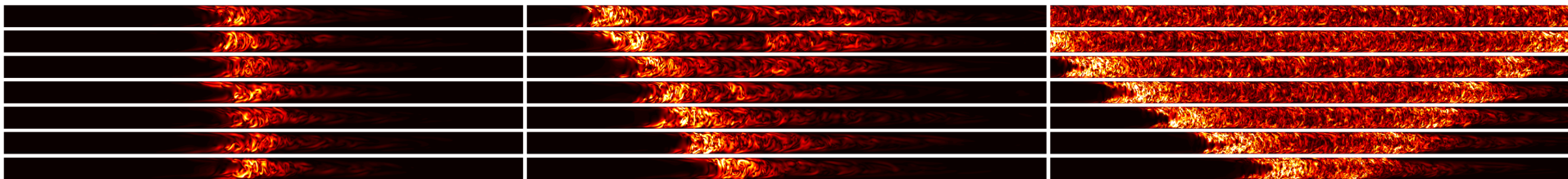


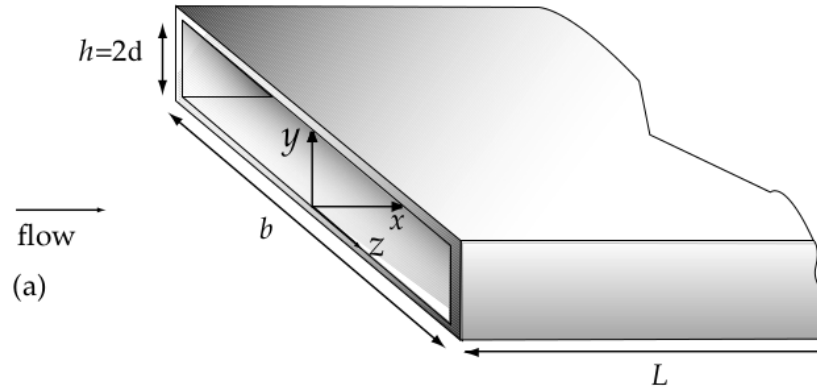
Physics and modeling of laminar-turbulent interfaces in pipe flow

Marc Avila

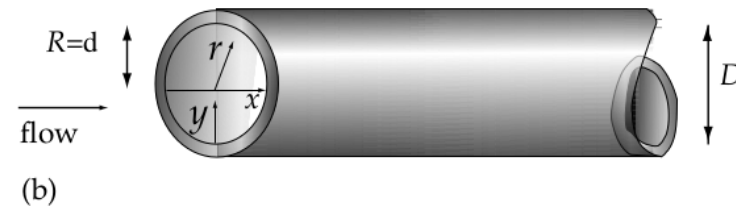


Why pipe flow?

Channel flow

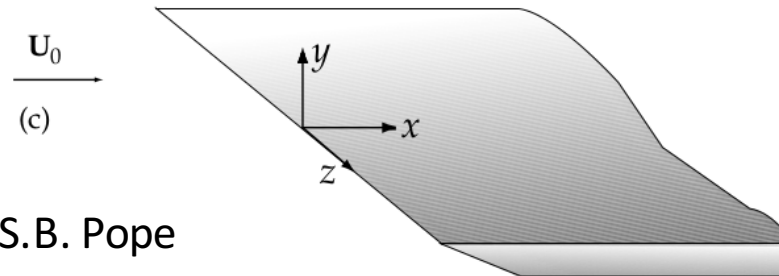


Pipe flow



Span-wise periodic
Quasi 1D

Boundary layer flow

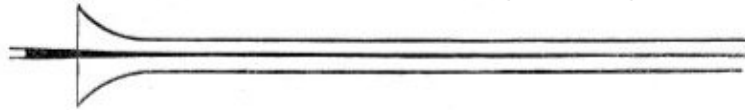


TURBULENT FLOWS, by S.B. Pope

Transitional pipe flow

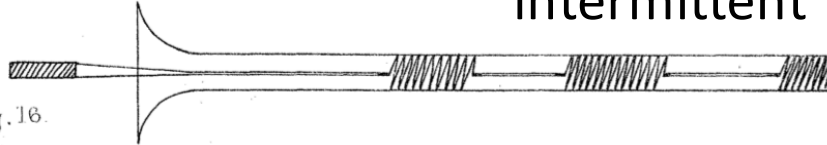
Reynolds *PTRSL* 1883

Fig. 3. laminar



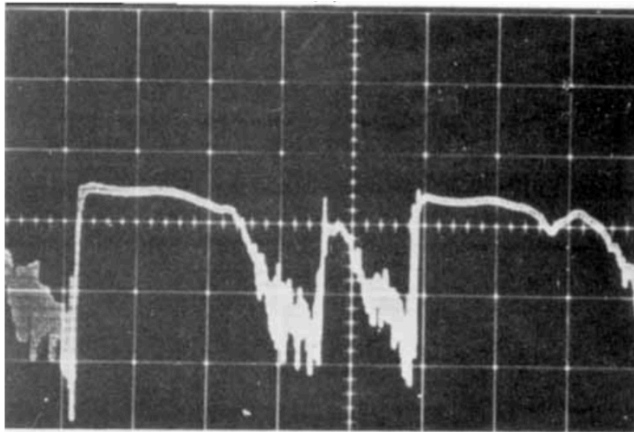
intermittent

Fig. 16.

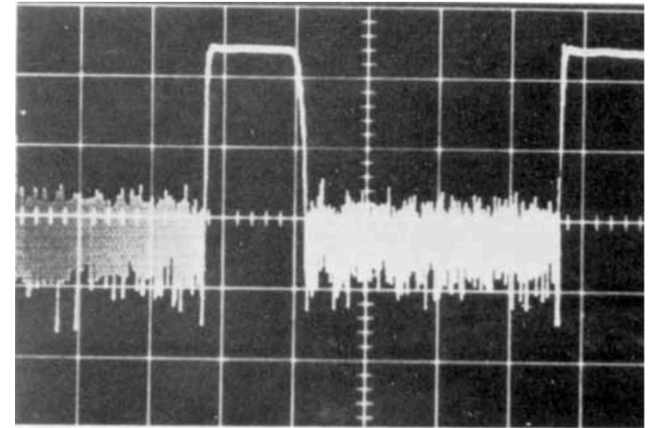


Wynanski & Champagne *JFM* 1973: obstacle at inlet

Re=2400 Spatio-temporal intermittency

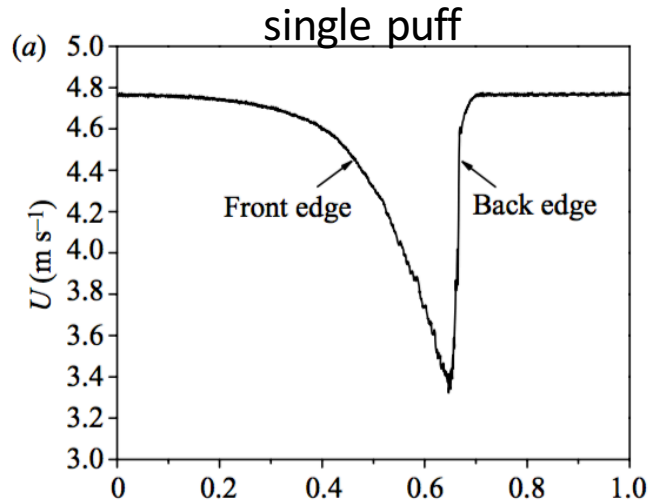


Re=4200 Slug sequence

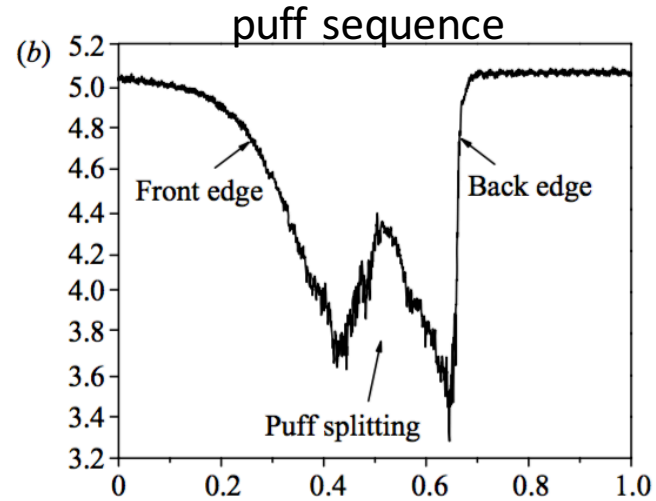


Laminar-turbulent fronts

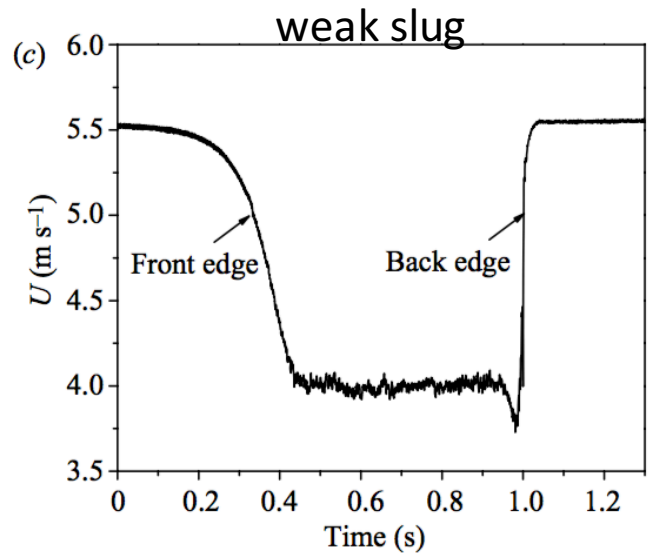
Re~2100



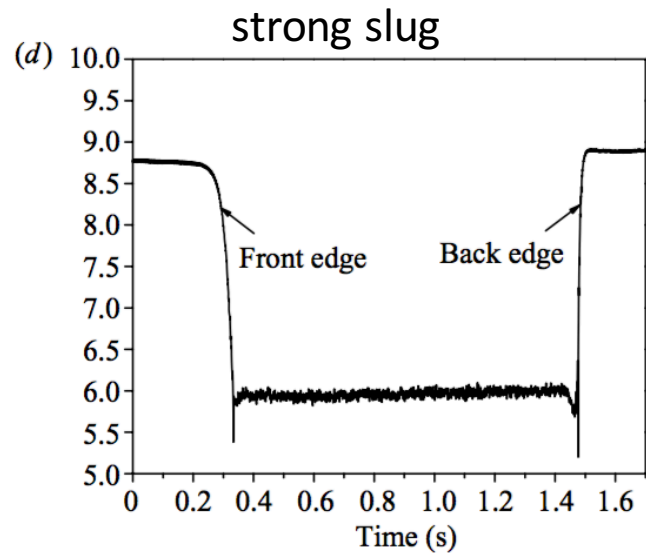
Re~2300



Re~3000



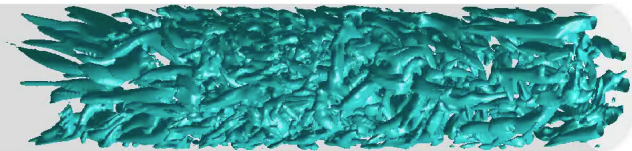
Re~4100



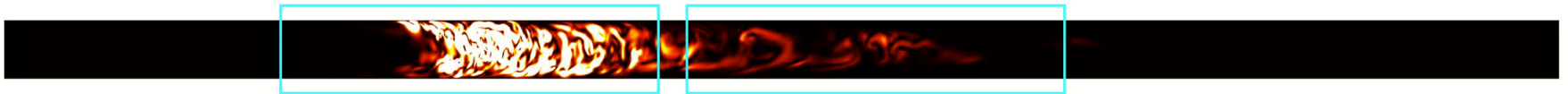
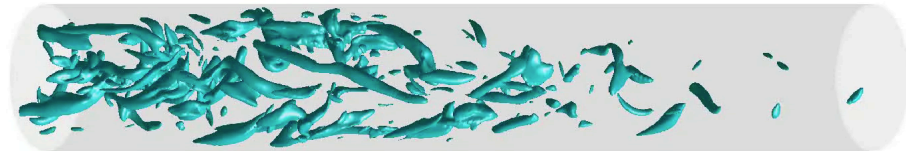
DNS: weak slug

Re=2600, weak downstream front

upstream



downstream



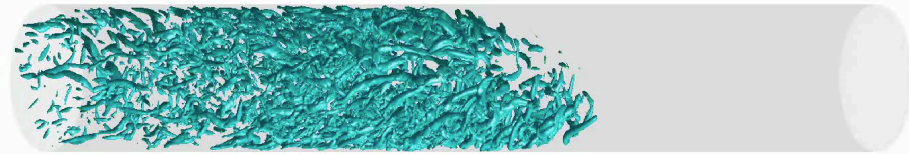
DNS: strong slug

Re=5000, strong downstream front

upstream



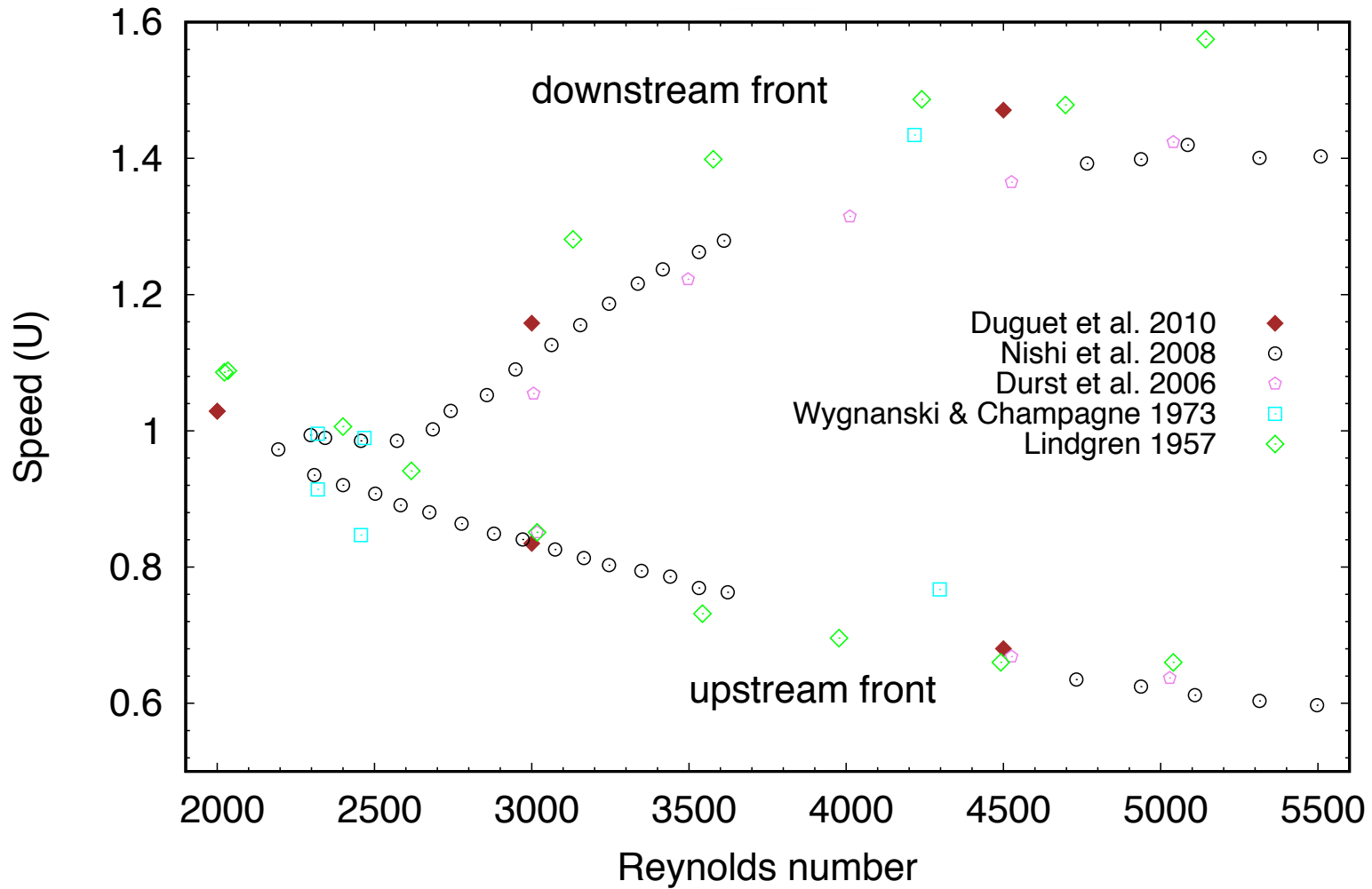
downstream



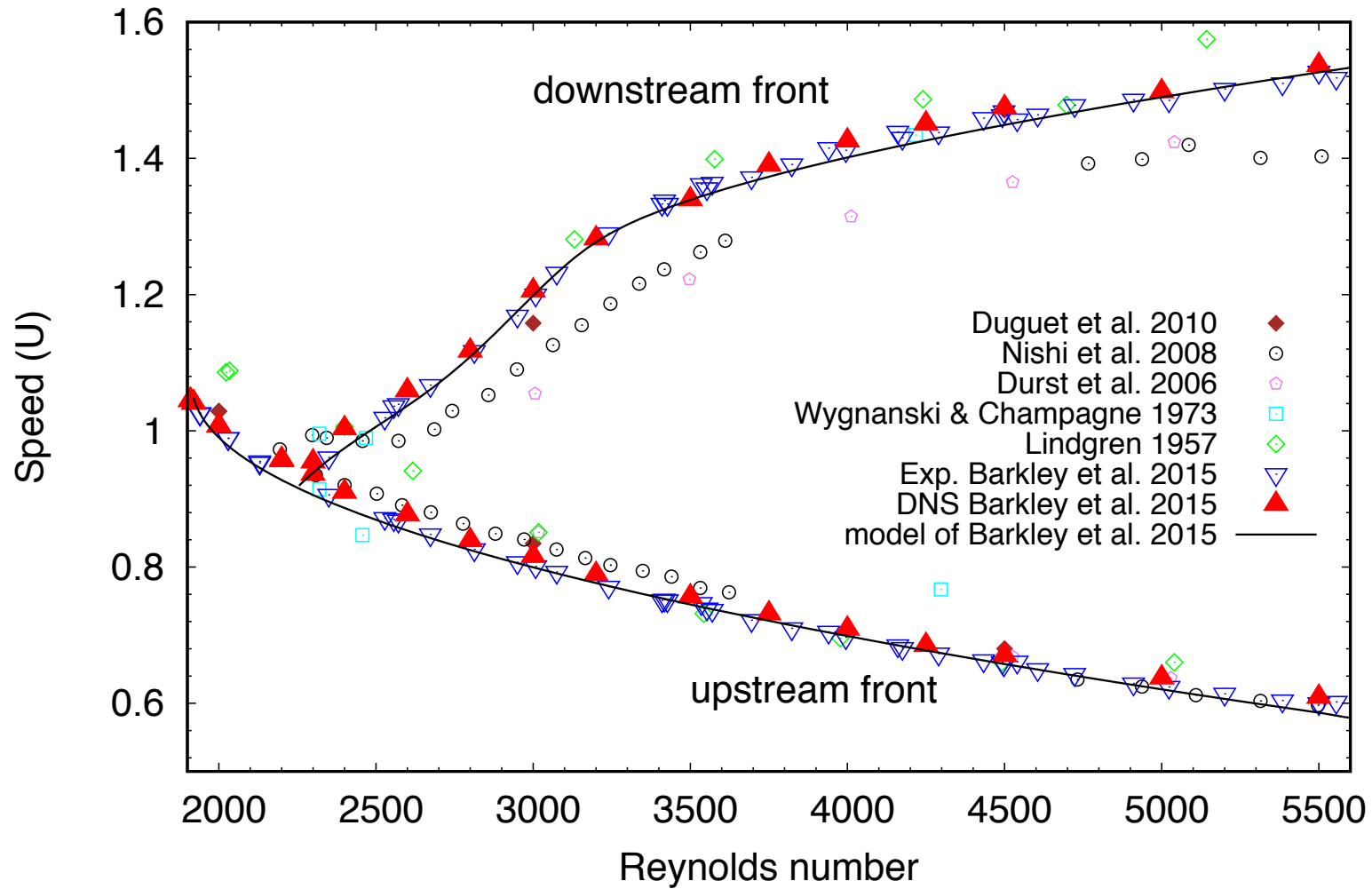
bulk



Front speeds



Front speeds

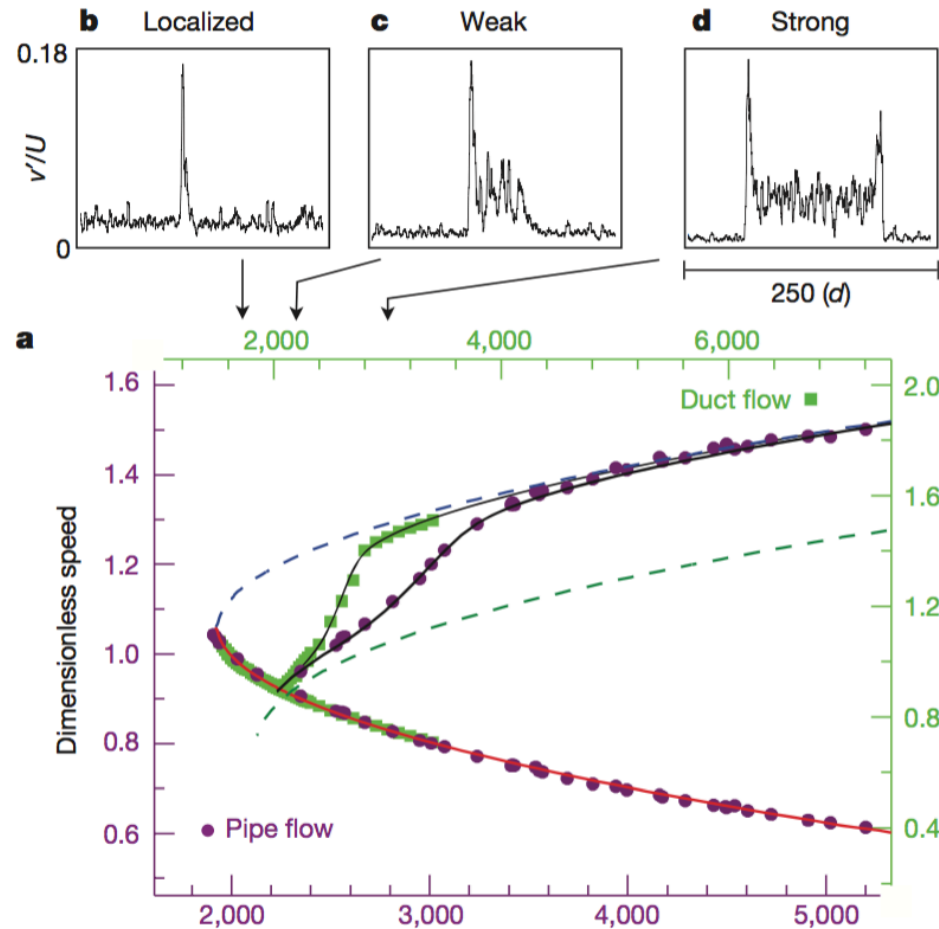
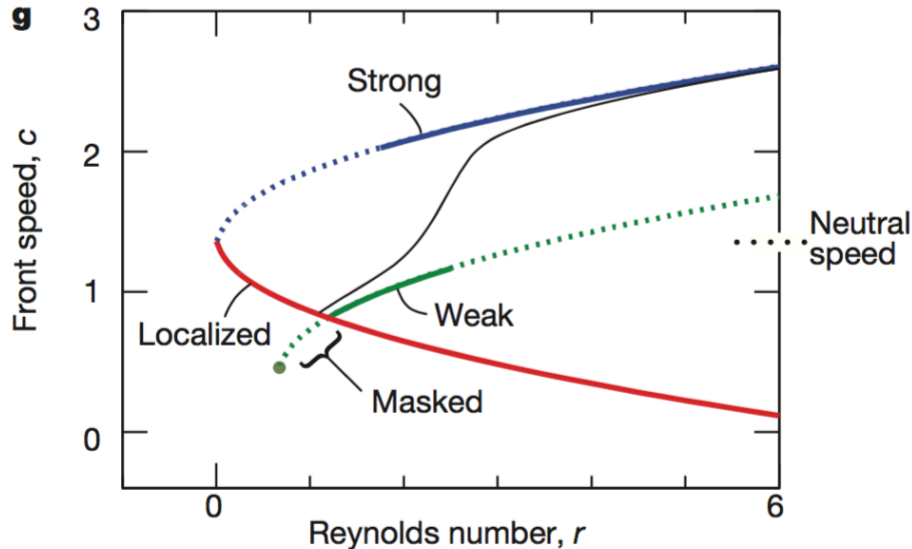


The rise of fully turbulent flow

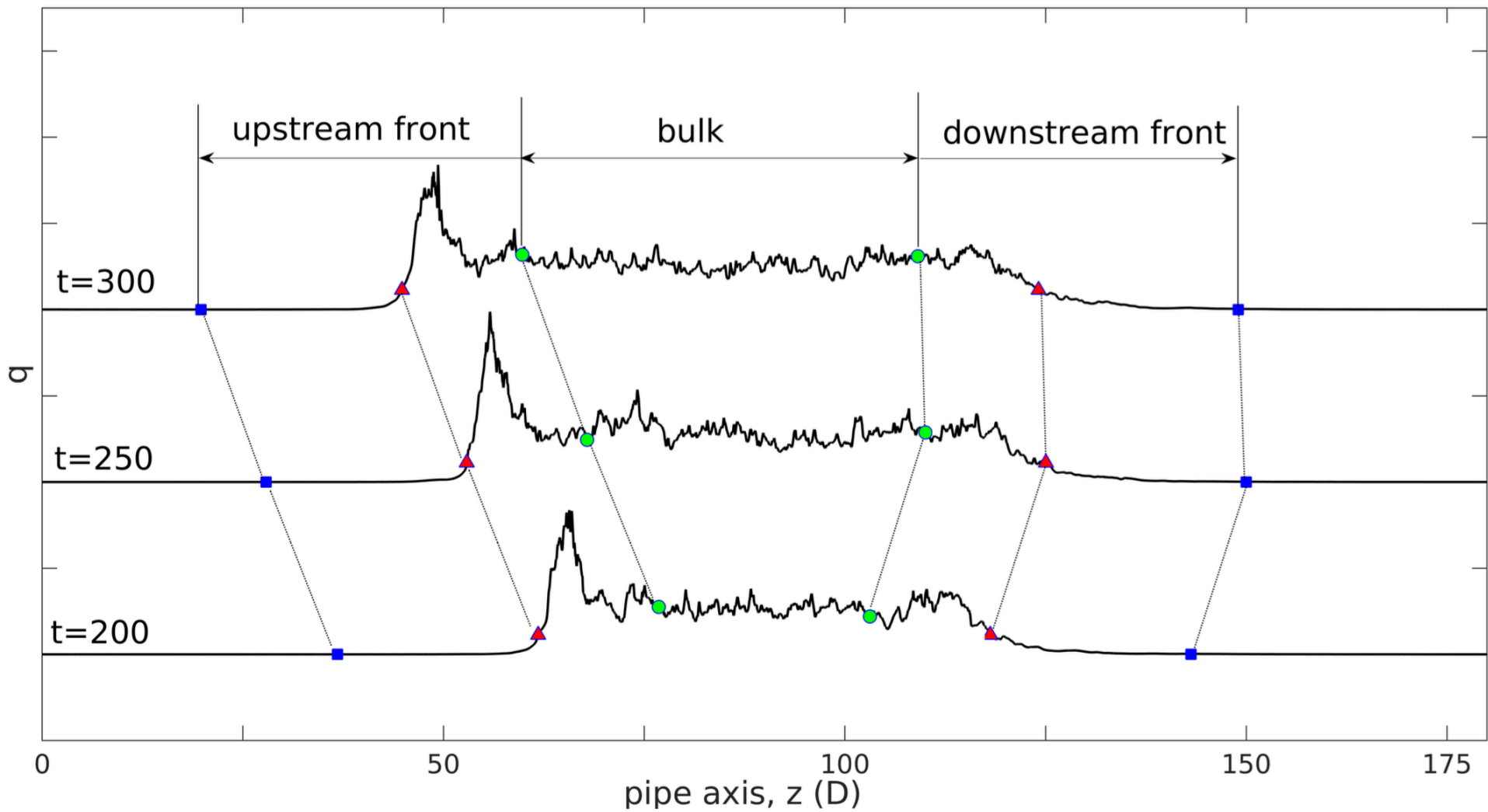
u : centerline velocity
 q : turbulence intensity

$$u_t + uu_x = \epsilon g(q, u)$$

$$q_t + (u - \zeta)q_x = f(q, u) + Dq_{xx}$$

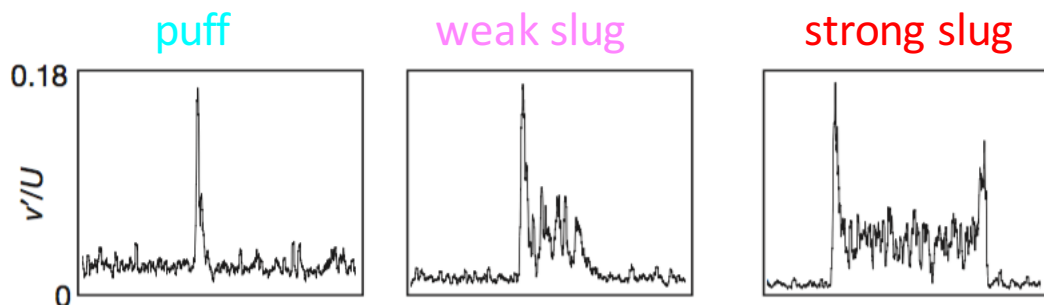
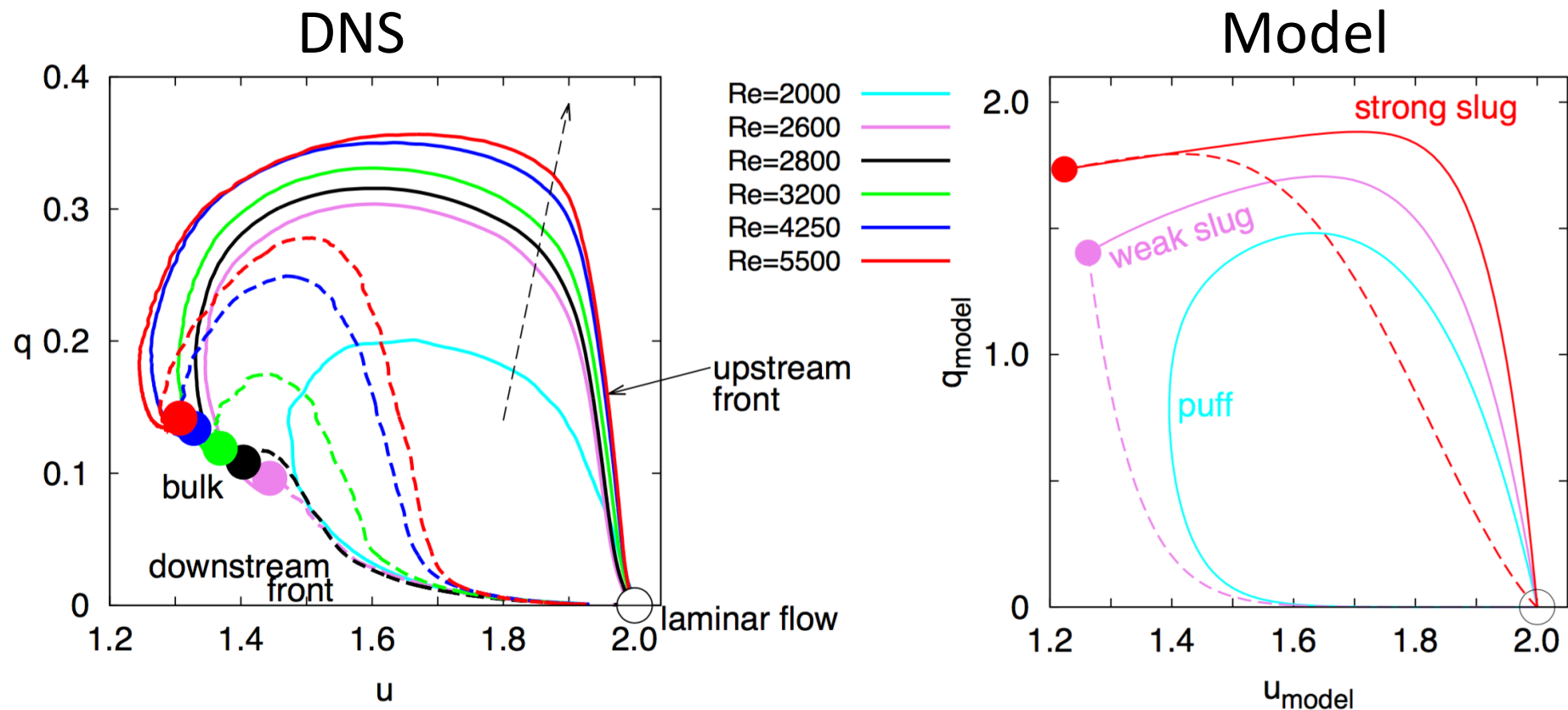


Front definitions

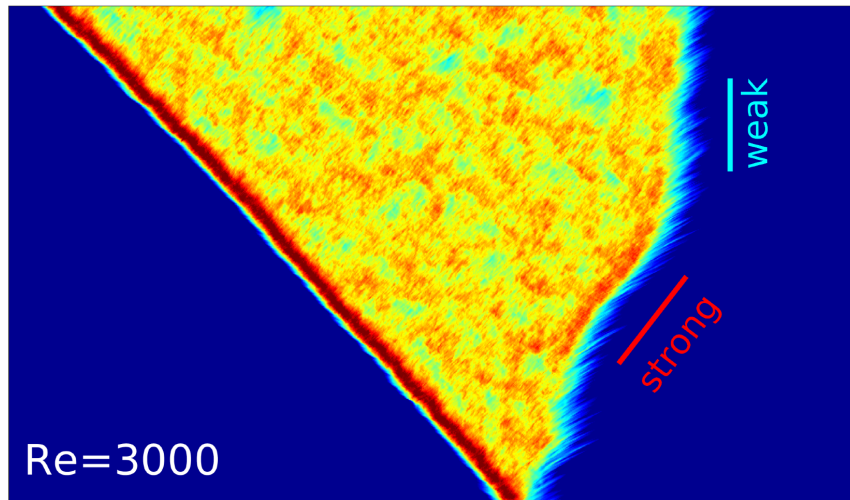
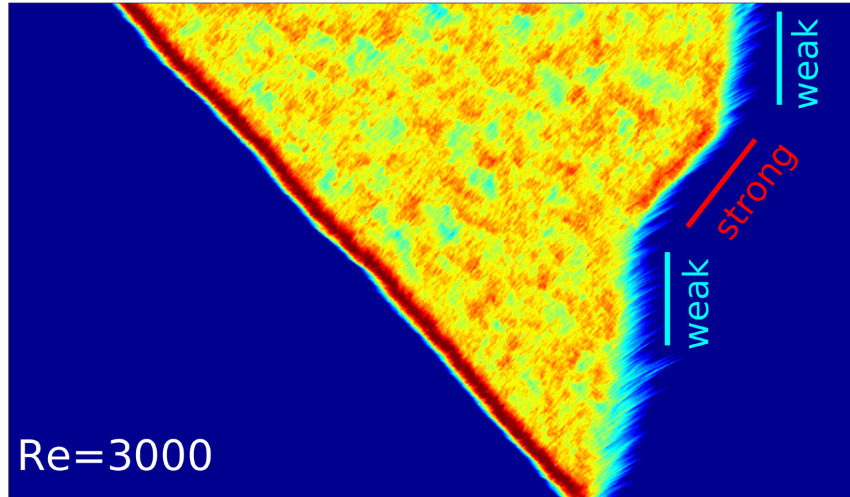


$$q(z) := \sqrt{\int \int (u_r^2 + u_\theta^2) r dr d\theta}$$

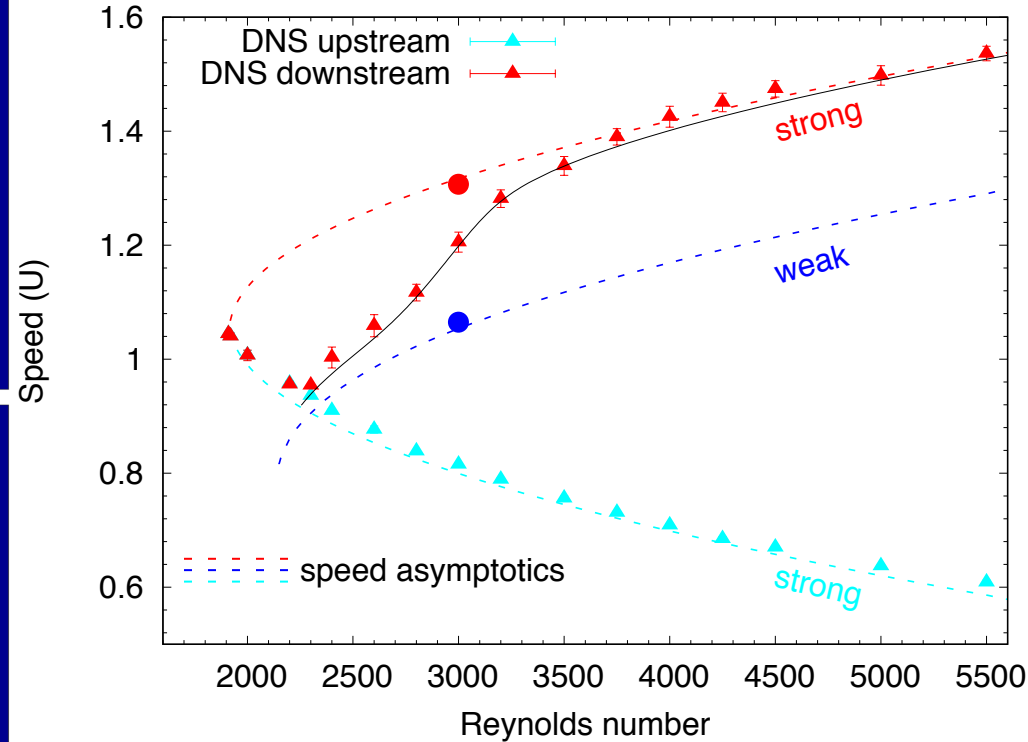
Front structure in (u, q) -space



Front switching



Pipe axis



Turbulent kinetic energy equation

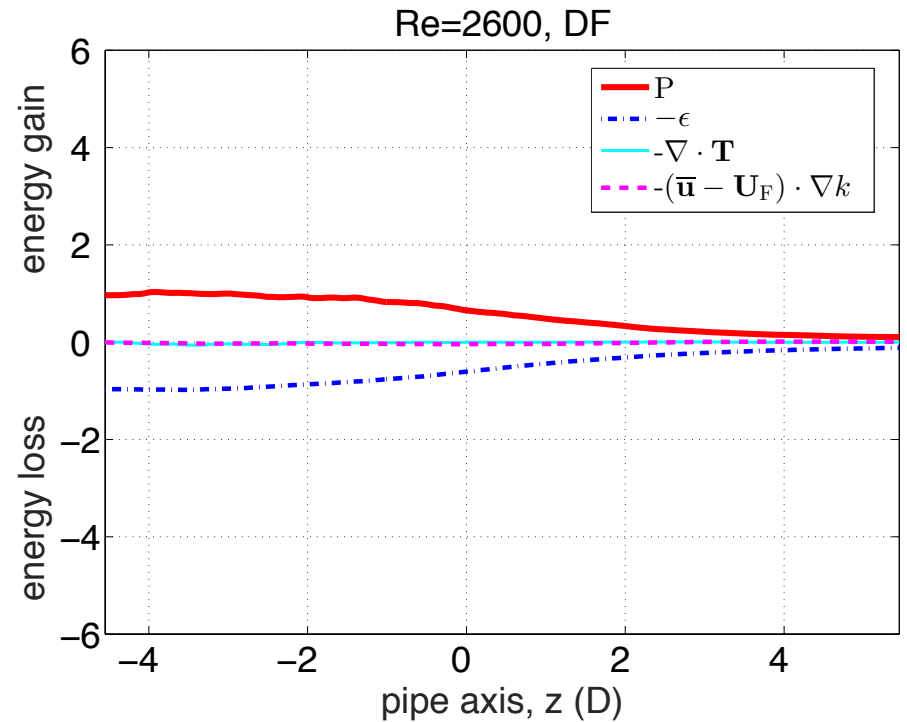
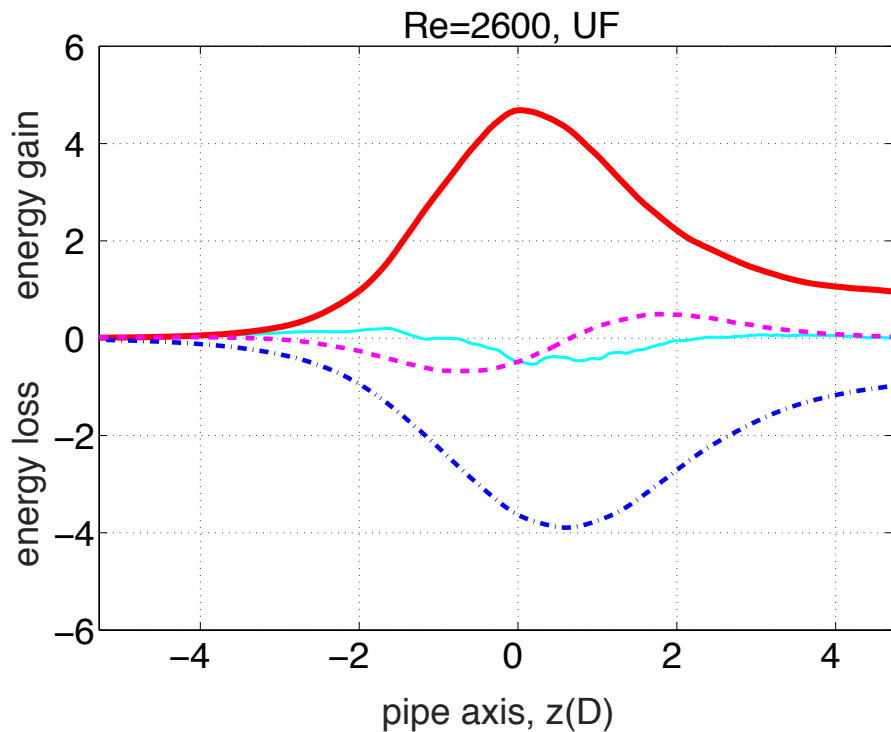
- Turbulent kinetic energy: $k = \frac{1}{2} \overline{\mathbf{u}' \cdot \mathbf{u}'}$

$$\frac{\bar{D}k}{\bar{D}t} + \nabla \cdot \mathbf{T} = P - \epsilon$$

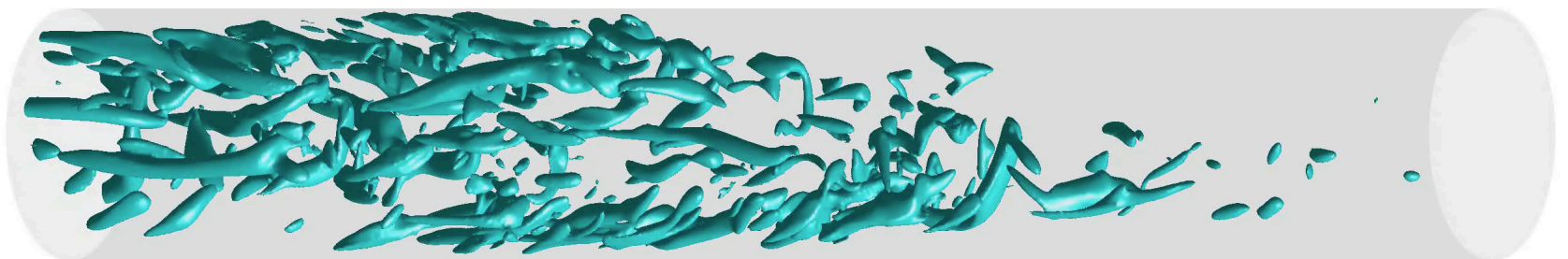
- Co-moving with the front:

$$\frac{\partial k}{\partial t} = P - \epsilon - (\bar{\mathbf{u}} - \mathbf{U}_F) \cdot \nabla k - \nabla \cdot \mathbf{T} = 0$$

Kinetic-energy budget at Re=2600

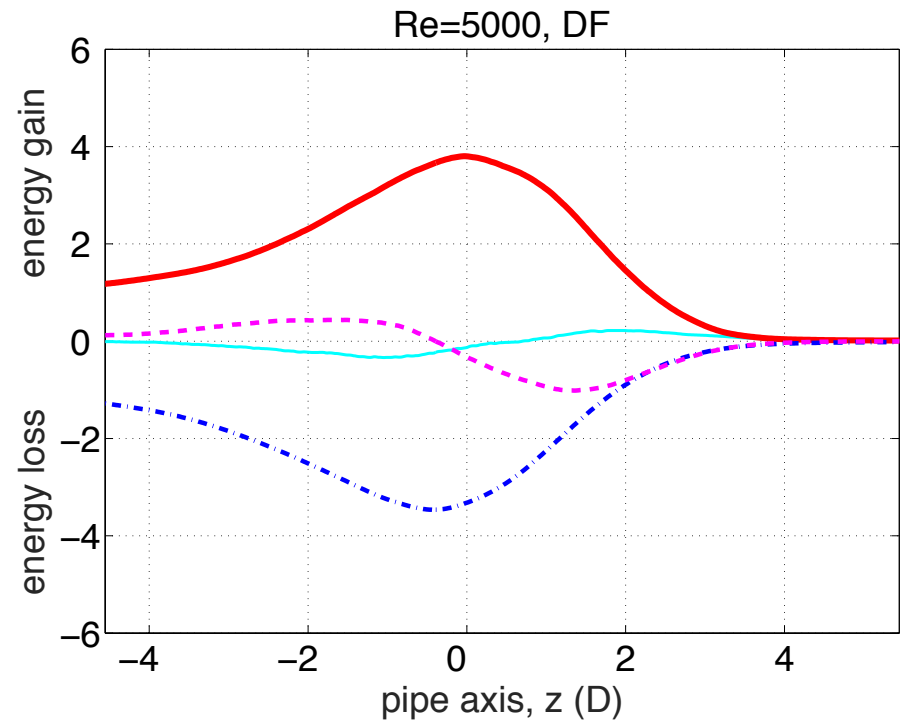
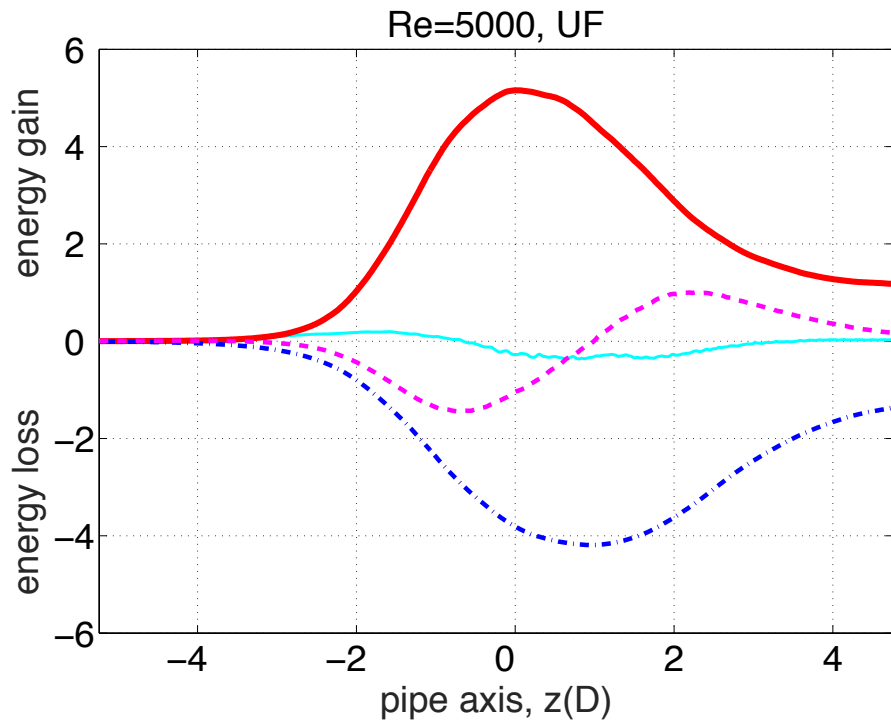


DF



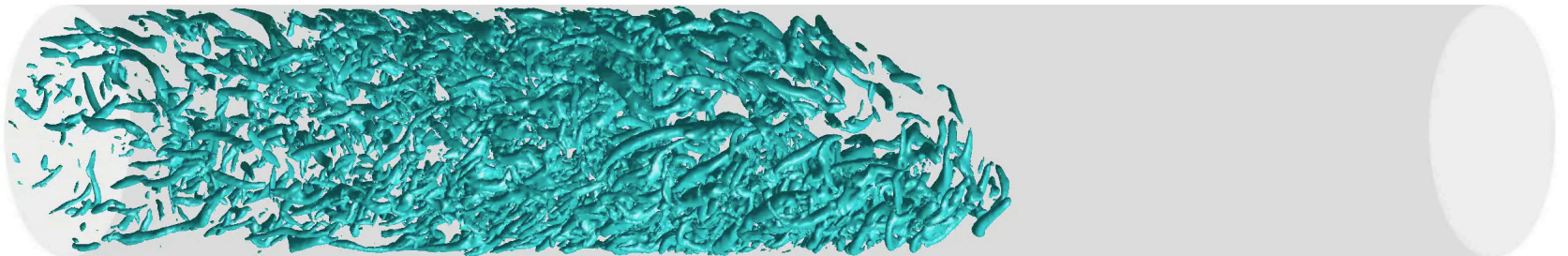
Song *et al.*, JFM (in press)

Kinetic-energy budget at $Re=5000$



UF

advection ←

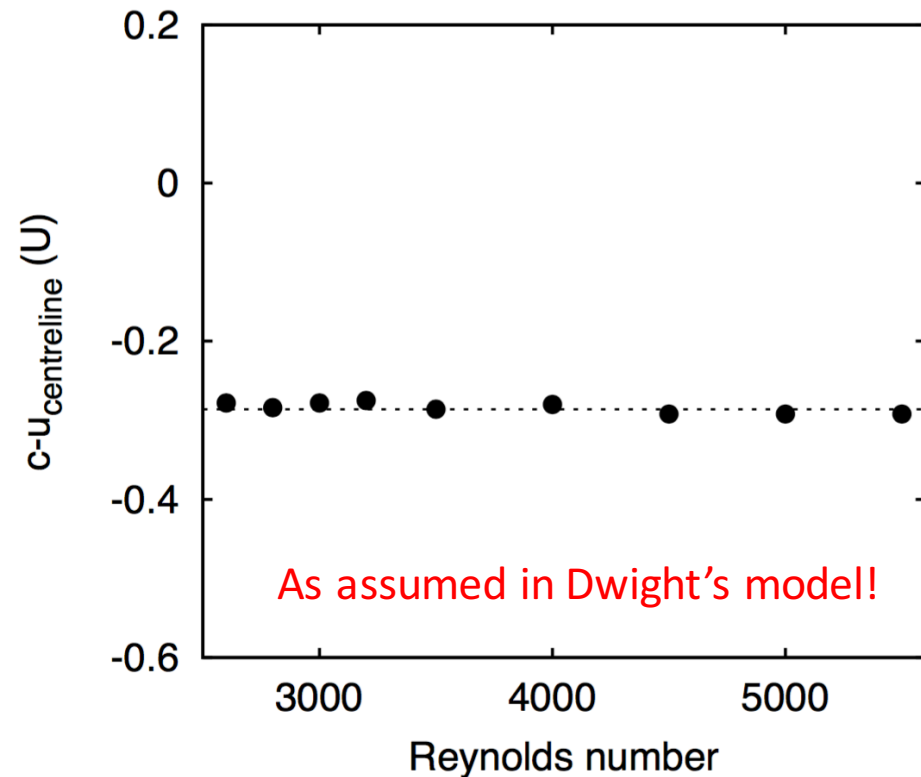
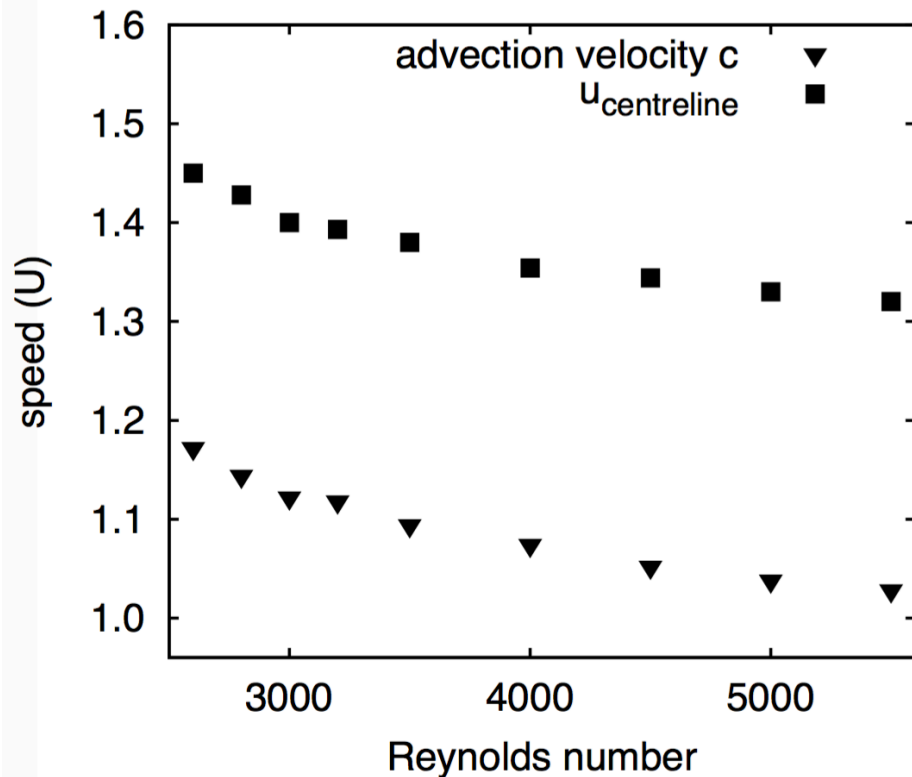


Song *et al.*, JFM (in press)

Mean turbulent advection speed

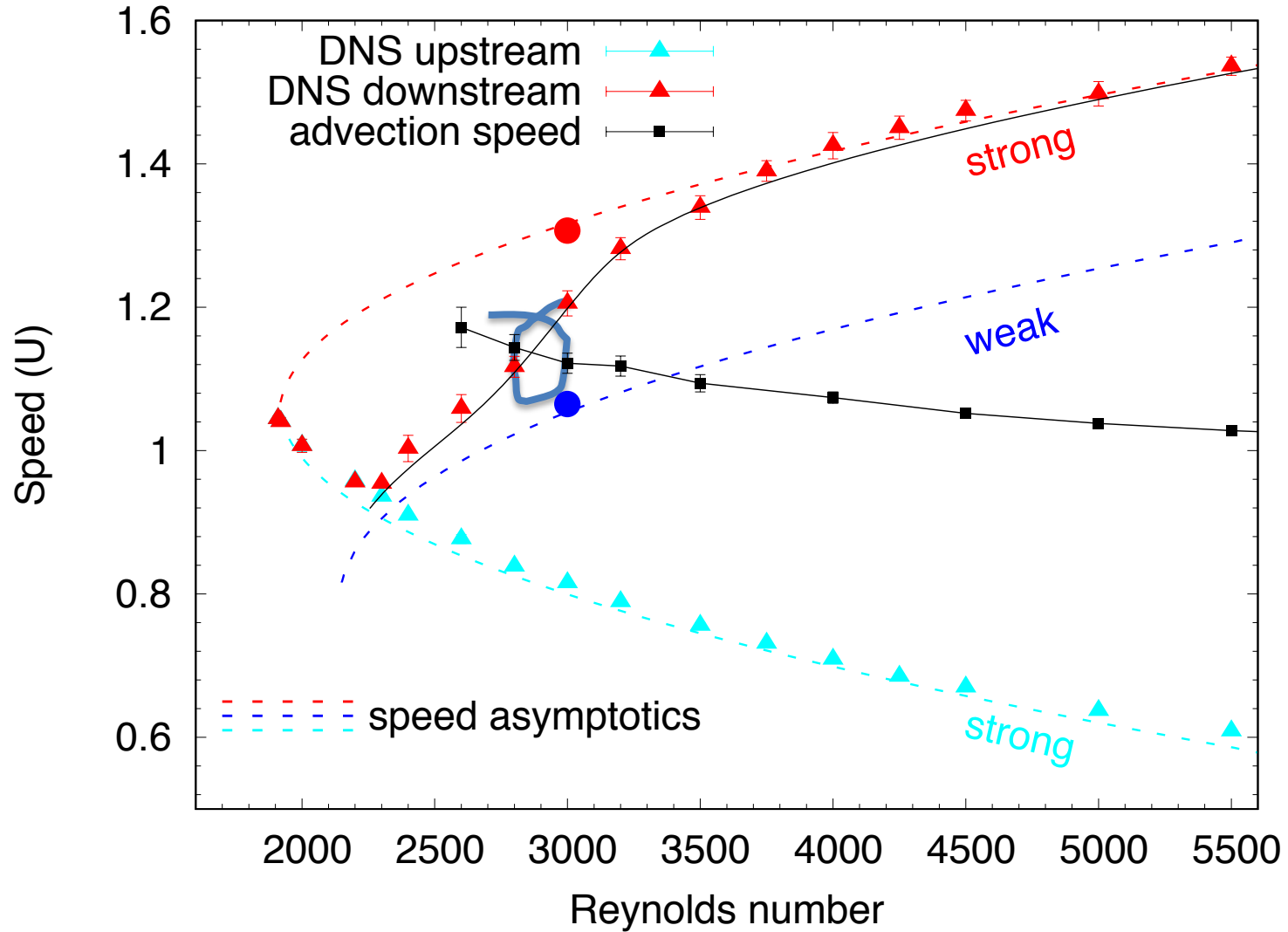
$$c = \frac{\langle \partial_z \mathbf{u}(r, \theta, z, t) \cdot \mathbf{f}(\mathbf{u}(r, \theta, z, t), p, t) \rangle}{\|\partial_z \mathbf{u}(r, \theta, z, t)\|^2}$$

Kreilos, Zammert & Eckhardt, JFM 2014

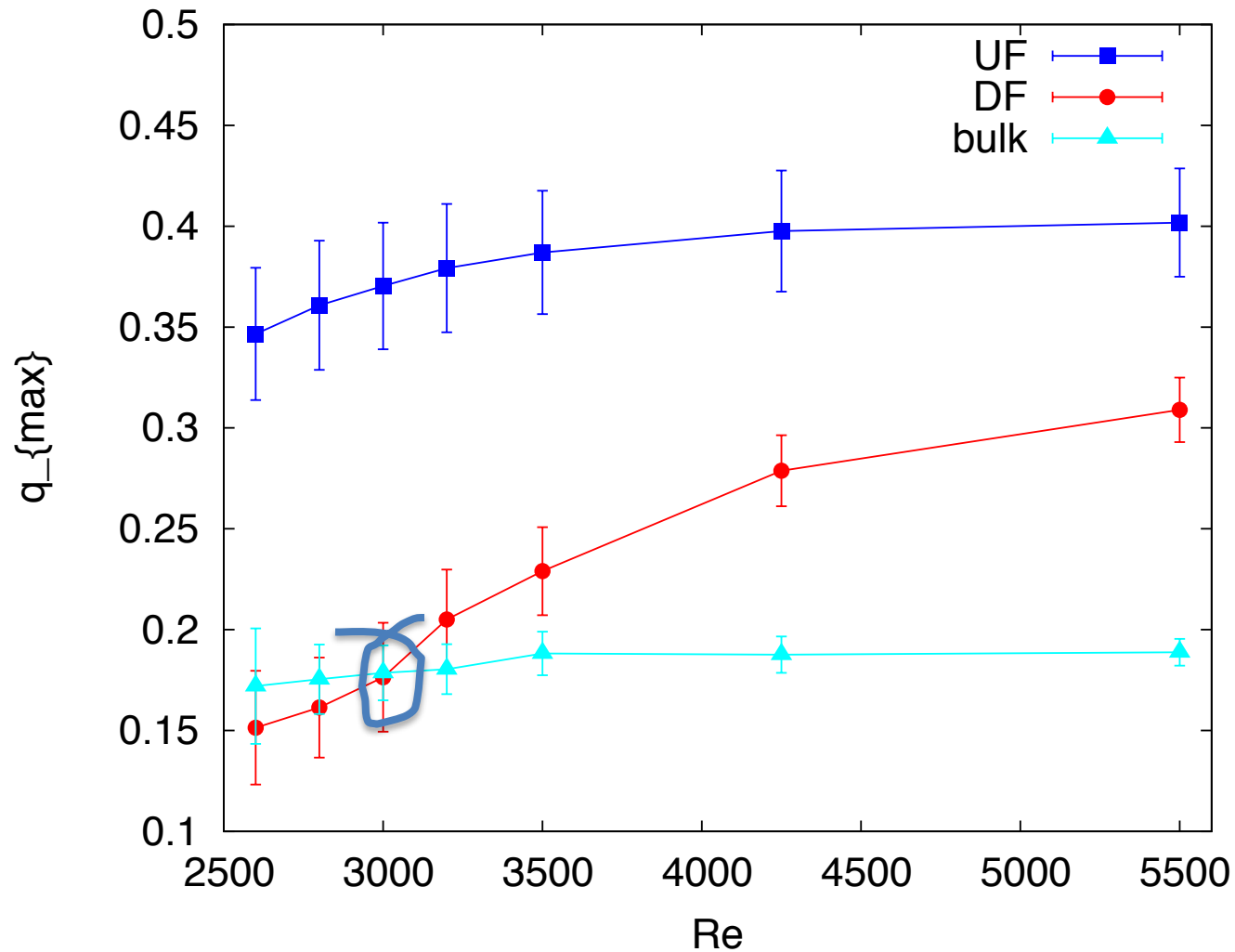


Song *et al.*, JFM (in press)

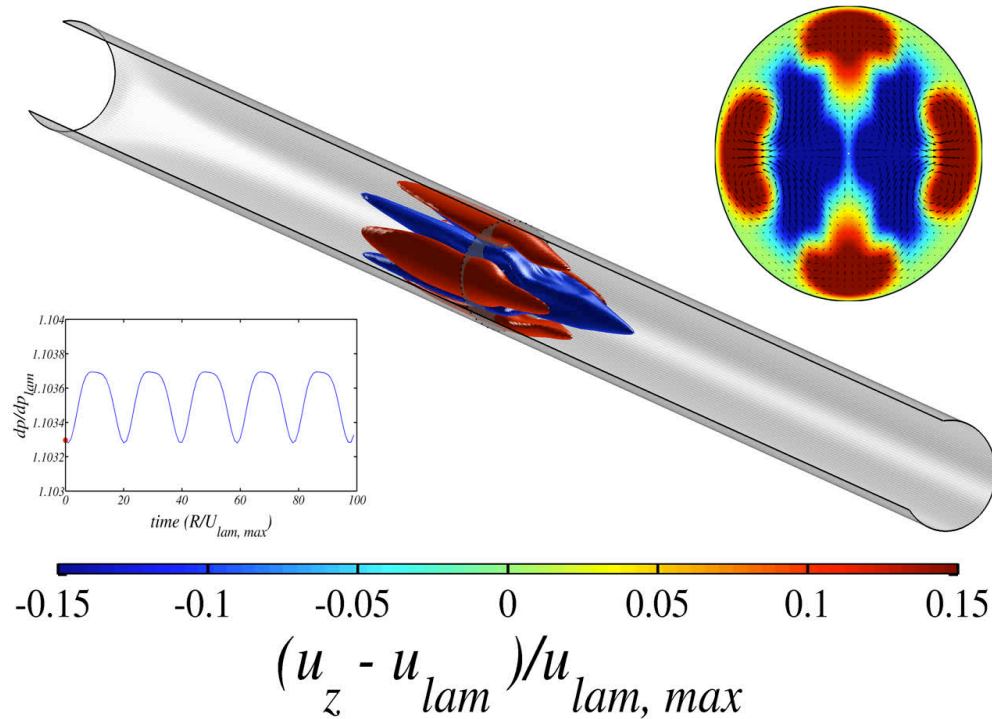
Crossover weak to strong fronts



Crossover weak to strong fronts



Fronts and ECS: localized RPO

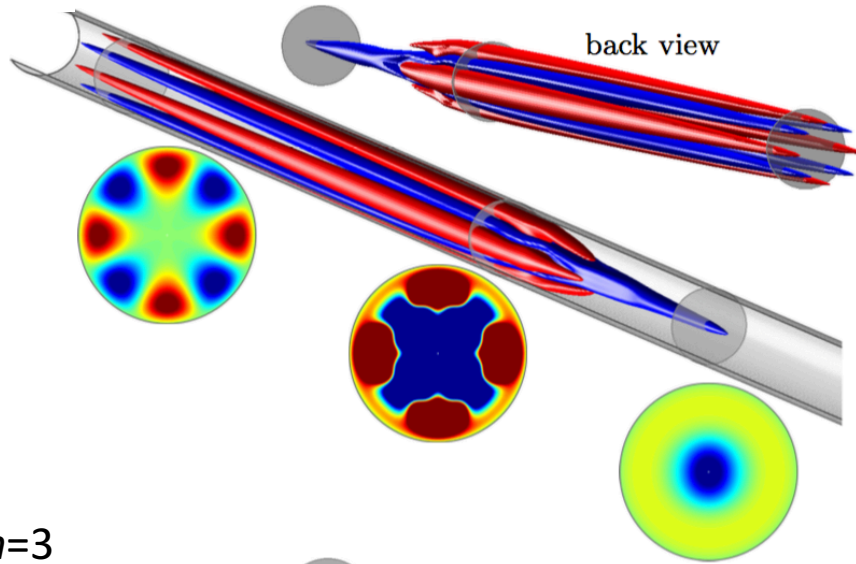


RPO* → torus → chaos → transient chaos → spatio-temporal chaos

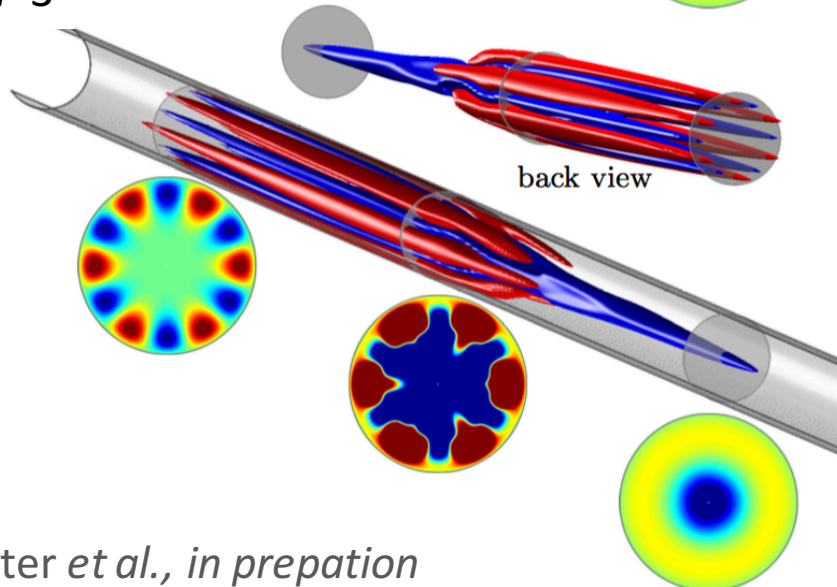
* Sub-harmonic bifurcation from TW: Chantry *et al.*, *PRL* 2014

Exponential localization

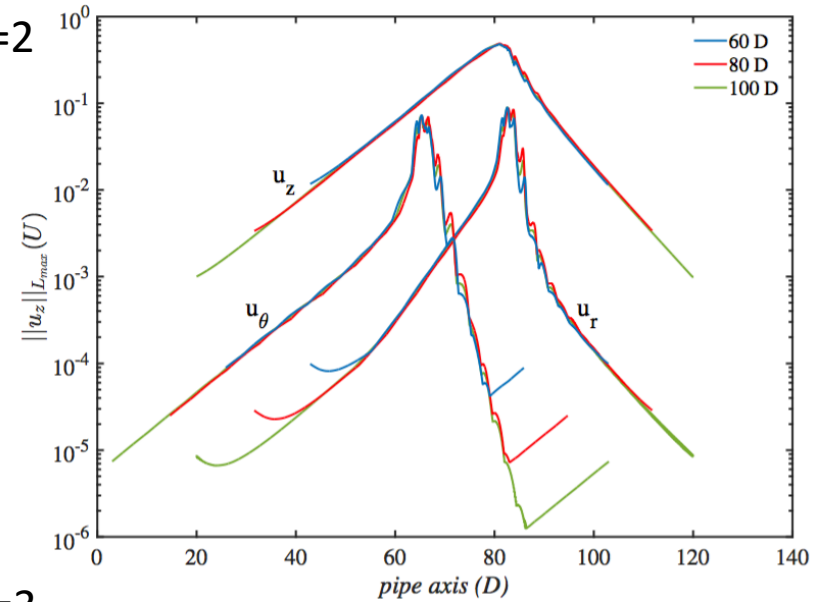
$m=2$



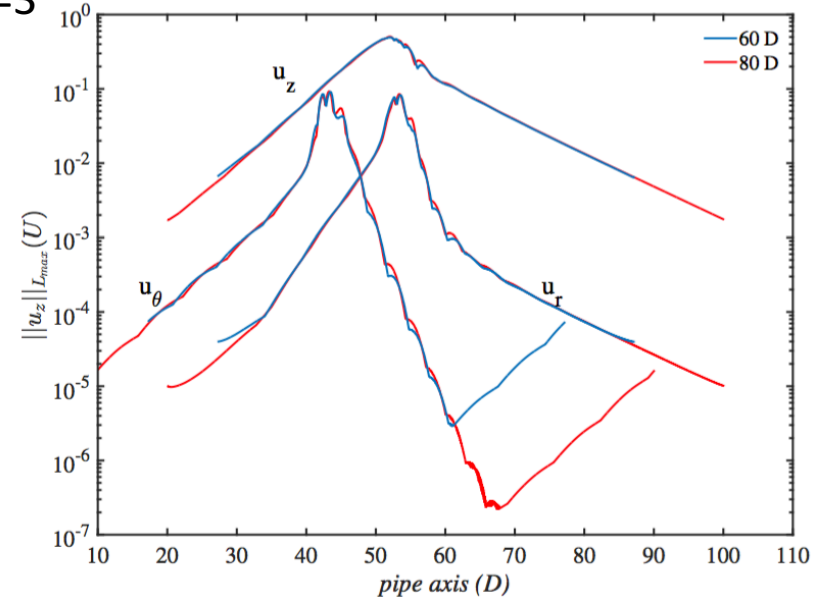
$m=3$



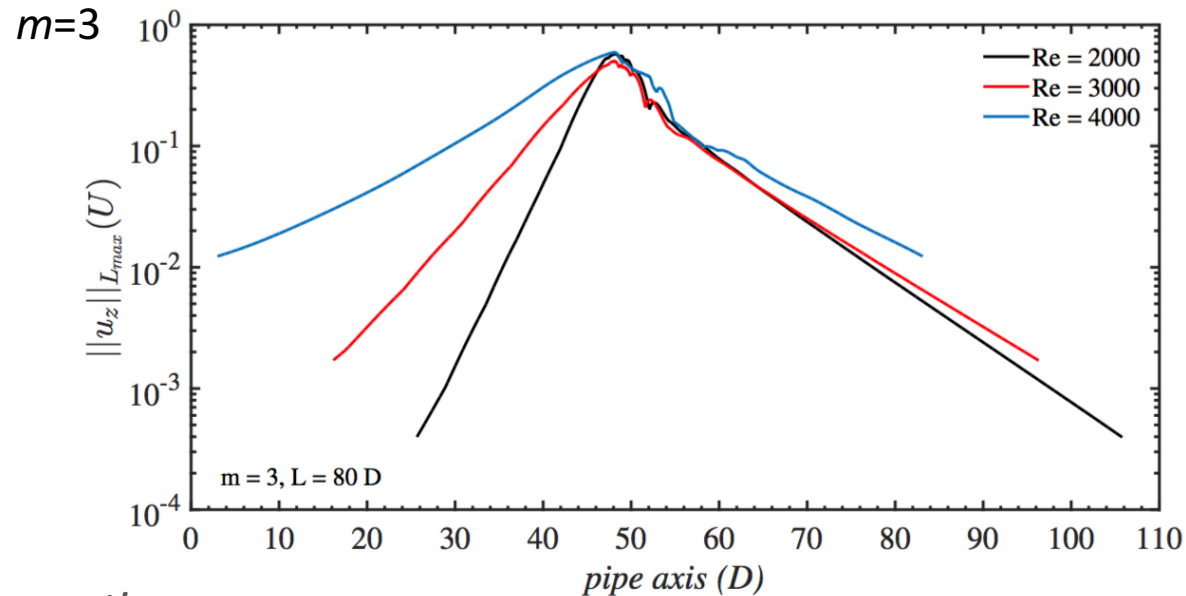
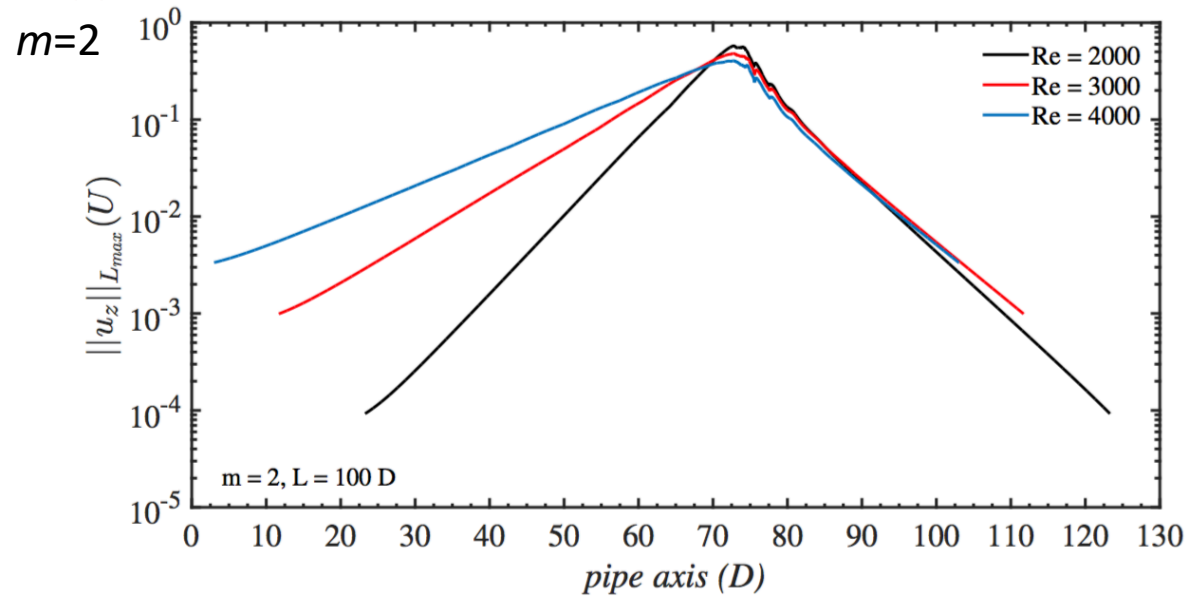
$m=2$



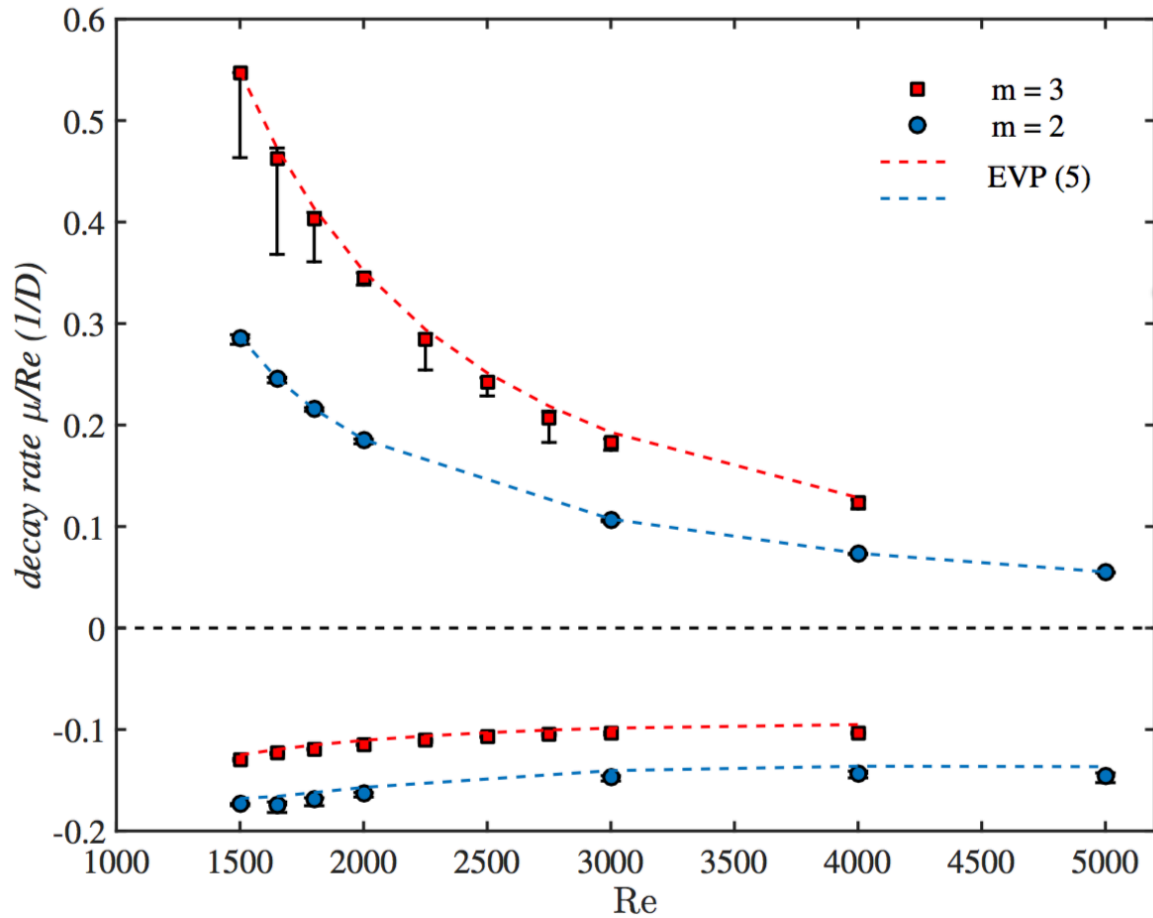
$m=3$



Localization: Re-dependence



Model vs. ECS

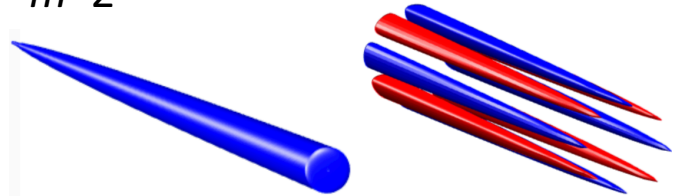


$$\mathbf{u} = \tilde{\mathbf{u}}(r) \exp \left[im\theta + \frac{\mu}{\text{Re}}(z - ct) \right]$$

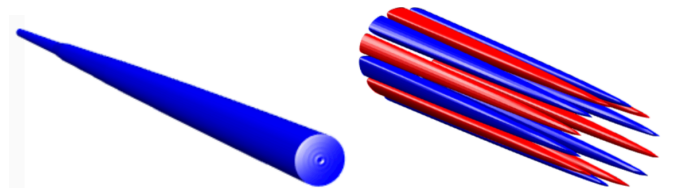
$$p = \tilde{p}(r) \exp \left[im\theta + \frac{\mu}{\text{Re}}(z - ct) \right]$$

Quadratic eigenvalue problem

$m=2$



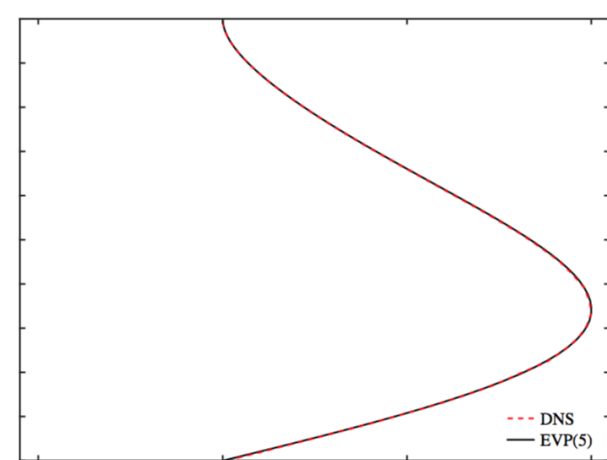
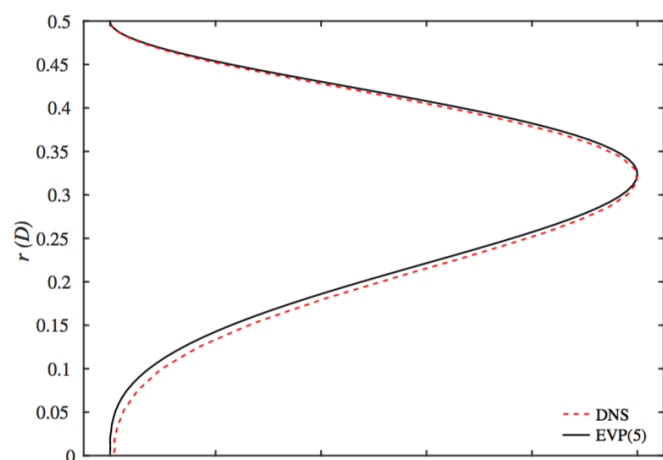
$m=3$



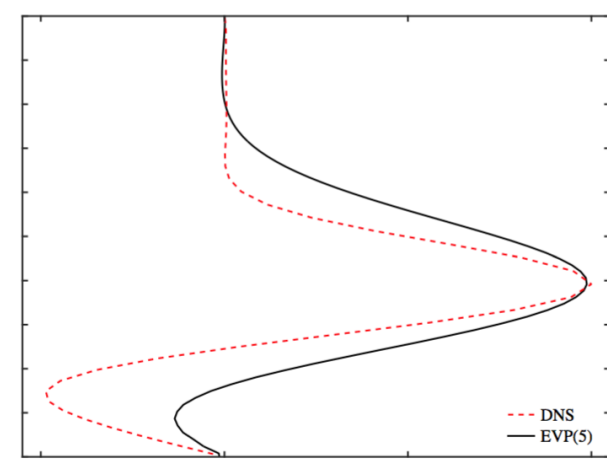
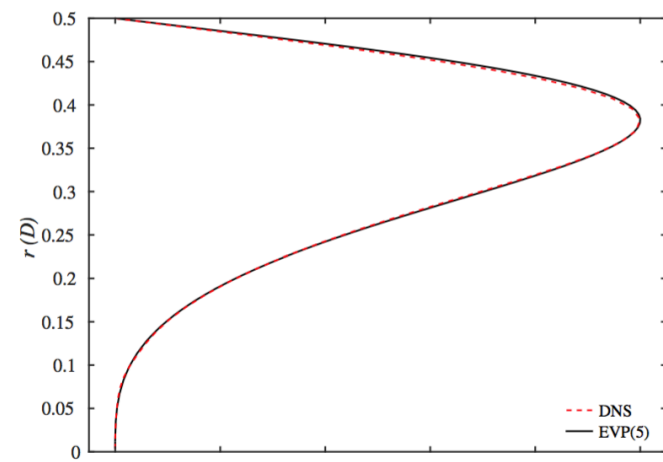
$m=2$
upstream

downstream

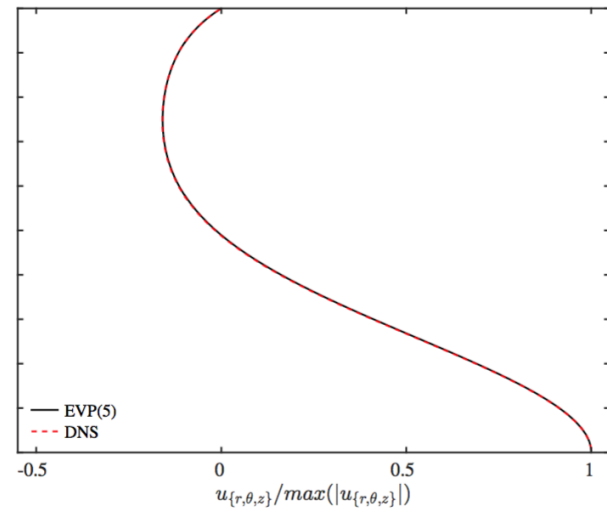
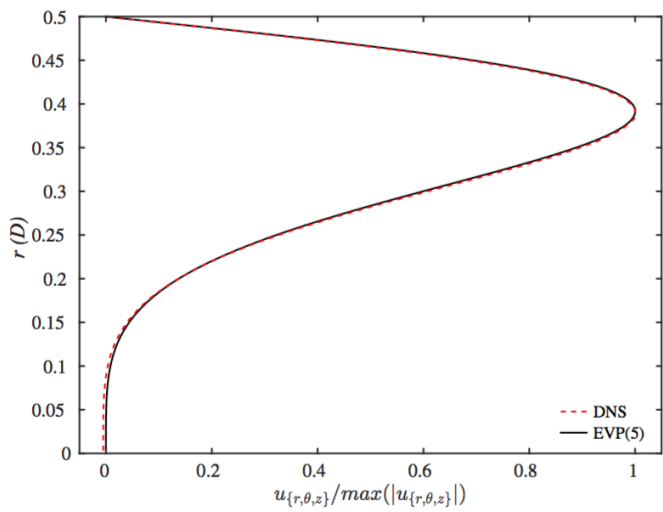
u_r



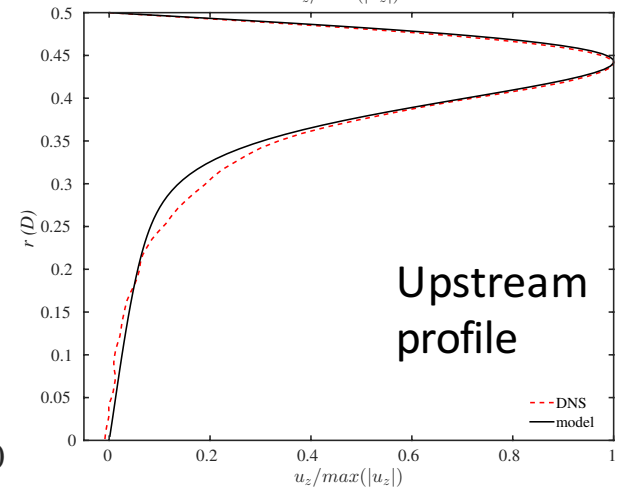
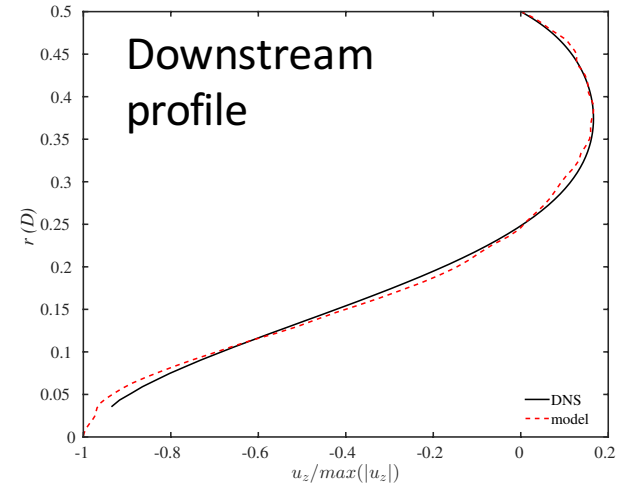
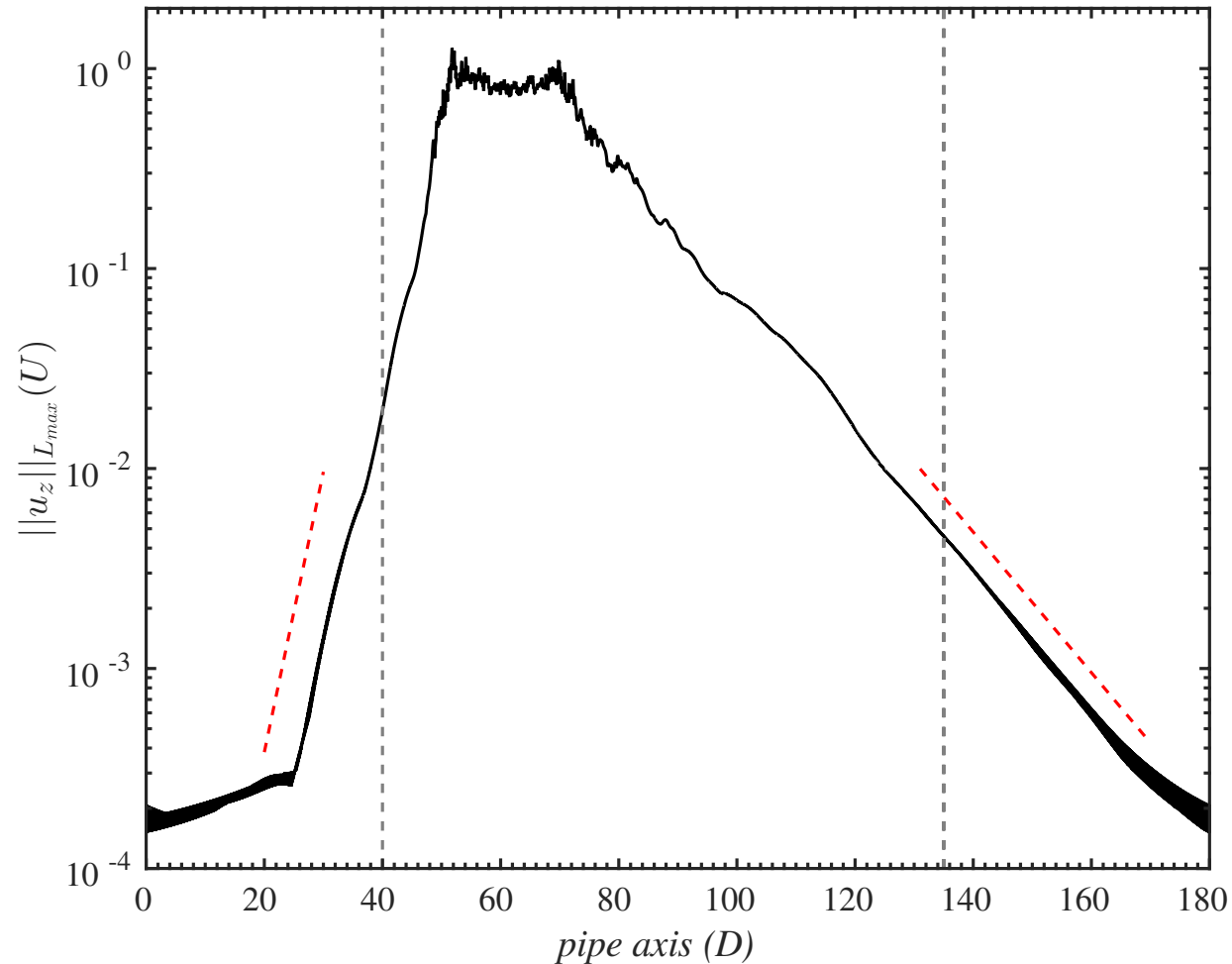
u_θ



u_z



Modeling fronts with linear model



Conclusions

- 1D-modeling (ARD) describes dynamics of interfaces
Barkley, Song, Mukund, Lemoult, Avila & Hof, *Nature* 2015
- Weak/strong fronts are physically distinct:
 - Weak: relaminarization at near (local) equilibrium
 - Strong: peak in production with delayed dissipation
 - Key role of mean advection speed of turbulence
Song, Barkley, Hof & Avila, *JFM* 2017 (in press)
- Fronts of ECS:
Ritter, Zammert, Eckhardt & Avila, in preparation

Further reading: Holzner, Song, Avila & Hof, *JFM* 2013:

“Lagrangian approach to laminar–turbulent interfaces in transitional pipe flow”