Magnetic semiconductor spintronics

Tomas Jungwirth

D. Fang, H. Kurebayashi, J. Wunderlich, K. Vyborny, Liviu P. Zarbo, R.P. Campion, A. Casiraghi, B.L. Gallagher, T. Jungwirth, A. J. Ferguson, *Nature Nanotech. 6 (2011) 413 – 417*

T. Jungwirth, J. Wunderlich, K. Olejník, Nature Mater. 11 (2012), 382-390.

P. Němec, E. Rozkotová, N. Tesařová, F. Trojánek, E. De Ranieri, K. Olejník, J. Zemen, V. Novák, M. Cukr, P. Malý, T. Jungwirth, *Nature Phys. 8 (2012) 411 - 415*.

P. Němec, V. Novák, N. Tesařová, E. Rozkotová, H. Reichlová, D. Butkovičová, F. Trojánek, K. Olejník, P. Malý, R. P. Campion, B. L. Gallagher, Jairo Sinova, and T. Jungwirth, *Nature Commun. 4 (2013) 1422(1) - 1422(8).*

N. Tesařová, P. Němec, E. Rozkotová, J. Zemen, F. Trojánek, K. Olejník, V. Novák, P. Malý, T. Jungwirth, *Nature Photonics* 7 (2013) 492 - 498.

E. De Ranieri, P. E. Roy, D. Fang, E. K. Vehsthedt, A. C. Irvine, D. Heiss, A. Casiraghi, R. P. Campion, B. L. Gallagher, T. Jungwirth, J. Wunderlich, *Nature Mater. 12 (2013) 808 – 814*

P. Wadley, V. Novak, R. P. Campion, C. Rinaldi, X. Marti, H. Reichlova, J. Zelezny, J. Gazquez, M. A. Roldan, M. Varela, D. Khalyavin, S. Langridge, D. Kriegner, F. Maca, J. Masek, R. Bertacco, V. Holy, A. W. Rushforth, K. W. Edmonds, B. L. Gallagher, C. T. Foxon, J. Wunderlich, and T. Jungwirth *, Nature. Commun.* 4 (2013) 2322(1) - 2322(6)

H. Kurebayashi, Jairo Sinova, D. Fang, A. C. Irvine, J. Wunderlich, V. Novak, R. P. Campion, B. L. Gallagher, E. K. Vehstedt, L. P. Zarbo, K. Vyborny, A. J. Ferguson, T. Jungwirth, *arXiv:1306.1893*

X. Marti, I. Fina, Di Yi, Jian Liu, Jiun-Haw Chu, C. Rayan-Serrao, S. Suresha, J. Železný, T. Jungwirth, J. Fontcuberta, R. Ramesh, *arXiv:1303.4704*

T. Jungwirth, J. Wunderlich, V. Novak, K. Olejnik, B. L. Gallagher, R. P. Campion, K. W. Edmonds, A. W. Rushforth, A. J. Ferguson, P. Nemec, *arXiv:1310.1944*

Magnetic semiconductor spintronics

Tomas Jungwirth



Institute of Physics ASCR





Univ. of Nottingham, UK Hitachi and Univ. Cambridge, UK & Japan



Charles Univ., Czech Rep.



Institut de Ciencia de Materials Univ. of California, Berkeley de Barcelona, Spain

Optimizing MBE growth and post-growth annealing of (Ga,Mn)As





Optimized (Ga,Mn)As



Wang et al. PRB '13

Dobrowolska et al. Nature Mater '12

Series of optimized Ga_{1-x}Mn_xAs

Fully reproducible and well behaved FM and degenerate semiconductor



Němec, Novák, TJ et al. Nature Commun. '13

Series of optimized Ga_{1-x}Mn_xAs

Fully reproducible and well behaved FM and degenerate semiconductor



.. and tuneable, and compatible with III-V heterostructure and fabrication techniques

... and strong exchange and spin-orbit ~ 100s meV \rightarrow strong disorder (even unintentional) is not detrimental to spintronics in (Ga,Mn)As

.. and the full range of spintronic effects described qualitatively or semiquantitatively by mutually consistent DFT, TB-Anderson, kinetic-exchange **k.p** models



GaMnAs in spintronics: arXiv:1310.1944





Electrical spin-tranfer torque



Optical spin-tranfer torque



Rossier, MacDonald, et al. '03 Němec, Novák, TJ et al. Nature Phys.'12

Electrical spin-orbit torque



Optical spin-orbit torque



Tesařová, Němec, Novák, TJ et al. Nature Photonics '13 Steady-state carrier spin polarization $\vec{s} \rightarrow \text{torque } \frac{dM}{dt}$ $\frac{d\langle \vec{\sigma} \rangle}{dt} = \frac{1}{i\hbar} \langle [\vec{\sigma}, H] \rangle \qquad \vec{s} = \langle \vec{\sigma} \rangle_{\text{QM averaging in non-equilibrium}}$

Electrical spin injection

Spin-transfer torque

$$\begin{split} H &= H_0 + H_{ex} \\ & \int H_{ex} = J \vec{M} \cdot \vec{\sigma} \\ \text{Steady state} \\ 0 &= \frac{d\vec{s}}{dt} = \frac{J}{\hbar} \vec{s} \times \vec{M} + P \hat{n} \\ & \text{External} \end{split}$$

$$\frac{d\vec{M}}{dt} = \frac{J}{\hbar}\vec{M}\times\vec{s} = P\hat{M}\times(\hat{n}\times\hat{M})$$



Optical spin injection

 $P\hat{n}$

Steady-state carrier spin polarization $\vec{s} \rightarrow \text{torque } \frac{dM}{dt}$ $\frac{d\langle \vec{\sigma} \rangle}{dt} = \frac{1}{i\hbar} \langle [\vec{\sigma}, H] \rangle \qquad \vec{s} = \langle \vec{\sigma} \rangle_{\text{QM averaging in non-equilibrium}}$

Electrical spin injection

Spin-orbit torque

$$\begin{split} H &= H_0 + H_{ex} + H_{so} \\ & \int & \int \text{Internal} \\ \text{Steady state} \\ 0 &= \frac{d\vec{s}}{dt} = \frac{J}{\hbar}\vec{s} \times \vec{M} + \frac{1}{i\hbar} \langle [\vec{\sigma}, H_{so}] \rangle \end{split}$$

$$\frac{d\vec{M}}{dt} = \frac{J}{\hbar}\vec{M} \times \vec{s} = \frac{1}{i\hbar} \langle [\vec{\sigma}, H_{so}] \rangle$$





Steady-state carrier spin polarization $\vec{s} \rightarrow$ torque $\frac{dM}{dt}$

$$\frac{d\langle \vec{\sigma} \rangle}{dt} = \frac{1}{i\hbar} \langle [\vec{\sigma}, H] \rangle$$

 $\vec{s} = \langle \vec{\sigma} \rangle$ Linear response: eigenstates of H & non-equilibrium distribution Electrical drift and relaxation

Spin-orbit torque

$$\begin{split} H &= H_0 + H_{ex} + H_{so} \\ & \int & \int \text{Internal} \\ \text{Steady state} \\ 0 &= \frac{d\vec{s}}{dt} = \frac{J}{\hbar}\vec{s} \times \vec{M} + \frac{1}{i\hbar} \langle [\vec{\sigma}, H_{so}] \rangle \end{split}$$

$$\frac{d\vec{M}}{dt} = \frac{J}{\hbar}\vec{M} \times \vec{s} = \frac{1}{i\hbar} \langle [\vec{\sigma}, H_{so}] \rangle$$



Optical generation and relaxation



Optical spin torques in ferromagnetic semiconductor (Ga,Mn)As



Magneto-optical parameters of (Ga,Mn)As



Tesařová, Němec, Novák, TJ et al. Appl. Phys. Lett. 12

Pump-and-probe magneto-optical signals in (Ga,Mn)As



Tesařová, Němec, Novák, TJ et al. Appl. Phys. Lett. 12

Pump-and-probe magneto-optical signals in (Ga,Mn)As







Optical spin-transfer torque - experiment



Němec, Novák, TJ et al. Nature Phys.'12

Optical spin-orbit torque - experiment



Optical spin-orbit torque - experiment



Optical spin-orbit torque - theory

Optical generation and relaxation



$$\begin{split} \frac{dM}{dt} &= \frac{J}{\hbar} \vec{M} \times \vec{s} = \frac{1}{i\hbar} \langle [\vec{\sigma}, H_{so}] \rangle & H_{ex} = J \vec{M} \cdot \vec{\sigma} \\ \vec{H}_{an} &= -\frac{\partial}{\partial \vec{M}} \sum_{a} \int d\mathbf{k} \epsilon_{a,\vec{k}} f_{a,\vec{k}} = -\sum_{a} \int d\vec{k} \langle a, \vec{k} | \frac{\partial H}{\partial \vec{M}} | a, \vec{k} \rangle f_{a,\vec{k}} \\ &= -\sum_{a} \int d\vec{k} \langle a, \vec{k} | J \vec{\sigma} | a, \vec{k} \rangle f_{a,\vec{k}} = -J \vec{s} \,. \end{split}$$

Photo-hole spin-density ↔ hole-density-dependent magnetic anisotropy field



Electrical spin-tranfer torque

Non-uniform magnetic structure



Optical spin-tranfer torque

Uniform magnet

Electrical spin-orbit torque

Broken inversion-symmetry magnet



Optical spin-orbit torque

Inversion-symmetric magnet



AFM semiconductors: prospect for room-T magnetic-semiconductor spintronics

TJ, Novák, et al. PRB '11, Cava Viewpoint, Physics '11, Máca, TJ et al. JMMM '12, Wadley, TJ, et al. Nature Commun. '13

II-VI	FM T _c (K)	AFM T _N (K)	III-V	FM T _c (K)	AFM
MnO		122	FeN		100
MnS		152	FeP		115
MnSe		173	FeAs		77
MnTe		323	FeSb		100-2
EuO	67		GdN	72	
EuS	16		GdP		15
EuSe		5	GdAs		19
EuTe		10	GdSb		27
					_
			11_\/_1\/_\/	EMT(K)	

I-VI-III-VI	FM T _c (K)	AFM T _N (K)
CuFeO ₂		11
CuFeS ₂		825
CuFeSe ₂		70
CuFeTe ₂		254

Fein		100
FeP		115
FeAs		77
FeSb		100-220
GdN	72	
GdP		15
GdAs		19
GdSb		27
II-V-IV-V	FM T _c (K)	AFM T _N (K)
MnSiN ₂		490
I-II-V	FM T _c (K)	AFM T _N (K)
Ia=Li, Na, Ib=Cu II=Mn V=Sb,As, P	Beleanu et al.	> room T arxiv:13076404

T_N (K)

Antiferromagnetic metals



Spin-orbit induced anisoropic electronic structure: DFTFerromagnetscalculationsAntiferromagnets



Spin-orbit-coupled Mott AFM semiconductor Sr₂IrO₄



Marti,TJ et al. arXiv:13034704

Field-rotation Ohmic AMR of Sr₂IrO₄ AFM semiconductor



Ferromagnets:

$\uparrow \uparrow \uparrow$

Ordered $M \neq 0$: good for manipulation by magnetic field and detection by stray fields,

Magnetic field not employed in advanced spintronics



perturbed by $<\sim$ T produces $<\sim$ T nearby stray field perturbation High T_c not well compatible with semiconductor band structure

Antiferromagnets:

 $\downarrow \uparrow \downarrow \uparrow$

Ordered M=0: bad for manipulation by magnetic field and detection by stray fields, insensitive to <~100T perturbation produces no stray field perturbation High T_N well compatible with semiconductor band structure

AFM-alone room-T memory resistor

Insensitive to magnetic field and no stray field



Marti, TJ et al. preprint '13

Laser-induced ultrafast spin reorientation in the antiferromagnet TmFeO₃

A. V. Kimel¹, A. Kirilyuk¹, A. Tsvetkov¹, R. V. Pisarev² & Th. Rasing¹

Nature '04, Th. Rasing Plenary Wed 16:15



VOLUME 93, NUMBER 11

PHYSICAL REVIEW LETTERS

week ending 10 SEPTEMBER 2004

Ultrafast Manipulation of Antiferromagnetism of NiO

N. P. Duong,¹ T. Satoh,^{1,2} and M. Fiebig^{1,*}