, M. Lukin, R. Schoelkopf and P. Zoller, preprint

- or: couple single molecules Yale – Harvard – Innsbruck collaboration, draft

## Spin Models with Optical Lattices

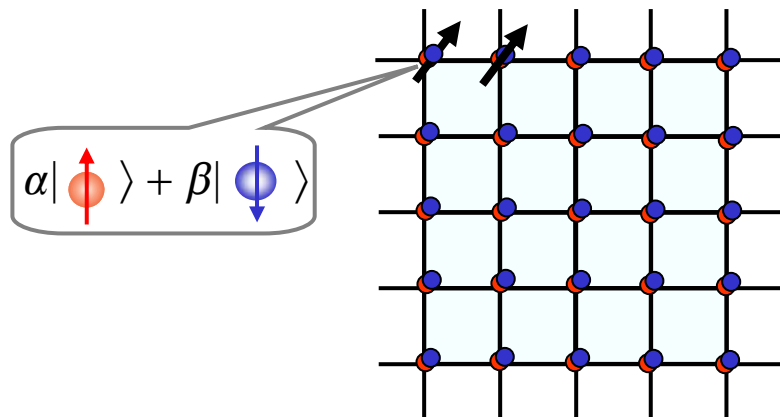
- quantum info relevance:
  - polar molecule realization of models for protected quantum memory (Ioffe, Feiguin et al.)
  - Kitaev model on honeycomb lattice

A. Micheli, G.K. Brennen, and P. Zoller, Nature Physics May 2006



# Polar Molecules in Optical Lattices: “Spin Lattice Toolbox“

- Mott insulator state of molecules in optical lattices



engineer spin models:

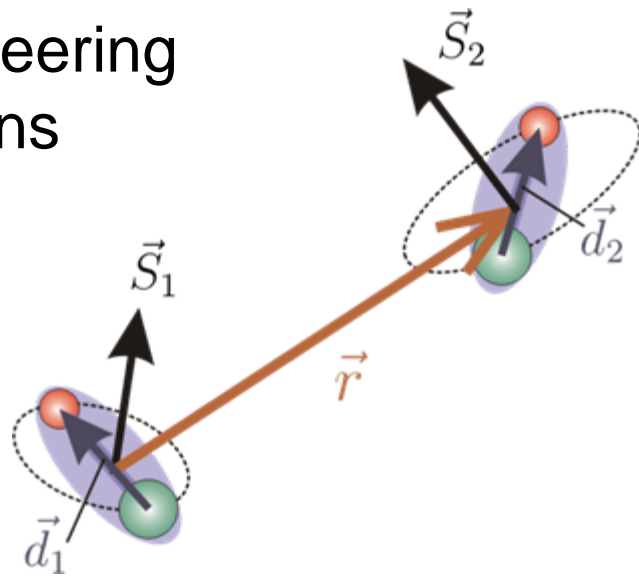
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- ✓ range
- ✓ spin / space anisotropy

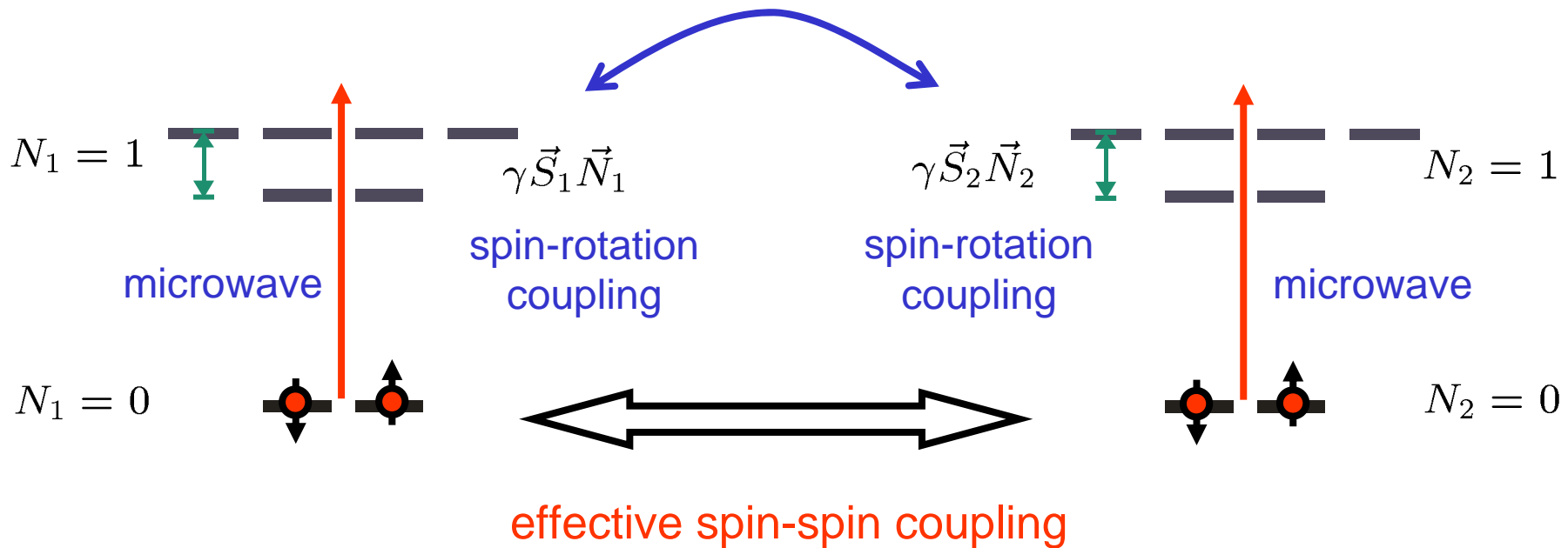
dipoles: large offsite couplings!

- [... and Hubbard models with large offsite interactions]

# Basic idea of engineering spin-spin interactions

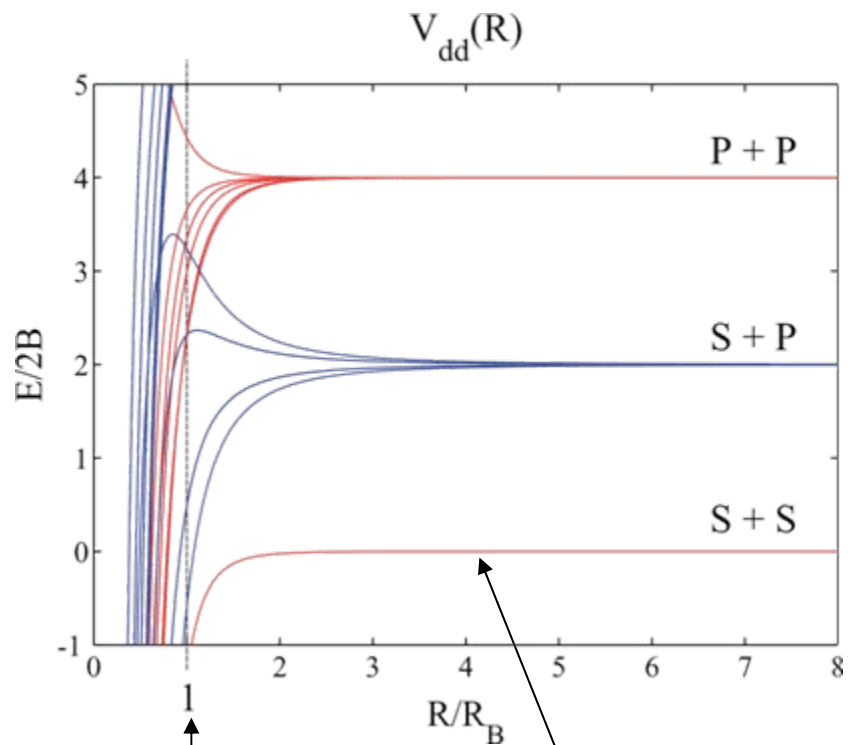
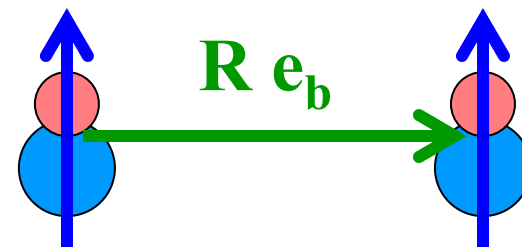


dipole-dipole:  
anisotropic + long range



# Adiabatic potentials for two (unpolarized) polar molecules

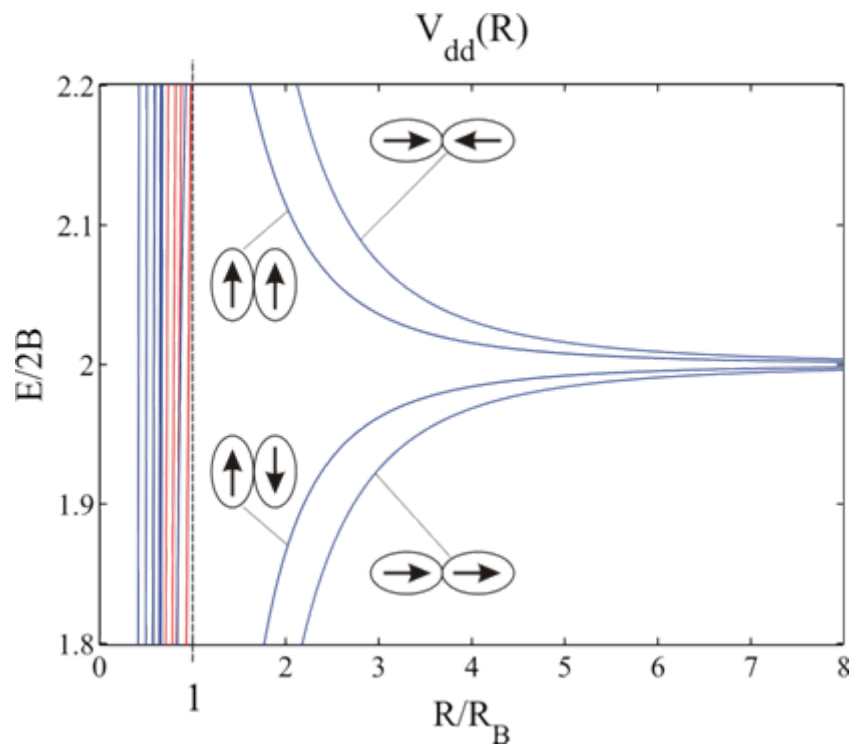
- Rotor**



$$R_B = \left( \frac{\mu^2}{3B} \right)^{1/3}$$

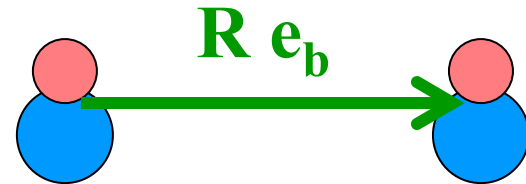
~ 30-60 nm

$$V_{\text{eff}}(R) = -\frac{C_6}{R^6}$$

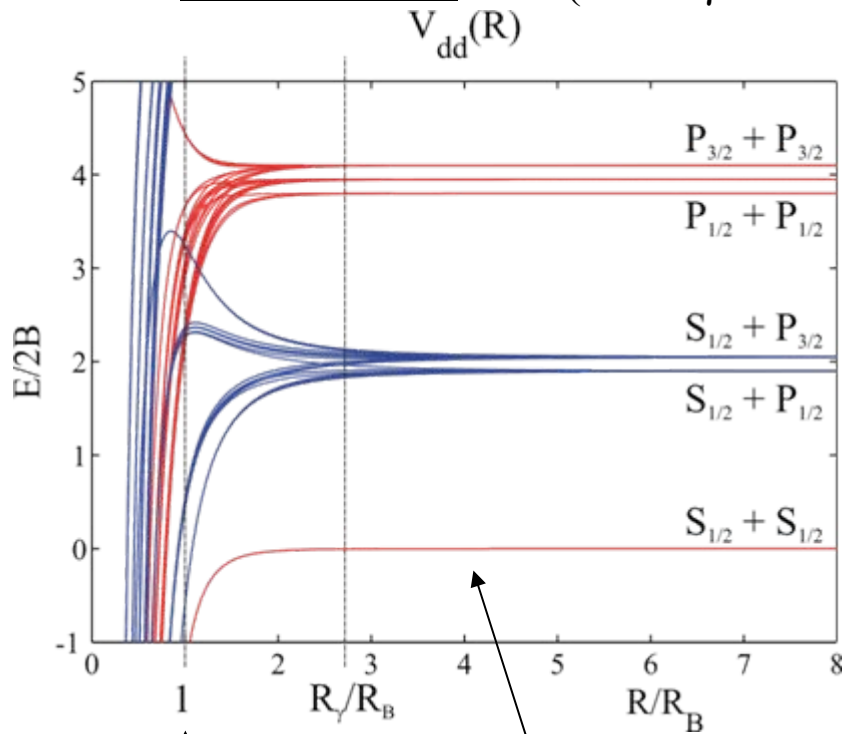


$$C_6 = \frac{\mu^4}{6B}$$

# Adiabatic potentials for two (unpolarized) polar molecules

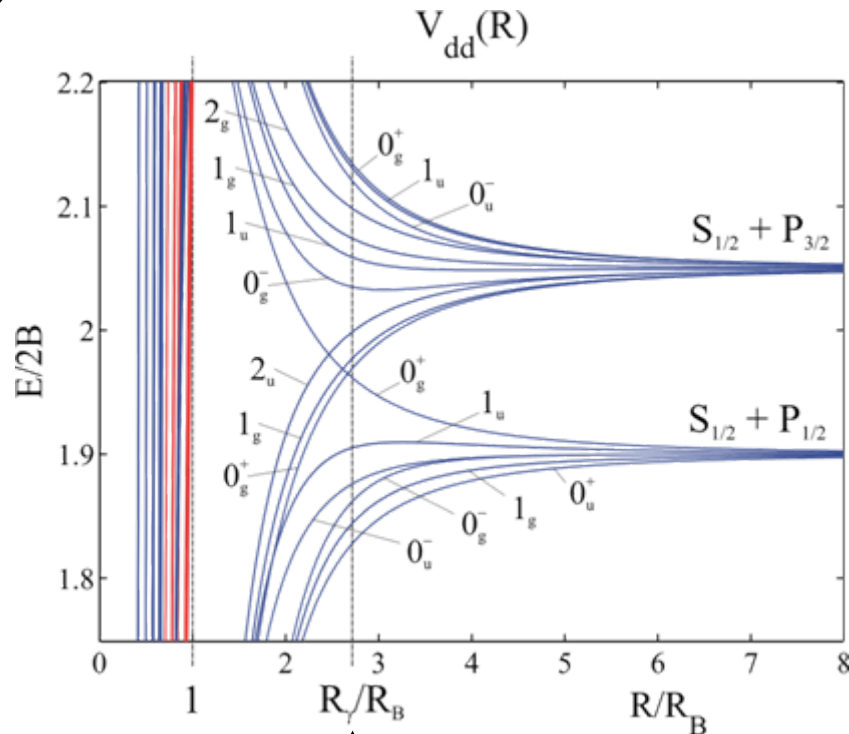


- **Spin Rotation** ( here:  $\gamma/B = 1/10$  )



$$R_B = \left( \frac{\mu^2}{3B} \right)^{1/3}$$

$$C_6 = \frac{\mu^4}{6B}$$

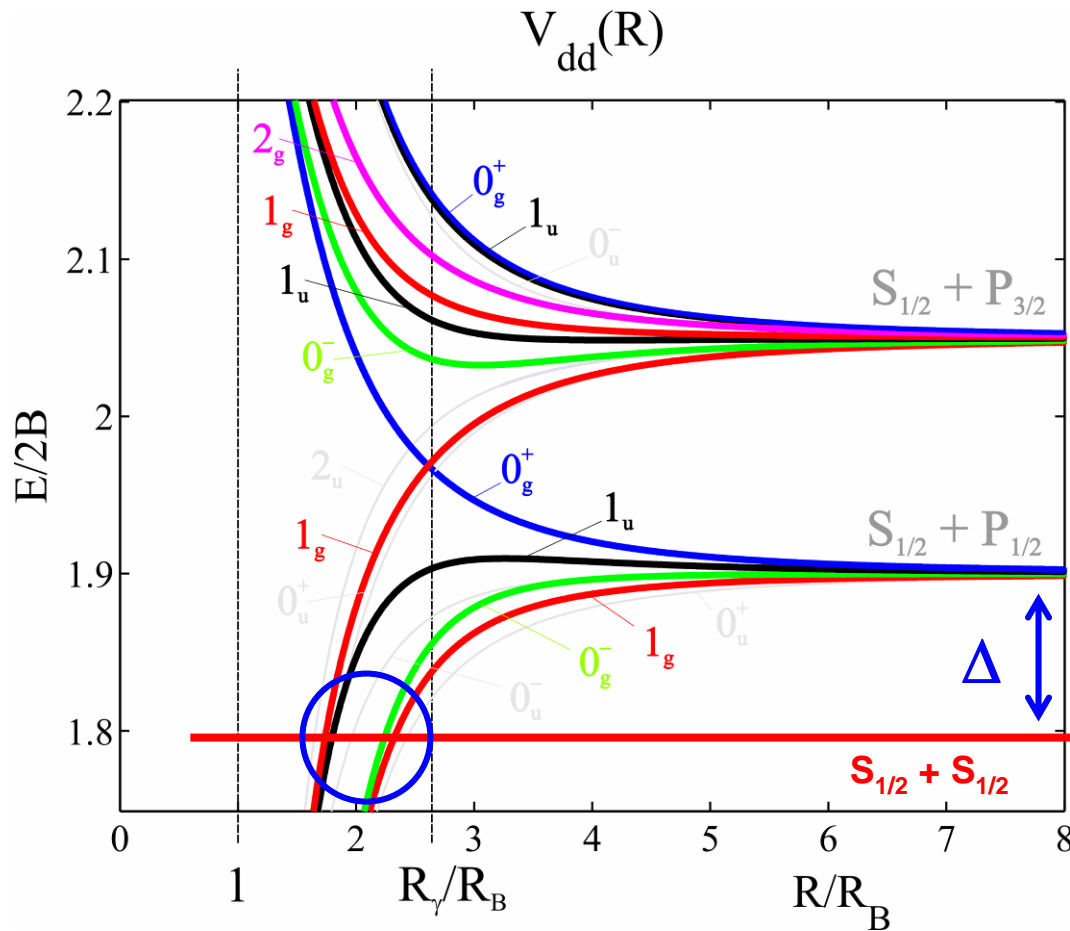


$$R_\gamma = \left( \frac{2\mu^2}{3\gamma} \right)^{1/3} \text{ typical: } \sim 300 \text{ nm}$$

$$V_{\text{eff}}(R) = -\frac{C_6}{R^6} \left[ 1 + \left( \frac{\gamma}{2B} \right)^2 (S_1 \cdot S_2 - S_1^b S_2^b) + O\left( \frac{\gamma}{2B} \right)^4 \right] \sim \text{small}$$

# Adiabatic potentials for two (unpolarized) polar molecules

- **Spin Rotation** ( here:  $\gamma/B = 1/10$  )



Induced effective interactions:

$$0_g^+ : + \mathbf{S}_1 \cdot \mathbf{S}_2 - 2 S_1^c S_2^c$$

$$0_g^- : + \mathbf{S}_1 \cdot \mathbf{S}_2 - 2 S_1^p S_2^p$$

$$1_g : + \mathbf{S}_1 \cdot \mathbf{S}_2 - 2 S_1^b S_2^b$$

$$1_u : - \mathbf{S}_1 \cdot \mathbf{S}_2$$

$$2_g : + S_1^b S_2^b$$

$$0_u : 0$$

$$2_u : 0$$

for  $\mathbf{e}_{\text{body}} = \mathbf{e}_x$  and  $\mathbf{e}_{\text{pol}} = \mathbf{e}_z$

$$0_g^+ : +XX - YY + ZZ$$

$$0_g^- : +XX + YY - ZZ$$

$$1_g : -XX + YY + ZZ$$

$$1_u : -XX - YY - ZZ$$

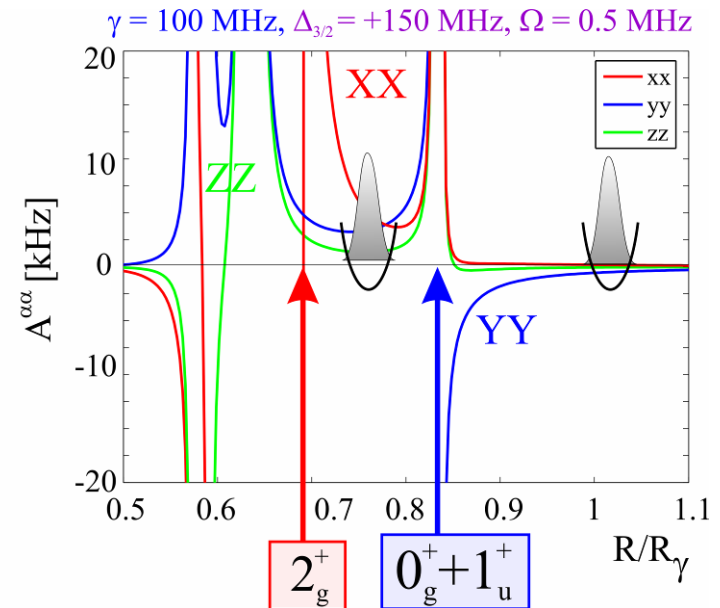
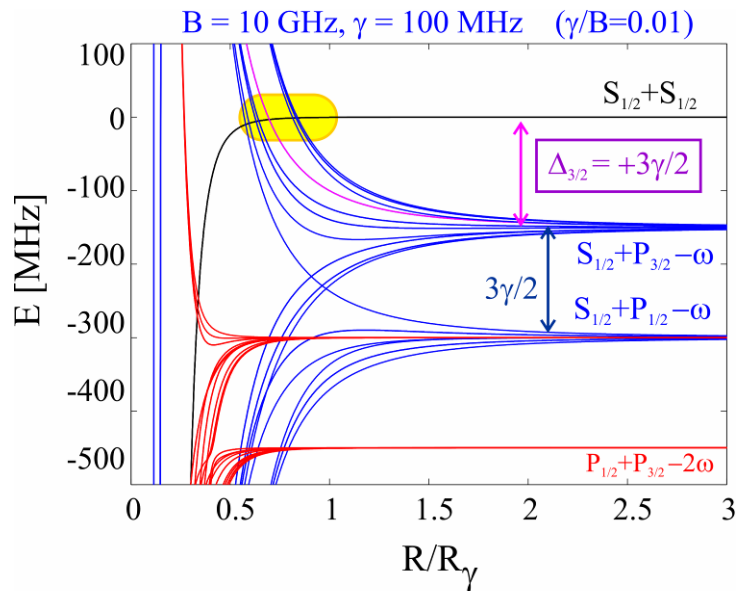
$$2_g : +XX$$

Feature 1. By tuning close to a resonance we can select a *specific spin texture*

# Example: "The Douçot-Feigel'man-Ioffe et al. Model"

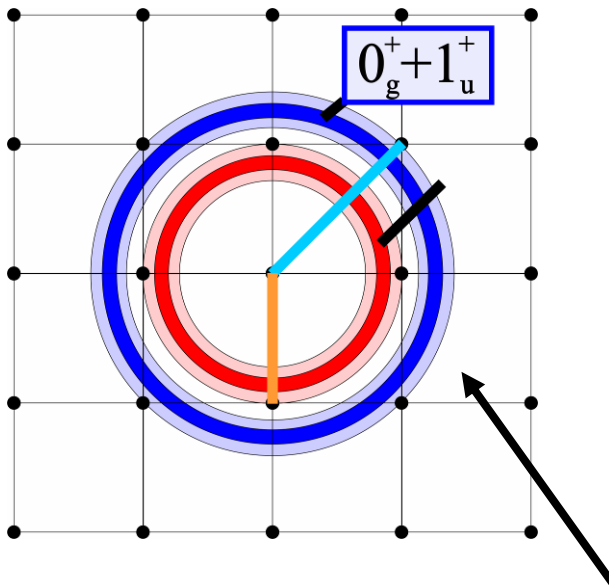
$$H = J_x \sum_{ij} \sigma_x^{(ij)} \sigma_x^{(i,j+1)} + J_z \sum_{ij} \sigma_z^{(ij)} \sigma_z^{(i+1,j)}$$

- Consider blue detuning from  $S+P_{3/2}$  transitions:



## Example: "The Douçot-Feigel'man-Ioffe et al. Model"

- Model is simple in terms of long-range resonances ...

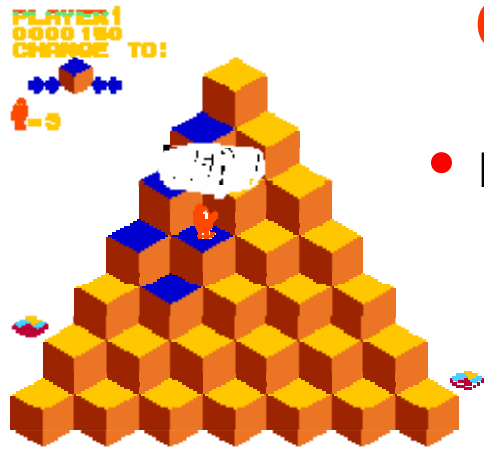


Rem.: for a multifrequency field we can *add* the corresponding spin textures.

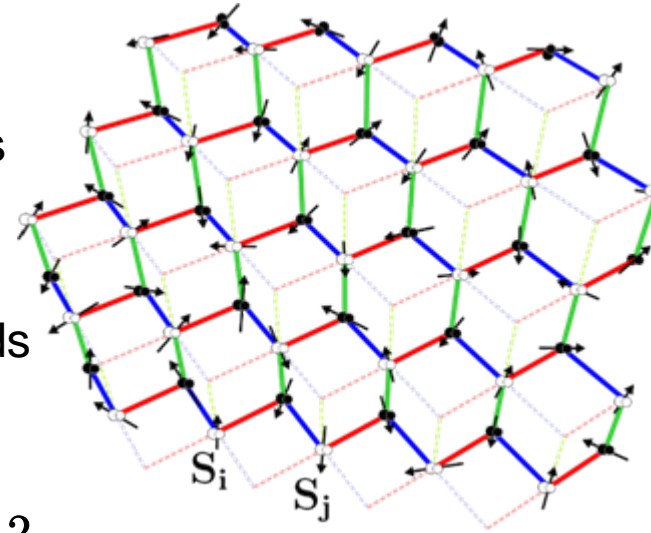
Feature 2. We can choose the *range* of the interaction for a given spin texture

Feature 3. for a *multifrequency* field spin textures are *additive*: toolbox

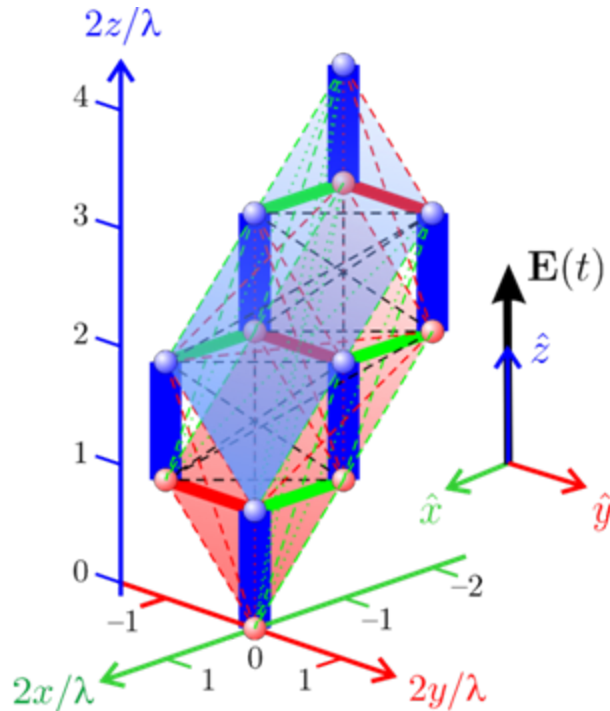
# Construction in an optical lattice



- Implementation in *Q\*bert lattice*:
  - Two staggered triangular lattices
  - Nearest neighbors give honeycombs
  - their edges form orthogonal triads



- Realization with 3 fields: (several possible choices) shown when all 3 being z polarized, resp. near  $0_g, 1_g, 2_g$



Spin pattern	Residual long range coupling strengths $ J_{lr} $
$\sigma^z \sigma^z$	
$\sigma^x \sigma^x$	
$\sigma^y \sigma^y$	$< 10^{-2}  J_z $
Other	$< 10^{-3}  J_z $
$ J_{\perp}  = 0.4  J_z $	

**Operator fidelity (on a 4 spin configuration)**

$$\sup [ \| |H_{\text{spin}} - H_{\text{spin}}^{(\text{II})} | \psi \rangle \|_2; \langle \psi | \psi \rangle = 1 ] = 10^{-4} |J_z|$$

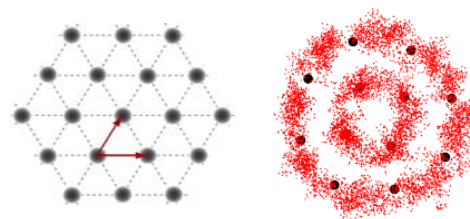
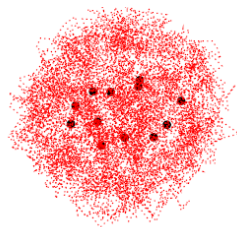
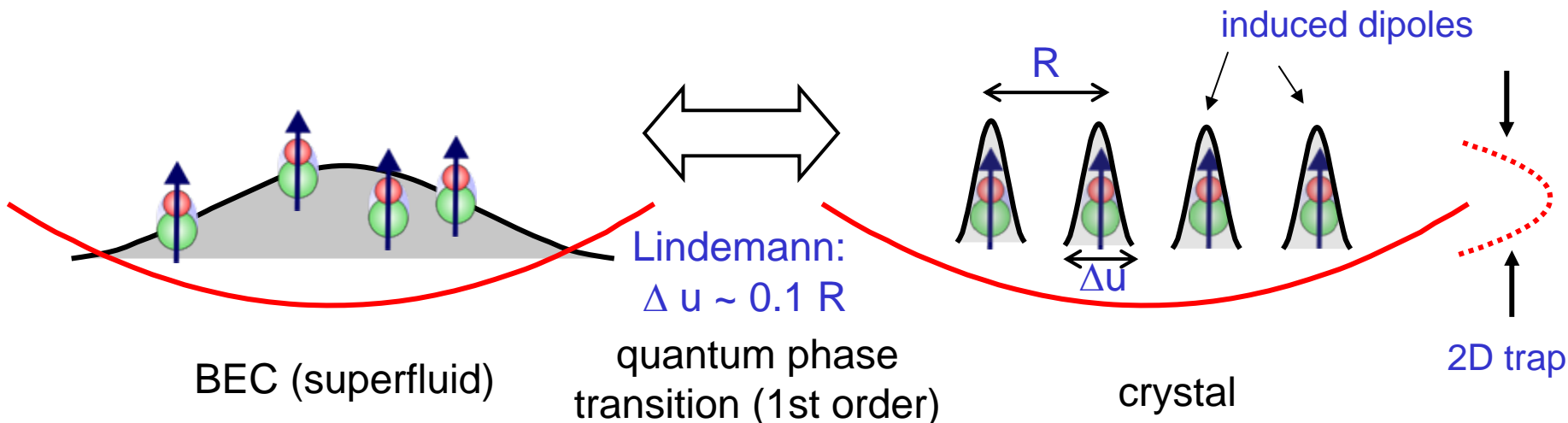


# Crystals / self-assembled Lattices with Polar Molecules (and Atoms)

G. Pupillo, A. Micheli, V. Steixner, H.P. Büchler, and P. Zoller (Innbruck)  
N.V. Prokof'ev (Umass, Amherst) & Lukin group (Harvard)

# Crystalline Phases with Molecules (and Atoms)

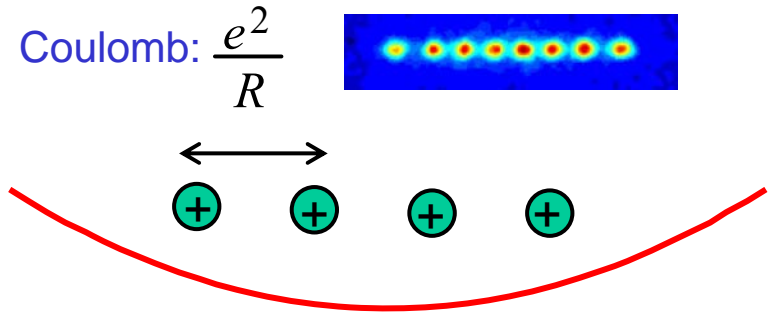
- requirement: engineering *long range repulsive* interactions, e.g.  $V \sim 1/r^\alpha$ 
  - e.g. (induced & aligned) electric dipole moments of molecules (or atoms) in 1D or 2D:  $V \sim \mu^2 / r^3$ ; or repulsive  $+C_6/r^6$
- superfluid – crystal quantum phase transition (at  $T=0$ )



self assembled lattice / shell structure

- how build, prepare, measure? properties? why? ... from an AMO point of view
- Ref.: many papers on dipolar gases, and *non-AMO* crystals (e.g. Lozovik et al.)

## Trapped Ions



- repulsive long range interaction

Coulomb / 3D

- crystal forms for ...

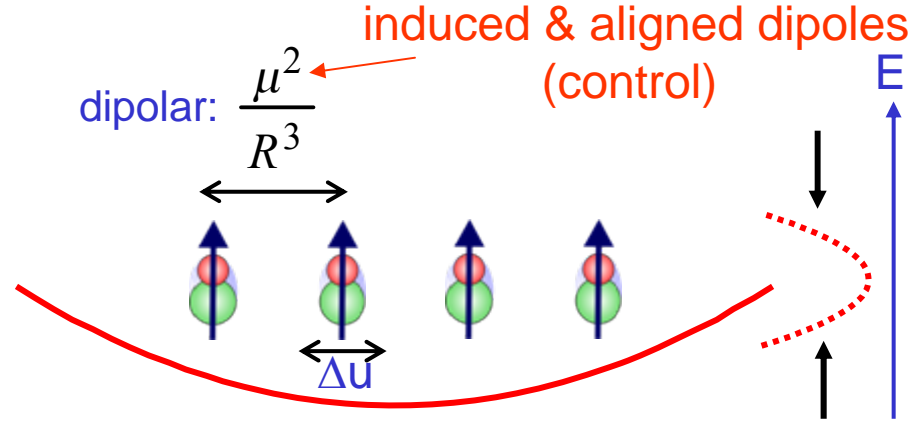
$$r_c = \frac{E_{\text{pot}}}{E_{\text{kin}}} \equiv \frac{e^2/R}{\hbar^2/mR^2} \sim R \gg 1$$

Wigner crystal @ low density

- melting (?)

~~BEC~~

## Trapped Polar Molecules



repulsive in 1D & 2D trap

stability!?

$$r_d = \frac{\mu^2/R^3}{\hbar^2/mR^2} \sim \frac{1}{R} \gg 1$$

crystal @ high density  
(high density ensembles / lattices)

$$\Delta u \sim 0.1 \dots 0.2R$$

Lindemann: dipolar BEC

# Hamiltonian

- Hamiltonian in 2D

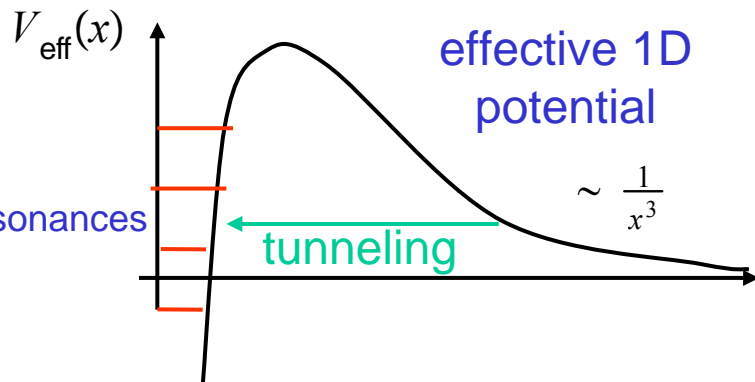
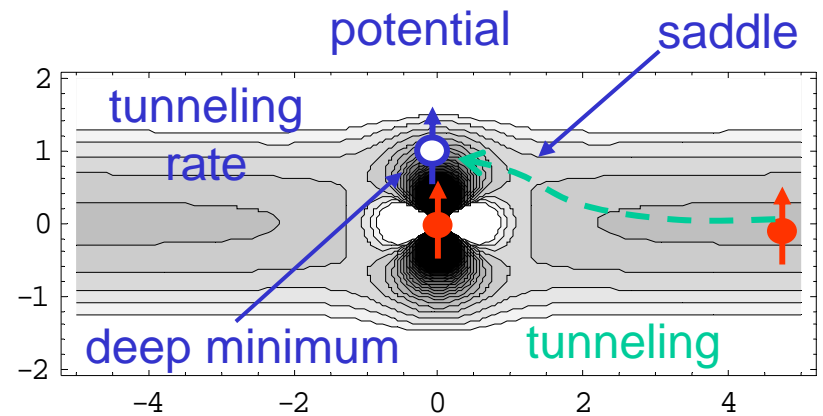
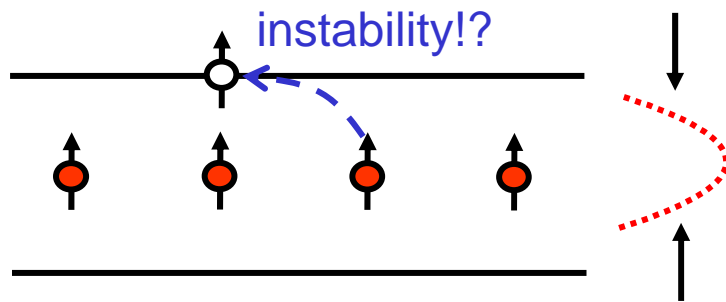
$$H_{2D} = \sum_i \frac{\vec{p}_i^2}{2m} + V_T(\vec{r}_i) + \sum_{ij} V_d(\vec{r}_i - \vec{r}_j) + U_s(\vec{r}_i - \vec{r}_j)$$

“In-plane” confinement  
(quadratic trap or box)

Dipole-dipole  
int.:  $\sim \mu^2/r^3$

Short-range

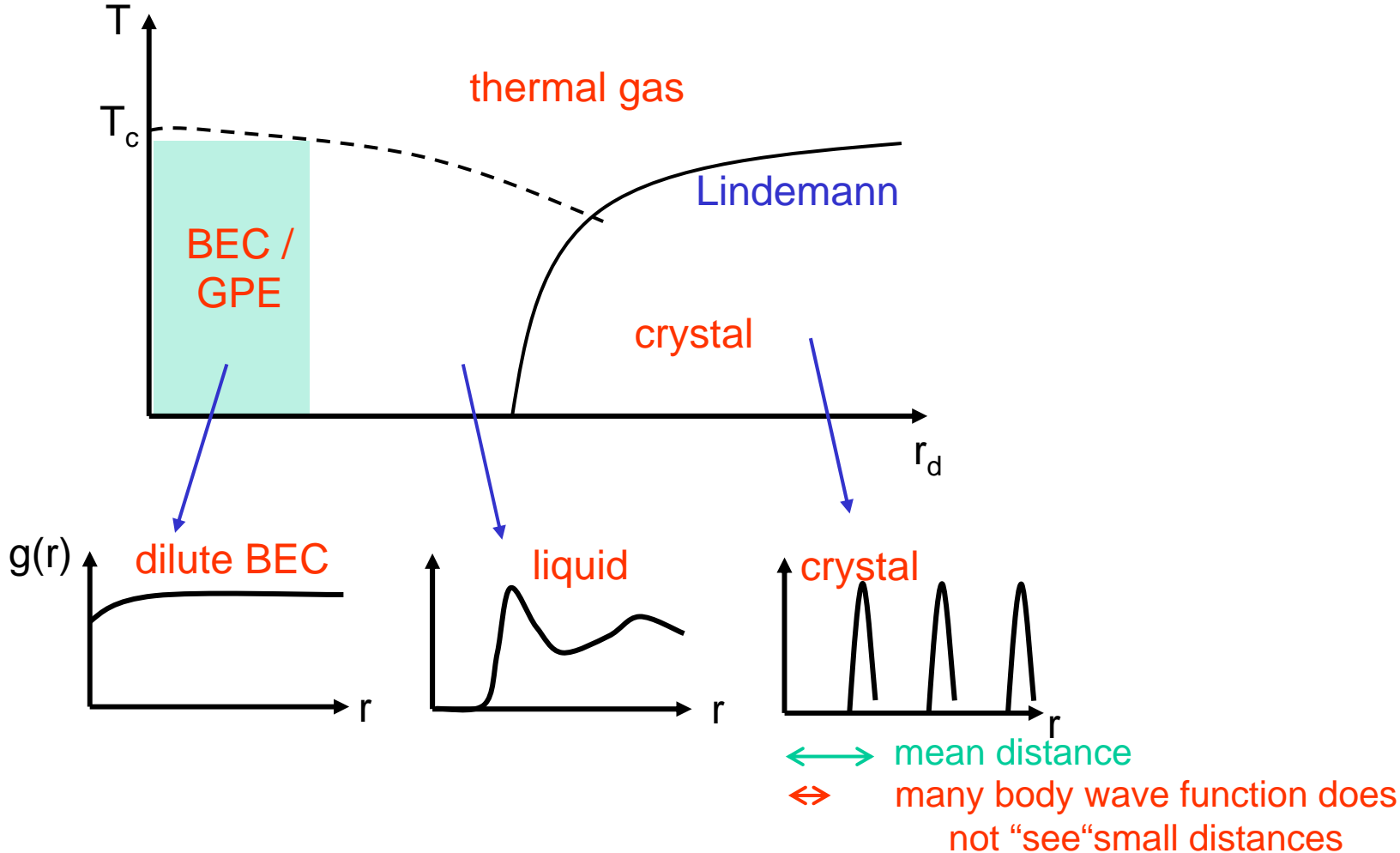
- tunneling / instability



$$\Gamma = \omega_p e^{-\left(\frac{r_d}{a_{\perp}}\right)^{2/5}}$$

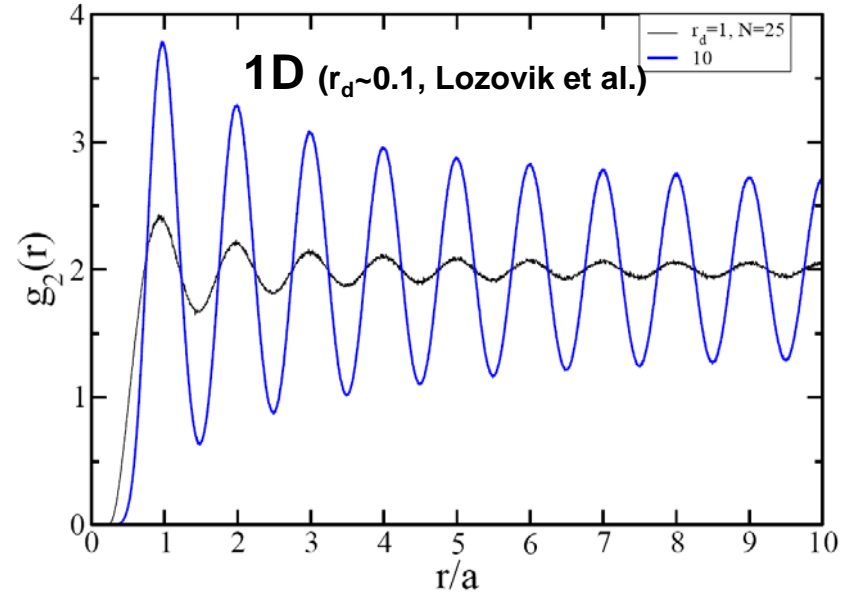
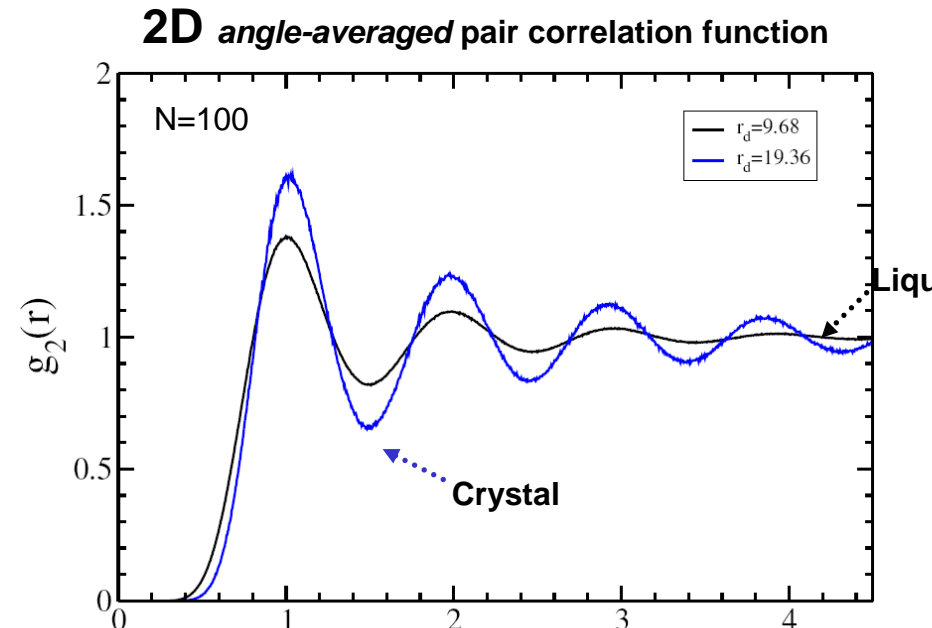
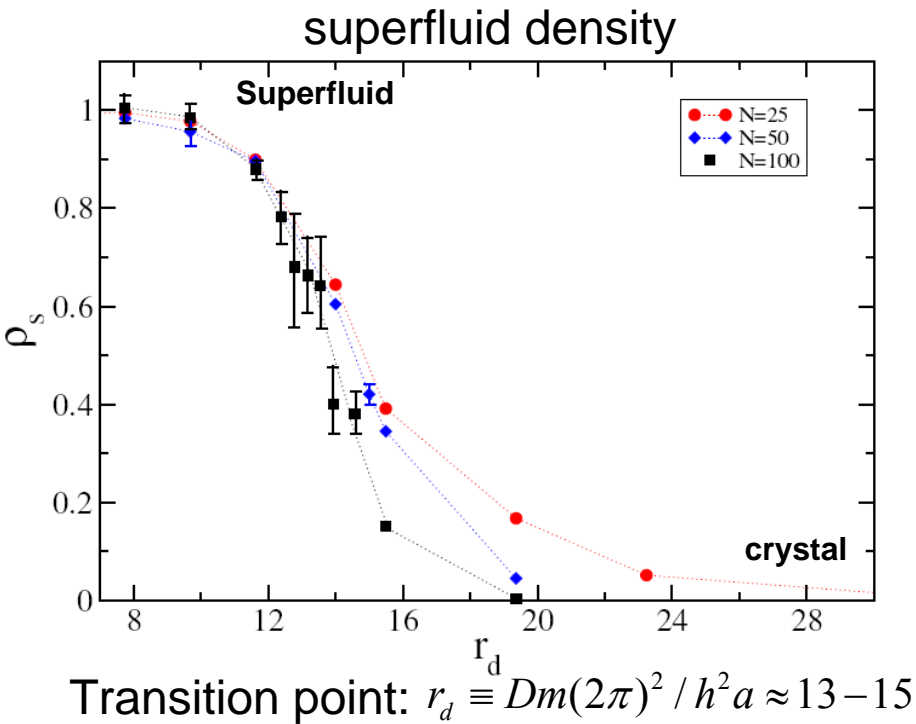
$\omega = (\mu^2/mR^5)^{1/2}$  (oscillation frequency)  
 $r_d > 15$   
 $a/a_{\perp} > 10$

# Sketch of Phase Diagram



# Homogeneous system in 2D

Methods: Boninsegni, Prokof'ev and Svistunov, PRL (2006)

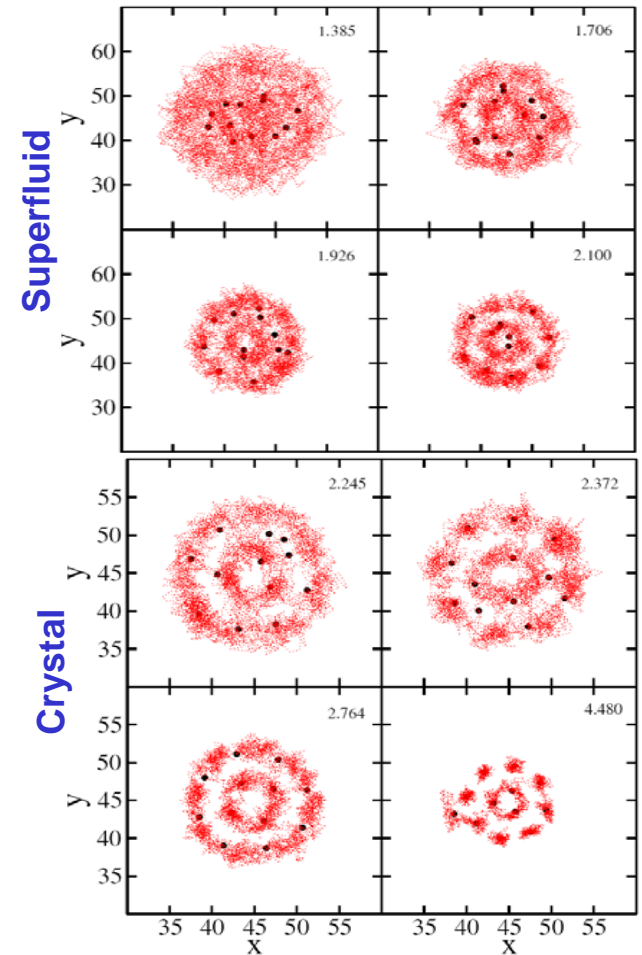


# Dipoles in 2D Harmonic Trap

$$H = \sum_{i=1}^N \left[ \frac{\mathbf{p}_i^2}{2m} + \frac{m\omega^2 \mathbf{r}_i^2}{2} \right] + \sum_{i<j} \frac{D}{|\mathbf{r}_i - \mathbf{r}_j|^3}$$
$$r_0 = (\bar{D}/m\omega^2)^{1/5} \quad l_{ho}^2 = \hbar/(m\omega)$$

$$r_s = r_0/l_{ho}$$

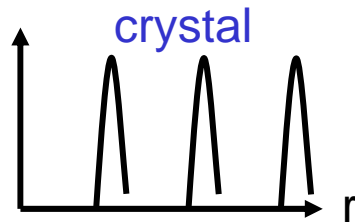
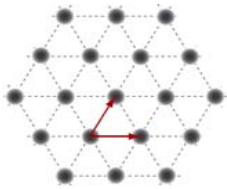
- New Features:
  - Shell structure
    - Indication of two-stage Crystal/SF melting transition
    - Orientational melting
  - Radial melting
  - “Magic Numbers”



# Why?

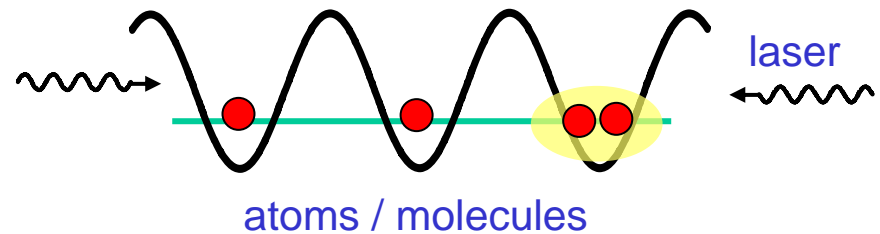
## Self-Assembled Crystals

- interest *per se*
  - quantum phase / transitions
- high-density stable ensembles



## Optical Lattices

- Mott insulator



switch off (destructive) short range interactions / collisions

- lattices parameters / dynamics
  - small lattice constants ☺
  - phonons

superradiance

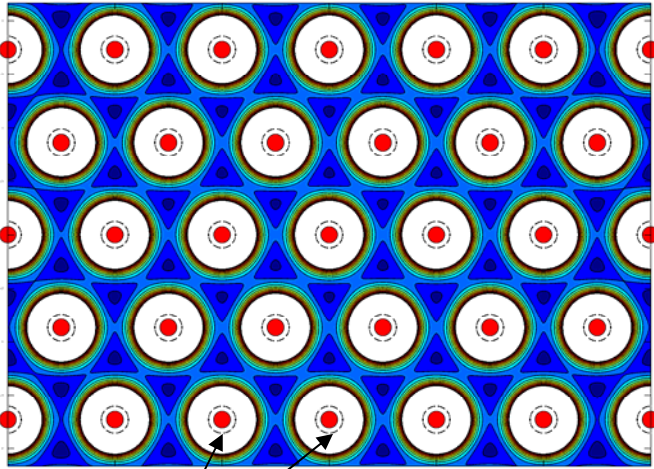
- ...
  - spacing: laser wave length  $\lambda$



# Potential future applications ...

↔ optical lattices

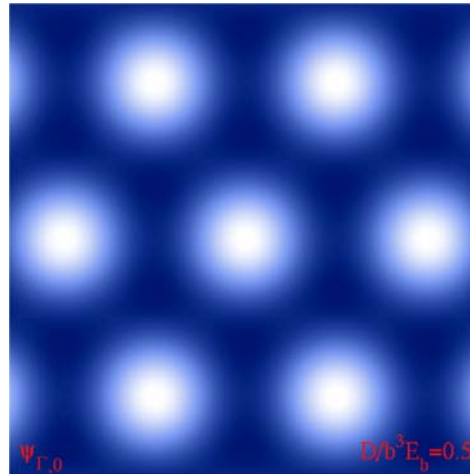
- spin lattice toolbox
  - add spin degree of freedom & engineer spin dependent interactions
- Hubbard models with the new lattices



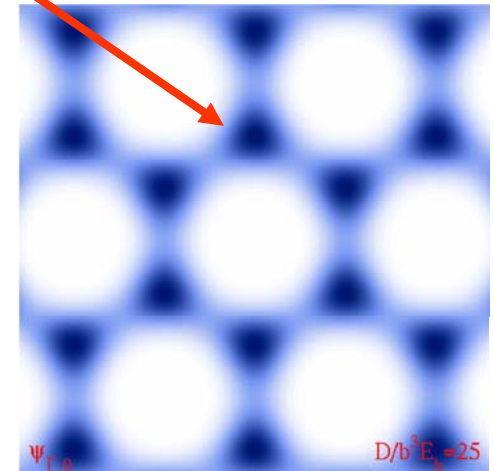
small lattice constants ☺

2. particles moving in this lattice will see a honey comb lattice= Hubbard model

1. dipoles generate a 2D periodic lattice = trap



shallow trap: quasi-free



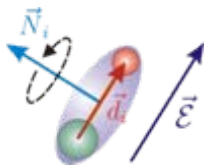
strong trap: tight binding

- bi- and multilayer systems

SF/C Transition point in homogeneous systems:

$$r_d^c \equiv Dm(2\pi)^2 / h^2 a$$

$$= d^2 m \pi / h^2 a \epsilon_0 \approx 13-15$$



## Polar molecules

Examples: CaF, CaCl, SrO

Obtain a crystal for

$$a < a_c = \frac{d^2 m \pi}{(h^2 \epsilon_0) r_d^c}$$

Requirements on confinement in z-direction:

$$E_{\text{int}}(a_c) < \hbar \omega_z$$

	CaF	CaCl	SrO
m[amu]	59	75 77	103
d[D]	3.07(7)	4.265(3)	8.9
$a_c$ [nm]	568.5	1326.5 1361.9	8180.3
$E_{\text{int}}(a_c)$ [KHz]	7.9	1.14 1.05	0.021

Ok!

Ok!

Preparation from optical lattice:  $E_{\text{int}}(\lambda/2) < \hbar \omega_z$

For CaF:  $\lambda/2 \sim 250$  nm,  
 $E_{\text{int}}(\lambda/2) \sim 93$  KHz  
 $r_d \sim 34 > r_d^c$

$$r_d(\lambda/2) > r_d^c$$

## Rydberg states

Example: Hydrogen Stark states:

• Stark splitting:  $\Delta E_{nqm} = \frac{3}{2} e a_0 \epsilon_0 \mathcal{E} n q$

with  $q = n-1, \dots, -(n-1)$

• effective dipole moment:

$$\mu_{\text{eff}} = \frac{3}{2} e a_0 \frac{\Omega^2}{4\Delta^2} n q \propto n^2$$

use  $q = n-1$

• Condition for crystal formation:

$$a_{\text{Ryd}} \ll a < a_c$$

$$a_{\text{Ryd}} = \frac{3}{2} a_0 n^2$$

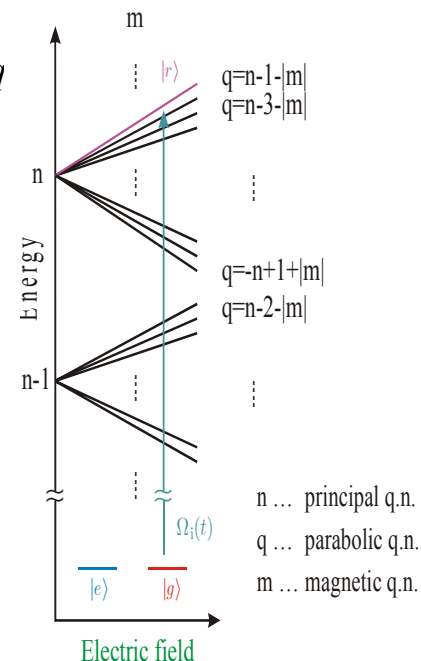
(size of Rydberg state)

$$a_c \equiv \frac{\mu_{\text{eff}}^2 m}{4\pi \epsilon_0 \hbar^2 \gamma_c}$$

(Kin. energy = pot. Energy)

• Example: (values for Hydrogen)

	$n = 20, \frac{\Omega}{\Delta} = 0.2$	$n = 20, \frac{\Omega}{\Delta} = 0.1$
$a_{\text{Ryd}}$ [μm]	31	31
$a_{c,2D}$ [μm]	209	13
$a_{c,1D}$ [μm]	$3.1 \times 10^4$	$2 \times 10^3$
$\mu_{\text{eff}}$ [D]	14.5	3.6
$\tau$ [μs]	50	200
$E_{\text{int}}(a_c)$	3.4 MHz	0.9 GHz



Example: Rb<sup>87</sup>  
Crystal for  $a < a_c \sim 1.1$  μm

Preparation from optical lattice:  
 $\lambda/2 = a = 390$  nm,  
 $E_{\text{int}}(\lambda/2) \sim 33$  KHz  
 $r_d \sim 44 \rightarrow$  Ok!

• Other requirements:  
dipole interaction < Level splitting

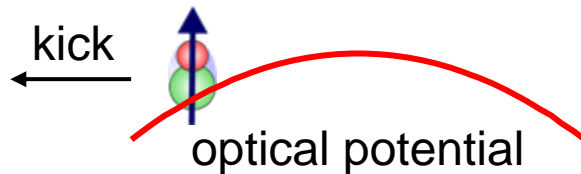
# Preparation of a Crystal

- Preparation from an optical lattice:
  - e.g. from a Mott state: switch optical lattice off, and dipoles on

## Detection of the shell-structure in real space

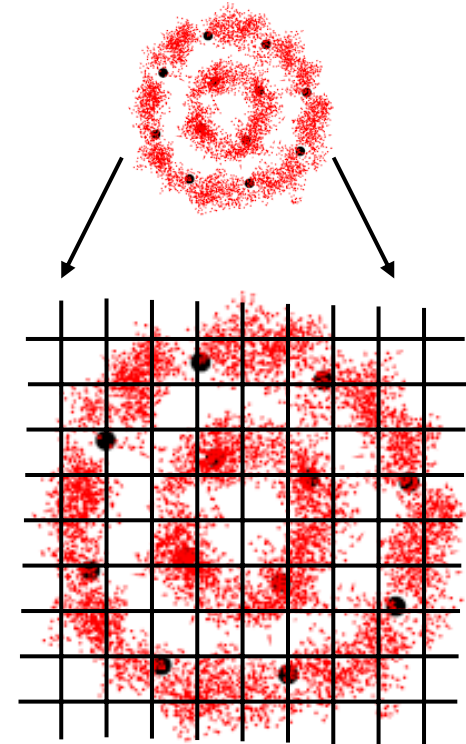
- “Magnifying lens” to see the shell structure

- defocusing „lens“

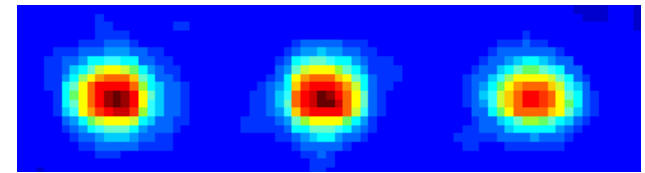


- expand

- focus = recapture, e.g. in optical lattice

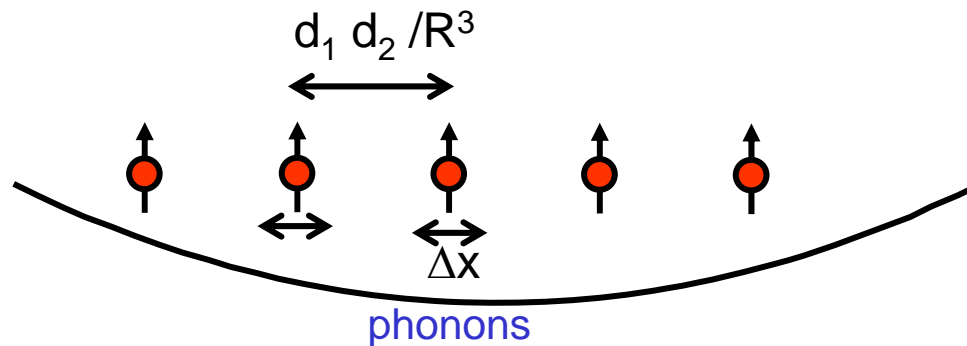


# Crystals for Quantum Computing



compare: ionic Coulomb crystal

- Ion trap like quantum computing with phonons as a bus; moving molecules

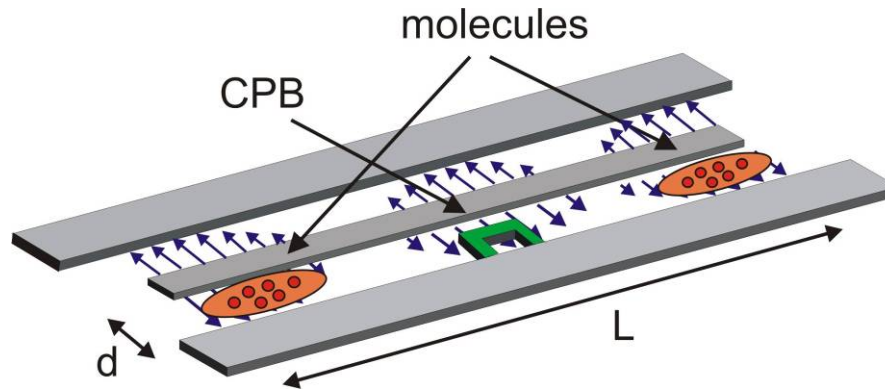


(breathing mode indep of # molecules)

- ✓ dipole moment variable
- ✓ spin dependent dipole
- ✓ qu melting / quantum statistics
- ✓ addressing (?)

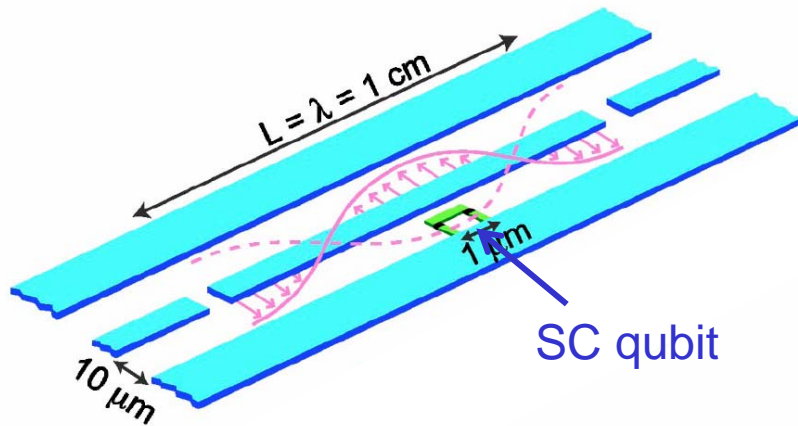
- Ensemble memory: dephasing / avoiding collision dephasing in a 1D and 2D WC
    - ensemble qubit in 2D configuration
    - [there is an instability: rotationally excited qubit -> spin waves]
- 
- ... vs. coupling *single* molecules to a strip line cavity
    - CQED type quantum computing

# Molecular ensembles + Circuit CQED



# 1. strong CQED with superconducting circuits

- Cavity QED



$$H = \omega_c a^\dagger a + \frac{1}{2} \omega_q(t) \sigma_z + g(a \sigma_+ + \text{h.c.})$$

Jaynes-Cummings

parameters:

cavity frequency  $\omega_c \sim 2\pi \times 10$  GHz

cavity damping  $\kappa \sim 2\pi \times 1 \dots 0.01$  MHz

SC qubit - cavity coupling  $g \sim 2\pi \times 30$  MHz

SC qubit damping  $\Gamma \sim 2\pi \times 1$  MHz

good cavity

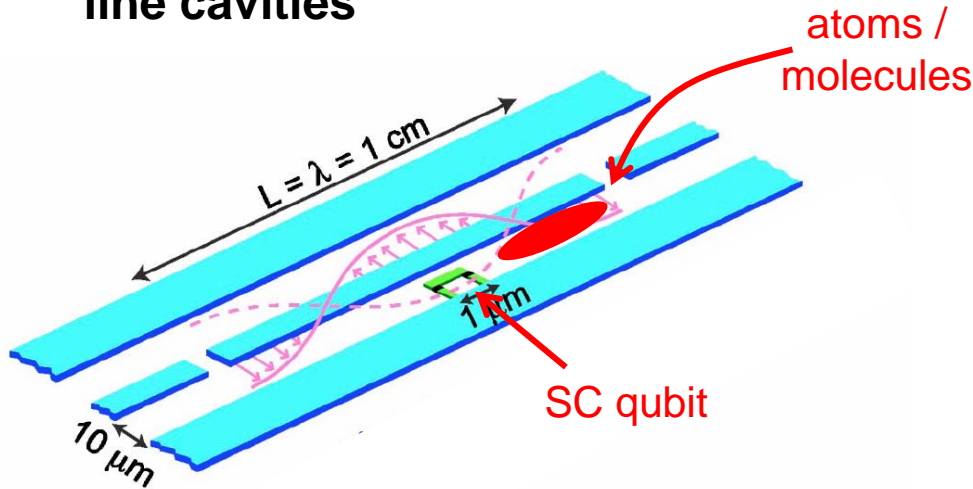
strong coupling!  
(mode volume  $V/\lambda^3 \approx 10^{-5}$ )

"not so great" qubits

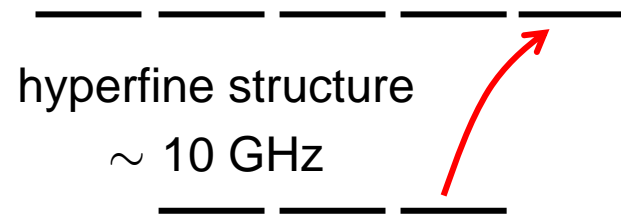
- [... similar results expected for coupling to quantum dots (Delft)]
- [compare with CQED with atoms in optical and microwave regime]

## 2. ... coupling atoms or molecules

- **superconducting transmission line cavities**



- **hyperfine excitation of BEC / atomic ensemble**



$$g \sim 2\pi \times 80 \text{ Hz} \sqrt{\#\text{atoms}}$$

- **rotational excitation of polar molecule(s)**



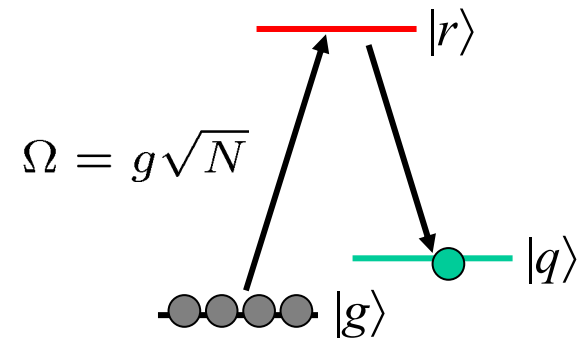
$$g \sim 2\pi \times 10 \text{ KHz} \sqrt{\#\text{molecules}}$$

$$\sim 2\pi \times 1 \dots 10 \text{ MHz ensemble}$$



- **Remarks:**
  - time scales compatible
  - laser light + SC is a problem: we must *move* atoms / molecules to interact with light (?)
  - traps / surface  $\sim 10 \mu\text{m}$  scale
  - low temperature: SC, black body...

### 3. Atomic / molecular ensembles: collective excitations as Qubits



- ground state

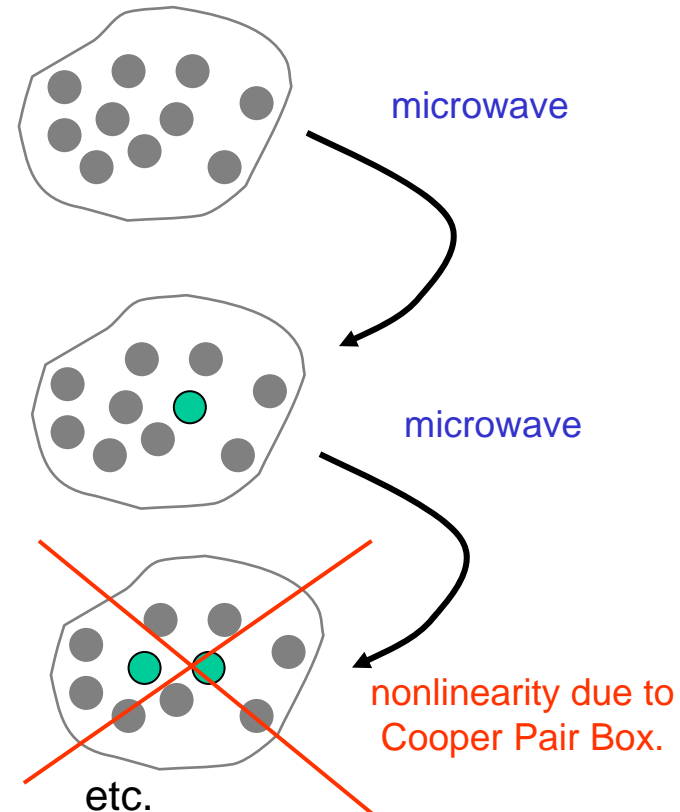
$$\begin{aligned}
 |g^N\rangle &= |g_1\rangle|g_2\rangle \dots |g_N\rangle \\
 &\equiv |\text{vac}\rangle \equiv |0\rangle_e
 \end{aligned}$$

- one excitation (Fock state)

$$\begin{aligned}
 |g^{N-1}q\rangle &= \frac{1}{\sqrt{N}} \sum_i^N |g_1\rangle \dots |q_i\rangle \dots |g_N\rangle \\
 &\equiv m^\dagger |\text{vac}\rangle \equiv |1\rangle_e
 \end{aligned}$$

$\curvearrowright [m, m^\dagger] \approx 1$  harmonic oscillator

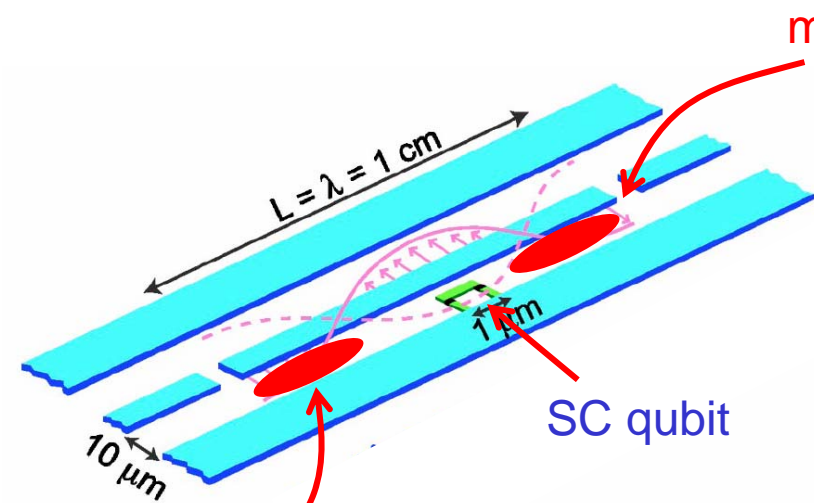
- two excitations ... eliminate?
  - in AMO: dipole blockade, measurements ...



- 
- also: ensembles as continuous variable quantum memory (Polzik, ...)
  - collisional dephasing (?)



# 4. Hybrid Device: solid state processor & molec mem



molecules:  
qubit 2

molecules:  
qubit 1

SC qubit

time independent

$$H = \sum_{i=1}^2 \omega_e(t) m_i^\dagger m_i + \omega_c a^\dagger a + \frac{1}{2} \omega_q(t) \sigma_z + g_c (a \sigma_+ + \text{h.c.}) + \sum_{i=1}^2 g_{i;e}(t) (a m_i^\dagger + \text{h.c.})$$

ensemble  
qubits

solid state system

swap molecule -  
cavity

+ dissipation (master equation)

## 5. Examples of Quantum Info Protocols

- SWAP

$$(\alpha|0\rangle_{cq} + \beta|1\rangle_{cq}) |0\rangle_c |0\rangle_e$$

Cooper Pair

$$\leftrightarrow |0\rangle_{cq} (\alpha|0\rangle_c + \beta|1\rangle_c) |0\rangle_e$$

cavity (bus)

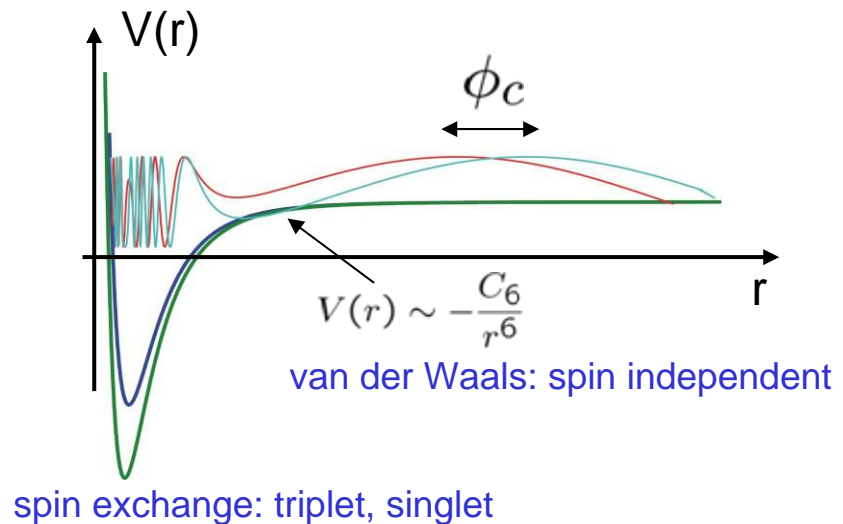
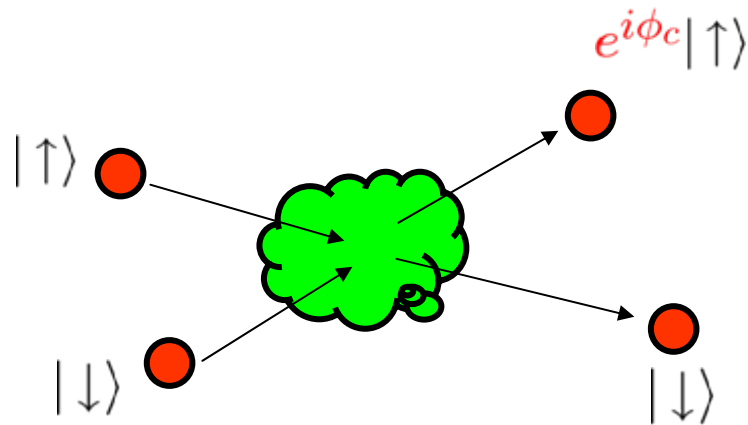
$$\leftrightarrow |0\rangle_{cq} |0\rangle_c (\alpha|0\rangle_e + \beta|1\rangle_e)$$

molec ensemble

- Single qubit rotations via SC qubit
- Universal 2-Qubit Gates via SC qubit
- measurement via ensemble / optical readout or SC qubit / SET

Atomic ensembles complemented by deterministic entanglement operations

# [Study of Decoherence of Molecular Ensemble Qubit : quantum memory & gate operations]



Collision rates (estimate): CaF,  $n=10^{12} \text{ cm}^{-3}$

$T=1 \text{ } \mu\text{K}$ : ( s-wave scattering )

$$\bar{a} = \sqrt[4]{mC_6/\hbar^2} \approx 450a_B$$

$$\gamma_{\text{col}} \approx 50 \text{ Hz}$$

$T=1 \text{ mK}$ : ( unitarity limit )

$$\sigma \leq \frac{4\pi}{k^2} \sum_{l=0}^{l_{\text{max}}} (2l + 1)$$

$$\gamma_{\text{col}} < 400 \text{ Hz}$$

Quantum kinetic theory: Bradley, Gardiner, J. Phys. B 35, 4299 (2002)  
detailed collision physics: J. Bohn, R. Krems, ...

# Summary: QIPC & Cond Mat with Polar Molecules

- single molecules / molecular ensembles
- coupling to optical & microwave fields
  - trapping / cooling
  - CQED (strong coupling)
  - spontaneous emission / engineered dissipation
- interfacing solid state / AMO & microwave / optical
  - strong coupling / dissipation
- collisional interactions
  - quantum deg gases / Wigner crystals (ion trap like qc)
  - WC / dephasing