Santa Barbara, February 17-20th 2016

Around the time-time correlations in KPZ growth models

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: $\ell(\pi) = \sum_{(i,j)\in\pi} \omega(i,j)$

• Maximal length: $L_{(m,n)} = \max_{\pi:(1,1)\to(m,n)} \ell(\pi)$



TASEP with step initial conditions

• TASEP: Totally Asymmetric Simple Exclusion Process

• Dynamics

Independently, particles jump on the right site with rate 1, provided the right is empty. \Leftrightarrow Waiting time $\exp(1)$ -distributed



- Particles are ordered: position of particle k is $x_k(t)$
- Initial condition: $x_k(0) = -k$, $k = 1, 2, \ldots$

• $\omega(i, j) \sim \exp(1)$ is the waiting time (once allowed) of the particle j to do its ith jump (= from -j + i - 1 to -j + i) $\Rightarrow L_{(m,n)}$ is the time when particle n reaches site -n + m

$$\mathbb{P}(L_{(m,n)} \le t) = \mathbb{P}(x_n(t) + n \ge m)$$

 Similarly, one can define LPP between two sets of points as well as TASEP with other initial conditions.

Three geometries

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- Point-to-point problem / step initial conditions
- Point-to-line problem / flat initial conditions
- Point-to-random walk line / stationary initial conditions



Different space-time cuts

- Cut at fixed vertical coordinate equal to t: TASEP configuration at time t
- Cut at j = n with n fixed: trajectory of a tagged particle
- Cut at j = i: integrated current at the origin



Figure by Michael Prähofer

• A simulation of tagged particle vs. integrated current for TASEP with initial half-flat initial conditions

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Space-time scaling and slow decorrelation

• The space-time is non-trivially fibred: spatial correlation length is $\mathcal{O}(t^{2/3})$, but in space-time, there are some directions (characteristic lines) with scaling exponent 1: slow decorrelations phenomenon

Ferrari'08;Corwin,Ferrari,Péché'10

• More generally (flat IC case): process to be studied



Characteristic lines for the step IC (left) and for flat IC (right)

- We consider flat and stationary with density 1/2 so that $\{(0,t),t\geq 0\}$ is a characteristic line
- From now on, we study the integrated current through the origin during time $[0,t],\ J(t)$
- Rescaled process

$$\tau \mapsto \mathcal{X}(\tau) = \lim_{t \to \infty} -t^{-1/3} \left(J(\tau t) - \frac{1}{4} \tau t \right).$$

One-point distribution

$$\mathbb{P}(\mathcal{X}^{\text{step}}(1) \le 2^{-4/3}s) = F_{\text{GUE}}(s),$$
$$\mathbb{P}(\mathcal{X}^{\text{flat}}(1) \le 2^{-1}s) = F_{\text{GOE}}(2s),$$
$$\mathbb{P}(\mathcal{X}^{\text{stat}}(1) \le 2^{-4/3}s) = F_{\text{BR}}(s).$$

Baik, Rains'99,'00; Johansson'00; Prähofer, Spohn'02

- Experimental results by Kazumasa Takeuchi Takeuchi, Sano'12 See also his talk of last week: http://online.kitp.ucsb. edu/online/randomkpz16/takeuchi/
- Formula for $\mathbb{P}\left(\mathcal{X}(\tau) \leq s_1, \mathcal{X}(1) \leq s_2\right)$ using replica approach in a polymer model

Dotsenko'13

• "Finite time" formula for the two-time distribution in a semi-directed polymer model and rigorous limiting formula for $\mathbb{P}(\mathcal{X}(\tau) \leq s_1, \mathcal{X}(1) \leq s_2)$

Johansson'15

• Covariance behavior as $\tau \to 0$:

$$C^{\text{step}}(\tau) = \mathcal{O}(\tau^{2/3}),$$

$$C^{\text{flat}}(\tau) = \mathcal{O}(\tau^{4/3}),$$

$$C^{\text{stat}}(\tau) = \mathcal{O}(\tau^{2/3}).$$

• Covariance behavior as $\tau \to 1$:

$$C^{\text{step}}(\tau) = C^{\text{step}}(1) - \mathcal{O}((1-\tau)^{2/3}),$$

$$C^{\text{flat}}(\tau) = C^{\text{flat}}(1) - \mathcal{O}((1-\tau)^{2/3}),$$

$$C^{\text{stat}}(\tau) = C^{\text{stat}}(1) - \mathcal{O}((1-\tau)^{2/3}).$$

• We are going to give a heuristic argument (no proofs) that explains the above behaviors.

Point-to-point problem

• Set
$$A_{\tau} = (\tau t/4, \tau t/4)$$
 and $I_{\tau}(u) = A_{\tau} + u(\tau t)^{2/3}(1, -1)$.

 $\int \pi_{\tau}$

(i) Step IC

Point-to-point problem

• Set $A_{\tau} = (\tau t/4, \tau t/4)$ and $I_{\tau}(u) = A_{\tau} + u(\tau t)^{2/3}(1, -1)$.

Then as $t \to \infty$ one has <code>Johansson'03+Corwin,Ferrari,Péché'10</code>

$$\frac{\frac{L_{0\to A_{\tau}} - \tau t}{t^{1/3}}}{t^{1/3}} \simeq \tau^{1/3} \mathcal{A}_2(0),$$

$$\frac{\frac{L_{0\to I_{\tau}(u)} - \tau t}{t^{1/3}}}{t^{1/3}} \simeq \tau^{1/3} \left(\mathcal{A}_2(u) - u^2\right),$$

$$\frac{L_{I_{\tau}(u)\to A_1} - (1-\tau)t}{t^{1/3}} \simeq (1-\tau)^{1/3} \left[\tilde{\mathcal{A}}_2(u\hat{\tau}^{2/3}) - (u\hat{\tau}^{2/3})^2\right],$$

where \mathcal{A}_2 and \mathcal{A}_2 are two independent Airy₂ processes. • We have

$$\mathcal{X}^{\text{step}}(\tau) = \tau^{1/3} \mathcal{A}_2(0)$$

and, using $L_{0 \to A_1} = \max_u \left(L_{0 \to I_\tau(u)} + L_{I_\tau(u) \to A_1} \right)$, also

 $\mathcal{X}^{\text{step}}(1) = \tau^{1/3} \max_{u \in \mathbb{R}} \left\{ \mathcal{A}_2(u) - u^2 + \hat{\tau}^{-1/3} \tilde{\mathcal{A}}_2(u \hat{\tau}^{2/3}) - u^2 \hat{\tau} \right\}$

with $\hat{\tau} = \tau/(1-\tau)$.

Point-to-point problem: $\tau \rightarrow 0$ limit

- Using $\hat{\tau}^{-1/3}\tilde{\mathcal{A}}_2(u\hat{\tau}^{2/3}) \simeq B(u)$ a Brownian motion Hägg'07,Corwin,Hammond'11
- For small τ ,

$$C^{\text{step}}(\tau) = \text{Cov}\left(\mathcal{X}^{\text{step}}(\tau), \mathcal{X}^{\text{step}}(1)\right)$$
$$\simeq \tau^{2/3} \text{Cov}\left(\mathcal{A}_2(0), \max_{u \in \mathbb{R}} \left\{\mathcal{A}_2(u) - u^2 + B(u)\right\}\right),$$

 $\bullet\,$ Conditioning on B

$$C^{\text{step}}(\tau) \simeq \tau^{2/3} \mathbb{E}\Big[\operatorname{Cov} \left(\mathcal{A}_2(0), \max_{u \in \mathbb{R}} \left\{ \mathcal{A}_2(u) - u^2 + B(u) \right\} \big| B \right) \Big].$$

• For typical realizations of *B*, the maximum is reached for *u* of order 1, where the last covariance if of order 1, leading to

$$C^{\operatorname{step}}(\tau) = \mathcal{O}(\tau^{2/3})$$

Point-to-point problem: $\tau \rightarrow 1$ limit

• In this case the maximum is reached for $u = \mathcal{O}((1-\tau)^{2/3})$. Set $u = v(1-\tau)^{2/3}/\tau^{2/3}$. As $\tau \to 1$, and conditioning on $\tilde{\mathcal{A}}_2$, $C^{\text{step}}(\tau) = \text{Cov}\left(\mathcal{X}^{\text{step}}(\tau), \mathcal{X}^{\text{step}}(1)\right)$ $\simeq \mathbb{E}\Big[\operatorname{Cov}\left(\mathcal{A}_2(0), \max_{v \in \mathbb{R}}\left\{\mathcal{A}_2(v(1-\tau)^{2/3}) + (1-\tau)^{1/3}(\tilde{\mathcal{A}}_2(v) - \tilde{\mathcal{A}}_2(0) - v^2)\right\} | \tilde{\mathcal{A}}_2 \Big)\Big]$

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 For typical realizations of the process A₂, the maximum is reached for v of order 1. Also, A₂ is locally Brownian, in particular

Cov
$$(\mathcal{A}_2(0), \mathcal{A}_2(v(1-\tau)^{2/3})) \simeq$$
Var $(\mathcal{A}_2(0)) - |v|(1-\tau)^{2/3}$

Prähofer, Spohn'02, Widom'03

• Using the independence of $\tilde{\mathcal{A}}_2$ and \mathcal{A}_2 , we thus expects that

$$C^{\operatorname{step}}(\tau) = \operatorname{Var}\left(\mathcal{A}_2(0)\right) - \mathcal{O}((1-\tau)^{2/3}).$$

Point-to-line problem

• Set

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Point-to-line problem

• Set $A_{\tau} = (\tau t/4, \tau t/4)$ and $I_{\tau}(u) = A_{\tau} + u(\tau t)^{2/3}(1, -1)$. Then as $t \to \infty$ one has

Borodin, Ferrari, Prähofer, Sasamoto'07+Corwin, Ferrari, Péché'10

$$\frac{L_{\mathcal{L}\to A_{\tau}} - \tau t}{t^{1/3}} \simeq \tau^{1/3} \mathcal{A}_{1}(0),$$

$$\frac{L_{\mathcal{L}\to I_{\tau}(u)} - \tau t}{t^{1/3}} \simeq \tau^{1/3} \mathcal{A}_{1}(\tilde{c}u),$$

$$\frac{L_{I_{\tau}(u)\to A_{1}} - (1-\tau)t}{t^{1/3}} \simeq (1-\tau)^{1/3} [\tilde{\mathcal{A}}_{2}(u\hat{\tau}^{2/3}) - (u\hat{\tau}^{2/3})^{2}],$$

where the Airy_1 process \mathcal{A}_1 is independent of the Airy_2 process $\tilde{\mathcal{A}}_2.$

We have

$$\mathcal{X}^{\mathrm{flat}}(\tau) = \tau^{1/3} \mathcal{A}_1(0)$$

and

$$\mathcal{X}^{\text{flat}}(1) = \max_{u \in \mathbb{R}} \left\{ \tau^{1/3} \mathcal{A}_1(u) + (1-\tau)^{1/3} \tilde{\mathcal{A}}_2(u\hat{\tau}^{2/3}) - u^2 \tau^{1/3} \hat{\tau} \right\}$$

Point-to-line problem: $\tau \rightarrow 0$ limit

- For small $\tau,$ the maximum over u is typically taken for $u\sim \mathcal{O}(\tau^{-2/3})$
- Since the covariance of the Airy₁ process decays superexponentially in *u* Bornemann,Ferrari,Prähofer'08 their contribution to

$$C^{\text{flat}}(\tau) = \text{Cov}(\mathcal{X}(\tau), \mathcal{X}(1))$$

is negligible.



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is negligible.

• With probability $\mathcal{O}(\tau^{2/3})$, the maximum is take for $u = \mathcal{O}(1)$. In this case the heuristic is the same as for the point-to-point case. Thus,

$$C^{\text{flat}}(\tau) = \mathcal{O}(\tau^{2/3}C^{\text{step}}(\tau)) = \mathcal{O}(\tau^{4/3})$$

Point-to-random walk line problem

• Set
$$A_{\tau} = (\tau t/4, \tau t/4)$$
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Point-to-random walk line problem

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Imamura, Sasamoto'05; Baik, Ferrari, Péché'10

$$\frac{\frac{L_{0 \to A_{\tau}} - \tau t}{t^{1/3}}}{t^{1/3}} \simeq \tau^{1/3} \mathcal{A}_{\text{stat}}(0),$$
$$\frac{\frac{L_{0 \to I_{\tau}(u)} - \tau t}{t^{1/3}}}{t^{1/3}} \simeq \tau^{1/3} \mathcal{A}_{\text{stat}}(u),$$
$$\frac{L_{I_{\tau}(u) - A_{1}} - (1 - \tau)t}{t^{1/3}} \simeq (1 - \tau)^{1/3} \left[\tilde{\mathcal{A}}_{2}(u\hat{\tau}^{2/3}) - (u\hat{\tau}^{2/3})^{2}\right],$$

where the processes \mathcal{A}_{stat} and $\tilde{\mathcal{A}}_2$ are independent.

We have

$$\mathcal{X}^{\mathrm{stat}}(\tau) = \tau^{1/3} \mathcal{A}_{\mathrm{stat}}(0)$$

and

$$\mathcal{X}^{\text{stat}}(1) = \max_{u \in \mathbb{R}} \left\{ \tau^{1/3} \mathcal{A}_{\text{stat}}(u) + (1-\tau)^{1/3} \tilde{\mathcal{A}}_2(u\hat{\tau}^{2/3}) - u^2 \tau^{1/3} \hat{\tau}^{-1} \right\}$$

Point-to-random walk line problem: $\tau \rightarrow 0$ limit

 For small τ, the maximum over u is typically taken for u ~ O(τ^{-2/3}), the maximizers to A_τ and A₁ uses then different noises (independent) except for the noise on the axis



- For small τ, the maximum over u is typically taken for u ~ O(τ^{-2/3}), the maximizers to A_τ and A₁ uses then different noises (independent) except for the noise on the axis
- Thus we expect

$$C^{\text{stat}}(\tau) = \text{Cov}(\mathcal{X}^{\text{stat}}(\tau), \mathcal{X}^{\text{stat}}(1))$$

$$\simeq \mathcal{O}(t^{-2/3}) \text{Cov}\left(L_{(-1,-1)\to C_{\tau}}, L_{(-1,-1)\to C_{1}}\right).$$

• The sums of random variables in the LPP problem between the origin and C_1 (and C_{τ}) are asymptoically Brownian motions. Thus implies that

$$C^{\mathrm{stat}}(\tau) = \mathcal{O}(\tau^{2/3})$$

Numerical results: step IC



Figure: Plot of $\tau \mapsto \text{Cov}(\mathcal{X}^{\text{step}}(\tau), \mathcal{X}^{\text{step}}(1)) / \text{Var}(\mathcal{X}^{\text{step}}(1))$. The top-left (resp. right-bottom) inset is the log-log plot around $\tau = 0$ (resp. $\tau = 1$).

Numerical results: flat IC



Figure: Plot of $\tau \mapsto \operatorname{Cov}(\mathcal{X}^{\operatorname{flat}}(\tau), \mathcal{X}^{\operatorname{flat}}(1)) / \operatorname{Var}(\mathcal{X}^{\operatorname{flat}}(1))$. The top-left (resp. right-bottom) inset is the log-log plot around $\tau = 0$ (resp. $\tau = 1$).

Numerical results: stationary IC



Figure: Plot of $\tau \mapsto \operatorname{Cov}(\mathcal{X}^{\operatorname{stat}}(\tau), \mathcal{X}^{\operatorname{stat}}(1))/\operatorname{Var}(\mathcal{X}^{\operatorname{stat}}(1))$. The top-left inset is the log-log plot around $\tau = 0$ and the right-bottom inset is the log-log plot around $\tau = 1$. The fit is made with the function $\tau \mapsto \frac{1}{2}(1 + \tau^{2/3} - (1 - \tau)^{2/3})$.

- TASEP with stationary initial conditions
- $\bullet~{\rm j}(t)$ is the empirical current across the bond (0,1) from which

$$J(t) = \int_0^t ds \, \mathbf{j}(s)$$

• Two-point function (stationary covariance)

$$S(j,t) = \mathbb{E}(\eta_j(t)\eta_0(0)) - \rho^2.$$

• A sum rule:

$$\operatorname{Var}(J(t)) = \sum_{j \in \mathbb{Z}} |j| S(j, t) - \sum_{j \in \mathbb{Z}} |j| S(j, 0)$$

Stationary case: current-current correlations

- $j(\rho) = \rho(1-\rho)$ is the expected current with respect to the stationary initial condition with density ρ
- A small perturbation of the steady state will propagate with velocity $v(\rho)=j'(\rho), \ v(1/2)=0.$
- The current-current covariance is then given by

 $\mathbb{E}(\mathbf{j}(t)\mathbf{j}(t')) - \mathbf{j}(\rho)^2 = \rho(1-\rho)\delta(t-t') + h(t-t')$

• The smooth part $h(t-t^{\prime})$ is given by

$$h(t) = -\langle (r_{0,1}^{\mathrm{R}} - j(\rho)) e^{L|t|} (r_{0,1} - j(\rho)) \rangle_{\rho},$$

where for TASEP $r_{0,1}(\eta) = \eta_0(1-\eta_1)$, $r_{0,1}^{\rm R}(\eta) = -(1-\eta_0)\eta_1$; $\langle \cdot \rangle_{\rho}$ is the average with respect to the stationary measure with density ρ ; L is the backwards generator of TASEP.

• According to the KPZ scaling theory:

$$S(j,t) \simeq \chi(\Gamma t)^{-2/3} f_{\rm KPZ}((\Gamma t)^{-2/3}j)$$

where in the special case of TASEP case $\chi = \rho(1-\rho)$, and $\Gamma = \chi^2$. Krug,Meakin,Halpin-Healy'92

Using the sum rule

$$\chi \int_{\mathbb{R}} dx |x| f_{\text{KPZ}}(x) (\Gamma t)^{2/3} \simeq -2 \int_0^t ds \int_s^\infty du h(u),$$

which implies

$$h(t) \simeq -c_0 t^{-4/3}$$
, $c_0 = \frac{1}{9} \Gamma^{2/3} \chi \int_{\mathbb{R}} dx |x| f_{\text{KPZ}}(x)$.

• Covariance of integrate current

$$\operatorname{Cov}(J(t), J(\tau t)) = -\int_0^t \int_0^{\tau t} ds ds' \Big(\int_{\mathbb{R}} du \, h(u) \delta(s - s') - h(s - s') \Big)$$
$$= -\int_0^{\tau t} ds \Big(2 \int_s^\infty ds' h(s') - \int_{\tau t - s}^{t - s} ds' h(s') \Big)$$

• Using

$$h(t) \simeq -c_0 t^{-4/3}$$

one gets

$$\operatorname{Cov}(J(t), J(\tau t)) \simeq (1 + \tau^{2/3} - (1 - \tau)^{2/3}) (\Gamma t)^{2/3} \chi \int_0^\infty dx |x| f_{\mathrm{KPZ}}(x)$$

Numerical results



Figure: The smooth part of the current-current correlations for TASEP. We plot -h(t) and the theoretical large time behavior $0.02013 \cdot t^{-4/3}$.

Numerical results



Figure: Log-log plot of the smooth part of the current-current correlation for TASEP.