

Several Issues in Color Superconductivity

First part: M. Buballa and M. Oertel

- 2SC versus CFL: influence of the strange quark mass
- Neutron stars: constraints by neutrality

Color superconductivity

Cold dense deconfined strongly interacting matter:

- Attractive interaction \rightarrow pair condensation
- Most important example (color $\bar{3}$ channel):

$$\langle \bar{\psi}_{\alpha a A} (\lambda^i)_{ab} (\tau^j)_{AB} (\gamma_5 C)_{\alpha \beta} \bar{\psi}_{\beta b B} \rangle = \Delta_{ij}$$

- λ^i and τ^j any antisymmetric $SU(3)$ matrix
- Two idealized cases:
 - $M_s \rightarrow \infty$:
Only u , d -quarks condense
 \rightarrow only Δ_{ij} with $j = 2$ nonzero
 - Only two out of three quark colors are “gapped”
 $\rightarrow SU(3)_c$ broken down to $SU(2)_c$
 - Chiral symmetry restored
- $M_s = M_u = M_d$:
 u , d and s -quarks condense
Most favored state: $\Delta_{22} = \Delta_{55} = \Delta_{77}$
- “Color-Flavor-Locked (CFL)”-state:
 $SU(3)_c \times SU(3)_L \times SU(3)_R$ broken to $SU(3)_{c+V}$
 \rightarrow color and chiral symmetry broken

Realistic case: $M_u = M_d < M_s$

Which condensation pattern? Step 1: $\mu_u = \mu_d = \mu_s = \mu$

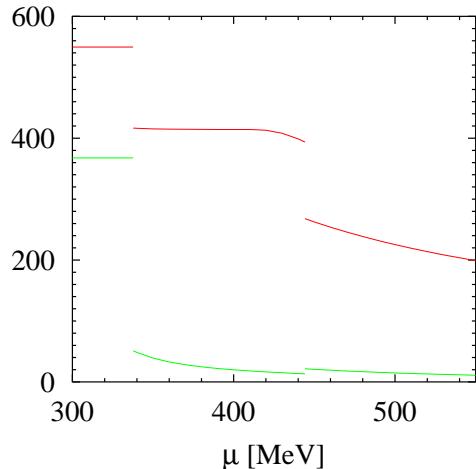
- $M_s \approx \mu$: $p_F^u \approx \mu \gg p_F^s = \sqrt{\mu^2 - M_s^2}$

$$\underline{M_s \ll \mu:} \quad p_F^u \approx p_F^s$$

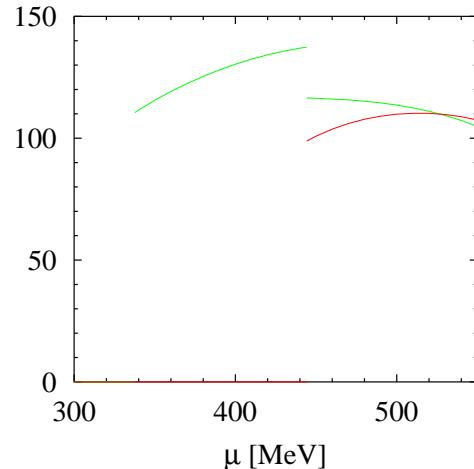
- Influence masses \Leftrightarrow phases:

NJL calculation: Masses related to $\langle \bar{u}u \rangle, \langle \bar{d}d \rangle, \langle \bar{s}s \rangle$

Masses M_u, M_s



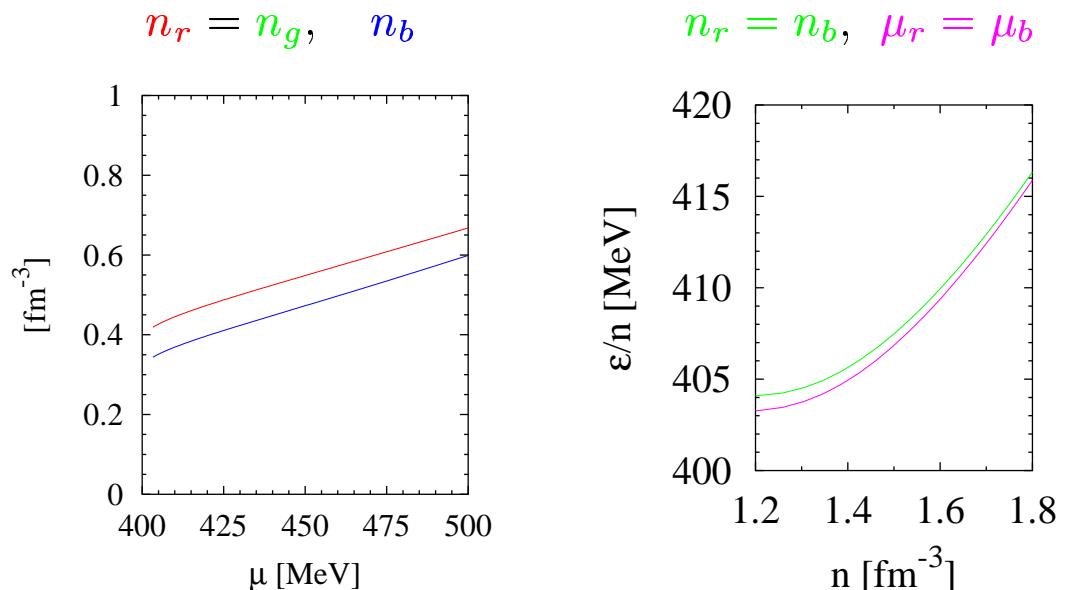
Condensates Δ_{22}, Δ_{55}



(M.B., M.O., NPA '02)

Bulk matter in compact stars

- Constraints:
 - β -equilibrium
 - (electric) charge neutrality
 - color singlet \Rightarrow color neutrality
- Color neutrality: $n_r = n_g = n_b$



(M.B., J. Hošek, M.O., PRD '01)

\Rightarrow minor effect

β -equilibrium + charge neutrality

- β -equilibrium:

$$d \leftrightarrow u + e^- + \bar{\nu}_e \leftrightarrow s \quad \Rightarrow \quad \mu_d = \mu_s = \mu_u + \mu_e$$

- charge neutrality:

$$\frac{2}{3}n_u - \frac{1}{3}n_d - \frac{1}{3}n_s - n_e = 0$$

- simple estimate: massless u and d , no s , no pairing

assumption: $n_e \approx 0$

$$\Rightarrow n_d \approx 2n_u \quad \Rightarrow \quad \mu_d \approx 2^{1/3} \mu_u$$

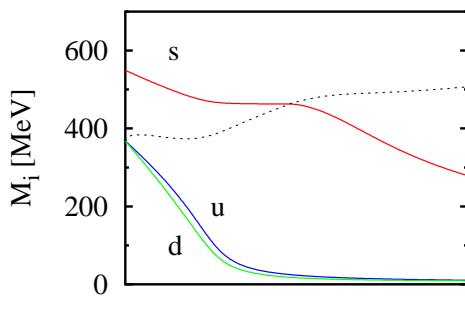
$$\Rightarrow \mu_e \approx (2^{1/3} - 1) \mu_u \simeq \frac{1}{4} \mu_u$$

$$\Rightarrow n_e \approx \frac{1}{3 \cdot 64} n_u \approx 0.005 n_u$$

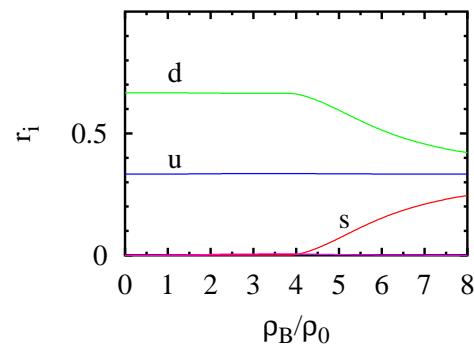
→ almost no electrons

- NJL-calculation: massive u , d , and s , no pairing

masses:



fractions:



(M.B., M.O., PLB '99)

Consequences for diquark pairing

- case 1: M_s large \Rightarrow no strange quarks

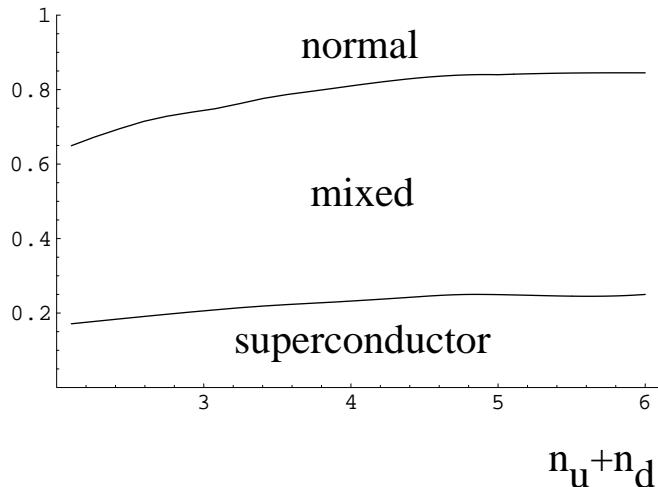
- $n_d \simeq 2 n_u \Rightarrow \mu_d \simeq 2^{1/3} \mu_u \simeq \frac{5}{4} \mu_u$,
- stability criterion: (Rajagopal, Wilczek, PRL '01)

$$\Delta > \sqrt{2} \delta \mu$$

example: $\mu_u = 400$ MeV $\Rightarrow \mu_d = 500$ MeV
 $\Rightarrow \Delta > 140$ MeV

- Model analysis:

$$(n_u - n_d) / (n_u + n_d)$$



(P. Bedaque, NPA '02)

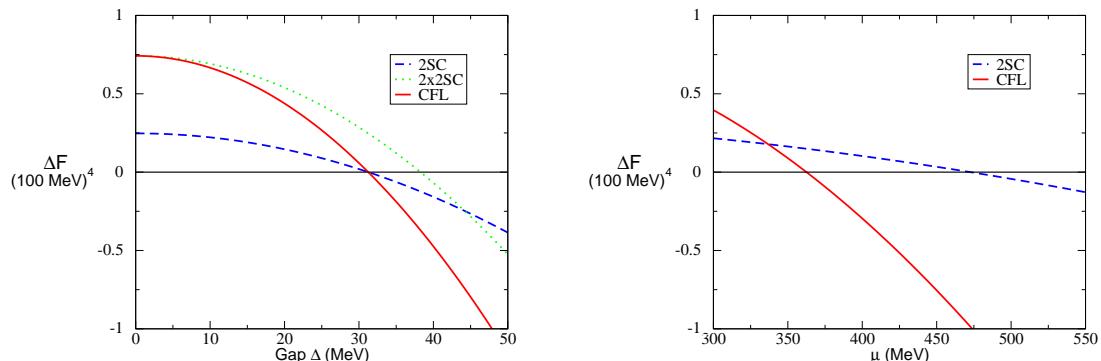
Consequences for diquark pairing

- case 2: M_s small \Rightarrow expansion in M_s

(Alford, Rajagopal, hep-ph/0204001)

$$\rightarrow p_F^d = p_F^u + \frac{M_s^2}{4\bar{\mu}} \quad p_F^s = p_F^u - \frac{M_s^2}{4\bar{\mu}}$$

\rightarrow equal probability for ud - and us -pairing



(Alford, Rajagopal, hep-ph/0204001)

Is the CFL phase charge neutral?

Argument (Rajagopal, Wilczek, PRL '01):

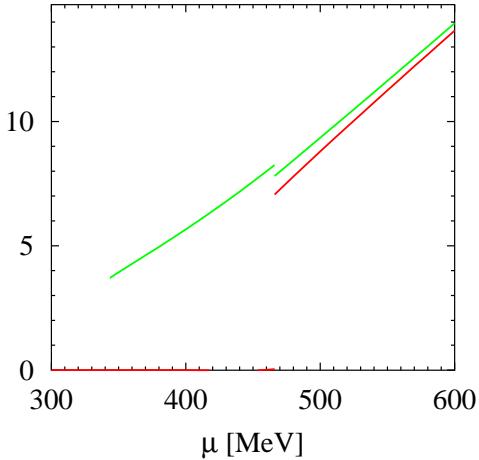
Densities equal \rightarrow charge neutrality enforced

Model calculation:

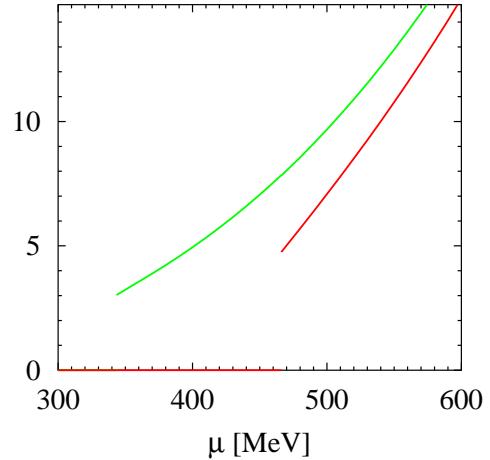
densities including pairing \Leftrightarrow “free ones”

$$-\frac{\partial \Omega}{\partial \mu_i} \quad \frac{1}{\pi^2} (\mu_i^2 - M_i^2)^{3/2}$$

$n_u/\rho_0, n_s/\rho_0$



free $n_u/\rho_0, n_s/\rho_0$



(M.B., F. Neumann, M.O., preliminary result)