

Gas, stars and gravitational waves

on the main driver of supermassive black hole binaries path to coalescence

Elisa Bortolas

Main Collaborators:

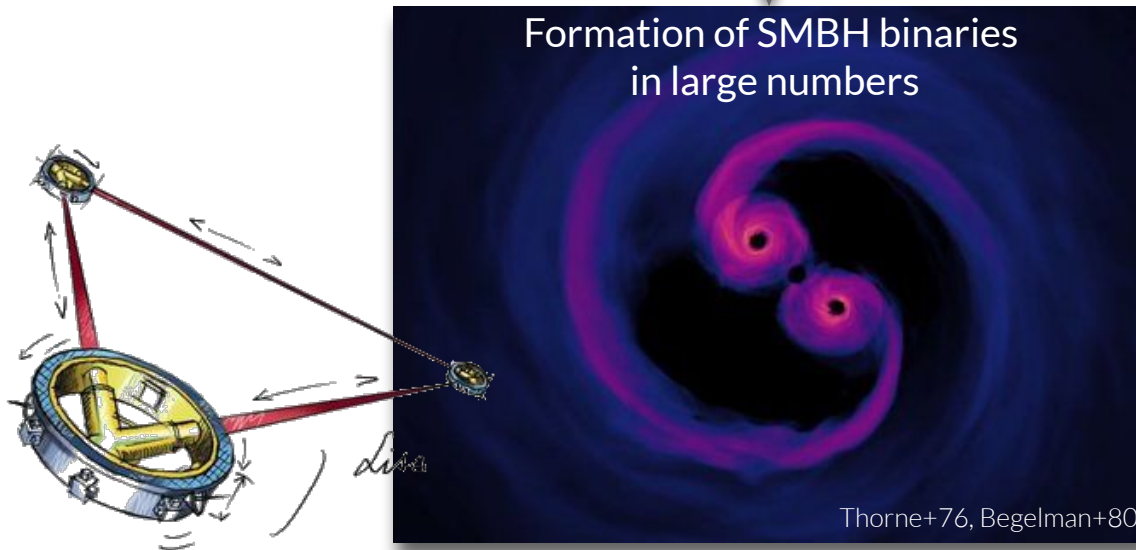
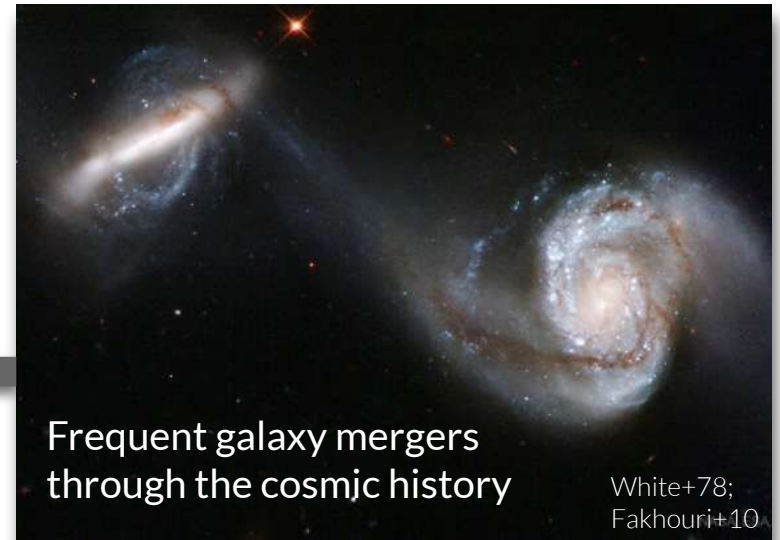
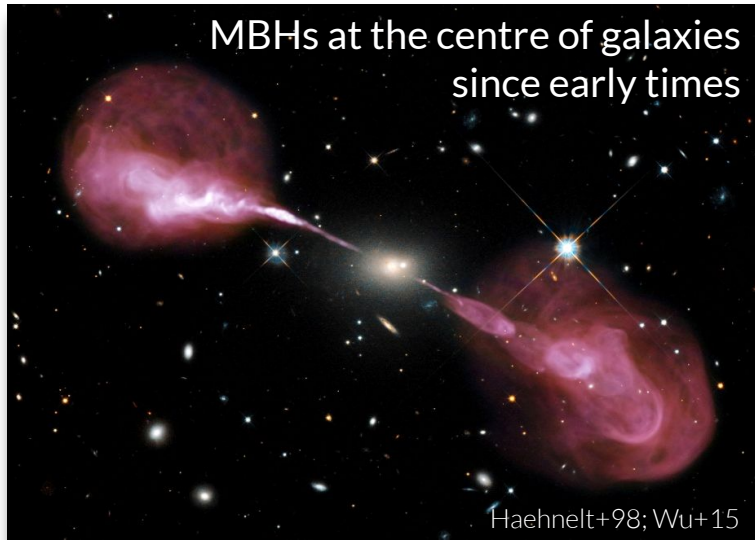
A. Sesana, A. Franchini, M. Bonetti,
A. Gualandris, A. Lupi, M. Dotti,
L.Mayer, P.Capelo, ...

**BINARY22 - Kavli Institute
for Theoretical Physics**

Santa Barbara, April 7th, 2022

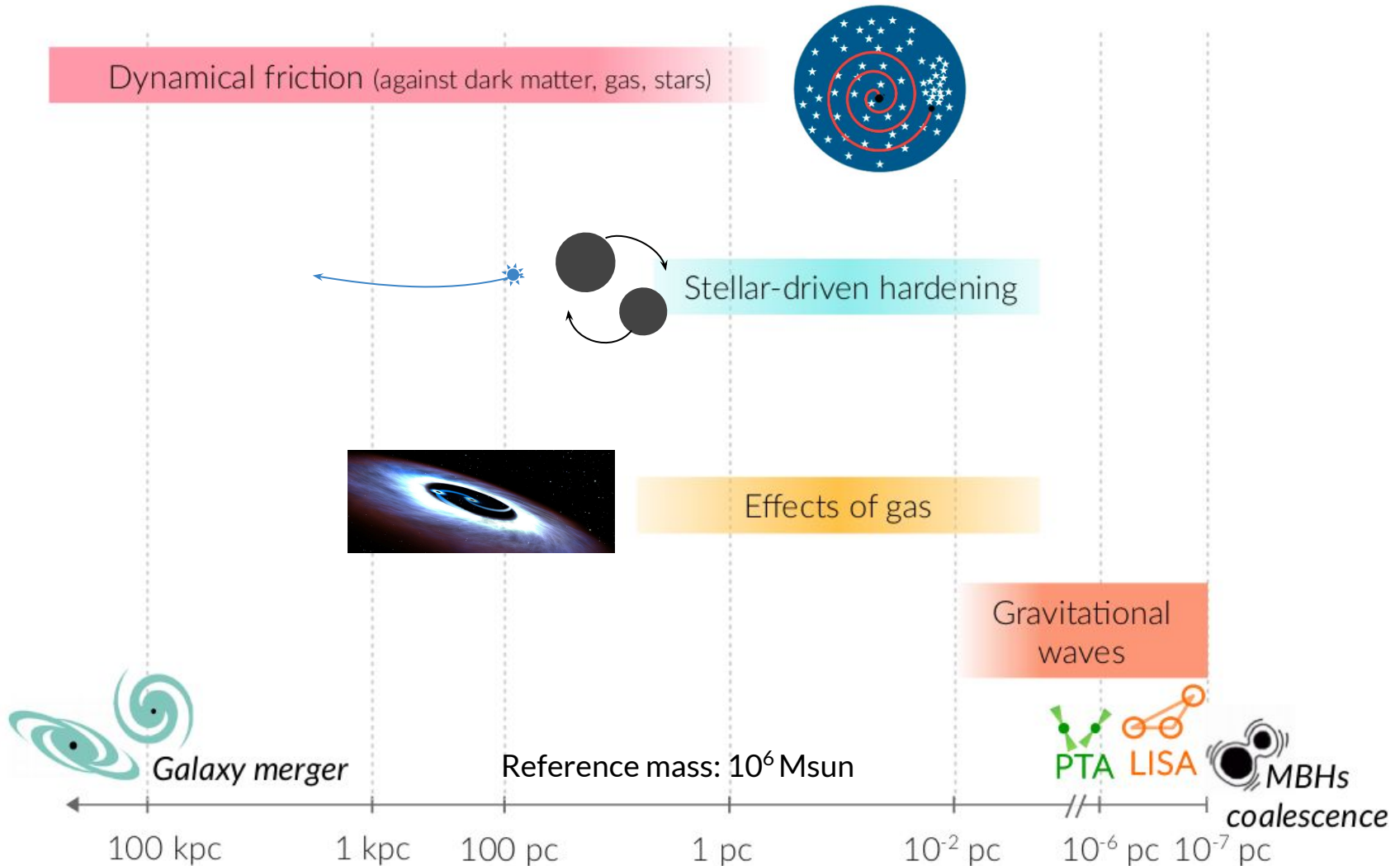


Massive black hole binaries



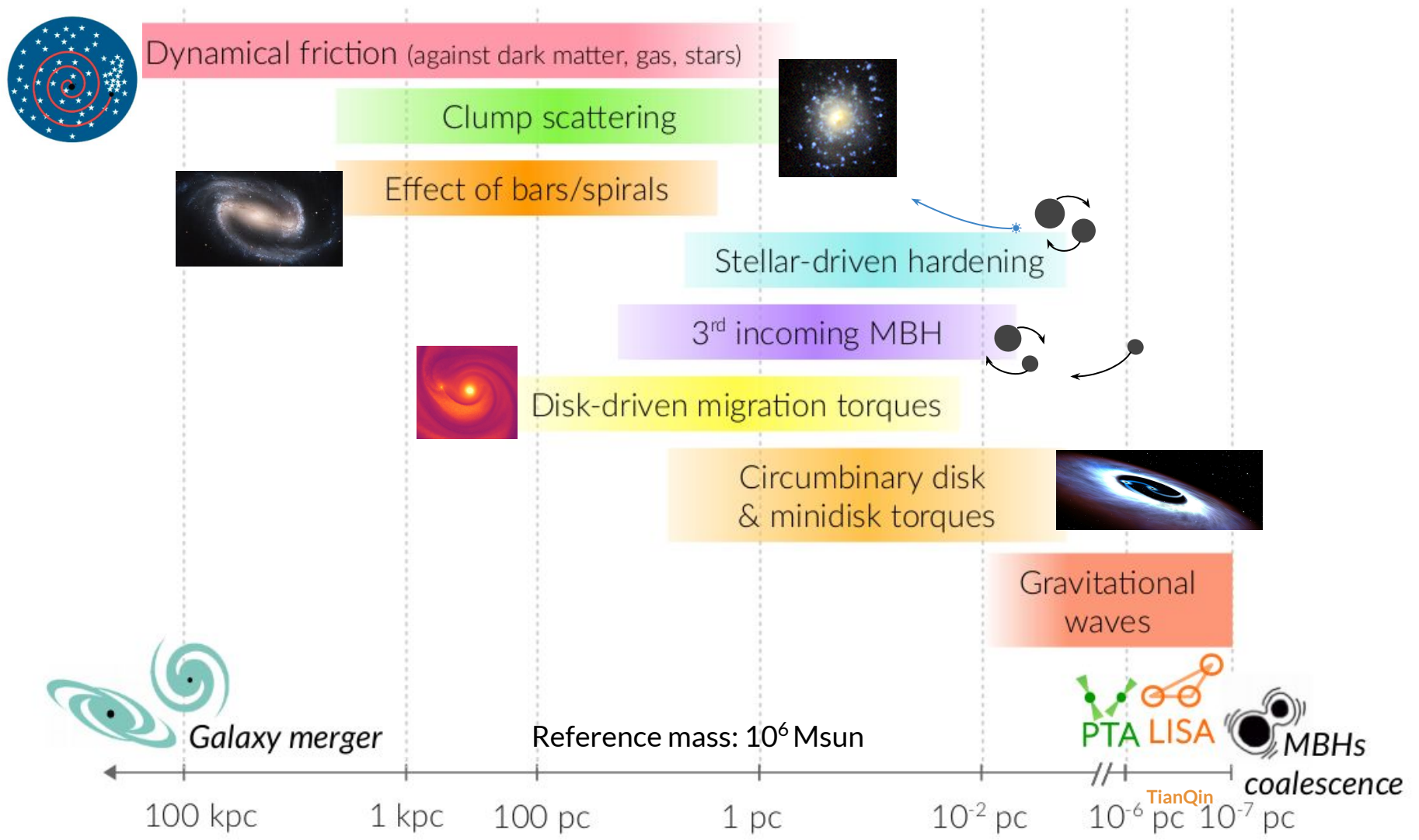
The path to coalescence ~40 years ago...

Begelman, Blandford and Rees 1980

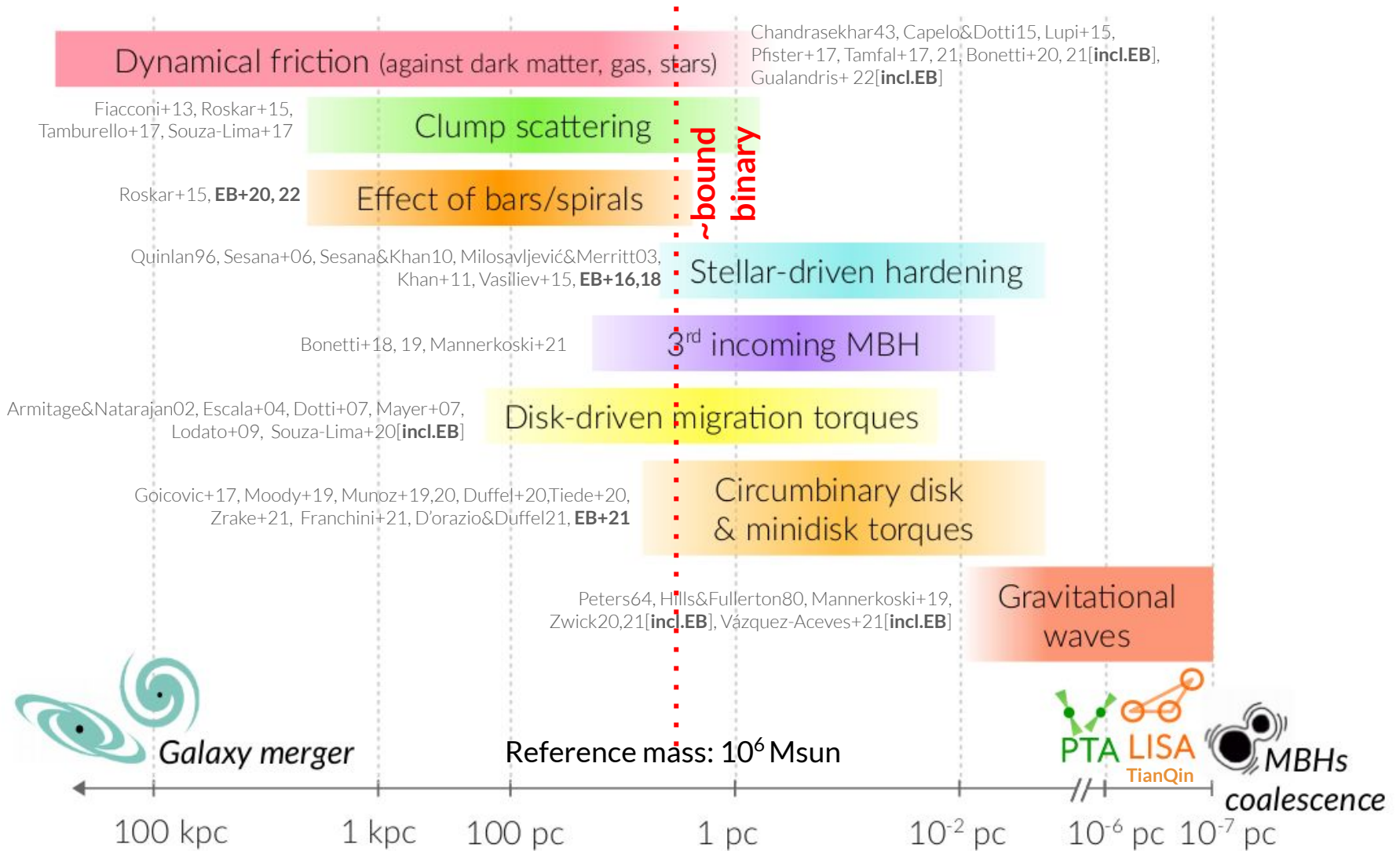


Reference timescales:
A few 100 Myr to a few Gyr

...and the path to coalescence today



...and the path to coalescence today



**Reference timescales:
A few 100 Myr to a few Gyr**

Real galaxies are complex environments

Dynamical friction (against dark matter, gas, stars)

Chandrasekhar43, Capelo&Dotti15, Lupi+15, Pfister+17, Tamfal+17, 21, Bonetti+20, 21[incl.EB], Gualandris+ 22[incl.EB]

Fiacconi+13, Roskar+15, Tamburello+17, Souza-Lima+17

Clump scattering

Roskar+15, EB+20, 22

Effect of bars/spirals

Quinlan96, Sesana+06, Sesana&Khan10, Milosavljević&Merritt03, Khan+11, Vasiliev+15, EB+16,18

Bonetti+18, 19, Mannerkoski+21

Armitage&Natarajan02, Escala+04, Dotti+07, Mayer+07, Lodato+09, Souza-Lima+20[incl.EB]

Disk-driven migration

Goicovic+17, Moody+19, Munoz+19,20, Duffel+20,Tiede+20, Zrake+21, Franchini+21, D'orazio&Duffel21, EB+21

Peters64, Hills Zwick20,21[incl.EB]

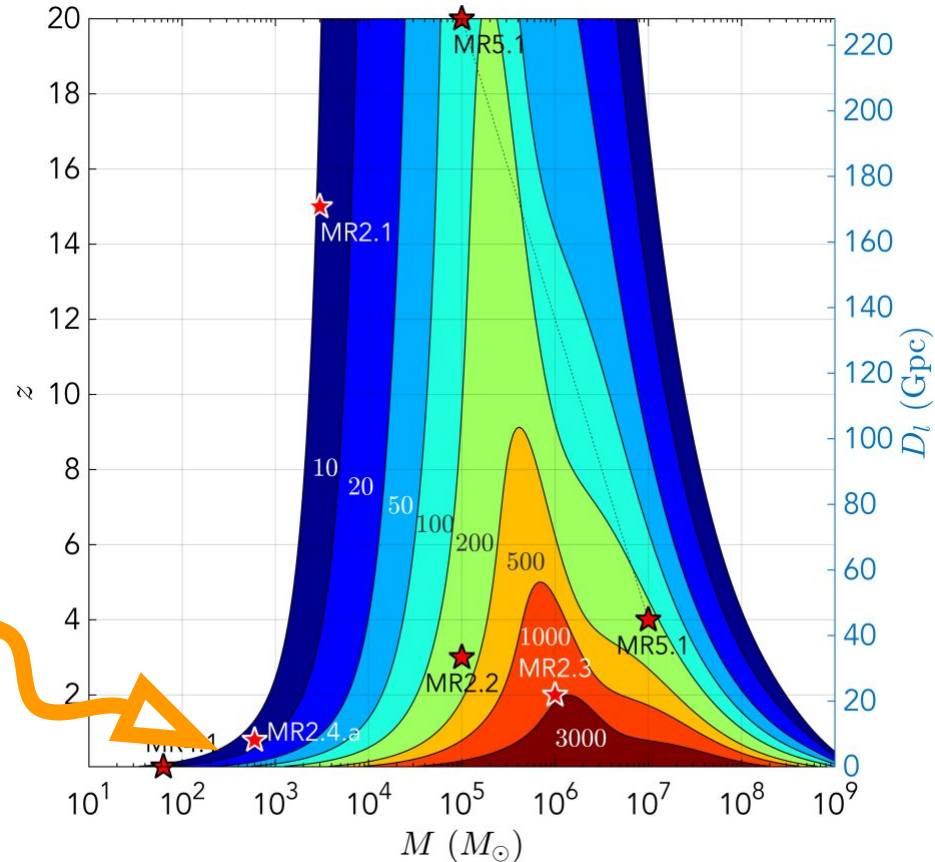
~bound binary

ance mas



Img credits: Wikipedia

LISA Consortium17



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Armit



-driven m

EB+21

Peters64, Hills

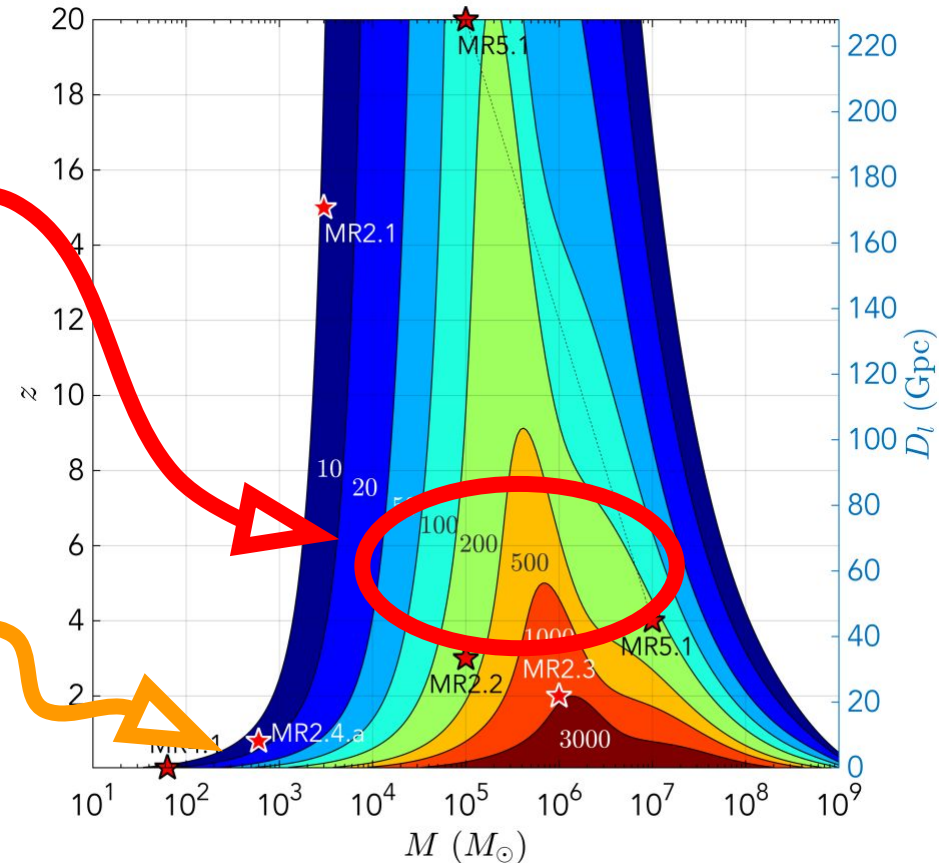
~bound binary

nce mas



Img credits: wikipedia

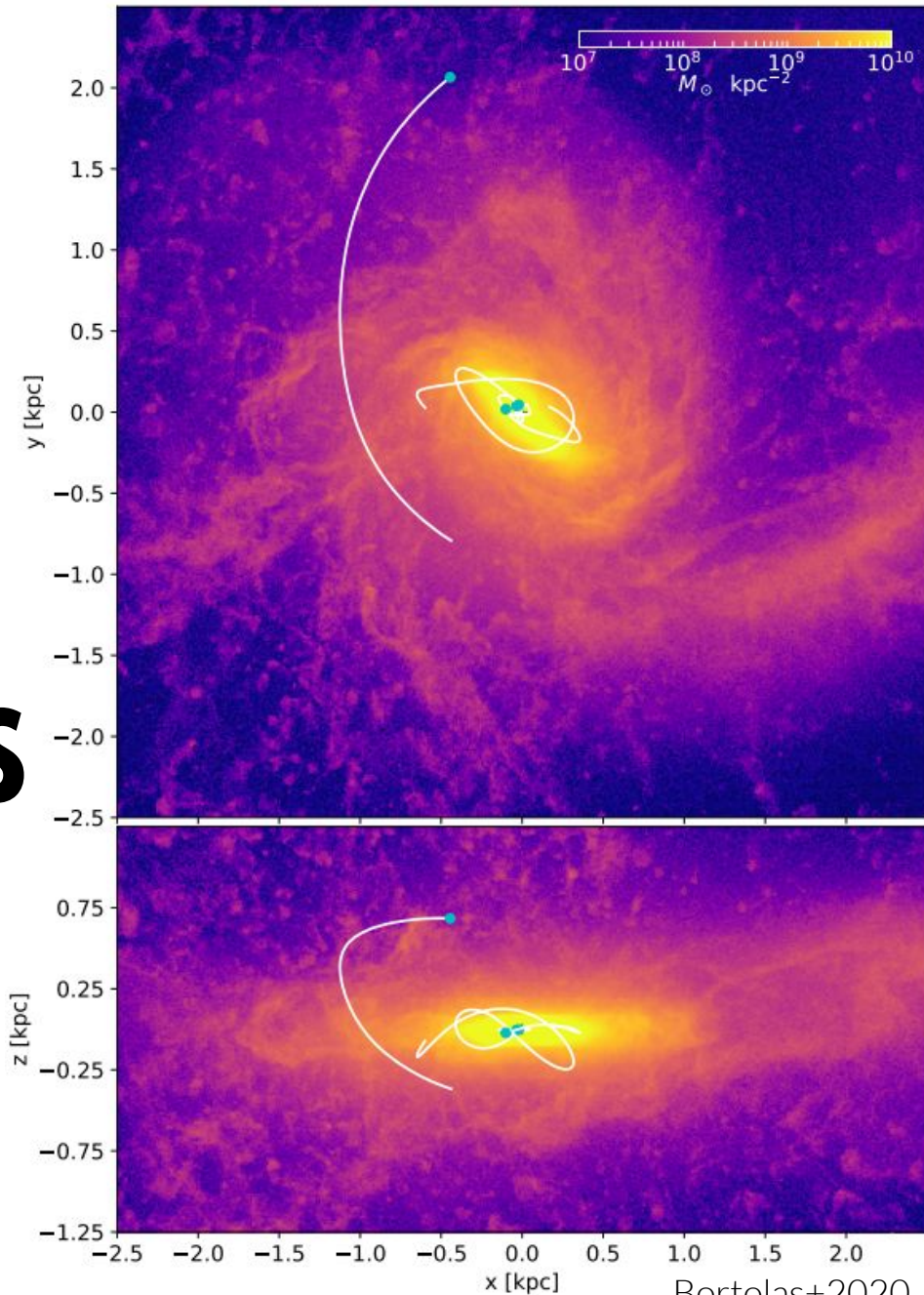
LISA Consortium17



The LARGE scale inspiral is the dynamical friction treatment good enough?

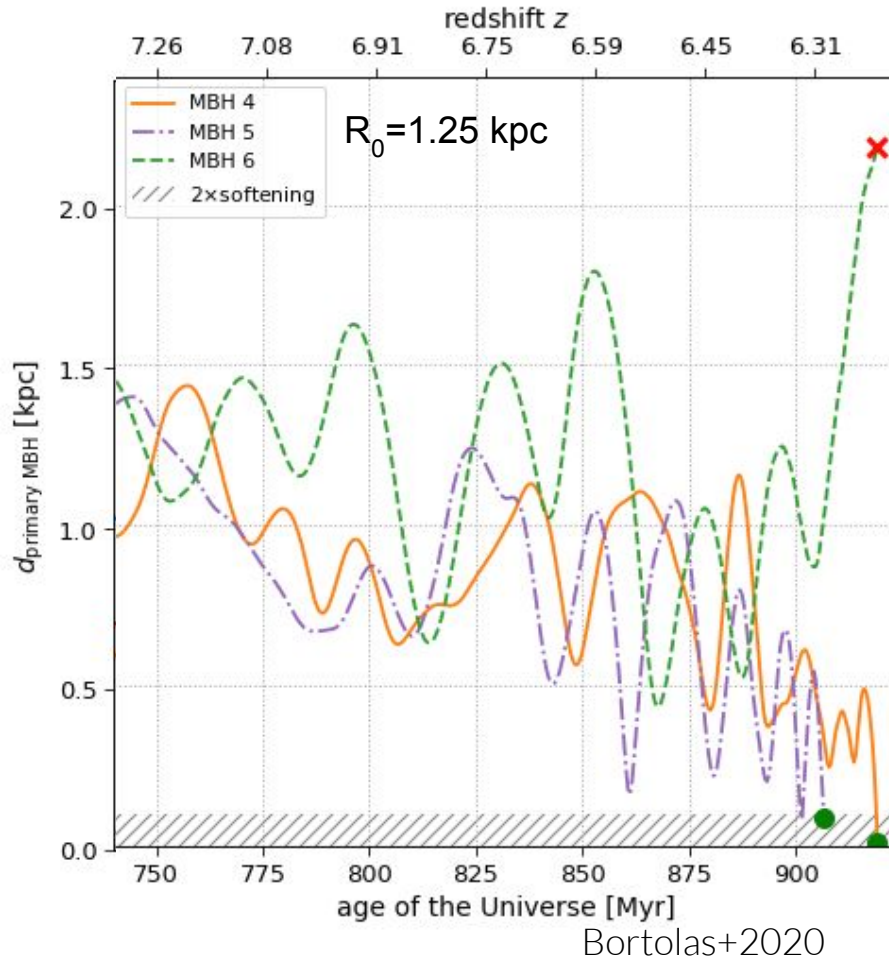


VS



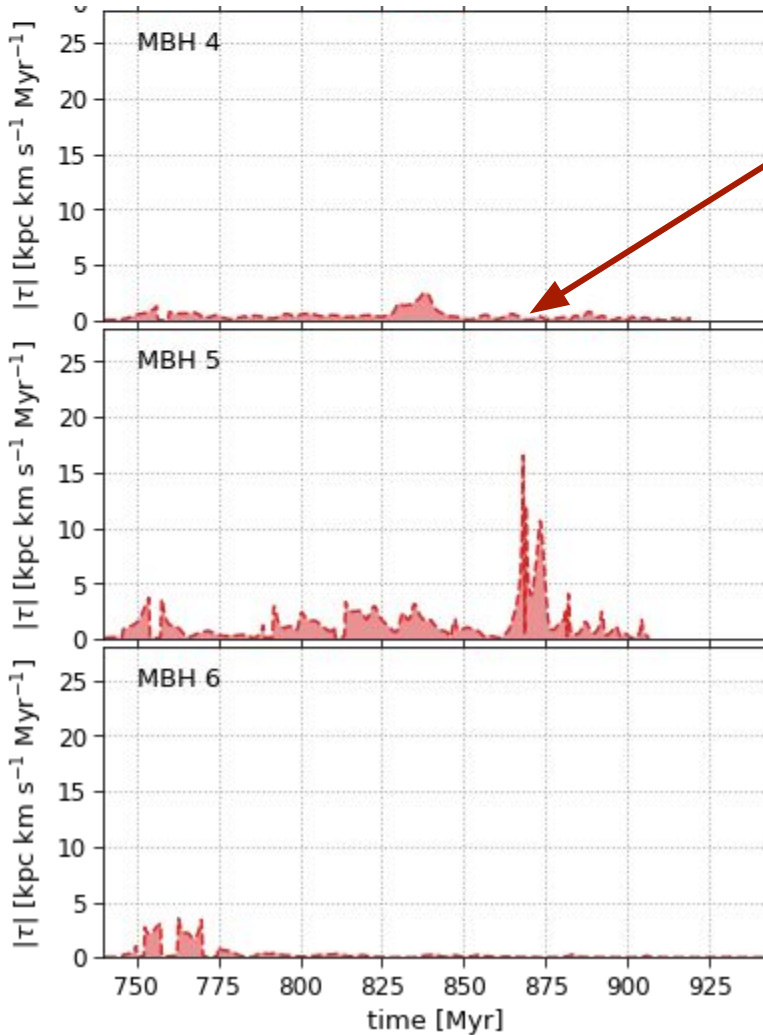
Bortolas+2020

BHs distance from the centre



TORQUES MATTER!

Also at large scale



Dynamical friction torque modulus

From all components (dark matter, stars, gas)
 Semi-analytical calculation in a sphere of 150 pc around each MBH

Gas

$$\mathbf{F}_{DF}^{gas} = -4\pi \ln \left[\frac{b_{max}}{b_{min}} \frac{(\mathcal{M}^2 - 1)^{1/2}}{\mathcal{M}} \right] G^2 M_{BH}^2 \rho_{gas} \frac{\mathbf{V}}{V^3}$$

Ostriker99

$b_{min} = GM/v^2 \sim 0.1 \text{ pc (physical)}$
 $b_{max} = 2 \text{ kpc (2 x disc scale radius)}$

Stars, DM

$$\mathbf{F}_{DF}^{stars, DM} = -4\pi \ln \Lambda G^2 M_{BH}^2 \rho_*$$

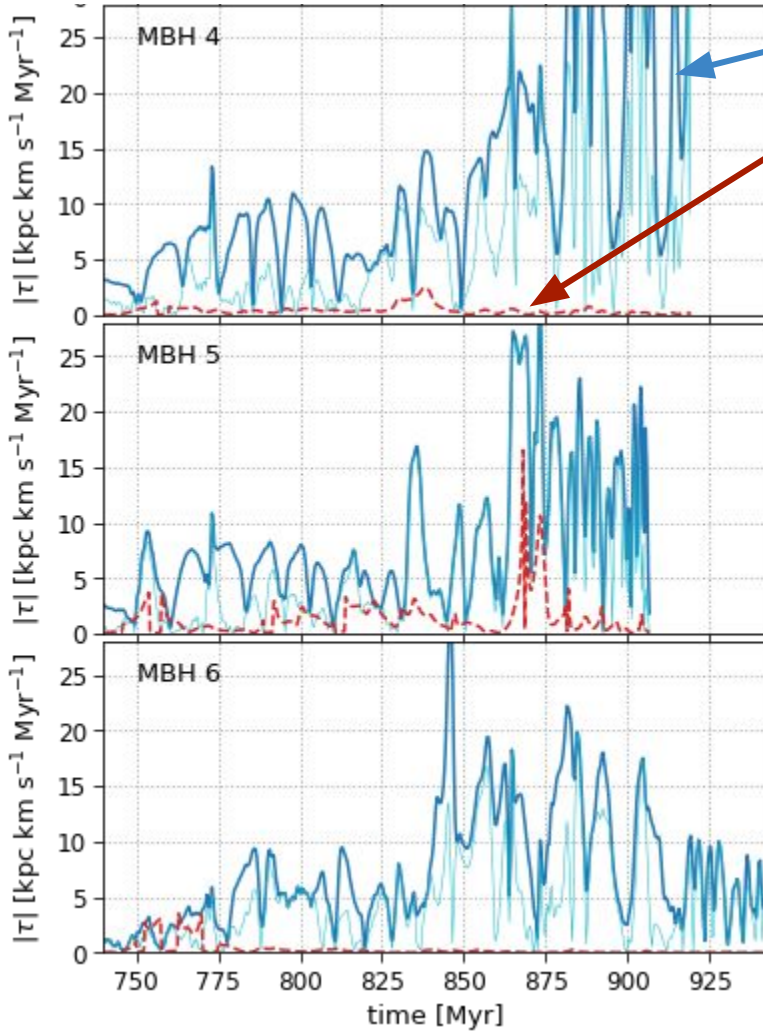
Chandrasekhar43

$$\left[\text{erf} \left(\frac{V}{\sqrt{2}\sigma} \right) - \left(\sqrt{\frac{2}{\pi}} \frac{V}{\sigma} \right) \exp \left(-\frac{V^2}{2\sigma^2} \right) \right] \frac{\mathbf{V}}{V^3}$$

Bortolas+2020

TORQUES MATTER!

Also at large scale



Global torque modulus

Dynamical friction torque modulus

From all components (dark matter, stars, gas)
Semi-analytical calculation in a sphere of 150 pc around each MBH

Reference: Bortolas et al. 2020, MNRAS, 493, 1017

Gas

$$\mathbf{F}_{\text{DF}}^{\text{gas}} = -4\pi \ln \left[\frac{b_{\text{max}} (\mathcal{M}^2 - 1)^{1/2}}{b_{\text{min}} \mathcal{M}} \right] G^2 M_{\text{BH}}^2 \rho_{\text{gas}} \frac{\mathbf{V}}{V^3}$$

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Chandrasekhar43

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Bortolas+2020, see also Bortolas+22

Doggy bag #1

The simple dynamical friction treatment for massive black holes inspiral may be poor in realistic galaxies especially:

- *At high z*
- *In irregular/barred galaxies*

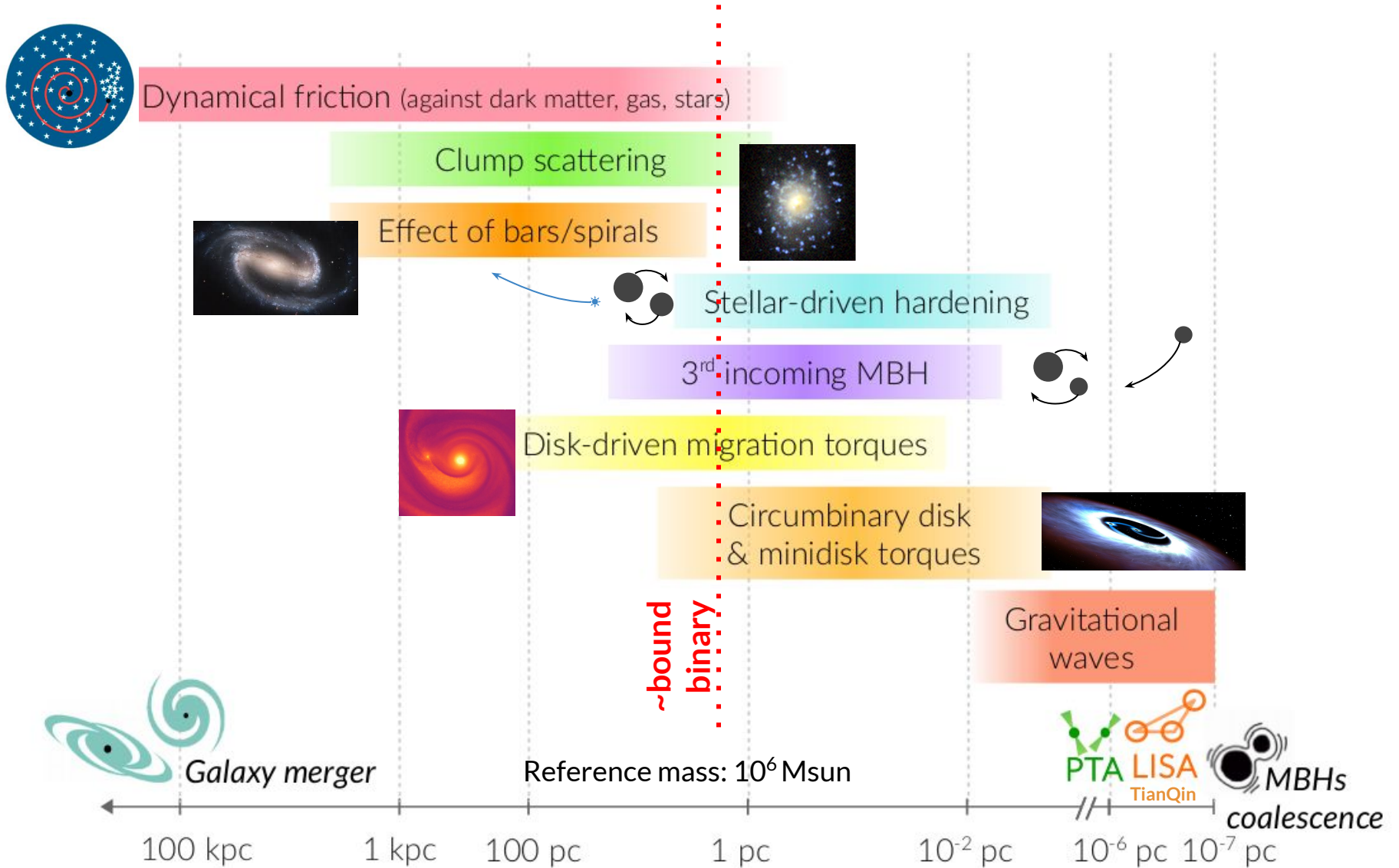
May I chip in...

5 cents for some discussion?

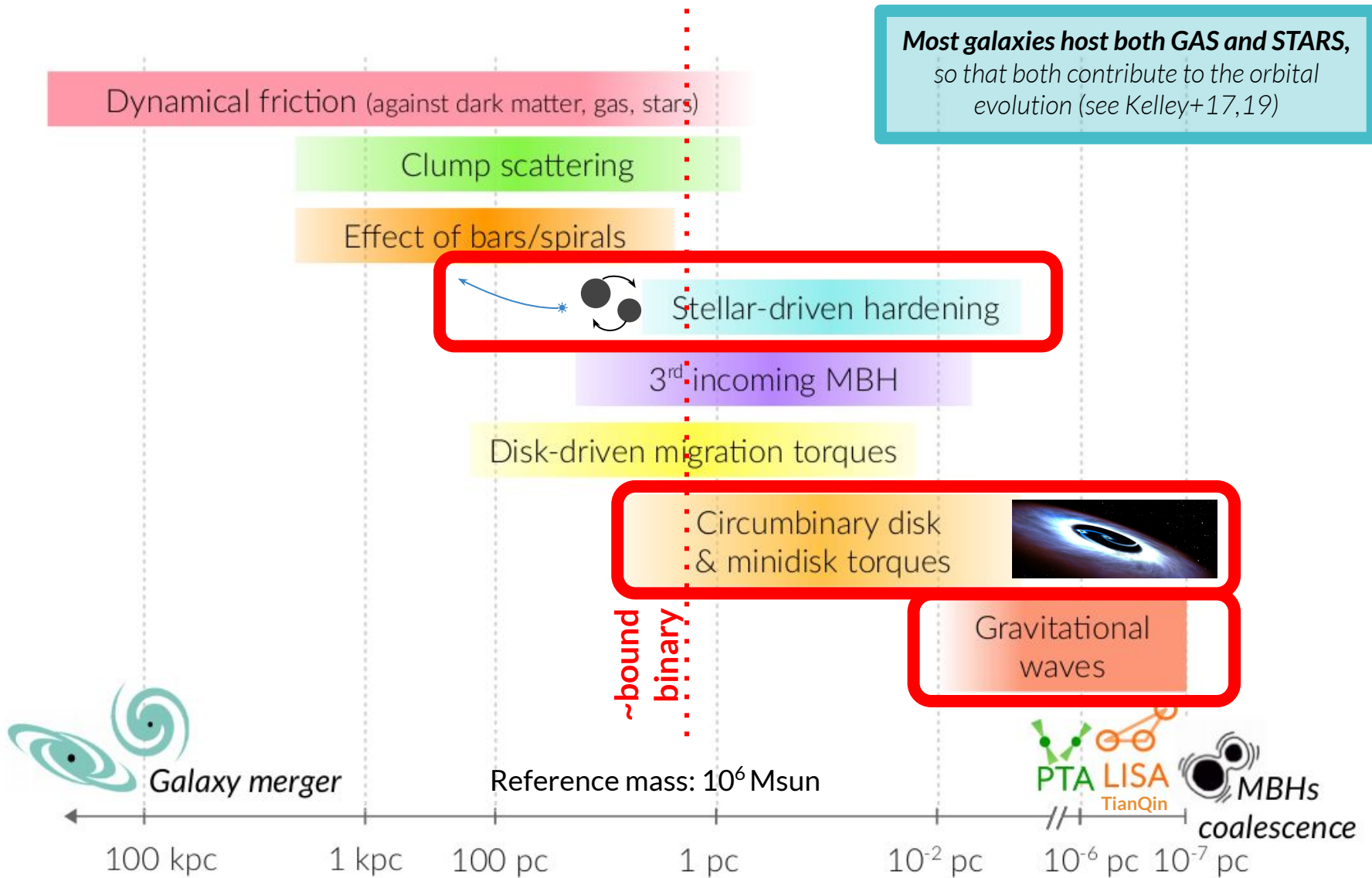
- *How to model this stochasticity in inexpensive semi-analytical models for studying the binary population and merger rates?*



Small scale (bound binary) evolution

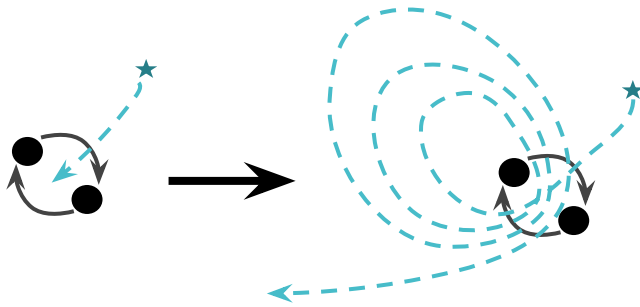


Small scale (bound binary) evolution

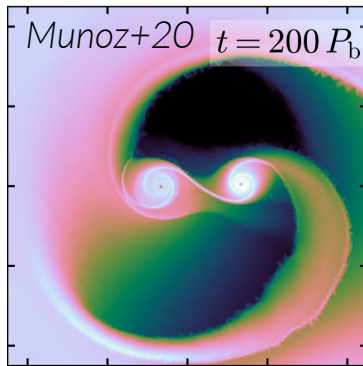


LET'S HAVE THE MOST PESSIMISTIC APPROACH

- **Stellar driven hardening** → Binary shrinking



- **Gas-driven evolution** → Shrinking or **expansion**?



Artymowicz&Lubow94, Moody+19,
Munoz+19,20, Duffel+20, Tiede+20,
Heath&Nixon20, Franchini+21,
D'orazio&Duffel21+.... Almost
everyone in the audience 😊



- **Gravitational wave emission** → Efficient small scale shrinking



• Stellar driven hardening

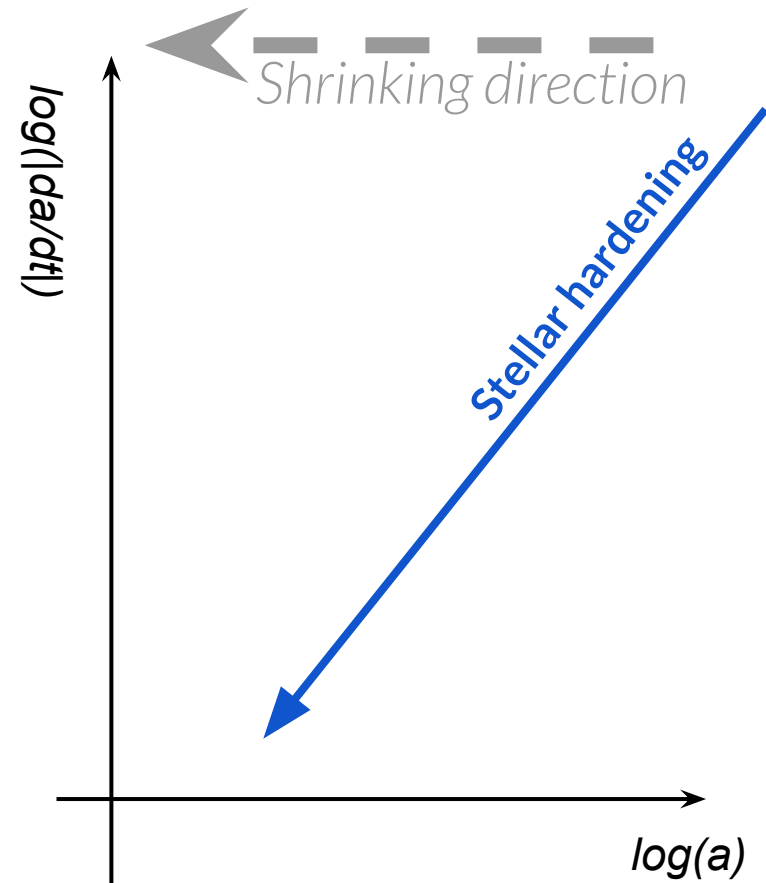
$$\frac{da_{\star}}{dt} \propto -a^2$$

No nuclear star cluster
M-sigma relation
(Kormendy&Ho13, Merrit+09)

Binary
shrinking

$$\dot{a}_{\star} = -\frac{HG\rho}{\sigma} a^2$$

Quinlan96, Sesana+06,
Sesana&Khan15



• Stellar driven hardening

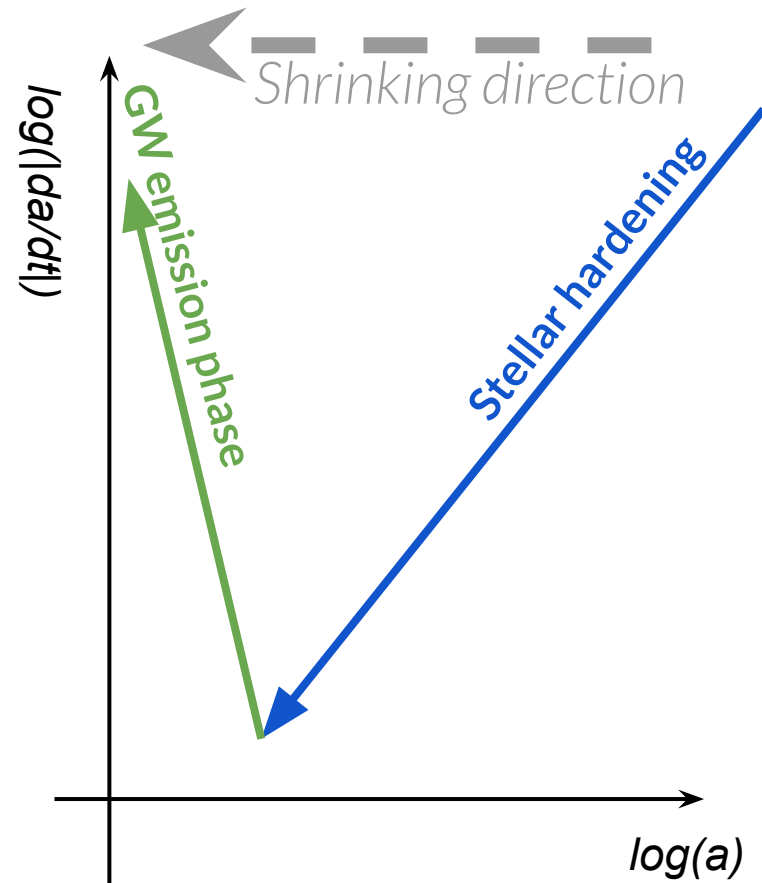
$$\frac{da_{\star}}{dt} \propto -a^2$$

Binary
shrinking

• Gravitational wave emission

$$\frac{da_{\text{GW}}}{dt} \propto -a^{-3}$$

Efficient
shrinking at
small scale



$$\dot{a}_{\text{GW}} = -\frac{64}{5} \frac{G^3}{c^5} \frac{q}{(1+q)^2} \frac{m^3}{a^3}$$

Peters64

- Stellar driven hardening

$$\frac{da_{\star}}{dt} \propto -a^2$$

Binary
shrinking

- Gas-driven evolution

$$\frac{da_{\text{gas}}}{dt} \propto +a$$

We assume
binary
expansion

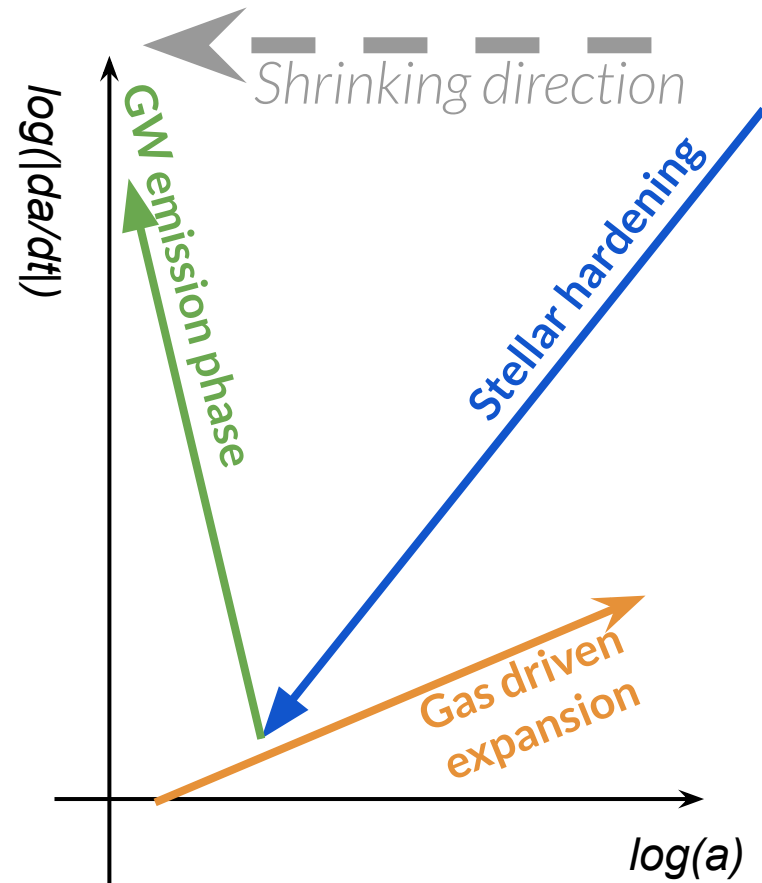
$$\dot{a}_{\text{gas}} = 2.68 \frac{\dot{m}}{m} a$$

Munoz+20

- Gravitational wave emission

$$\frac{da_{\text{GW}}}{dt} \propto -a^{-3}$$

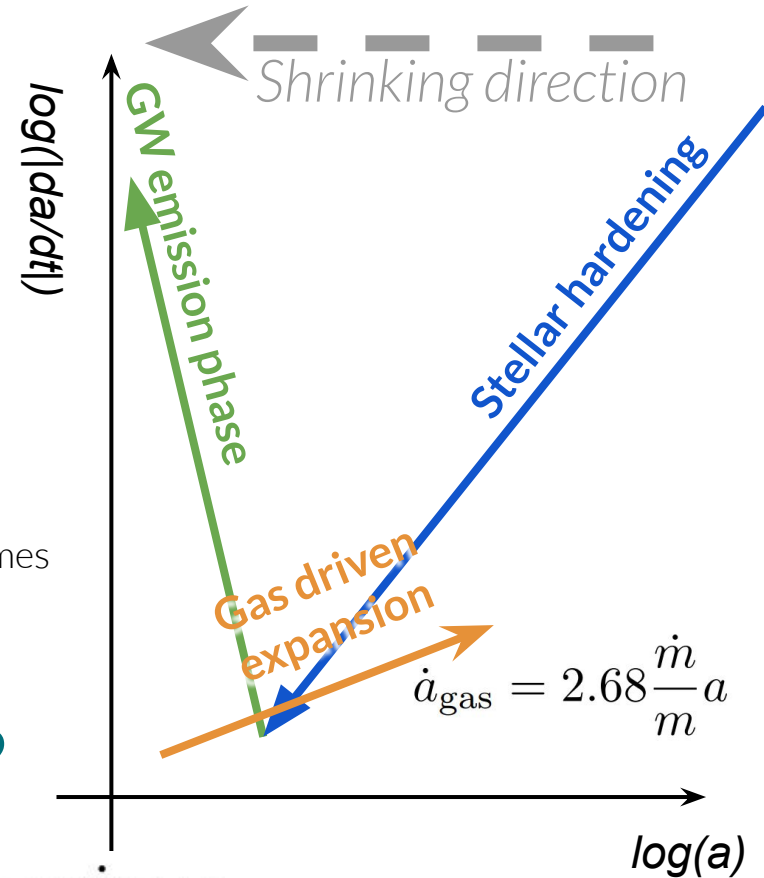
Efficient
shrinking at
small scale



When is the gas effective?

$$\left\{ \begin{array}{l} \dot{a}_* = -\frac{HG\rho}{\sigma} a^2 \\ \dot{a}_{\text{gas}} = 2.68 \frac{\dot{m}}{m} a \\ \dot{a}_{\text{GW}} = -\frac{64 G^3}{5 c^5} \frac{q}{(1+q)^2} \frac{m^3}{a^3} \end{array} \right.$$

Some more assumptions:
 Equal mass binary;
 Binary eccentricity = 0 at all times

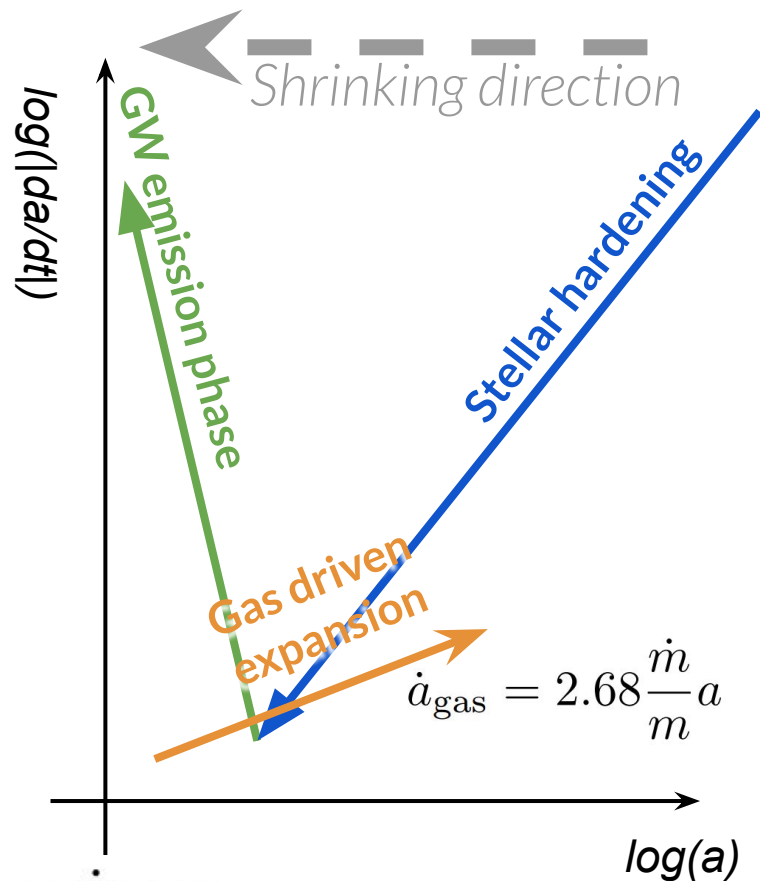
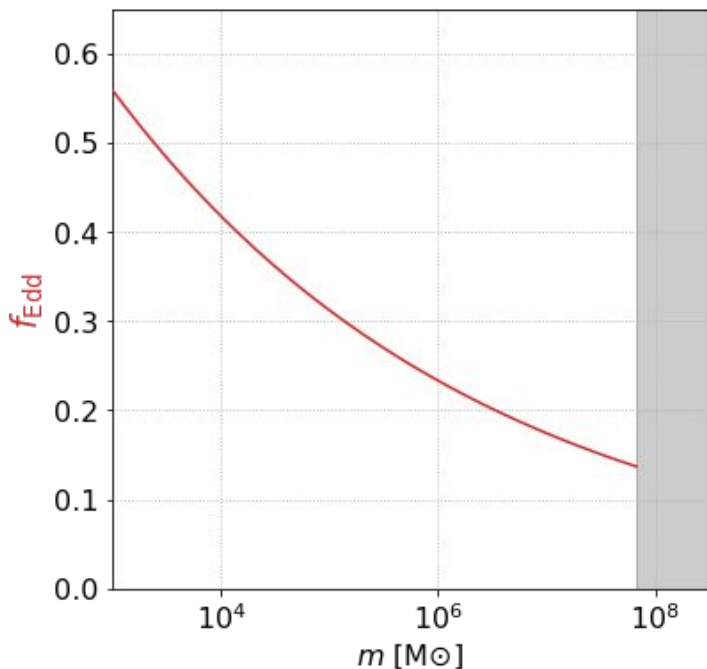


➔ How do we model accretion?

- **Fixed eddington ratio** f_{Edd} so that $\dot{m} = f_{\text{Edd}} \dot{m}_{\text{Edd}} \propto m$
 (the accretion rate grows linearly with the binary mass, and remains a fixed fraction of the Eddington accretion rate)

NOTE THAT DIFFERENT ASSUMPTIONS (FIXED MDOT) FOR THE MASS ACCRETION RATE RESULT IN AN EVEN LESS EFFICIENT GAS-DRIVEN EXPANSION

When is the gas effective?



➔ How do we model accretion?

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NOTE THAT DIFFERENT ASSUMPTIONS (FIXED MDOT) FOR THE MASS ACCRETION RATE RESULT IN AN EVEN LESS EFFICIENT GAS-DRIVEN EXPANSION

Can the binary expand indefinitely?

NO!

THE SELF GRAVITATING RADIUS

Franchini+21

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THE SELF GRAVITATING RADIUS

Franchini+21

$$R_{\text{sg}} \propto f_{\text{Edd}}^{-22/45} m^{-7/45}$$

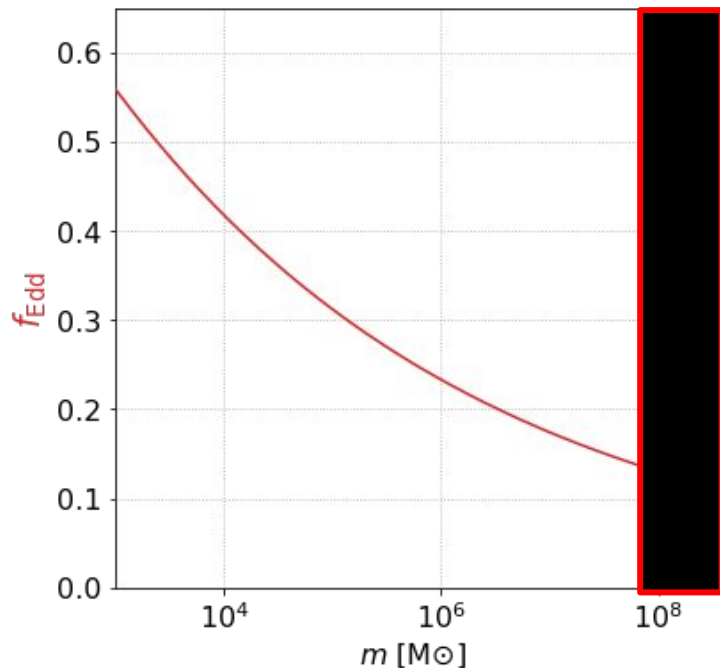
Perego+09

$R_{\text{sg}} \sim 10^{-2}$ pc for a 10^6 solar mass
binary accreting at Eddington

Can the binary expand indefinitely?

NO!

THE SELF GRAVITATING RADIUS

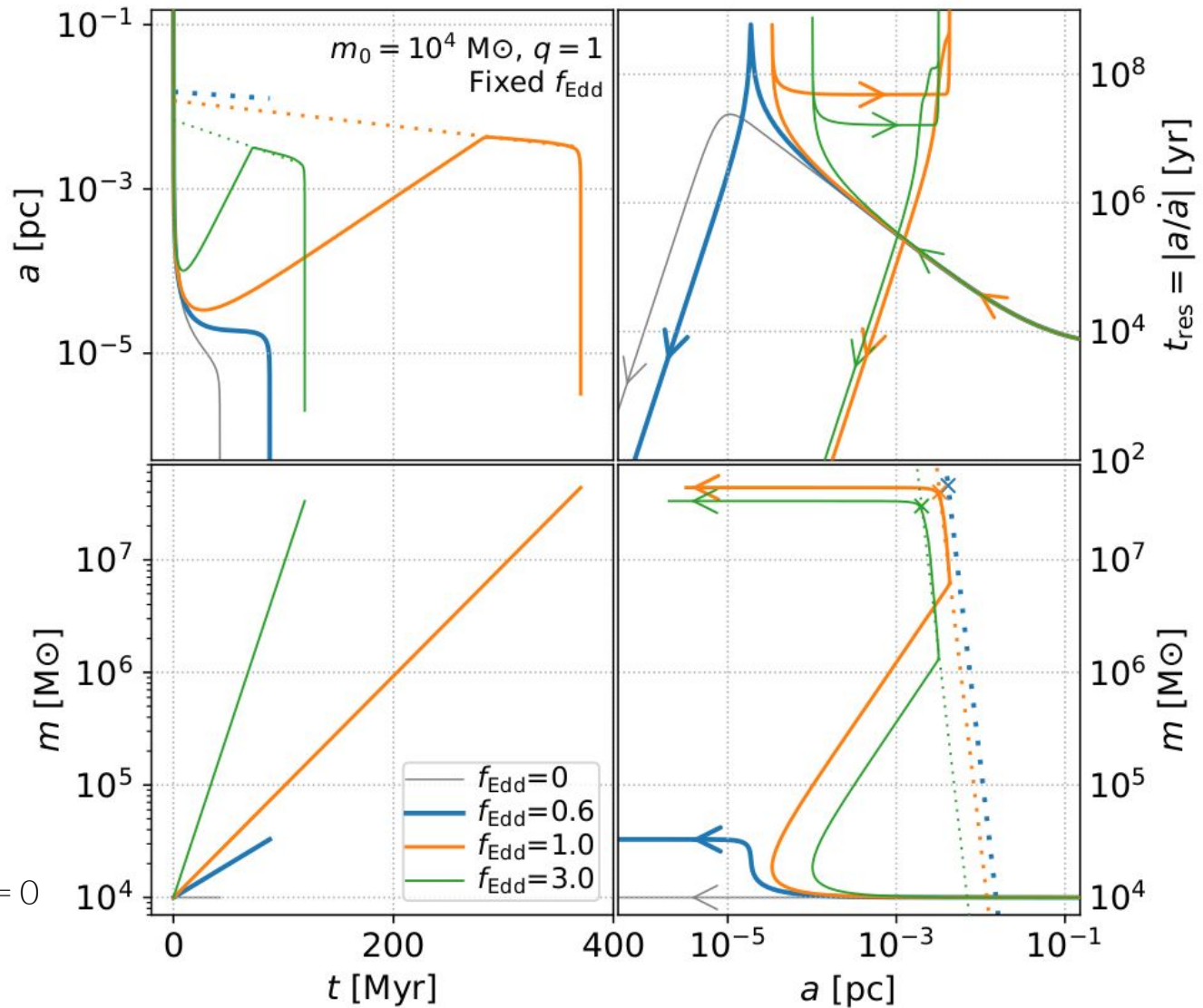


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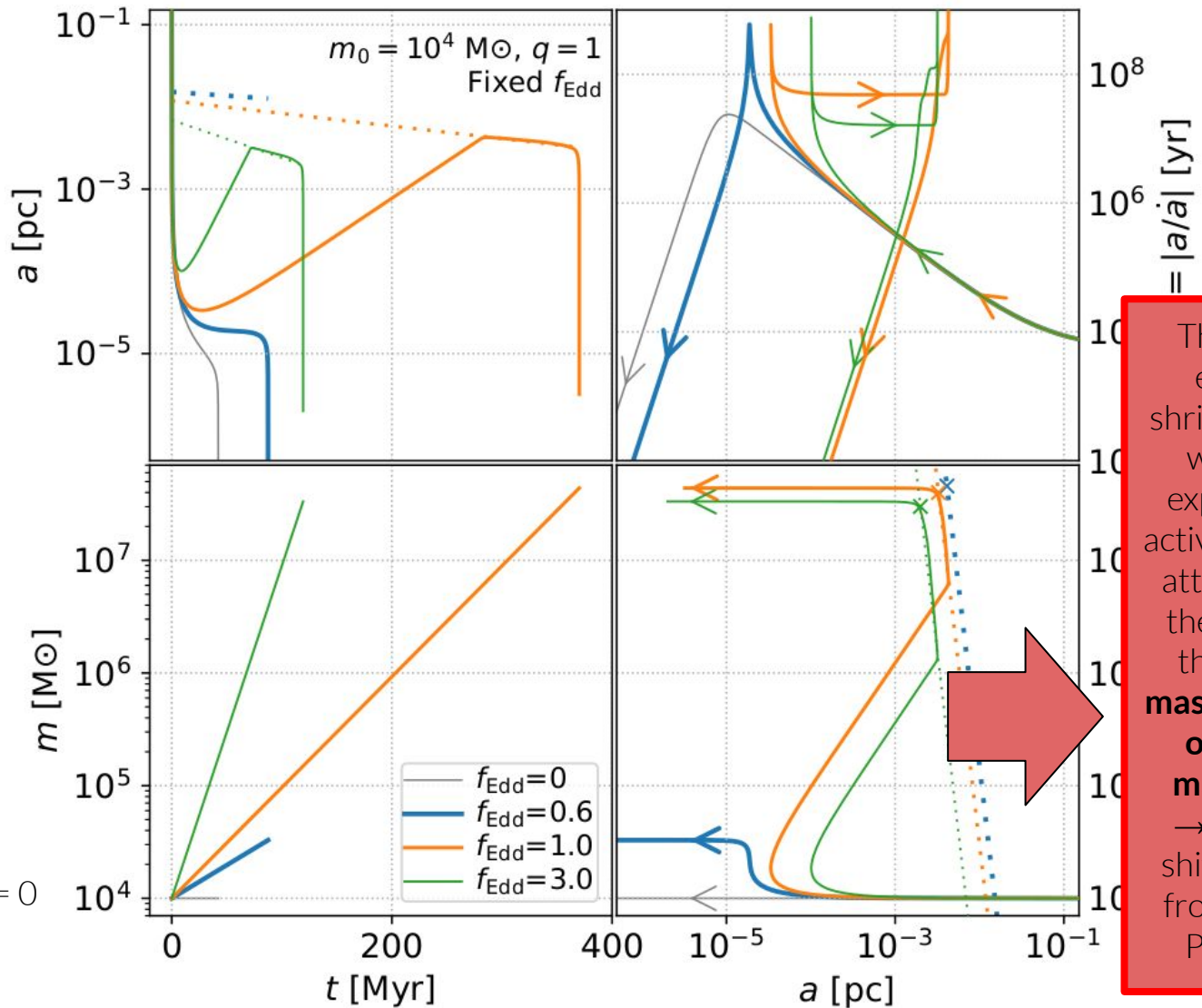
Here the disk is self gravitating down to the scales at which GW always dominate

Evolution for fixed f_{Edd}



Equal mass binary;
 Binary eccentricity = 0
 at all times

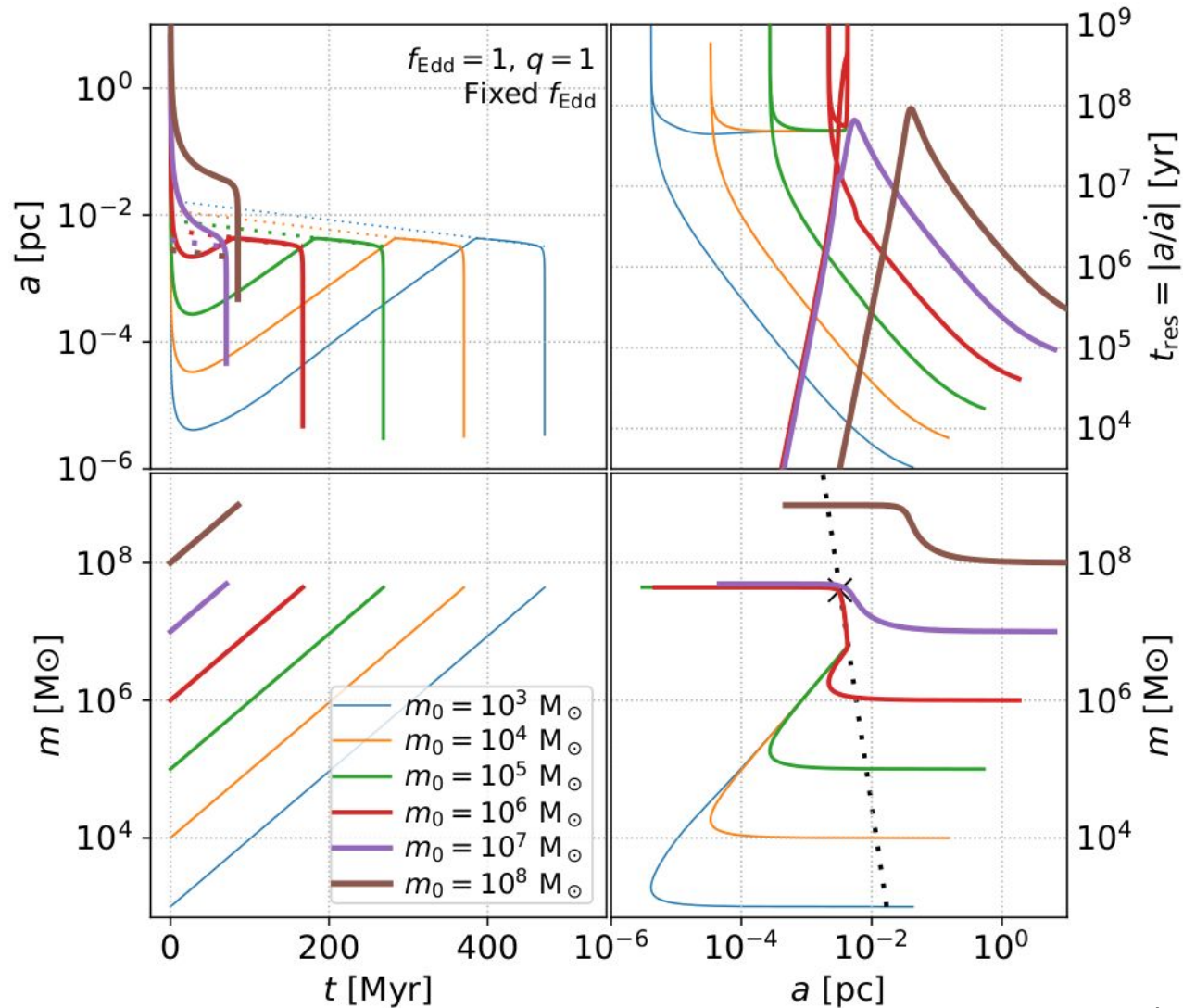
Evolution for fixed f_{Edd}



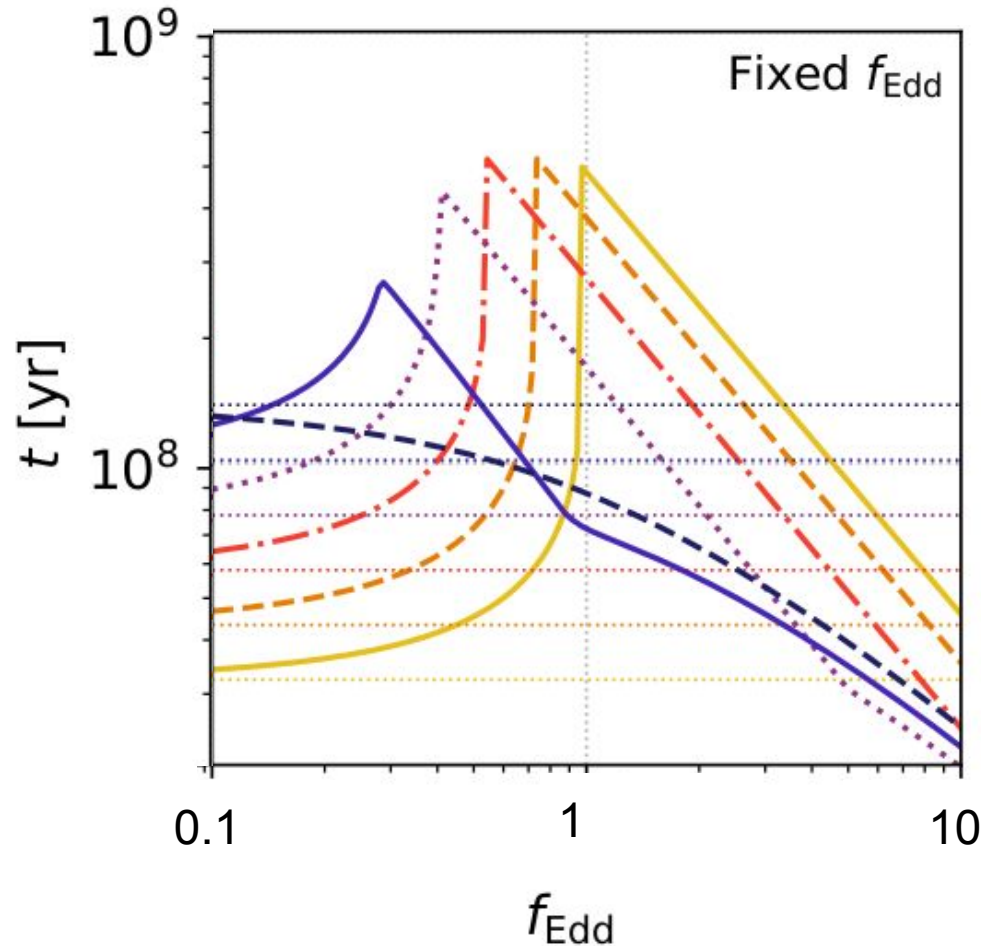
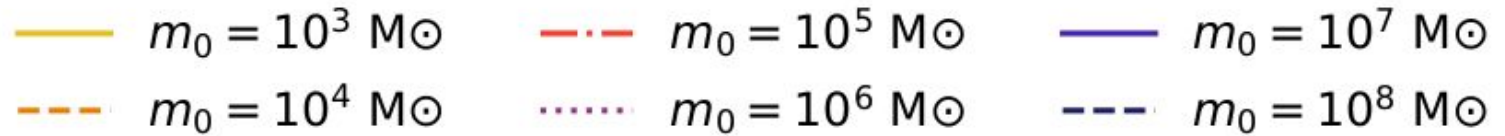
The binary efficient shrinking even when gas expansion is active has to be attributed to the fact that the **binary mass grows by orders of magnitude** → this may shift binaries from LISA to PTA band

Equal mass binary;
 Binary eccentricity = 0
 at all times

Different initial binary mass

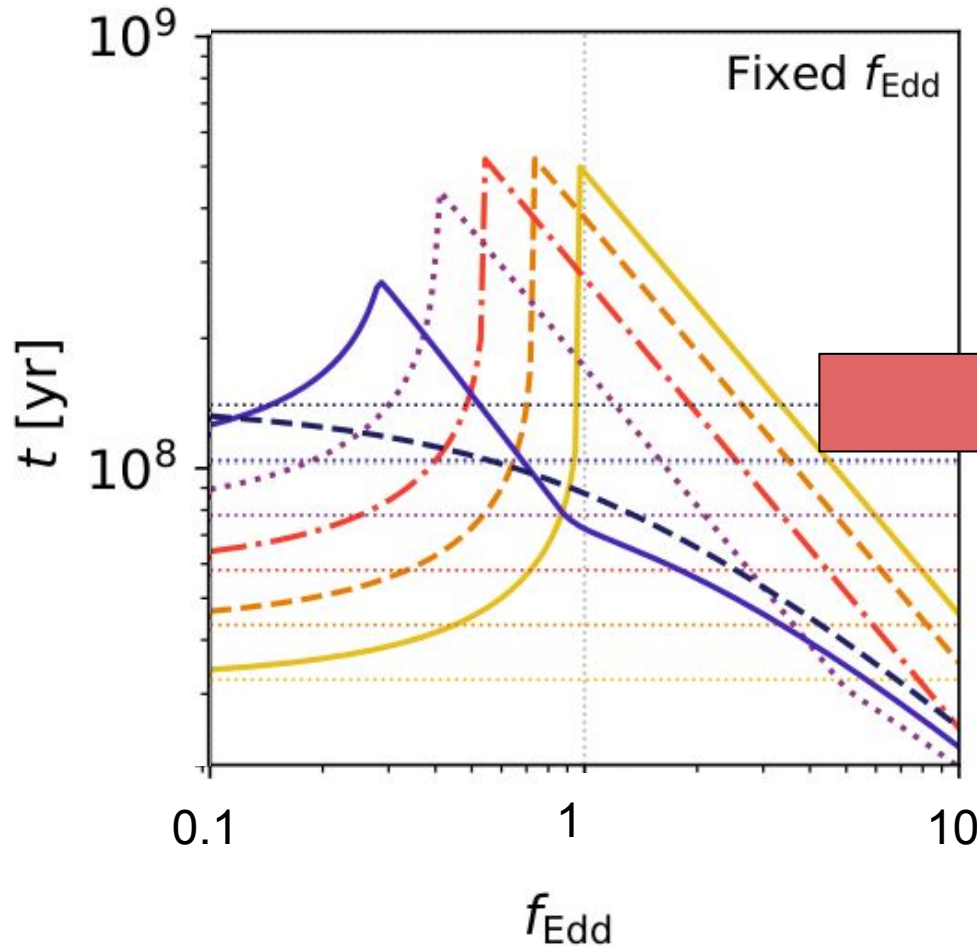
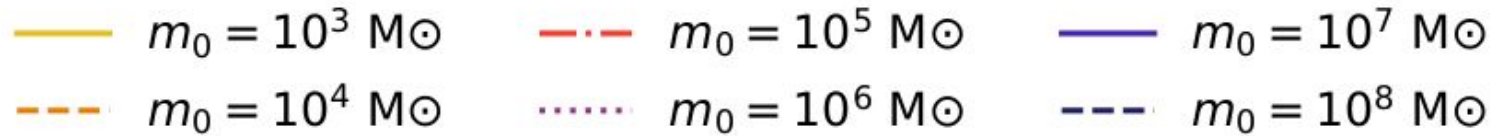


Inspirational timescale



*To be compared with
the timescale of large
scale inspiral, which can
be much longer!*

Inspiral timescale



This can imply that accreting, **expanding binaries can be observable for a longer time** via electromagnetic surveys, but they anyways eventually coalesce!

Doggy bag #2

- **GAS AND STARS GENERALLY COEXIST!!**
- **Gas driven expansion does not dramatically impact the coalescence time of massive binaries**

→ **Expansion necessarily reverts into shrinking** when the binary mass gets large enough owing to accretion: at that point, gravitational wave emission becomes dominant.

May I chip in...

another 5 cents for some discussions/ideas?

- *This would imply expanding binaries would be observed at lower frequencies (PTAs)*
- *We should think of better simulations/works accounting for concurrent effects of stars and gas*
- *Would it be possible to use a similar approach in the framework of stellar binaries? [e.g. tides instead of stellar hardening and so on...]*

