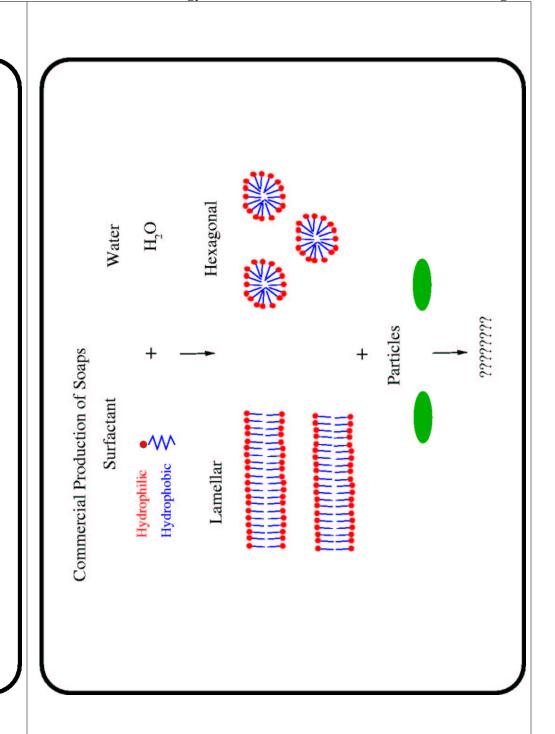
STRESS TRANSMISSION BY A DEFECT NETWORK IN A LIQUID CRYSTALLINE PHASE

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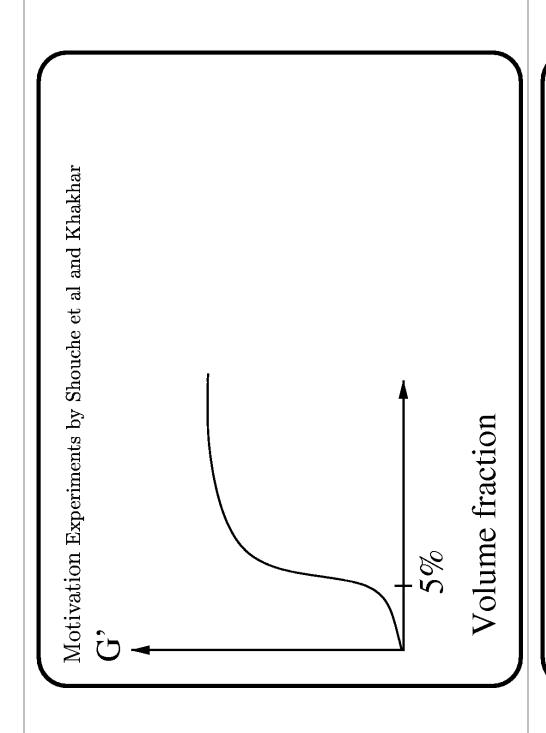


Manufacture of Personal Wash Bars

- Raise yield stress and shear modulus (to provide rigidity to the bar) with as little solids content as possible.
- Develop constitutive models for the rheology of particle laden liquid crystalline phases for use in the process design.
- Presently used rheological models for soft solids extremely inadequate for reliable, predictive equipment design.

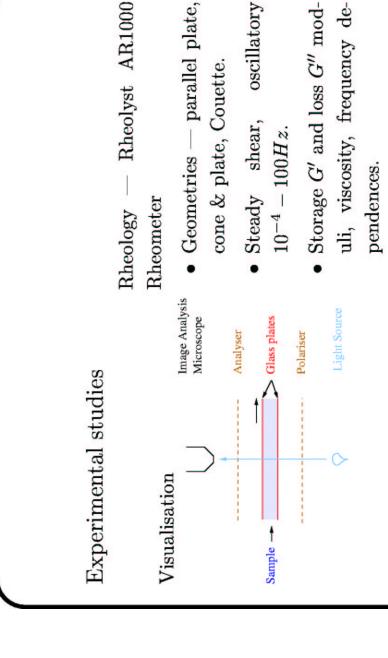
Objectives

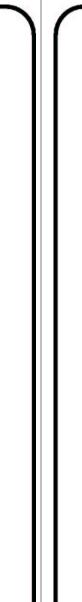
- Develop a microscopic understanding of the relation between the structure and rheology of particle laden liquid crystalline phases.
- macroscopic rheology which can be used for design of materials with appropriate phase composition and particle concentration Develop the links between microscopic structure and to provide required macroscopic properties. ö
- models for liquid crystalline phases with suspended particles Incorporate this understanding into macroscopic rheological which can be used for flow computations. ന :

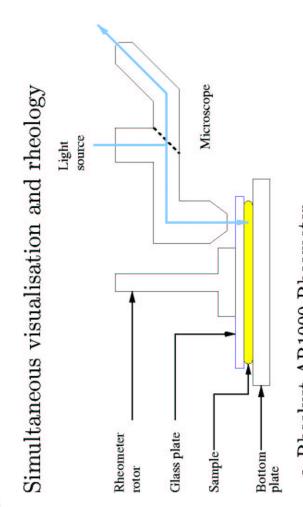


Outline of presentation

- Rheological measurements of lamellar phases with and without particles.
- Visualisation of the structures formed during shear of the lamellar phase. જાં
- 3. Effect of added particles.
- Estimation of the elasticity from the structures formed. 4.
- Dynamical evolution of structures under shear. ъ.
- 6. Conclusions.







- Rheolyst AR1000 Rheometer.
- Top plate modified by sticking a glass plate on the metal disk.
- Bottom reflecting plate.
- Custom built microscope.

Rheological Measurements

System

Sodium dodecyl ether sulphate (SLES) + Water

$$CH_3 - (CH_2)_{10} - [OC_2H_4] - SO_4^-Na^+$$

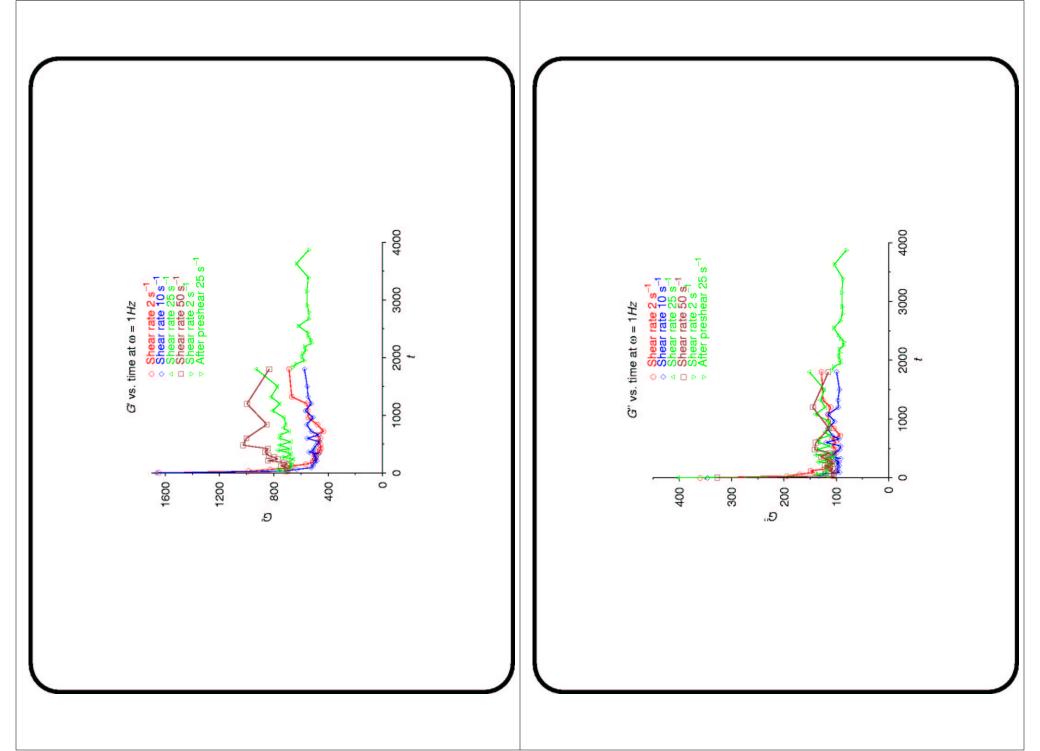
• Phases (% of SLES)

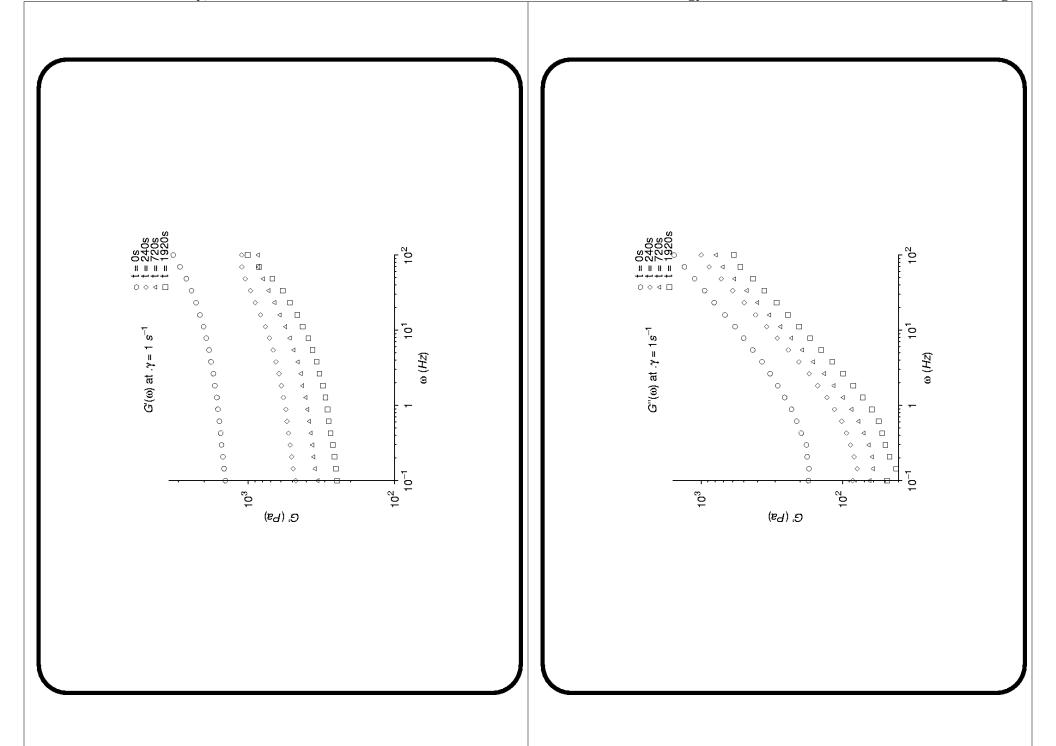
$$I \sim 37
ightarrow H \sim 58
ightarrow H + L \sim 63
ightarrow L \sim 81
ightarrow L + K$$

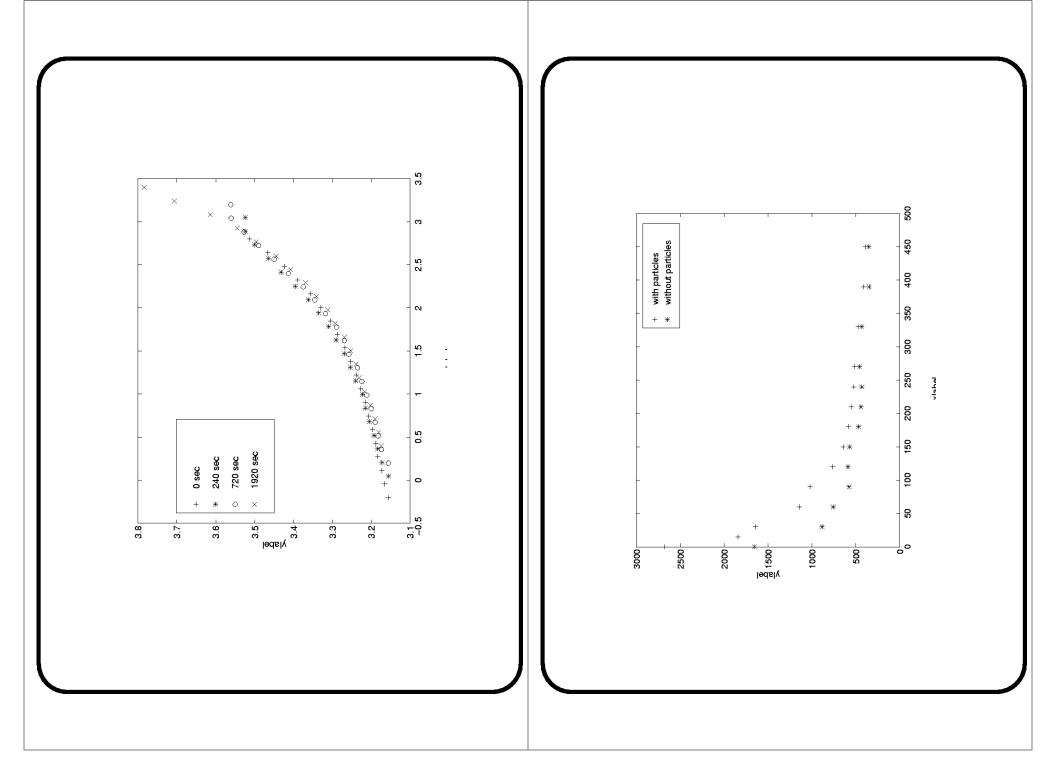
- Operating concentration SLES: $H_2O = 72.3:26.8$
- Dispersed particles: 19 μ silica particles, 9.5 μ polystyrene particles.

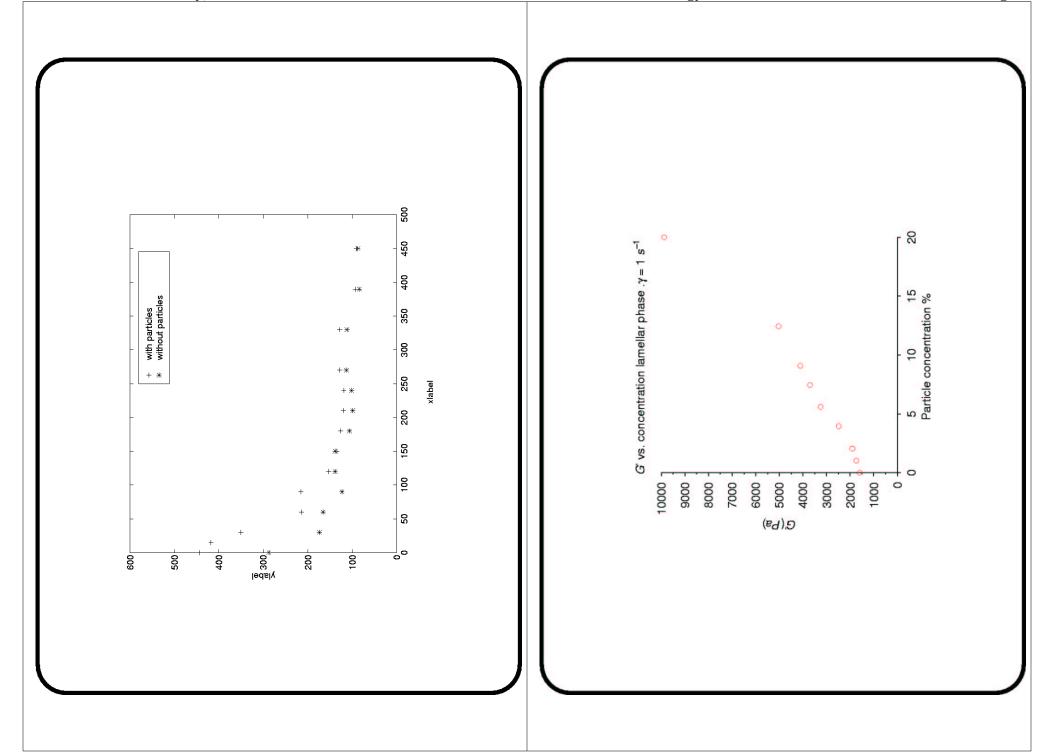
Protocol for rheology experiments

- Place sample in the parallel plate geometry of the rheometer (in disordered state.
- Steady shear for time interval Δt .
- Stop steady shear and determine storage and loss modulus using oscillatory measurements in the rheometer.
- Steady shear for next interval Δt followed by oscillatory measurements.
- Plot storage and loss moduli versus total time of shear.



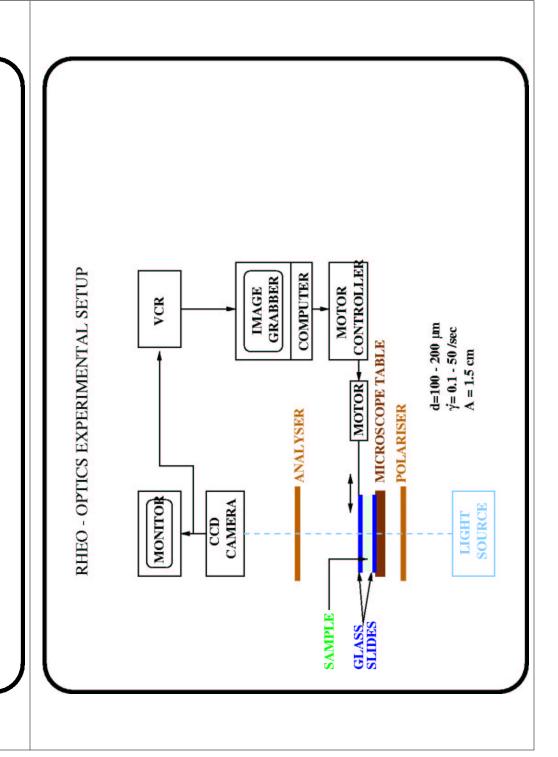


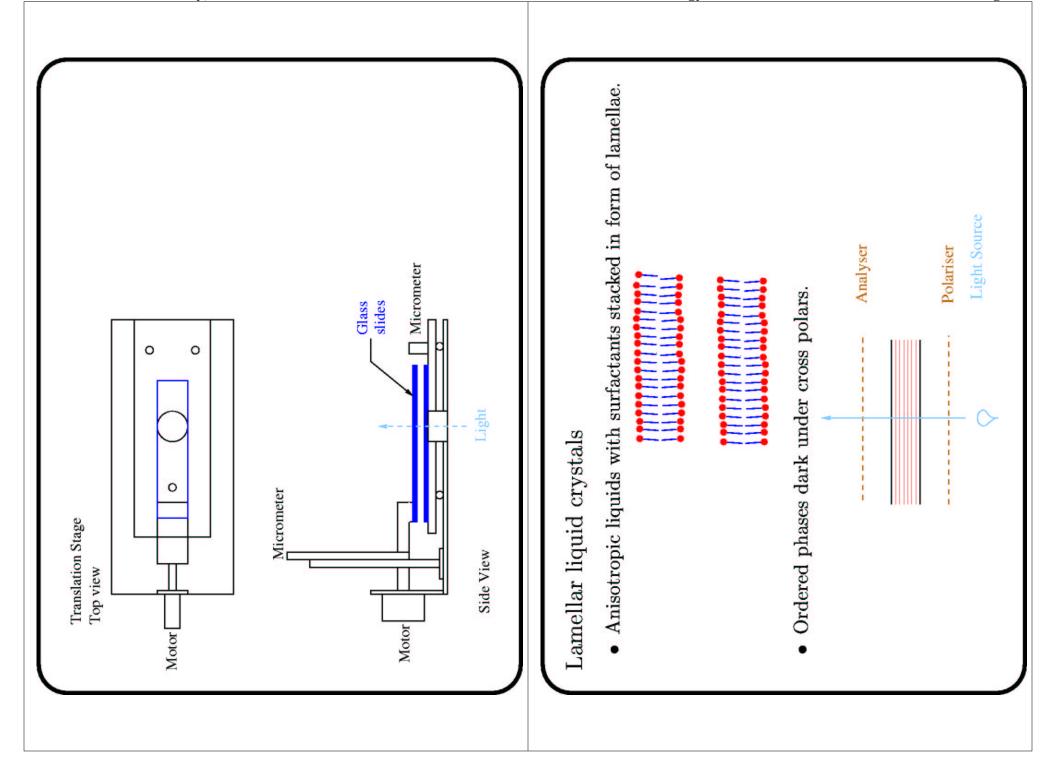


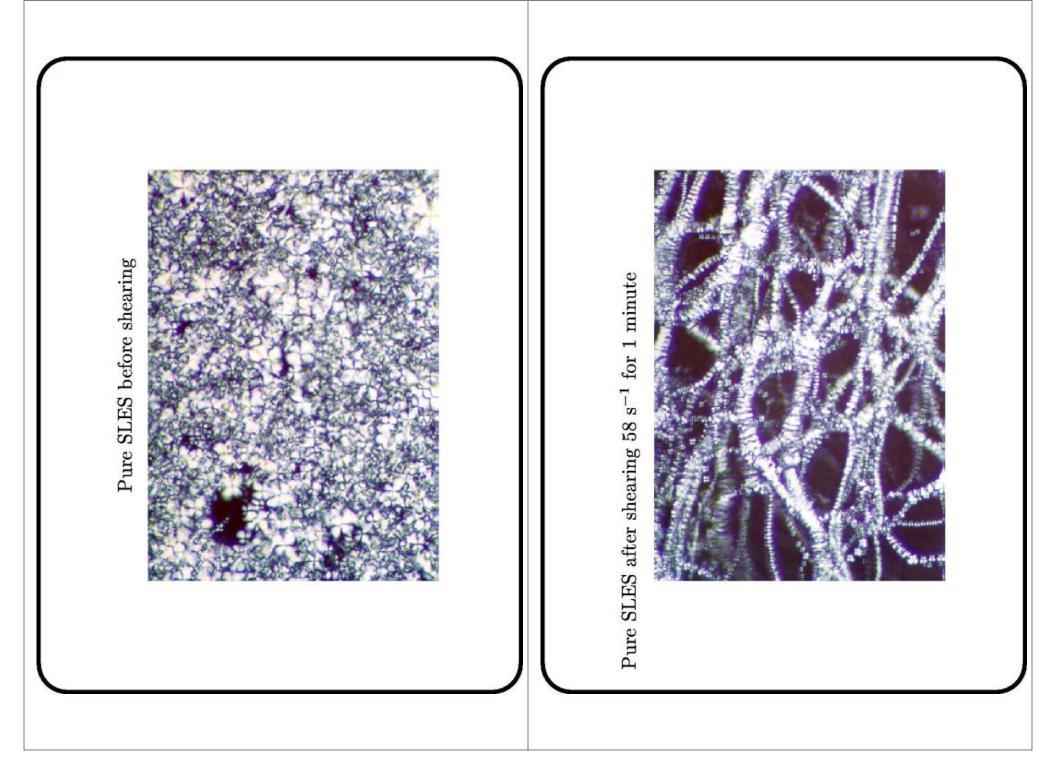


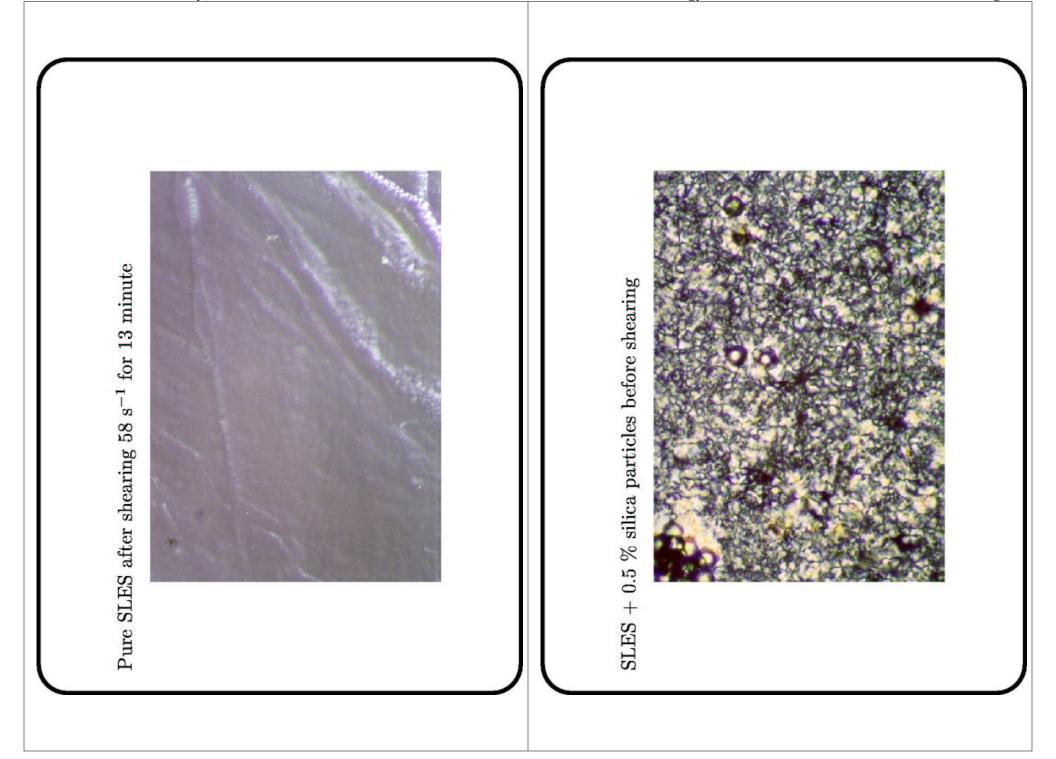
Conclusions (Rheological Measurements)

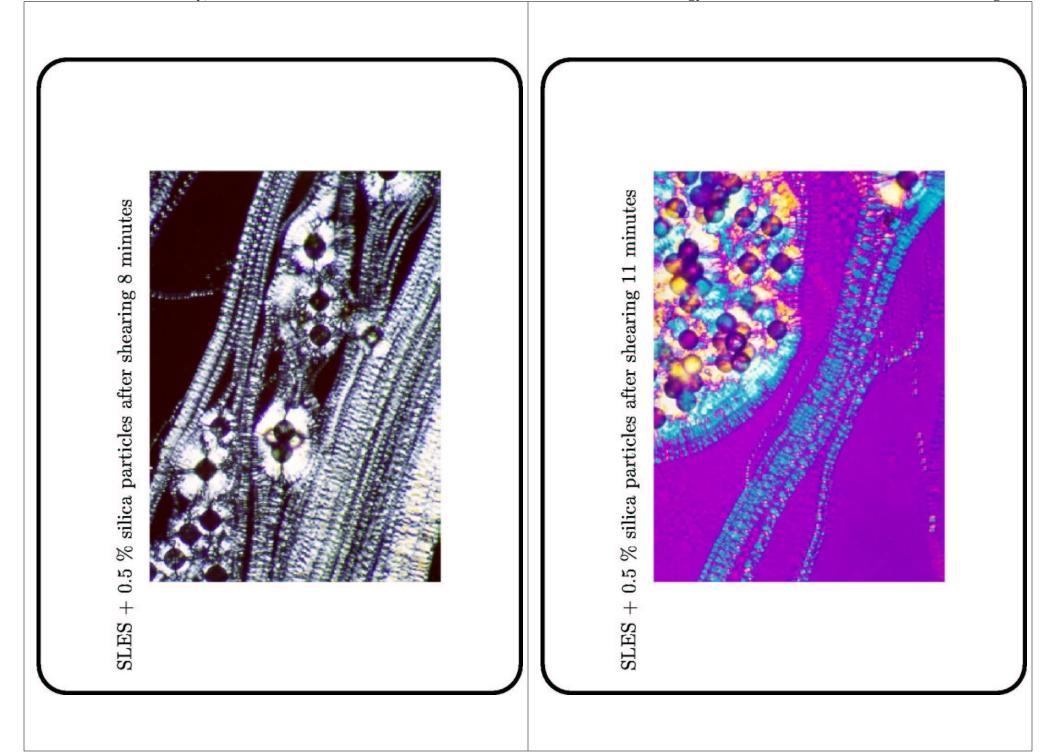
- Storage and loss moduli decrease as system is sheared, reach what appears to be a steady state value, and then appear to increase on further shearing.
- At any point in time, frequency dependence of moduli are similar to those seen in typical soft solids. a
- Transient dynamical properties depend on history of shear. က
- Shear treatment affects microstructure, which in turn affects properties. Working in of defects at late times? 4.
- but qualitative decrease of moduli with Presence of particles increases storage and loss moduli at early shear treatment maintained and intermediate times, ъ.

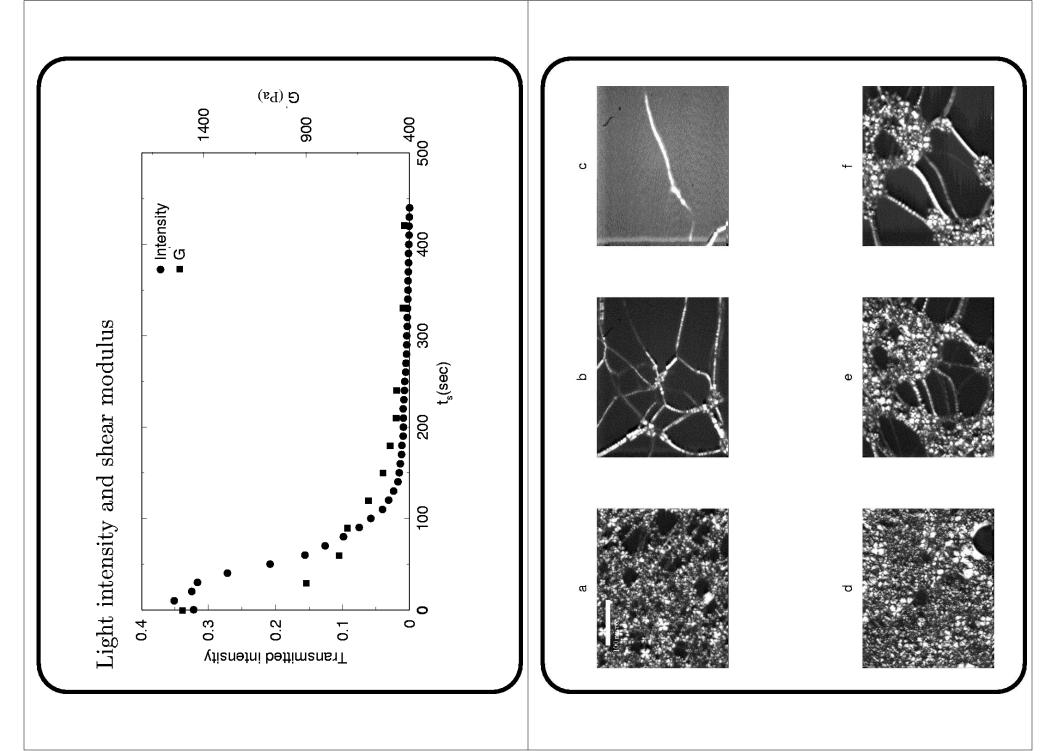


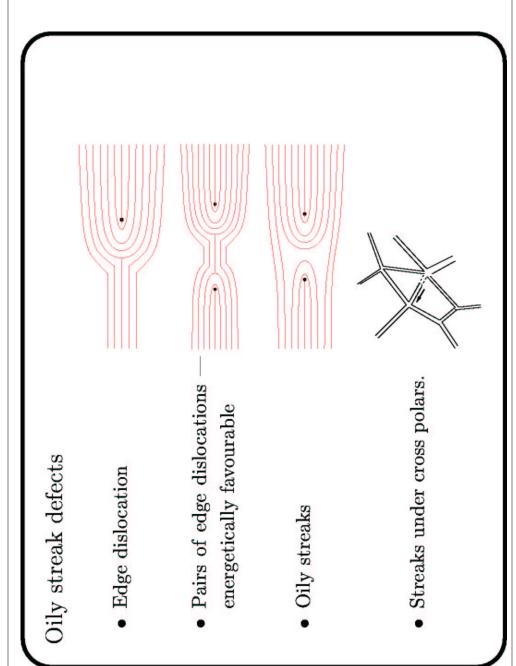


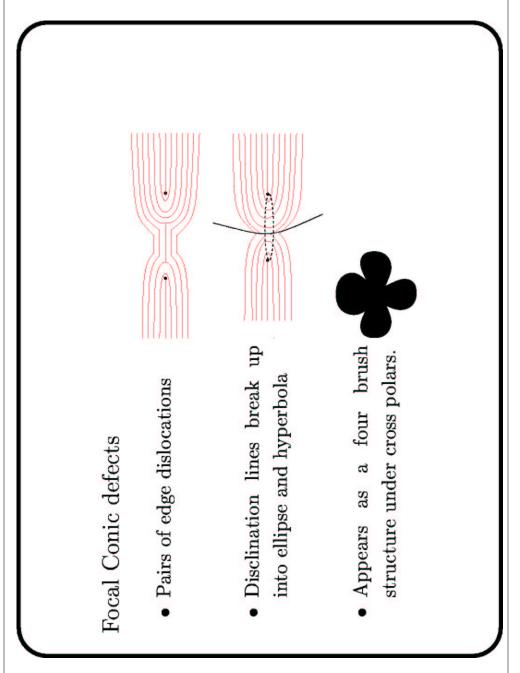


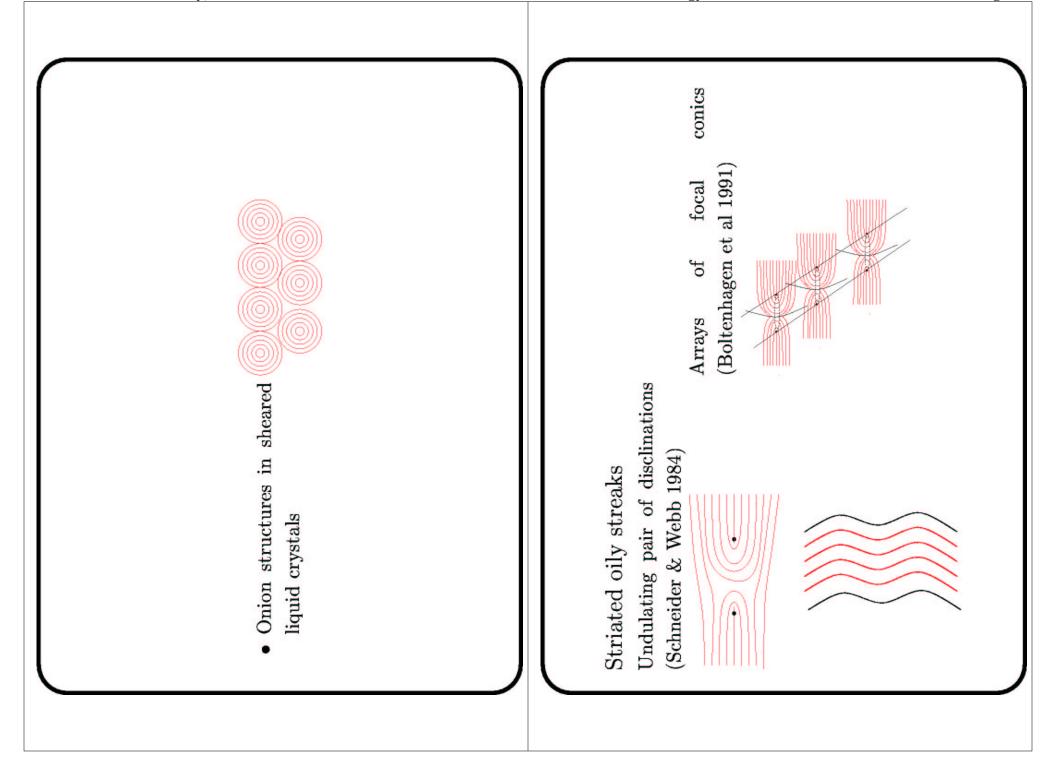




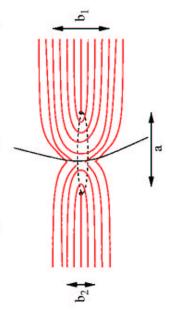








Energetics of defects (Boltenhagen et al 1991)



- Ellipse with major and minor axes $a = (b_1 + b_2)/2$ and
 - $b = \sqrt{b_1 b_2}.$
- Eccentricity $e = (b_1 b_2)/(b_1 + b_2)$.

Energetics of defects (Boltenhagen et al 1991)

• Bending energy

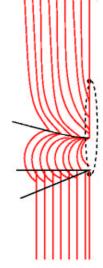
$$f_B = \frac{1}{2}K\left(\frac{1}{R_1} + \frac{1}{R_2}\right)^2 + \frac{\bar{K}}{R_1R_2}$$

Bending energy of edge dislocations (per length) of major axis a, eccentricity e

$$F_B = 4\pi K(1 - e^2)a \left[\log (a/r_c) + 1 - 2\frac{\bar{K}}{K} \right]$$

Bending energy of edge dislocations of length l is $\sqrt{KB}bl$, where b is the Burgers vector.

Energetics of defects



Compression energy of focal conics

$$E_c = Bh^2 \left[\frac{e}{\sqrt{1 - e^2}} - \frac{a}{h} \right]^5$$

- Network, compression screened at length l.
- proportional to e^5 as eccentricity $e = b/(b_1 + b_2)$ increases Combining the two, we get a compression energy increase
- Stabilises oily streaks at higher values of Burgers vector b.
- For $B \sim 10^5 Pa$, $l \sim 50 \mu m$, the minimum b for stable oily streaks is $O(10\mu m)$.

Stress transmission by oily streak network

- ullet Network of oily streaks with mesh size l and tension Γ has shear modulus $G' = (\Gamma/l^2)$.
- Estimation of line tension Γ .
- Compression energy in an oily streak in a sample of width h

$$E = Bh^2 \left[\frac{e}{\sqrt{1 - e^2}} - \frac{a}{h} \right]^5$$

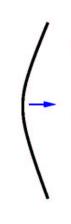
- Network, compression screened at length l.
- For $B \sim 10^5 Pa$, $l \sim 50 \mu m$ and $a \sim 5 10 \mu m$, $G' \sim 10^2 - 10^3 Pa$.
- Correct order of magnitude.

Rheological measurements

- function of the applied strain rate, and not just of the material Shear modulus curve suggests that the rheological state is a properties.
- If stress is transmitted by defects, defect density depends on strain rate.
- It is known that shear anneals defects, but shear should also work in defects.
- Defect density higher at higher strain rates.
- Mechanism?

Undulation instability in lamellar liquid crystals





- Scalar field u describes displacement of layers from their equilibrium position.
- Expansion or contraction perpendicular to the layers

 $S = \left(\partial_z u - rac{1}{2} (
abla_\perp u)^2
ight)$

Bending of layers, the curvature is given by

$$\kappa =
abla_{\perp}^2 u$$

Undulation instability in lamellar liquid crystals

Free energy
$$F = \int dV B \left[\left(\partial_z u - \frac{1}{2} (\nabla_\perp u)^2 \right)^2 + \frac{K}{2} (\nabla_\perp^2 u)^2 \right]$$

Dilation $\partial_z u = D$

$$F = \int dV \left[-\frac{B}{2} D(\nabla_{\perp} u)^2 + \frac{K}{2} (\nabla_{\perp}^2 u)^2 \right]$$

Unstable if

$$k_{\perp} = \sqrt{rac{BD}{K}}$$

Undulation instability in lamellar liquid crystals

Upper cut off of instability based on system size:

- ullet A perturbation of length L_{\perp} along the layers propagates distance (L_{\perp}^2/λ) perpendicular to the layers.
- For a thickness h perpendicular to the layers, the cut off for waves along the layers is $L_{\perp} = \sqrt{h\lambda}$.
- Cut off for wave vector $k_{\perp} = \sqrt{BD/K} > 1/\sqrt{h\lambda}$
- Minimum dilation $D > (K/B\lambda h)$.

Undulation instability in lamellar liquid crystals

- Undulation caused by a normal stress due non uniformity in the boundaries as the system is sheared.
- Normal stress $\propto \eta \dot{\gamma}$ (viscosity \times strain rate).

$$D = \frac{a\eta\dot{\gamma}}{B}$$

• Unstable if $\dot{\gamma} > \dot{\gamma}_c$, where

$$\dot{\gamma_c} \sim rac{K}{\eta \lambda h}$$

Using numerical values used earlier, we find $\dot{\gamma}_c$ is in the range

Conclusions

- network of oily 1. Evolution and annealing of defect structure – streaks under shear.
- striated oily streaks and lines Two different types of defects of focal conic defects. જાં
- Rheology and visualisation \rightarrow stress transmitted due to defect lines, not background lamellar phase. က
- Further shearing works in defects, possibly due to working in of defects by the undulation instability. 4
- Particles move to nodes of defect structure, decrease rate of coarsening of defects. ъ.

Requirement of simulations

- Should predict anisotropic properties.
- Alignment under shear.
- Variation of storatge and loss moduli under shear.
- Annealing and introduction of defects under shear.
- Flow around particles.
- Flow through complex geometries.

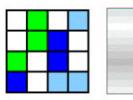
Simulation and Analysis

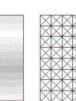


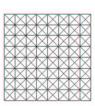






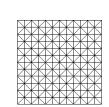


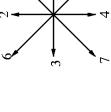




Lattice Boltzmann simulations

• Define 'distribution function' $f_i(\mathbf{x},t)$ for discrete velocity directions on a lattice.





Relation to macroscopic properties

$$\rho = \sum_i f_i$$

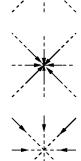
$$u_{\alpha} = \sum_{i} f_{i} e_{i\alpha}$$

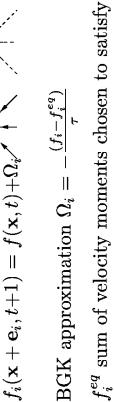
$$ho u_{lpha} u_{eta} - T_{lphaeta} = \sum_i f_i e_{ilpha} e_{ieta}$$

Streaming, collision and re-

distribution steps.







$$\sum_i f_i^{eq} = \rho$$

$$\sum_i f_i^{eq} e_{i\alpha} = \rho u_\alpha$$

$$\sum_{i} f_{i}^{eq} e_{i\alpha} e_{i\beta} = \rho u_{\alpha} u_{\beta} + \rho C_{s}^{2} \delta_{\alpha\beta}$$

Can show above scheme correctly reproduces macroscopic Stokes equations with kinematic viscosity Navier -

 $0.5)C_s^2$.

Lattice Boltzmann for lamellar phase

Free energy functional $F[\psi]$

$$F[\psi] = \int dV \left[\frac{r}{2} \psi^2 + \frac{u}{4} \psi^4 + c[(\nabla^2 + q_0^2) \psi]^2 \right]$$
• Two distributions $f(\mathbf{x},t)$ for total density and $g(\mathbf{x},t)$ for order

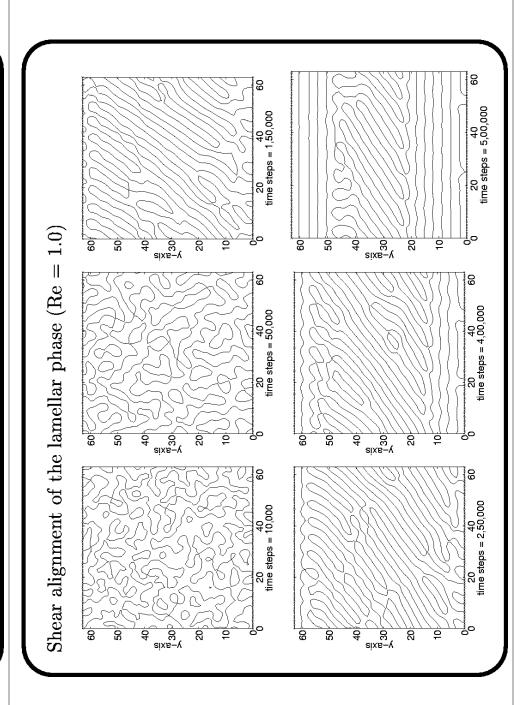
parameter.

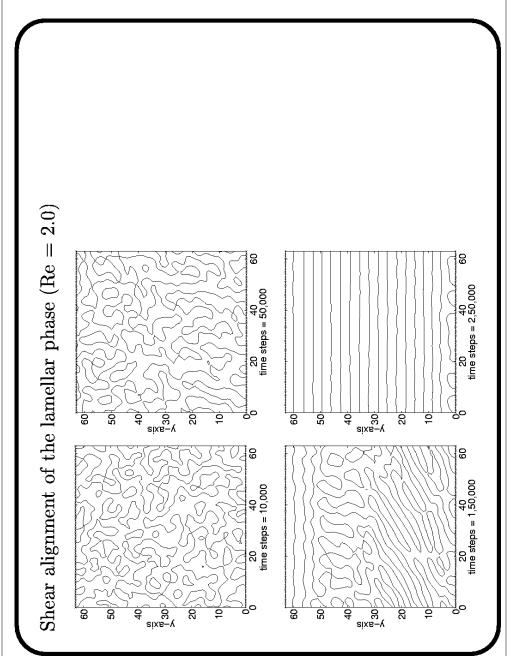
 $\sum g_i^{eq} e_{i\alpha} = \psi u_\alpha$

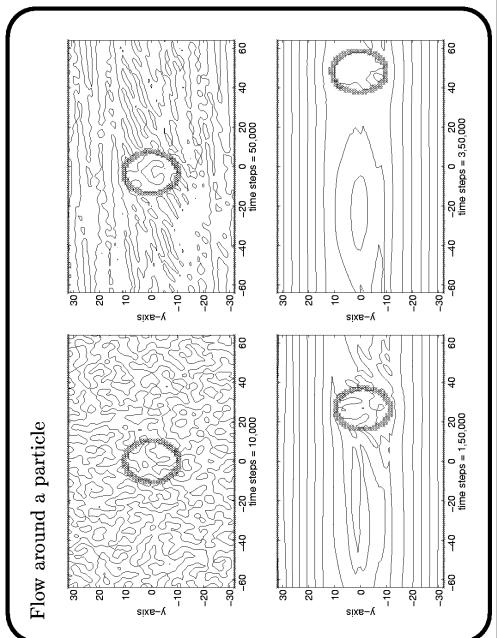
Second moment relations

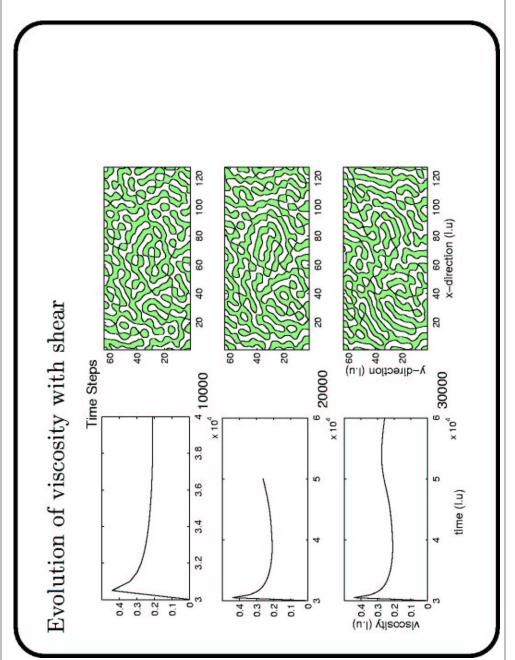
$$\sum f_i^{eq} e_{i\alpha} e_{i\beta} = \rho u_\alpha u_\beta + \mathcal{F}[f(\psi)]$$

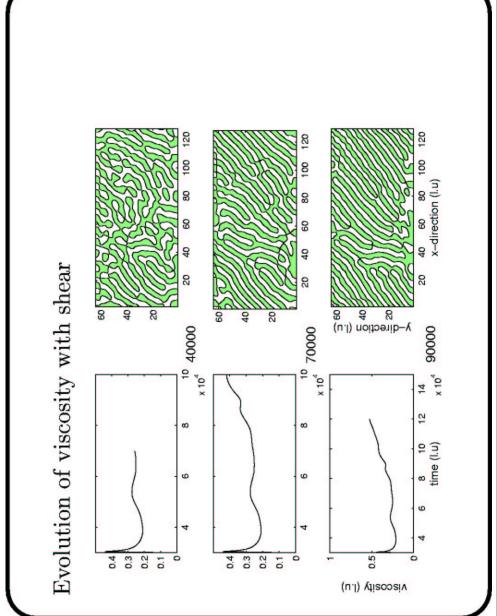
$$\sum_{i}g_{i}^{eq}e_{ilpha}e_{ieta}=\psi u_{lpha}u_{eta}+\mathcal{G}[f(\psi)]$$

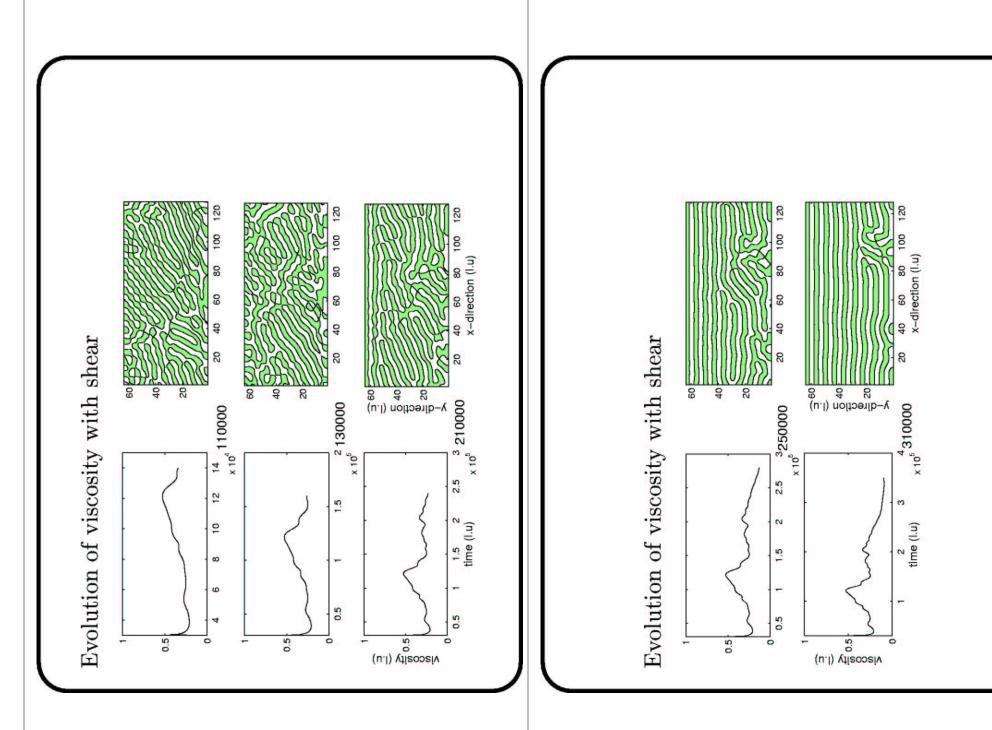


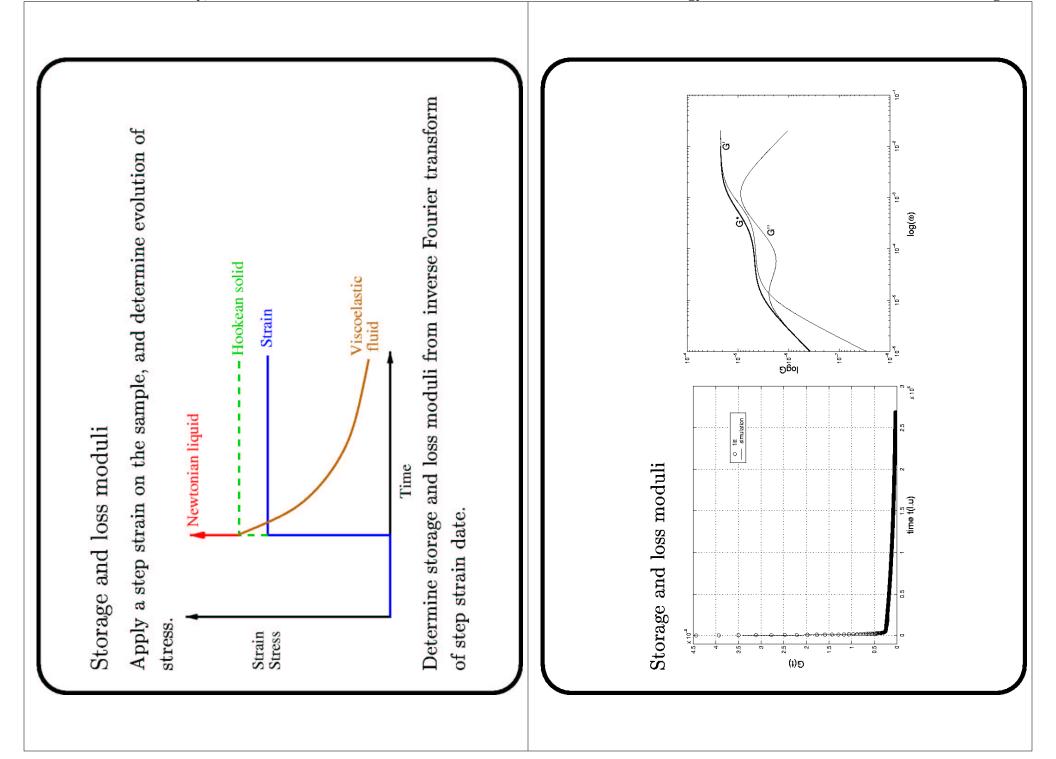


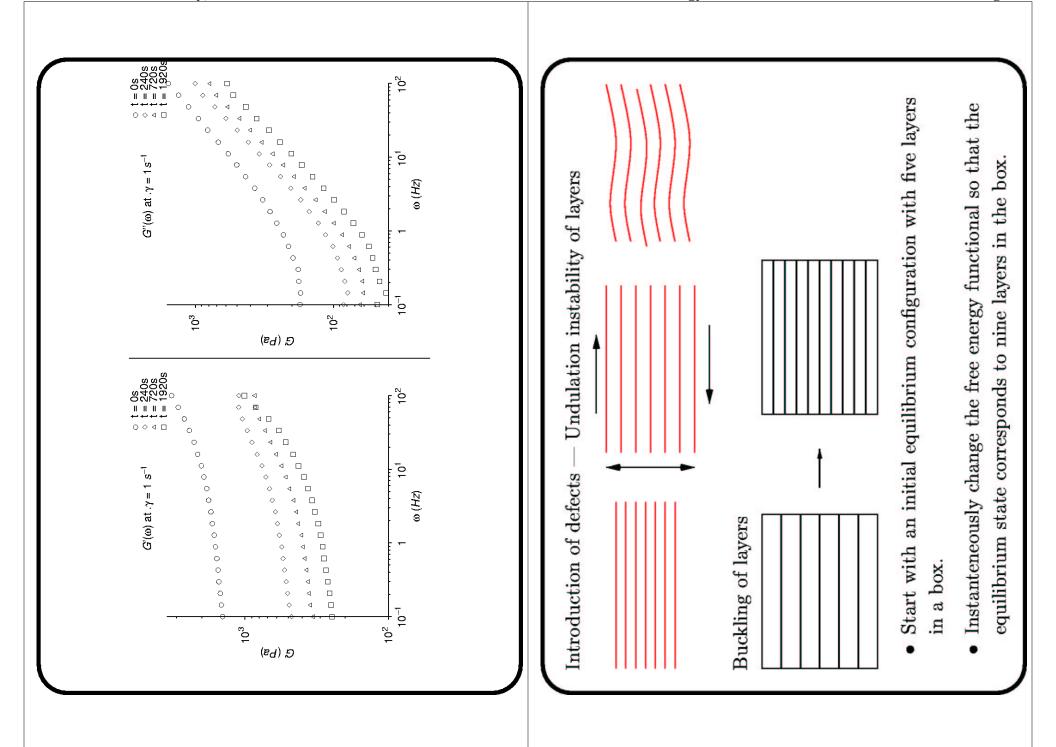


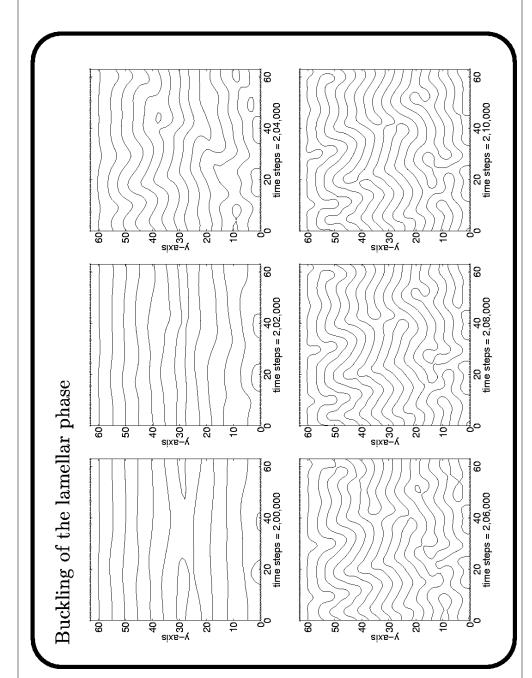


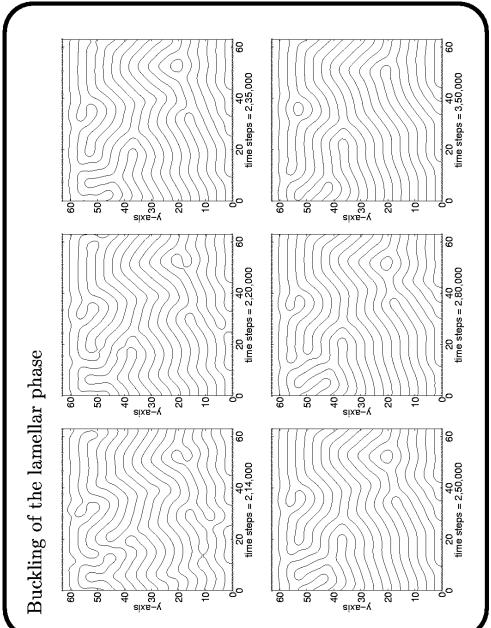




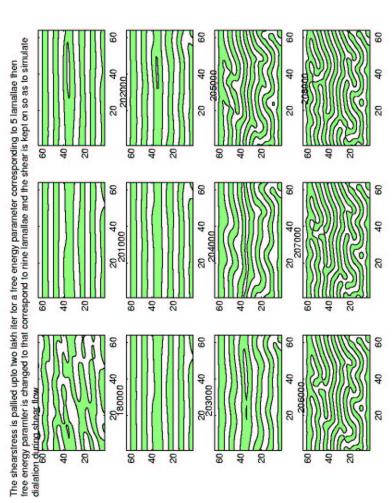


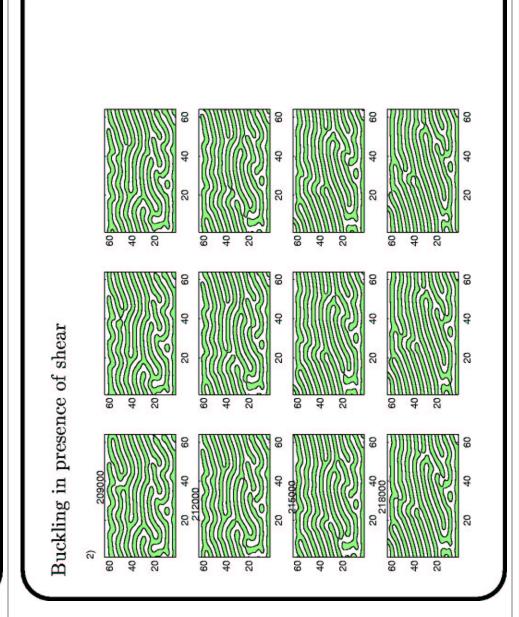


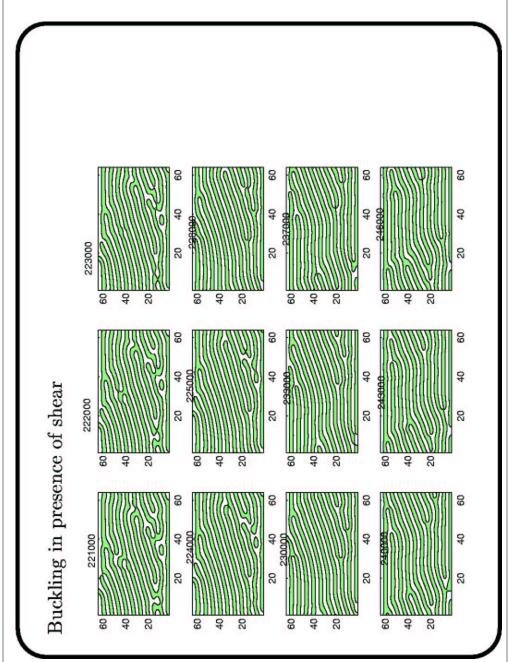


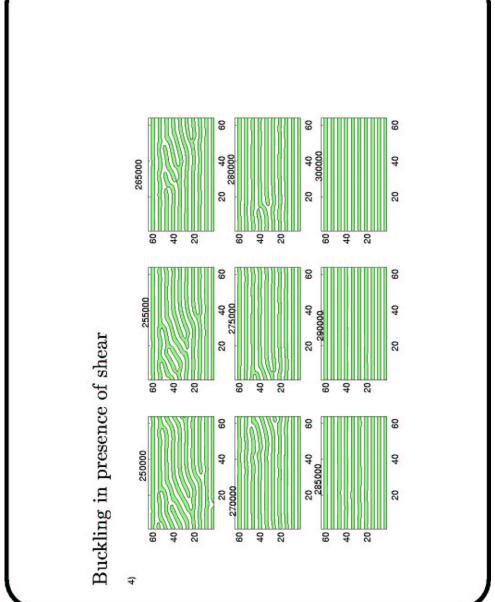


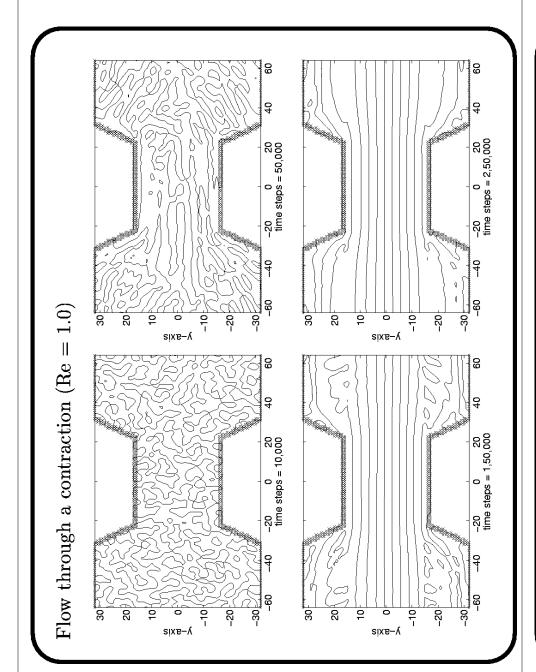
Buckling in presence of shear

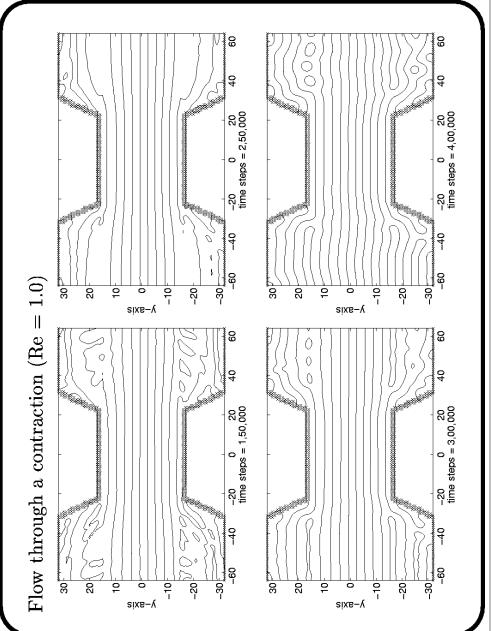






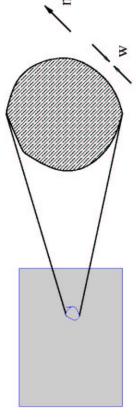






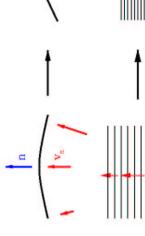
Continuum macroscopic model

Variables

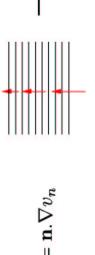


- Fluid velocity $\mathbf{u}(\mathbf{x})$
- Local unit normal direction $\mathbf{n}(\mathbf{x})$
- Local layer spacing $w(\mathbf{x})$
- Surfactant concentration along layers c
- Velocity of surfactants $\mathbf{v}(\mathbf{x})$ with components v_n and \mathbf{v}_t along normal and tangential directions.

Evolution equation for n, w and



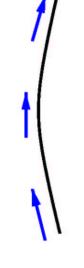
 $\frac{D\mathbf{n}}{Dt}$



$$\frac{w}{t} = \mathbf{n} \cdot \nabla v_n$$

$$\frac{w}{t} = \mathbf{r} \cdot \nabla v_n$$

$$\frac{\partial c}{\partial t} + \nabla_s (\cdot \mathbf{v}_t c) = 0$$



Stress balance for fluid phase (normal direction)

$$-\mathbf{n}\cdot\nabla p + \chi(u_n - v_n) = 0$$

Stress balance for fluid phase (tangential direction)

$$(\mathbf{I} - \mathbf{n}\mathbf{n}).(-\nabla p + \eta \nabla^2 \mathbf{u} + \chi_t (\mathbf{u} - \mathbf{v})) = 0$$

Stress balance for surfactant (normal direction)

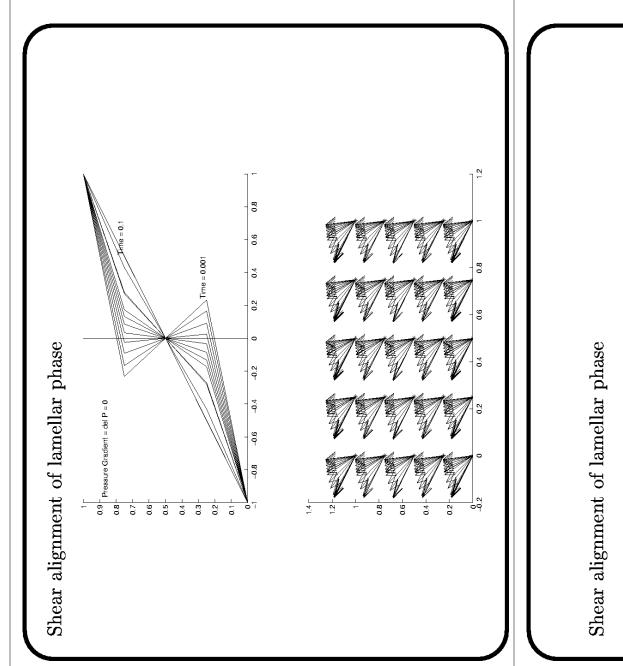
Stress balance for surfactant (normal direction)
$$B\nabla_{s}.\mathbf{n} + K\kappa - \chi(u_{n} - v_{n}) = 0$$

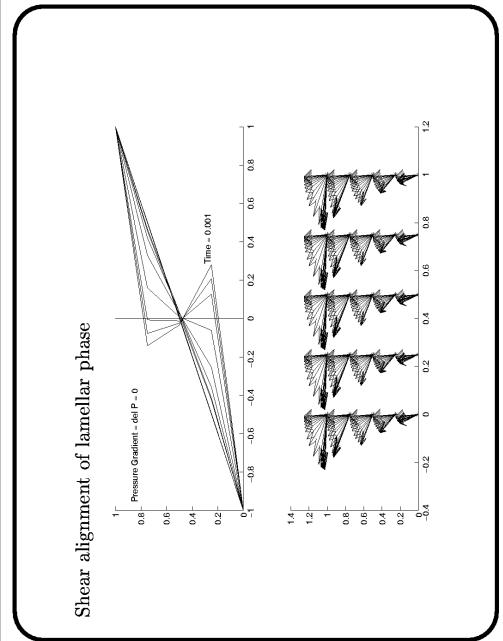
Stress balance for surfactant (tangential direction)

$$-\nabla_s \Pi + \chi_t (\mathbf{u}_t - \mathbf{v}_t) = 0 \qquad [\Pi]$$

Continuum Macroscopic model

- Given initial configurations for w, \mathbf{n} and c, \mathbf{v} and \mathbf{u} , equations can be solved for time evolution of these parameters.
- Equations simplify in the case c is a constant, and $\mathbf{v}_t = \mathbf{u}_t$.
- Predicts simple effects observed, such as alignment of layers under shear.
- Particles can be incorporated using boundary conditions at the surface for unit normal and concentration fields.
- Necessary to incorporate defects within this description!





Conclusions

- Formulation of Lattice Boltzmann simulations for determining dynamical properties.
 - Predicts all qualitative effects observed in lamellar phase rheology in two dimensions.
- Larger 3D computation carried out using parallel computer.
- Microscopic parameters such as local unit normal and layer Coarse grained macroscopic model for the lamellar phase. spacing incorporated in the model.
- Work in progress to incorporate defects in this Macroscopic model correctly predicts effects such as shear alignment. model.