The Black Hole Interior in SYK

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Outline

- SYK Black Hole Microstates. Atypical and Typical States.
- Necessity of State Dependence.
- AdS/CFT as a QEC.
- A Dictionary for the Interior from QEC.
- Comments on the Papadodimas-Raju proposal.

[Kourkoulou, Maldacena]

- N Majorana Fermions = N/2 Spins. S_{i}
- $S_k = 2i\psi^{2k-1}\psi^{2k}$

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• Lower Energy via Euclidean Ev.

$$|B_s^\beta\rangle = e^{-\frac{\beta}{2}H}|B_s\rangle$$

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[Kourkoulou, Maldacena]

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- $\mathcal{O}(N)$ Symmetry, ~ $\mathcal{O}(1/N^{q-1})$. "Flip" subgroup: Flip sign of any even fermion. Relates any two states $|B_s\rangle$ and $|B_{s'}^{\beta}\rangle$

e.g.
$$\psi^2 \to -\psi^2 \implies |\uparrow\uparrow \dots \uparrow\rangle \to |\downarrow\uparrow \dots \uparrow\rangle$$

 $S_k = 2i\psi^{2k-1}\psi^{2k}$

[Kourkoulou, Maldacena]

• **Diagonal Correlator:** $_{R}\langle B_{s}^{\beta}|\psi_{R}^{1}\psi_{R}^{1}|B_{s}^{\beta}\rangle_{R} = {}_{LR}\langle\beta|\Big[|B_{s}\rangle_{L}\langle B_{s}|\otimes\psi_{R}^{1}\psi_{R}^{1}\Big]|\beta\rangle_{LR}$

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$$= 2i \times {}_{LR}\langle\beta|\Big[|B_{s}\rangle_{L}\langle B_{s}|\psi_{L}^{1}(0)\psi_{L}^{2}(0)\otimes\psi_{R}^{1}\psi_{R}^{2}\Big]|\beta\rangle_{LR}$$
$$\begin{cases} S_{k}|B_{s}\rangle = s_{k}|B_{s}\rangle_{LR} \\ S_{k} = 2i\psi^{2k-1}\psi^{2k} \\ S_{k} = 2i\psi^{2k-1}\psi^{2k} \end{cases}$$

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• Bulk is deduced from the correlation functions (a) $1 \ll \beta J$

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• These are atypical BH microstates: Simple observables have not thermalized.

- Begin with wormhole with OTO shockwaves: $|W\beta\rangle_{LR} \equiv W_L|\beta\rangle_{LR}$
- Then project: $|B_s^{W_\beta}\rangle_R \equiv {}_L\langle B_s|W\beta\rangle_{LR}$

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 $\sim e^{-n_{Shocks}}$



[AA, Zhenbin Yang - WIP]

- SYK: $H_{SYK} \sim \sum_{i_1...i_q}^{I_1} J_{i_1...i_q} \psi^{i_1}...\psi^{i_q}$
- Raising and lowering operators: $\chi_k^{\sigma} = \psi^{2k-1} + i\sigma_k \psi^{2k}$
- Modified SYK: $H_{SYK}^{Mod} \sim \sum_{\sigma_1...\sigma_q}^{\pm 1} \sum_{k_1...k_q}^{N/2} J_{k_1...k_q}^{\sigma_1...\sigma_q} \chi_{k_1}^{\sigma_1} ... \chi_{k_q}^{\sigma_q} \Big|_{\prod_k^q \sigma_k = +1} \langle J_{a_1,...a_q}^{\sigma_1,...\sigma_q} J_{b_1,...b_q}^{\tilde{\sigma}_1,...\tilde{\sigma}_q} \rangle = \mathcal{J}^2 \delta_{a_1,b_1} ... \delta_{a_q,b_q} \delta_{\sigma_1+\tilde{\sigma}_1} ... \delta_{\sigma_q+\tilde{\sigma}_q}$

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- Modified SYK: $H_{SYK}^{Mod} \sim \sum_{\sigma_1...\sigma_q}^{\pm 1} \sum_{k_1...k_q}^{N/2} J_{k_1...k_q}^{\sigma_1...\sigma_q} \chi_{k_1}^{\sigma_1} ... \chi_{k_q}^{\sigma_q} \Big|_{\prod_k^q \sigma_k = +1} \langle J_{a_1,...a_q}^{\sigma_1,...\sigma_q} J_{b_1,...b_q}^{\tilde{\sigma}_1,...\tilde{\sigma}_q} \rangle = \mathcal{J}^2 \delta_{a_1,b_1} ... \delta_{a_q,b_q} \delta_{\sigma_1+\tilde{\sigma}_1} ... \delta_{\sigma_q+\tilde{\sigma}_q}$

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 $|\uparrow\uparrow\dots\uparrow\rangle+O(1)$ spin flips

 $|\uparrow\downarrow\uparrow\ldots\downarrow\rangle$, O(N) spin flips

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- [AA, Zhenbin Yang WIP]
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Different Microstates

• Atypical Microstates

 $|B_s^\beta\rangle_R = {}_L \langle B_s |\beta\rangle_{LR}$

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- Both are over-complete bases of BH microstates of temperature β
- Does this have consequences for interior reconstruction?

• Suppose the existence of a linear operator N_R^F that measures whether there is a shockwave behind the horizon.

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R

• Suppose the existence of a linear operator N_R^F that measures whether there is a shockwave behind the horizon.

$$2^{-N/2} \sum_{s} {}_{R} \langle B_{s}^{\beta} | N_{R}^{F} | B_{s}^{\beta} \rangle_{R} \sim \mathcal{O}(1/N^{q-1})$$
$$2^{-N/2} \sum_{s} {}_{R} \langle B_{s}^{W_{\beta}} | N_{R}^{F} | B_{s}^{W_{\beta}} \rangle_{R} \sim \mathcal{O}(1)$$





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• Contradiction! Can't be the same by assumption!

Take Away

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- Can we find such a dictionary?
- Yes! By assuming the AdS/CFT dictionary to be a QEC code, one can generate such a dictionary for the interior!
- Key: the dictionary is fluid, responding to the projection in a way that maps the interior operators to the remaining boundary in a state dependent way.



AdS/CFT as QEC

- I will describe the duality using a circuit diagram.
- Consider Isometry *T* from systems *a* & *H* into *L*.



Operator Mapping [HPPY, Hayden et al]

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$$S(\rho_R^{\widetilde{\psi}}) = \ln \mathcal{D}_H + S(\rho_b^{\psi})$$





[Harlow]

Subregion-Subregion Duality



Subregion-Subregion Duality



Subregion-Subregion Duality





Projected BH Code

- Projected BH has 'same' geometry.
- Project on original tensor network.
- We want to understand the conditions on *P* such that operators acting on *a* map to *R*





























Properties of the Dictionary

 ${}_{L}\langle P|\widetilde{\mathcal{O}}_{a}|\widetilde{\psi}\rangle_{LR} = \mathcal{O}_{R}^{P} {}_{L}\langle P|\widetilde{\psi}\rangle_{LR}$

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Relation to PR

- Problem with PR: Consider Typical Microstates
- All simple operators have thermalized
- Mirror construction would predict 'nothing' behind the horizon.



 $_L\langle B_s^*|W_L|\beta\rangle_{LR}$

Relation to PR

- Problem with PR: Consider Typical Microstates
- All simple operators have thermalized





 ${}_L\langle B_s^*|W_L|\beta\rangle_{LR}$

- The QEC construction is not a statement about typical states, but states constructed in a special way.
- It IS sensitive to what happens behind the horizon.

Thank You!

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