Mesoscopic Physics of Motile Protrusions in Eukaryotic Cells

Garegin Papoian



Department of Chemistry and Biochemistry

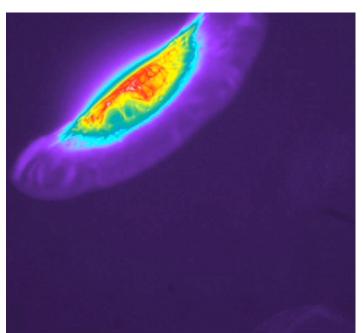
&

Institute for Physical Science and Technology

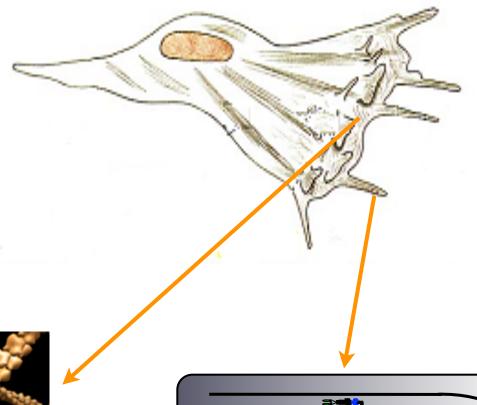
University of Maryland

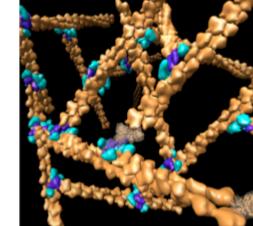
Cell's Cytoskeleton at the Leading Edge

Wadsworth lab at the University of Massachusetts



Motile Cell

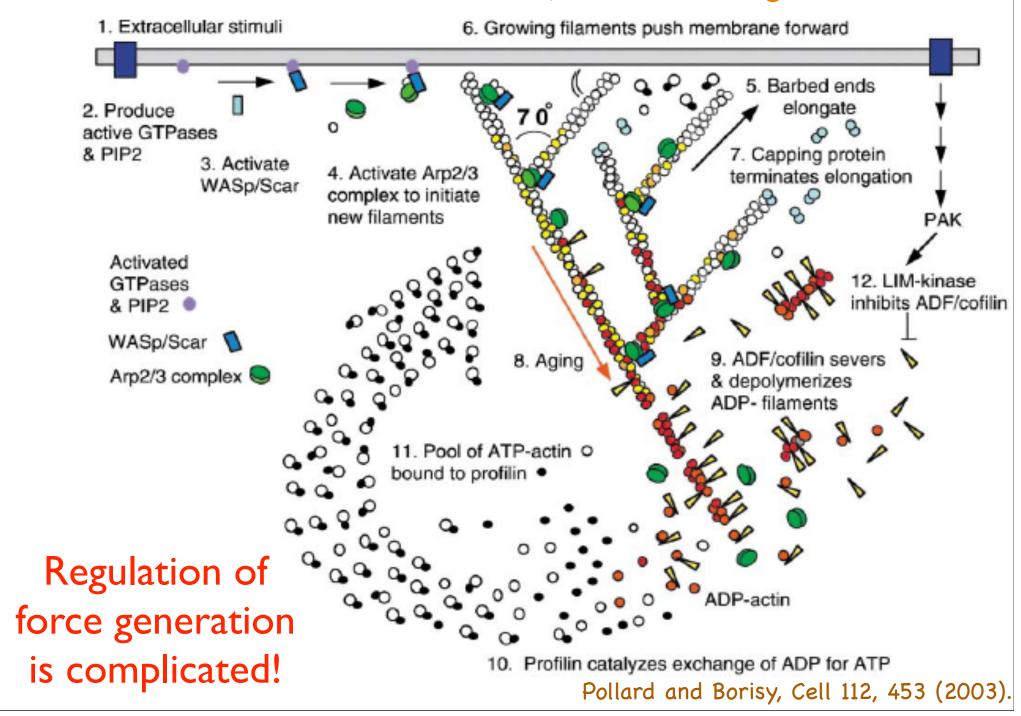




Filopodial Actin (1D)

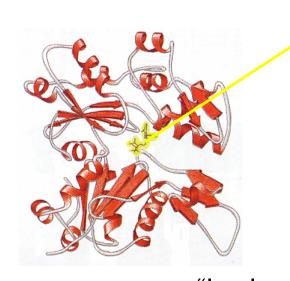
Lamellipodial Actin (3D)

Dendritic nucleation/Array treadmilling model



Actin filaments and monomer diffusion

G-actin 45kDa



"pointed end" end" +

F-actin

37nm

 $2\delta = 5.4$ nm

ATP

Some parameter values

Persistence length $15\mu m$

Buckling length 120nm(10pN)

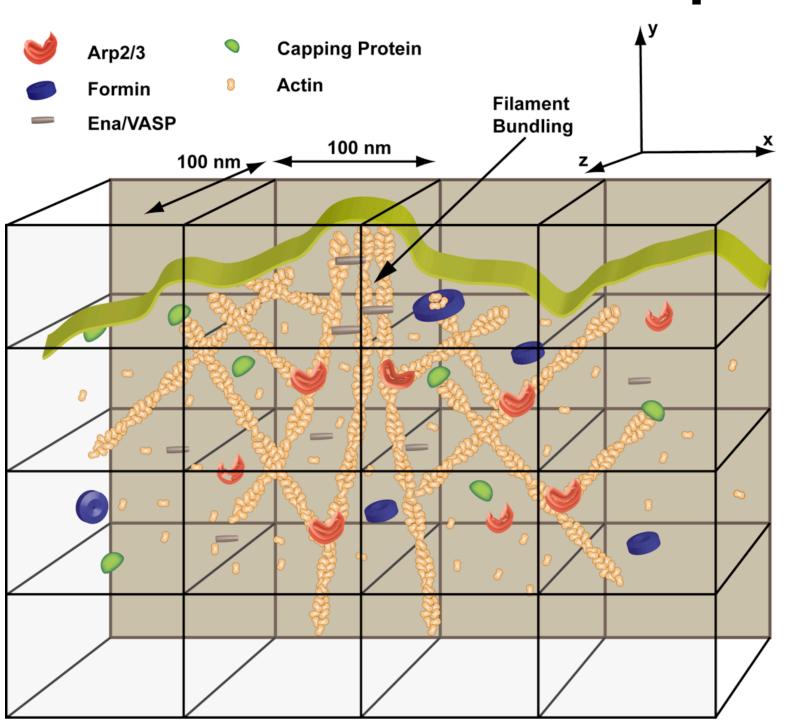
Diffusion rate $5\mu m^2/s$

Bulk concentration $10\mu M$

Polymerization $10 \mu M^{-1} s^{-1}$ rate

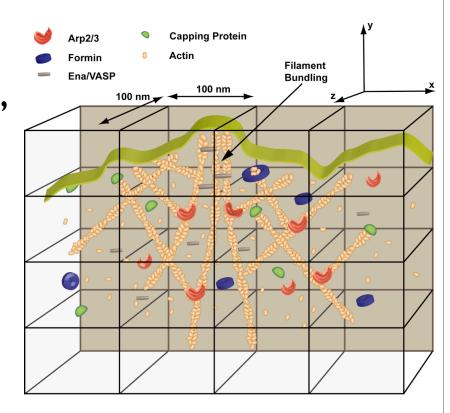
Mogilner and B. Rubinstein, Biophys. J. 89, 782

3D Active Mesh within Lamellipodia



Stochastic simulations of lamellipodia protrusion

- Simulation region is divided into compartments.
- Diffusion (Actin, Capping protein, Arp2/3) between compartments.
- Chemical reactions in compartments:
 - Polymerization, Depolymerization, Capping, Branching...
- Monte Carlo algorithm to generate stochastic trajectories



L. Hu and G. A. Papoian,
 Biophys. J.; 2010, 98,1375

Filaments and membrane

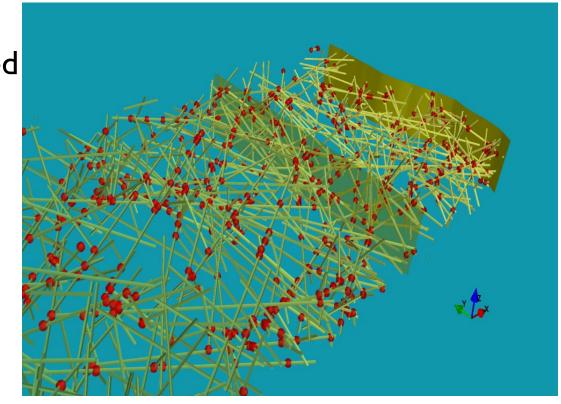
Filaments are assumed to be straight and rigid.

Leading edge membrane is modeled as a 1-D curve x=x(y,z)=h(y) due to the flatness of the lamellipodia:
 ~200nm (z) compared to micron size of the other direction (v)

direction (y).

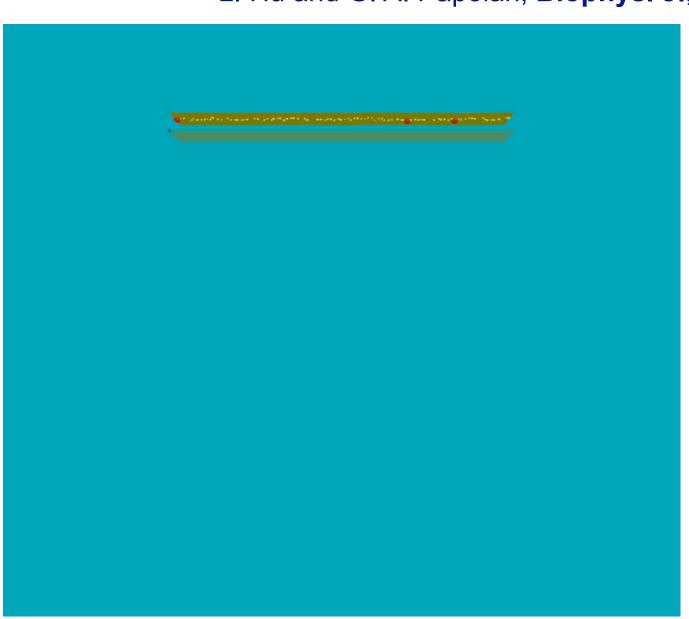
The membrane is modeled as an elastic sheet under tension which also resists bending

Steric repulsion between the filament tips and the membrane



Stochastic Growth of a Lamellipodium

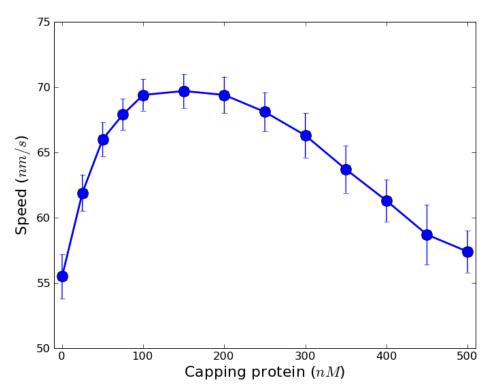
L. Hu and G. A. Papoian, Biophys. J.; 2010, 98,1375–1384



- Tubes indicate growing actin filaments
- Red spheres indicateArp2/3 nucleationpoints
- Diffusing monomeric species are not shown (actin, Arp2/3, and capping proteins)
- Mechano-chemical couplings between the membrane and the filament growth

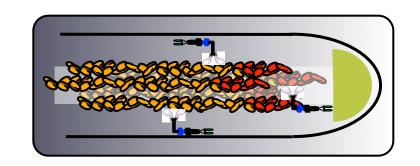
Capping protein enhances motility

- Capping proteins block the polymerization of actin filaments.
- However, capping protein can enhance motility. The mechanism?

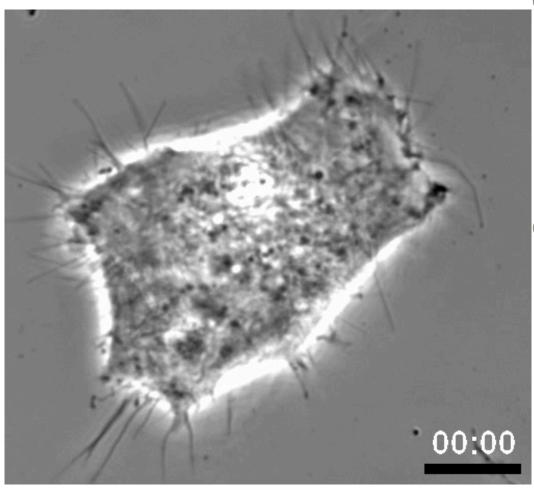


Experimental studies on capping protein promotes motility: Carlier and Pantaloni, JMB 269, 459 (1997). Loisel et al., Nature 401, 613 (1999). Akin and Mullins, Cell 133, 841 (2008).

Numerous Long Filopodia Grow in HeLa Cells

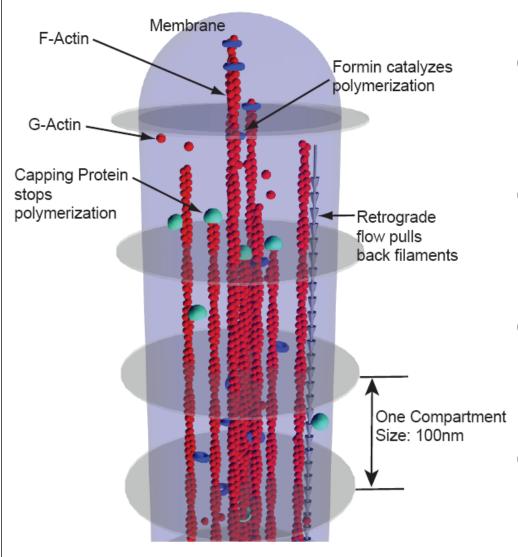


Richard E. Cheney and coworkers at UNC-CH, Proc. Natl. Acad. Sci. USA (2006) v 103, pp 12411



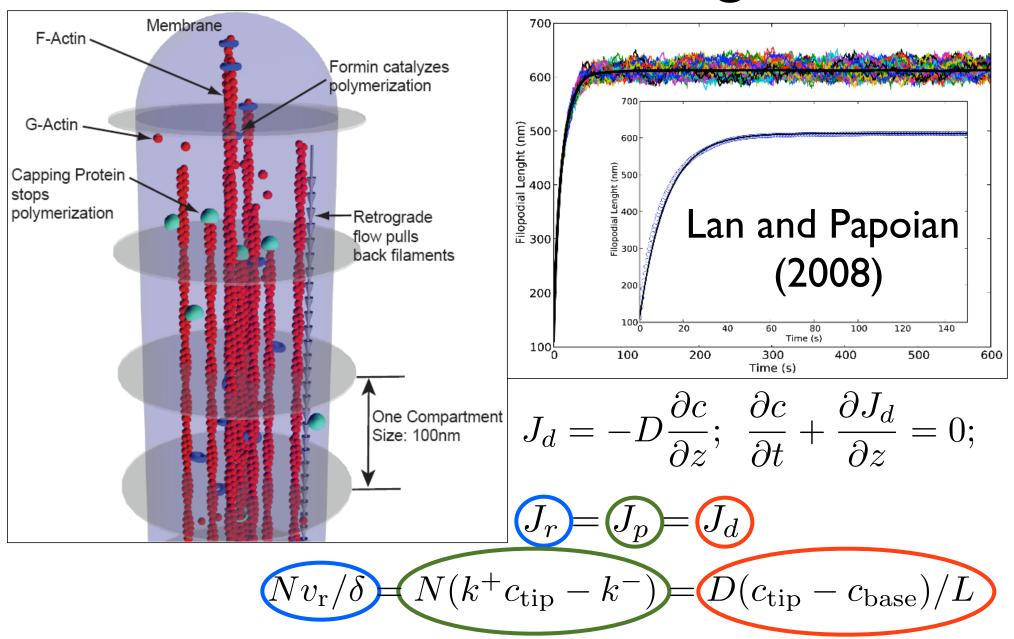
- Biological role of filopodia
 - Guiding cell movement
 - Embryonic development, wound healing, cancer spreading
- Filopodial structure
 - Parallel actin filaments, enveloped by a membrane
 - Cross-linking proteins
 - Tip protein complex

Our Computational Model for Filopodia



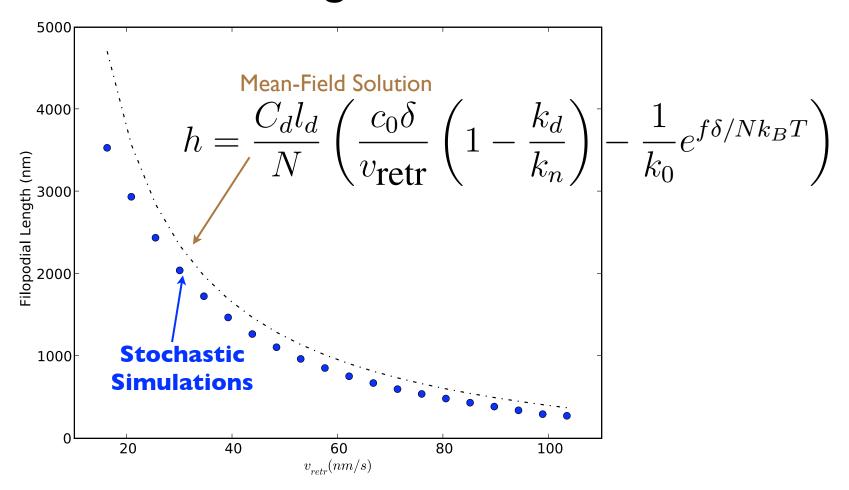
- Stochastic model for both the polymerization and diffusion
- The filopodial tube is partitioned into 50 nm compartments
- G-actin molecules "hop" between neighboring compartments
- Uneven loading of the membrane force among the filaments in the bundle
- Retrograde flow pulls back the actin filaments from the filopodial tube to the cell's body

What limits the length?



Y. Lan and G.A. Papoian. The stochastic dynamics of filopodial growth. Biophys J (2008) vol. 94 (10) pp. 3839-3852

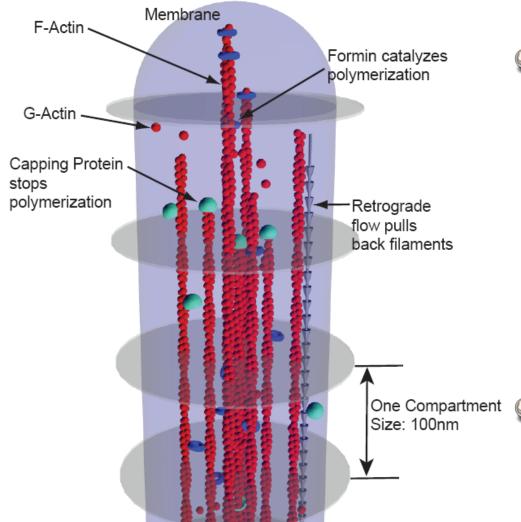
Retrograde Flow



- Filopodia can grow very long at low retrograde flow speeds
- The discrepancy between the mean-field and stochastic solutions can be quite large at low retrograde flow rates

Y. Lan and G. A. Papoian, **Biophys. J.**, (2008) 94, 3839-3852

Adding Capping Proteins & Formins

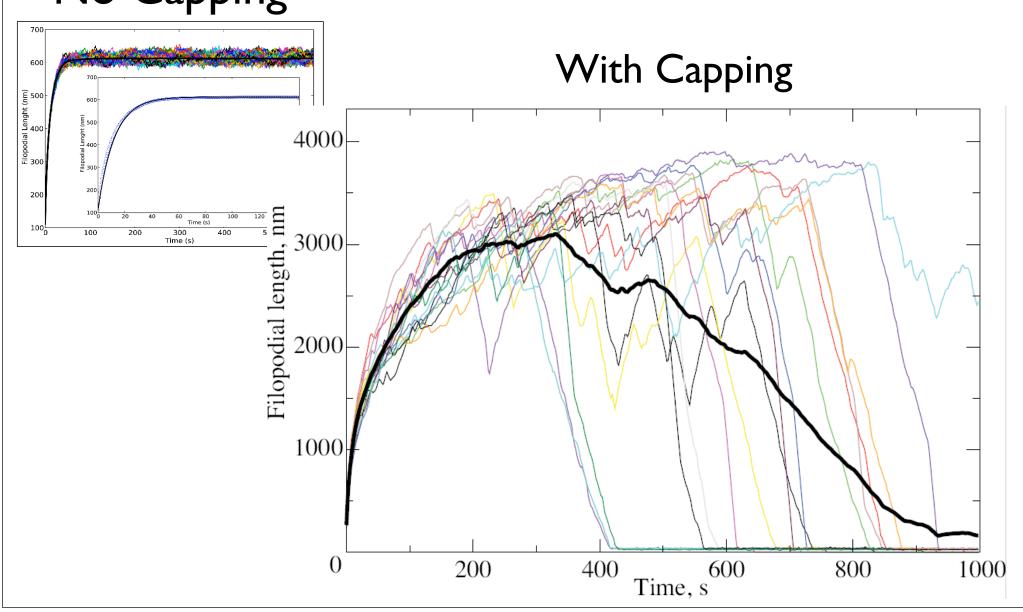


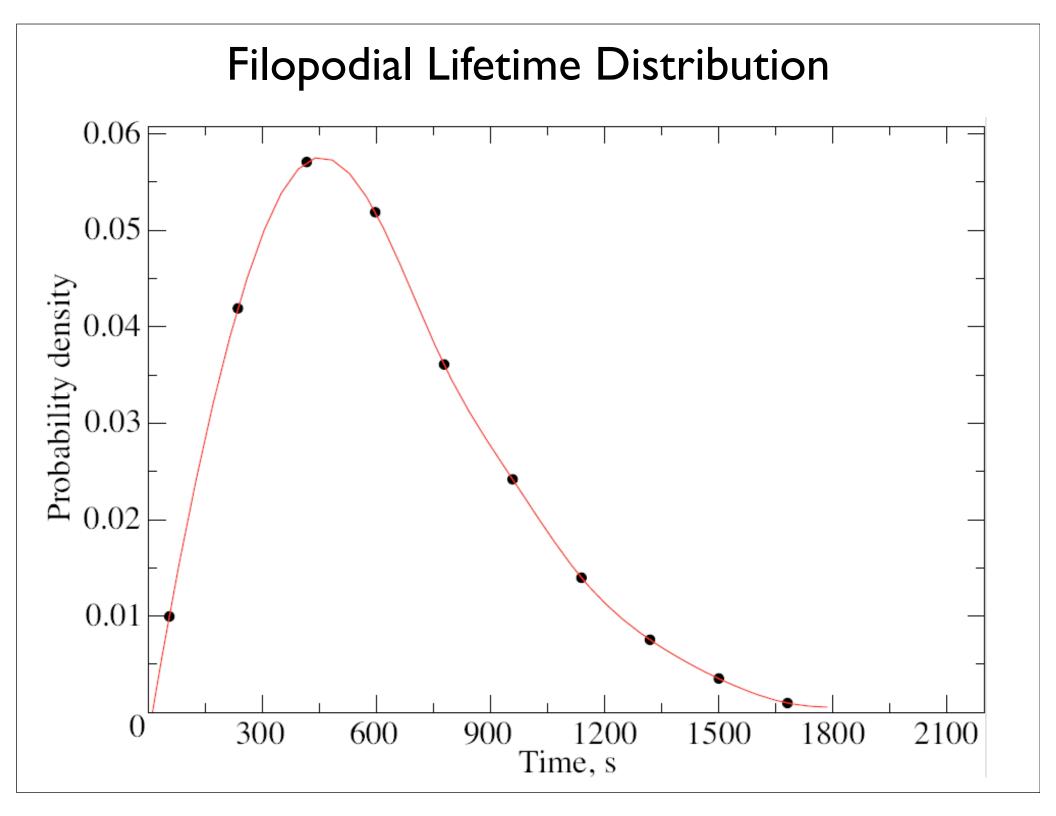
- Capping proteins, when bound to filament tips, arrest polymerization
 - Consequently, the filament starts to retract
 - It may even completely disappear
- Formins accelerate polymerization

□ P. I. Zhuravlev and G. A. Papoian, Proc. Natl. Acad. Sci. USA; 2009, 106, 11570-11575

Macroscopic oscillations of filapodial length induced by low concentration of a capping protein



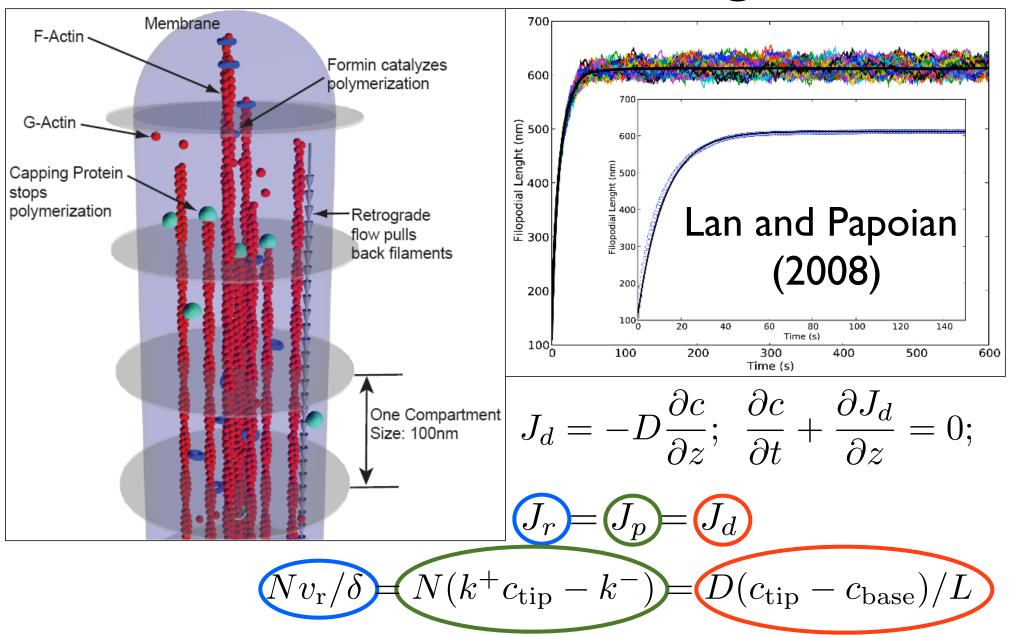




The filopodial length

- Modeling:
 - 0.5 2 μm length,
 - growth speed on µm/min scale
 - stationary
- Experiment:
 - typically | | 0 μm,
 - up to 100 μm with 10 μm/min growth speed
 - growth retraction cycles

What limits the length?



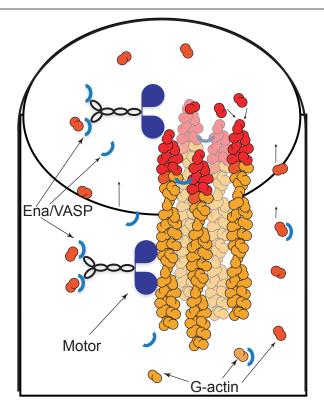
Y. Lan and G.A. Papoian. The stochastic dynamics of filopodial growth. Biophys J (2008) vol. 94 (10) pp. 3839-3852

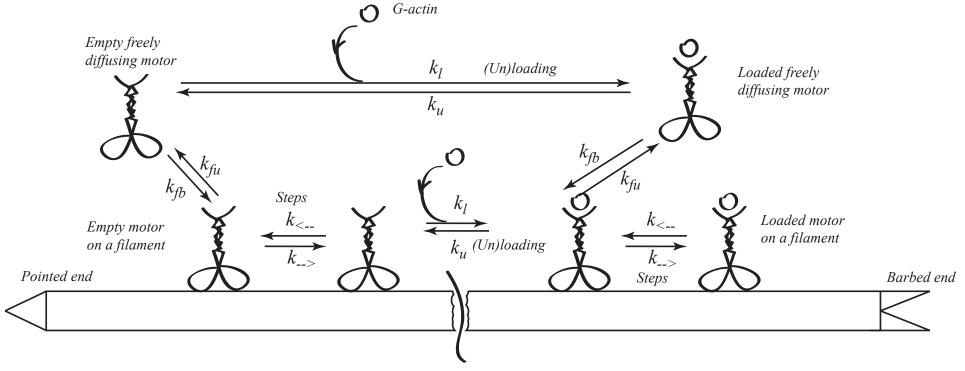
Myosin X

Zhuravlev, Der and Papoian (2010):

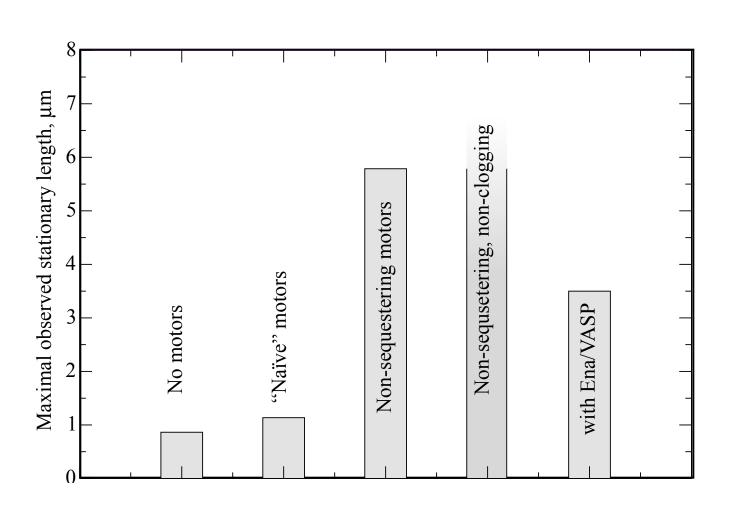
Myosin X transports G-actin?

P. I. Zhuravlev, B. Der and G. A. Papoian, **Biophys. J.**; 2010, 98, 1439–1448





Effect of active transport on filopodial length



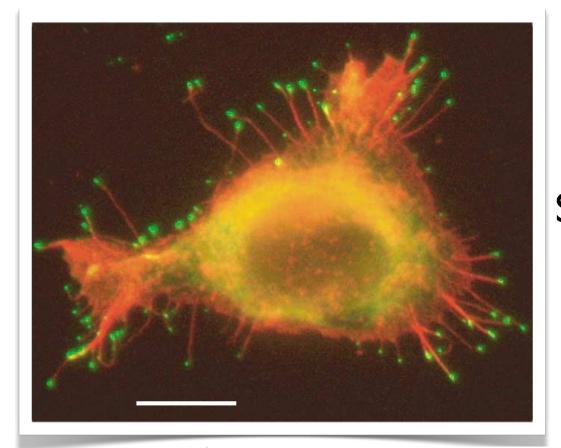
Active Transport Conclusions

- Naïve addition of motors does not work
- There are rules for efficient active transport:
 - Prevent sequestration
 - Clear the "rails"
- Ena/VASP scheme may be a way to achieve these requirements
- Flux balance dramatically affects growth dynamics

Motor concentration profiles

Upper: activ X green.

Lower: myo shining, you concentrati and individu traveling

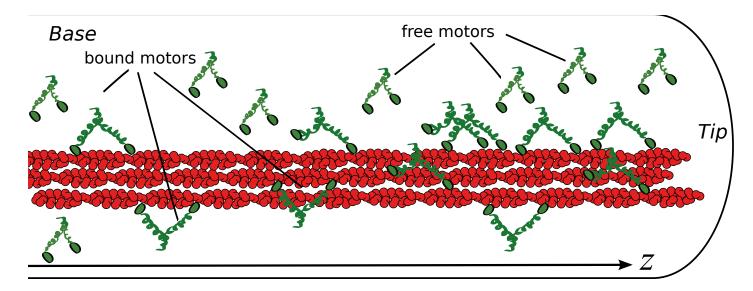


Sousa and Cheney (2005)



Kerber et al. (2009)

Motor Distributions



TASEP + diffusion:

$$\frac{\partial c_{\rm f}}{\partial t} + \frac{\partial J_{\rm f}}{\partial z} = k_{\rm off} c_{\rm b} - k_{\rm on} c_{\rm f}$$
$$\frac{\partial c_{\rm b}}{\partial t} + \frac{\partial J_{\rm b}}{\partial z} = k_{\rm on} c_{\rm f} - k_{\rm off} c_{\rm b}$$

$$J_{\rm f}(z) = -D \frac{\partial c_{\rm f}}{\partial z}$$

Boundary conditions for stationary solution:

$$J_{\rm f}(0) = c_{\rm f}(0), (J_{\rm f} + J_{\rm b}) \Big|_{tip} = 0, J_{\rm b}(0) = 0$$

$$J_{
m f}+J_{
m b}$$
 is an integral



The solution does not depend on the filopodial length!

Master Equation: Neglecting Correlations Between Sites

- Discrete hopping for bound motors
- Continuous diffusion for free motors

$$\dot{b_n} = k_{\to} b_{n-1} + k_{\leftarrow} b_{n+1} - (k_{\to} + k_{\leftarrow}) b_n - k_{\text{off}} b_n + k_{\text{on}} c_f(z_n)$$

$$\left(\frac{\partial c_{\rm f}}{\partial t} = \frac{\partial}{\partial z} \left(D_{\rm m} \frac{\partial c_{\rm f}}{\partial z} \right) + k_{\rm off} c_{\rm b} - k_{\rm on} c_{\rm f}, \right)$$

$$z_{n} = n\varepsilon \qquad v = (k_{\rightarrow} - k_{\leftarrow})/\varepsilon$$

$$b_{n} = b(z_{n}) = c_{b}(z)$$

$$b_{n-1} = c_{b}(z - \varepsilon) = c_{b}(z) - \varepsilon c'_{b}(z) + \dots$$

$$b_{n+1} = c_{b}(z + \varepsilon) = c_{b}(z) + \varepsilon c'_{b}(z) + \dots$$

$$\left(\frac{\partial c_{\rm b}}{\partial t} = -\frac{\partial}{\partial z} \left(v c_{\rm b}\right) - k_{\rm off} c_{\rm b} + k_{\rm on} c_{\rm f}\right)$$

More Accurate Semi-Mean-Field Equation

$$\dot{b_n} = k_{\to} b_{n-1} (1 - b_n) + k_{\leftarrow} b_{n+1} (1 - b_n) - k_{\to} (1 - b_{n+1}) b_n - k_{\leftarrow} (1 - b_{n-1}) b_n - k_{\text{off}} b_n + k_{\text{on}} (1 - b_n) c_f(z_n)$$

$$\dot{b_n} = k_{\to} b_{n-1} + k_{\leftarrow} b_{n+1} - (k_{\to} + k_{\leftarrow}) b_n - b_n (k_{\to} - k_{\leftarrow}) (b_{n-1} - b_{n+1}) - k_{\text{off}} b_n + k_{\text{on}} (1 - b_n) c_f(z_n)$$

$$\frac{\partial c_{\rm b}}{\partial t} = -\frac{\partial}{\partial z} (vc_{\rm b}) + 2vc_{\rm b} \frac{\partial c_{\rm b}}{\partial z} - k_{\rm off} c_{\rm b} + k_{\rm on} (1 - c_{\rm b}) c_{\rm f}$$

$$\frac{\partial c_{\rm b}}{\partial t} = -\frac{\partial}{\partial z} \left(v c_{\rm b} (1 - c_{\rm b}) \right) - k_{\rm off} c_{\rm b} + k_{\rm on} (1 - c_{\rm b}) c_{\rm f}$$

...and then we start neglecting terms

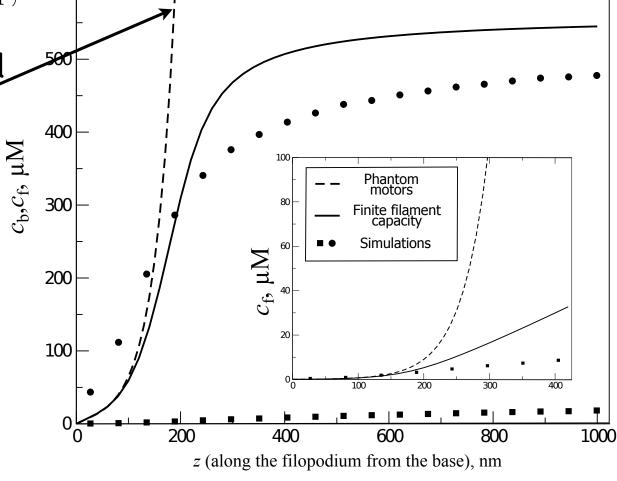
Phantom Motors: Neglecting Next Neighbor Correlations

Naoz et al., Biophys J, 2008 (stereocilia):

$$\begin{cases} -D_{\rm m} \frac{\partial^2 c_{\rm f}}{\partial z^2} = k_{\rm off} c_{\rm b} - k_{\rm on} c_{\rm f}, \\ v \frac{\partial c_{\rm b}}{\partial z} = -k_{\rm off} c_{\rm b} + k_{\rm on} c_{\rm f}, \end{cases}$$

Analytical exponential solution blows up-

near the tube base

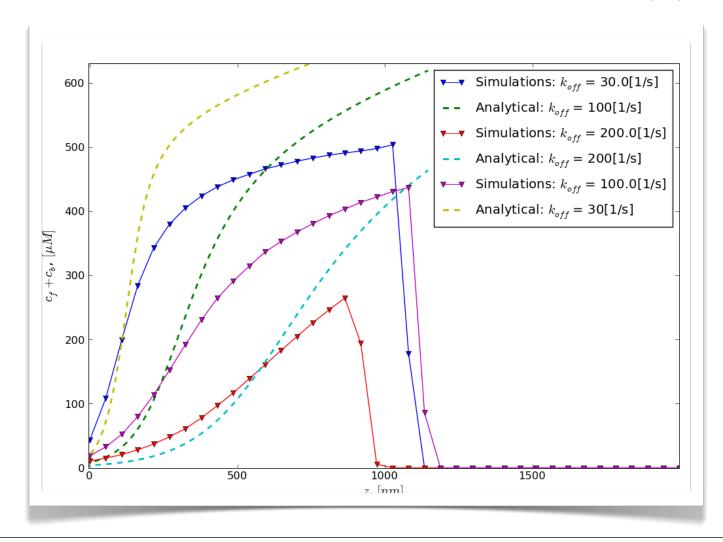


Finite Filament Capacity: Saturation Effect

$$\begin{cases} -D_{\rm m} \frac{\partial^2 c_{\rm f}}{\partial z^2} = k_{\rm off} c_{\rm b} - k_{\rm on} (1 - c_{\rm b}) c_{\rm f}, \\ v \frac{\partial c_{\rm b}}{\partial z} = -k_{\rm off} c_{\rm b} + k_{\rm on} (1 - c_{\rm b}) c_{\rm f}, \end{cases}$$

"on" rate diminishes with the fraction of bound filaments:

$$k_{\rm on}(c_{\rm b}) = k_{\rm on}(1 - c_{\rm b})$$



Jammed Motor Model: Traffic Jam + Saturation Effect

$$\begin{cases} -D_{\rm m} \frac{\partial^2 c_{\rm f}}{\partial z^2} = k_{\rm off} c_{\rm b} - k_{\rm on} (1 - c_{\rm b}) c_{\rm f}, \\ \frac{\partial}{\partial z} \left(v c_{\rm b} (1 - c_{\rm b}) \right) = -k_{\rm off} c_{\rm b} + k_{\rm on} (1 - c_{\rm b}) c_{\rm f}, \end{cases}$$

motor speed slows down due to neighbors

