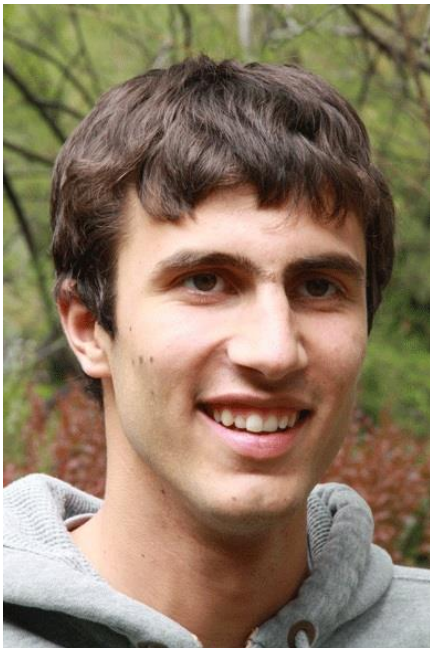




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Microscopic dynamics and failure precursors during the creep of a colloidal gel



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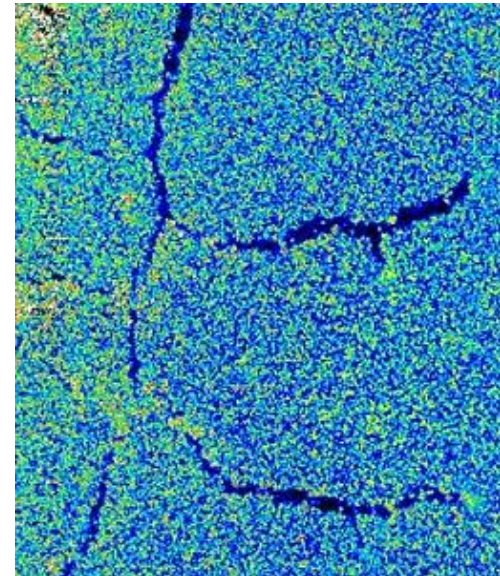
Material failure



12th conference on sinkholes



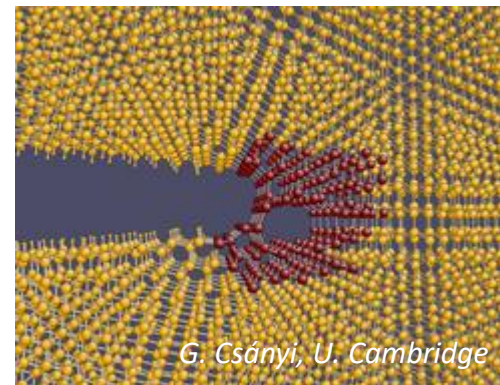
I-35W Mississippi River Bridge



Cracks in cornstarch, L2C



Vasquez Canyon Road, CA

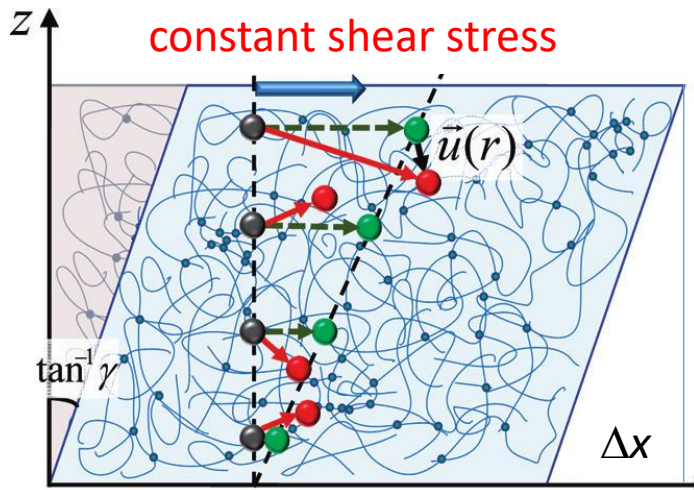


G. Csányi, U. Cambridge

Crack at atomistic level

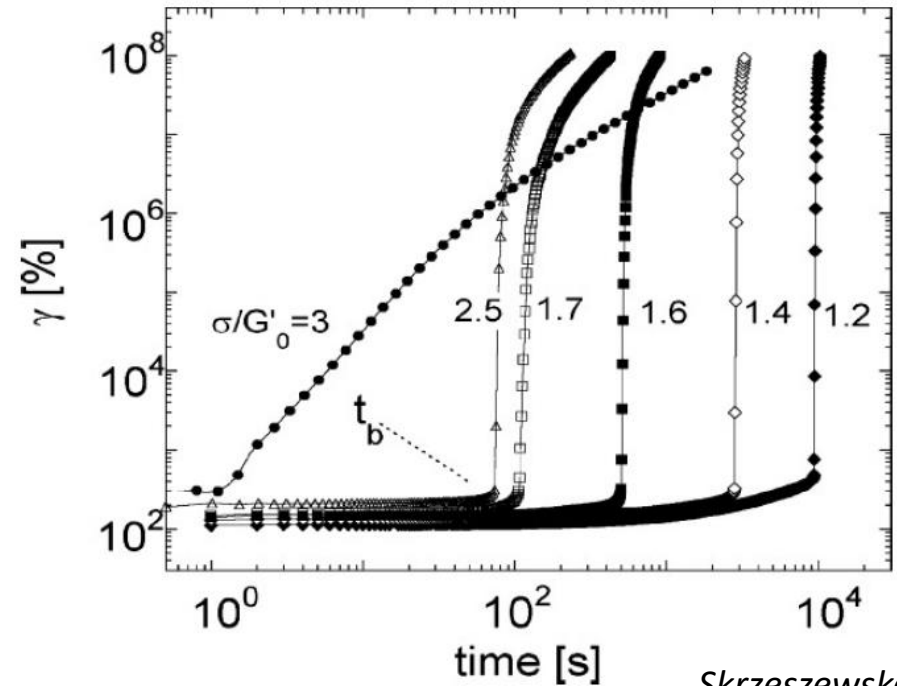
Material failure

- Highly non-linear
- Abrupt
- Unpredictable



Adapted from Basu et al, 2011

Creep test on a polymer network

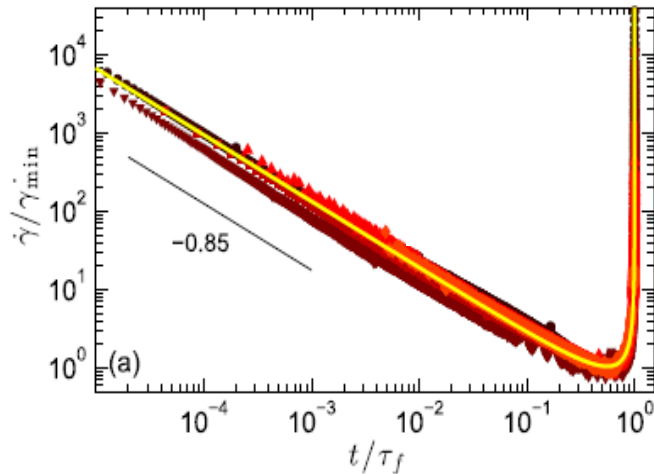
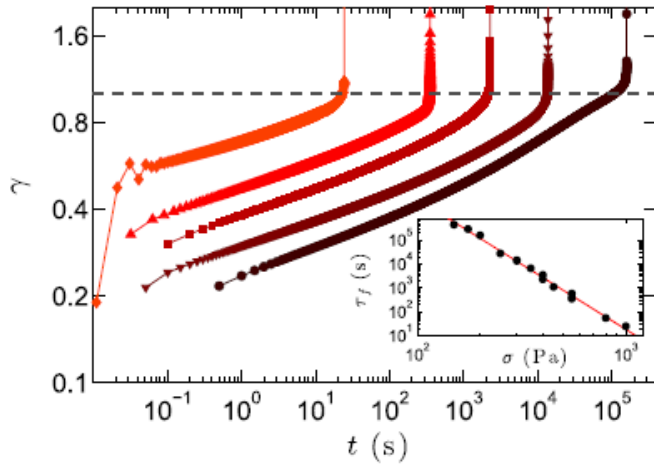


Skrzeszewska
et al., 2010

shear deformation $\gamma = \Delta x/e$

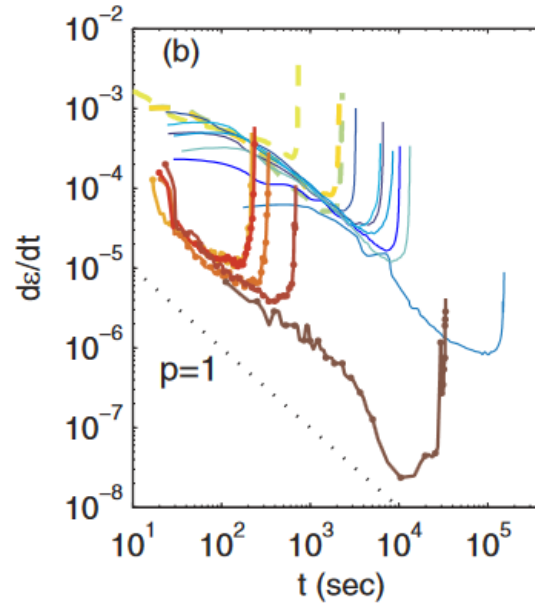
'Delayed' failure ubiquitous

Biopolymer gel under *shear* stress

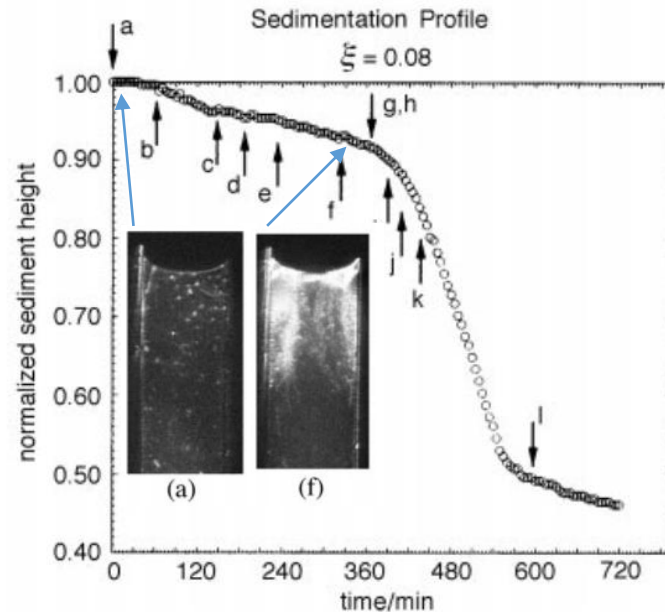


Leocmach et al, 2014

Composite fibers under *tensile* stress



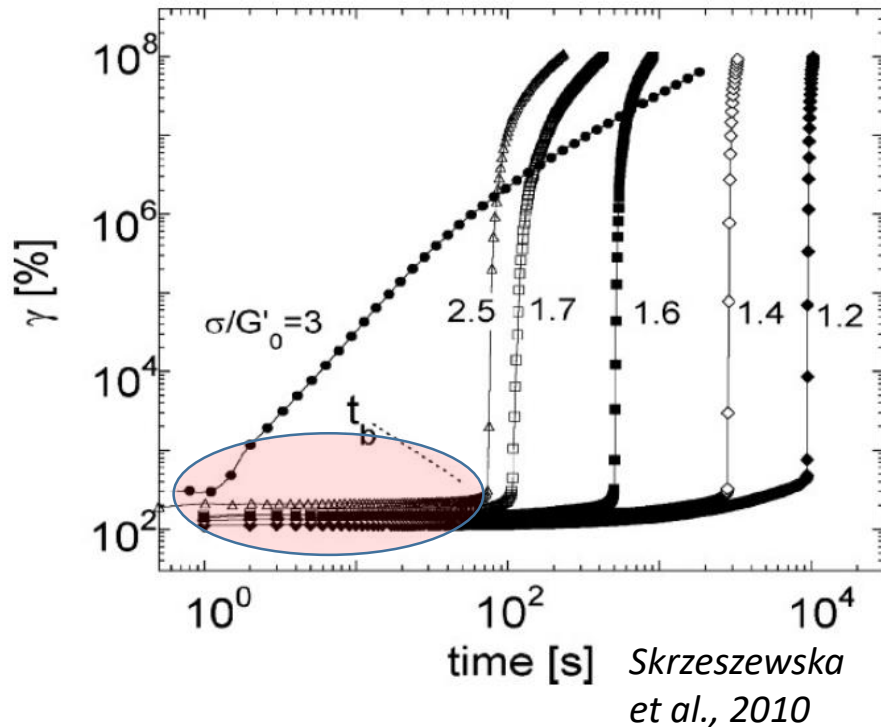
Nechad et al, 2005



Colloidal gel under *compressive* (gravitational) stress

Poon et al, 1999

Key questions



1. Mechanisms at the **microscopic level** leading to catastrophic failure?
2. Relation between **microscopic** evolution and **macroscopic** (mechanical) properties?

Experiments

- Soft matter (tunable, length & time scale accessible)
- Couple rheology and microscopic measurements (structure, dynamics)

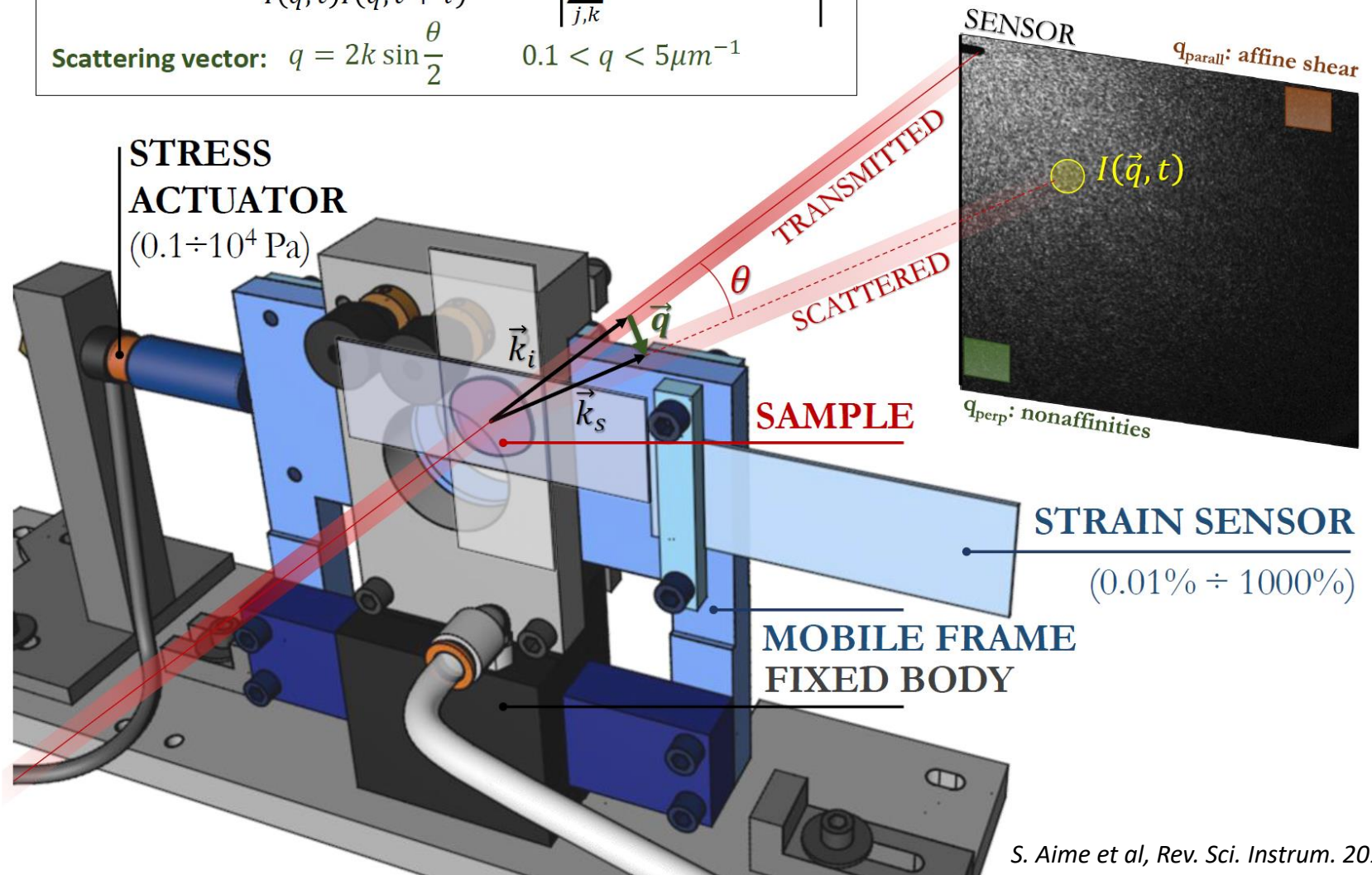
Custom setup: σ controlled shear cell + Small Angle Light Scattering

single scattering

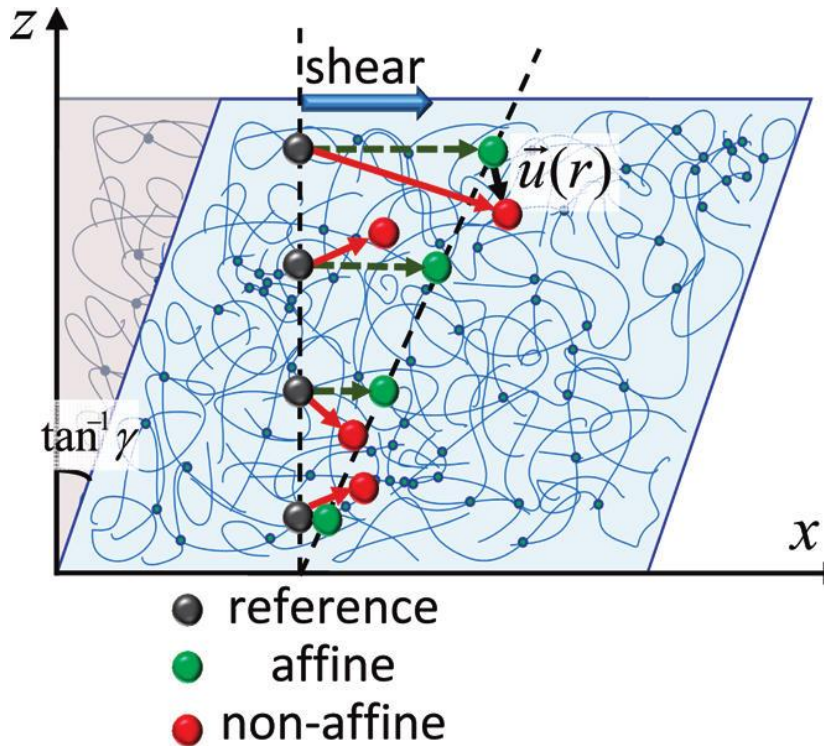
Dynamic light scattering: probing particle displacements

$$g_2(\vec{q}, t, \tau) - 1 = \frac{I(\vec{q}, t + \tau)I(\vec{q}, t)}{I(\vec{q}, t)I(\vec{q}, t + \tau)} - 1 \propto \left| \sum_{j,k} e^{i\vec{q} \cdot [\vec{r}_j(t) - \vec{r}_k(t + \tau)]} \right|^2$$

Scattering vector: $q = 2k \sin \frac{\theta}{2}$ $0.1 < q < 5 \mu\text{m}^{-1}$



Affine vs non-affine particle displacements



Ideal solid: **affine** displacements only

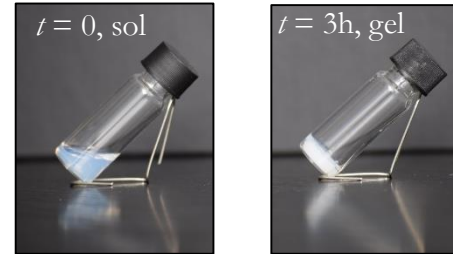
Non-affine displacements:
 G' heterogeneity,
rearrangements...

Basu et al, 2011

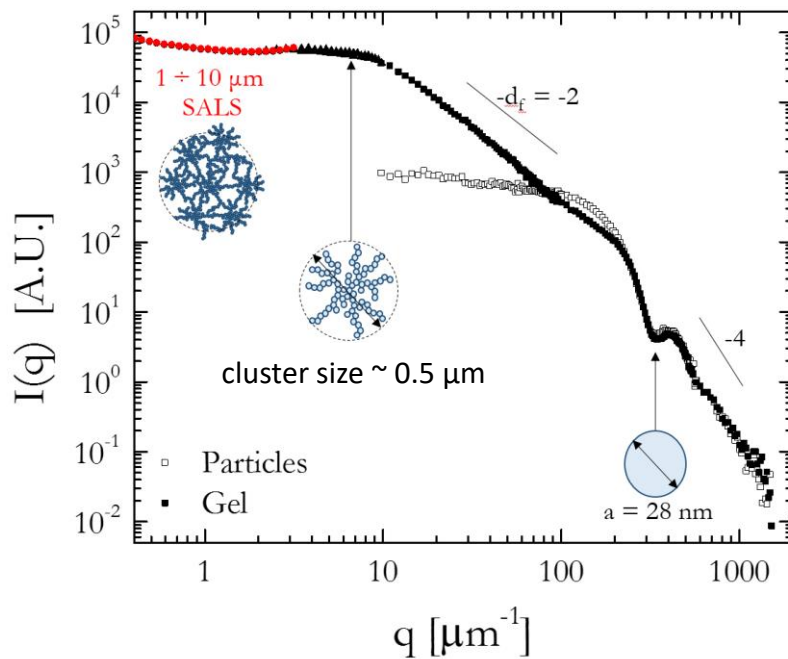
Light scattering: sensitive to $\vec{q} \cdot \vec{\Delta r} \rightarrow$ discriminate **affine**/**non-affine**

Colloidal gel: structure & rheology

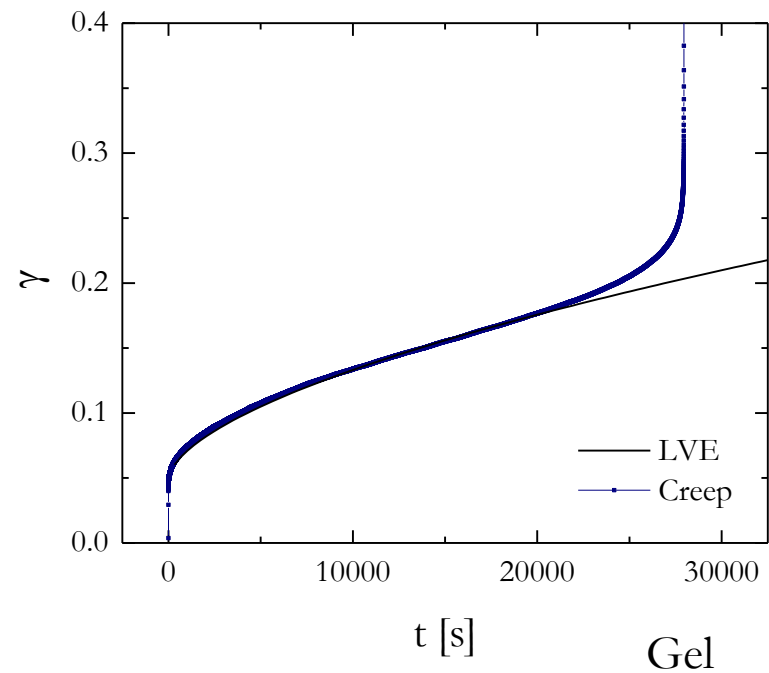
Attractive colloidal particles, diam ~ 30 nm
Diluted gel: $\varphi = 5\%$, $G' \sim 5$ kPa



Sample structure



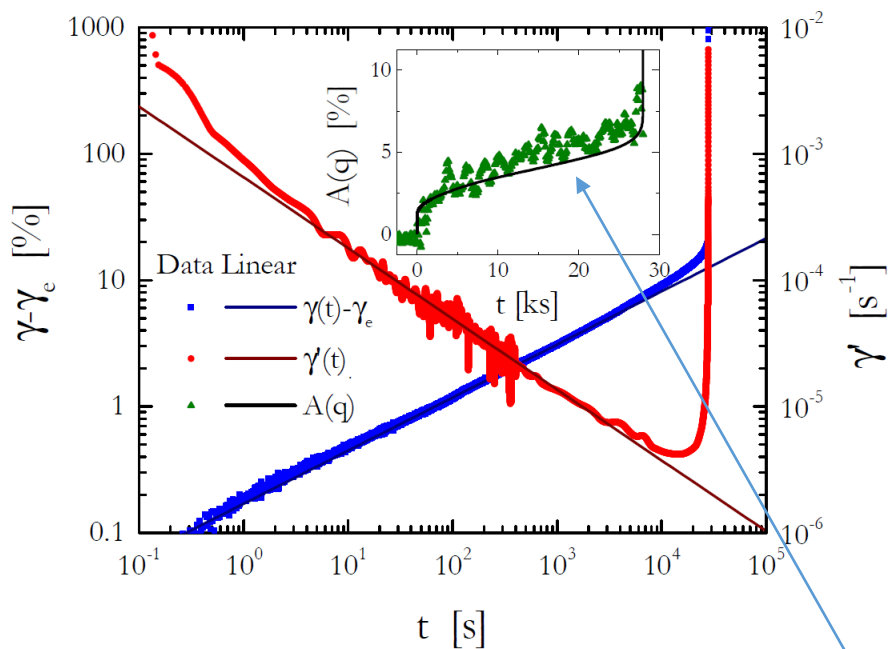
Rheology: creep test



Colloidal gel: structure & rheology

Attractive colloidal particles, diam ~ 30 nm
 Diluted gel: $\phi = 5\%$, $G' \sim 5$ kPa

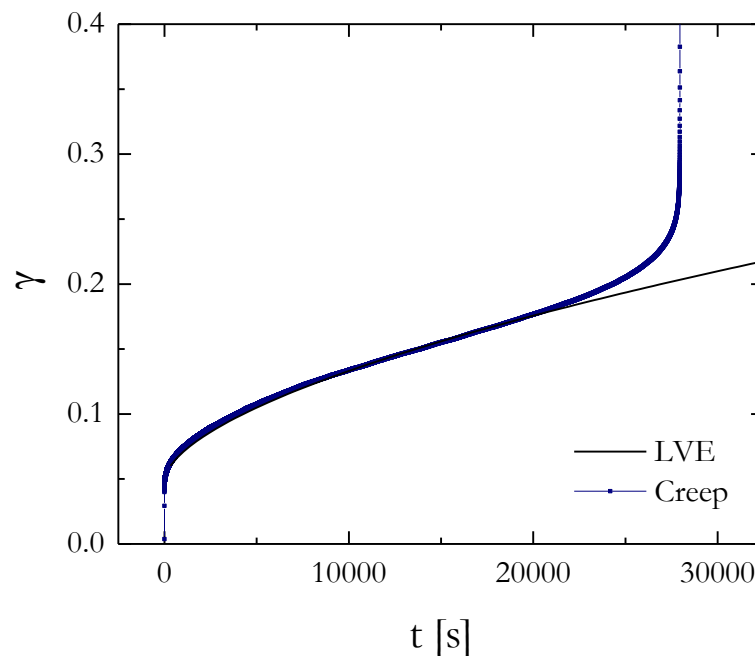
Creep



$$\dot{\gamma} \sim t^{-0.57}$$

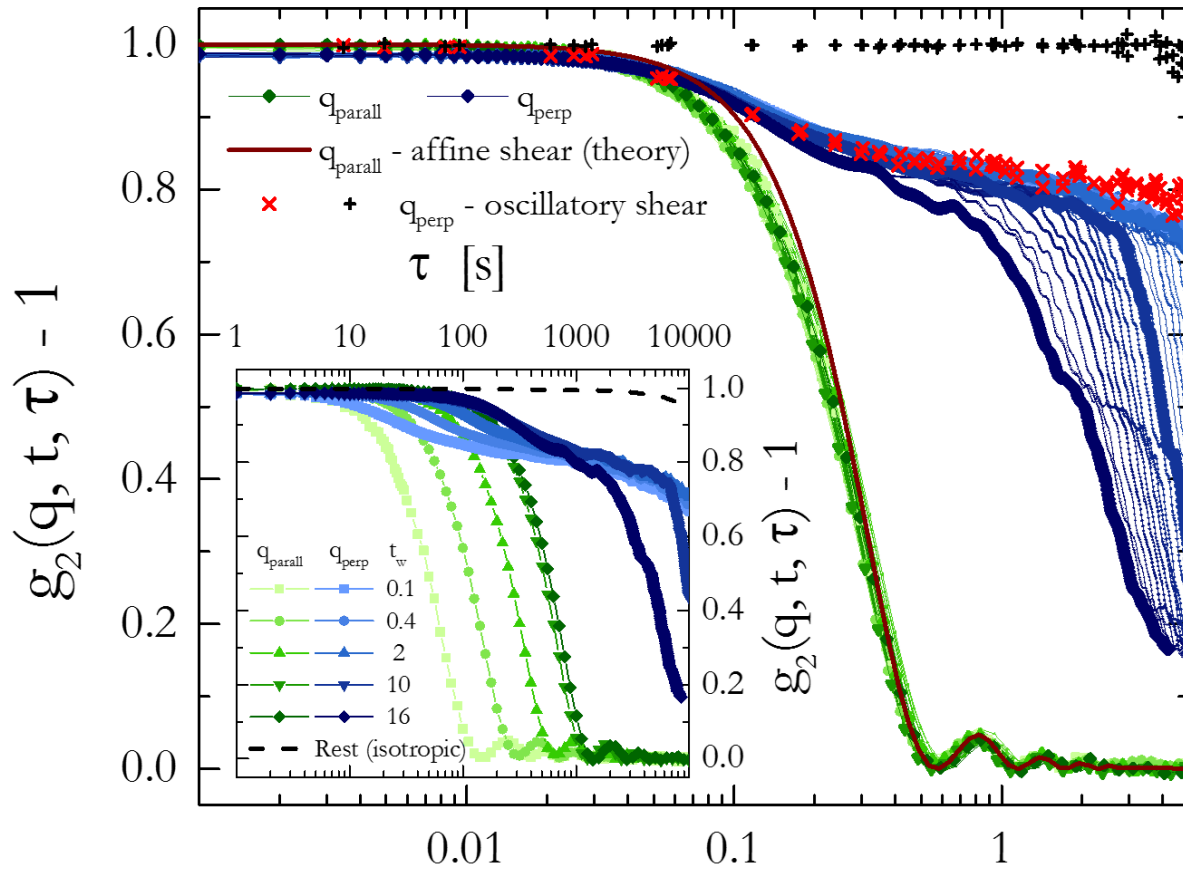
Anisotropy of
 scattered intensity

Rheology: creep test



Microscopic dynamics under creep

Intensity correlation function $\propto \left| \sum_{j,k} e^{i\vec{q} \cdot [\vec{r}_j(t) - \vec{r}_k(t+\tau)]} \right|^2$

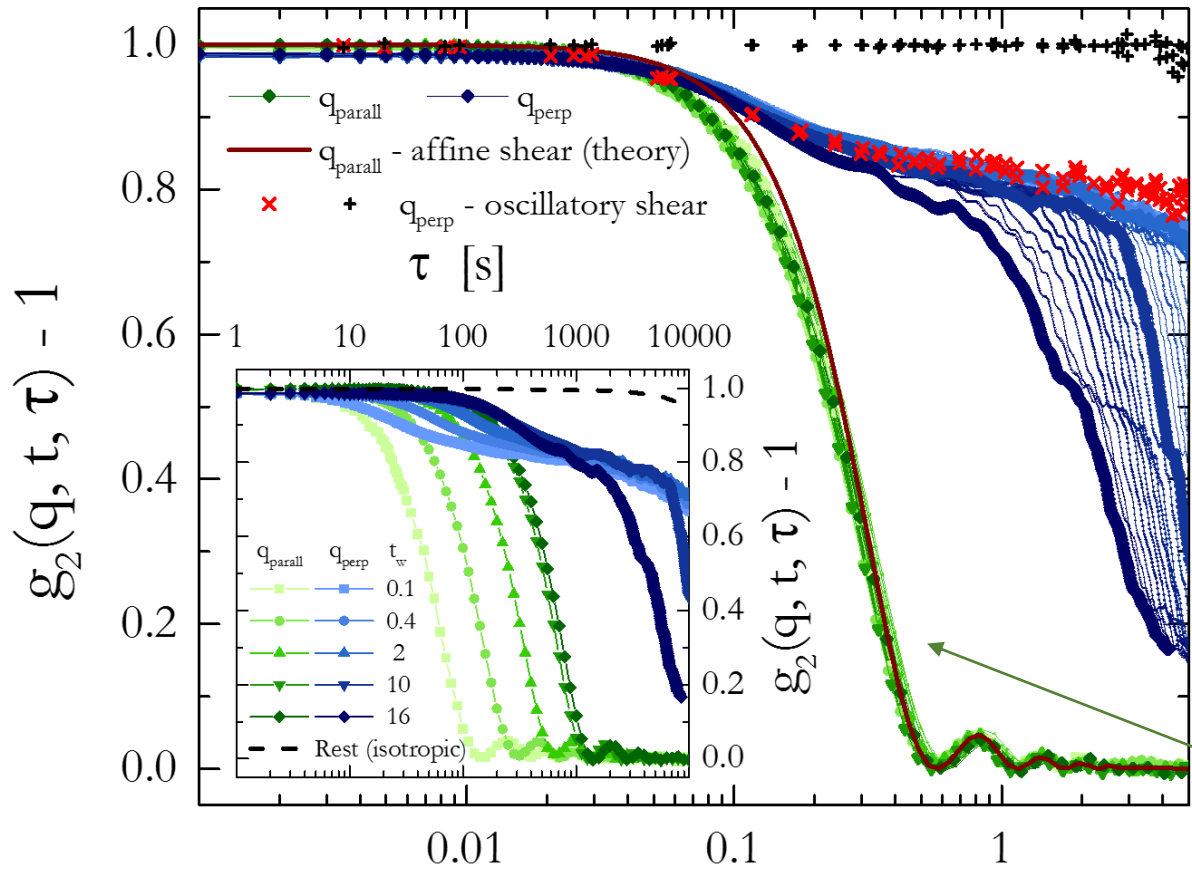


fixed $q = 3.1 \mu\text{m}^{-1}$

$$\Delta\gamma = \gamma(t+\tau) - \gamma(t) \text{ [%]}$$

Microscopic dynamics under creep

Intensity correlation function $\propto \left| \sum_{j,k} e^{i\vec{q} \cdot [\vec{r}_j(t) - \vec{r}_k(t+\tau)]} \right|^2$

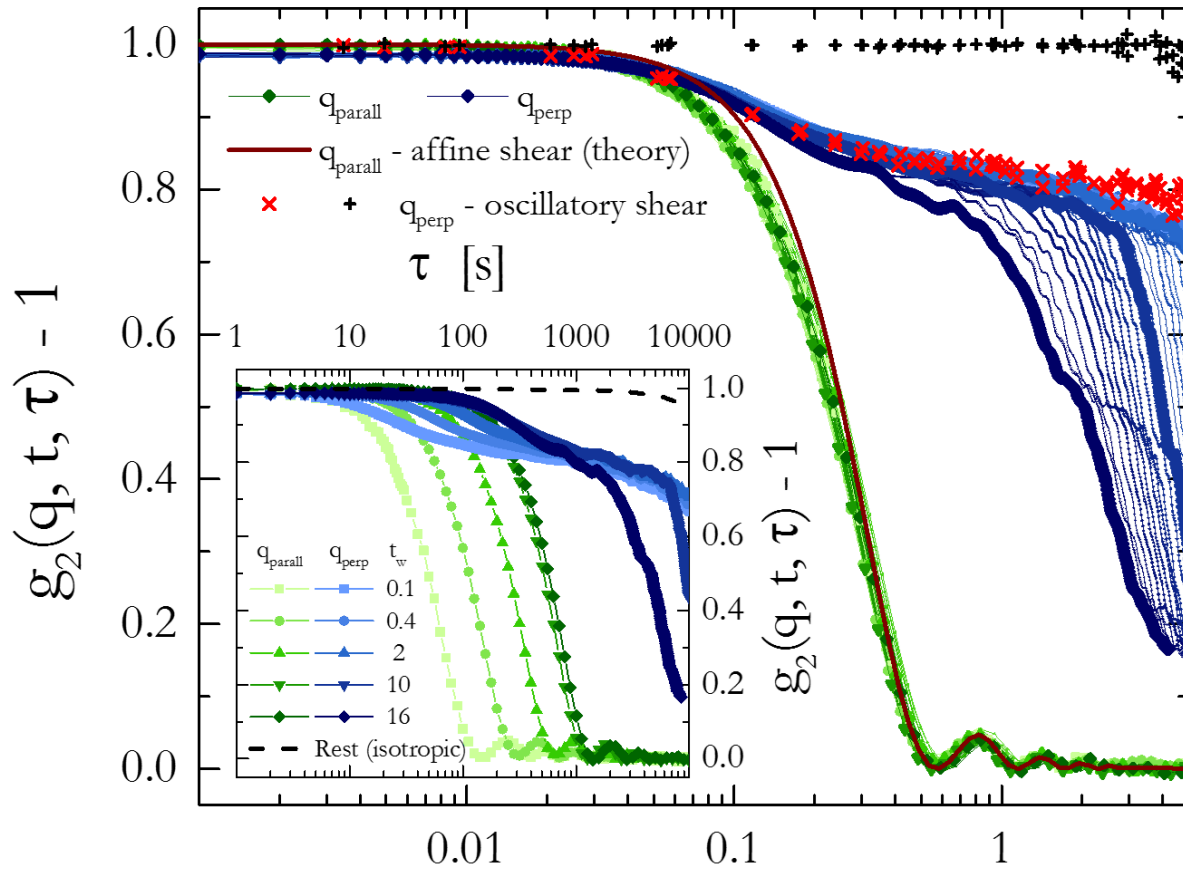


perpendicular to shear direction:
non-affine deformation,
plasticity

parallel to shear direction:
affine deformation

$$\Delta\gamma = \gamma(t+\tau) - \gamma(t) \text{ [}\% \text{]}$$

Initial regime ($\gamma \leq 5\%$): reversible?



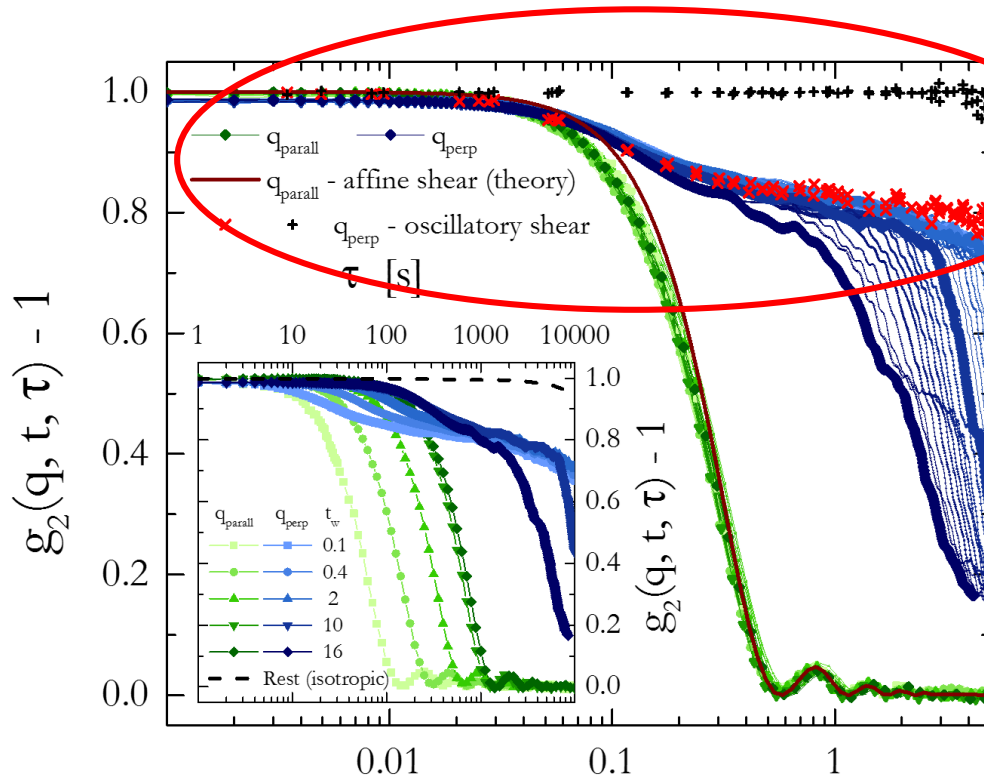
Master curve: **non-affine displacements** independent of strain history



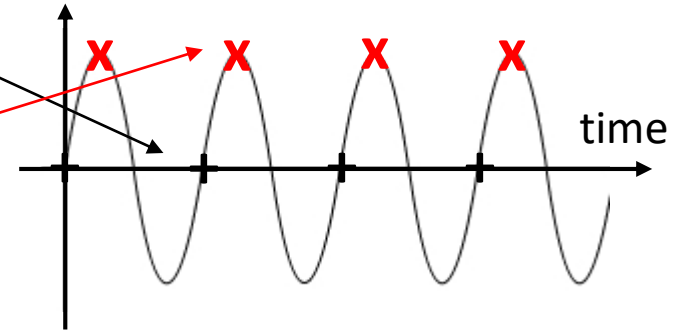
Reversible non-affinities?

$$\Delta\gamma = \gamma(t+\tau) - \gamma(t) \text{ [%]}$$

Initial regime: reversible non-affine elasticity



strain



Oscillatory experiments:
Confirm **reversibility** for $\gamma \leq 5\%$



$$\Delta\gamma = \gamma(t+\tau) - \gamma(t) \text{ [%]}$$

Non-affinity stems from **G'** heterogeneity

Reversible regime: length scale of G' heterogeneity

$$g_2(\gamma) - 1 \approx \exp\left(\frac{-R_{NA}^2(\gamma)}{3} q^2\right)$$

Non-affine msd

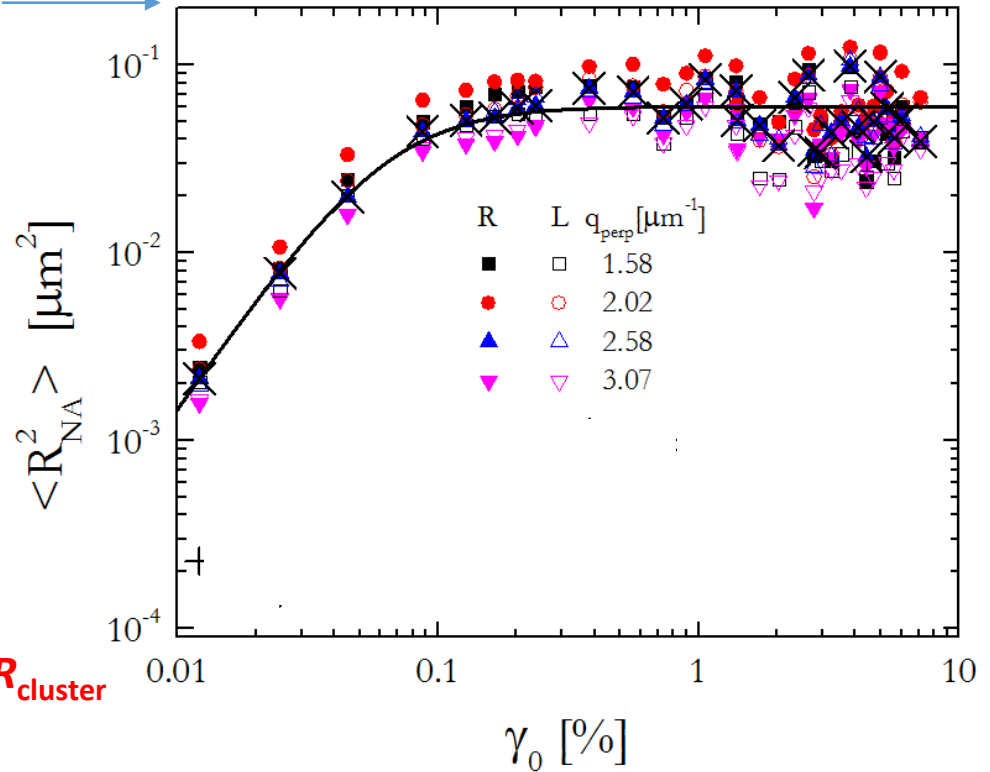
low γ : $R_{NA}^2(\gamma) \sim \gamma^2$ Hook's law
(see also Basu et al. 2011)

R_{NA}^2 saturates close to cluster size



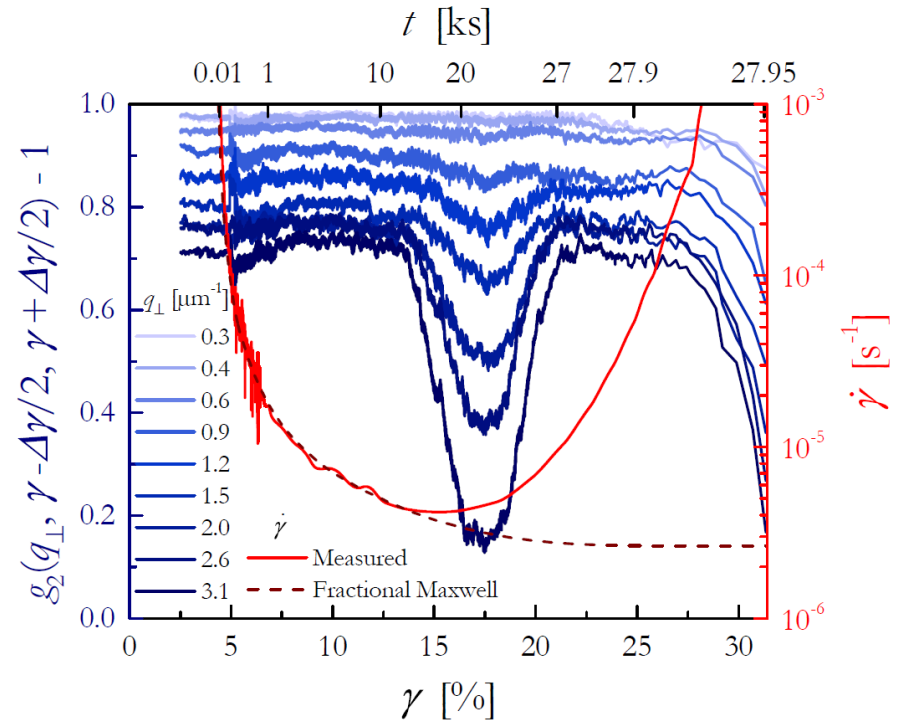
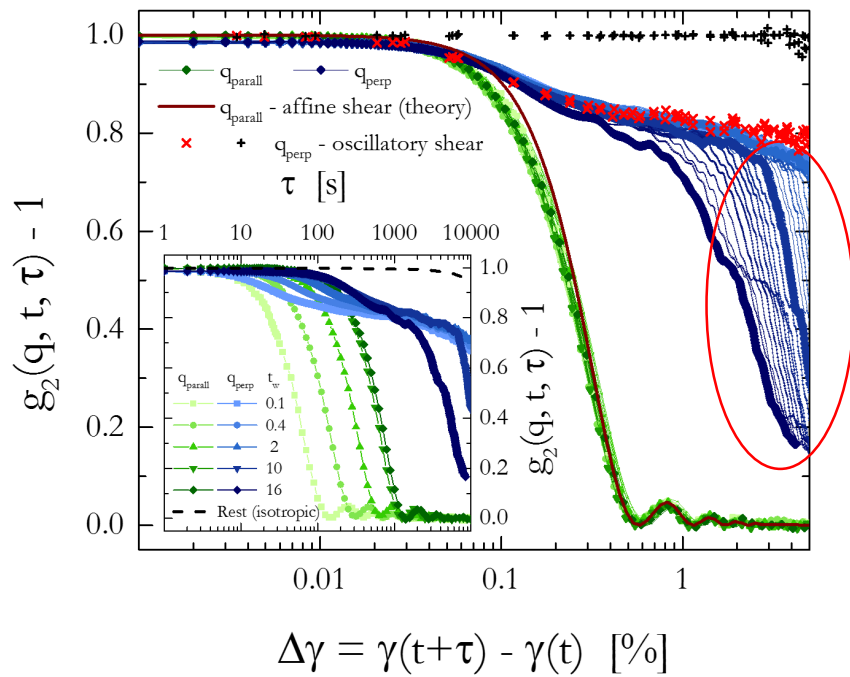
Gel elastically homogeneous beyond $\sim R_{cluster}$

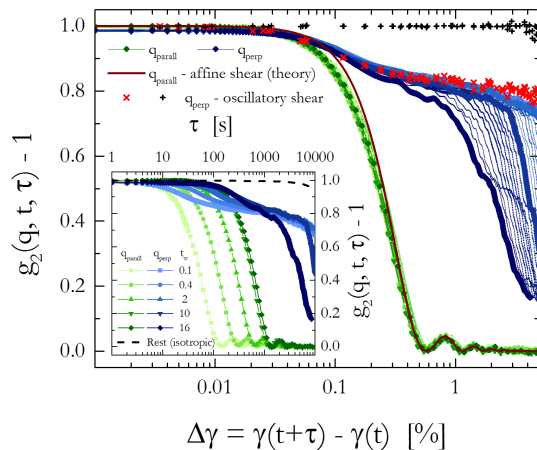
cluster size \rightarrow



$\gamma \geq 5\%$: plastic (irreversible) rearrangements

Dynamics over a fixed strain increment





Plastic activity

Non-affine dynamics:

reversible

plastic

$$g_2(q, \gamma, \Delta\gamma) - 1 = R(q, \Delta\gamma) \times P(q, \gamma, \Delta\gamma)$$

1. Brownian motion: $\exp[-2q^2 D\tau]$

Plastic activity

2. Stationary shear-induced diffusion: $\exp[-q^2(\mathbf{A} \cdot \Delta\gamma)]$

exponent

3. Generalize to non-diffusive dynamics: $\exp[-f(q)(\mathbf{A} \cdot \Delta\gamma)^p]$

scaling factor

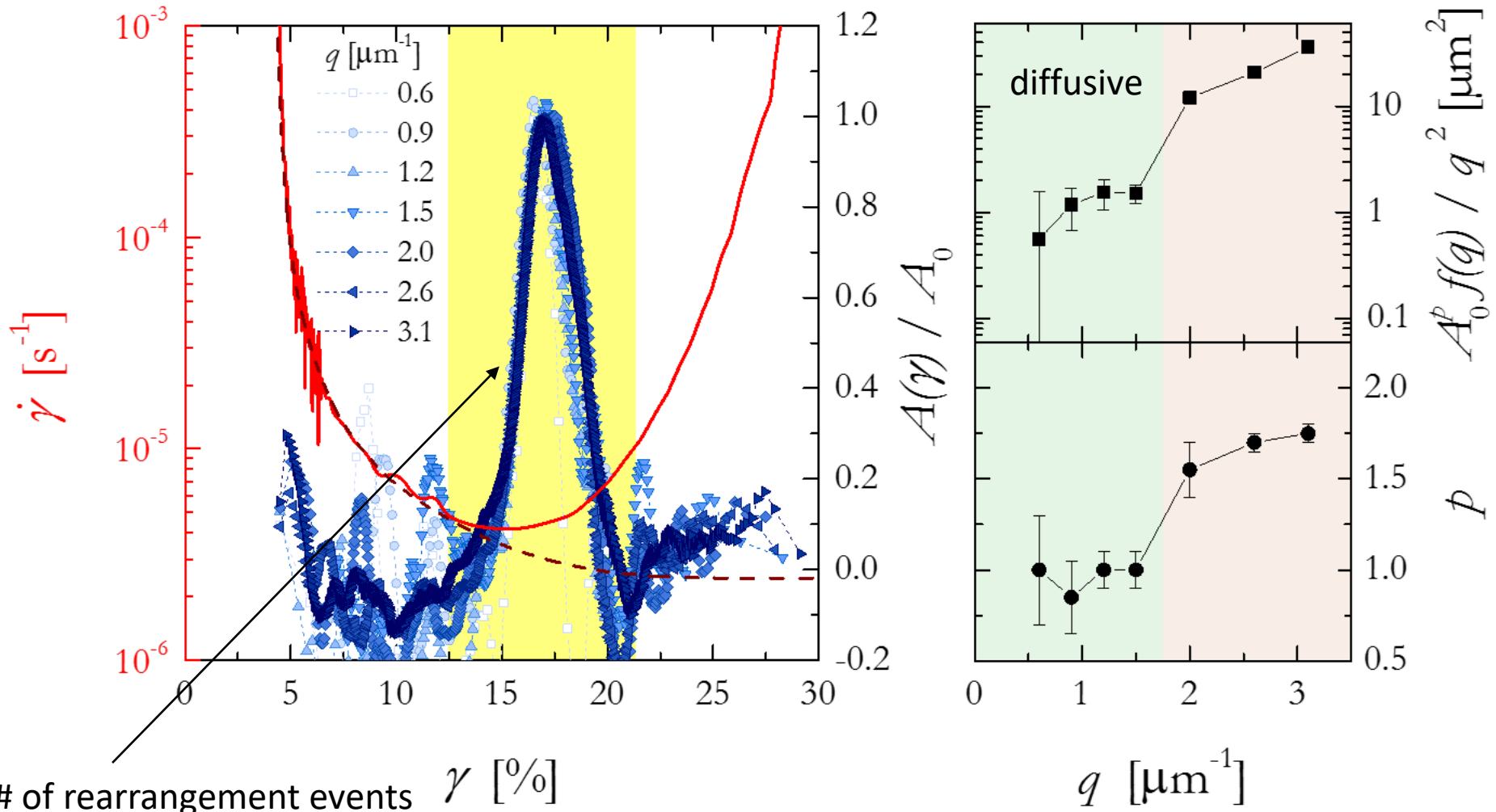
4. Generalization to strain dependent plasticity: $A \mapsto A(\gamma)$

$$A \cdot \Delta\gamma \mapsto \int_{\gamma}^{\gamma+\Delta\gamma} A(\gamma') d\gamma'$$

MODEL:
$$P(q, \gamma, \Delta\gamma) = \exp \left\{ -f(q) \left[\int_{\gamma}^{\gamma+\Delta\gamma} A(\gamma') d\gamma' \right]^p \right\}$$

Failure precursors in the dynamics of a colloidal gel under creep

Burst of plastic activity precedes failure



of rearrangement events per unit volume & unit γ increment

Microscopic dynamics of a colloidal gel under creep: Recap

- Initial regime: creep is (**microscopically!**) **reversible**, non-affinity due to heterogeneity of G' .
- **Dynamic precursor** of failure
- Plastic rearrangements: **diffusive** at low q , additional dynamics at larger q

Thanks to...

C. Ligure, D. Vlassopoulos, T. Divoux, K. Martens



... you all for your attention!