
Toy Models and Fast Scrambling (I)

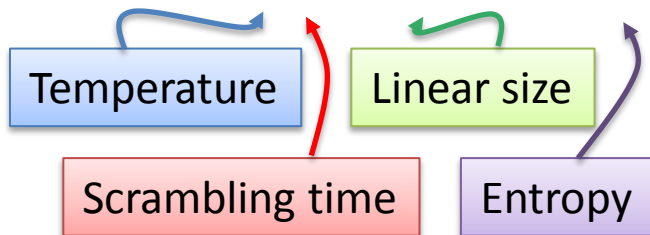
Patrick Hayden with Nima Lashkari, Douglas Stanford,
Tobias Osborne and Matt Hastings



Scrambling

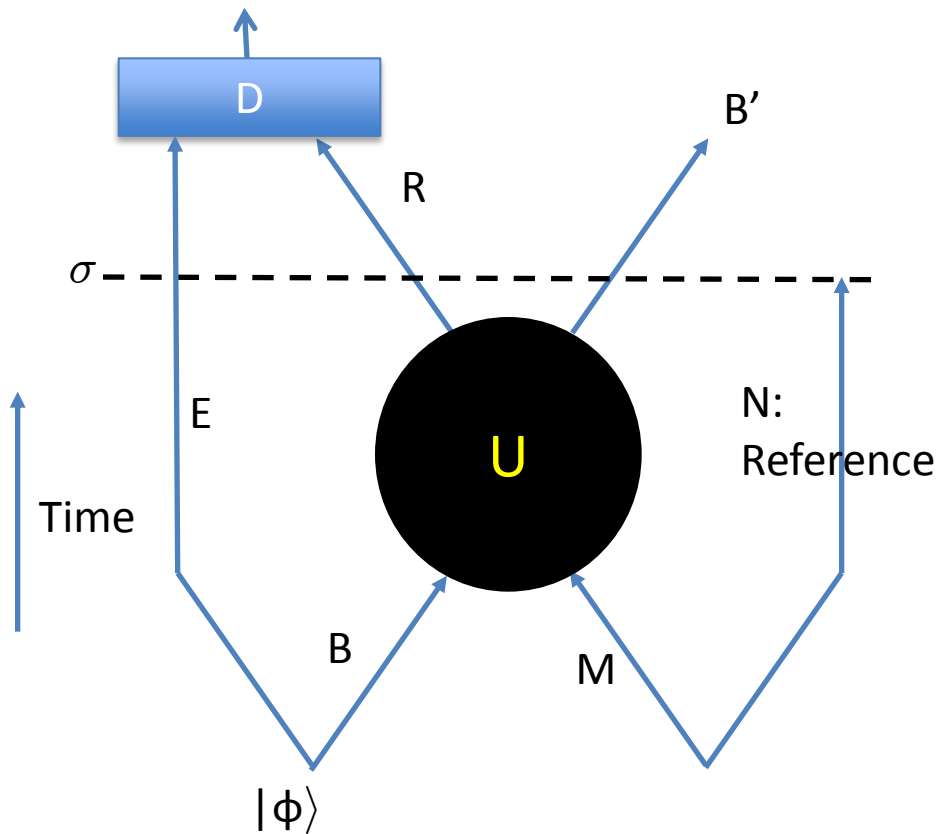
- Minimum time for “localized” information to become inaccessible without measuring fraction $O(1)$ of the whole system
- Normal systems: geometrical locality

$$- \quad Tt^* \sim L^{\text{const}} \sim S^{\text{const}/d}$$



- Schwarzschild black holes
 - $Tt^* \sim \log S$ (estimate based on charge spreading)
- Conjecture: this is correct, and no system can scramble its degrees of freedom faster [Sekino-Susskind'08, Susskind'11]
 - Motivation: black hole complementarity principle

Scrambling and quantum error correction



Sending arbitrary states from M to R is *equivalent* to establishing entanglement between N and R

Establishing entanglement between N and R is *equivalent* to eliminating all correlations between N and B'

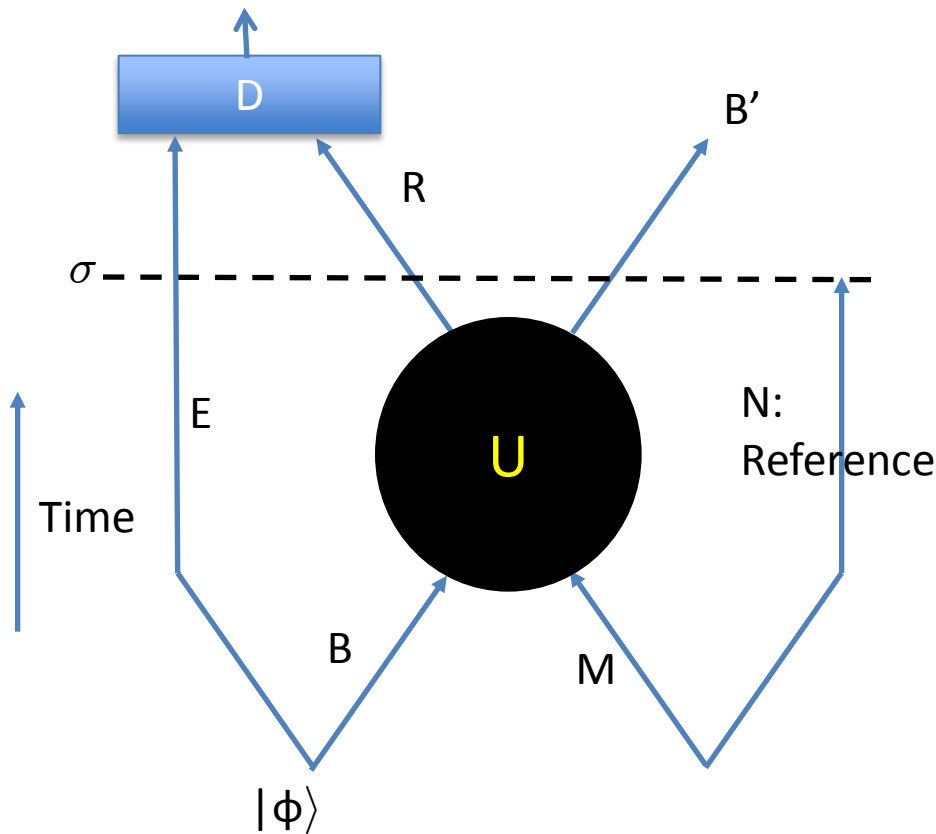
$$\text{Tr}_R \sigma_{NB'R} = \text{Tr}_N \xi_{B'}$$

$$\Rightarrow$$

$$|\sigma_{NB'R}\rangle = (\text{id}_{NB'} \otimes U_R) |\phi_{NR1}\rangle |\psi_{B'R2}\rangle$$

If U scrambles systems of size $|B'|$, then the message M can be decoded from R. (No-cloning requires $|R| > |B'|$.)

Scrambling and quantum error correction



Sending arbitrary states from M to ER is *equivalent* to establishing entanglement between N and ER

Establishing entanglement between N and R is *equivalent* to eliminating all correlations between N and B'

$$\text{Tr}_{ER} \sigma_{NB'R} = T_N \xi_{B'}$$

\Rightarrow

$$|\sigma_{NB'ER}\rangle = (\text{id}_{NB'} \otimes U_{ER}) |\phi_{NR1}\rangle |\psi_{B'R2}\rangle$$

Scrambling time controls information release. Faster than $\log S$ leads to problems for black hole complementarity.

If U scrambles systems of size $|B'|$, then the message M can be decoded from ER. (No cloning requires $|R| > |B'|$.)

Big picture versus toy examples

String theory descriptions of black holes couple degrees of freedom nonlocally.

e.g. BFSS Matrix theory:
$$L = \sum_a \text{tr} \dot{M}^a \dot{M}^a - \sum_{ab} \text{tr} [M^a, M^b]^2$$

Every pair of matrix entries appears together in at least one term

Would like to show by direct analysis of the system that it is a fast scrambler

HARD



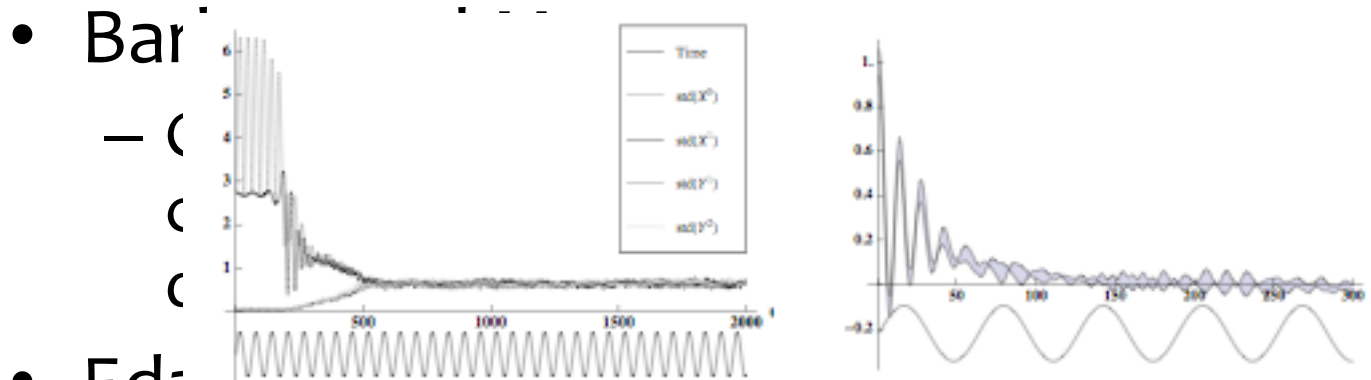
Goal of this hour is more modest:

- 1) Find examples of toy systems that scramble quickly
Want time independent, 2-body interactions, unengineered
Should scramble a whole subspace of initial states
- 2) Prove general lower bounds on scrambling times



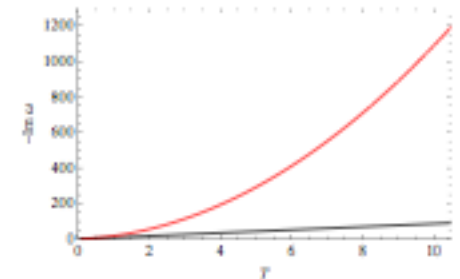
Related work

- Asplund, Berenstein, Trancanelli
 - Numerical simulation of BMN matrix model (classical)



hyperbolic
:

- Edar ...
 - Use AdS/CFT to study thermalization in strongly coupled noncommutative gauge theories

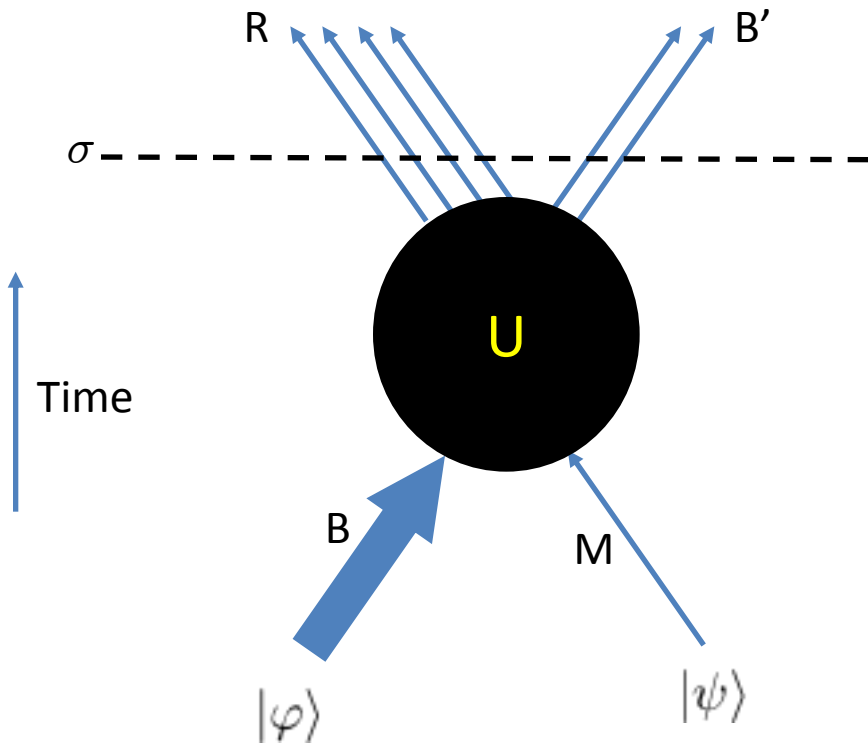


Outline

- Part I
 - Scrambling and quantum error correction ✓
 - Definitions and calibration
 - Brownian quantum circuits
- Part II (Douglas Stanford)
 - Ising interaction on random graphs
 - Lieb-Robinson bounds for nonlocal interactions
 - Comments on AdS/CFT

Some formality

Scrambling n subsystems:
any $n/3$ should be independent of ψ



There should exist a single ψ_0 such that for all valid ψ

$$\| \text{tr}_R \sigma(\psi) - \text{tr}_R \sigma(\psi_0) \|_1 < \epsilon$$

Schatten- l_1 norm measures statistical distinguishability

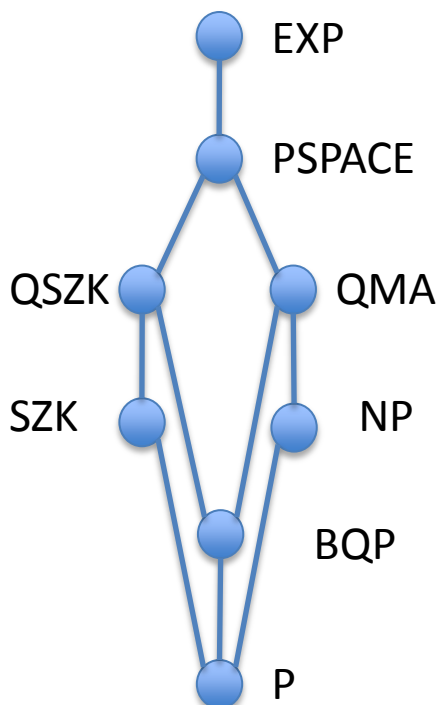
In our toy models, we will simply compare $\text{tr}_R \sigma(\psi)$ to the unique maximum entropy state on B' .

The computer scientist's cop-out

How hard is it to determine if an efficiently specified $U(t)$ scrambles a specific B' ?

Good news! (maybe)

Graph isomorphism



Entropy of time-averaged quantum state

Gnd state energy of 2-body quantum spin systems

Gnd state energy of classical spin glass

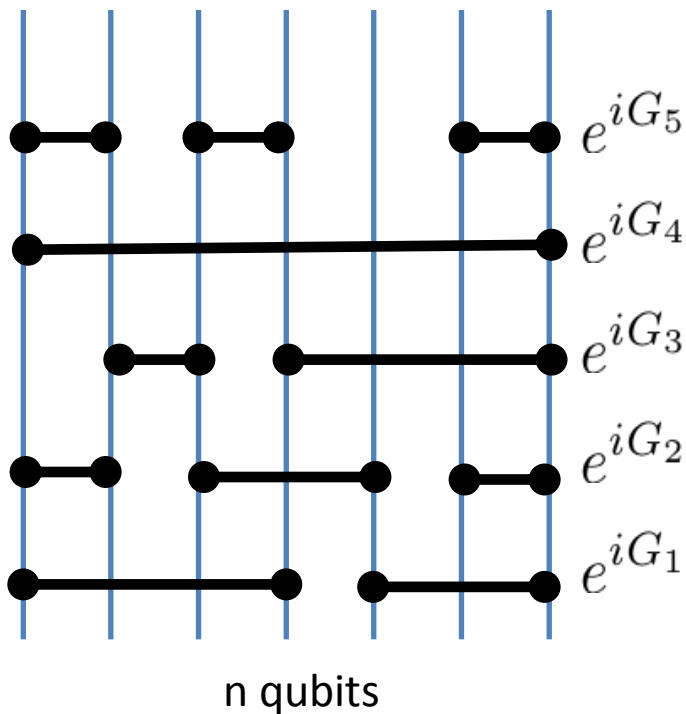
Scattering amplitudes in ϕ^4 theory

Gnd state energy of 1-d Ising model, arb couplings

Determining whether a noisy quantum evolution is correctable is QSZK-complete

[with Brian Swingle]

Brownian circuits



G_i random pairwise interaction

$$G_i = \sum_{\langle j,k \rangle} \sum_{\alpha_j, \alpha_k} \sigma_{\alpha_j}^{(j)} \otimes \sigma_{\alpha_k}^{(k)} g_{ijk\alpha_j\alpha_k}$$

Location \rightarrow $\sigma_{\alpha_j}^{(j)}$ and $\sigma_{\alpha_k}^{(k)}$

Operator \rightarrow $g_{ijk\alpha_j\alpha_k}$

i.i.d. Gaussian $N(0, \varepsilon)$ \rightarrow $g_{ijk\alpha_j\alpha_k}$

Limit of infinitesimal ε :

$$U(t) = \sum_{\langle j,k \rangle} \sum_{\alpha_j, \alpha_k} \sigma_{\alpha_j}^{(j)} \otimes \sigma_{\alpha_k}^{(k)} U(t) dW_{jk\alpha_j\alpha_k}(t) - \frac{1}{2} U(t) dt$$

Weiner process \rightarrow $dW_{jk\alpha_j\alpha_k}(t)$



Dankert et al.: Construction of circuit scrambling in time $O(\log n)$

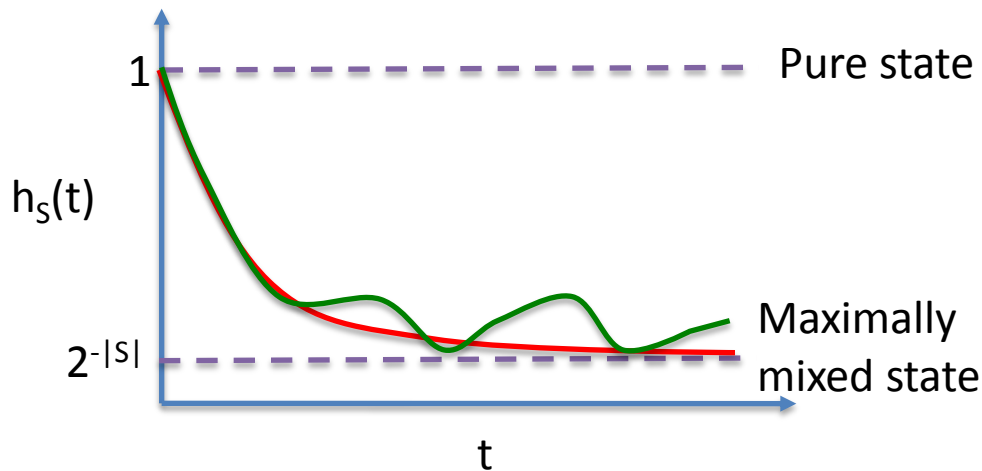
$$j=1$$

$$\sigma_0 = I \quad \sigma_1 = \sigma_x \quad \sigma_2 = \sigma_y \quad \sigma_3 = \sigma_z$$

Subsystem entropies

State $\Psi(t)$. Density operator for S subset of $\{1,2,\dots,n\}$: $\Psi_S(t) = \text{tr}_{[n]\setminus S} \Psi(t)$.

Interested in *purity* $h_S(t) = \text{tr} \Psi_S(t)^2$



Smooth out fluctuations by averaging over trajectories: $\langle h_S(t) \rangle$

Simplify by choosing product pure input state $|\psi(0)\rangle = |\psi_1\rangle |\psi_2\rangle \dots |\psi_n\rangle$

Gives $\langle h_S(t) \rangle = \langle h_{|S|}(t) \rangle = \langle h_k(t) \rangle$

Small miracle: system of linear ODE closes and is (almost) solvable

$$\frac{d\langle h_k \rangle}{dt} = k(n-k) \left[2\langle h_{k-1} \rangle - 5\langle h_k \rangle + 2\langle h_{k+1} \rangle \right]$$

Analysis of ODE

$$|\psi(0)\rangle = |\psi_1\rangle|\psi_2\rangle\dots|\psi_n\rangle$$

$$\text{Purity } h_k(t) = \text{tr } \psi_{|S|=k}(t)^2$$

$$\frac{d\langle h_k \rangle}{dt} = k(n-k) \left[2\langle h_{k-1} \rangle - 5\langle h_k \rangle + 2\langle h_{k+1} \rangle \right]$$

Quick and dirty analysis

Let t_k be time at which $\langle h_k(t) \rangle = (1+\delta)2^{-k}$

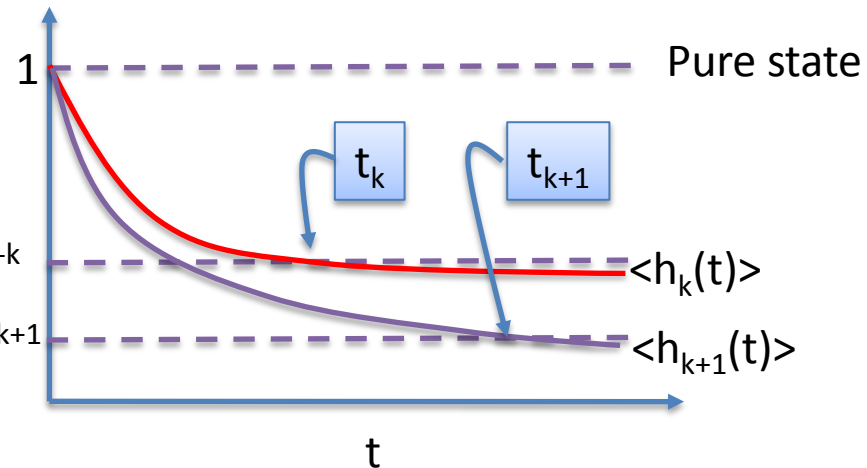
For $t > t_{k-1}$

$$\begin{aligned} \frac{d\langle h_k \rangle}{dt} &\sim \leq kn \left[2 \frac{1+\delta}{2^{k-1}} - 5\langle h_k \rangle + 2\langle h_{k+1} \rangle \right] \\ &\leq kn \left[2 \frac{1+\delta}{2^{k-1}} - 3\langle h_k \rangle \right] \end{aligned}$$

Exponential decay with rate proportional to k .

So $t_k - t_{k-1} \leq O(1/k)$

$$t_k \sim \sum_{j=1}^k \frac{1}{j} \sim \log(k)$$



Careful analysis

Solve using Gauss hypergeometric functions

$$\langle h_k(t) \rangle \sim \sum_{j=1}^n \alpha_{jk} e^{-3jt} n 2^{-n-1+2j} {}_2F_1 \left(n+1, 1-m; 2; \frac{3}{4} \right)$$

Brownian circuits: take-home

- Scramble very effectively: subsystems of size smaller than half become almost maximally mixed
- Scramble quickly: $t^*/t_1 = O(\log n)$
- But:
 - Time-dependent
 - Not very physical
 - Lots of randomness