Homogeneous isotropic turbulence with polymer additives

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Work done in collaboration with

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References

- Manifestations of Drag Reduction by polymer additives in Decaying, Homogeneous, Isotropic Turbulence, <u>P. Perlekar</u>, D. Mitra, and R. Pandit, Phys. Rev. Lett., **97**, 264501 (2006).
- Numerical studies of three dimensional turbulence with polymer additives and two dimensional turbulence in thin films, <u>Prasad Perlekar</u>, Ph.D. Thesis, Indian Institute of Science, Bangalore,
- Direct numerical simulations of statistically steady, homogeneous, isotropic fluid turbulence with polymer additives, <u>P. Perlekar</u>, D. Mitra, and R. Pandit, Phys. Rev. E., **82**, 066313 (2010).
- Statistics of polymer extensions in turbulent channel flow, F. Bagheri, D. Mitra, P. Perlekar, and L. Brandt, arXiv:1011.3766.

Outline

- Overview of turbulence with polymers.
- Modelling polymer solutions.
- Direct numerical simulations(DNS): Decaying and Forced Turbulence.
- Conclusions.

Drag Reduction

 Toms (1946): Monochlorobenzene with 0.25% (by weight) of polymethylmethacrylate



- Reduction in the pressure gradient across the pipe, on the addition of polymers, for the same volumetric flow rate
- ► Drag Reduction(in percentage) $DR \equiv \left(\frac{\Delta P_s \Delta P_p}{\Delta P_s}\right) \times 100$

Reduction of small scale structures

 Turbulent jet of water with 50ppm polyethylene oxide at *Re* ~ 225
 [Turbulence structure in a water jet discharging in air, J.W. Hoyt and J.J. Taylor, Phys. Fluids, **20**, *S*253 (1977).]



Energy spectra



[Effect of polymer additives on the small-scale structure of grid-generated turbulence, W.D. McComb, J. Allan, and C.A. Greated, Phys. Fluids, **20**, 873 (1977).]

- Grid Reynolds number $Re_M = 7.6 \times 10^3$;
- For low polymer concentrations (50 and 100 ppm) there is no significant change in the energy spectrum; at somewhat higher concentrations(500 and 1000ppm) the spectra fall more steeply.

Eigenvalues of the strain tensor

[A. Liberzon, et al., Phys. Fluids, 17, 031701 (2005).]



- Length:140mm, Width:120mm, Disk Dia.:40mm, Observation volume:10 × 10 × 10mm, Re_λ = 38.
- Regions of large strains reduced on the addition of polymers.

Structure function: $S_2(r)$

N.T. Ouellette, H. Xu, and E. Bodenschatz, ICTR website, (2007).



- $c = 5ppm, Re_{\lambda} = 290, Wi = 3.5,$
- Small scale structures are modified on the addition of polymers.

Typical drag-reducing polymer: Polyethylene oxide N×[-CH₂-CH₂-O-]

- \blacktriangleright Degree of polymerization (N) $\simeq 10^4$
- Molecular weight \simeq 4 imes 10⁶ amu
- Zimm relaxation time $\simeq 10^{-4} s$
- RMS end-to-end distance at maximal extension $\simeq 34 \mu m$

Modelling polymer solutions

Navier-Stokes(NS) with Polymer Additives:
 3D, unforced, incompressible, NS with additional stress because of polymers:

$$\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla)\mathbf{u} = -\nabla p + \nu \nabla^2 \mathbf{u} + \nabla \cdot \boldsymbol{\mathcal{T}},$$

where

- u(x, t): fluid velocity; point x; time t;
- ν : Kinematic viscosity of the fluid;
- \mathcal{T} : polymer contribution to the fluid stress;

 $\nabla \cdot \mathbf{u} = \mathbf{0}$ enforces incompressibility.

Modelling polymer solutions

Finitely Extensible Nonlinear Elastic-Peterlin(FENE-P) model

$$rac{\partial \mathcal{C}_{lphaeta}}{\partial t} + (u_\gamma\partial_\gamma)\mathcal{C}_{lphaeta} ~=~ (\partial_\gamma u_lpha)\mathcal{C}_{\gammaeta} + \mathcal{C}_{lpha\gamma}(\partial_\gamma u_eta) - rac{1}{\mu}\mathcal{T}_{lphaeta}.$$

["Dynamics of polymeric liquids", Bird, et al.]

- $c = \mu/(\nu + \mu)$; $c = 0.1 \simeq 100 ppm$ of PEO
- $We = \tau_{poly} \sqrt{(\epsilon(t_m)/\nu)}$; t_m is the time corresponding to the peak in ϵ for c = 0

[Vaithianathan, et al., JCP, 187, 1 (2003).]

Direct Numerical Simulations

Direct Numerical Simulation(DNS)

Solve NS and FENE-P numerically

$$\begin{aligned} \frac{\partial u_{\alpha}}{\partial t} + (u_{\gamma}\partial_{\gamma})u_{\alpha} &= -\partial_{\alpha} p + \nu \partial_{\gamma\gamma} u_{\alpha} + \partial_{\gamma} \mathcal{T}_{\alpha\gamma}, \\ \partial_{\gamma} u_{\gamma} &= 0, \end{aligned}$$

$$rac{\partial \mathcal{C}_{lphaeta}}{\partial t} + (u_\gamma\partial_\gamma)\mathcal{C}_{lphaeta} ~=~ (\partial_\gamma u_lpha)\mathcal{C}_{\gammaeta} + \mathcal{C}_{lpha\gamma}(\partial_\gamma u_eta) - rac{1}{\mu}\mathcal{T}_{lphaeta}.$$

Decaying Turbulence

Results: Initial Condition

- Start from an initial energy spectrum with energy concentrated in the first few Fourier modes and the polymers unstretched
- Monitor the decay of the energy dissipation rate and the energy spectrum for the fluid with and without polymer additives.

Energy Dissipation Rate

$$N = 256$$
, $\nu = 10^{-3}$, $\tau_{poly} = 1$



- The energy dissipation rate ε(t) as a function of time t for different values of c.
- The peak in $\epsilon(t)$ decreases as c increases.

Dissipation Reduction(DR)

$$N = 96, \nu = 10^{-2}$$



- ► Natural definition of dissipation-reduction % $DR = \left(\frac{\epsilon^{f,m} - \epsilon^{\rho,m}}{\epsilon^{f,m}}\right) \times 100;$
- f and p stand, respectively, for the fluid without and with polymers.
- An increase in c enhances the dissipation reduction DR (cf., earlier shell-model study).
- ► DR decreases marginally with an increase in We.

Fluid energy spectrum

$$N=192,~
u=10^{-2},~ au_{poly}=1$$



- ► $E_f(k) = \sum_{k-1/2 < k' < k+1/2} |\mathbf{u}(k')|^2$ at t_m for polymer concentrations c = 0(o-), c = 0.1(-) and c = 0.4(-).
- Energy spectrum at cascade completion changes significantly for large Fourier modes.
- This had not been resolved by earlier, high-Re simulations!

Scale-dependent viscosity

$${\sf N}=192,~
u=10^{-2},~ au_{poly}=1$$



► The change in the spectra and ϵ can be understood in terms of an additional, effective, scale-dependent viscosity $\Delta \nu(k) \equiv -\mu \sum_{k=1/2 < k' \le k+1/2} \mathbf{u}_{k'} \cdot (\nabla \cdot \mathcal{J})_{-k'} / [\tau_{poly} k'^2 E^{p,m}(k')].$

• Since $\Delta \nu$ becomes negative, polymers pump energy into the fluid around $k \simeq 10$.

Order-*p* equal-time, longitudinal velocity structure function.

$$\begin{split} \mathcal{S}_{p}(r) &\equiv \langle \delta u(r,t)^{p} \rangle, \\ \delta u_{||}(r,t) &\equiv \left[\vec{u}(\vec{x}+\vec{r},t) - \vec{u}(\vec{x},t) \right] \cdot (\vec{r}/r). \end{split}$$

Second order structure function $S_2(r)$

Experiments(Ouellette et al.)



Our DNS

Figure: c = 5ppm, $Re_{\lambda} = 290$, and We = 3.5

Figure: N = 128, $\nu = 0.01$, and $\tau_P = 1.5$

PDF of $|\omega|$



- ▶ Probability distribution of the modulus of the vorticity(P(|ω|)) at cascade completion(c=0, c=0.4).
- Addition of polymers leads to a decrease in the regions of large vorticity.

Isosurfaces of $|\omega|$

$$N = 256, \ \nu = 10^{-3}, \ \tau_{poly} = 1$$



- ▶ Iso- $|\omega|$ surfaces for $|\omega| = \langle |\omega| \rangle + 2\sigma$ for c = 0(left) and c = 0.4(right) at t_m .
- Small-scale structures are suppressed on the addition of polymers.

Stretching of Polymers:Cumulative distribution(CDF)

$${\it N}=256$$
, $u=10^{-3}$, $au_{\it poly}=1$



- ► c = 0.1(dashed line), c = 0.4(line).
- ▶ An increase in *c* leads to a decrease in the polymer extension.

Summary of Results: Decaying turbulence

- ▶ Polymer additives lead to a decrease in small-scale structures.
- Polymers decrease the energy of the turbulent fluid at intermediate length scales and increase it at small scales.
- Dissipation reduction is the analogue in homogeneous, isotropic turbulence of drag-reduction in wall-bounded turbulence.
- An effective scale-dependent viscosity leads to a natural explanation of our results.
- This points toward an increase in the effective viscosity, but one that is scale-dependent.
- Ref: Manifestations of Drag Reduction by polymer additives in Decaying, Homogeneous, Isotropic Turbulence, <u>P. Perlekar</u>, D. Mitra, and R. Pandit, Phys. Rev. Lett., **97**, 264501 (2006).

Forced Turbulence

Deterministic forcing of M.A. Taylor, S. Kurien, and G. Eyink, Phys. Rev. E, 68, 26310, (2003).

Time evolution of E and ϵ

 $N = 256, Re_{\lambda} \simeq 80, c = 0.1$



- Time averaged E decreases with an increase in We
- Time averaged ϵ decreases with an increase in We
- We = 3.5 (blue circles), We = 7.1 (black dashed line), NS (red)

PDF of $|\omega|$ and ϵ_{loc}

$${\it N}=256,~{\it Re}_\lambda\simeq 80,~c=0.1$$



•
$$\omega \equiv |\sqrt{\sum_{i,j} \omega_{ij} \omega_{ij}}|, \ \epsilon_{loc} = \nu s^2 \equiv \sum_{i,j} s_{ij} s_{ij},$$

 $s = (\nabla \mathbf{u} + (\nabla \mathbf{u})^T)/2, \ \omega = \nabla \times \mathbf{u}$

 Regions of large strain and vorticity decrease on the addition of polymers PDF of $|\omega|$ and ϵ_{loc}

 $N=256,~Re_\lambda\simeq 80,~c=0.1$



•
$$\omega \equiv |\sqrt{\sum_{i,j} \omega_{ij} \omega_{ij}}|, \ \epsilon_{loc} = \nu s^2 \equiv \sum_{i,j} s_{ij} s_{ij}$$

 $s = (\nabla \mathbf{u} + (\nabla \mathbf{u})^T)/2, \ \omega = \nabla \times \mathbf{u}$

 Regions of large strain and vorticity decrease on the addition of polymers

Isosurfaces of $|\omega|$

 $N=256,~Re_\lambda\simeq 80,~c=0.1$



- ▶ Iso- $|\omega|$ surfaces for $|\omega| = \langle |\omega| \rangle + 2\sigma$ for c = 0(left) and c = 0.1, We = 7.1(right).
- Small-scale structures are suppressed on the addition of polymers.

QR plots

 $N=256,~Re_{\lambda}\simeq 23,~We=7.1,~c=0.1$



• Left: NS; Right: Polymer (c = 0.1, We = 7.1)

Energy spectrum

(Left)
$$N = 256$$
, $Re_{\lambda} = 80$
(Right) $N = 512$, $Re_{\lambda} = 20$



Polymer extensions

 $N=256,~Re_{\lambda}\simeq23,~c=0.1$



► We = 7.1, c = 0.1(line); We = 3.5, c = 0.1(dashed line);

- Polymer extensions larger in comparison to decaying turbulence
- At fixed c, polymer extension increases with an increase in We

Conclusions

- Our simulations show that the addition of polymers to flows that display homogeneous isotropic turbulence leads to dissipation reduction in both decaying and statistically steady turbulence; this dissipation reduction is the analogue of drag reduction in wall-bounded flows.
- Our numerical results agree with the experimental results of

 (a) Liberzon, et al., op. cit. and (b) Ouellette et al., op. cit.
- Polymers decrease the energy of the turbulent fluid at intermediate length scales and increase it at small scales; a scale-dependent viscosity provides a natural means of understanding our results.