

InAs/GaSb – A New Quantum Spin Hall Insulator

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Rice University

1. “Old Material” for New Physics
2. Quantized Edge Modes
3. Andreev Reflection
4. Summary

KITP Workshop on
Topological Insulator/Superconductor

December 13, 2011

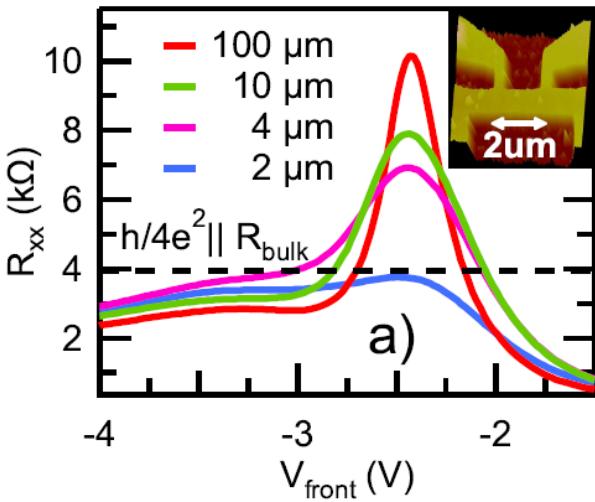
Transport



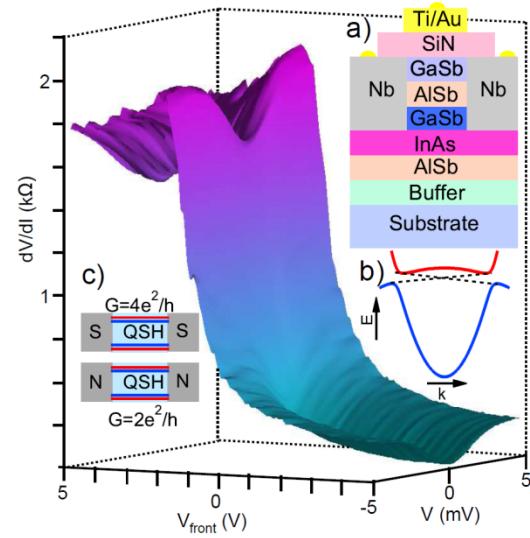
Ivan Knez
R. R. Du
Rice

MBE

Gerry Sullivan
*Teledyne Scientific
and Imaging*



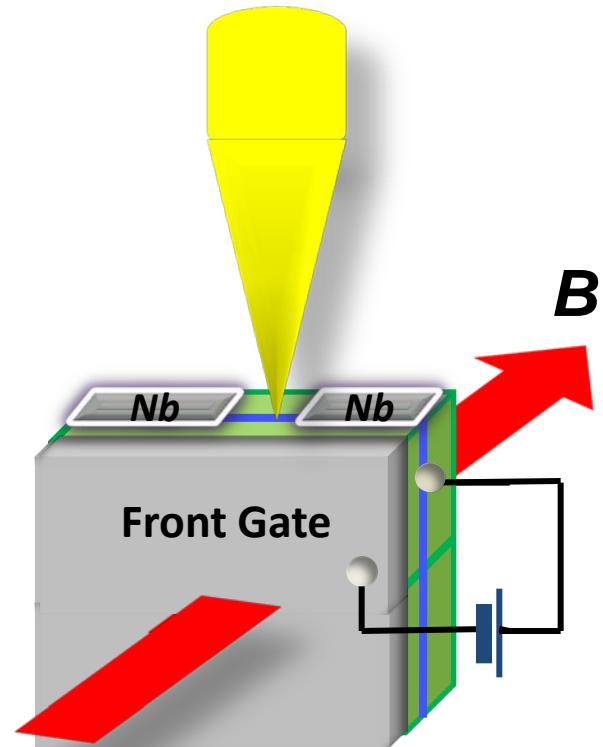
PRL (2011)



ArXiv:1106.5819 (2011)

Acknowledge

S.-C. Zhang, X.-L. Qi,
C.-X. Liu (*Stanford*)
F. Liang (*MIT*)



Majorana Bound State

Quantum Spin Hall Effect in Hg/Te

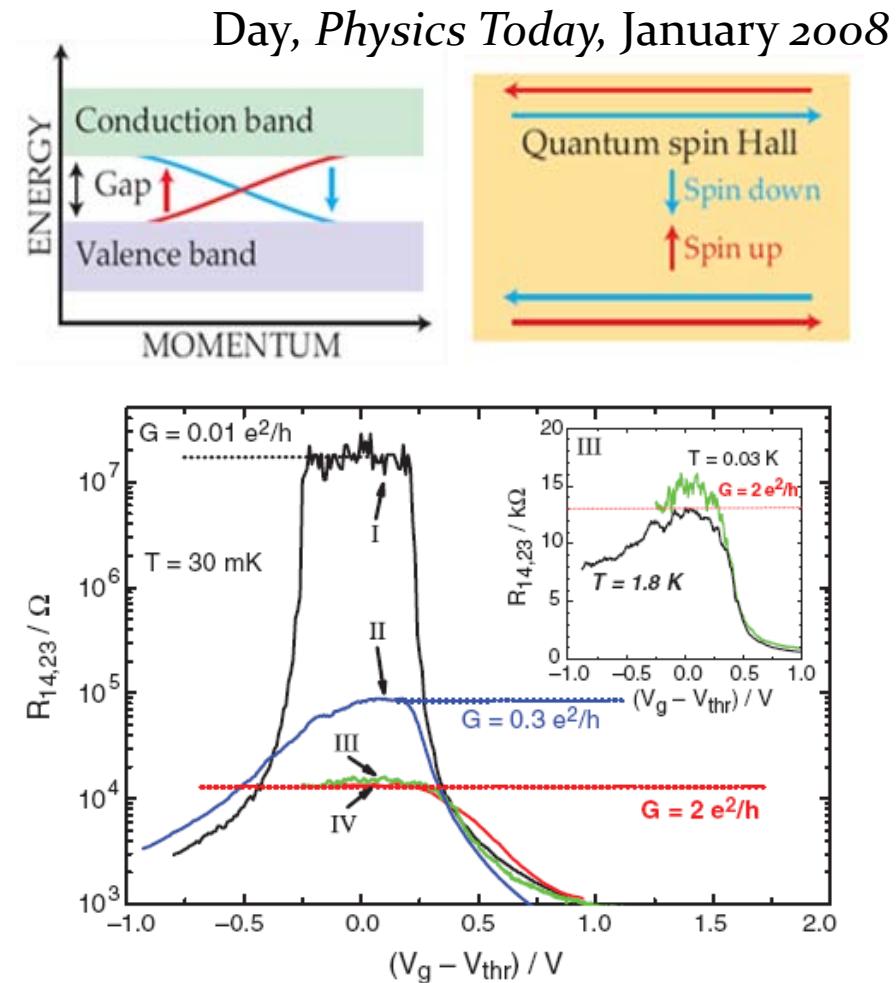
Quantum Spin Hall Effect

* Bulk charge energy gap

* Gapless edge states – odd number of Kramers pairs

* Helical edge states – four terminal conductance of $2e^2/h$

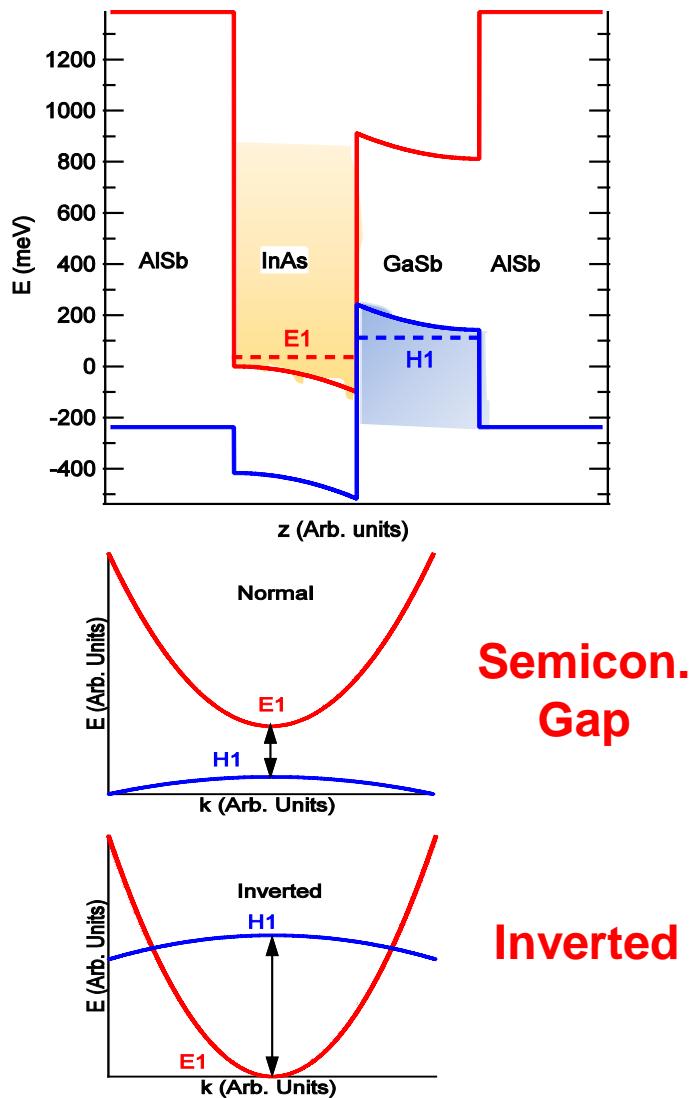
* Observed in HgTe/CdTe QWs in mesoscopic samples



Bernevig et al, *Science*, 2006

Konig et al, *Science*, 2008

Coupled Electron-Hole Double Quantum Well



1. Band parameter

$$\Delta E = E_1 - H_1$$

$Al_xGa_{1-x}Sb$

can be tuned by width, x, gates

2. $\Delta E < 0$ inverted -0.15 eV to 0

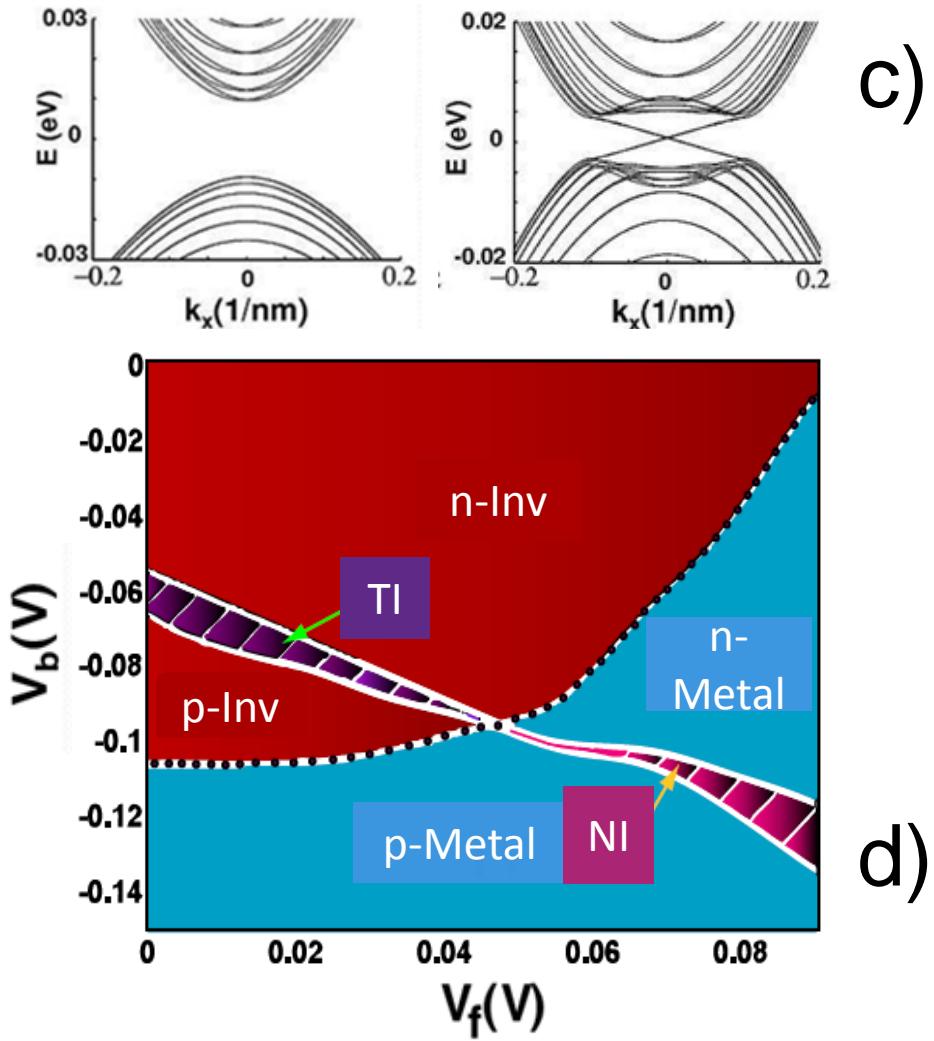
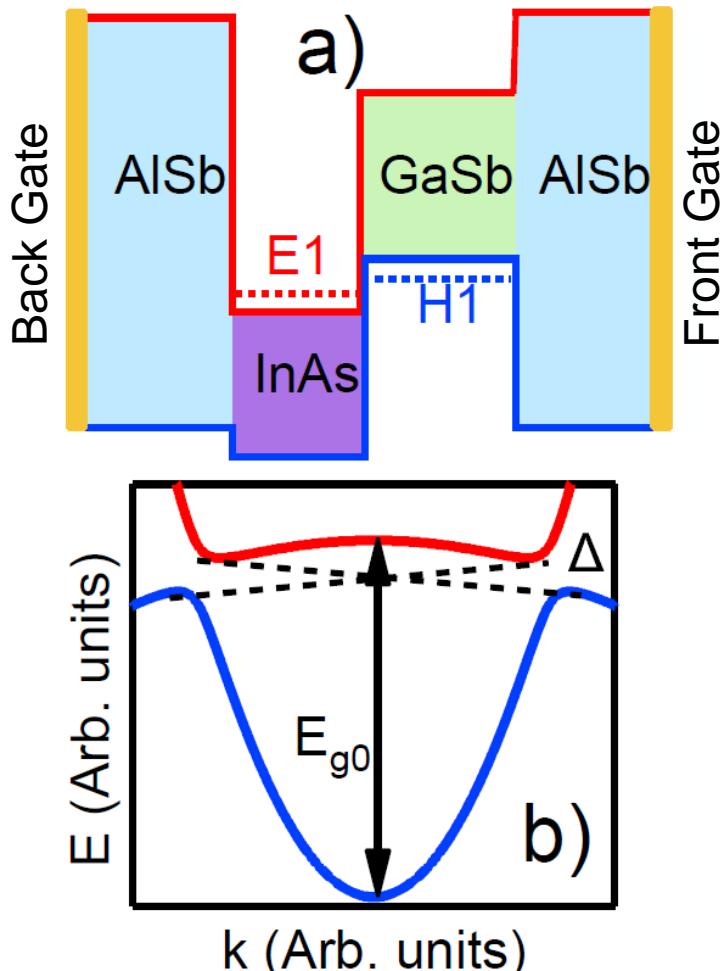
Tunneling mix e-h and opens a gap at finite k-vector

3. Used as infrared detector and night vision CCD

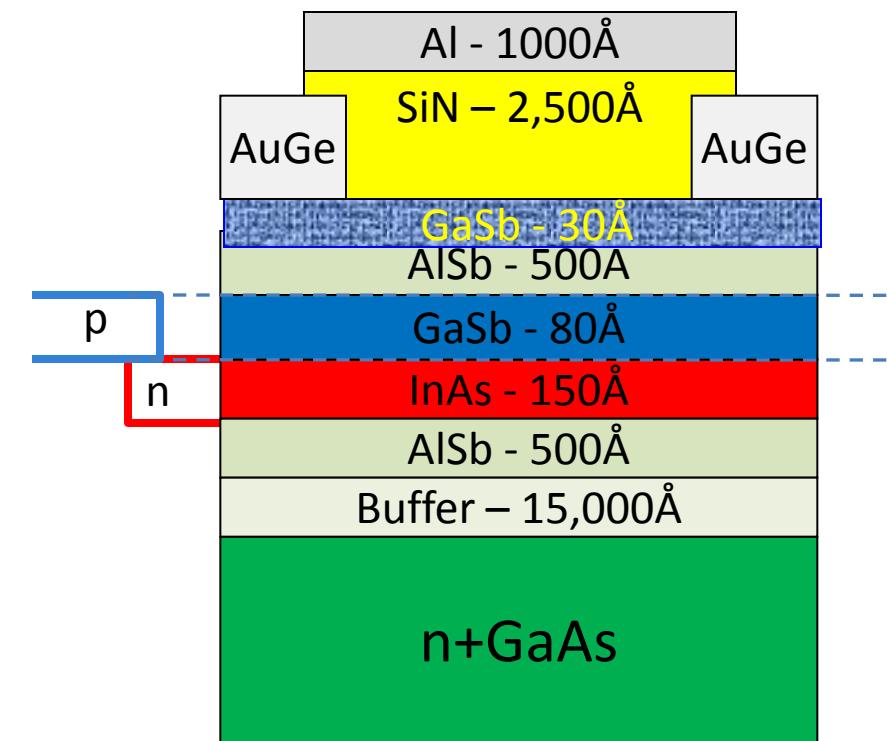
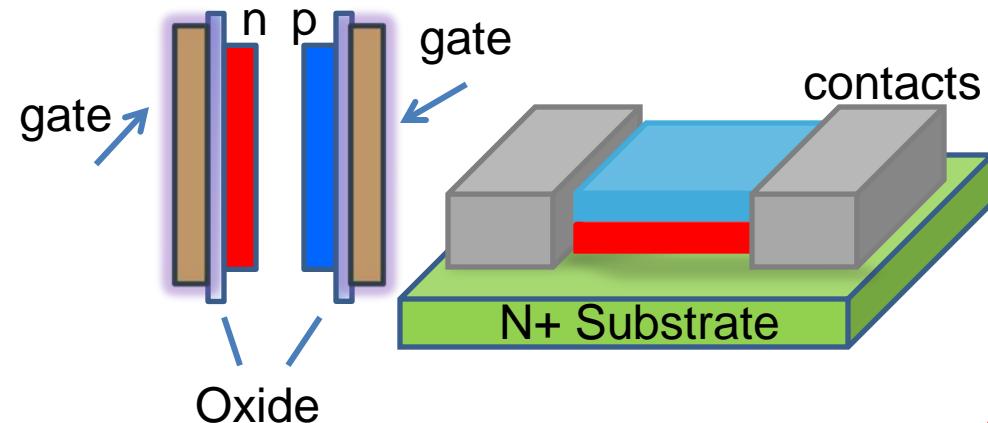
4. High mobility MBE wafers produced in best industrial labs ($1\sim 5 \times 10^5 \text{ cm}^2/\text{Vs}$)

InAs/GaSb as a QSH Insulator

Liu *et al*, PRL 100, 236601 (2008)



Band-Gap Engineering



	Wafer A	B	C
GaSb	8 nm	5nm	8nm
InAs	15nm	12.5 nm	8nm
$\Delta E =$ $E_f - E_1$	-33 meV	-18 mev	~ 0

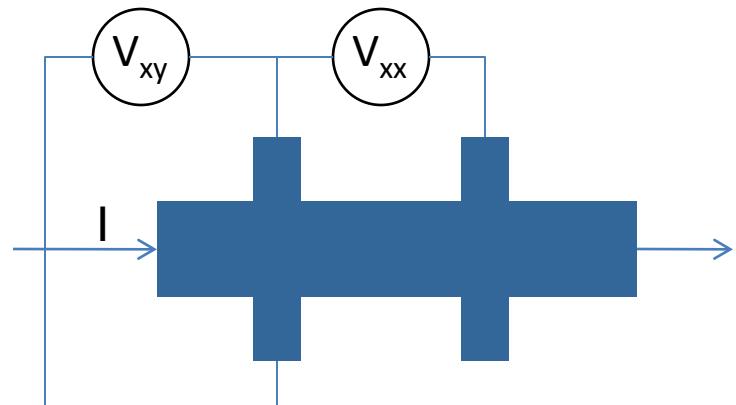
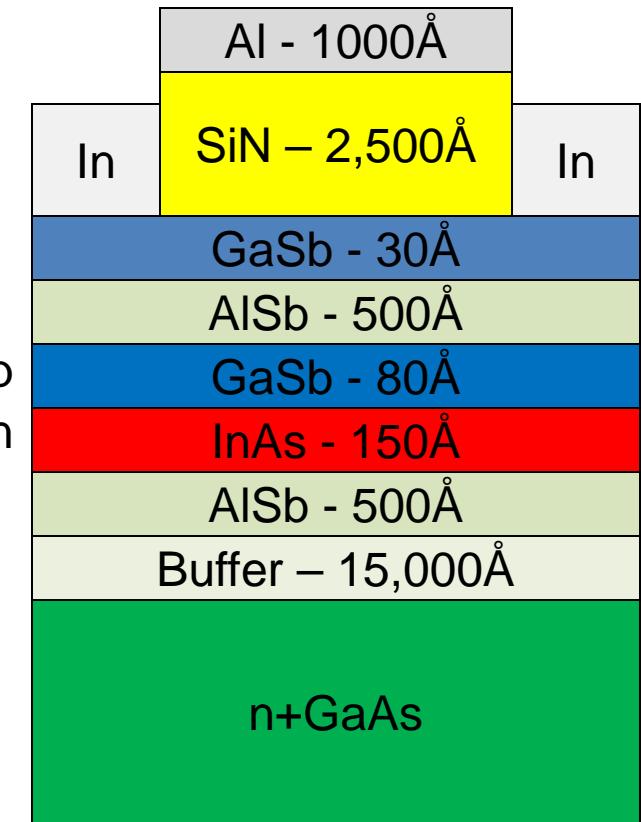
E_f , ΔE can be readily tuned by gates →

Simiconductor-TI transition
QSHE
InAs/GaSb/InAs/GaSb... Superlattice
- Helical surface states?
Re-entrant Quantum Hall Effect
Bound Excitons - BEC

Experimental details

- **Wafer A:** 150Å InAs and 80Å GaSb composite QW – inverted regime :
 $E_c - E_v \approx -33\text{meV}$
- **Hall bar**
 - Sample a: 10 μm by 20μm
 - Sample b: 0.5 μm by 1.5μm
- **Frontgate** – Ti/Al or Ti/Au on SiN
Backgate – n+GaAs substrate
- **Transport**

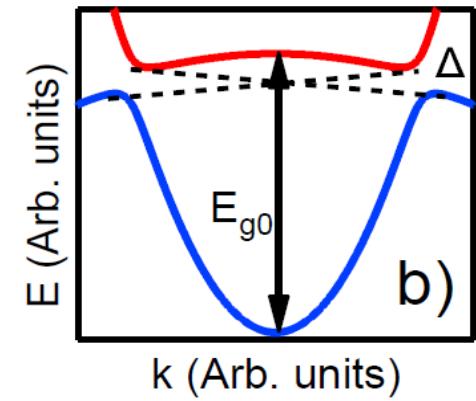
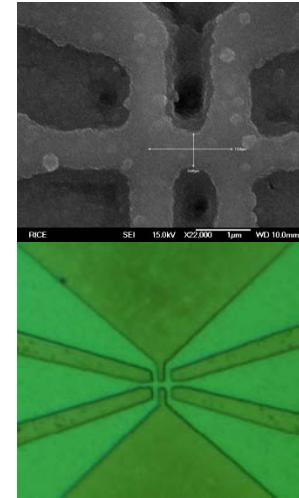
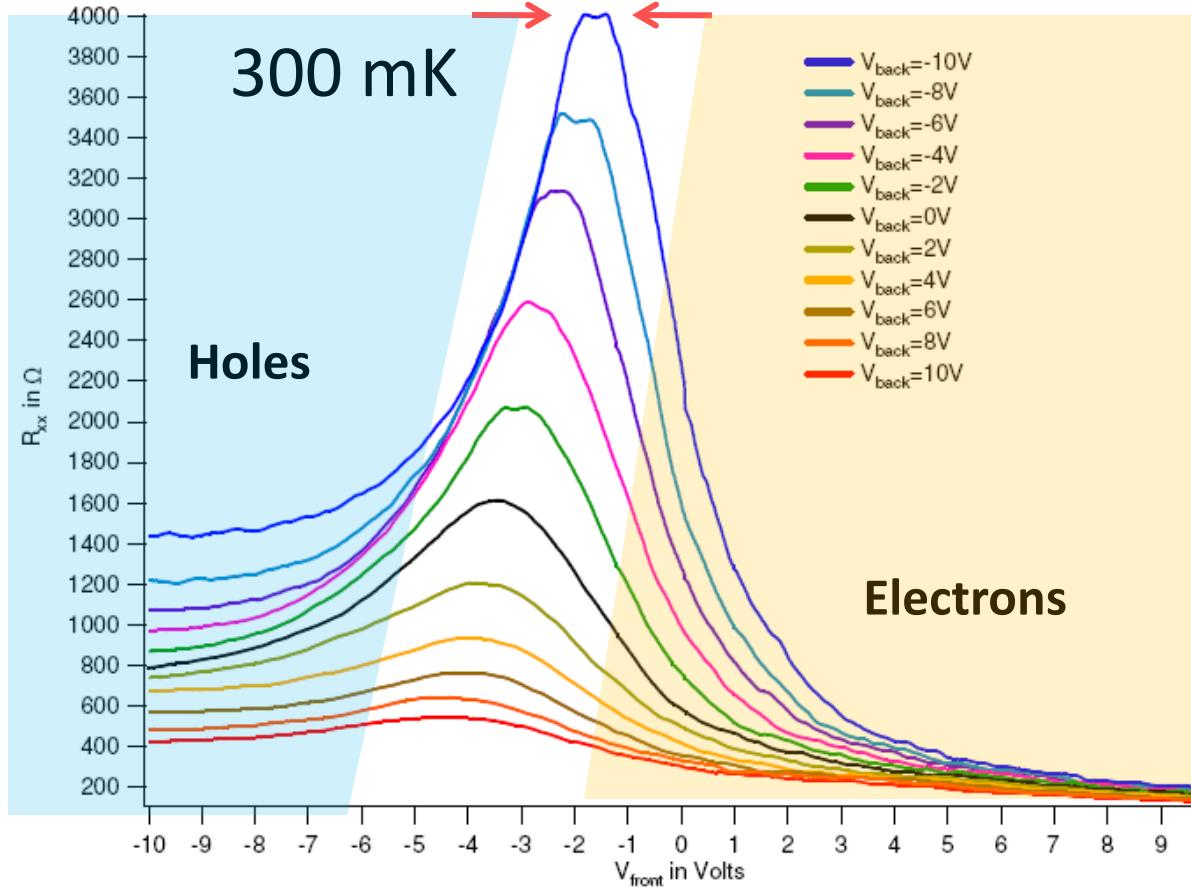
T=20mK, B=18 Tesla, NHMFL
T=0.3K, B= 4/2 Tesla, Rice



1 Tuning Through Mini Gap in Meso Samples

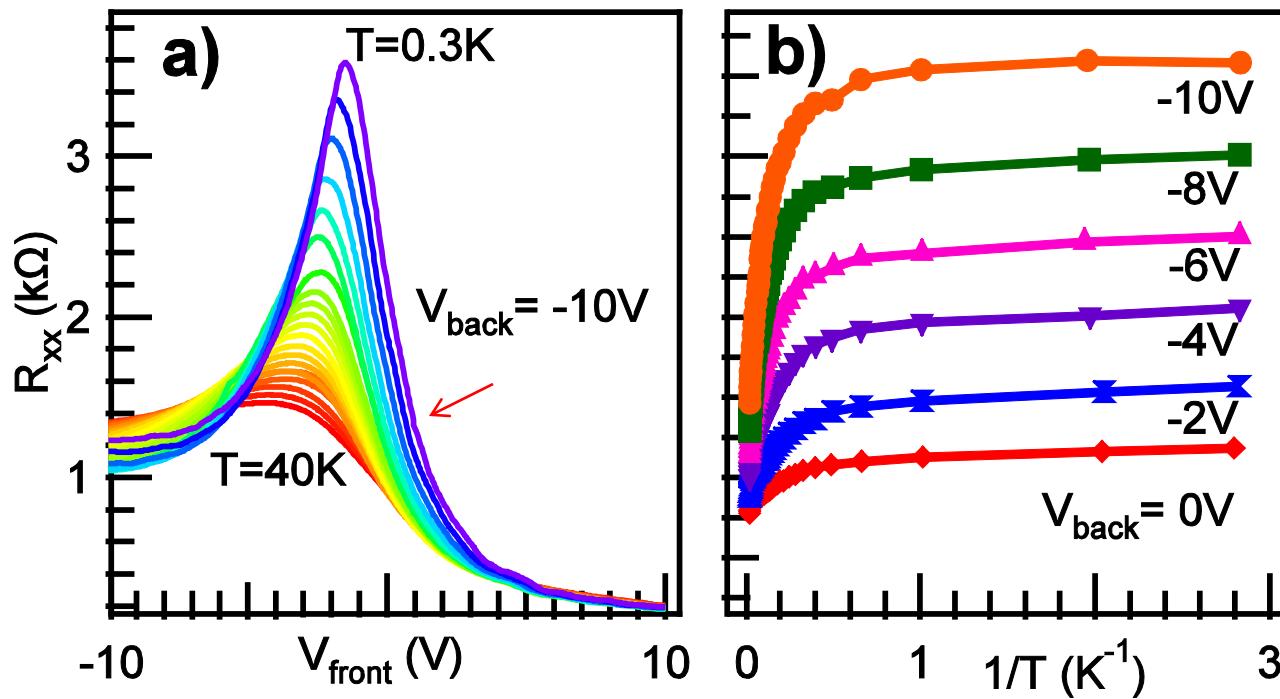
12.9 k ?

(Wafer A 1.5 by 0.5um Hall bar)



Temperature dependence

Knez *et al*, Phys. Rev. B 81, 201301(R) (2010)



- E_f sweeps the minigap
- Gap energy $\sim 40\text{ K}$
- Bulk conductivity $g_0 > e^2/h$ (for wafer A)
- Band-gap engineering of g_0

Understanding Residual Conductivity

Naveh & Laikhtman,

Euro. Phys. Lett. 55, 545(2001) Hebrew U

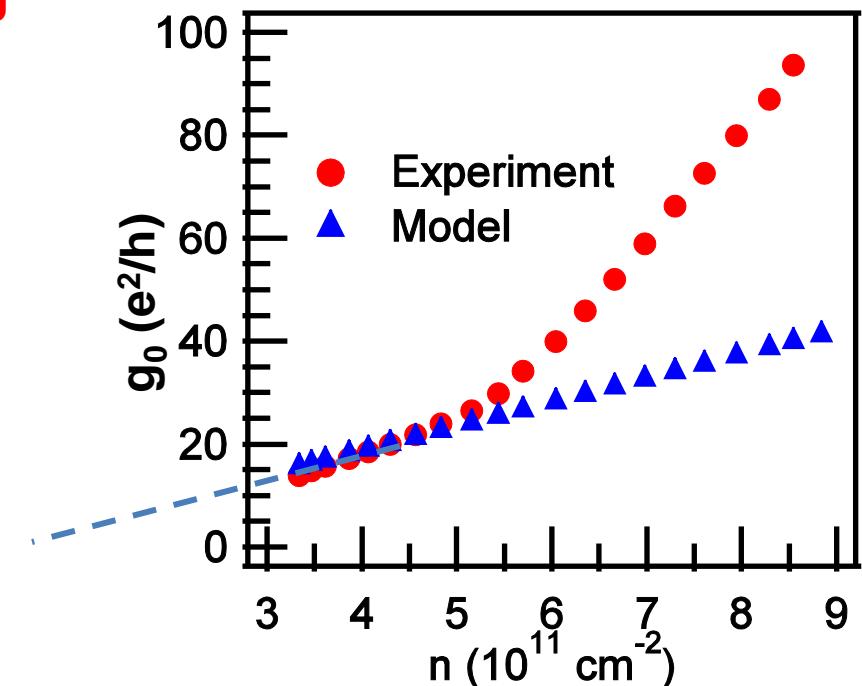
- Hybridized e-h not necessary charge neutral
- Hybridized e-h break due to impurities
- Relevant time scale: $\tau_t = \frac{\hbar}{2\Delta}$
- Number of scattered carriers:

$$n_0 \cdot \left(1 - e^{-\tau_t / \tau_r}\right) = n_0 \cdot \left(1 - e^{-\Gamma / \Delta}\right)$$

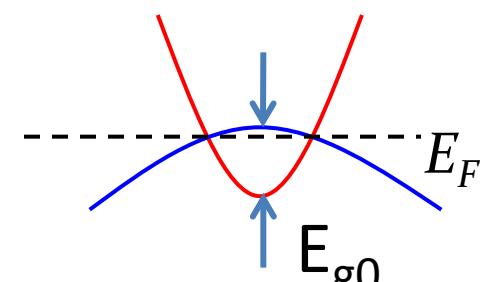
- Mobility: $\mu = \frac{e}{m} \tau = \frac{e}{m} \frac{\hbar}{2\Gamma}$

- Conductivity:
$$g_0 = ne\mu \sim 3 \frac{e^2}{h} \frac{E_{g0}}{\Delta}$$

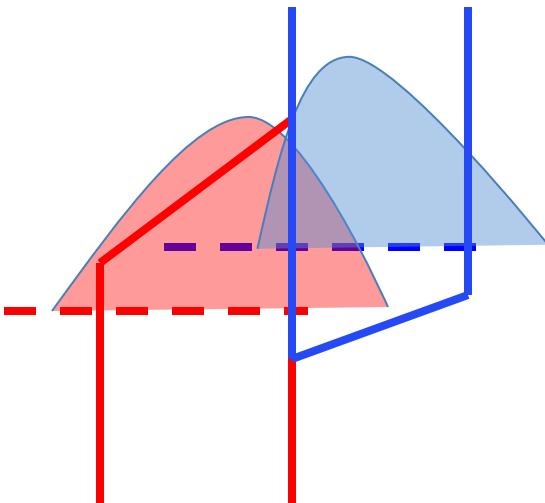
$$E_{g0} = \frac{\pi \hbar^2}{m^*} n_0$$



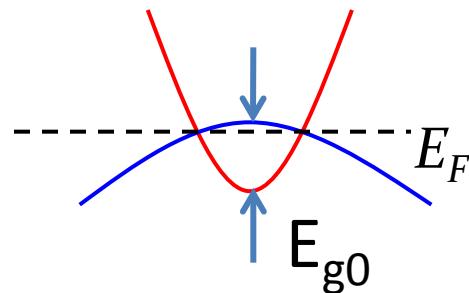
$$\Gamma \ll \Delta,$$



Band-Gap Engineering of g_0



$$\sigma_{on} \approx \frac{3e^2}{h} \frac{E_{g0}}{\Delta}$$



1. Less inverted
2. Narrow QWs
3. Higher mobility



Wafer A $E_{g0} \sim -33$ meV, σ (bulk) > 12 e²/h

Wafer B $E_{g0} \sim -18$ meV, σ (bulk) << 10 e²/h

Wafer B $E_{g0} \sim 0$ meV, σ (bulk) << 0.5 e²/h

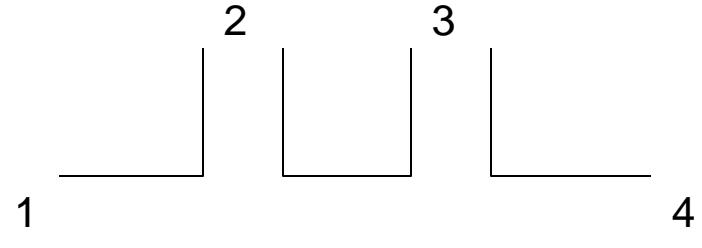
σ (bulk) ~ 10 e²/h is a “magic” number

Edge modes de-couple from bulk transport

Helical edge modes in four probe geometry

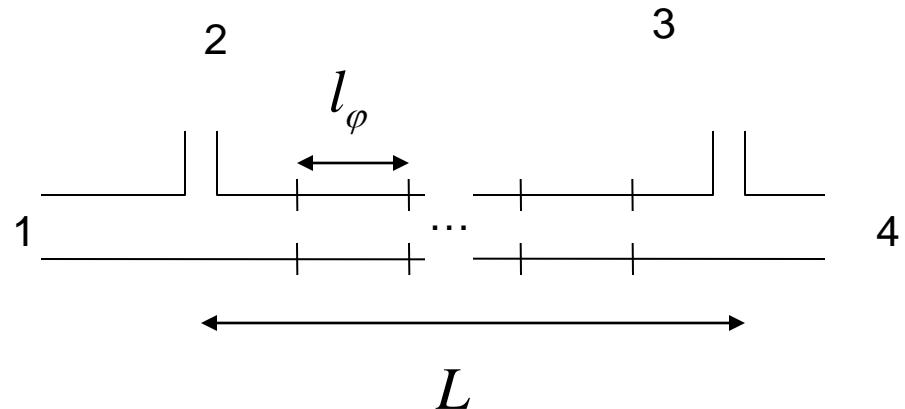
- For phase coherent samples Landauer-Buttiker formula gives conductance value:

$$G_{14,23} = \frac{4e^2}{h}$$

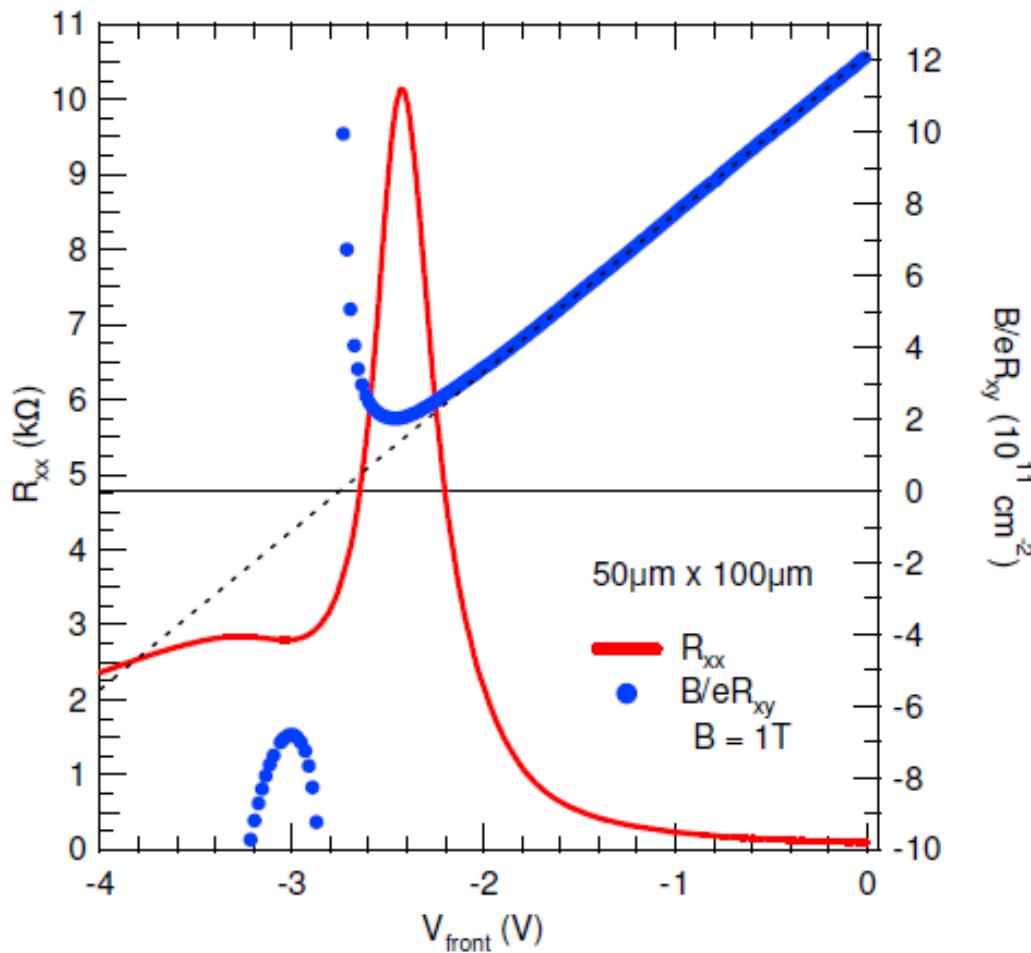


- Longer samples can be modeled by inserting phase breaking probes and applying LB-formula:

$$G_{14,23} = \frac{2e^2}{h} \left(\frac{l_\phi}{L} + \left(\frac{l_\phi}{L} \right)^2 \right)$$



Transport in the Mini-gap



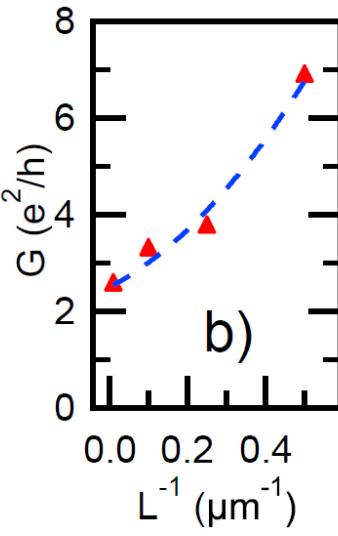
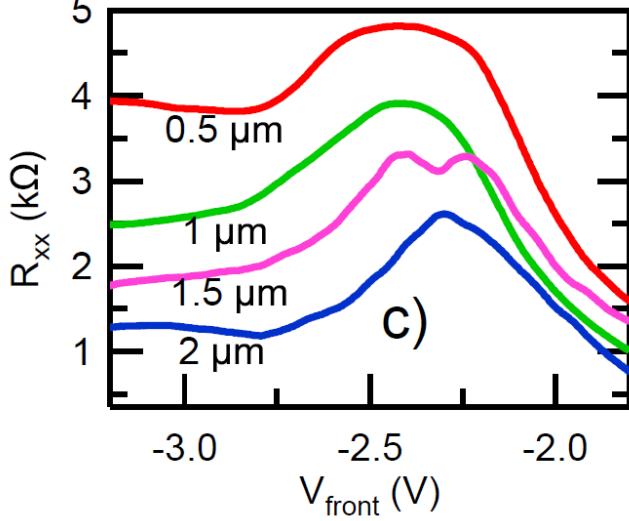
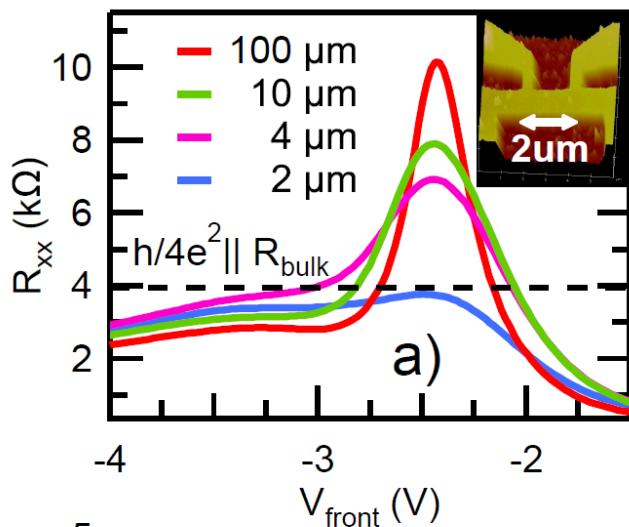
- Size of the gap can be estimated from the width of the peak in front gate bias:
 $\Delta = 4 \text{ meV}$, while
 $E_{g0} = 16 \text{ meV}$
- Resistance peak value agrees well with what is expected for bulk transport

$$g_0(\text{expected}) \approx \frac{e^2}{h} \frac{E_{g0}}{\Delta} = \frac{4e^2}{h}$$

$$g_0(\text{measured}) = 5.05 \frac{e^2}{h}$$

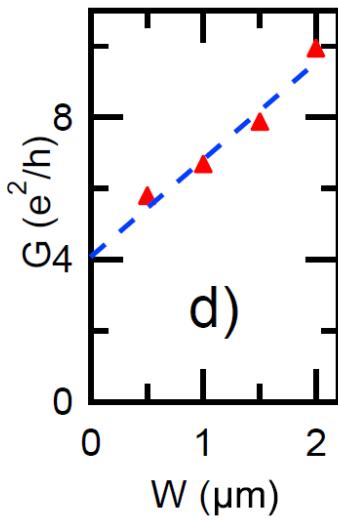
Scaling Evidence for The Edge Modes

Wafer B



L-Dependence, keep $W = 1 \mu\text{m}$

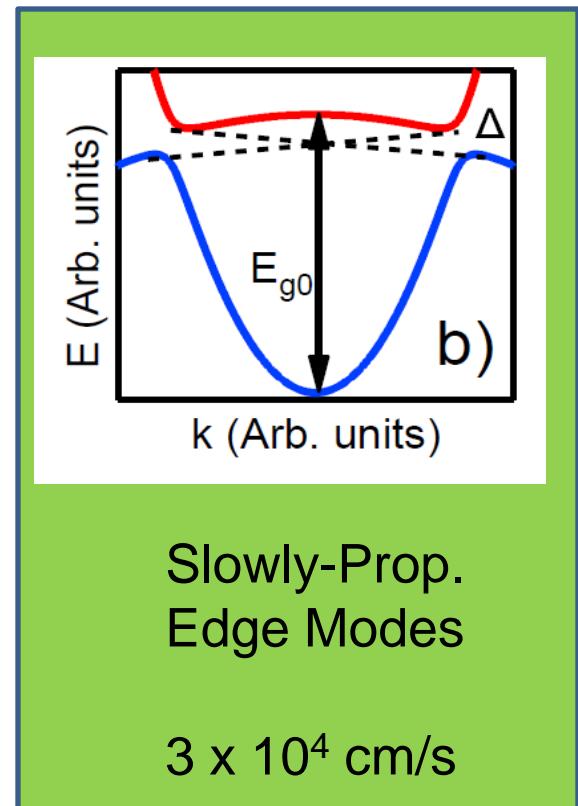
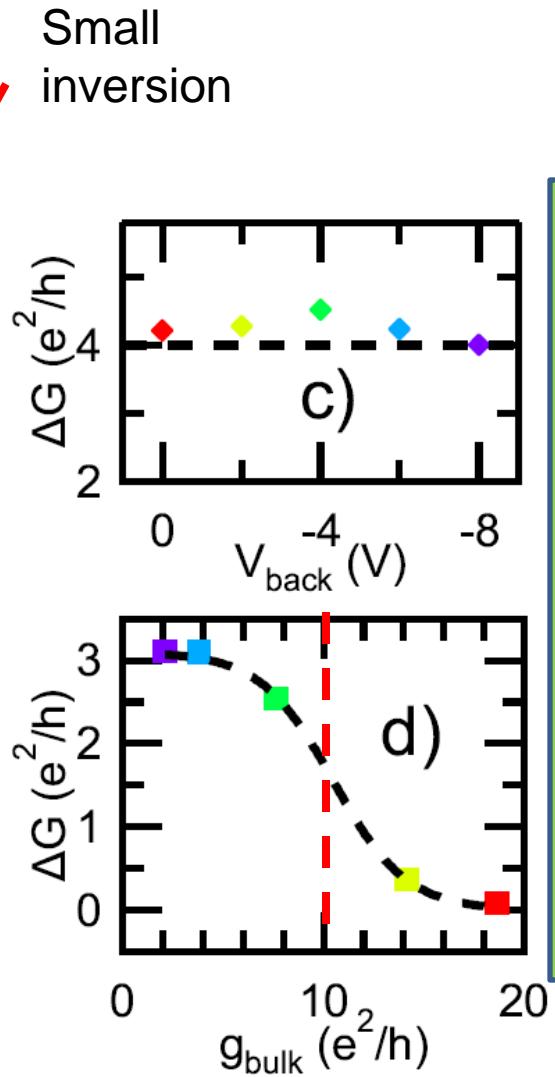
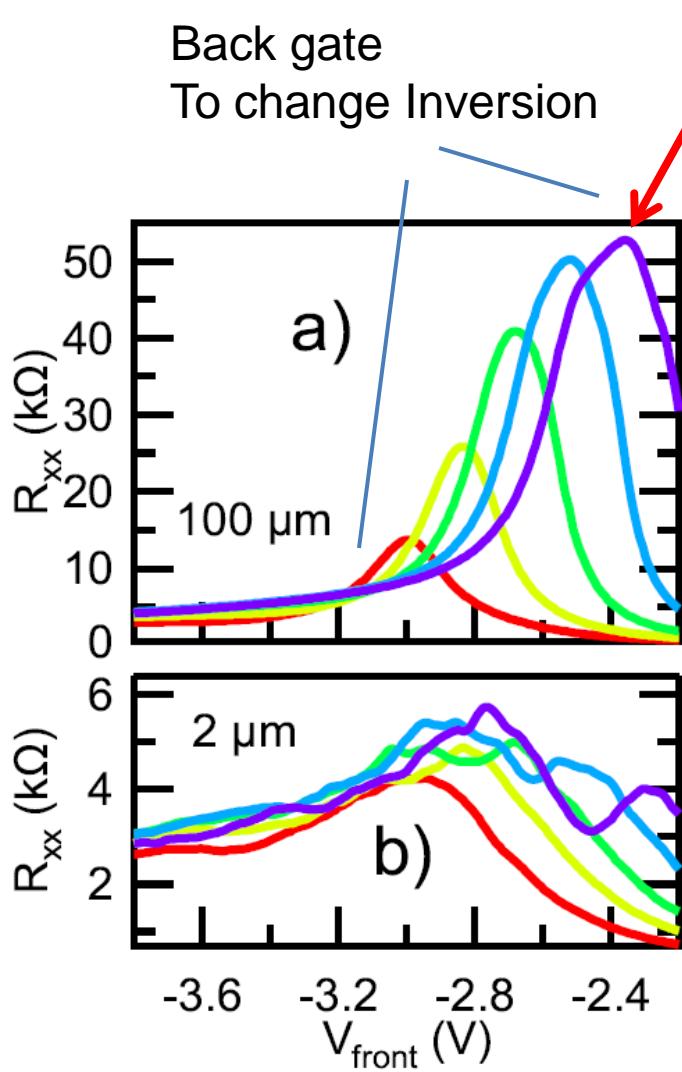
$$G_{14,23} = \frac{2e^2}{h} \left(\frac{l_\phi}{L} + \left(\frac{l_\phi}{L} \right)^2 \right)$$



W-Dependence, Keep $L/W=2$

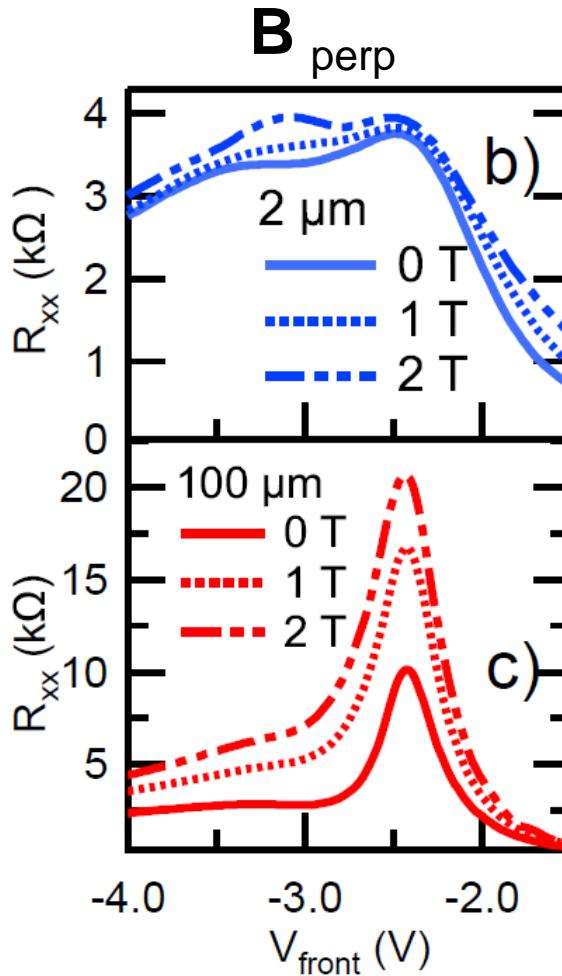
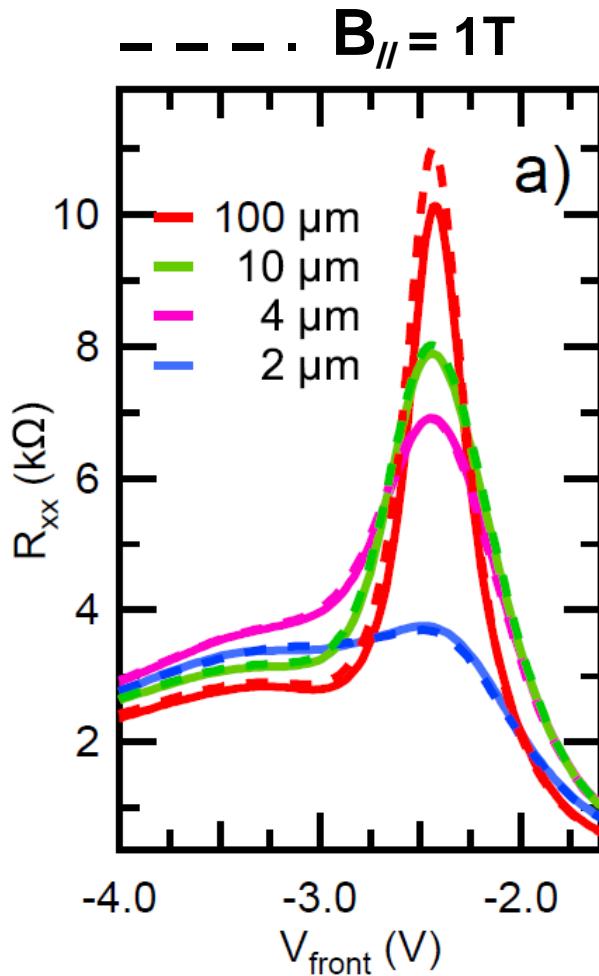
Conclusion For π geometry, the edge conductivity is \sim quantized to $4 e^2/h$

Robustness of Edge Quantization vs Bulk Conduct.

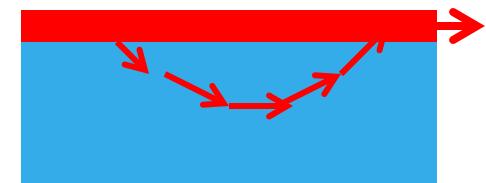


“Anomalous” Magnetic Field Dependence

All 300 mK

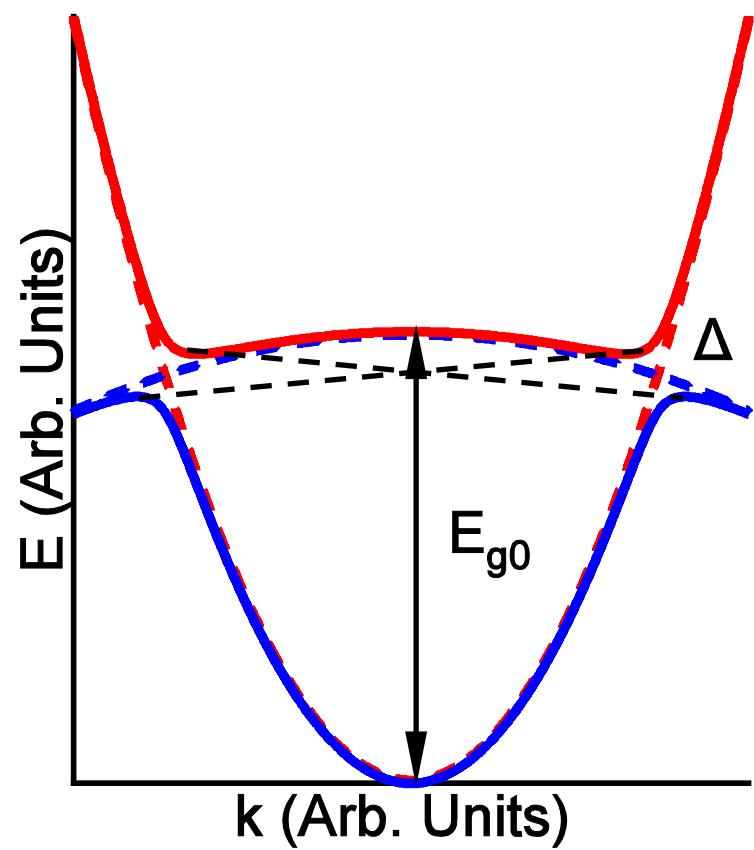
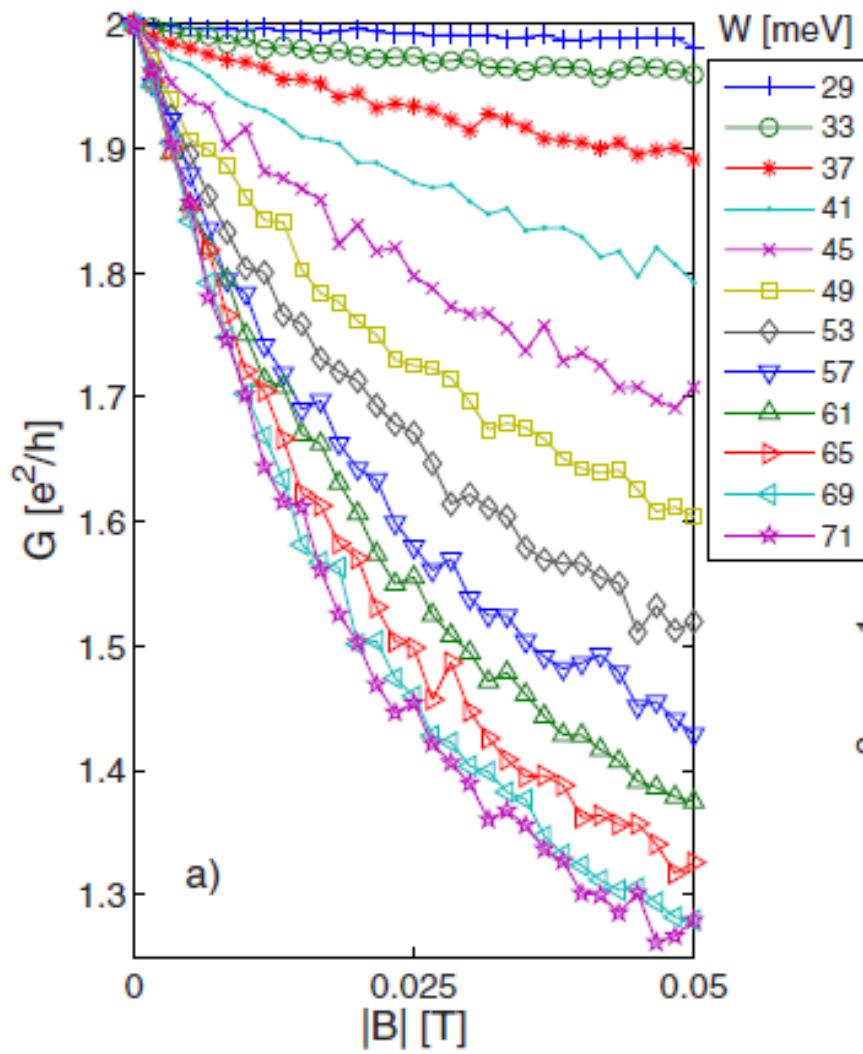


1. Long samples show stronger positive MR –bulk effect
2. μm samples weaker depend.
3. Edge states does not scatt. into and out of bulk – lack of loop



Decouple from bulk
Maciejko *et al*, *Phys. Rev. B*,
83, 155310 (2010)

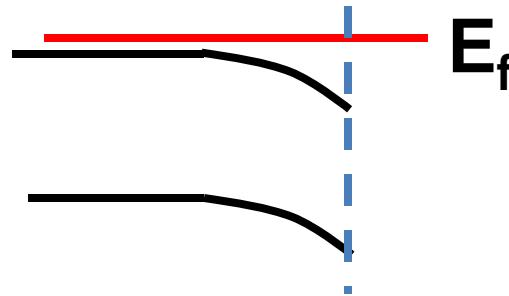
Magneto-conductance of QSH



Opens a Number of Exciting Experiments

1. Edge states extend all the way to the material edge

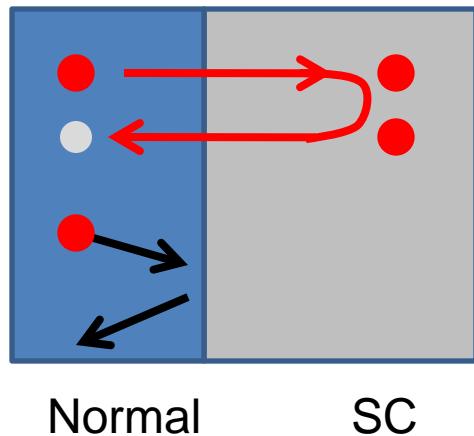
(E_f pinned above E_c for InAs)



2. Good interface with superconductors due to low Schottky barrier --- SC Hybrids

Examples

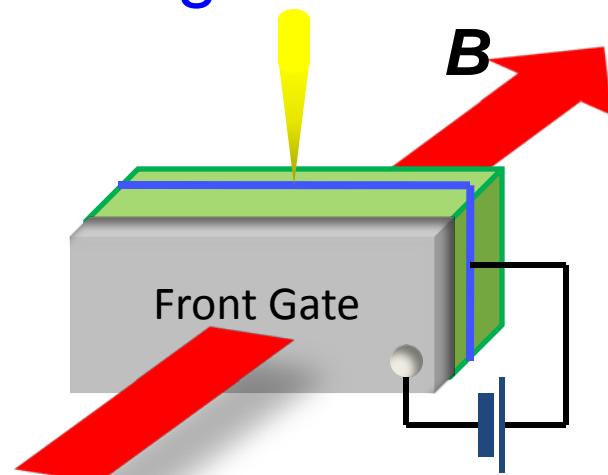
Andreev Reflection



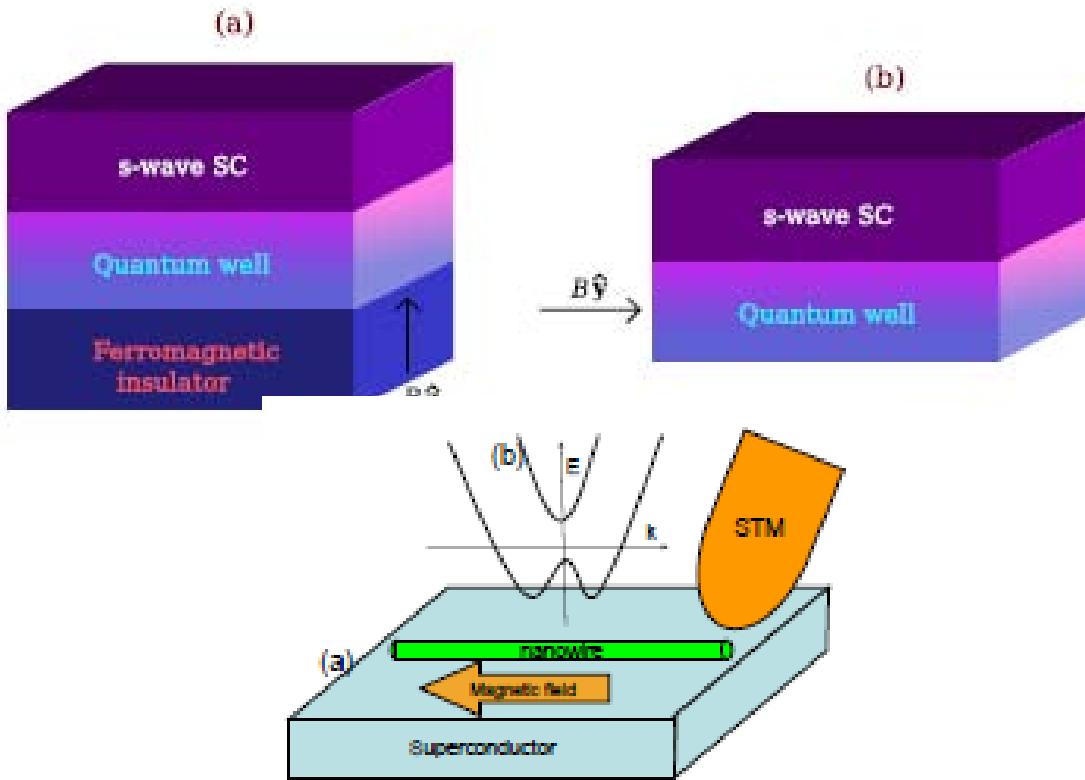
Majorana
Bound State
(MBS)

Non-Equilibrium
Josephson Effect
Through Helical
Edge State
arXiv: 1108.3870

Cleavage STM/STS



Majorana Fermions in SC/InAs/SC



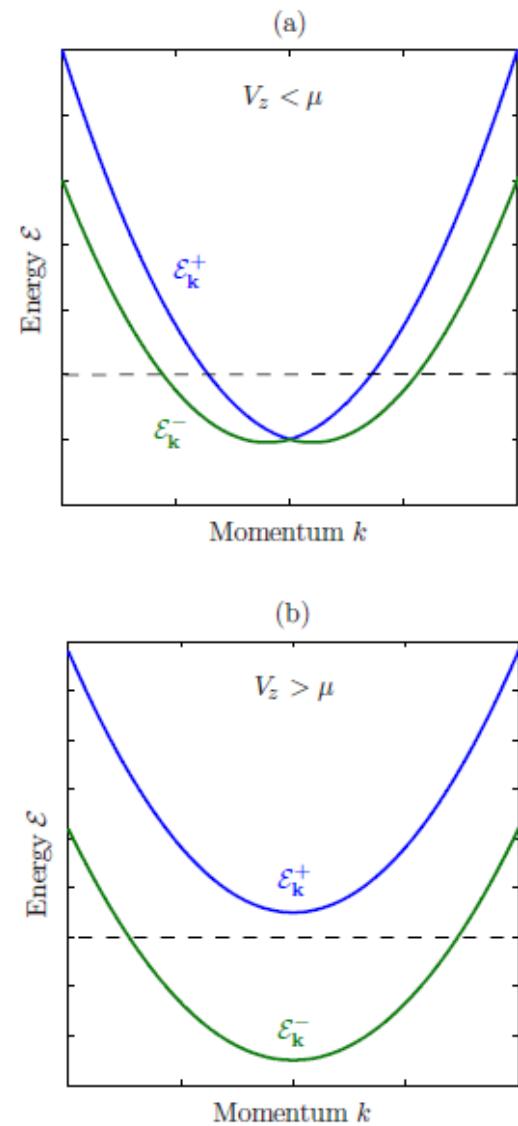
Fu & Kane, Phys. Rev. Lett, **100**, 096407 (2008)

Sau *et al*, Phys. Rev. Lett. **104**, 040502 (2010)

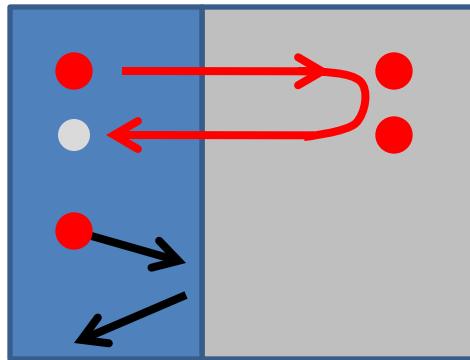
Alicea, Phys. Rev. B **81**, 125318 (2010)

Linder and Sudbo, Phys. Rev. B **82**, 085314 (2010)

Sau *et al*, arXiv:1006.2829 (2010)



Andreev Reflection (Single N-S)



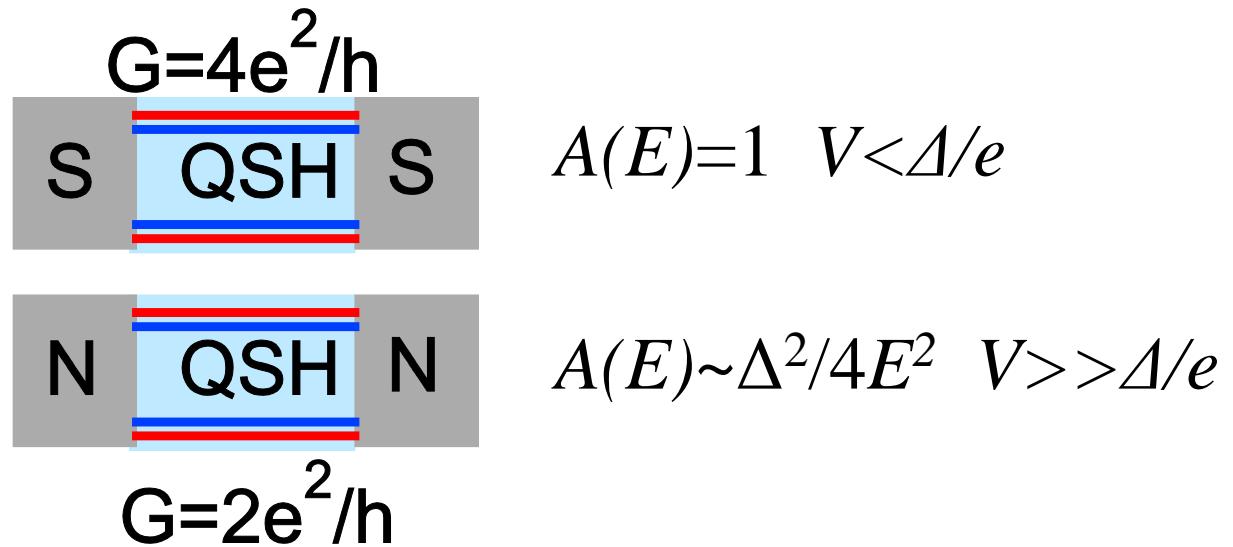
Normal **SC**

the interface. The electrical current through a single S-N interface can be calculated using the Blonder-Tinkham-Klapwijk (BTK) model:²²

$$I = \frac{N \cdot e}{h} \int [f(E + eV) - f(E)] [1 + A(E) - B(E)] dE \quad (1)$$

where N is the number of modes in the normal conductor, $f(E)$ is the equilibrium Fermi distribution function, V is the voltage drop at the interface, and $A(E)$ and $B(E)$ are probabilities for Andreev and normal reflection (NR) of the electron at the interface. In the case of

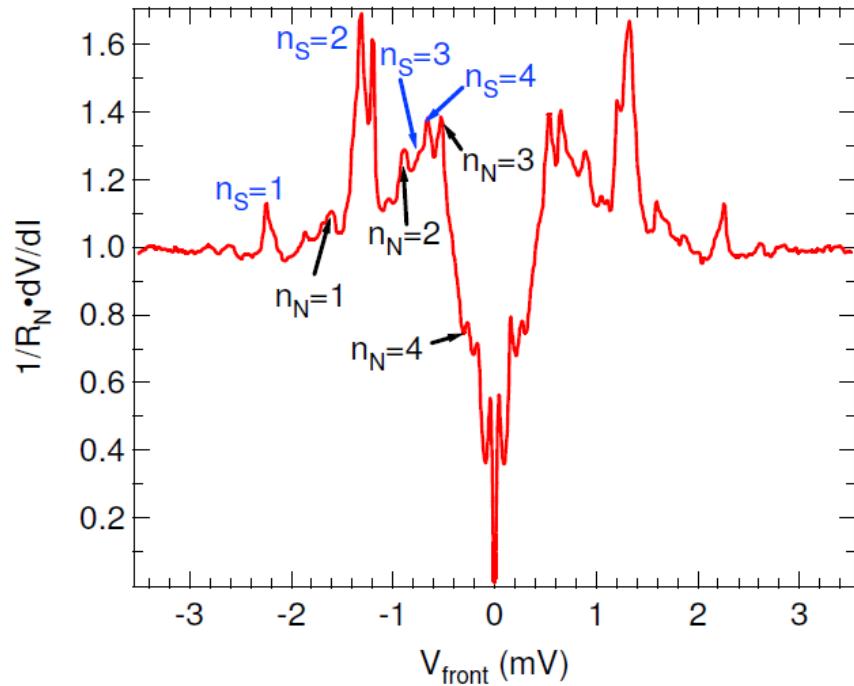
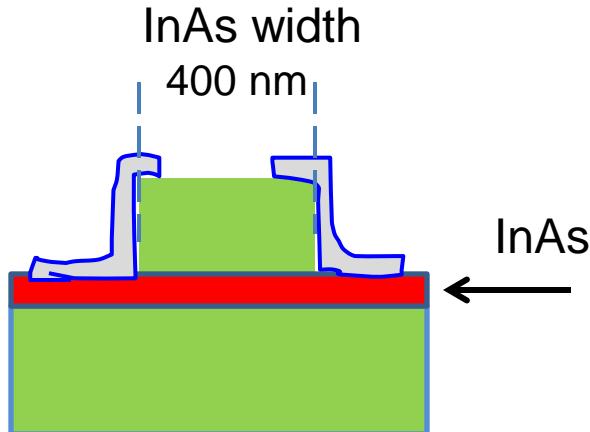
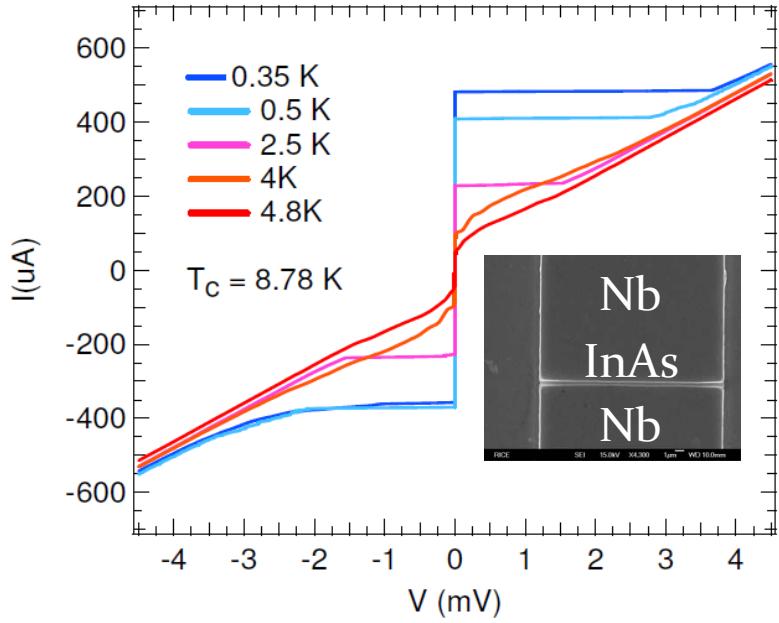
Andreev Reflection of QSH Edges



$$I = \frac{N \cdot e}{h} \int [f(E + eV) - f(E)] [1 + A(E) - B(E)] dE$$

1. Can we make Nb/InAs Junction - Test ✓

2. Can we make Nb/InAs-GaSb Junction - New



Multiple Adreev Reflection:

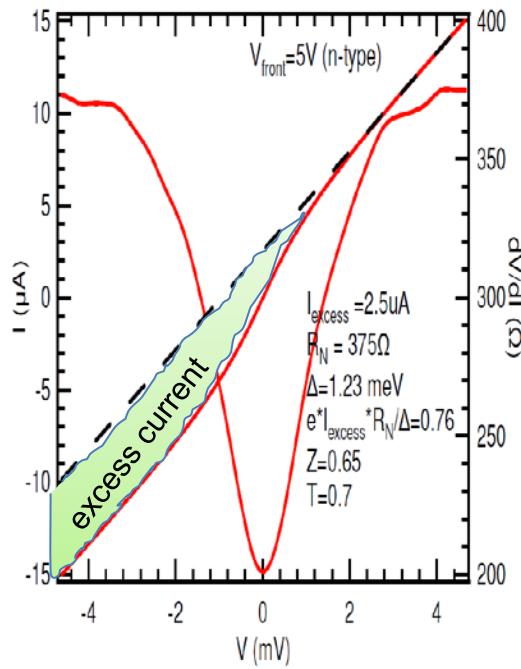
3 sets of peaks: Δ_N , Δ_S , $\Delta_N + \Delta_S$

$$V_{SGS} = \frac{2 \cdot \Delta_S}{n \cdot e}, \frac{2 \cdot \Delta_N}{n \cdot e}, \frac{2 \cdot (\Delta_S \pm \Delta_N)}{n \cdot e}$$

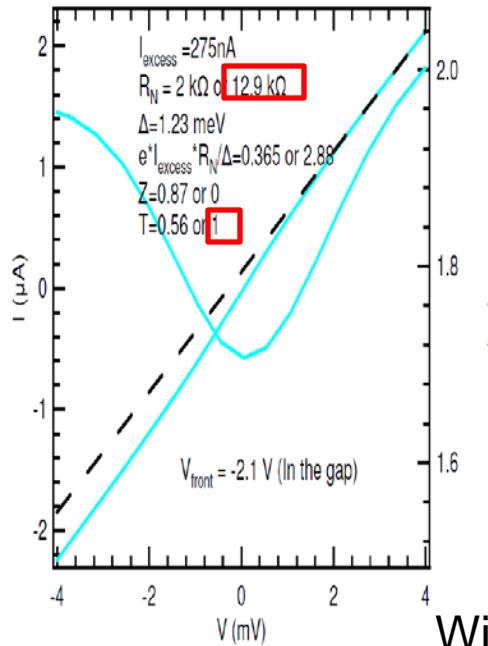
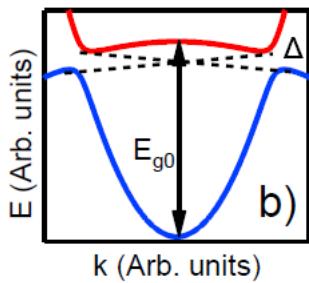
$$\Delta_N / \Delta_S \sim 0.67$$

$$\Delta_S = 1.31 \text{ meV}; \Delta_N = 0.87 \text{ meV}$$

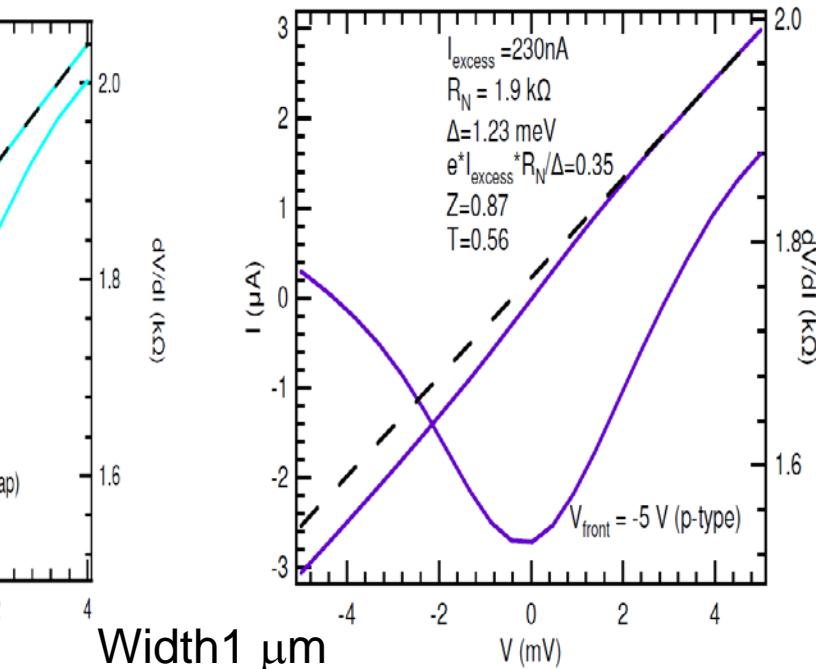
Andreev Reflection in 3 Regimes



electron side

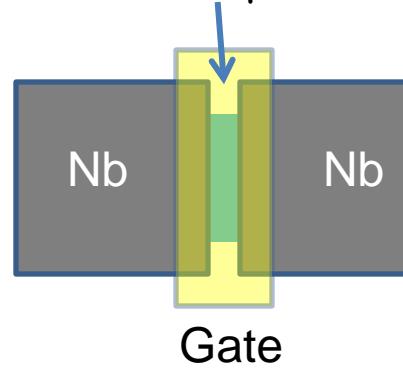


In the gap



Width 1 μm

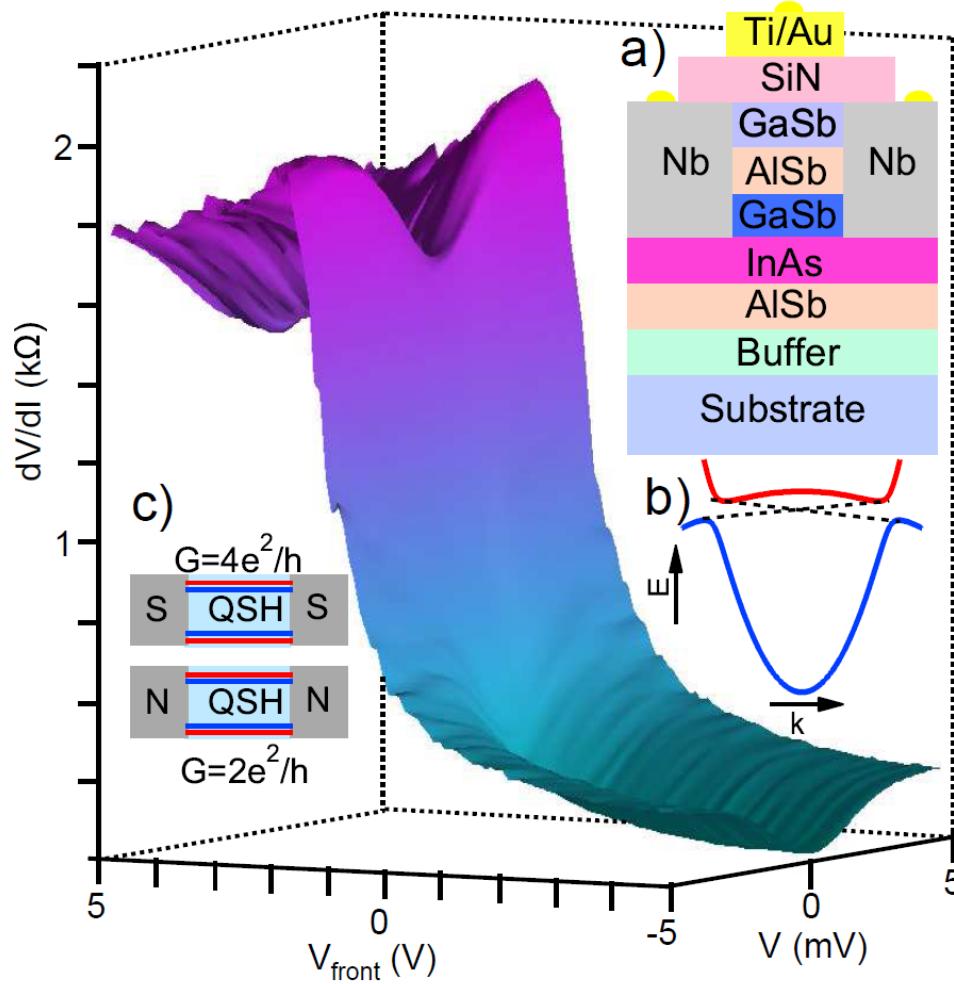
hole side



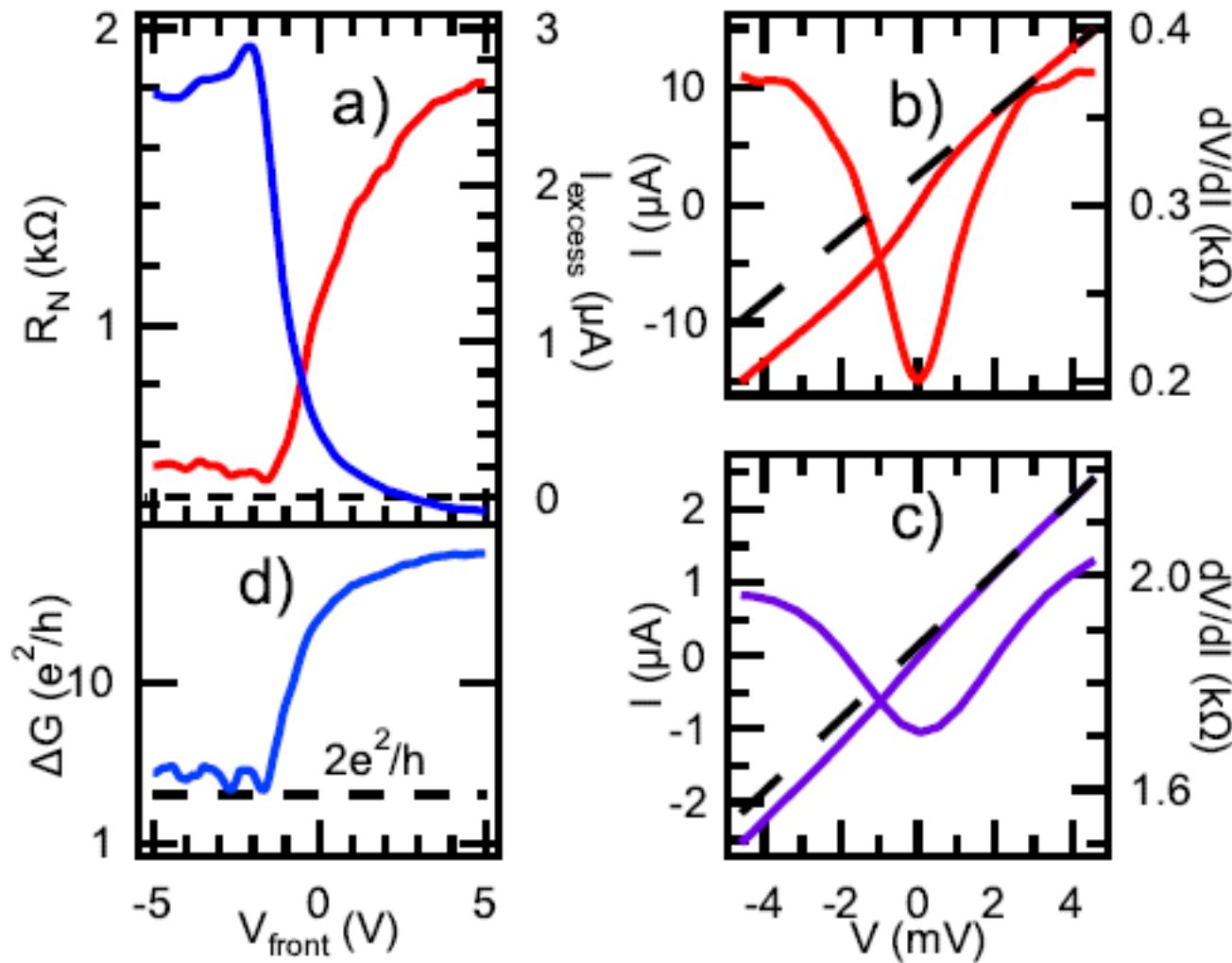
Andreev Reflection in Nb/(InAs-GaSb)/Nb

Knez, Du, Sullivan

ArXiv:1106.5819 (2011)

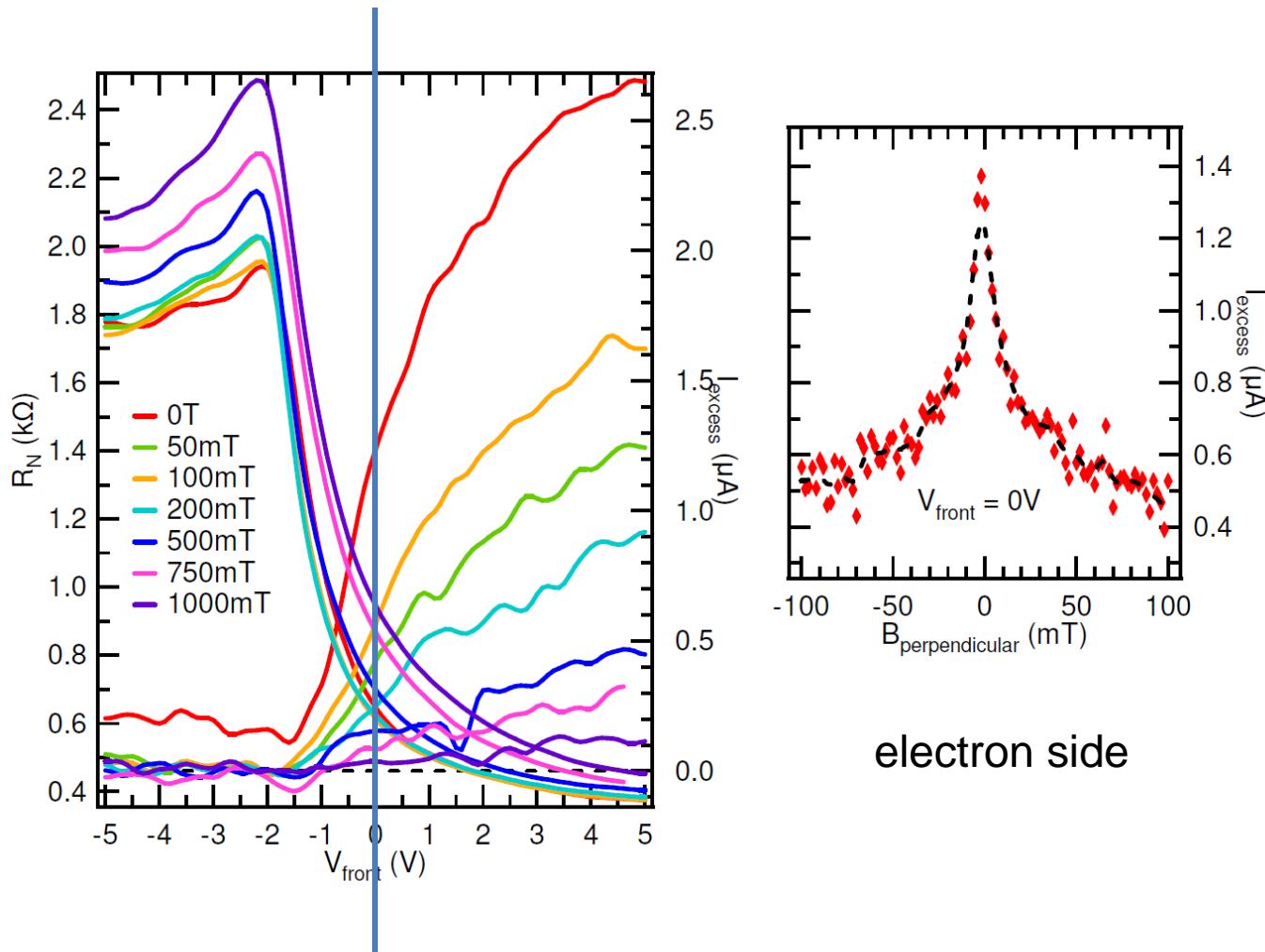


Perfect Andreev Reflection



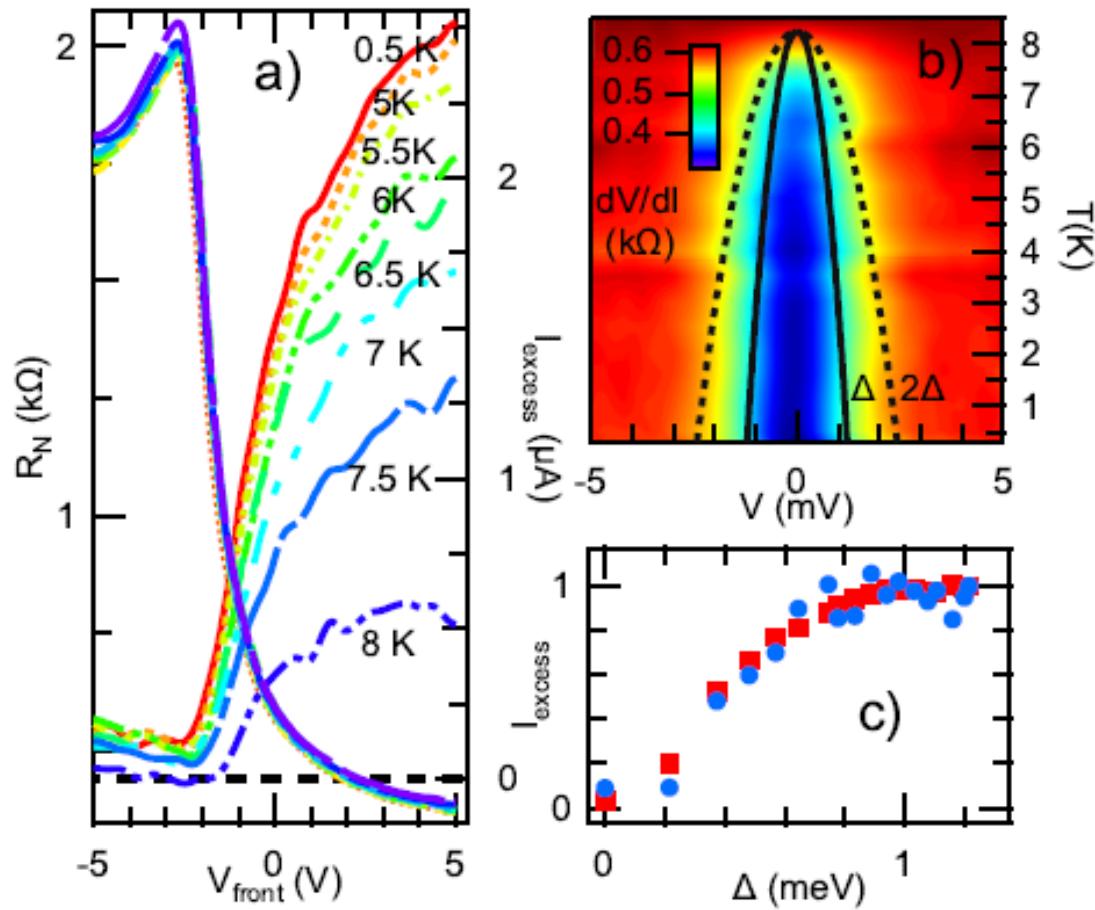
3

Excess Current Suppressed by Magnetic Fields



electron side

Temperature dependence

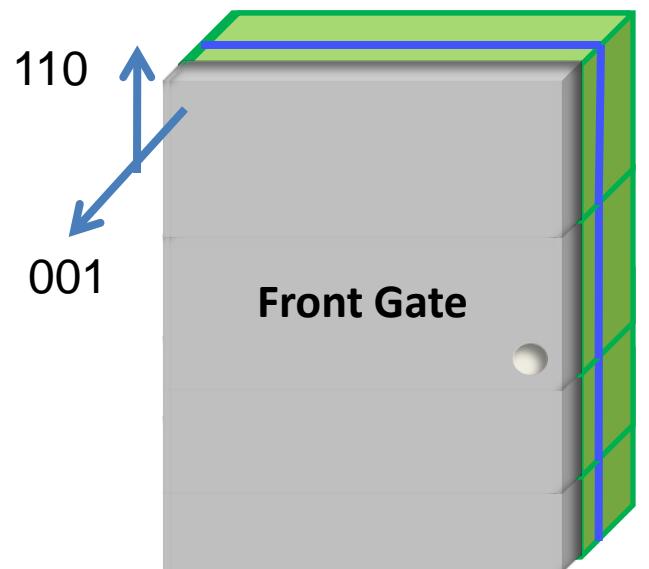


Challenges

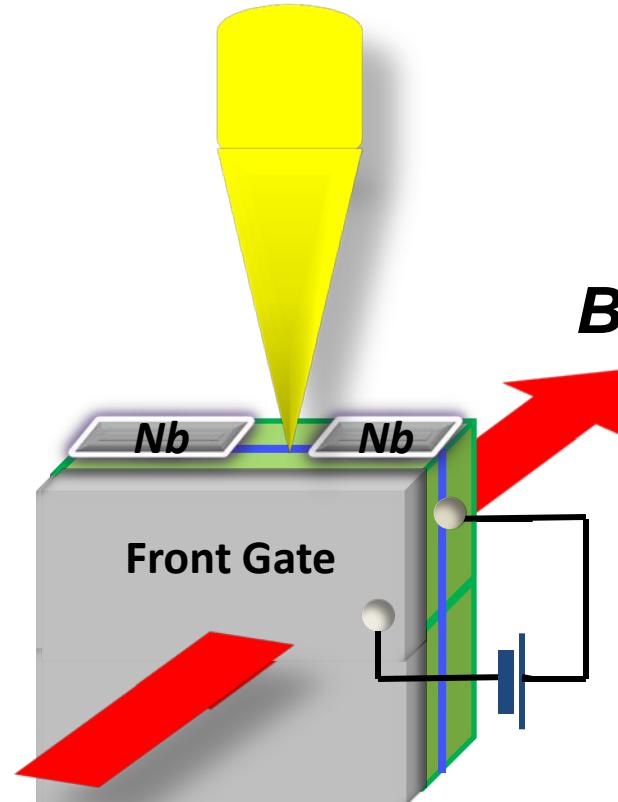
1. Geometry: SNS difficult to analyze
2. Supercurrent
3. Interface quality

Example: STM/STS

Nb/InAs-GaSb Junction on cleaved edge



UHV
Cleavage
Deposition



STM/STS
400 mK
9/2/2 T

SUMMARY

1. We confirmed existence of hybridization gap in inverted InAs/GaSb.
2. Length and width dependence of transport coefficients in in inverted InAs/GaSb QWs suggests existence of helical edge modes
3. Edge and Bulk interaction shows critical behavior
4. Edge modes show only weak dependence on magnetic fields
5. Preliminary data for InAs/GaSb –SC Andreev reflection

THANKS