

# Kinetic equation approach to describe dynamics of irradiated samples

**Beata Ziaja**

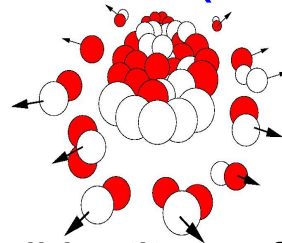
CFEL, DESY and INP, Kraków

in collaboration with

C. Bostedt (SLAC), H. Chapman (DESY),

S. Hau-Riege (LLNL)

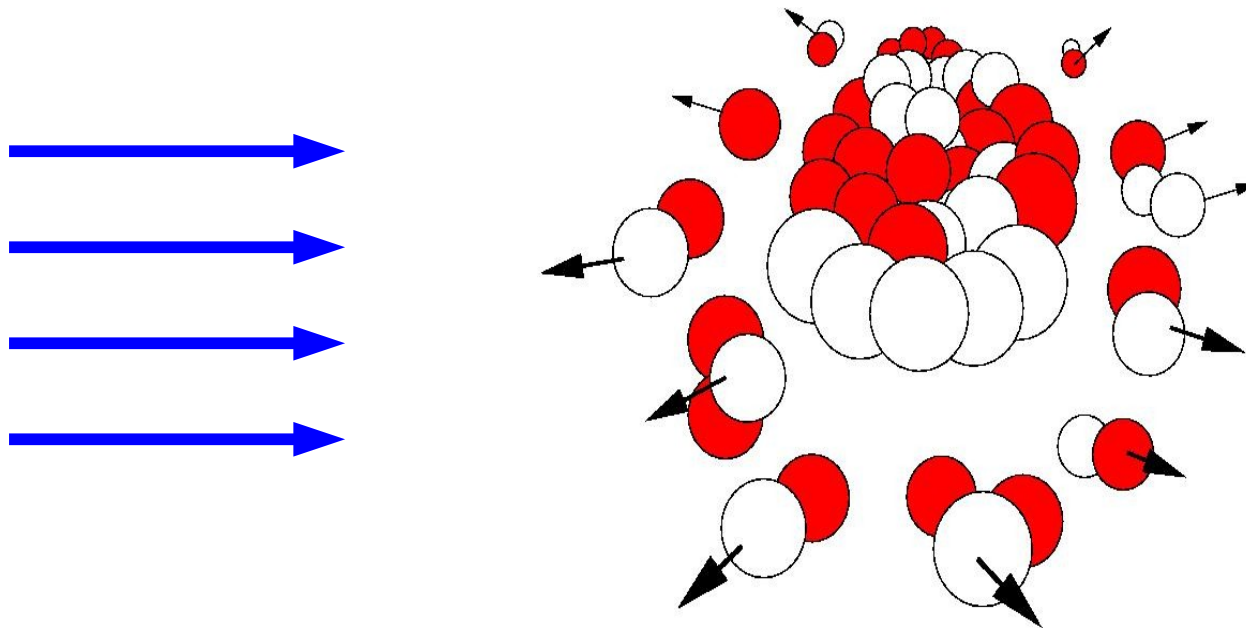
T. Laarmann, R. Santra, F. Wang, E. Weckert (DESY)  
and T. Möller (TU Berlin)



X-Ray Science in the 21<sup>st</sup> Century, Kavli Institute of Theoretical Physics, 2-6 August 2010

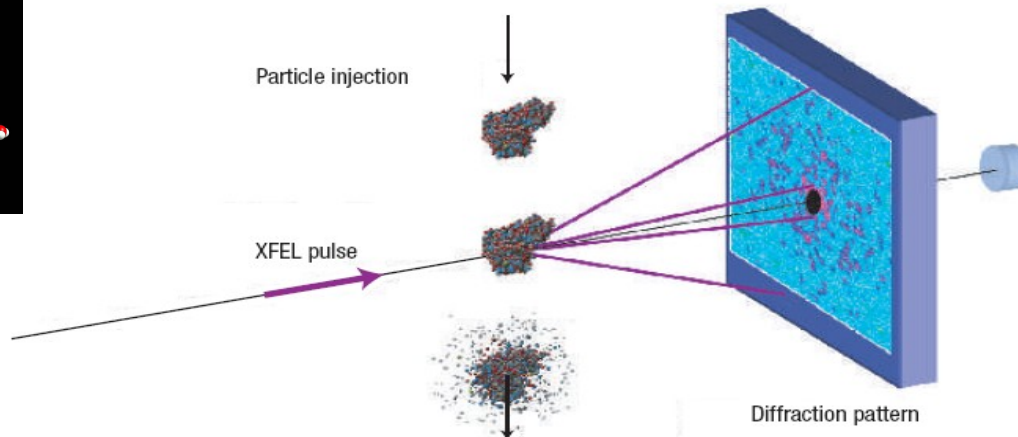
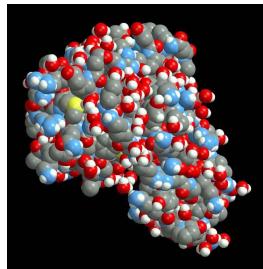
# Motivation for theoretical study

- Important for:
- (i) experiments with FEL → cluster experiments, single particle imaging, warm dense matter etc.
  - (ii) construction of the laser → FEL optics,
  - (iii) test of various theoretical models

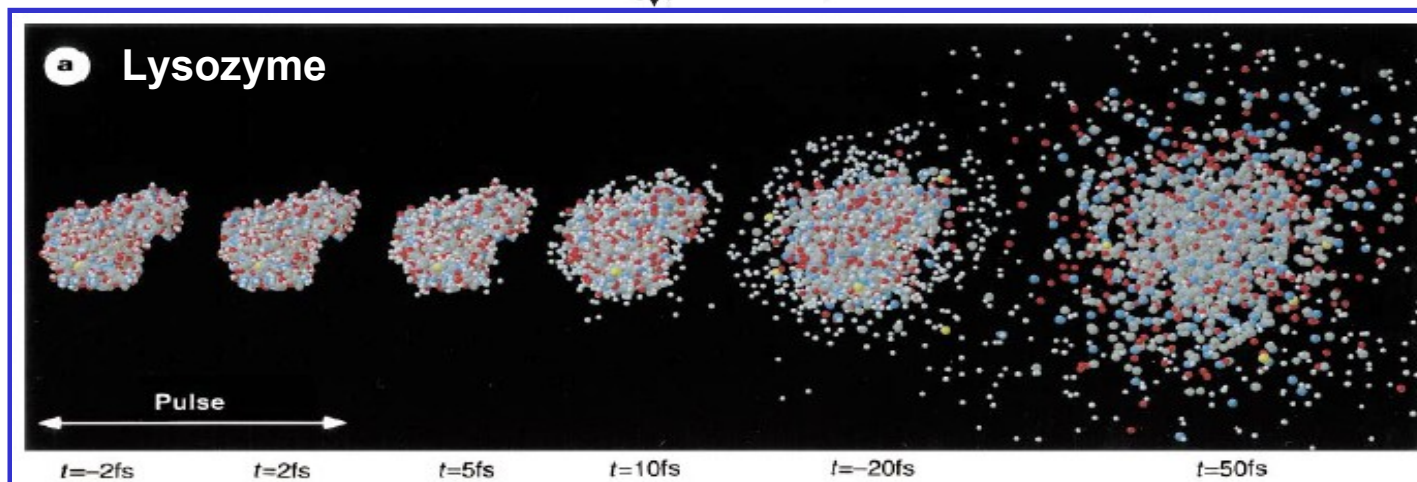
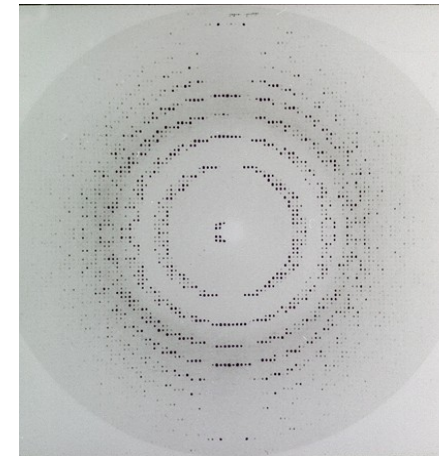


# Structure determination through single particle diffraction imaging?

## Molecules at atomic resolution



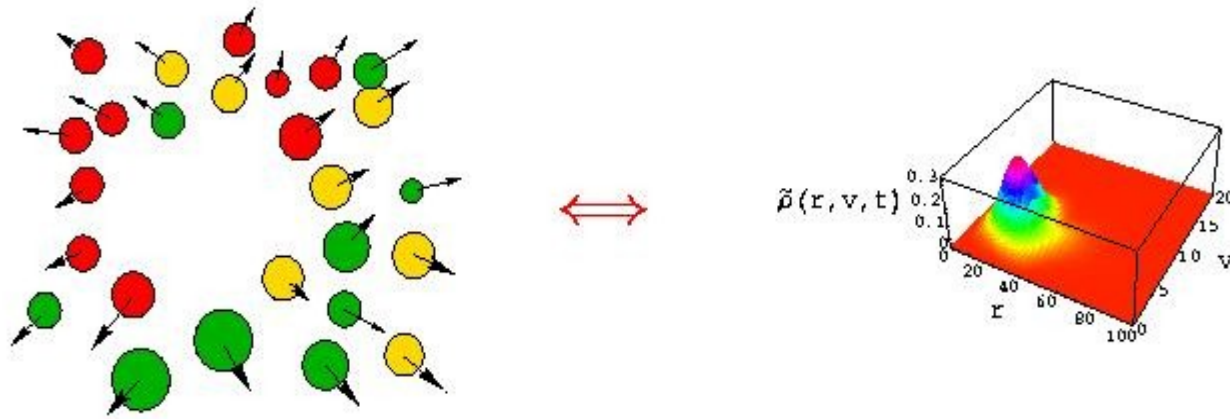
## Crystal



R. Neutze,  
R. Wouts,  
D. van der  
Spoel,  
E. Weckert,  
J. Hajdu  
Nature 406,  
752 (2000)  
**Radiation  
damage  
and Coulomb  
explosion**

# Tool: statistical Boltzmann approach

Evolution of larger systems described in terms of statistical density function,  $\rho(r,v,t)$ , in phase space:

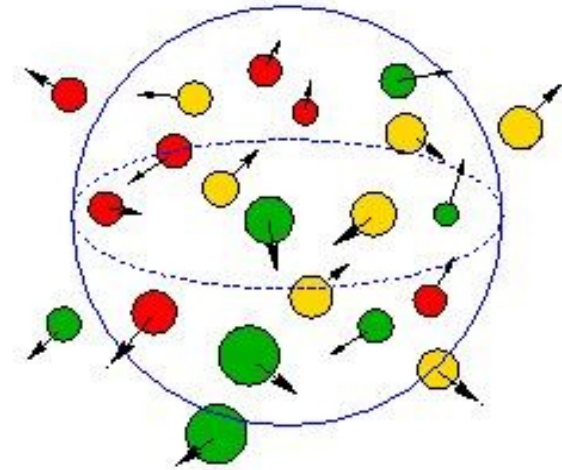


# Tool: statistical Boltzmann approach

- first-principle approach
- single-run method
- computational costs do not scale with number of atoms

Boltzmann equations are able to follow

non-equilibrium processes



Difficulty:

- requires advanced numerical methods



# Solving Boltzmann equations

The general coupled **Boltzmann equations** for electron,  $\rho^{(e)}(\mathbf{r}, \mathbf{v}, t)$ , and ion densities,  $\rho^{(i)}(\mathbf{r}, \mathbf{v}, t)$ , where  $i = 0, 1, \dots, N_J$  denotes the ion charge, and  $N_J$  is the maximal ion charge in the system are:

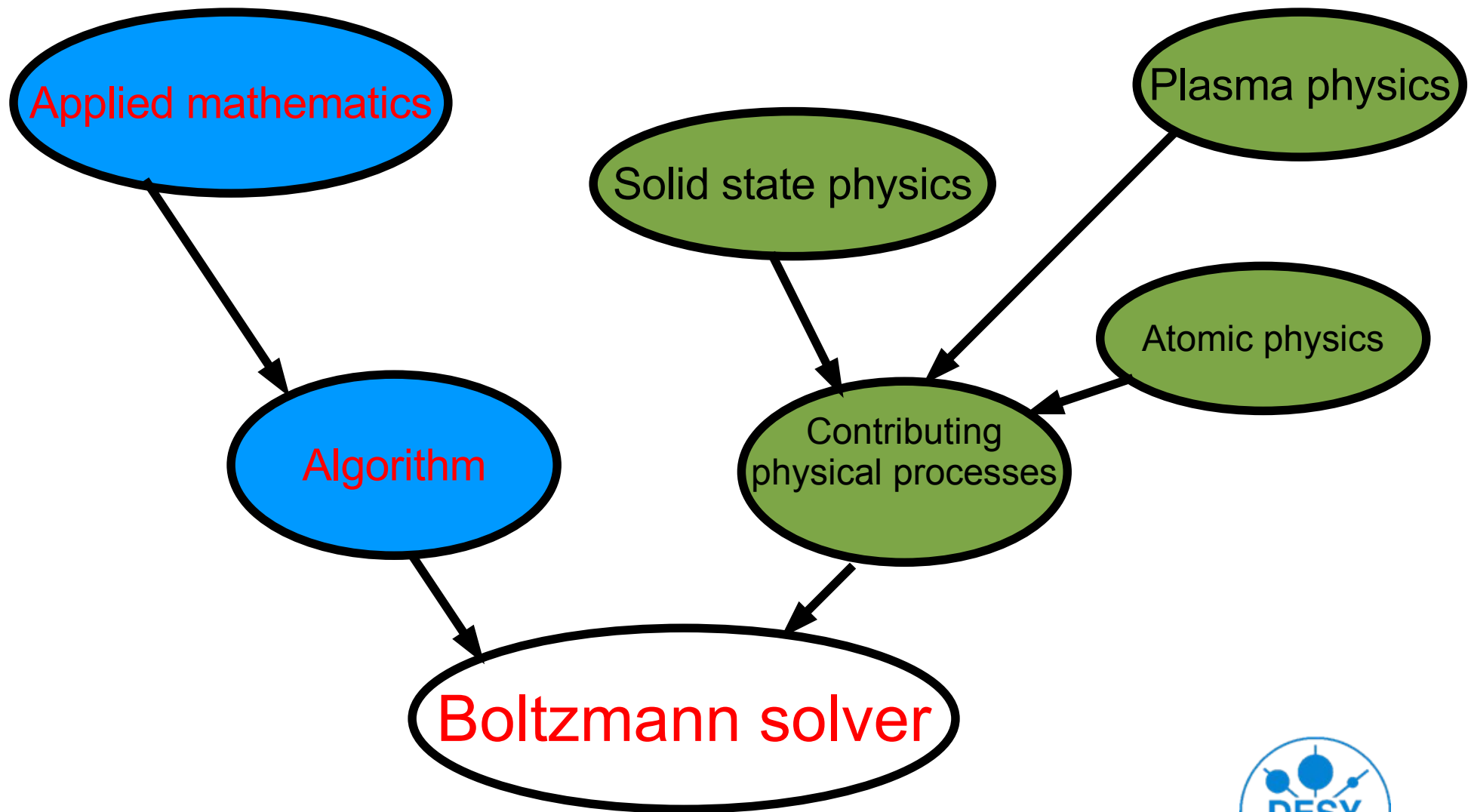
$$\partial_t \rho^{(e)}(\mathbf{r}, \mathbf{v}, t) + \mathbf{v} \cdot \partial_{\mathbf{r}} \rho^{(e)}(\mathbf{r}, \mathbf{v}, t) + \frac{e}{m} (\mathbf{E}(\mathbf{r}, t) + \mathbf{v} \times \mathbf{B}(\mathbf{r}, t)) \cdot \partial_{\mathbf{v}} \rho^{(e)}(\mathbf{r}, \mathbf{v}, t) = \Omega^{(e)}(\rho^{(e)}, \rho^{(i)}, \mathbf{r}, \mathbf{v}, t),$$

$$\partial_t \rho^{(i)}(\mathbf{r}, \mathbf{v}, t) + \mathbf{v} \cdot \partial_{\mathbf{r}} \rho^{(i)}(\mathbf{r}, \mathbf{v}, t) - \frac{ie}{M} (\mathbf{E}(\mathbf{r}, t) + \mathbf{v} \times \mathbf{B}(\mathbf{r}, t)) \cdot \partial_{\mathbf{v}} \rho^{(i)}(\mathbf{r}, \mathbf{v}, t) = \Omega^{(i)}(\rho^{(e)}, \rho^{(i)}, \mathbf{r}, \mathbf{v}, t).$$

These equations include the **total electromagnetic force** acting on ions and electrons. Collision terms,  $\Omega^{(e,i)}$ , describe the changes of the electron/ion densities of velocities  $(\mathbf{v}, \mathbf{v} + d\mathbf{v})$  measured at the positions  $(\mathbf{r}, \mathbf{r} + d\mathbf{r})$  with time. These changes are due to short-range processes, e. g. collisions, photoabsorptions. The **number of processes involved** in the sample dynamics depends on the radiation wavelength.



# Solving Boltzmann equations



# Boltzmann solver

- Investigates the non-equilibrium phase of evolution of an irradiated sample until thermalization of electrons and saturation of ionization is reached

It uses the angular moment expansion for density function,  $\rho$ :

$$\rho_e \sim \rho_0 + \rho_1 \cdot \cos \theta_{vr} + \rho_2 \cdot \cos \theta_{v\varepsilon}$$

with dominating isotropic component of  $\rho$ . This is appropriate for the non-equilibrium phase of sample evolution

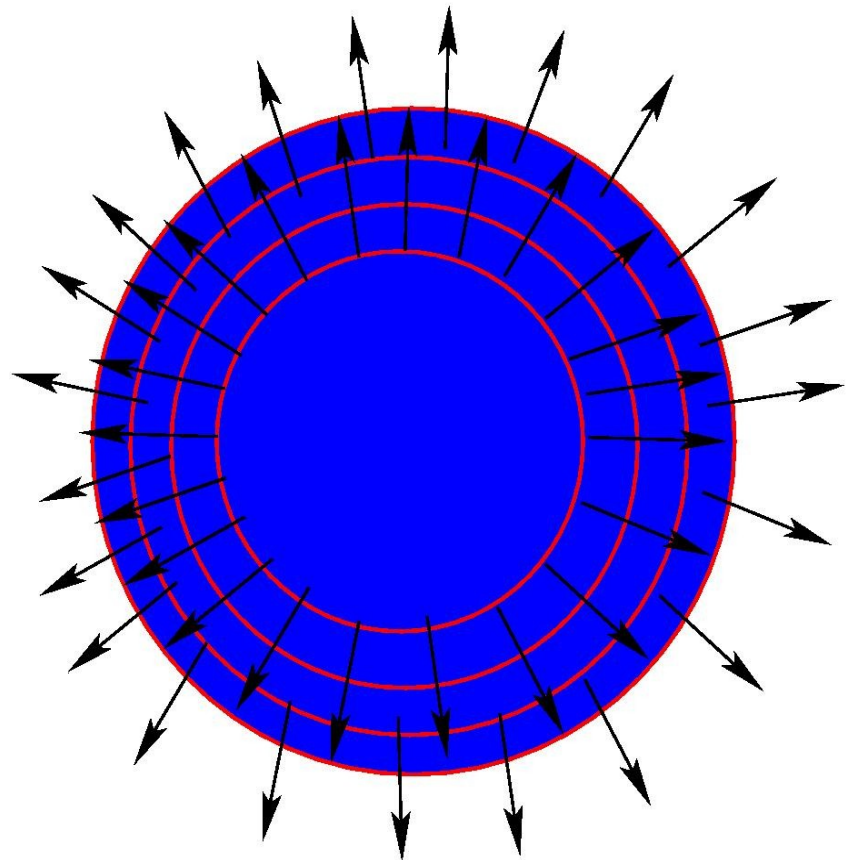




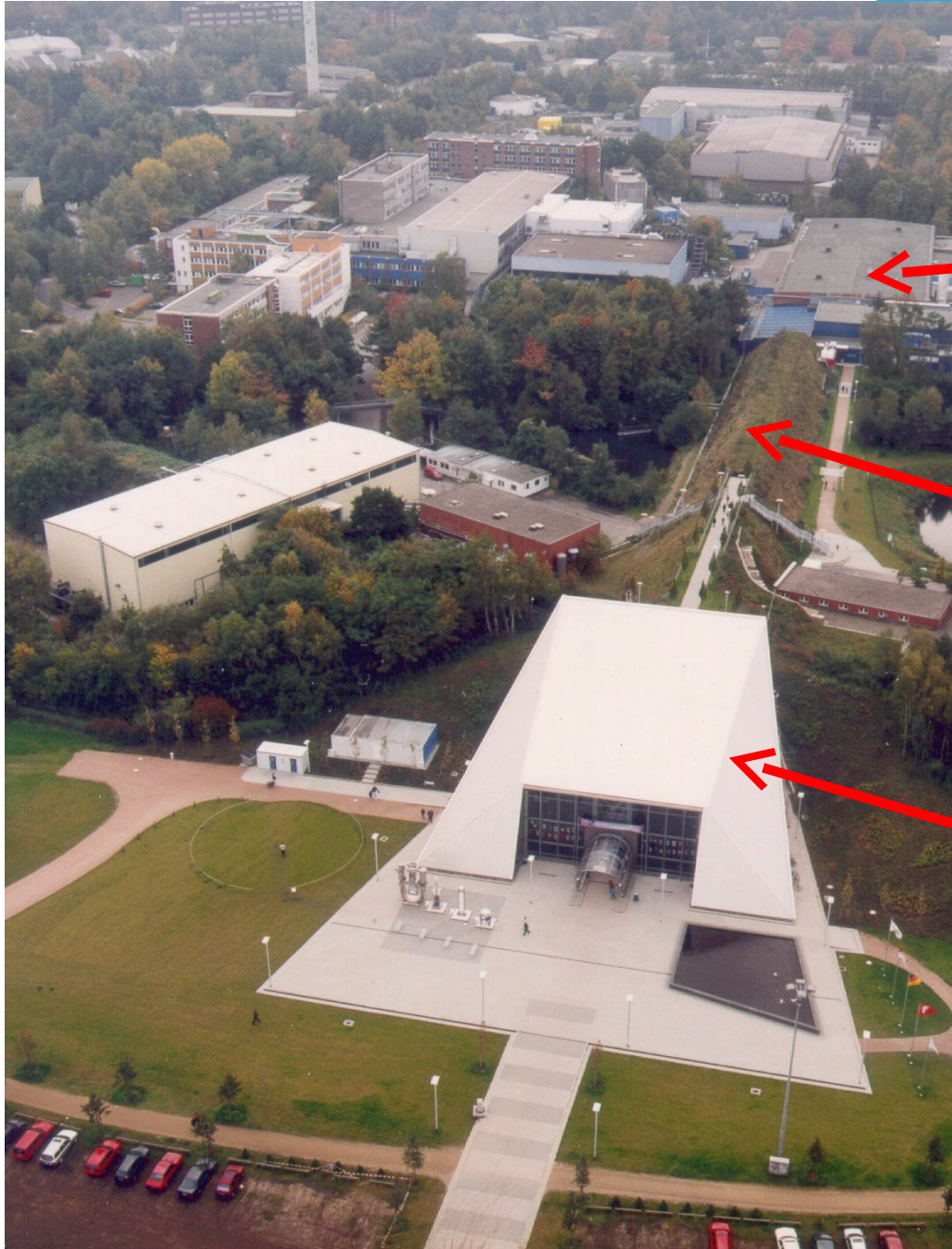
# Boltzmann solver: accuracy tests

- **Flow control in real space** → check how many particles and how much energy escaped from the simulation box → **accuracy < 5%**
- **Energy and number of particles conserved** with a **good accuracy (< 1%)** in collisionless case
- **Collisionless motion** simulated with Boltzmann solver checked with an analytical model
- **Accuracy of pseudospectral approximation** checked with an independent method

## Four spheres of flow control



# FLASH FEL at DESY



Electron gun

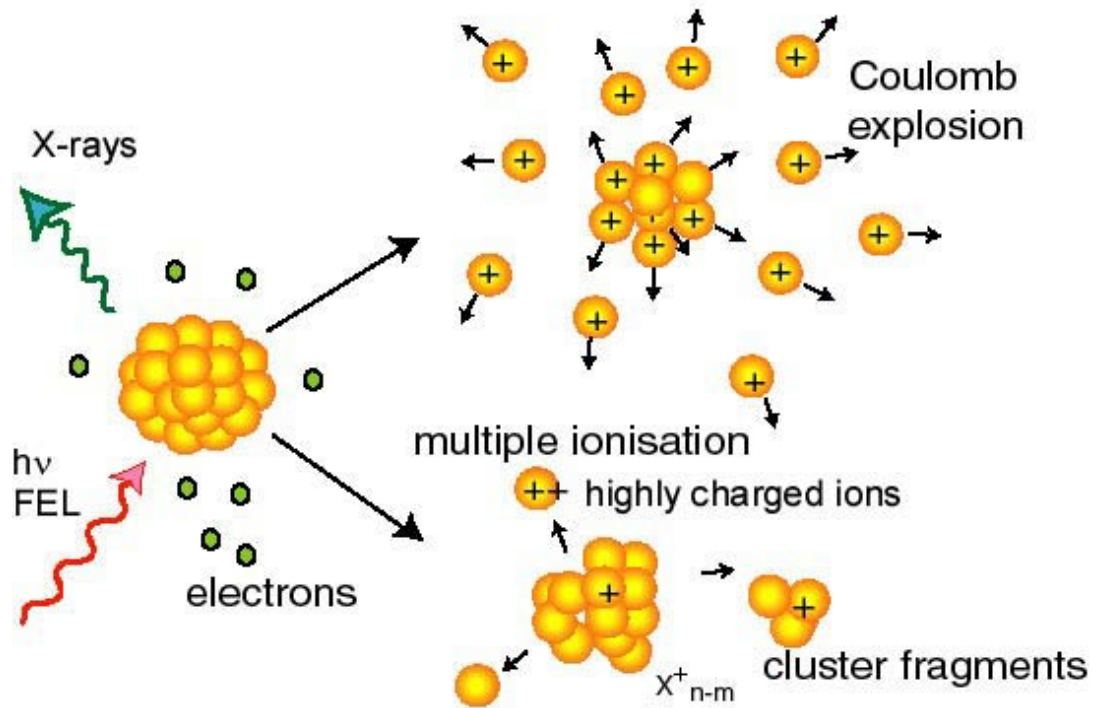
Linac and FEL  
undulator

Experimental hall  
(User Facility  
started July 2005)

- 6,5- 50 nm
- 10-100  $\mu$ J
- 1 GW<sub>peak</sub>
- 10-100 fs



# Experimental studies on clusters irradiated with intense FEL pulses



- Mechanisms of energy absorption and ionization
- Non-linear / multi-photon processes observed?
- Timescales of electron emission and of ion motion
- New processes identified?

$\lambda = 100\text{nm}$  (2002)  
valence electrons

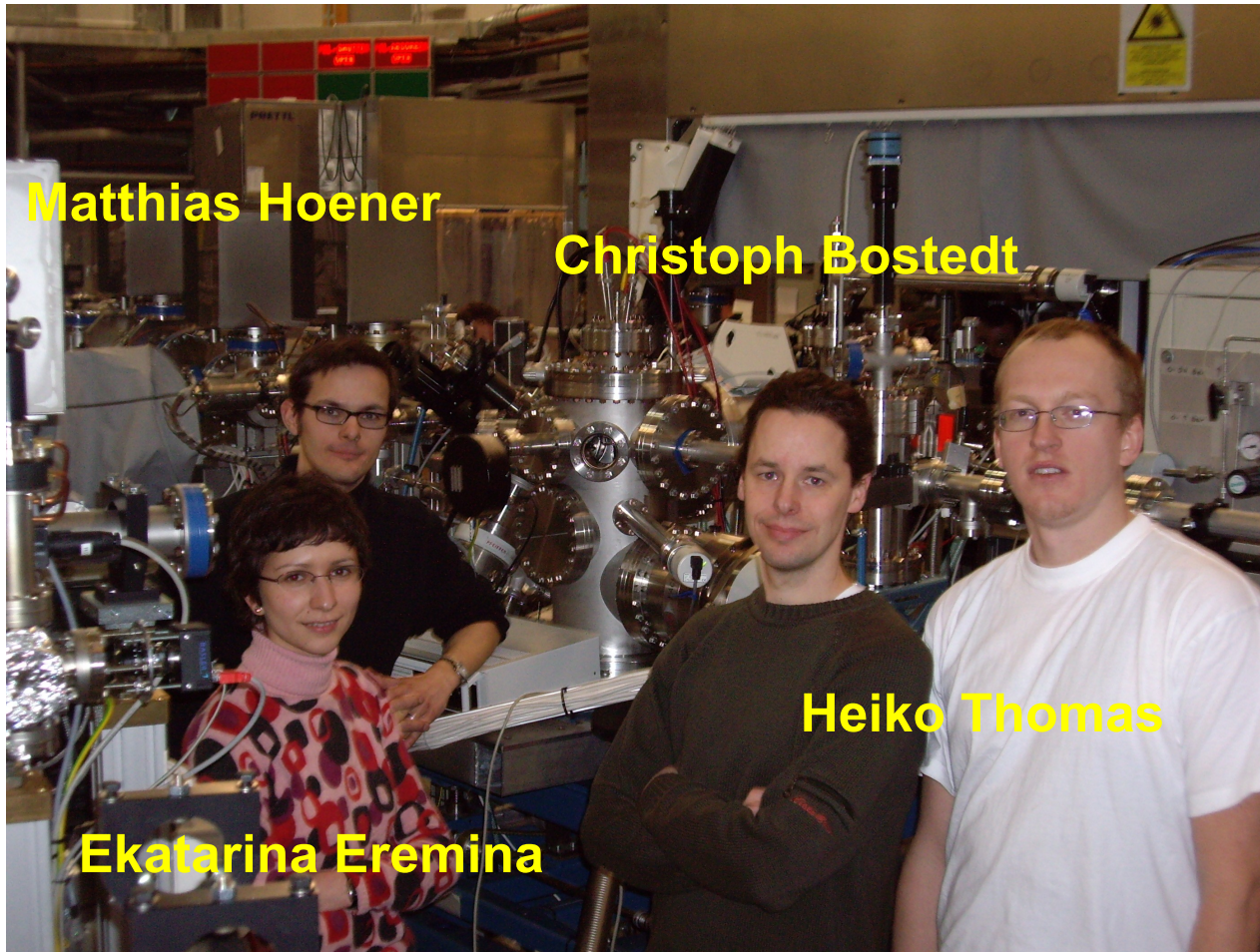
$32\text{ nm}/13\text{ nm}$  (2007-2009)  
valence/inner-shell electrons

future  $6 - 0.1\text{ nm}$   
atomic resolution

# The experimental group



Technische Universität Berlin



**Matthias Hoener**

**Christoph Bostedt**

**Heiko Thomas**

**Ekatarina Eremina**

**Thomas Möller**

**Daniela Rupp  
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Sebastian Schorb**

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**Funding:  
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**Special thanks to  
R. Treusch, S. Düsterer, J. Feldhaus  
and the FLASH control room Team**

**Collaboration  
with the groups of J. Hajdu ( Uppsala, Stanford)  
and H. Chapman (CFEL)  
R. Hartmann, C. Reich, L. Strüder, MPG Halbleiterlabor**

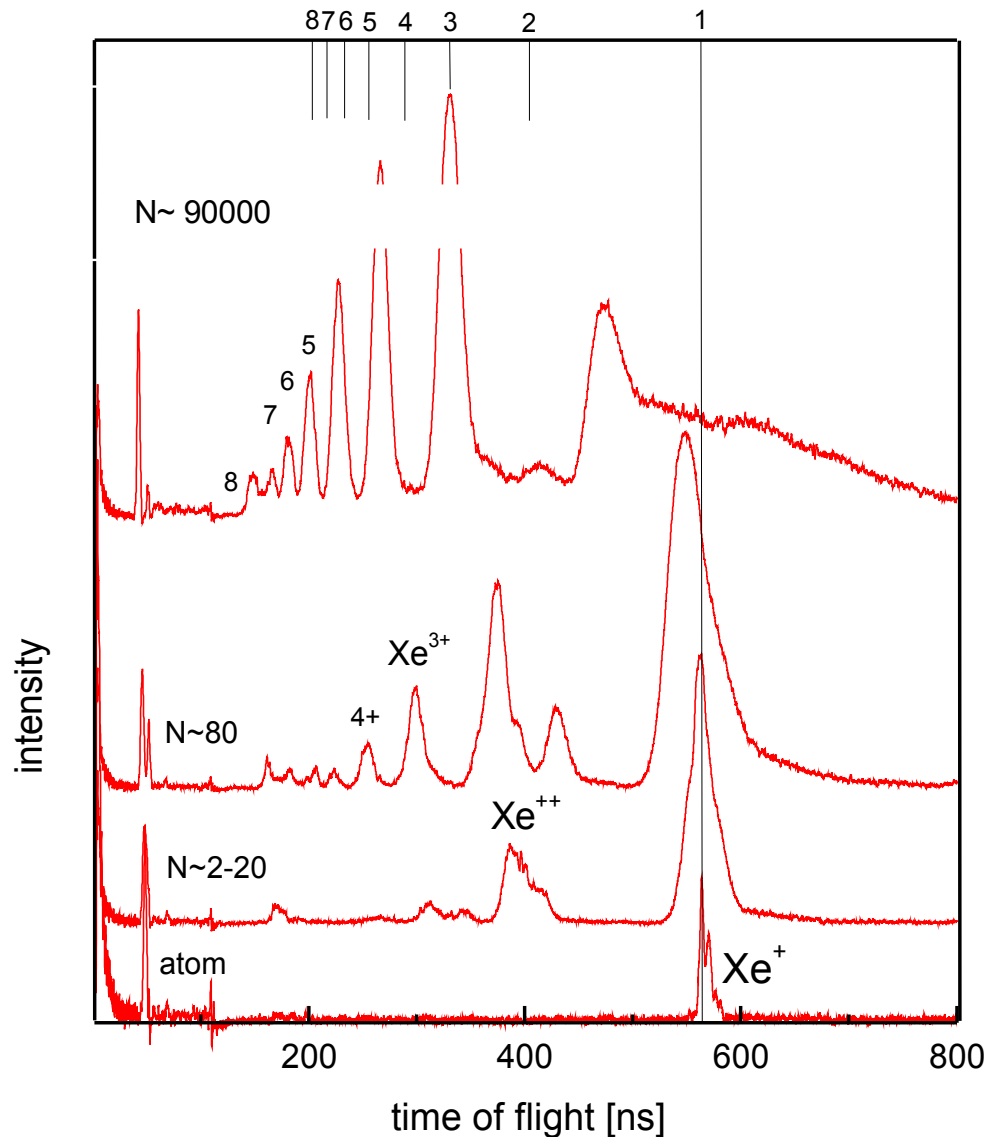
# Evolution of irradiated samples

## Two phases:

- **non-equilibrium ionization phase:** starts with the photon irradiation and lasts until thermalization of electrons and saturation of ionizations from ground state is reached
- **semi-equilibrium expansion phase:** electron-ion plasma in local thermal equilibrium, ions and electrons slowly escaping from outer shells → expansion of the sample



# Time of flight mass spectra of Xe atoms and clusters at radiation wavelength 100 nm



$10^{12}-10^{14} \text{ W/cm}^2$   $I_{p_{Xe}} = 12.1 \text{ eV}$   
[H. Wabnitz et al,  $E_{\text{phot}} = 12.8 \text{ eV}$   
Nature 420, 482 (2002)]

- Multiply charged ions from clusters, keV energy
- Only singly charged ions from atoms



**Dedicated** theoretical study needed to explain the enhanced energy absorption

pulse duration  $\sim 50 \text{ fs}$



# Theoretical models proposed

- Enhanced inverse bremsstrahlung heating of quasi-free electrons within the cluster [Santra, Greene]. Enhanced heating rate obtained with effective atomic potential. High charge states produced during collisional ionizations [R. Santra, Ch. H. Green, PRL 91, 233401 (2003)]
- High charge states within clusters are produced by single photon absorptions due to the suppression of the interatomic potential barriers within the cluster environment [Georgescu, Saalman, Siedschlag, Rost] [C. Siedschlag, J. M. Rost, PRL 93, 43402 (2004)]
- Heating of quasi-free electrons through many-body recombination [Jungreuthmayer, Ramunno, Zanghellini, Brabec]. High charge states produced during collisional ionizations.

[C. Jungreuthmayer et al., J. Phys. B 38, 3029 (2005)]

# Theoretical modelling

## What happens if all enhancement factors are included in one model?

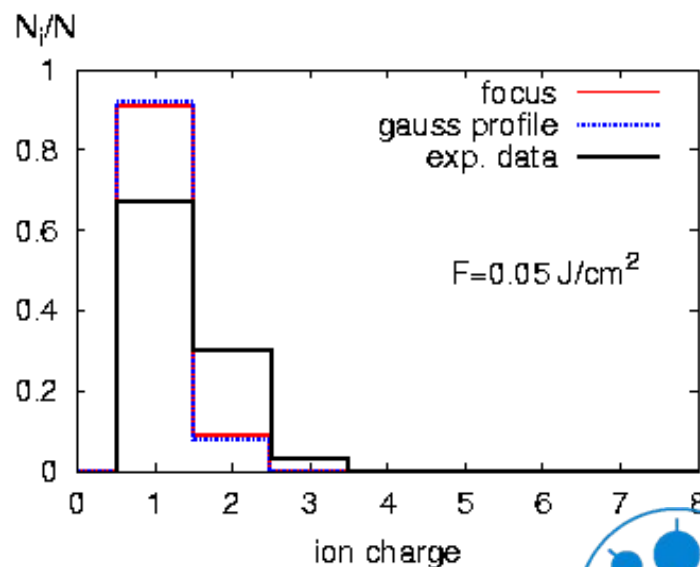
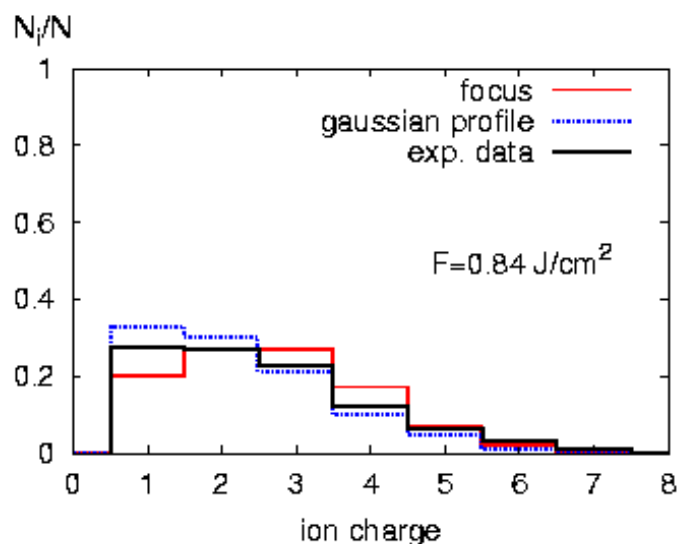
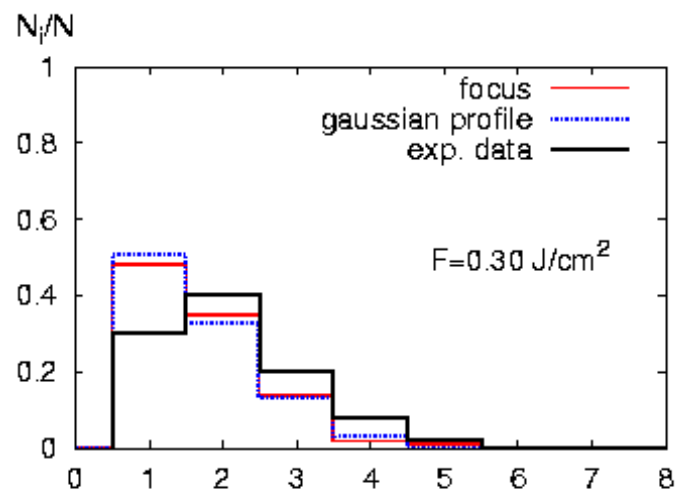
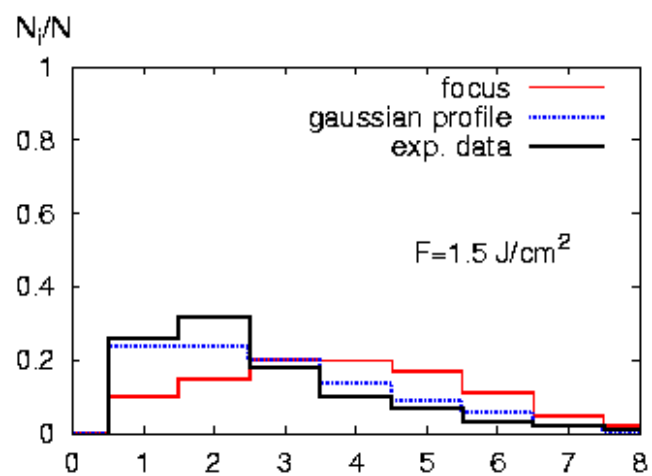
- Creation of high charges through: photoionizations → IB heating of quasi-free electrons as proposed by [Santra, Greene] → collisional ionizations / recombinations;
- Modification of atomic potentials by electron screening and ion environment tested
- Plasma regime tested → possible contribution of many-body effects (many-body recombination)

↑ ↑ ↑ ↑

- Independent cross-check with MD simulation successful [F.Wang]



# Clusters of 2500 xenon atoms irradiated with 100nm FEL pulses of various energies



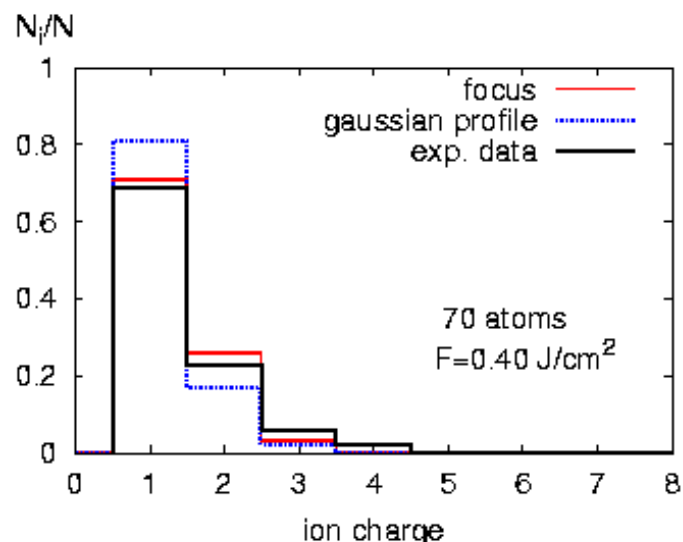
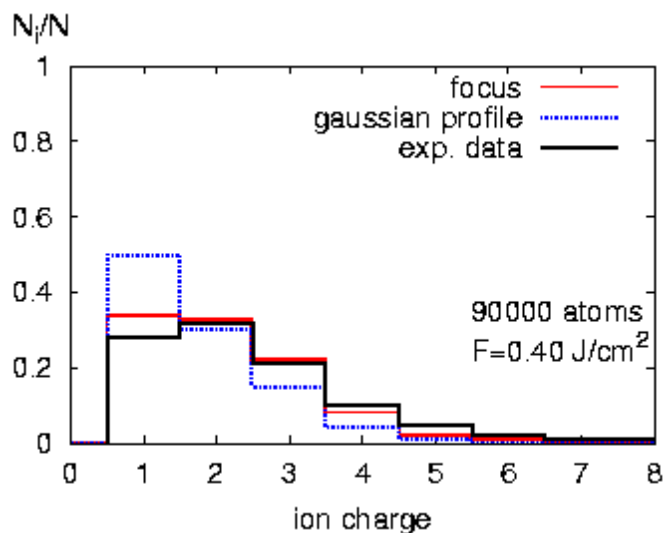
Ion fractions  
at various  
pulse fluences

TOF detects only ions !

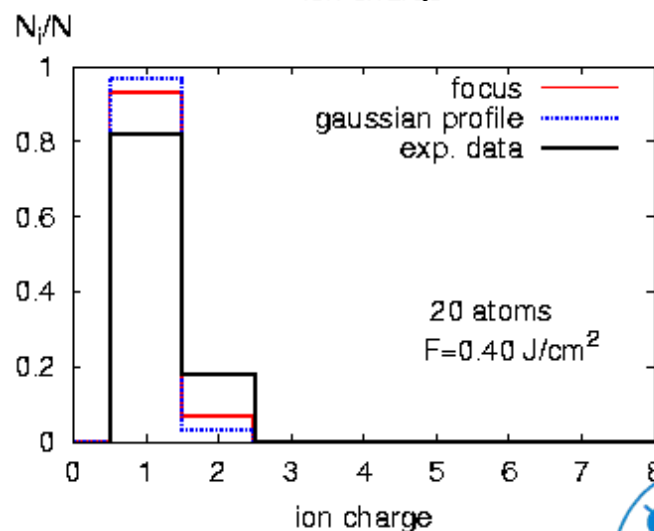
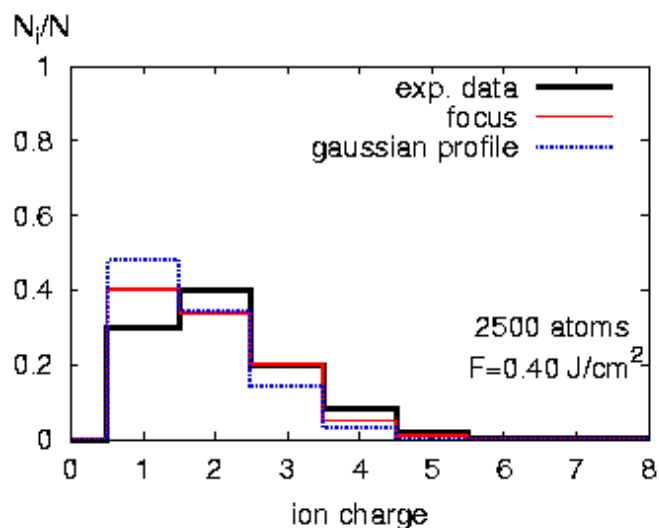
[Ziaja et al., Phys. Rev. Lett. 102 (2009) 205002]



# Clusters of various size irradiated with 100nm FEL pulses of a fixed flux, $F=0.4 \text{ J/cm}^2$



Ion fractions  
at various  
cluster size

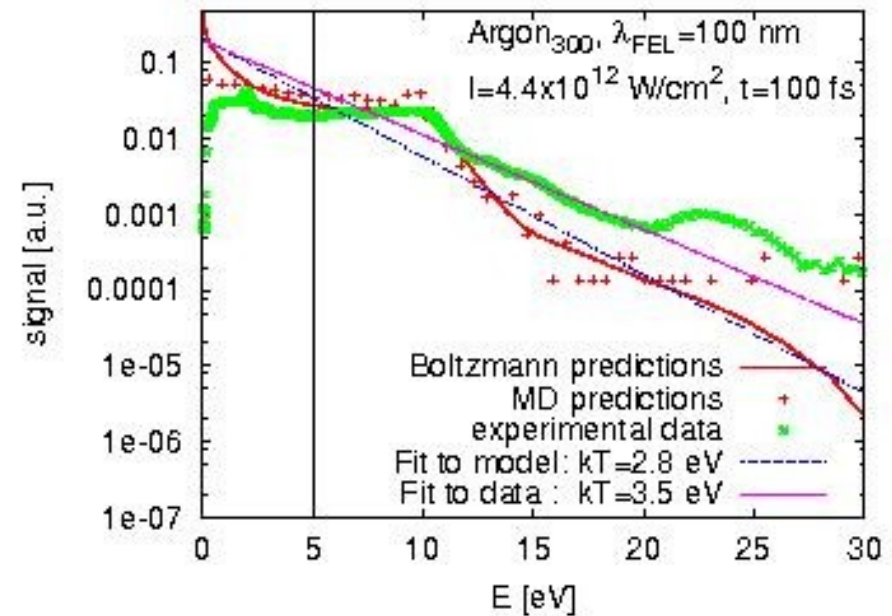
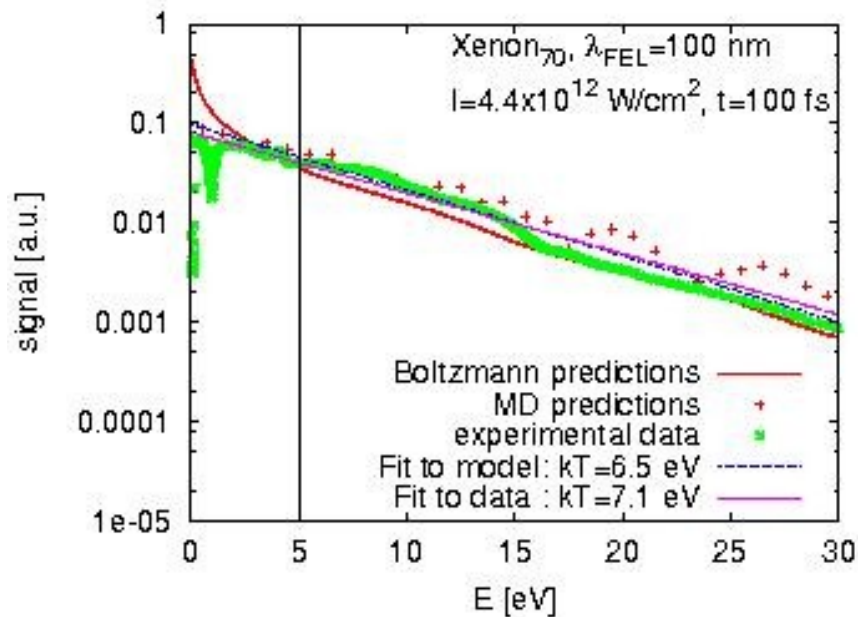


TOF detects only ions !



# Electron spectroscopy at 100 nm: indication of energy absorption mechanisms

Electron emission spectra at 100 nm for Xe (70) and Ar (300):

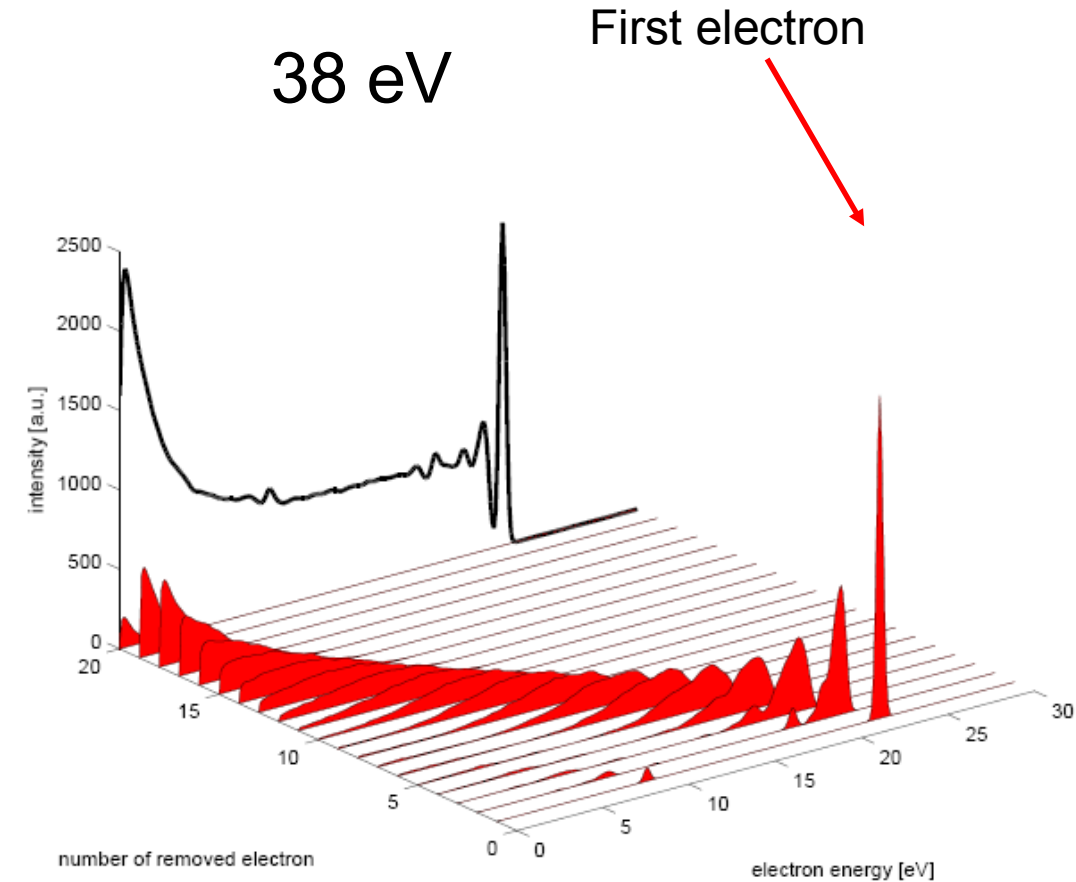
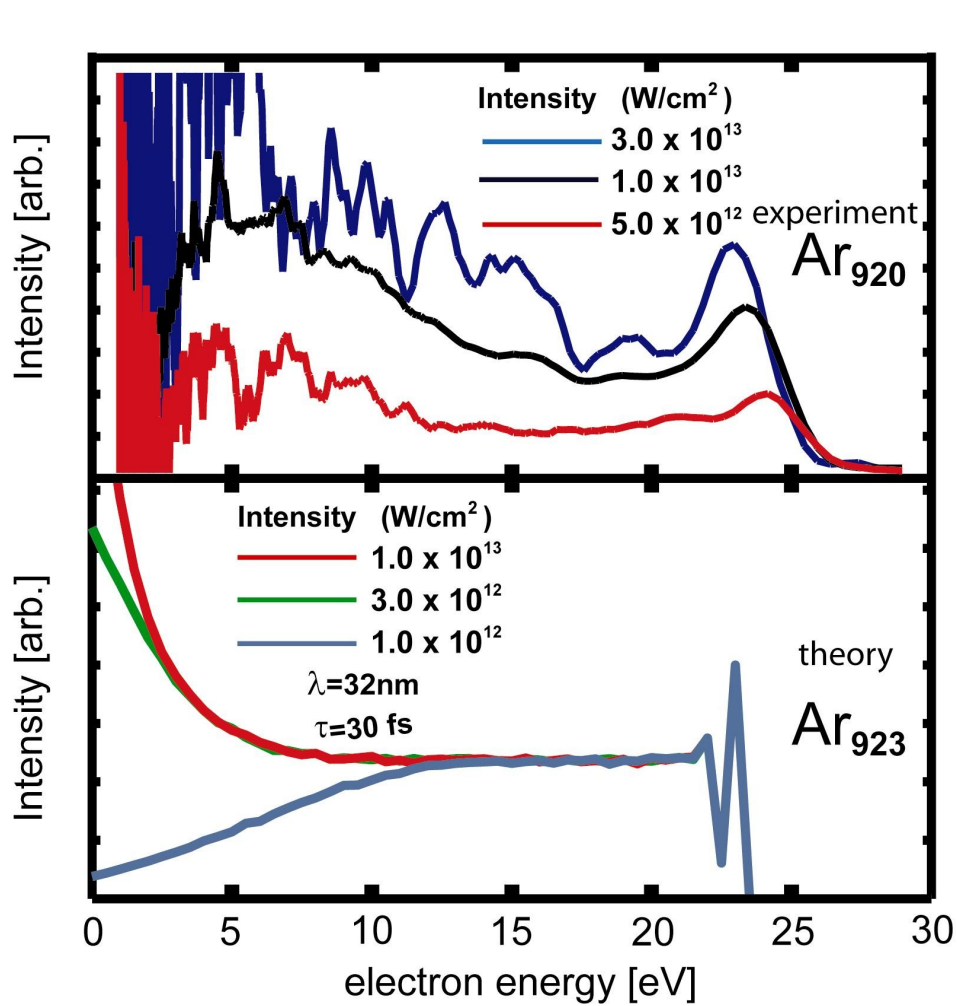


Plasma formation, intense heating of quasi-free electrons

[Ziaja et al., New J. Phys. 11 (2009) 103012]



# Electron spectroscopy at 32 nm: indication of energy absorption mechanisms



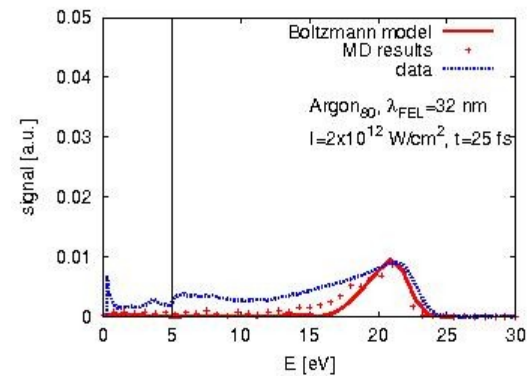
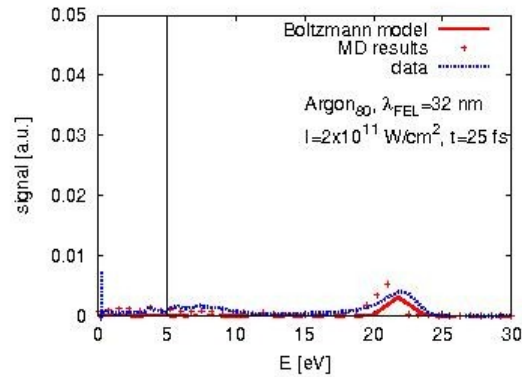
**no thermionic electron emission  
no plasma absorption [T. Fennel et al.]**

[C. Bostedt et al. Phys. Rev. Letters 100, 133401 (2008)]

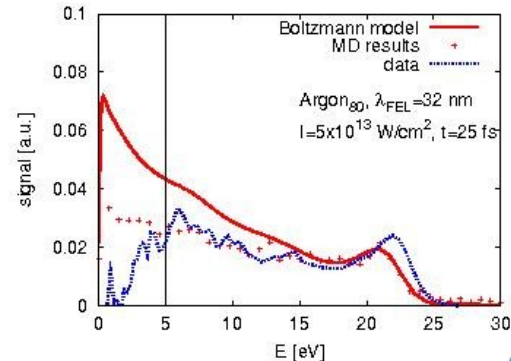
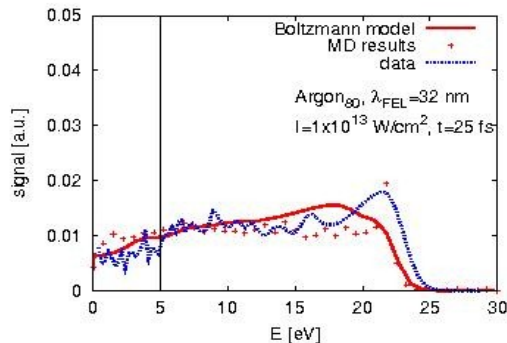


# Electron spectroscopy at 32 nm indication of energy absorption mechanisms

Electron emission spectra at 32 nm for Ar(80) and Ar (150):  
**sequential ionization**  
[Ziaja et al., New J. Phys. 11 (2009) 103012]



FEL pulse length=25 fs

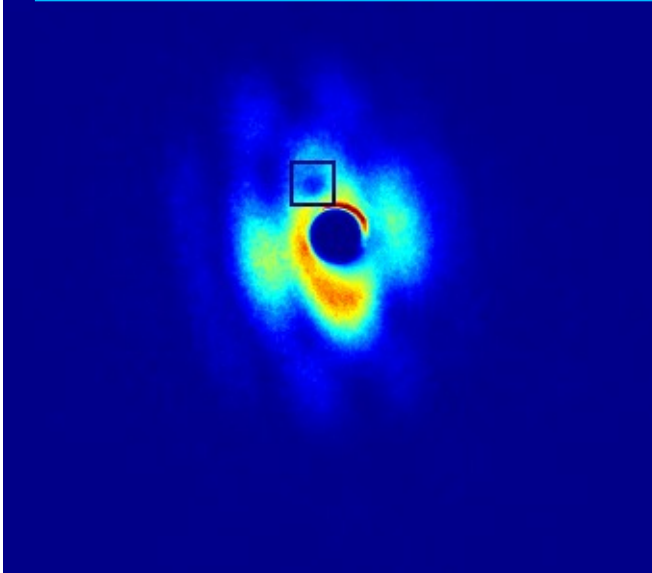


Good agreement of the results with MD simulations by the Rostock group [T. Fennel et al.] and in-house MD simulations [F. Wang]



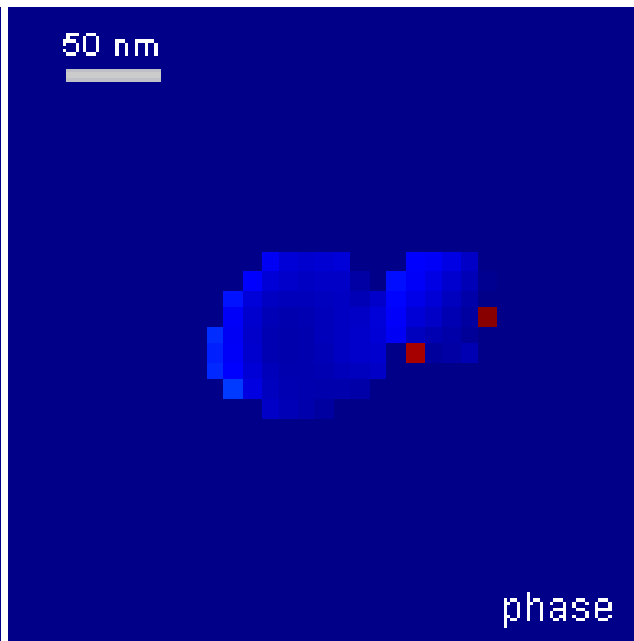
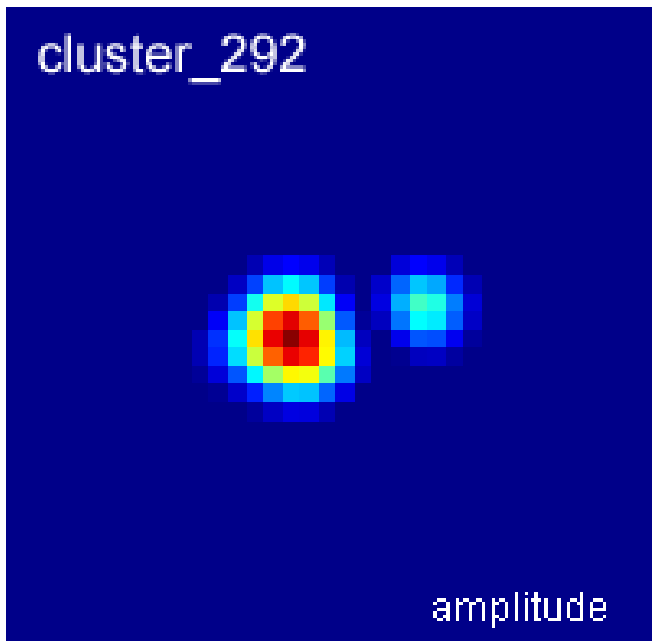
# Single shot scattering and imaging of large noble gas clusters at wavelength $\sim 13$ nm

[C. Bostedt, H. Chapman, F. Wang and T. Möller]



scattering pattern

**Wavelength 13.7 nm**

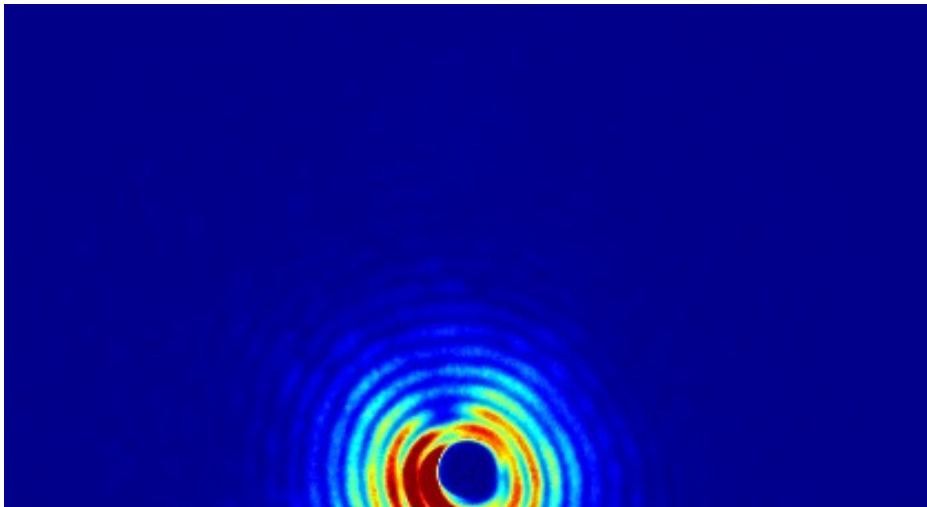


reconstructed image

two clusters in  
direct contact

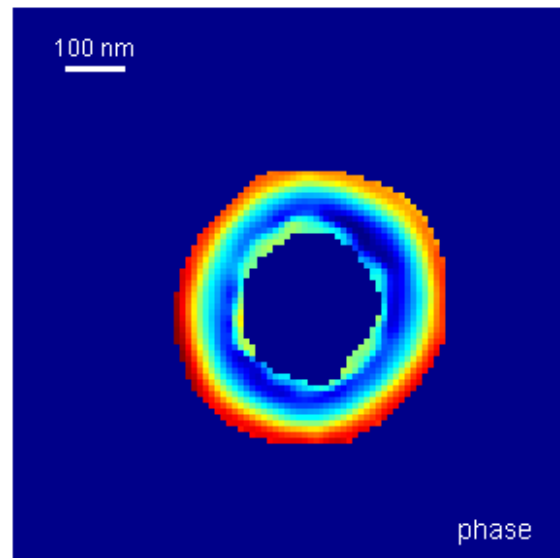
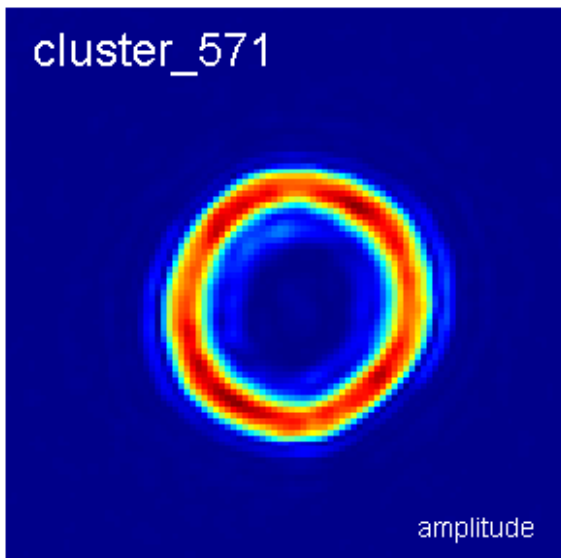


# Single shot scattering and imaging of large clusters: reconstruction



scattering pattern

**Wavelength 13.7 nm**




reconstructed image

large cluster

penetration depth of light





# Soft X-Ray Thomson Scattering in **Warm Dense Hydrogen**

R.R. Fäustlin, S. Toleikis et al.

[Phys. Rev. Lett. 104 (2010), 125002]





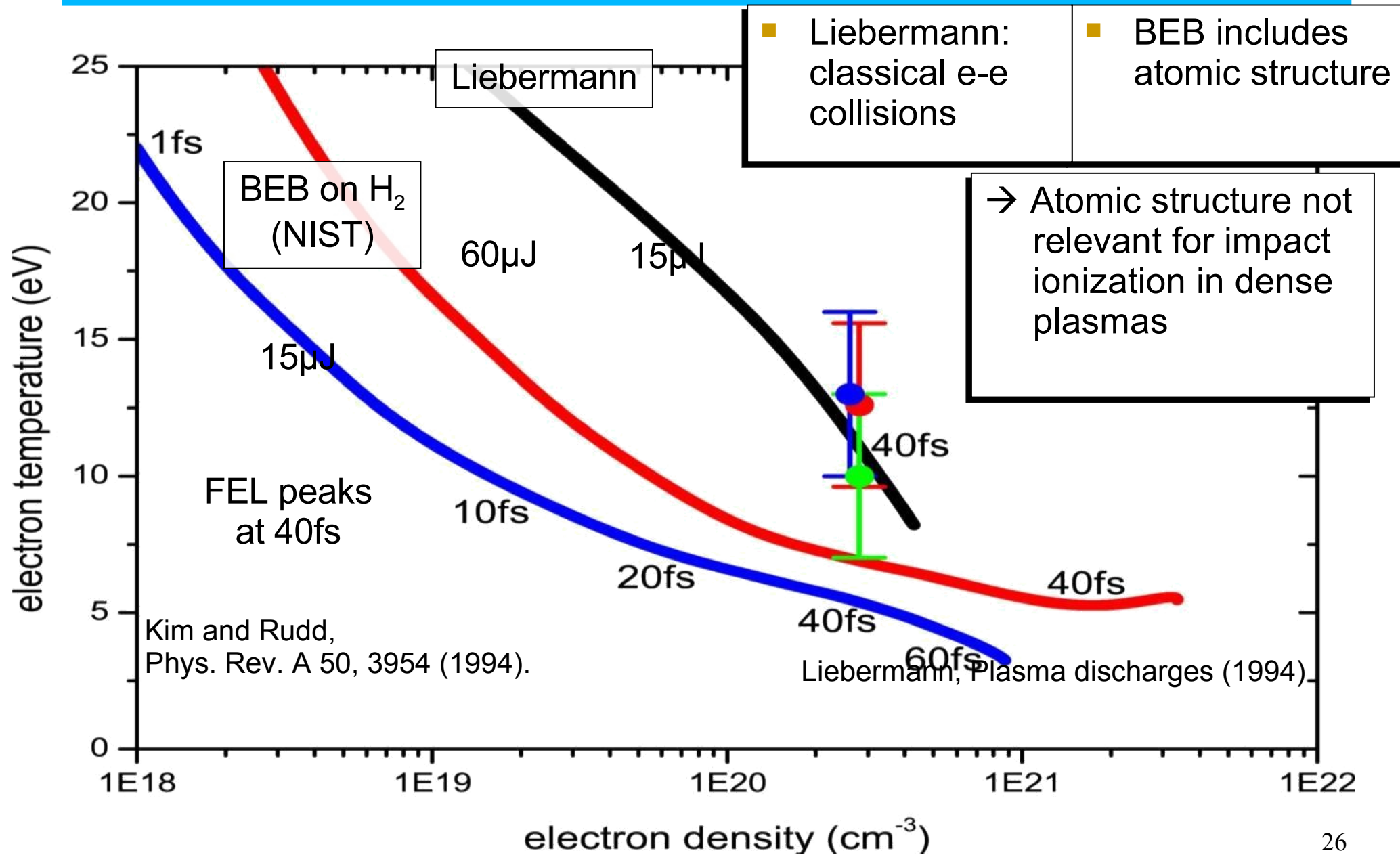
# Collaboration

- **DESY, Hamburg**  
S. Düsterer, R.R. Fäustlin, T. Laarmann, H. Redlin,  
N. Stojanovic, F. Tavella, S. Toleikis, T. Tschentscher
- **University of California, Berkeley**  
H.J. Lee
- **University of Jena**  
E. Förster, I. Uschmann, U. Zastra
- **LLNL, Livermore**  
T. Döppner, S.H. Glenzer
- **University of Oxford / RAL, Chilton, Didcot**  
G. Gregori, B. Li, J. Mithen, J. Wark
- **University of Rostock**  
T. Bornath, C. Fortmann, S. Göde, R. Irsig,  
K.-H. Meiwes-Broer, A. Przystawik, R. Redmer,  
H. Reinholz, G. Röpke, R. Thiele, J. Tiggesbäumker



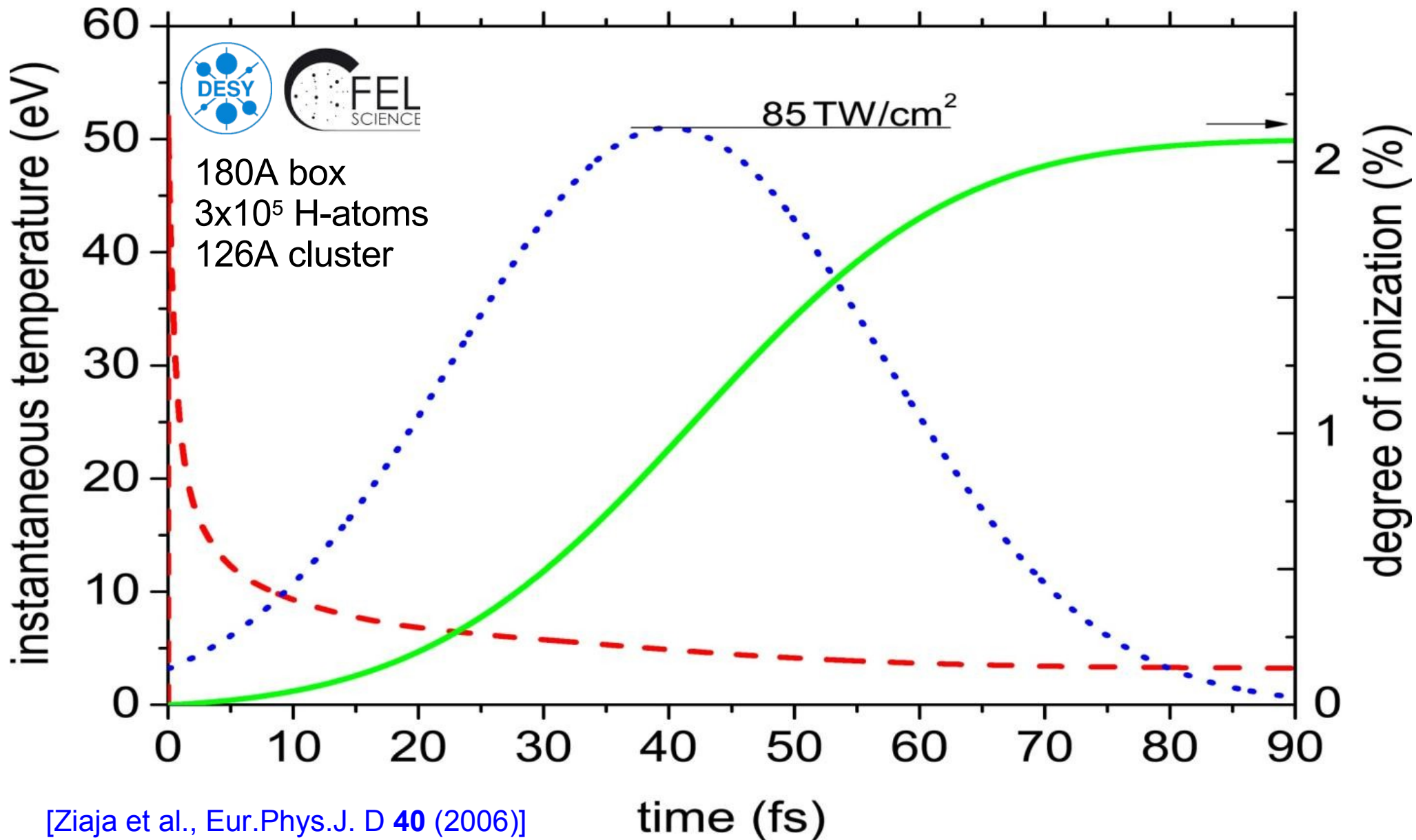
**Helmholz Association:** VH-VI-104; **Science and Technology Facilities Council UK;**  
**German Research Foundation:** GRK 1355, SFB 652, LA 1431/2-1; **DOE:** DE-AC52-07NA27344,  
LDRDs 08-ERI-002, 08-LW-004; **German Ministry for Education and Research:** FSP 301-FLASH;  
**European Community:** RII3-CT-2004-506008 (IA-SFS)

# We Validate Impact Ionization Models



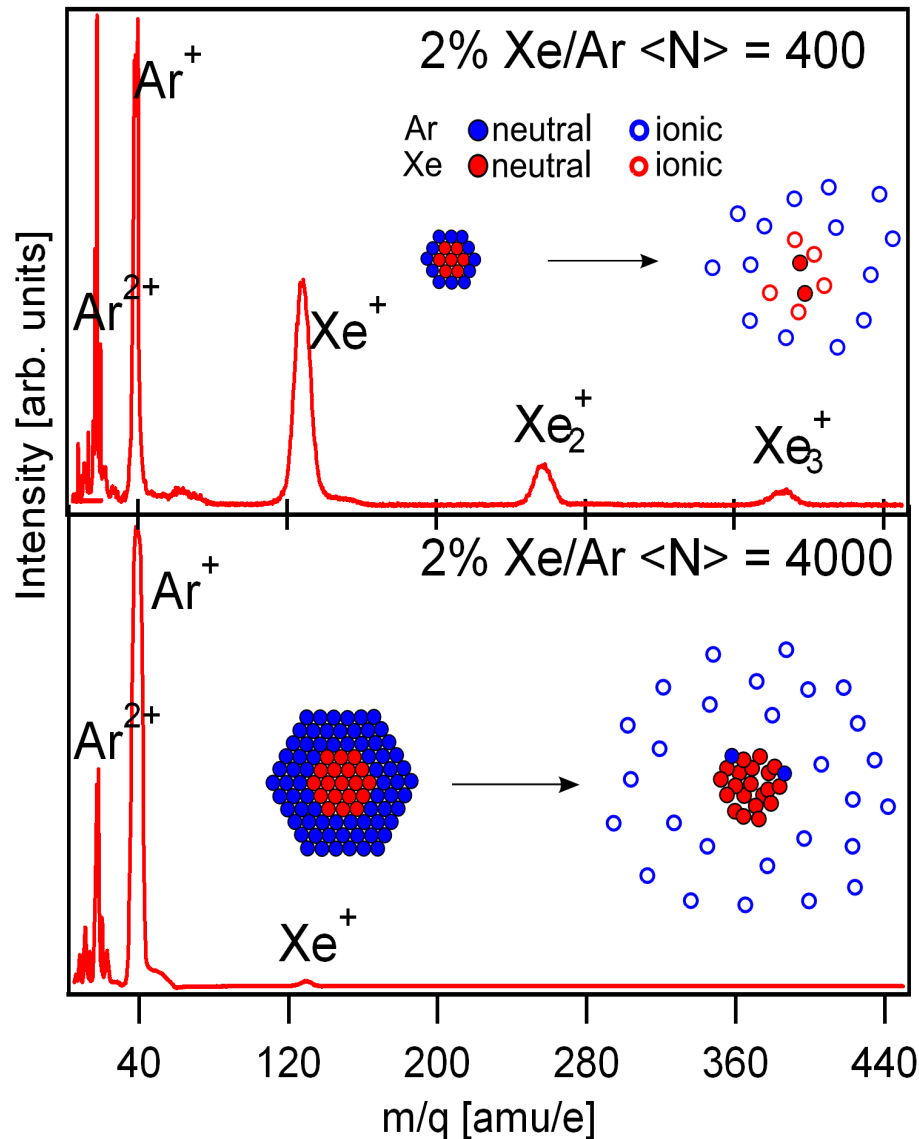
Simulations performed for clusters consisting of 0.5 million atoms

# STS Probes fs Electron Thermalization



[Ziaja et al., Eur.Phys.J. D **40** (2006)]

# Perspectives for imaging: delayed ionization and expansion through tampering.

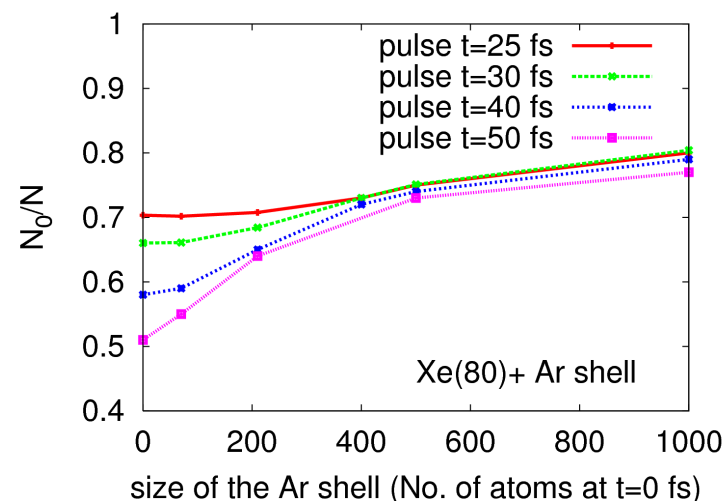
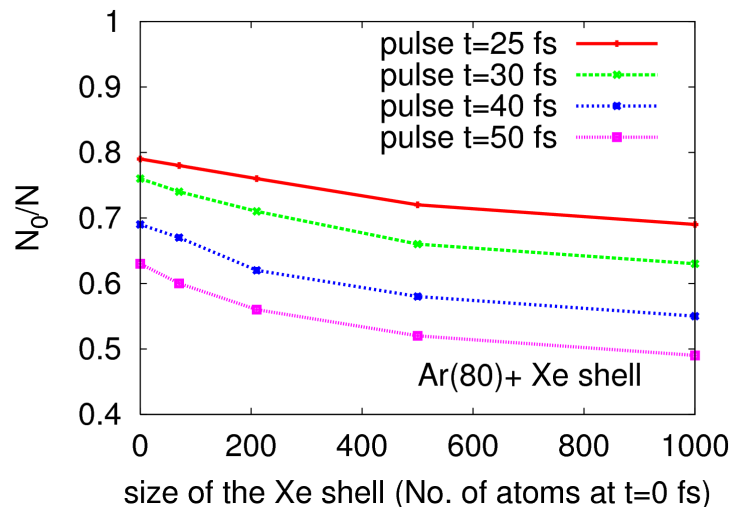


**Experiment:**  
**13 nm/ 92 eV,  $10^{14}$  W/cm<sup>2</sup>**

- **singly charged Ar ions from the surface**
- **strong size effect**
- **Xe plasma in the interior recombines, **neutral atoms**?**
- **delayed cluster expansion**



# Perspectives for imaging: increased/decreased ionization in mixed Xe/Ar clusters at 32 nm



32 nm/ 40 eV,  $10^{13}$  W/cm<sup>2</sup>

- Increased core ionization for Ar/Xe
- Decreased core ionization for Xe/Ar

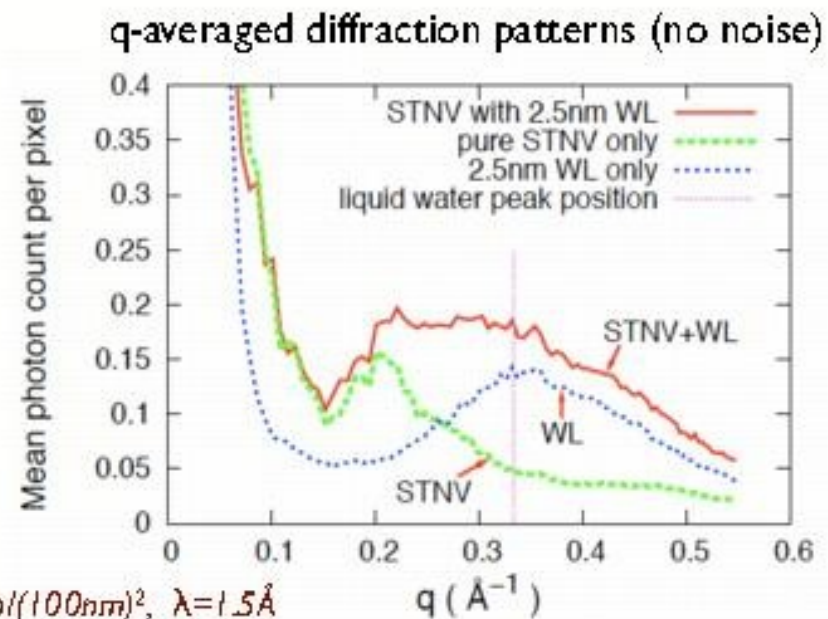
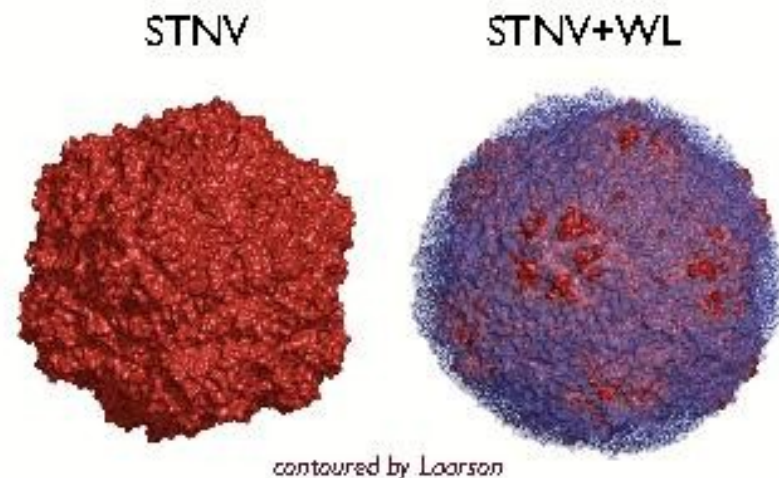
**PRELIMINARY!**



# Perspectives for imaging: loss of structural and orientational information due to the water layer around the imaged object

Satellite Tobacco Necrosis virus (STNV, PDB ID: 2BUK)  
object size ~17 nm, ~0.18M atoms, icosahedral symmetry.

- \* **dynamic** simulations with GROMACS (Larrson & van der Spoel)
  - no radiation damage considered
  - generate random water layers
- \* **diffractive** simulations with obtained atom coordinates ( use Moltrans code)



# Perspectives for imaging: loss of structural and orientational information due to the water layer around the imaged object

## Correlation analysis

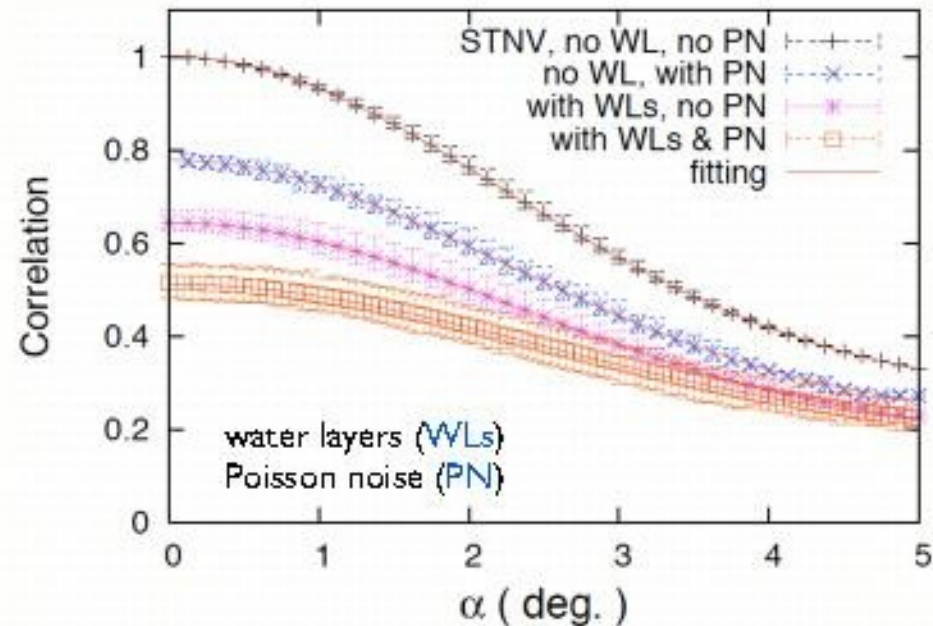
examples of correlation curves (given  $q$ , between different views)

$$G(\alpha, q) \equiv \frac{\tilde{I}(\alpha, q) \cdot \tilde{I}(0, q)}{\sqrt{\langle \tilde{I}(\alpha, q)^2 \rangle_q} \cdot \sqrt{\langle \tilde{I}(0, q)^2 \rangle_q}}$$

correlations in given  $q$  range:  $[0.069, 0.104] \text{ \AA}^{-1}$

- \* Suppression of correlation
- \* value fluctuation
- \* averaged curve perfectly fitted by Gaussian-like function:

$$G(\alpha, q) = a(q) e^{-\alpha^2/2b(q)^2} + c(q)$$

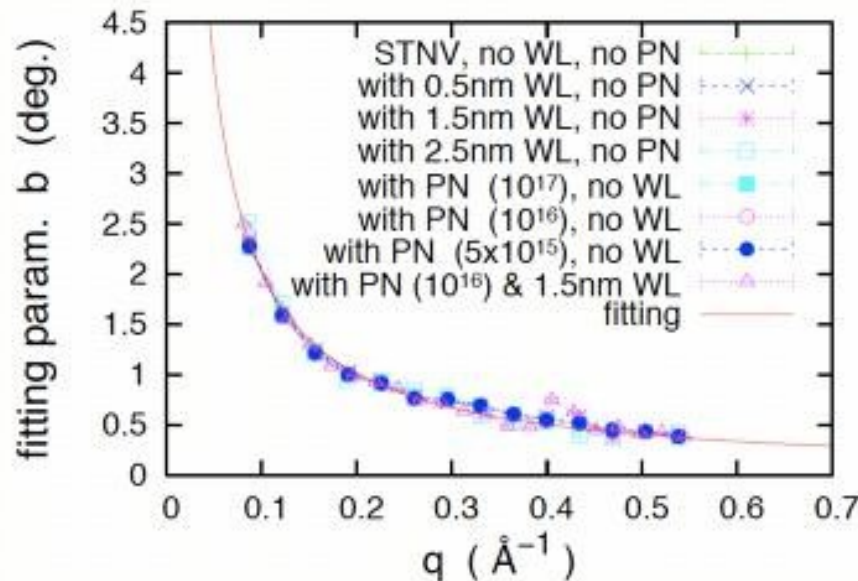


# Perspectives for imaging: loss of structural and orientational information due to the water layer around the imaged object

Fitting results of correlation curves (all  $q$  ranges)

$$G(\alpha, q) = a(q) e^{-\alpha^2/2b(q)^2} + c(q)$$

fitting parameter  $b(q)$  & its fitting



\* Orientation information remained

\* Fitting  $q$ -dependence by

$$b(q) = (4 R_0 q)^{-1}$$

(half-Shannon angle)

fitting value  $R_0 \sim 71 \text{ \AA}$



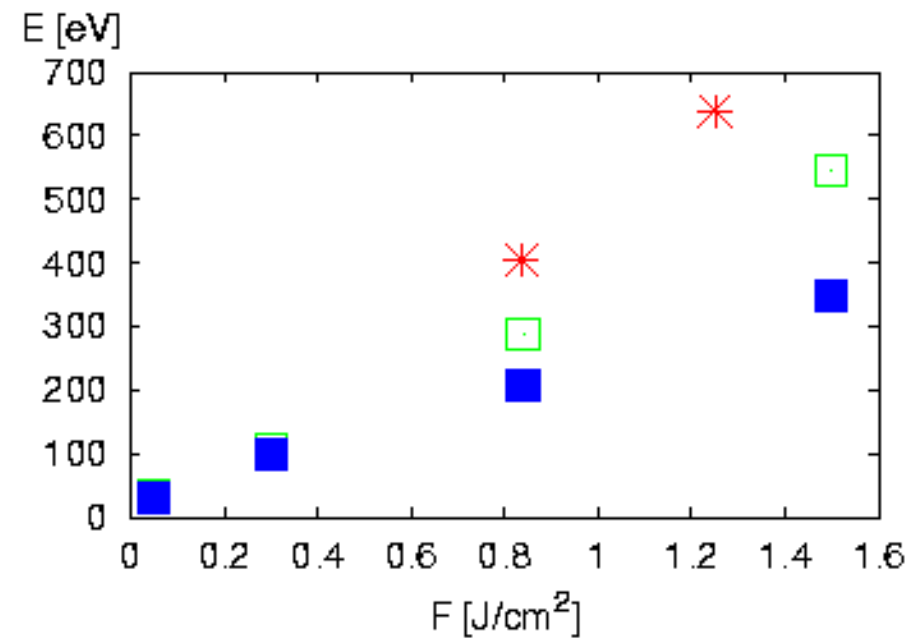
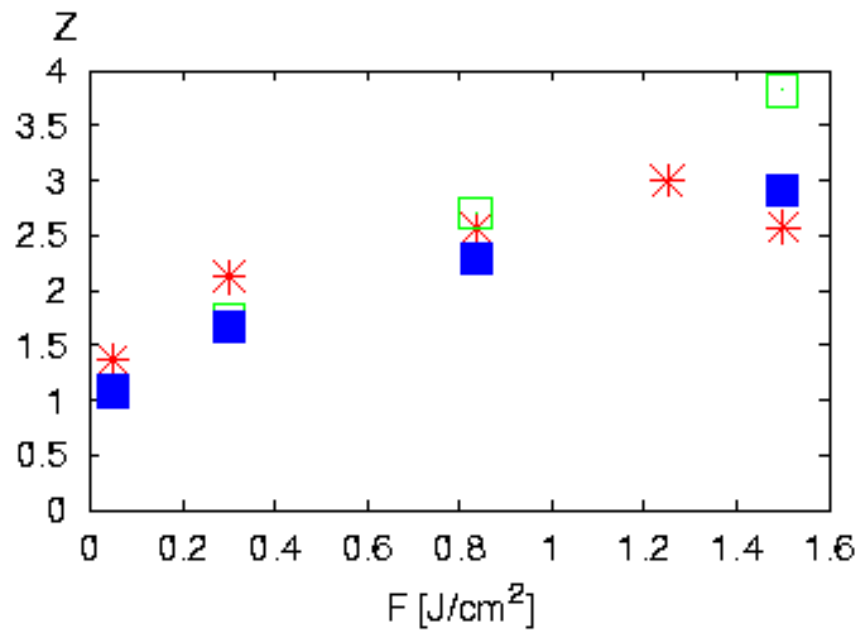


# Summary and outlook

- **We constructed a useful tool for studying the evolution of FEL irradiated samples (computationally efficient treatment of large samples)**
- **Our model is so far the only one that gives an accurate description of all of the experimental data collected from atomic clusters at 100 nm and 32 nm wavelength**
- **Good agreement with data from warm dense matter hydrogen experiment**
- **Several problems can be studied in the next future:**
  - **evolution of clusters irradiated with XUV and X-rays**
  - **clusters of various structures**
  - **mechanisms of slowing down the cluster explosion**
  - **samples exposed to ultra short and/or ultra intense pulses**



# Clusters of 2500 xenon atoms irradiated with 100nm FEL pulses of various energies



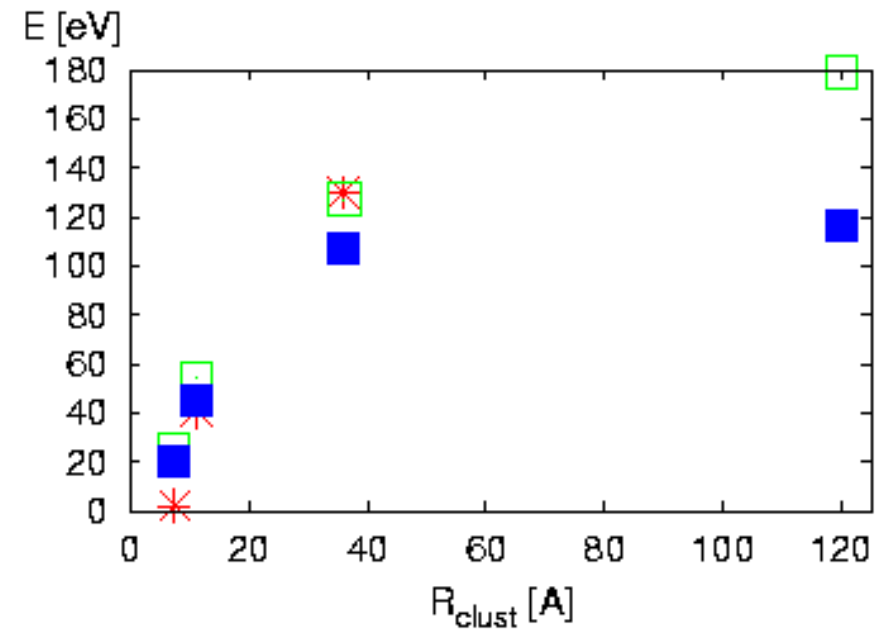
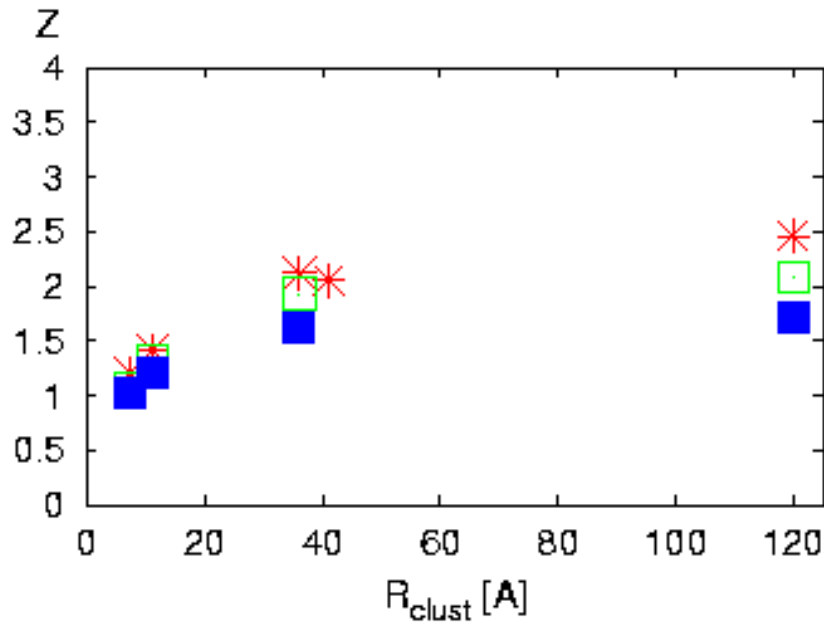
Results from model □ (in focus)

□ (averaged over spatial pulse profile)

Experimental data \*



# Clusters of **various size** irradiated with **100 nm** FEL pulses of a fixed flux, $F=0.4 \text{ J/cm}^2$



Results from model □ (in focus)      Experimental data \*

□ (averaged over spatial pulse profile)

