

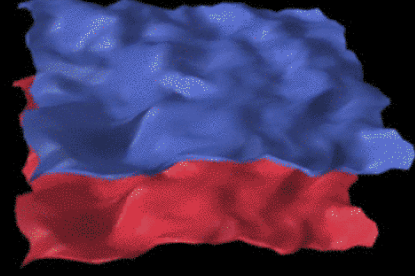
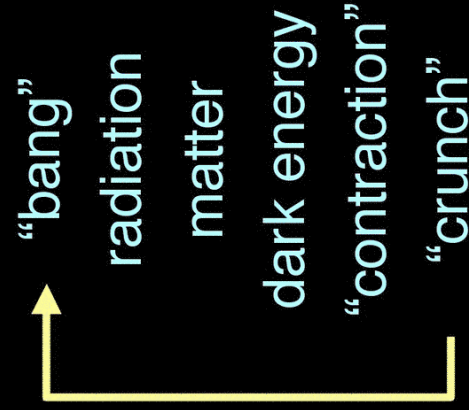
# Recent Progress on Cyclic Universe Models

w/ Neil Turok

Andrew Tolley, Justin Khoury, Latham Boyle, Joel Erickson,  
Daniel Wesley, Steven Gratton

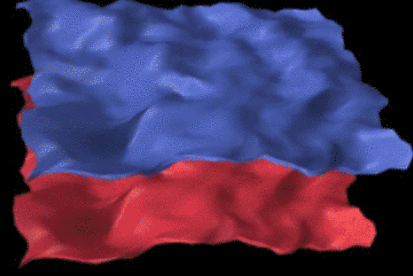
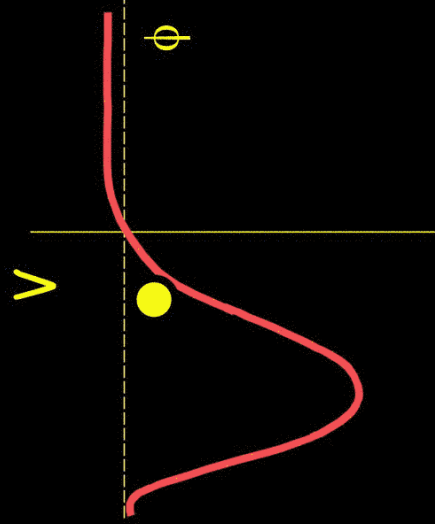
## Cyclic Model

5d Brane Picture



# Cyclic Model

4d Effective Picture



*Many interesting issues*

*is cycling better/worse than inflation?*

*finite or infinite number of cycles ?*

*eternal ?*

*holographic bounds ?*

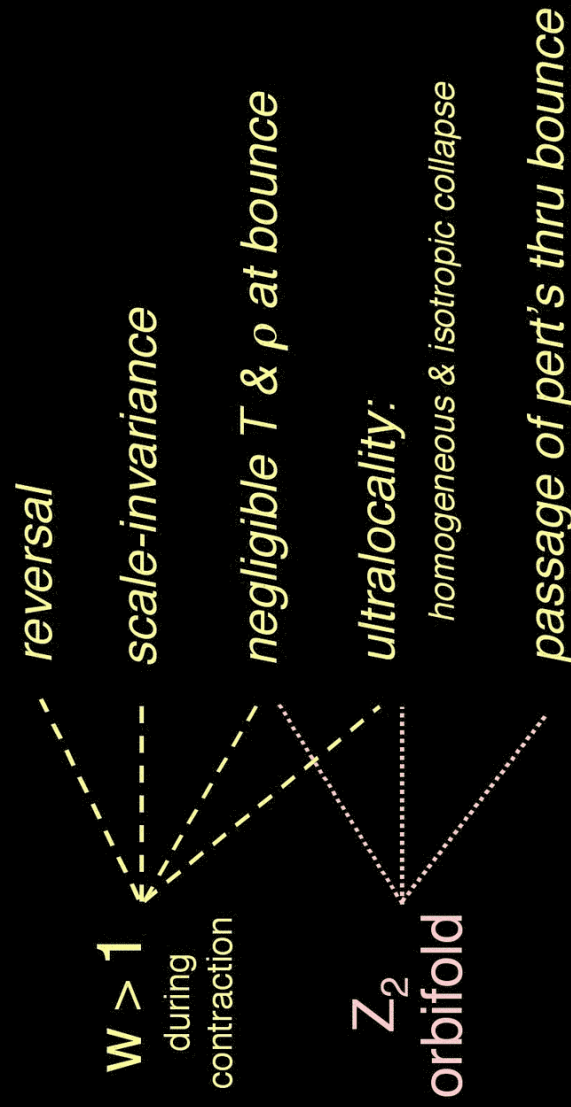
*distinctive predictions ?*

*What I want to talk about today:*

*can there be a bounce ?*

*can perturbations produced  
before the bounce pass thru ?*

## Contraction and Crunch: Building Block Concepts

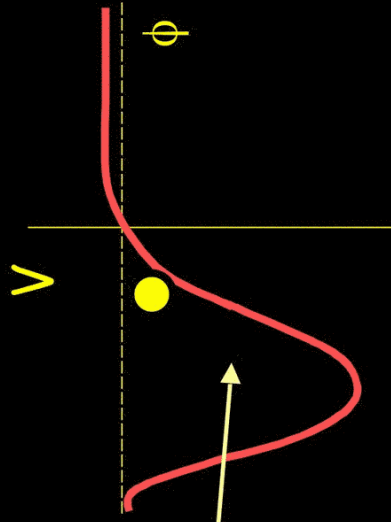


# reversal from expansion to contraction in a flat universe

$$H^2 = \frac{1}{3}(\frac{1}{2}\dot{\phi}^2 + V)$$

reversal  $\rightarrow H = 0 \rightarrow V < 0$

$$w = \frac{\frac{1}{2}\dot{\phi}^2 - V(\phi)}{\frac{1}{2}\dot{\phi}^2 + V(\phi)} > 1$$



# scale-invariance

approach: quantum fluct. exit horizon & re-enter later

$$\epsilon \equiv \frac{3}{2}(1+w)$$

$$a(t) \sim t^{\frac{1}{\epsilon}} \sim (H^{-1})^{\frac{1}{\epsilon}}$$

expanding

$$\epsilon < 1$$

$$n_s \approx 1: \quad \epsilon \ll 1$$

$$n_s - 1 = -2\epsilon + \frac{d \ln \epsilon}{dN}$$

contracting

$$\epsilon > 1$$

$$\epsilon \gg 1$$

$$n_s - 1 = -\frac{2}{\epsilon} - \frac{d \ln \epsilon}{dN}$$

# scale-invariance

approach: quantum fluct. exit horizon & re-enter later

$$\epsilon \equiv \frac{3}{2}(1+w)$$

$$a(t) \sim t^{\frac{1}{\epsilon}} \sim (H^{-1})^{\frac{1}{\epsilon}}$$

expanding

$$\epsilon < 1$$

contracting

$$\epsilon > 1$$

$$n_s \approx 1: \quad \epsilon \ll 1$$

$$\epsilon \gg 1$$

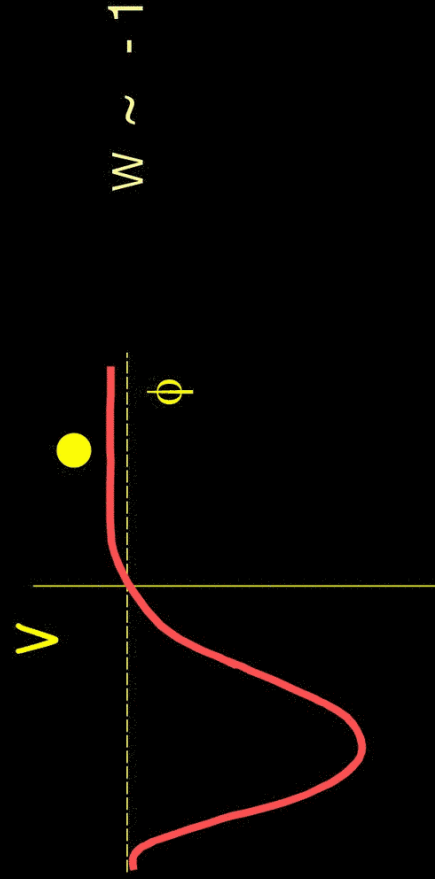
$$n_s - 1 = -2\epsilon + \frac{d \ln \epsilon}{dN}$$

"dual"  
 $\epsilon \rightarrow 1/\epsilon$

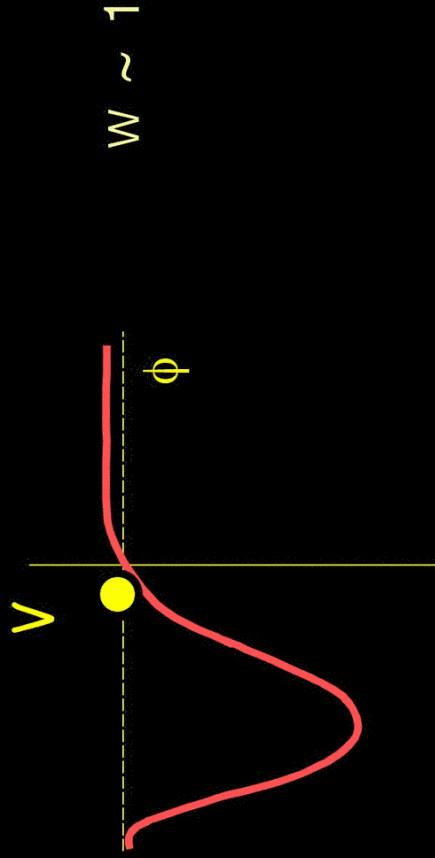
$$n_s - 1 = -\frac{2}{\epsilon} - \frac{d \ln \epsilon}{dN}$$

Now we turn to what happens closer to the crunch

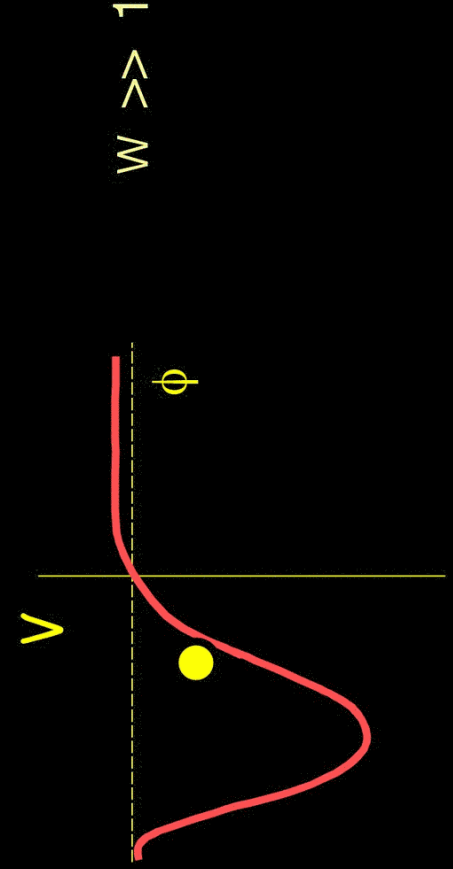
Recall:



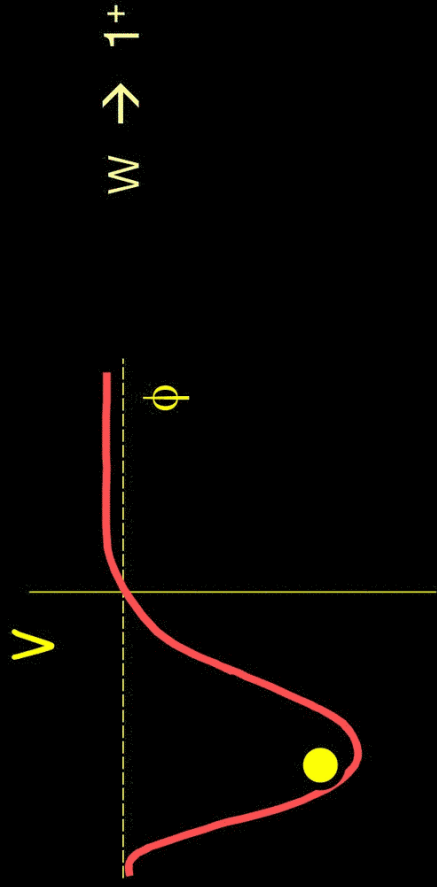
Recall:



Recall:



Recall:



$\rho$  &  $T$  finite

$$S = \int d^4x a^4 \left[ \frac{1}{16\pi G} R - \frac{1}{2} (\partial\phi)^2 + \beta^4(\phi) \rho_M \right]$$

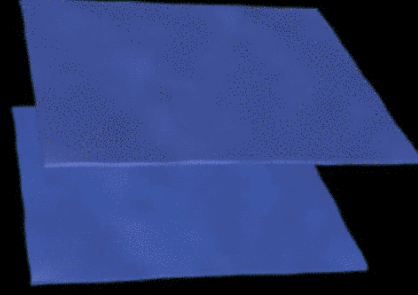
$$a \rightarrow 0 \quad \beta \rightarrow \infty \quad a\beta \rightarrow \text{const.}$$

$$\rho_M \sim 1/(a\beta)^4 \rightarrow \text{const.}$$

... hence, geometry simple

$$a_+ = a \cosh \phi \quad a_- = -a \sinh \phi$$

$$\rightarrow S \sim \int d^4x \left[ \mathcal{L}_+^2 - \mathcal{L}_-^2 \right] \quad a^2 = a_+^2 - a_-^2$$



*ultralocality / “no-hair”*

inflation

$$H^2 = \frac{8\pi G}{3} \frac{\rho_m^0}{a^3} + \frac{8\pi G}{3} \frac{\rho_r^0}{a^4} + \frac{\sigma^2}{a^6} + \dots - \frac{k}{a^2} + \Lambda$$

*ultralocality / “no-hair”*

cyclic

$$H^2 = \frac{8\pi G}{3} \frac{\rho_m^0}{a^3} + \frac{8\pi G}{3} \frac{\rho_r^0}{a^4} + \frac{\sigma^2}{a^6} + \dots - \frac{k}{a^2} + \Lambda$$

$$+ \frac{8\pi G}{3} \frac{\rho_\phi^0}{a^{3(1+w)}} \quad \downarrow \quad w > 1$$

*also suppresses chaotic mixmaster behavior*



*ultralocality / “no-hair”*

*cyclic*

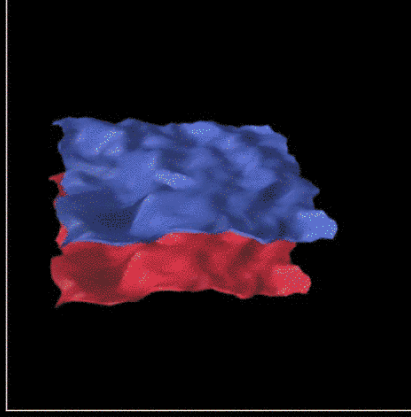
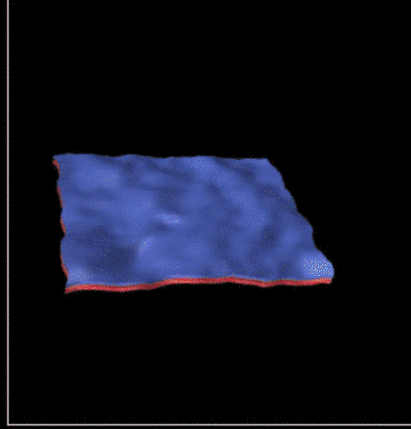
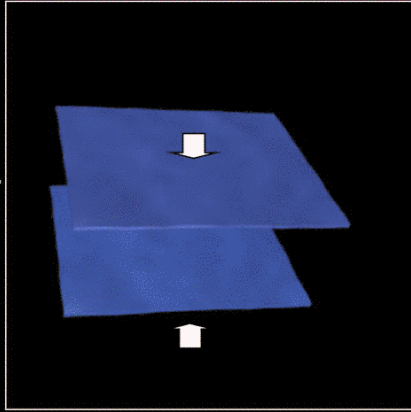
$$H^2 = \frac{8\pi G}{3} \frac{\rho_m^0}{a^3} + \frac{8\pi G}{3} \frac{\rho_r^0}{a^4} + \frac{\sigma^2}{a^6} + \dots - \frac{k}{a^2} + \Lambda$$

$$+ \frac{8\pi G}{3} \frac{\rho_\phi^0}{a^{3(1+w)}} \quad \leftarrow \quad w > 1$$

*Important lesson: quantities diverge...  
but evolution becomes simpler !*

*passage of perturbations  
thru the bounce*

What’s the problem?



4D: contracting

$a \rightarrow 0$

$\beta \rightarrow \infty$

expanding

$$S = \int d^4x a^4 \left[ \frac{1}{16\pi G} R - \frac{1}{2} (\partial\phi)^2 + \beta^4(\phi) \rho_M \right]$$

Newtonian  
potential

$$\Phi_k$$

is nearly scale invariant

Curvature

$$\zeta_k \equiv \frac{1}{\epsilon} \frac{1}{a^2} \left( \frac{\Phi_k}{a'^3} \right)'$$

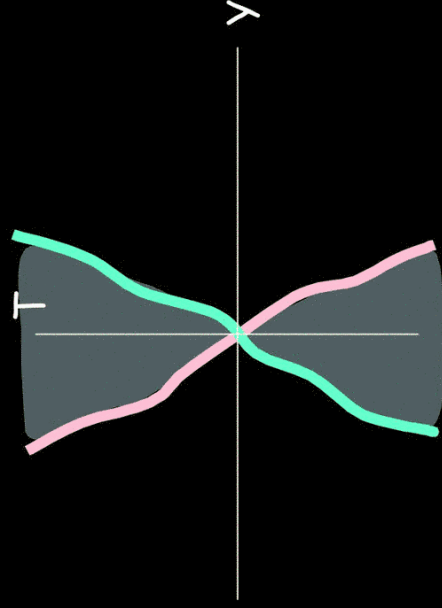
is nearly scale invariant

$\zeta_k$  is conserved for modes outside horizon

blue spectrum !

$\zeta_k$  is not conserved at the bounce

nearly scale invariant spectrum !



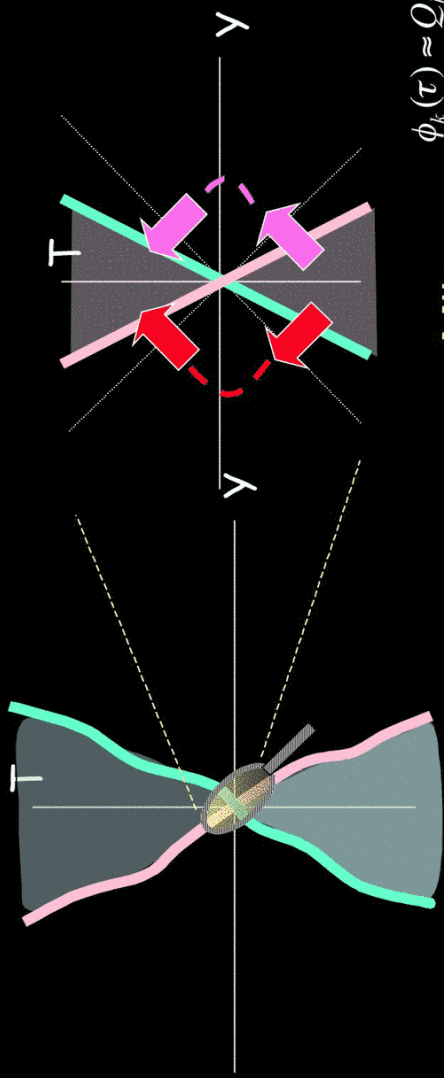
$\zeta_{5d} ??$

$\zeta_{4d} \rightarrow \zeta_{\pm}$

*bounce does not occur along const.- $\zeta$  surface !*

*...hence,  $\zeta$  cannot be continuous at bounce:  
must switch to slicing where bounce is simultaneous ...*

*...small deviation from const.- $\zeta$  implies small amplitude!*



**Milne**

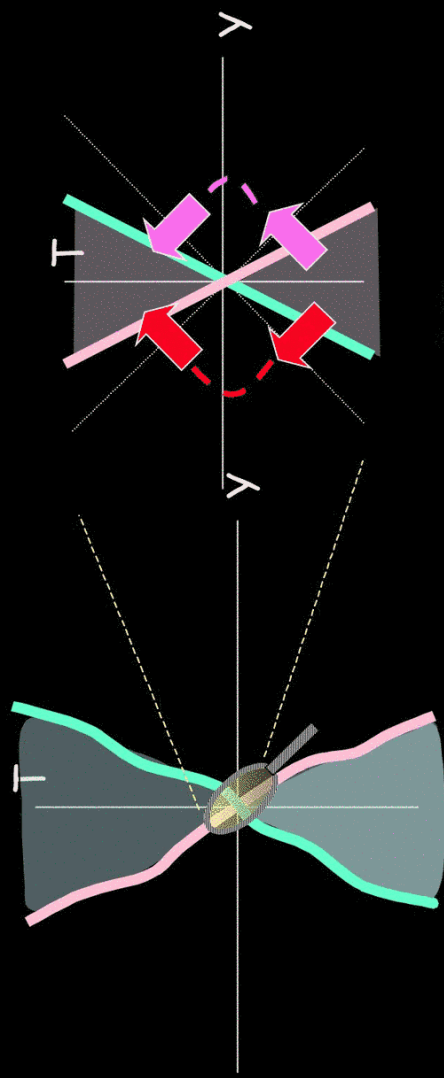
$$\phi_k(\tau) \approx Q_k + P_k \ln |t|$$

$$\rightarrow \phi_k(\tau) \approx -Q_k + P_k \ln |\tau|$$

**unique, unitary, Hadamard, ....**

See also:  
 Berkooz, Craps, Kutasov, Rajesh, hep-th/0212215  
 Craps & Ovrut, hep-th/0308057  
 Cornalba & Costa, hep-th/031009

Tolley &, Turok, hep-th/0204091  
 Tolley, Turok, PJS hep-th/0306109



**Milne**

**unique, unitary, Hadamard, ....**

**tilt unaffected (nearly scale-invariant)**

amplitude depends on  
orbifold/warp/kinematics of 5d!

*Final Speculations:*

## *Beyond the Linear Regime*

$$\text{long wavelengths: } a^2(1 + \Psi) \approx t^{\frac{2}{3}}(1 + \varepsilon \ln t)$$

$$\rightarrow t^{\frac{2}{3} + \varepsilon} \quad (\text{Kasner})$$

*short wavelengths ( $< 10^{-25}$  cm)?*

Papers

and

Mathematica Files

available at

[feynman.princeton.edu/~steinh](http://feynman.princeton.edu/~steinh)